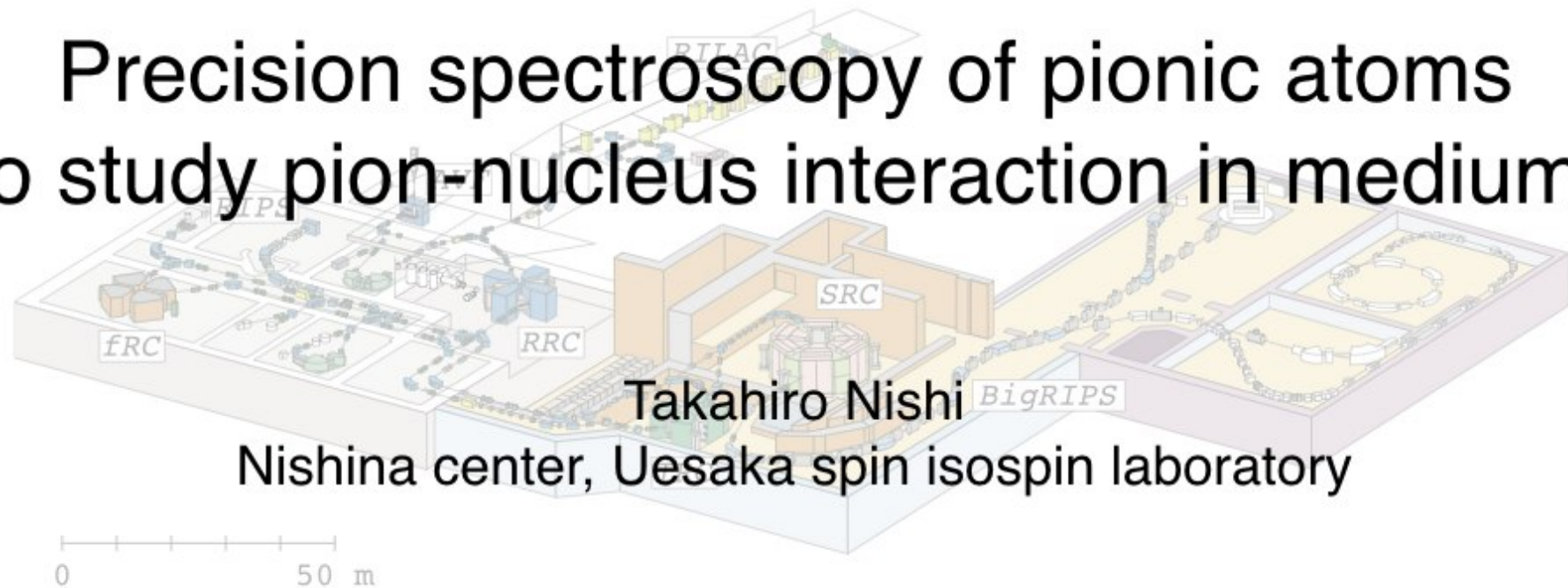


# Precision spectroscopy of pionic atoms to study pion-nucleus interaction in medium



# Collaborators

DeukSoon Ahn, Georg P.A. Berg, Masanori Dozono, Daijiro Etoh, Hiroyuki Fujioka, Naoki Fukuda, Nobuhisa Fukunishi, Hans Geissel, Emma Haettner, Tadashi Hashimoto, Ryugo S. Hayano, Satoru Hirenzaki, Hiroshi Horii, Natsumi Ikeno, Naoto Inabe, Kenta Itahashi\* , Sathoshi Itoh, Masahiko Iwasaki, Daisuke Kameda, Shouichiro Kawase, Keichi Kisamori, Yu Kiyokawa, Toshiyuki Kubo, Kensuke Kusaka, Hiroaki Matsubara, Masafumi Matsushita, Shin'ichiro Michimasa, Kenjiro Miki, Go Mishima, Hiroyuki Miya, Daichi Murai, Yohei Murakami, Hideko Nagahiro, Masaki Nakamura, Megumi Niikura, Takahiro Nishi\*\*, Shumpei Noji, Kota Okochi, Shinsuke Ota, Naruhiko Sakamoto, Kimiko Sekiguchi, Hiroshi Suzuki, Ken Suzuki, Motonobu Takaki, Hiroyuki Takeda, Yoshiki K. Tanaka, Koichi Todoroki, Kyo Tsukada, Tomohiro Uesaka, Yasumori Wada, Yuni N. Watanabe, Helmut Weick, Hiroyuki Yamada, Hiroki Yamakami, Yoshiyuki Yanagisawa and Koichi Yoshida



\*spokesperson, \*\* co-spokesperson

University of Tokyo, RIKEN, Nishina Center, University of Notre Dame, Tohoku University, Kyoto University, GSI Helmholtzzentrum für Schwerionenforschung GmbH, Nara Women's University, Osaka University, Stefan Meyer Institute

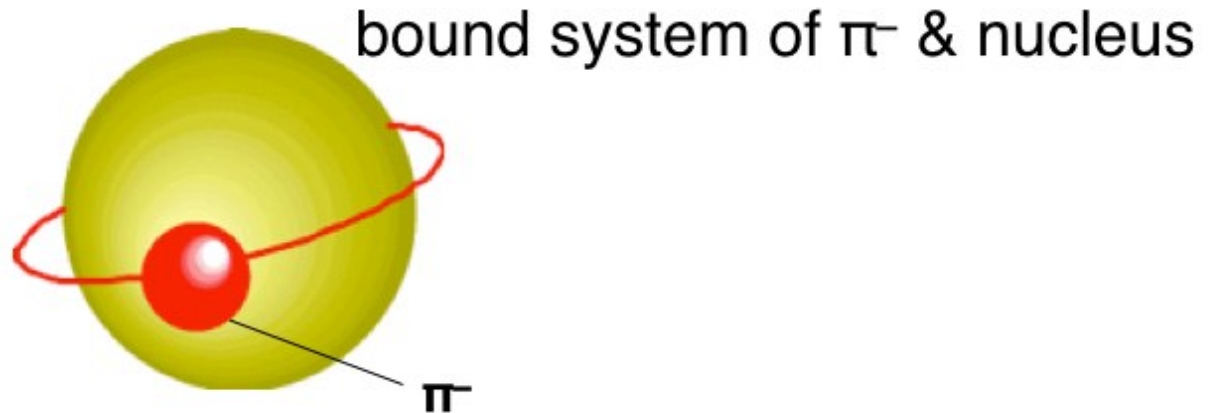
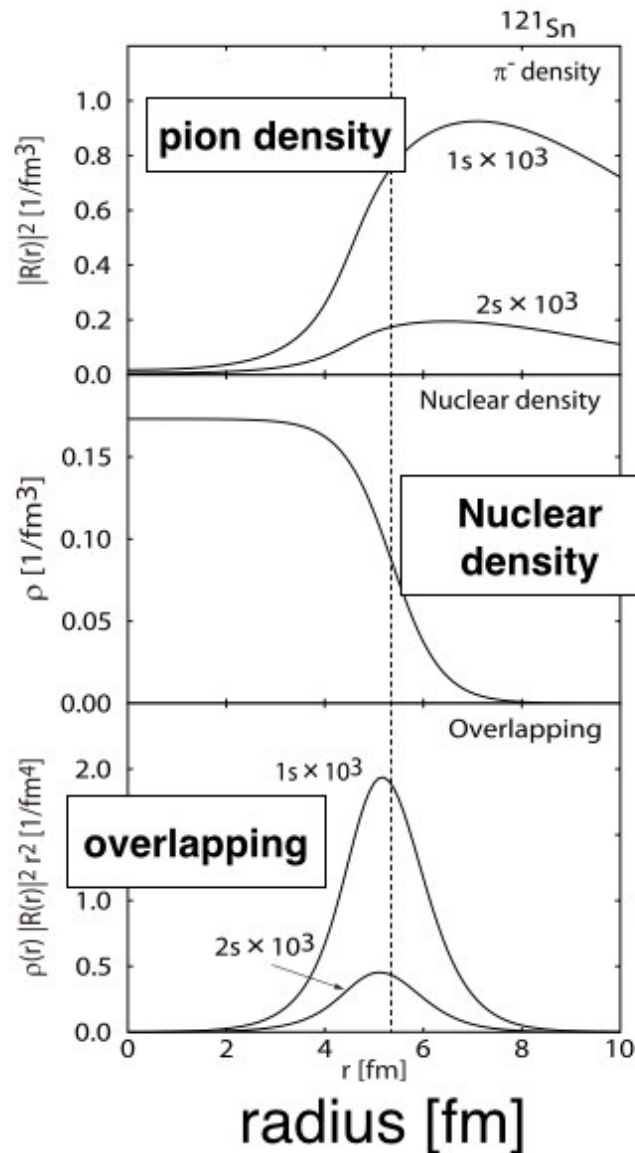
# Index

- (i) Motivation and prior research
- (ii) Experiment performed in 2014 at RIKEN
- (iii) Data analysis (pID / optics / decomposition)
- (iv) Summary and future plan

(i) Motivation and prior research

(i) Motivation and prior research

# Deeply bound pionic atoms



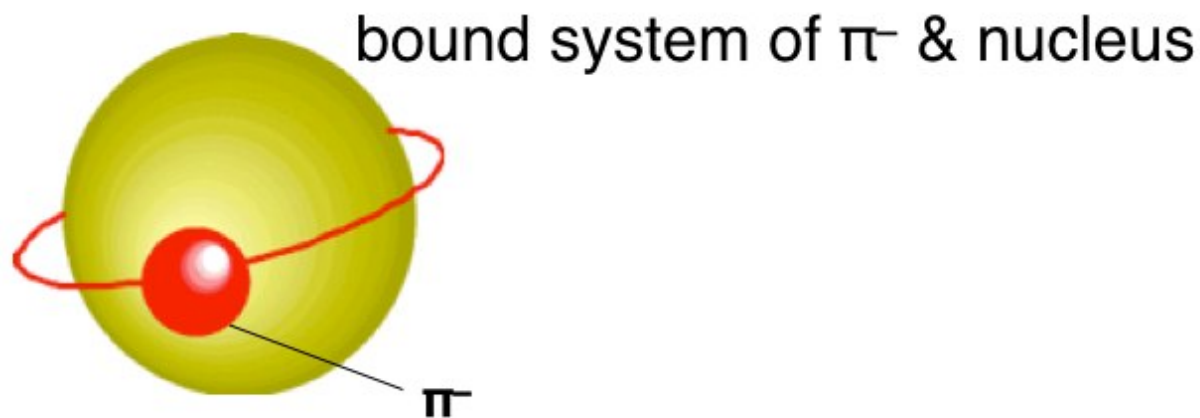
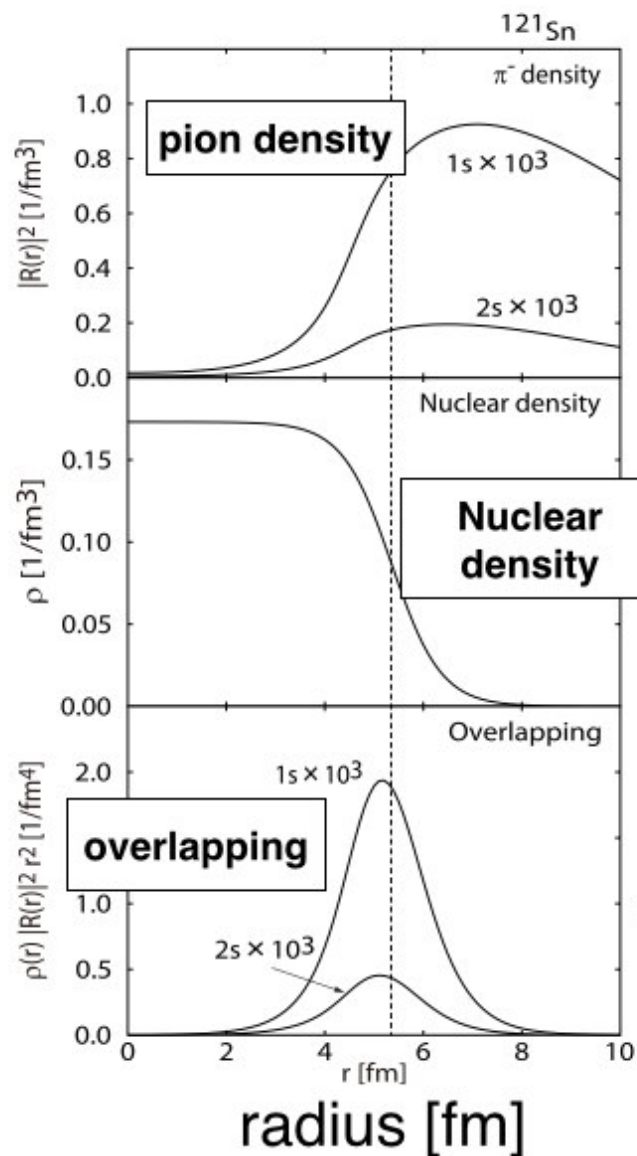
very unique objects in terms of

- 1) Bound system of meson-nucleus
- 2) Quasi-stable state



(i) Motivation and prior research

# Deeply bound pionic atoms



very unique objects in terms of

- 1) Bound system of meson-nucleus  
**No other "meson in nucleus" system established**
- 2) Quasi-stable state  
**Quantitative evaluation of strong interaction in medium**

(i) Motivation and prior research

# Deeply bound pionic atoms

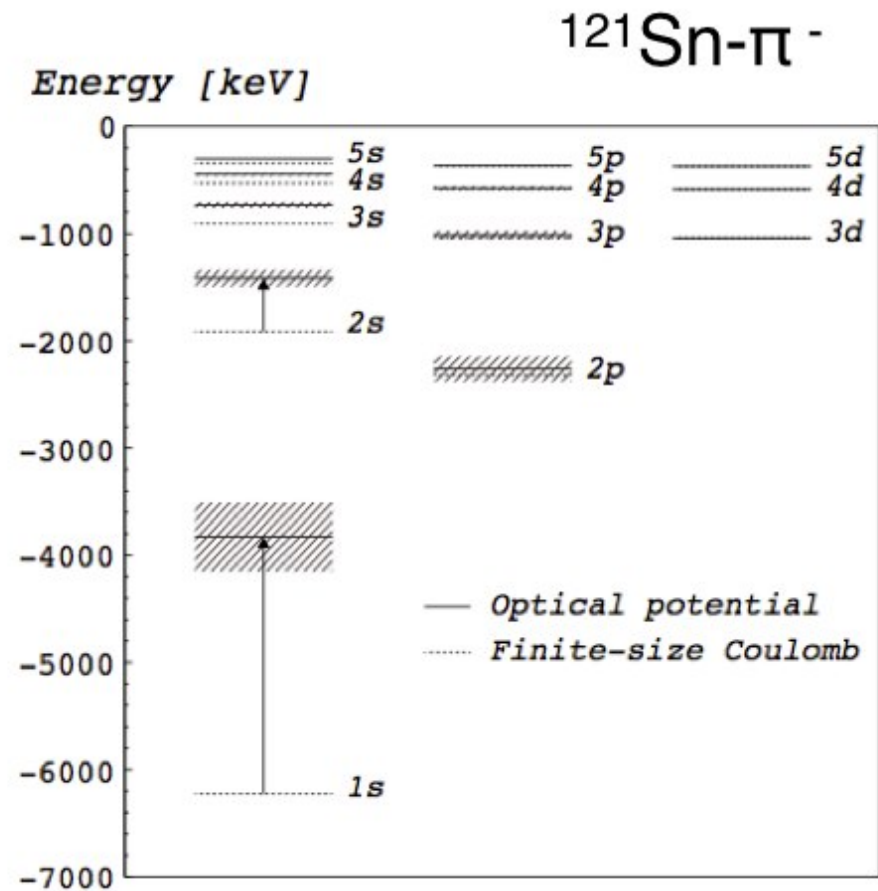
**BE,  $\Gamma$  of 1s pionic state  
 $\Leftrightarrow$  strong interaction effect**

$\pi$ -A s-wave optical potential (s-wave)

$$V_s(r) = -\frac{2\pi}{\mu} [\epsilon_1 \{b_0\rho + b_1\delta\rho\} + \epsilon_2 B_0\rho^2]$$

$$\rho = \rho_p + \rho_n$$

$$\delta\rho = \rho_p - \rho_n$$



N. Ikeno et al., Prog. Theor. Phys. 126 (2011) 483.  
S. Itoh, Doctoral Dissertation, Univ. of Tokyo (2011)

(i) Motivation and prior research

# Deeply bound pionic atoms

**BE,  $\Gamma$  of 1s pionic state  
 $\Leftrightarrow$  strong interaction effect**

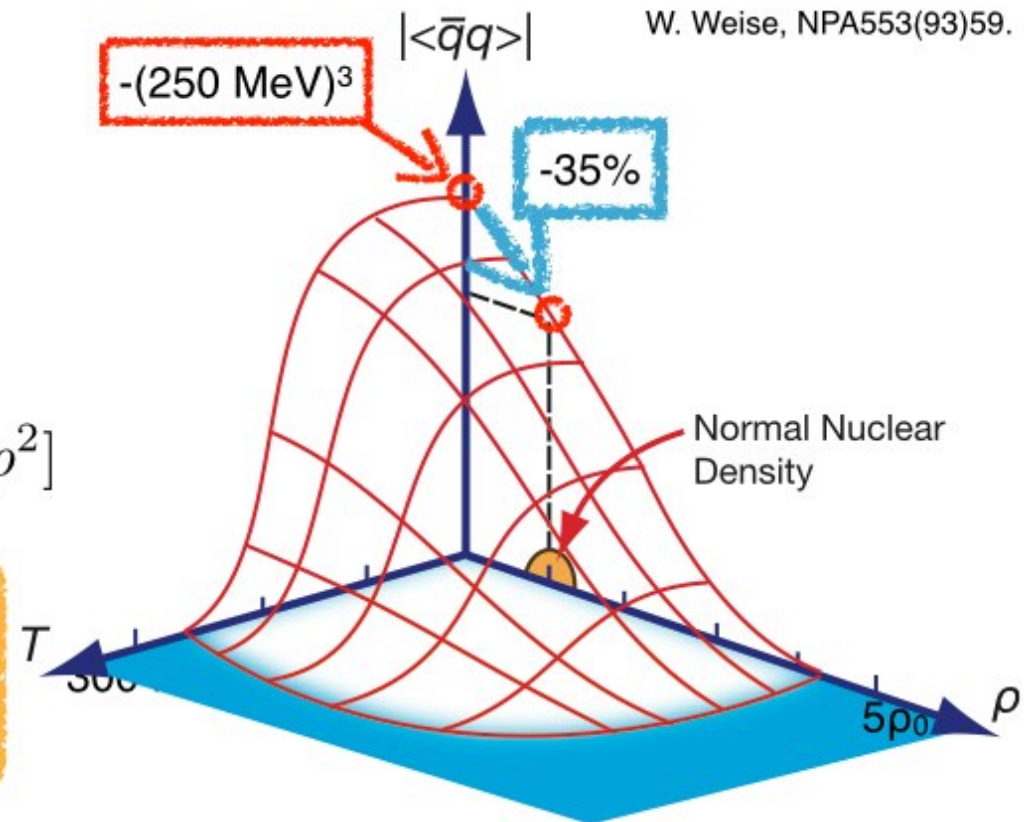
$\pi$ -A s-wave optical potential (s-wave)

$$V_s(r) = -\frac{2\pi}{\mu} [\epsilon_1 \{b_0\rho + b_1\delta\rho\} + \epsilon_2 B_0\rho^2]$$

$$\frac{\langle \bar{q}q \rangle_{\rho=\rho_0}}{\langle \bar{q}q \rangle_{\rho=0}} \simeq \left( \frac{m_\pi^*}{m_\pi} \right)^2 \frac{b_1(\rho=0)}{b_1(\rho=\rho_0)}$$

Quantitative evaluation of  $\langle \bar{q}q \rangle$

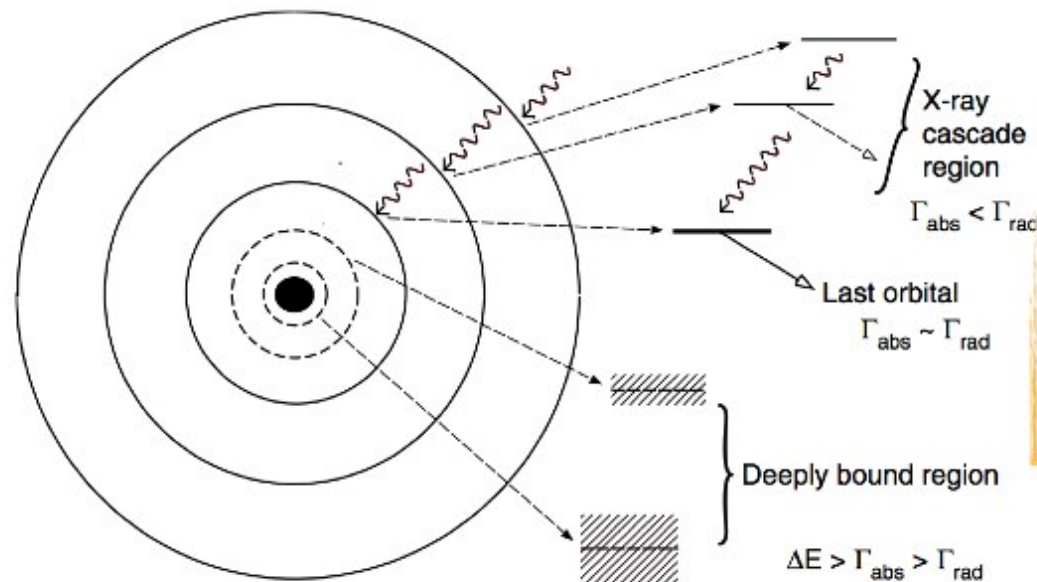
**Order parameter of  
chiral symmetry breaking**





(i) Motivation and prior research

## Conventional method: X-ray spectroscopy



Yamazaki *et al*, Phys. Rep. 514, 1(2012)

x rays during atomic cascade  
→ higher orbits / light nuclei  
( $\sim {}^{24}\text{Mg}$  for 1s)

pionic 1s state in H  
→  $b_1$  in vacuum

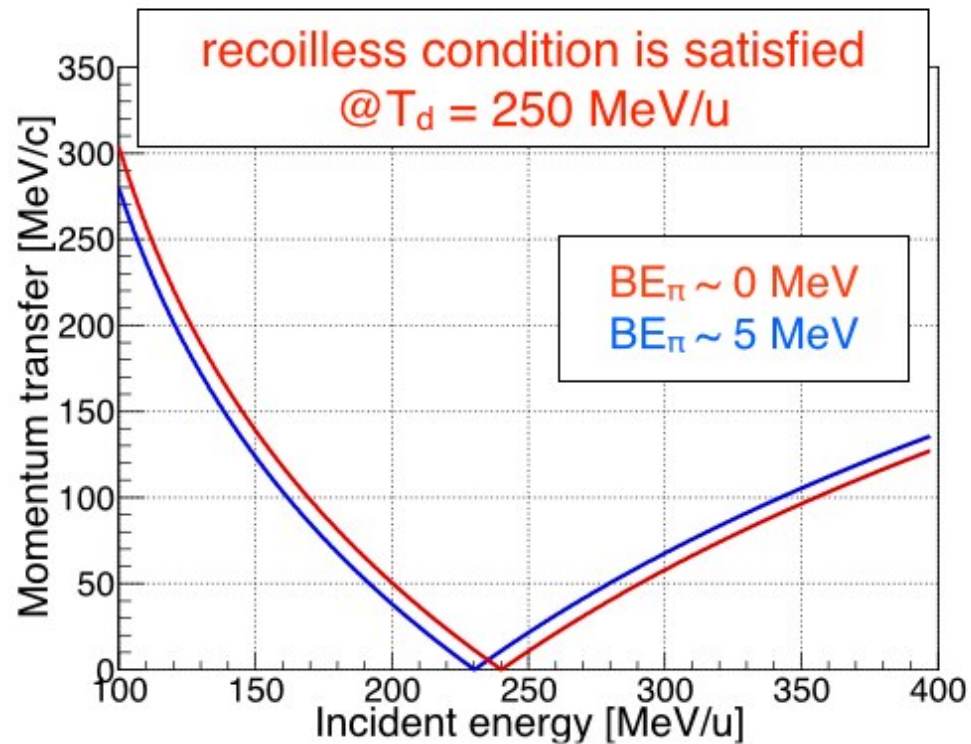
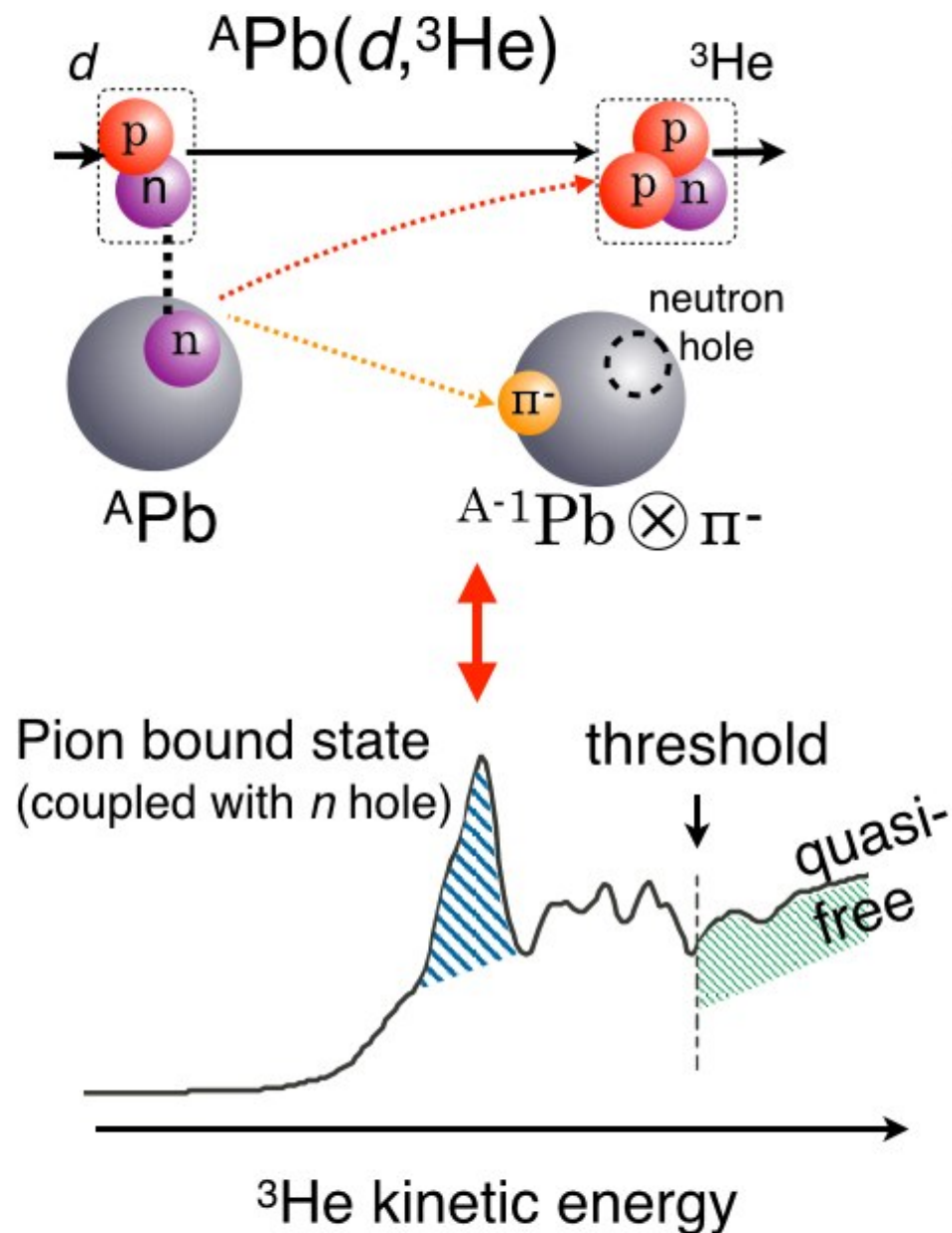
H. Schröder *et al.*, Eur. Phys. J. C 21, 473 (2001).

for deep orbit in heavy nuclei  
absorption is faster

**This method cannot produce “deeply-bound” pionic atom...**

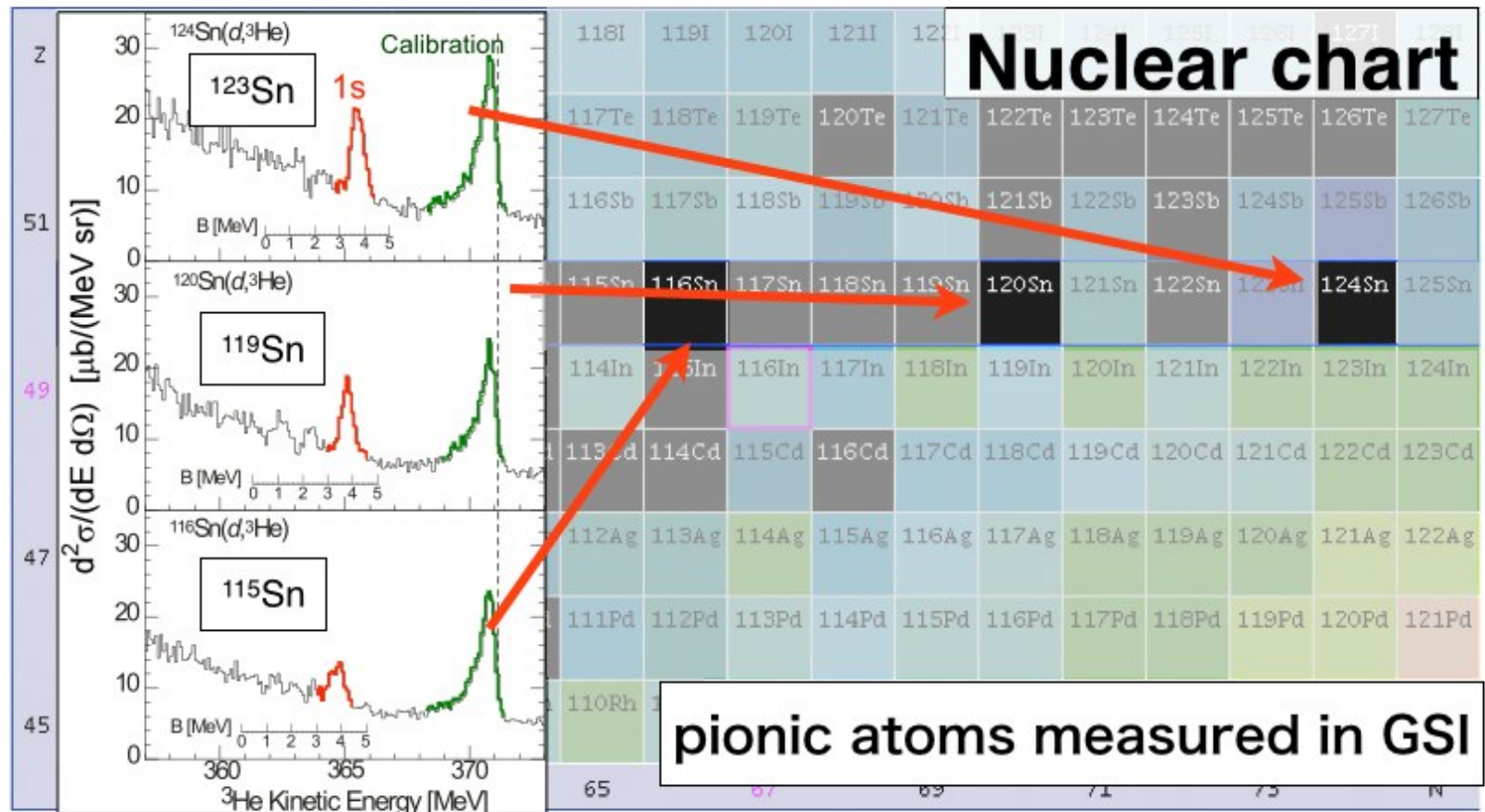
(i) Motivation and prior research

# Our method: Missing mass spectroscopy



(i) Motivation and prior research

## Prior research: Experiment at GSI



K. Suzuki et al., PRL92 072302 (2004)

NuDat

## Systematic study of pionic Sn isotopes

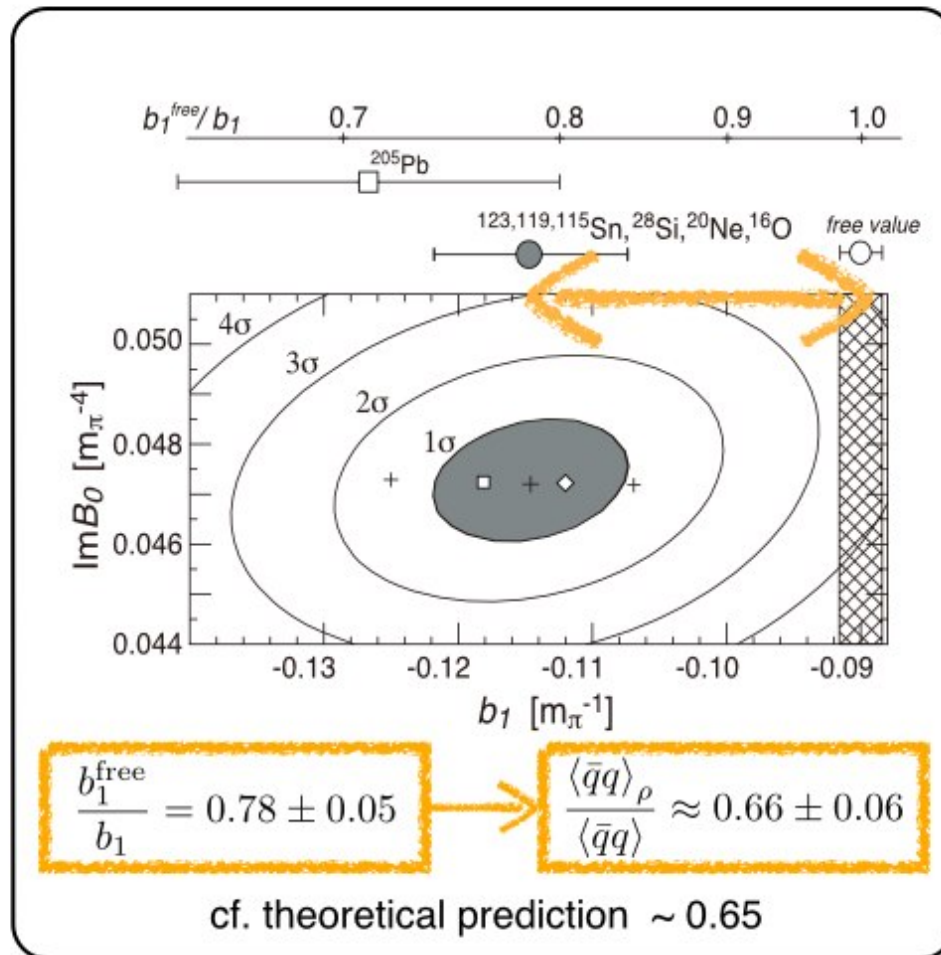
~ **3 month** measurement for 3 isotopes



(i) Motivation and prior research

# Extract $b_1$ from experimental data at GSI

Contour plot of  $\chi^2$



$\pi$ -A s-wave optical potential

$$V_s(r) = -\frac{2\pi}{\mu} [\epsilon_1 \{b_0 \rho + b_1 \delta \rho\} + \epsilon_2 B_0 \rho^2]$$

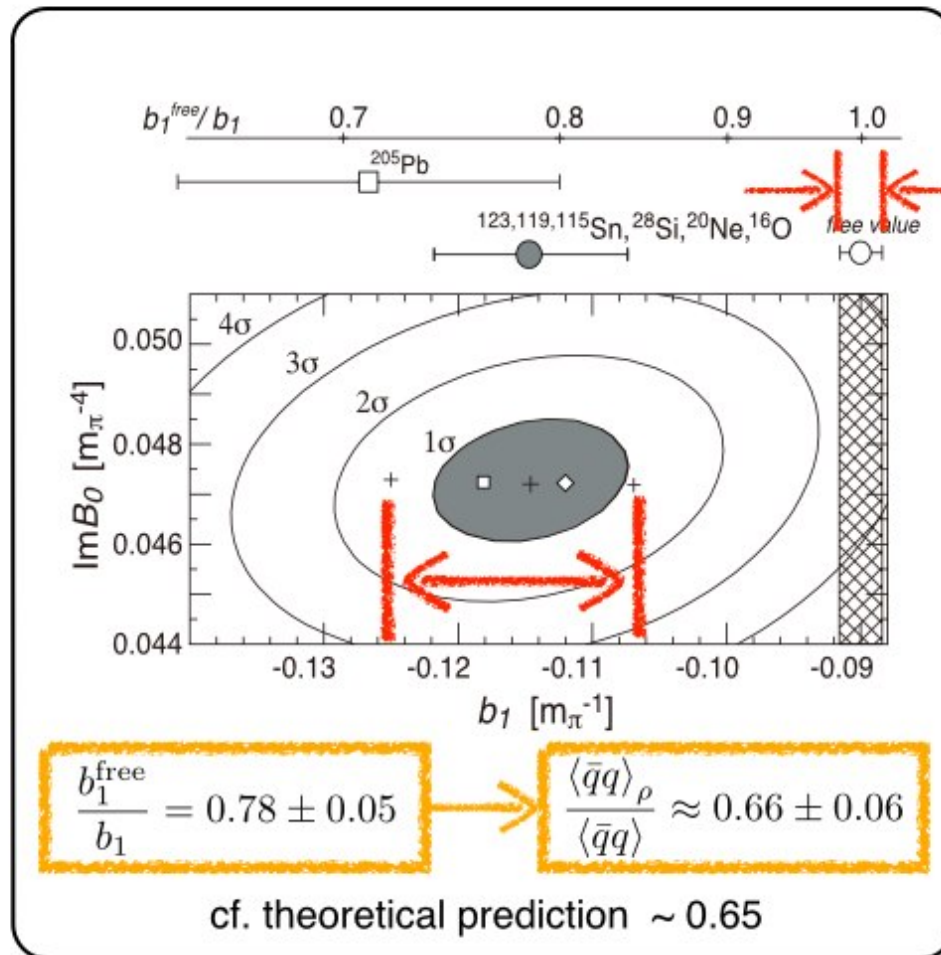
\*  $b_0$ ,  $\text{Re}B_0$  are deduced from data of light / symmetric pionic atoms



(i) Motivation and prior research

# Extract $b_1$ from experimental data at GSI

Contour plot of  $\chi^2$



$\pi$ -A s-wave optical potential

$$V_s(r) = -\frac{2\pi}{\mu} [\epsilon_1 \{b_0 \rho + b_1 \delta \rho\} + \epsilon_2 B_0 \rho^2]$$

error of  $b_1$  in medium is still large compared with that in vacuum!!  
two main sources are

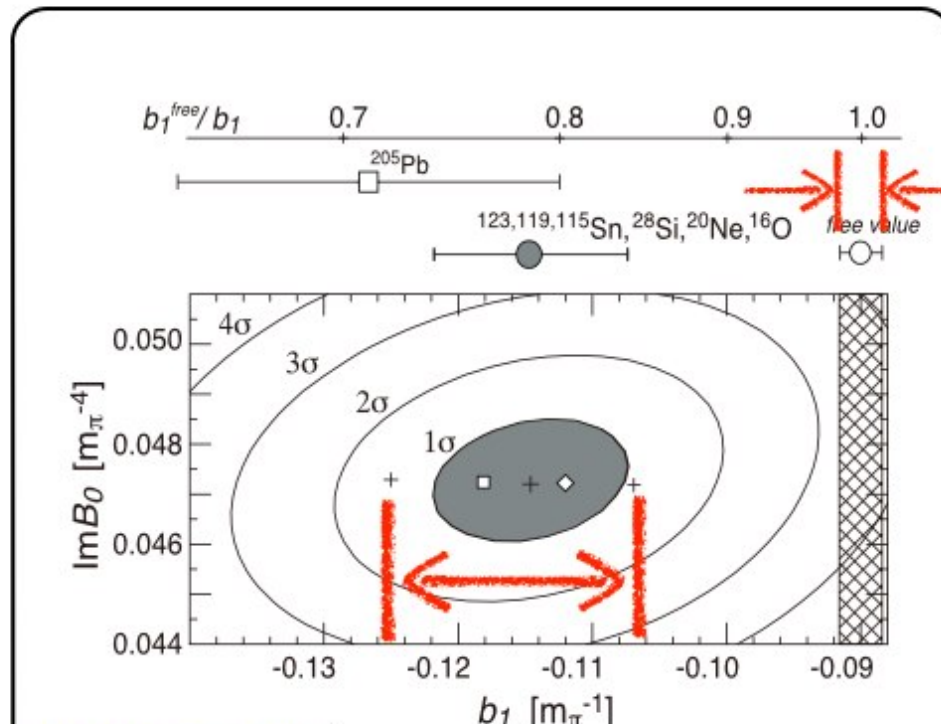
- **experimental error**
- **neutron distribution ambiguities**

\*  $b_0$ ,  $\text{Re}B_0$  are deduced from data of light / symmetric pionic atoms

(i) Motivation and prior research

# Extract $b_1$ from experimental data at GSI

Contour plot of  $\chi^2$



$\pi$ -A s-wave optical potential

$$V_s(r) = -\frac{2\pi}{\mu} [\epsilon_1 \{b_0\rho + b_1\delta\rho\} + \epsilon_2 B_0\rho^2]$$

error of  $b_1$  in medium is still large compared with that in vacuum!!  
two main sources are

- experimental error
- neutron distribution ambiguities

$$\frac{b_1^{\text{free}}}{b_1} = 0.78 \pm$$

cf. the

**To extract  $b_1$  with higher precision**  
improve resolution / calibration  
More isotopes

(ii) Experiment performed in 2014 at RIKEN

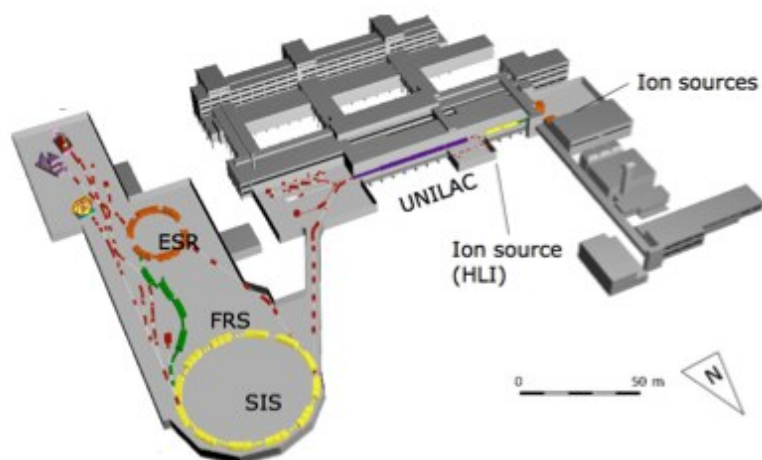
## First production experiment in 2014 @ RIKEN (11 days)



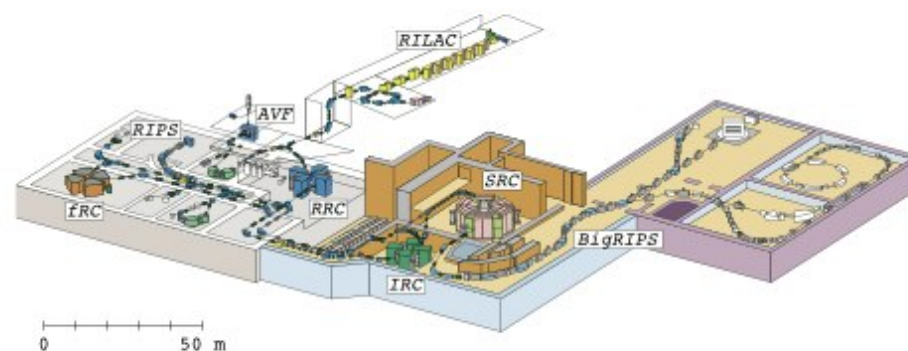


# Comparison between RIBF and GSI

GSI facilities

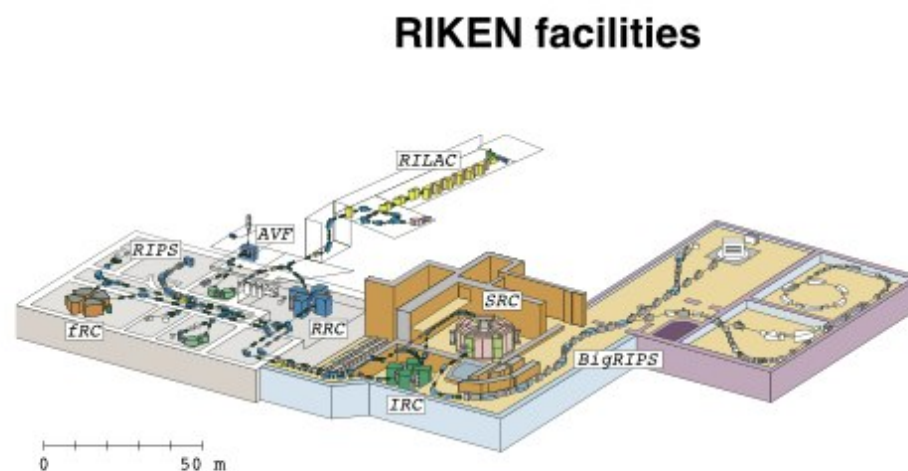
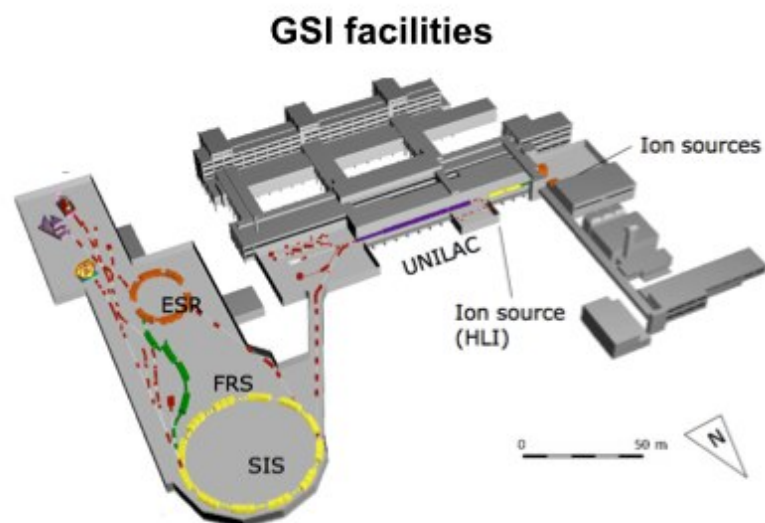


RIKEN facilities



	GSI (FRS)	RIBF (BigRIPS)	Improvement
intensity	$\sim 10^{11} / 6 \text{ s (1 spill)}$	$\sim 10^{12} / \text{s}$	<b><math>\times 60</math></b>
angular acceptance (H / V)	15 / 10 mrad	40 / 60 mrad	$\times 16$
resolution (FWHM)	400 keV	$\sim 600 \text{ keV}$	$\searrow$ factor 1.5 ~ 2

# Comparison between RIBF and GSI

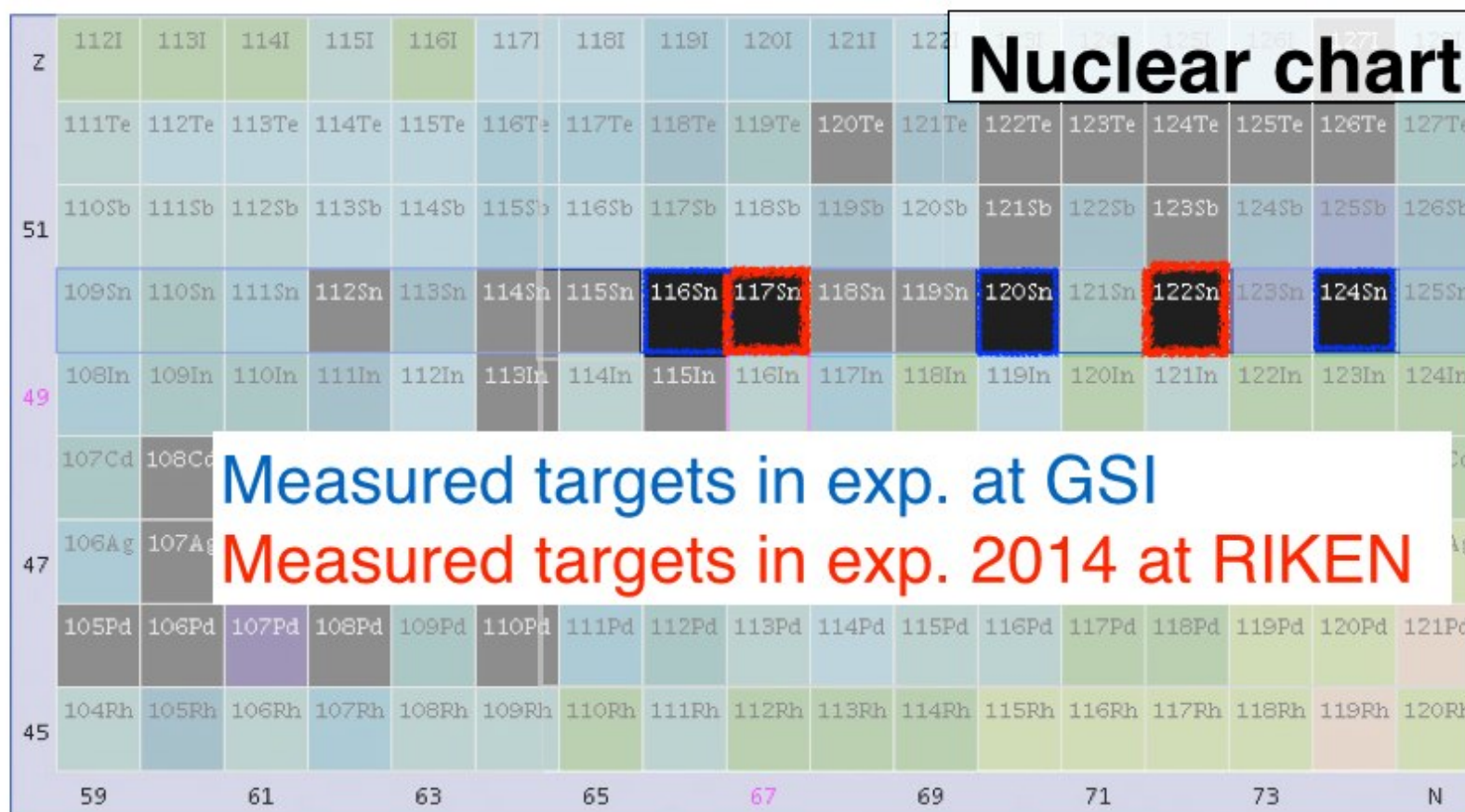


	GSI (FRS)	RIBF (BigRIPS)	Improvement
intensity	$\sim 10^{11} / 6 \text{ s (1 spill)}$	$\sim 10^{12} / \text{s}$	<b><math>\times 60</math></b>
angular acceptance (H / V)	15 / 10 mrad	40 / 60 mrad	$\times 16$
resolution (FWHM)	400 keV	<b>200 ~ 300 keV</b>	improve

by dispersion matching optics

# Targets in the experiment at RIKEN

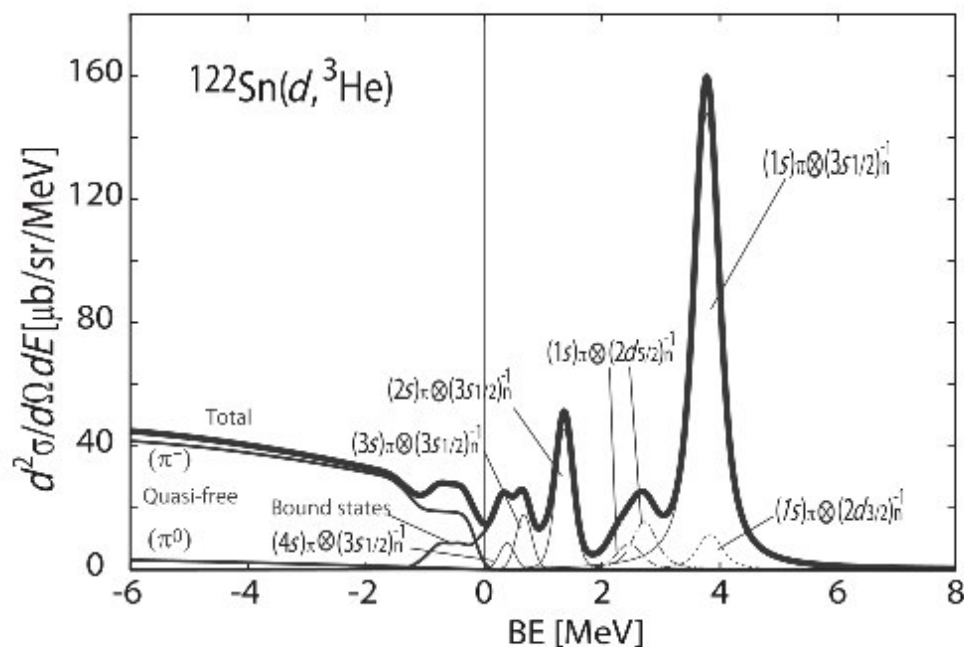
NuDat



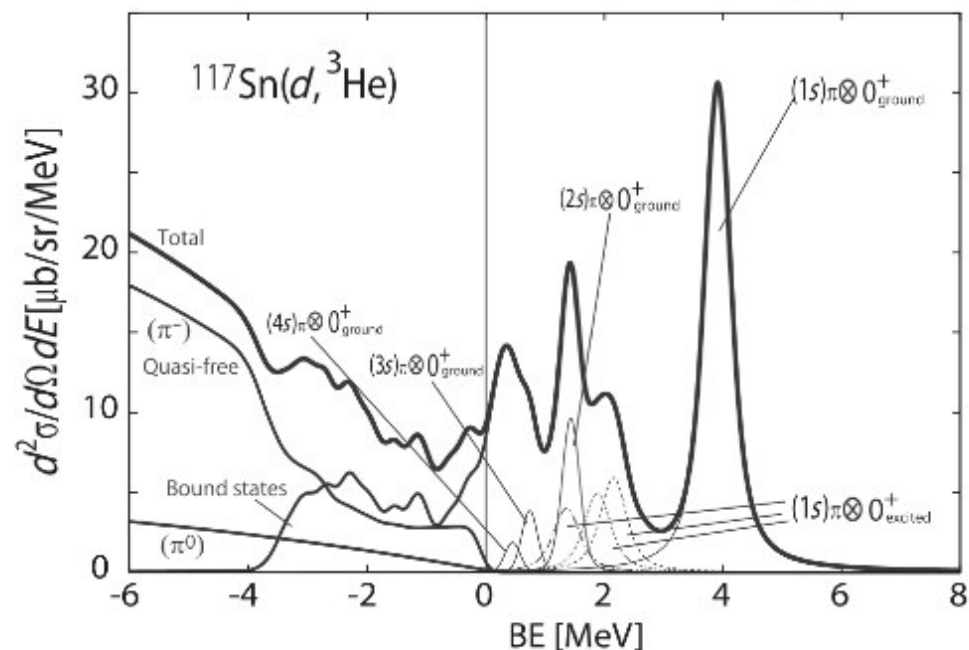


# Expected excitation spectra

$^{122}\text{Sn}$ : relatively large cross section



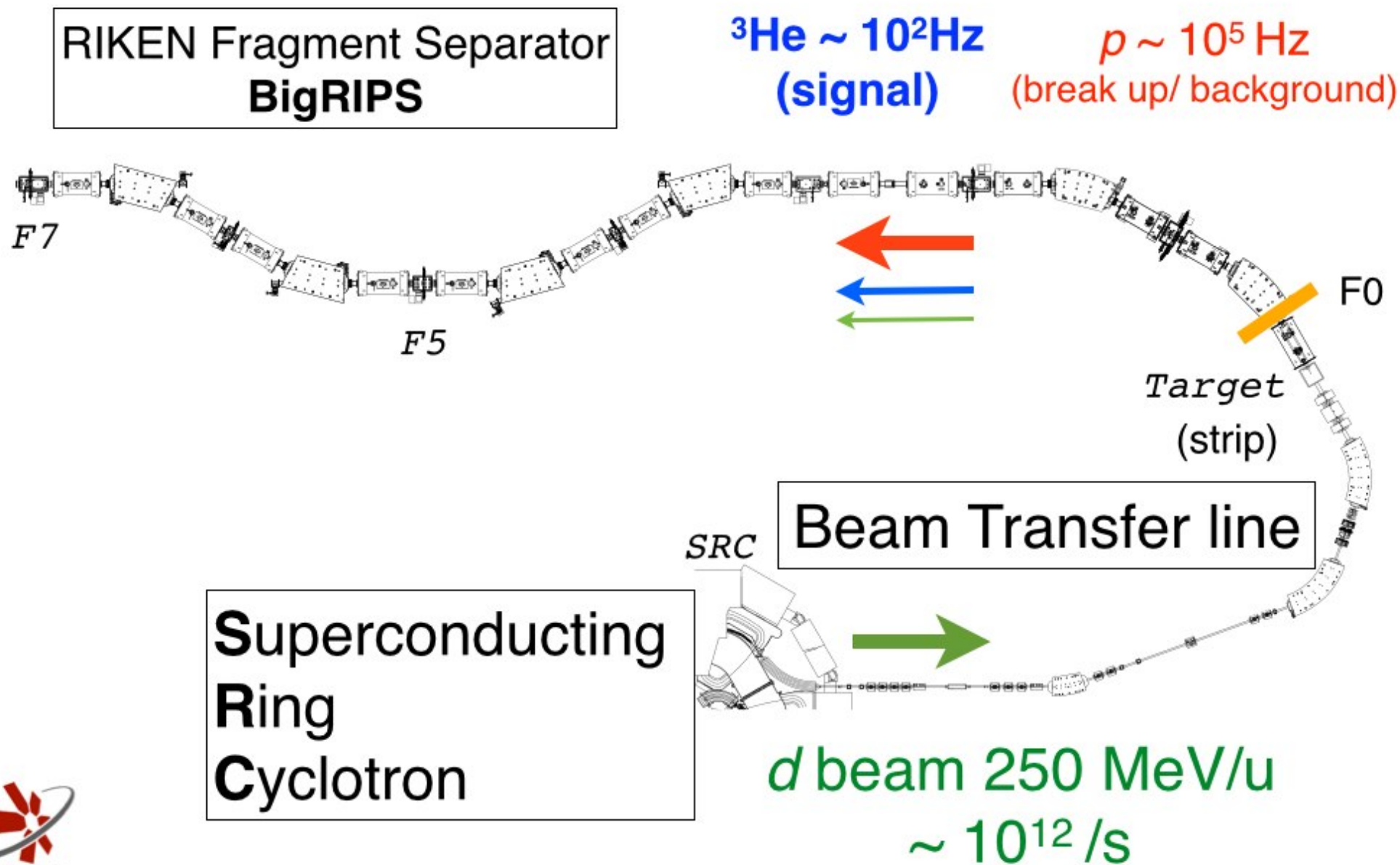
$^{117}\text{Sn}$ : first odd-A target



N. Ikeno et al., Prog. Theor. Exp. Phys. (2013) 2013 (6): 063D01.



# Experimental setup

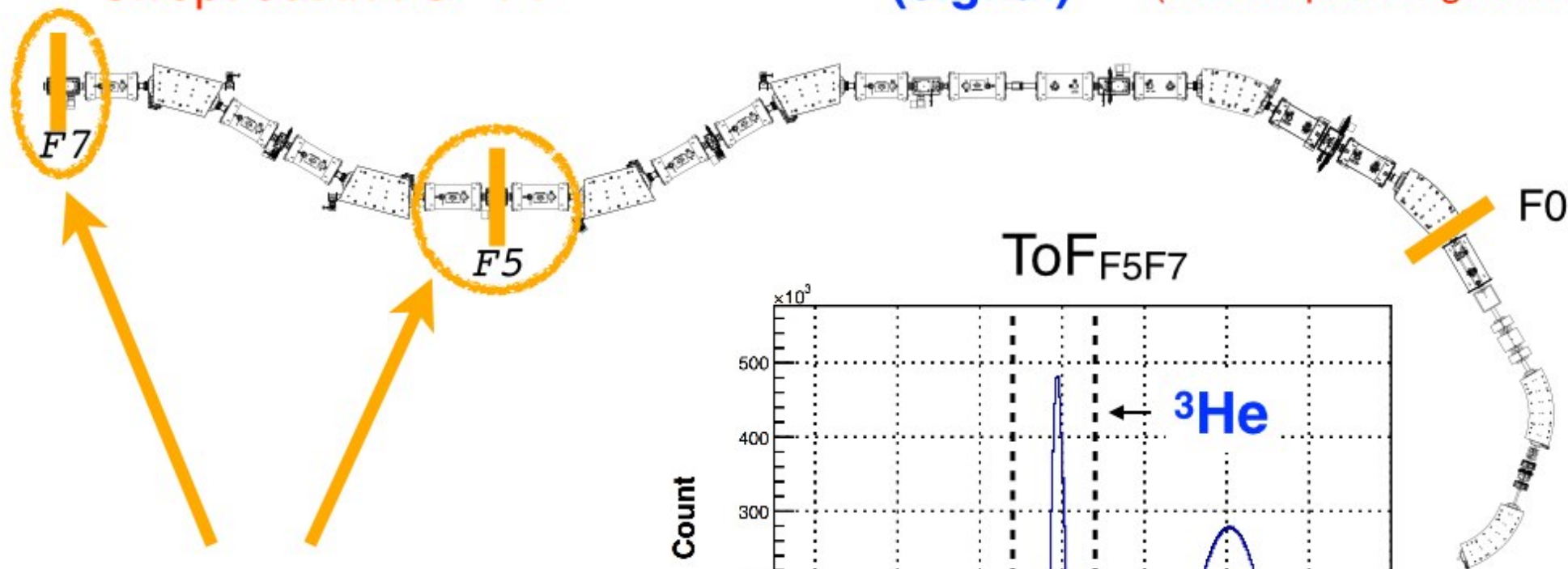


# Experimental setup

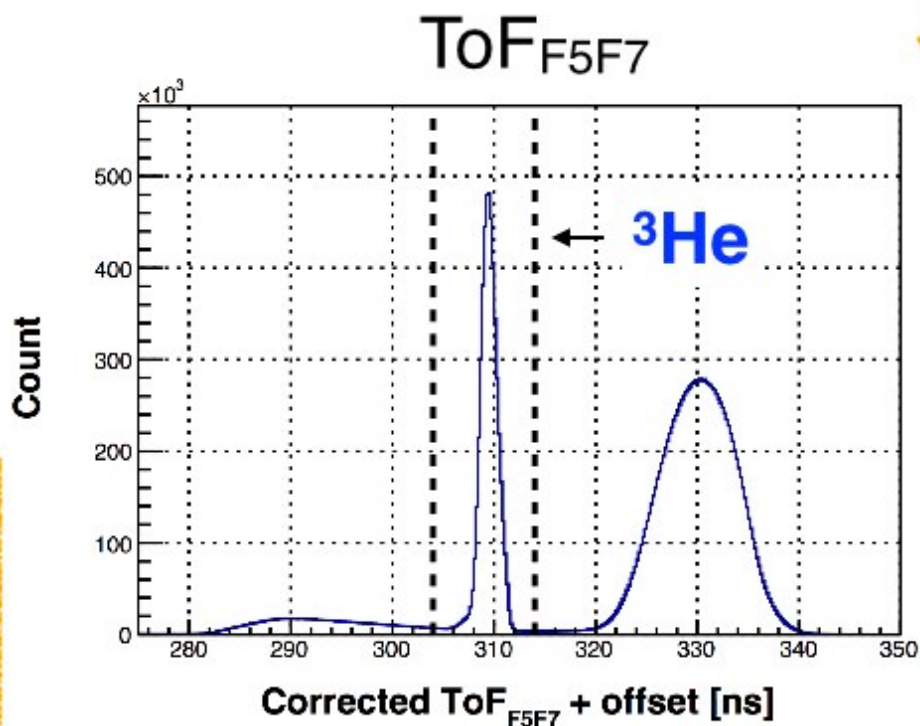
large part of  $p$ :  
swept out in F5 - F7

${}^3\text{He} \sim 10^2 \text{ Hz}$   
(signal)

$p \sim 10^5 \text{ Hz}$   
(break up/ background)



$\Delta E / \text{TOF}$  measurement  
(pID, hardware trigger)  
by **plastic scintillator**

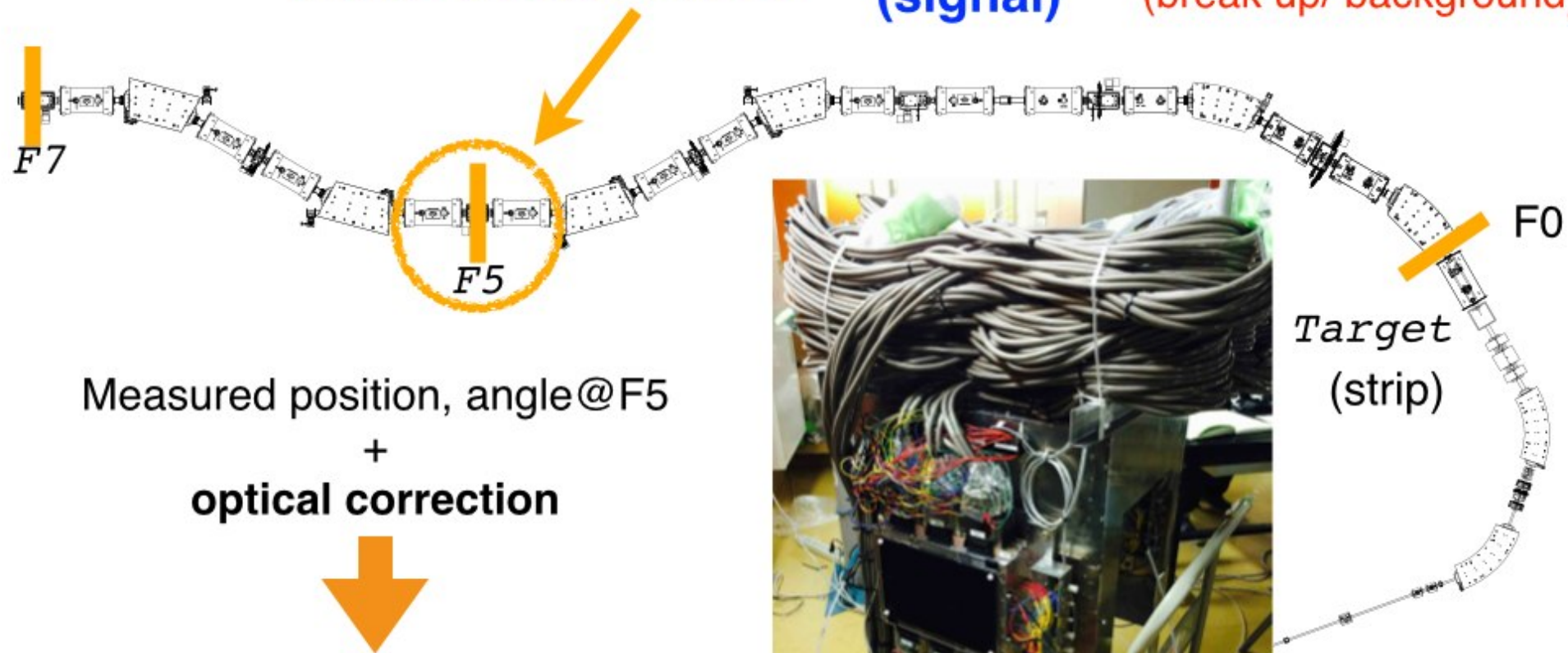


# Experimental setup

Tracking by MWDC

${}^3\text{He} \sim 10^2 \text{ Hz}$   
(signal)

$p \sim 10^5 \text{ Hz}$   
(break up/ background)



Measured position, angle@F5  
+  
optical correction



$P_{{}^3\text{He}}$  + reaction angle at target



Multi Wire Drift Chamber



# Ion optics (design)

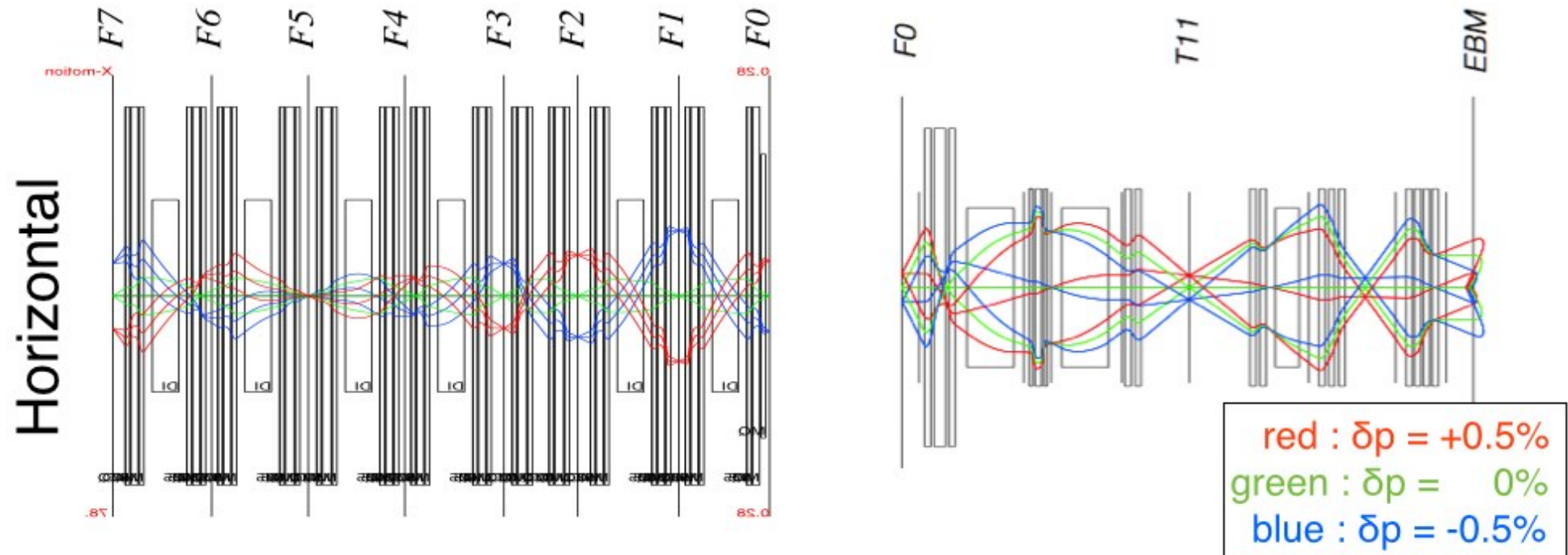
Dispersion matching:  
Eliminate contribution of beam momentum spread

$$\begin{array}{c}
 \begin{array}{|c|} \hline \text{Spectrometer} \\ \hline \text{(BigRIPS)} \\ \hline \end{array}
 \quad
 \begin{array}{|c|} \hline \text{reaction} \\ \hline \end{array}
 \quad
 \begin{array}{|c|} \hline \text{Analyzer} \\ \hline \text{(Beam Transfer Line)} \\ \hline \end{array}
 \\
 \left( \begin{array}{c} x_{fp} \\ \theta_{fp} \\ \delta p_{fp} \end{array} \right) = \begin{pmatrix} S_{11} & S_{12} & S_{16} \\ S_{21} & S_{22} & s_{26} \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & C \end{pmatrix} \begin{pmatrix} A_{11} & A_{12} & A_{16} \\ A_{21} & A_{22} & A_{26} \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x_0 \\ \theta_0 \\ \delta p_0 \end{pmatrix}
 \\
 \text{F5} \qquad \qquad \qquad *C: \text{ kinematical factor} = 1.31 \qquad \qquad \qquad \text{inside SRC}
 \end{array}$$

$$x_{fp} = \dots + \underbrace{(S_{11} A_{16} + C S_{16})}_{\text{red}} \delta p_0$$

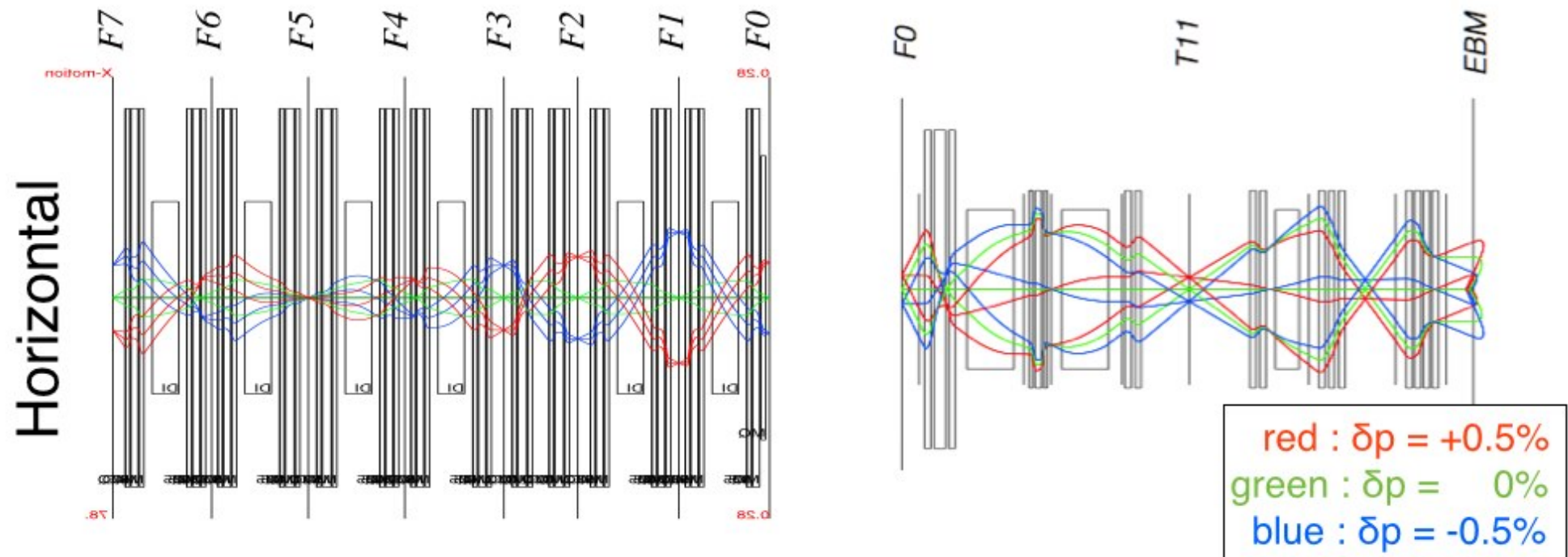


# Ion optics (design)



$$x_{fp} = \dots + \underbrace{(S_{11} A_{16} + C S_{16})}_{\text{red}} \delta p_0$$

# Ion optics (design)

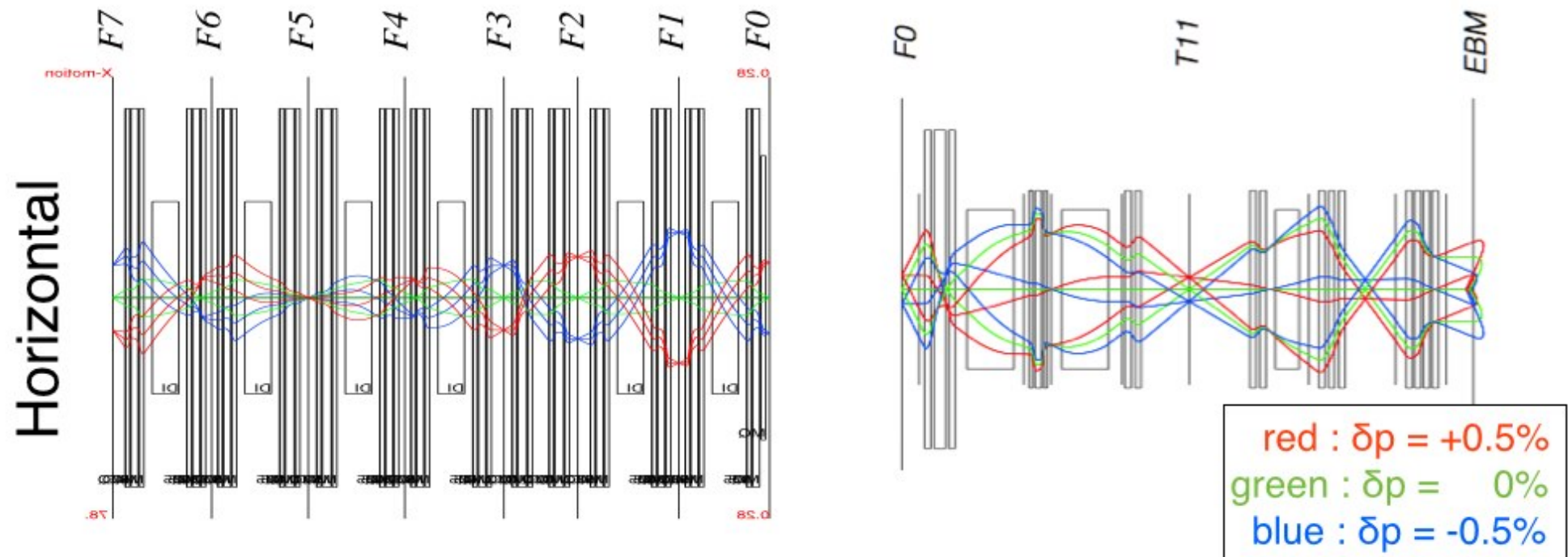


$$x_{fp} = \dots + \underbrace{(S_{11}A_{16} + CS_{16})}_{\text{magnification of BigRIPS}} \underbrace{\delta p_0}_{\text{position at F0}}$$

magnification of  
BigRIPS

position at F0

# Ion optics (design)



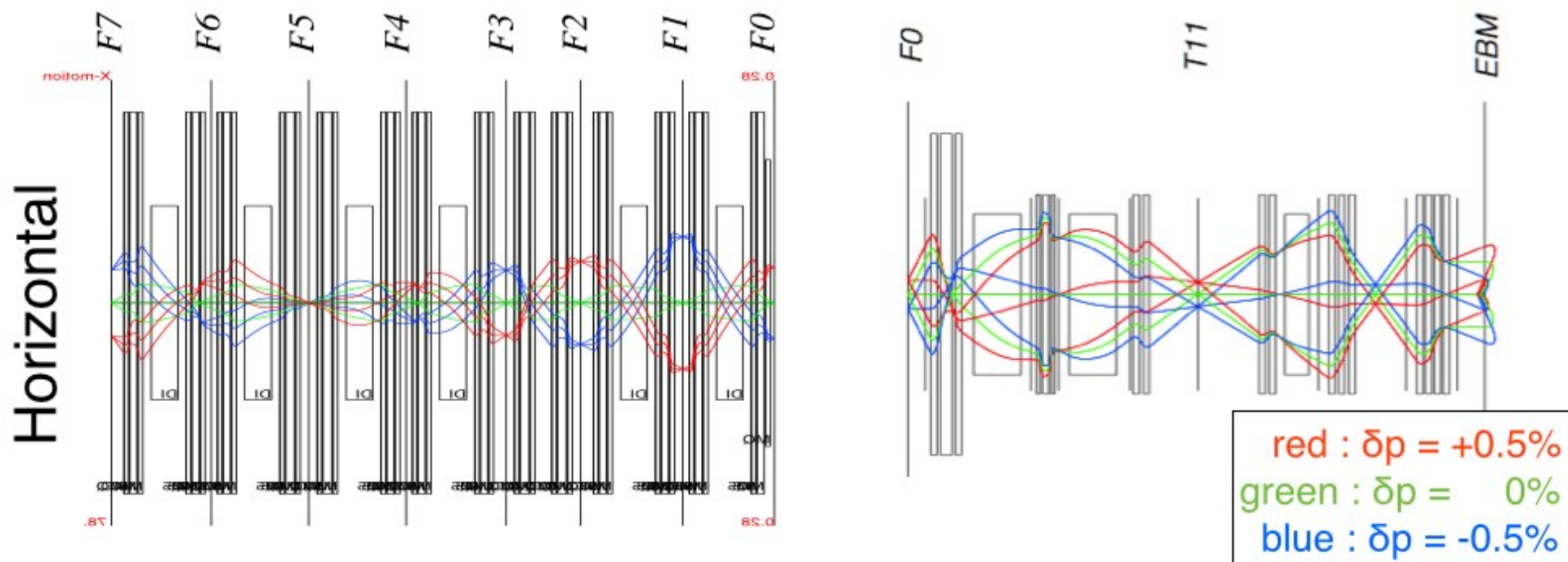
$$x_{fp} = \dots + \underbrace{(S_{11}A_{16} + CS_{16})}_{\text{dispersion of BigRIPS}} \delta p_0$$

$\delta$  at F0

dispersion of BigRIPS



# Ion optics (design)



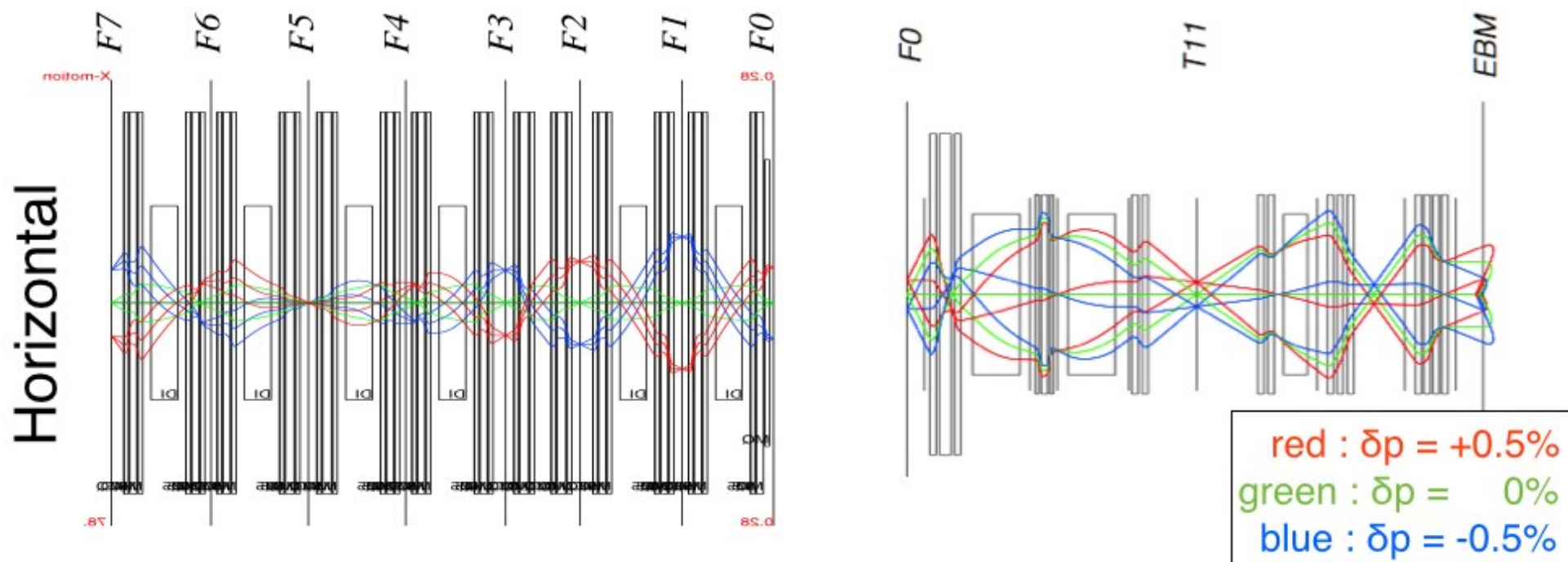
$$x_{fp} = \dots + \underbrace{(S_{11}A_{16} + CS_{16})}_{\text{Cancel out}} \delta p_0$$

← → Cancel out

same Q value  
different primary beam Energy → same position at F5



# Ion optics (design)

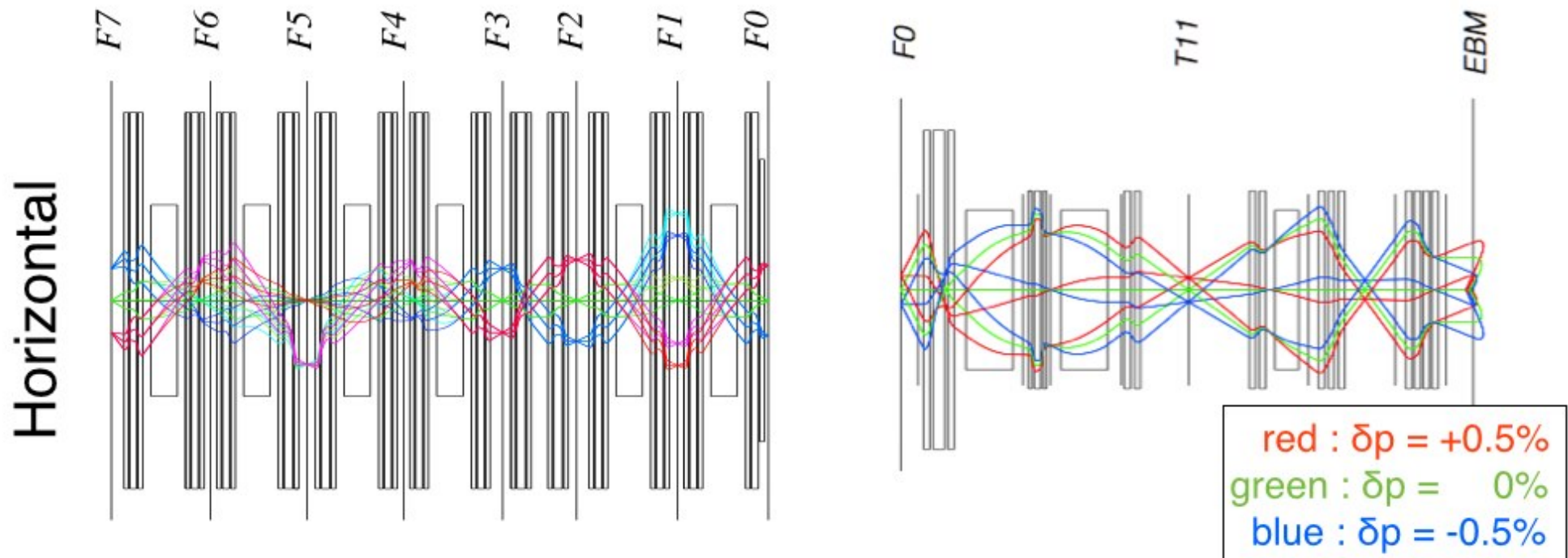


$$x_{fp} = \cdots + \underbrace{(S_{11} A_{16} + C S_{16})}_{\text{Cancel out}} \delta p_0$$

Cancel out

design values:  $S_{11} = -1.8$   $A_{16} = 44.6 \text{ mm}/\%$   
 $C = 1.31$   $S_{16} = 62 \text{ mm}/\%$

# Ion optics (design)

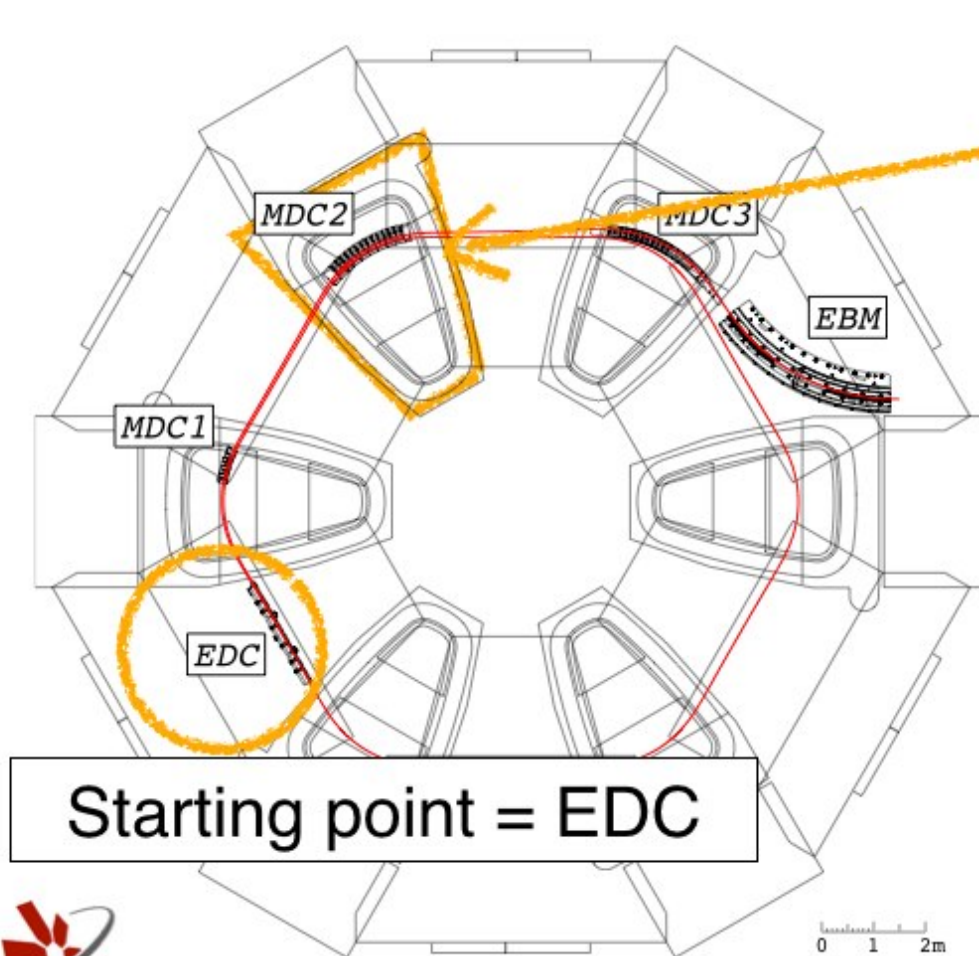


$$x_{fp} = \dots + \underbrace{(S_{11}A_{16} + CS_{16})}_{\text{red}} \underbrace{\delta p_0}_{\text{blue}}$$

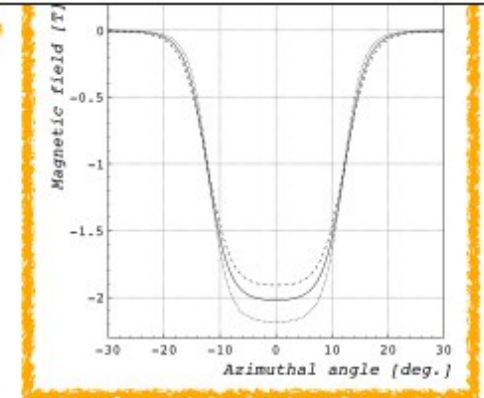


# Ion Optics inside SRC (calc.)

Dispersion matching using **primary beam**



magnetic field in the magnet



calculate the transfer matrix using Runge-Kutta method

$$\begin{pmatrix} (x|x) & (x|a) & (x|y) & (x|b) & (x|\delta) \\ (a|x) & (a|a) & (a|y) & (a|b) & (a|\delta) \\ (y|x) & (y|a) & (y|y) & (y|b) & (y|\delta) \\ (b|x) & (b|a) & (b|y) & (b|b) & (b|\delta) \end{pmatrix}_{\text{EDC} \rightarrow \text{EBM}} = \begin{pmatrix} -1.00 & -3.35 & 0.0 & 0.0 & 76.9 \\ 0.30 & -0.01 & 0.0 & 0.0 & -25.4 \\ 0.0 & 0.0 & -1.03 & -1.75 & 0.0 \\ 0.0 & 0.0 & -0.09 & -1.12 & 0.0 \end{pmatrix}$$

Starting point = EDC



SRC



# Key points of the experiment for better resolution

① Tuning of the beam transfer line optics  
→ realize dispersion matching

② Reduce the primary beam emittance /  $\delta p$

①  $\times$  ②  $\Rightarrow$  contribution of  $\delta p$  on resolution



# ① Tuning of the beam transfer line optics

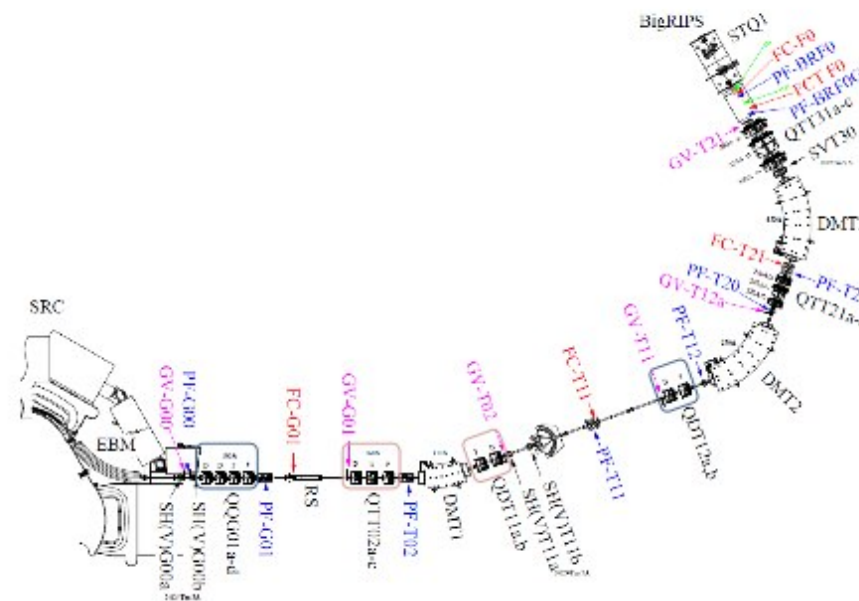
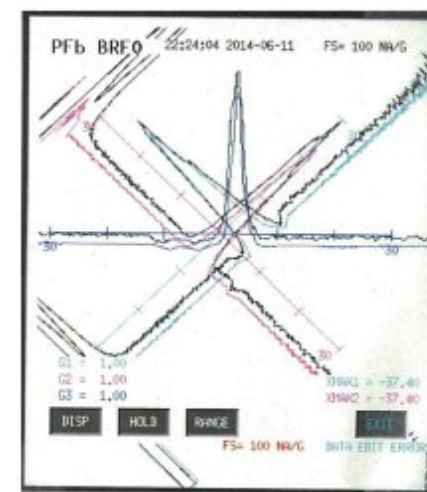
beam envelope measured by profile monitor

## Diagnostics in BT line

- ~ 10 profile monitors (wires)
- **No tracker**



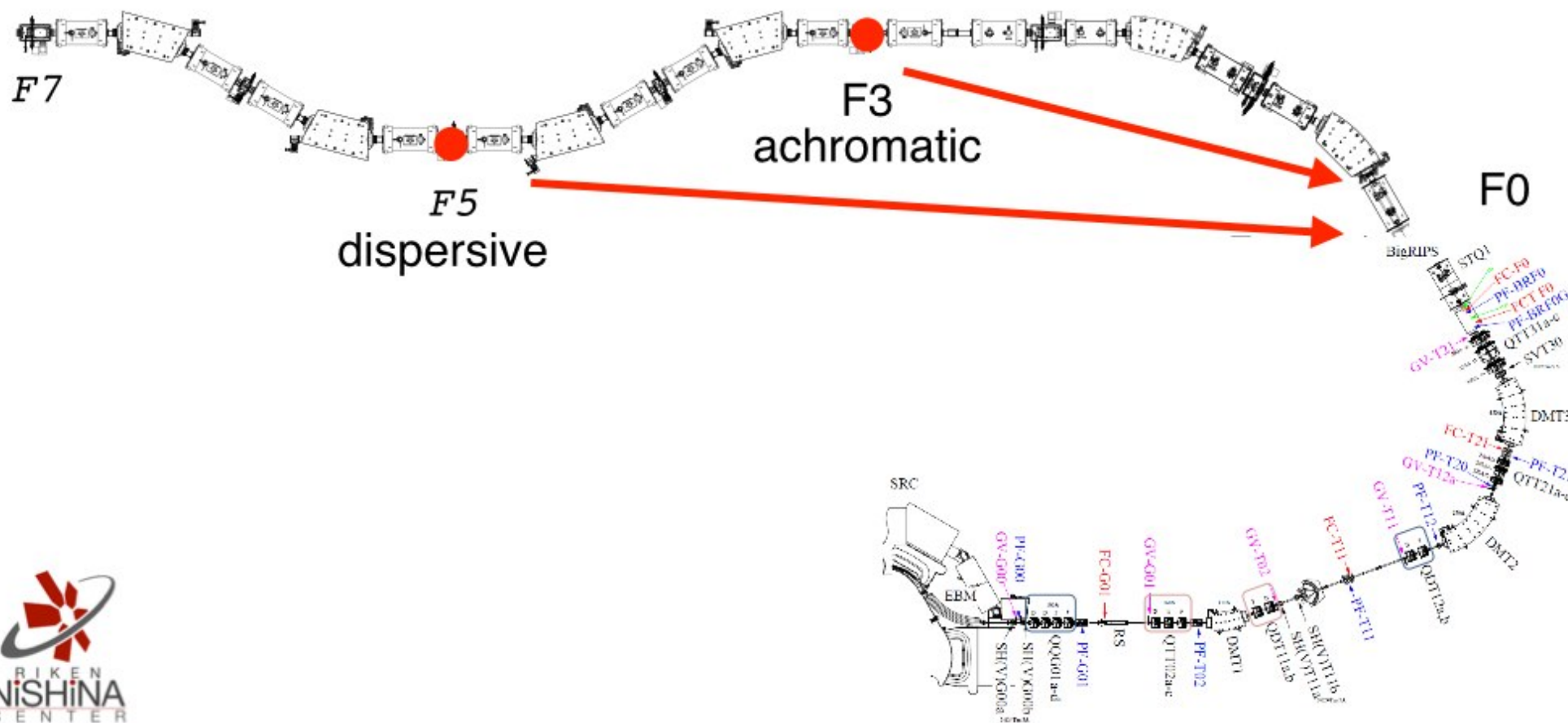
not enough to measure dispersion / other transfer matrix



# ① Tuning of the beam transfer line optics

Use BigRIPS as a diagnostics of BT line

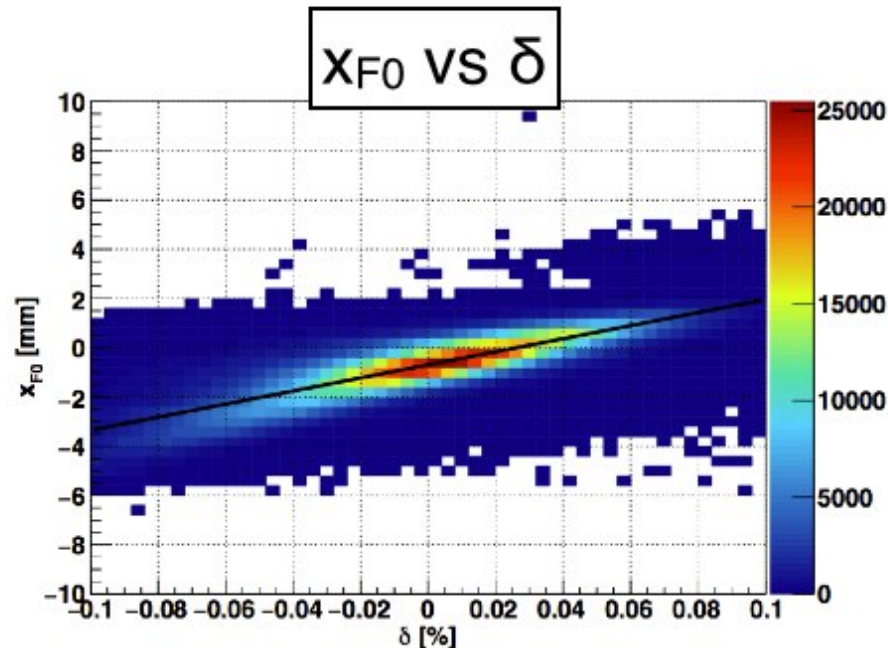
- measure the position / angle at F3 and F5 focal plane
- trace back to F0



# ① Tuning of the beam transfer line optics

Use BigRIPS as a diagnostics of BT line

- measure the position / angle at F3 and F5 focal plane
- trace back to F0

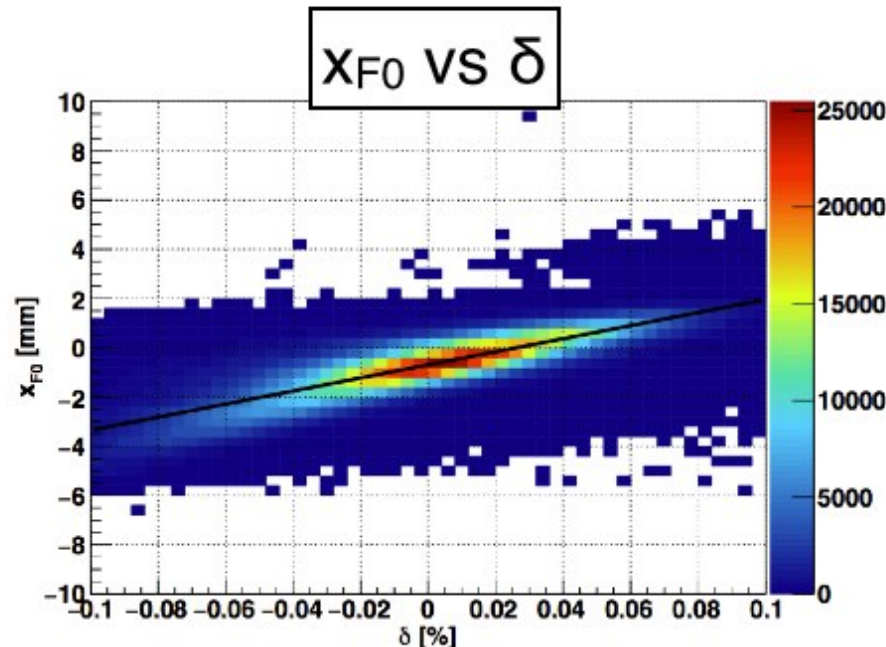


measured dispersion in BT line

# ① Tuning of the beam transfer line optics

Use BigRIPS as a diagnostics of BT line

- measure the position / angle at F3 and F5 focal plane
- trace back to F0



measured dispersion in BT line

→ 24 mm/% !!!

( cf. designed value: 44.6 mm/%)

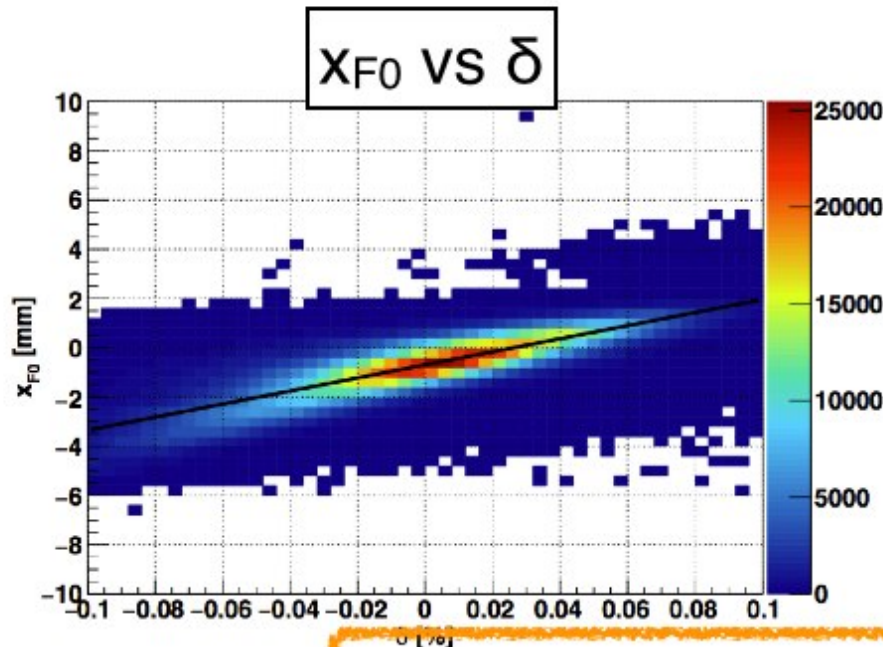
The difference may be caused by the optics inside SRC



# ① Tuning of the beam transfer line optics

Use BigRIPS as a diagnostics of BT line

- measure the position / angle at F3 and F5 focal plane
- trace back to F0



measured dispersion in BT line

→ **24 mm/% !!!**

( cf. designed value: 44.6 mm/%)

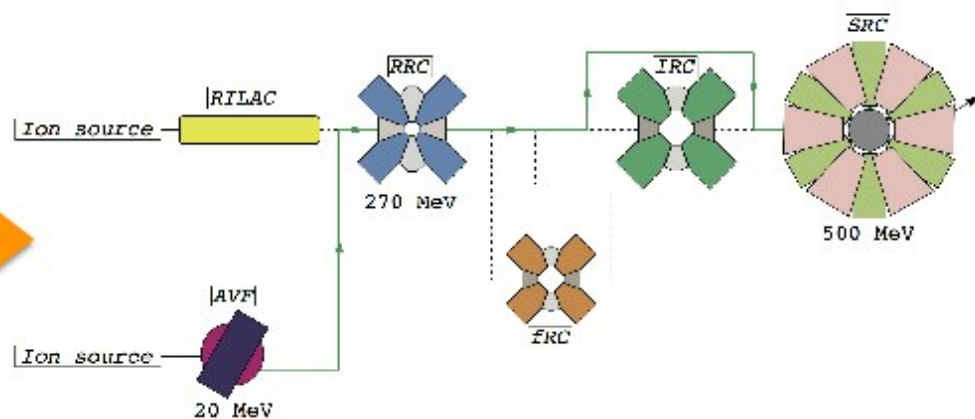
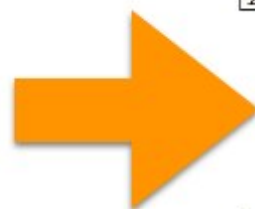
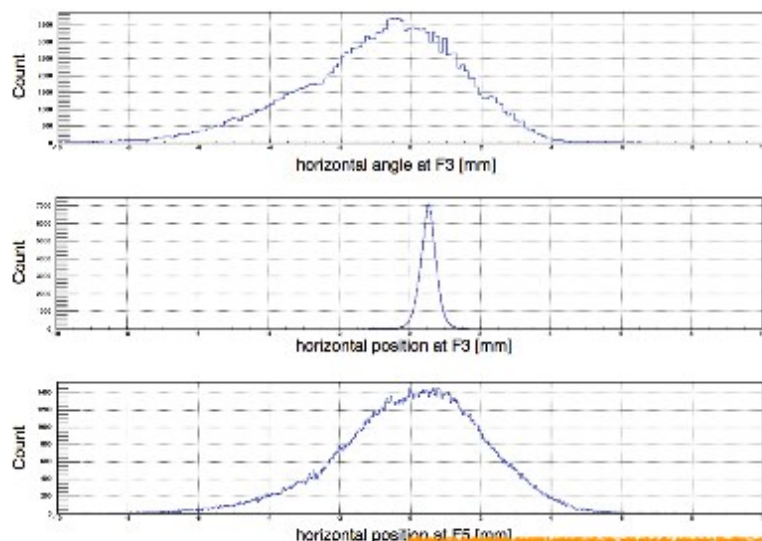
The difference may be caused by the optics inside SRC

After phenomenological tuning,  
dispersion is improved to be **28 mm/%**  
⇔ reduce the contribution of  $\delta p$  by **~ 50%**

## ② Reduce the primary beam emittance / $\delta p$

Use BigRIPS as a diagnostics of Primary beam

- position and angle distribution at achromatic focal plane (F3)  
→ beam emittance
- position distribution at dispersive focal plane (F5)  
→  $\delta p$



feedback to the accelerator team  
to tune accelerator parameters / slit condition  
achieve  **$0.2 \times 2.0 \text{ mm} \cdot \text{mrad} / 0.03 \% \text{ (RMS)}$**

## Contribution of $\delta p$ on resolution

- ① Tuning of the beam transfer line optics  
→ dispersion **28 mm/%**  $\Leftrightarrow$  reduce the contribution by  **$\sim 50\%$**
- ② Reduce the primary beam emittance /  $\delta p$   
→ achieve  **$0.2 \times 2.0 \text{ mm} \cdot \text{mrad} / 0.03 \%$  (RMS)**

$$\text{①} \times \text{②} \Rightarrow \sim \mathbf{220 \text{ keV (FWHM)}}$$

acceptable compared with  $\Gamma_{1s} \sim 300 \text{ keV}$

※ With other contributions, resolution is estimated to be 250  $\sim$  400 keV (position dependent)

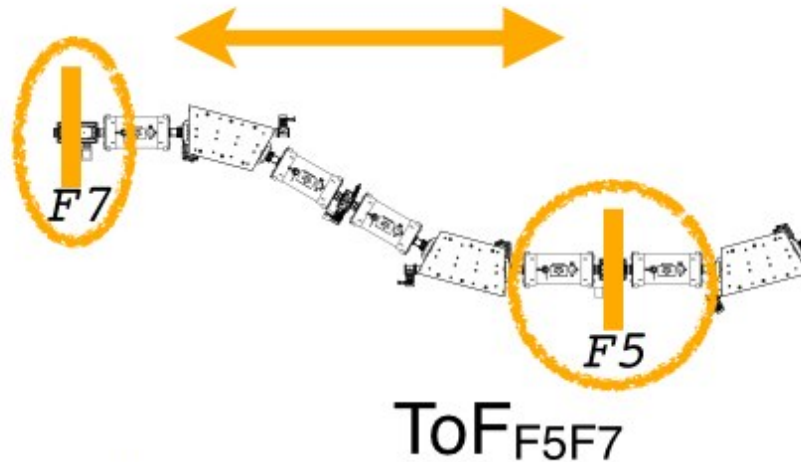


### (iii) Data analysis (pID / optics / decomposition)

- Particle identification
- Optics analysis
- Decomposition of Eex spectra



# Particle identification



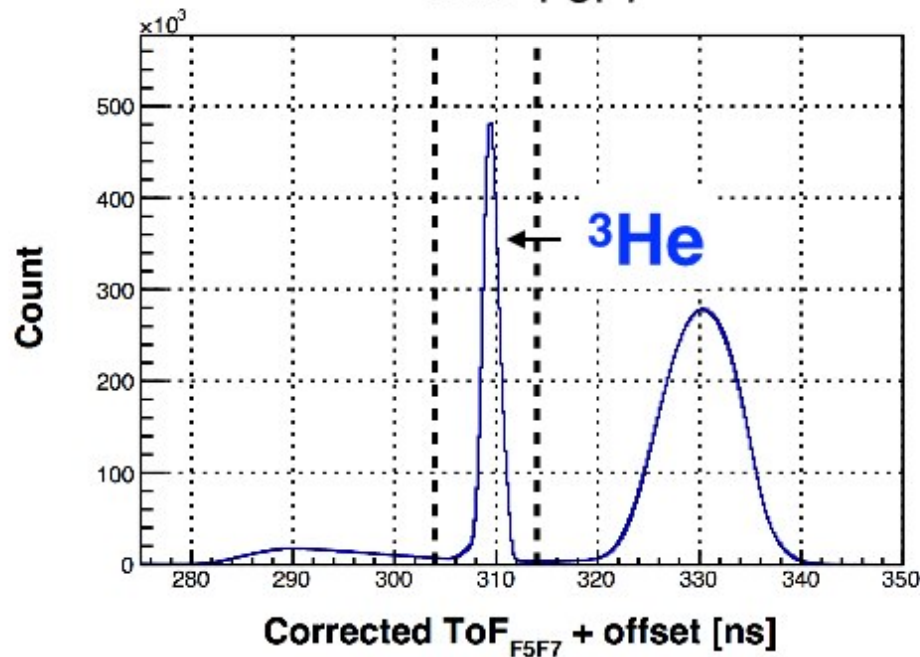
<sup>3</sup>He is well separated only with ToF.  
Using additional information of

- Eloss in MWDC (as ToT)
- relative timing with RF of SRC

We achieve

efficiency :  $\geq 98 \%$

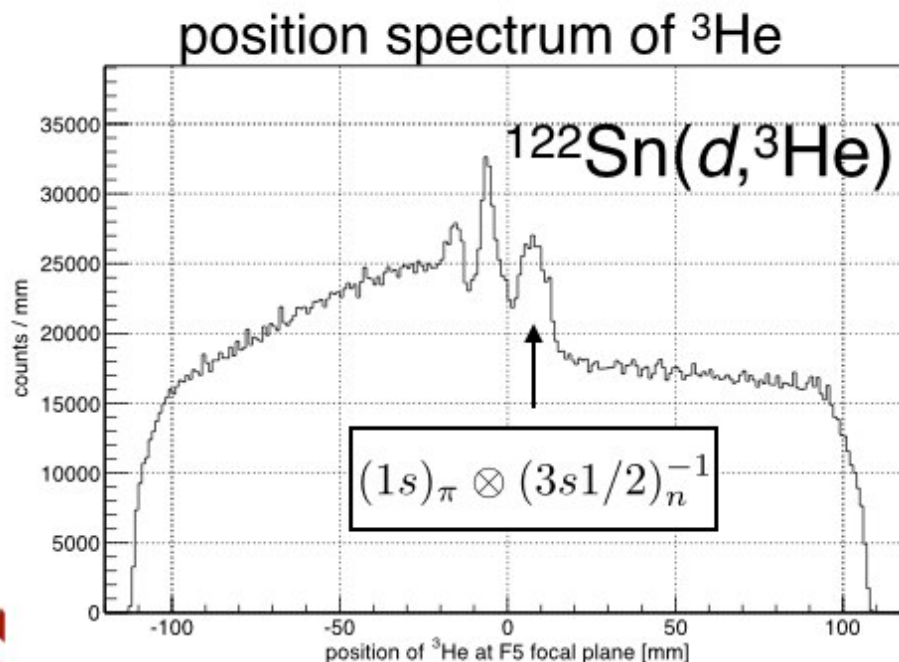
purity :  $\geq 99.5 \%$



# Optics analysis

To convert the position spectrum to Eex spectra, we should

- ① measure the **dispersion**
- ② correct the **optical higher order aberration**



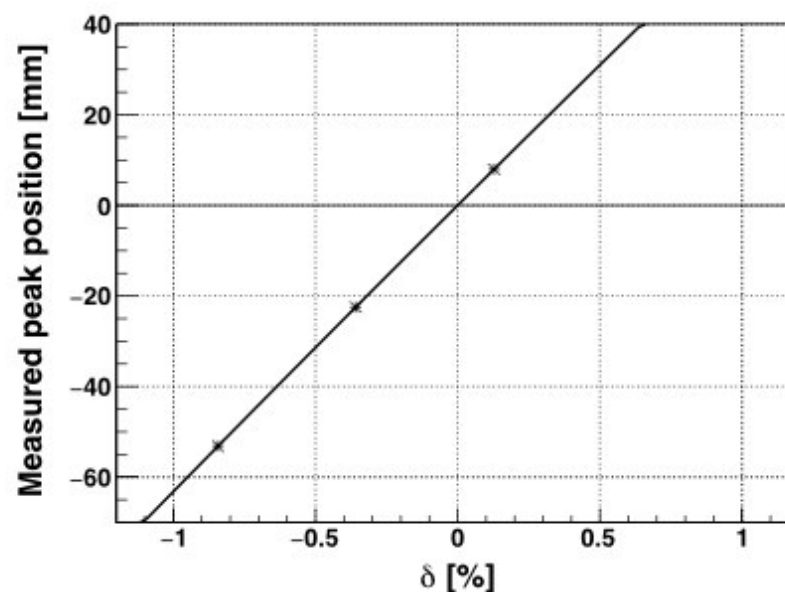
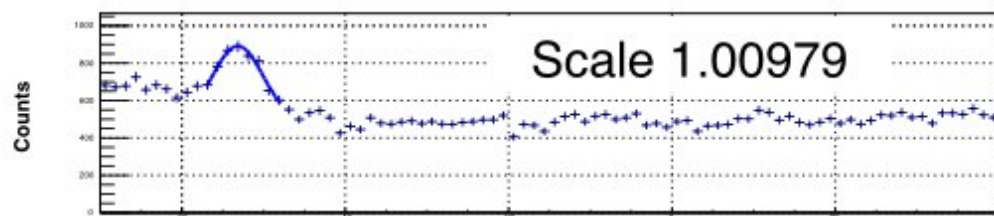
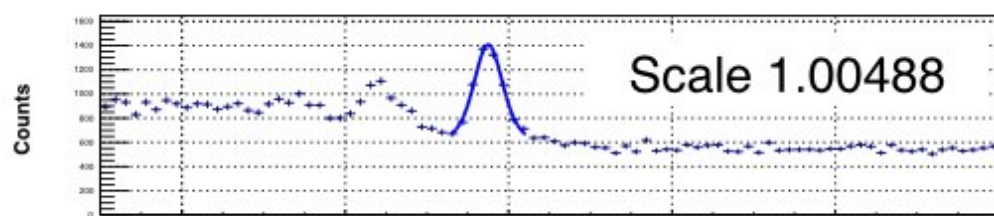
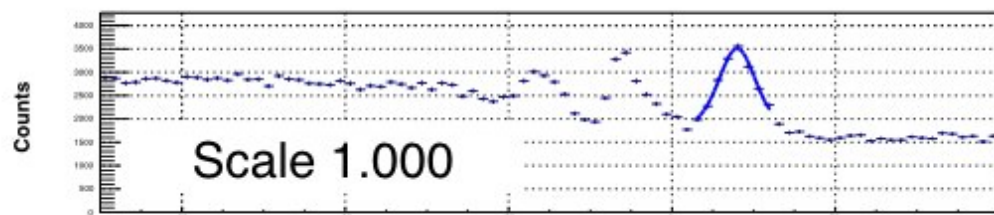
To analyze the optics,  
we used the peak of  $(1s)_\pi \otimes (3s1/2)_n^{-1}$

- sufficient strength in forward  $\theta$
- no angle dependence of Eex

# ① Dispersion of BigRIPS

$$\delta = (P_{3He}^{1s} - P_{\text{BigRIPS}}) / P_{\text{BigRIPS}} * 100.$$

↑ ↑  
 scale BigRIPS by 1.000 (default), 1.00488, 1.00979

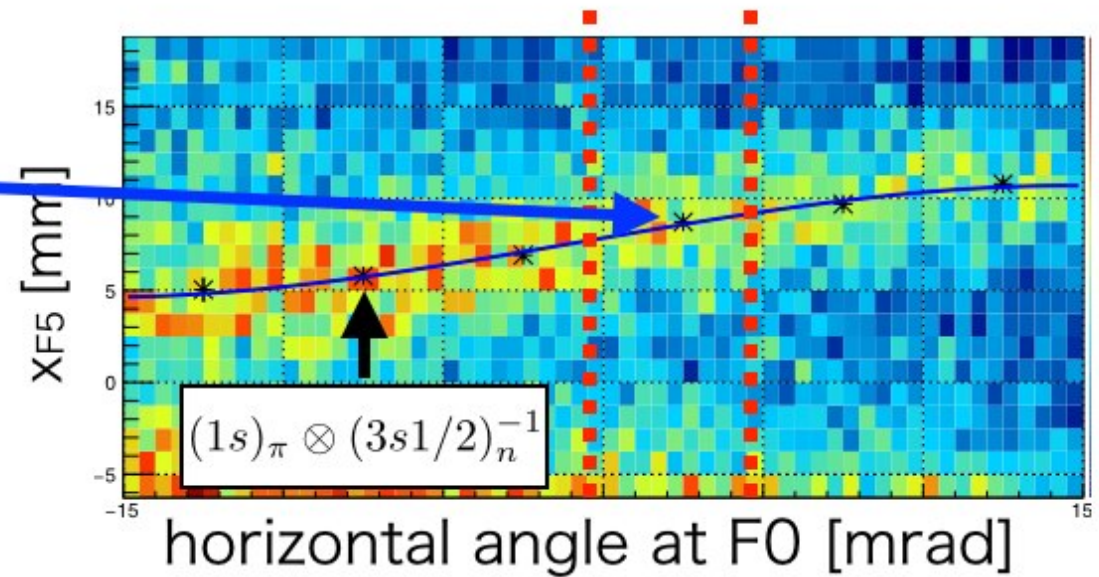
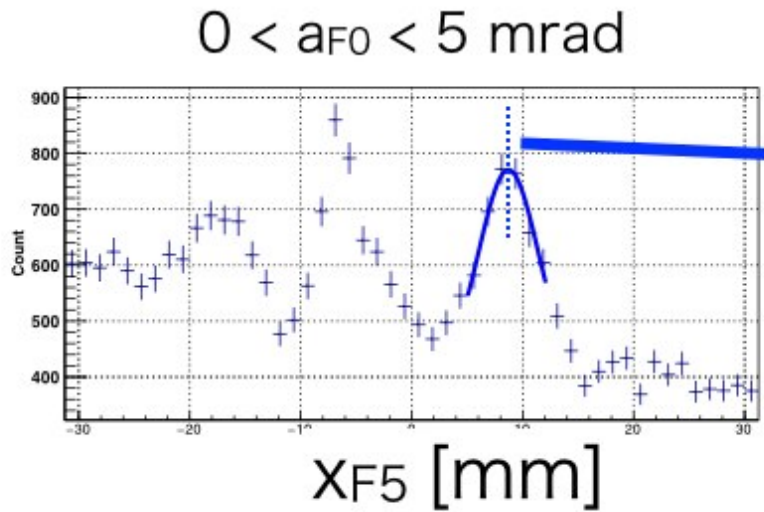


$$X_{F5} = \alpha \delta + \beta \delta^2$$

Dispersion

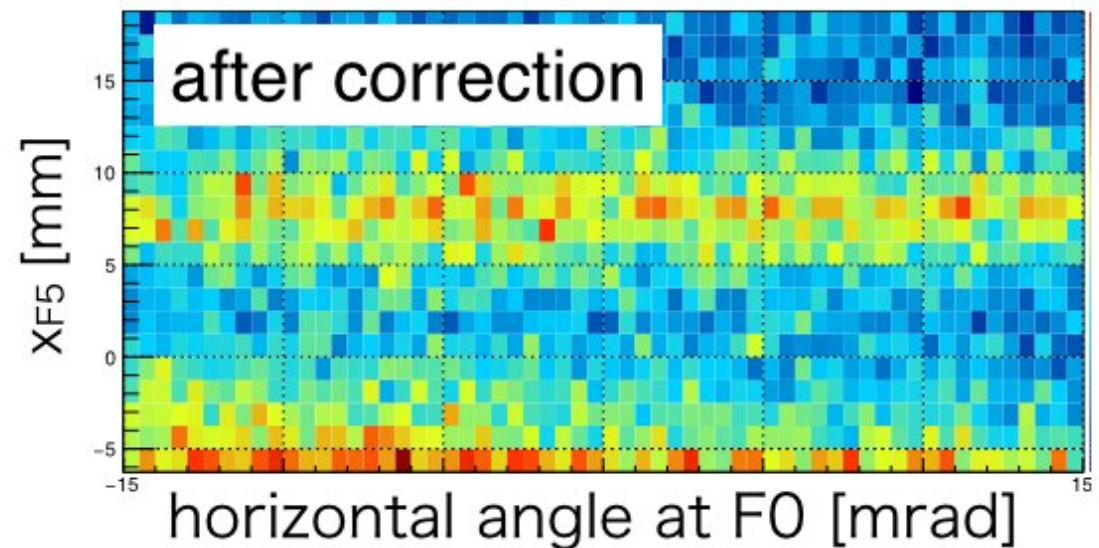


## ② Correct optical higher order aberration



vertical angular aberration is also corrected

\*data for the evaluation separated from the production run

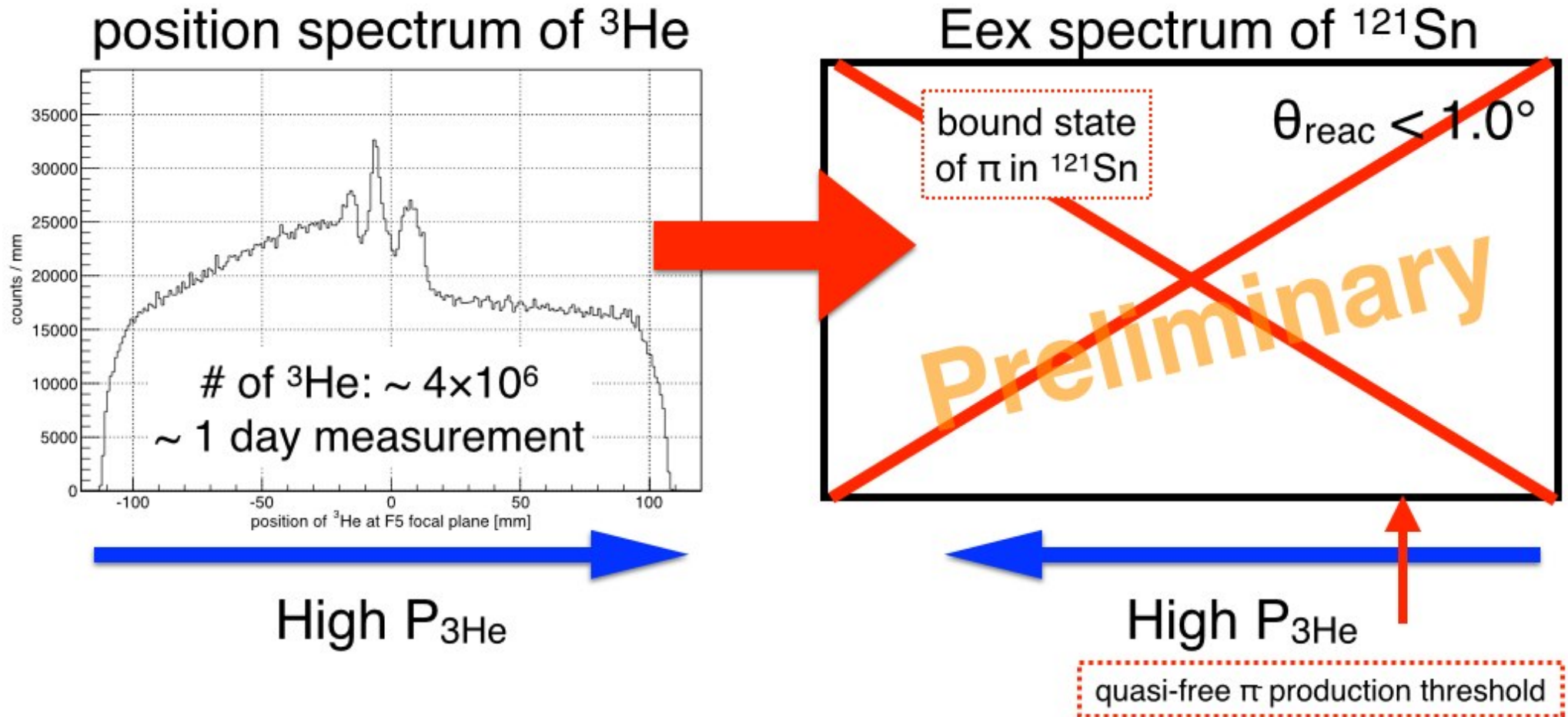




Finally, we include 15 elements of transfer matrix for the analysis.

matrix element	deduced values from exp. data	calculated values
$(x_{F5} a_{F0})$ [mm/mrad]	$(+2.75 \pm 0.2) \times 10^{-1}$	$+1.93 \times 10^{-1}$
$(x_{F5} a_{F0}^2)$ [mm/mrad <sup>2</sup> ]	$(-9.44 \pm 8.0) \times 10^{-4}$	$-1.37 \times 10^{-3}$
$(x_{F5} a_{F0}^3)$ [mm/mrad <sup>3</sup> ]	$(-4.41 \pm 0.9) \times 10^{-4}$	$-7.10 \times 10^{-4}$
$(x_{F5} b_{F0})$ [mm/mrad]	$(+4.87 \pm 6.1) \times 10^{-3}$	0.00
$(x_{F5} b_{F0}^2)$ [mm/mrad <sup>2</sup> ]	$(-8.45 \pm 4.0) \times 10^{-4}$	$-2.15 \times 10^{-4}$
$(x_{F5} a_{F0}b_{F0}^2)$ [mm/mrad <sup>3</sup> ]	$(-2.53 \pm 0.5) \times 10^{-4}$	$-3.01 \times 10^{-4}$
$(x_{F5} a_{F0}\delta)$ [mm/mrad/%]	$(+7.83 \pm 0.2) \times 10^{-1}$	$8.27 \times 10^{-1}$
$(x_{F5} a_{F0}^2\delta)$ [mm/mrad <sup>2</sup> /%]	$(+6.74 \pm 1.9) \times 10^{-3}$	$+1.28 \times 10^{-3}$
$(x_{F5} \delta)$ [mm/%]	$+62.42 \pm 0.27$	+61.20
$(x_{F5} \delta^2)$ [mm/% <sup>2</sup> ]	$(-7.84 \pm 4.7) \times 10^{-1}$	$-8.34 \times 10^{-1}$
$(a_{F5} a_{F0})$ [mrad/mrad]	$(-5.19 \pm 0.5) \times 10^{-1}$	$-5.48 \times 10^{-1}$
$(b_{F5} b_{F0})$ [mrad/mrad]	$(-5.40 \pm 0.5) \times 10^{-1}$	$-6.39 \times 10^{-1}$
$(a_{F5} a_{F0}\delta)$ [mrad/mrad/%]	-	$-3.72 \times 10^{-2}$
$(b_{F5} b_{F0}\delta)$ [mrad/mrad/%]	-	$-6.79 \times 10^{-2}$
$(b_{F5} b_{F0}\delta^2)$ [mrad/mrad/% <sup>2</sup> ]	-	$+5.29 \times 10^{-2}$

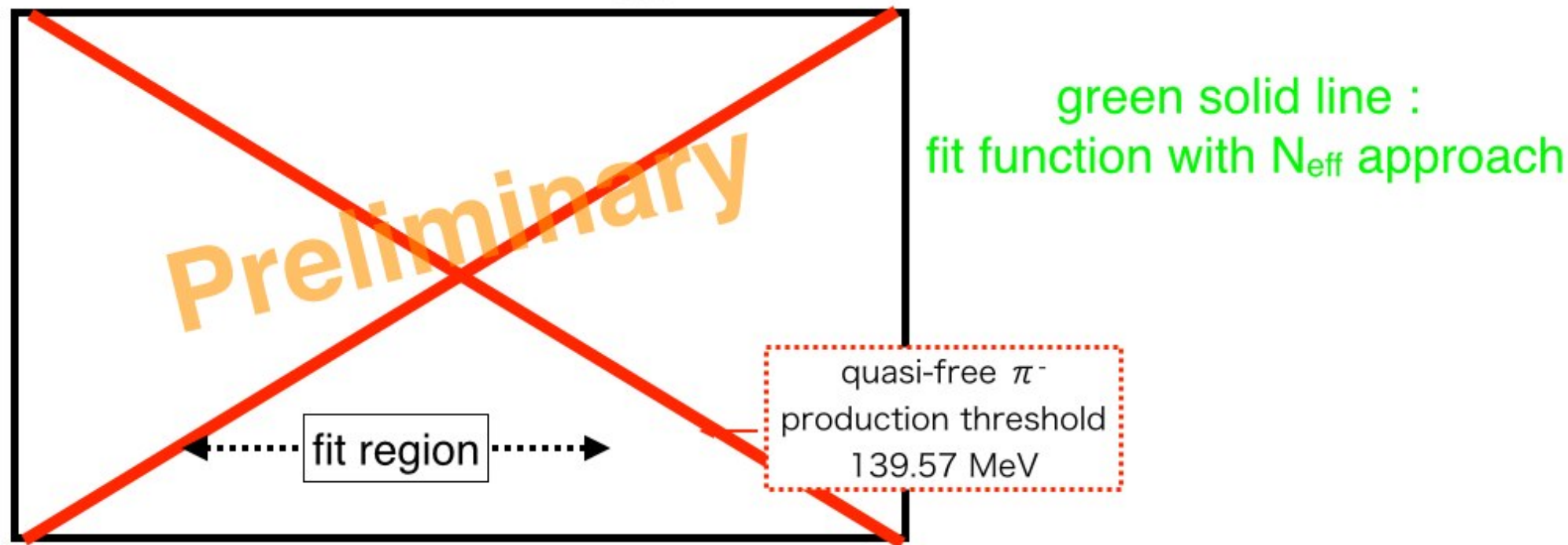
## Conversion from position to Eex



The spectrum seems to achieve the best resolution among the past deeply-bound pionic atom experiment.

## Decomposition of the $E_{ex}$ spectrum : $^{122}\text{Sn}$ target

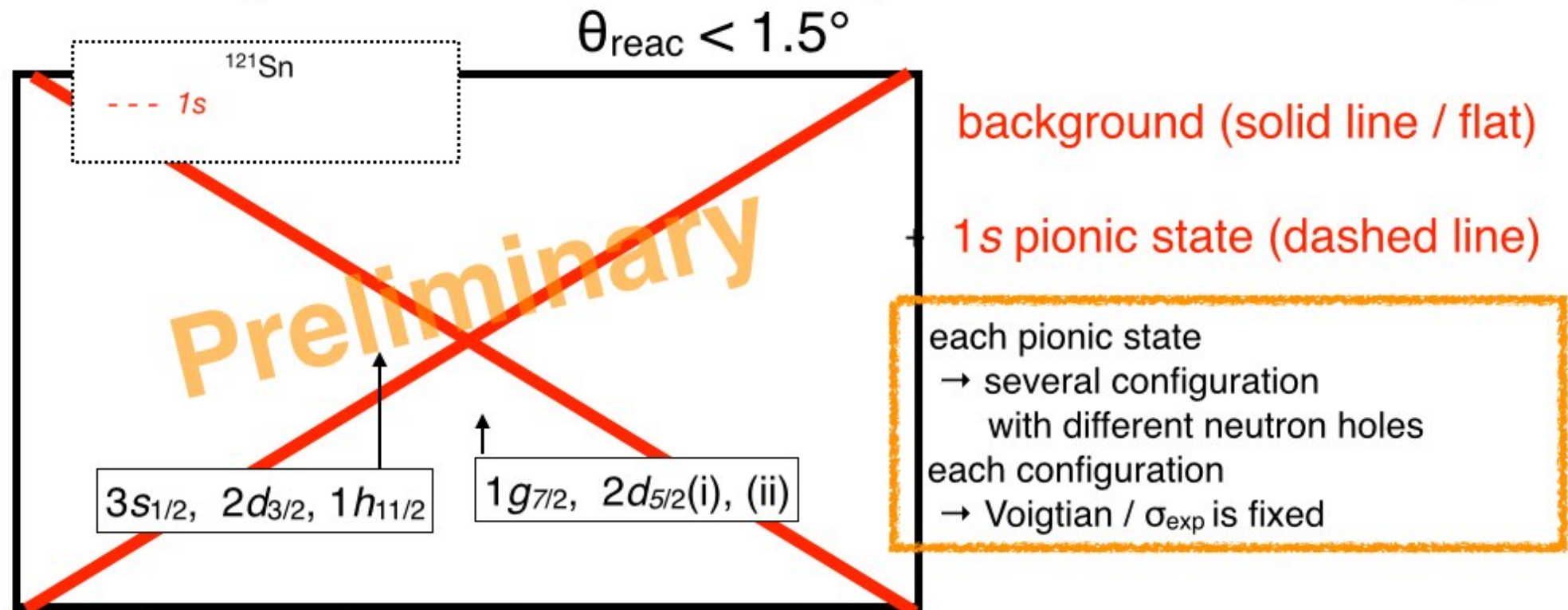
$$\theta_{\text{reac}} < 1.5^\circ$$



The  $E_{ex}$  spectrum is fit by the function with several components  
→ deduce binding energies and widths of pionic states



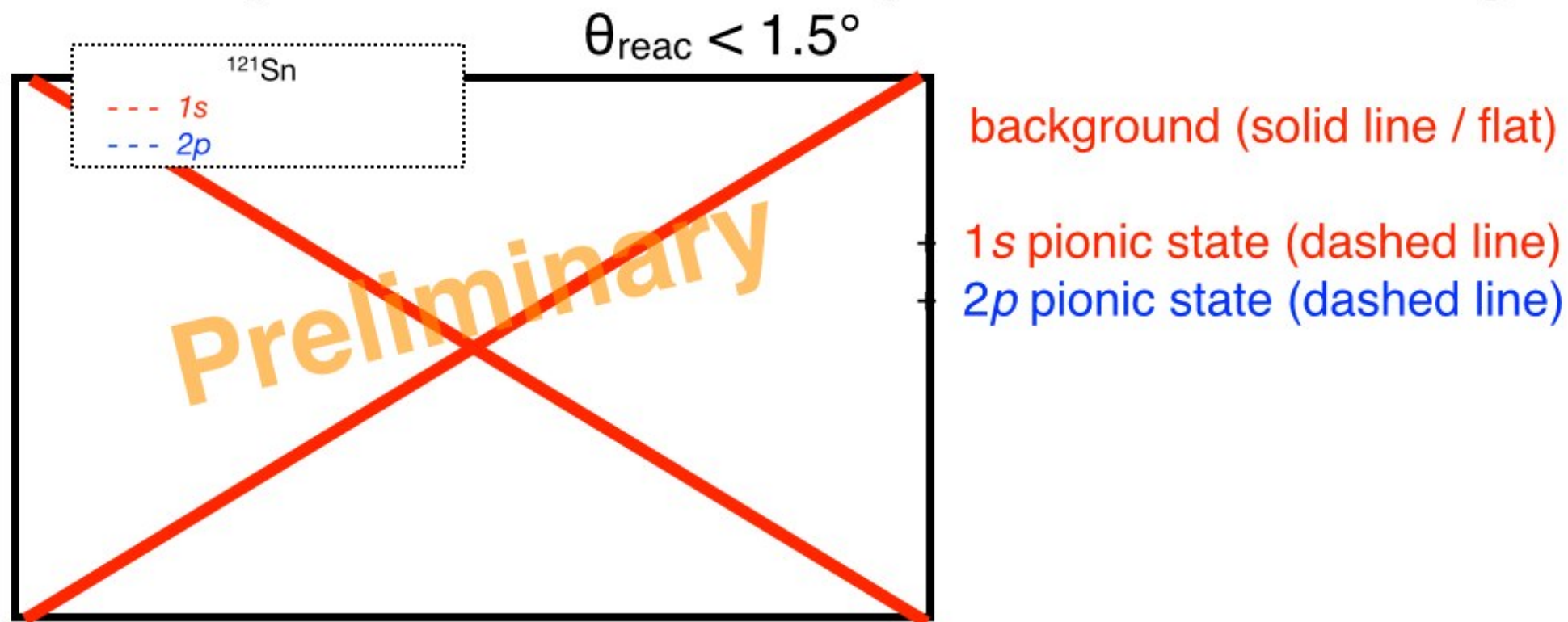
# Decomposition of the $E_{ex}$ spectrum : $^{122}\text{Sn}$ target



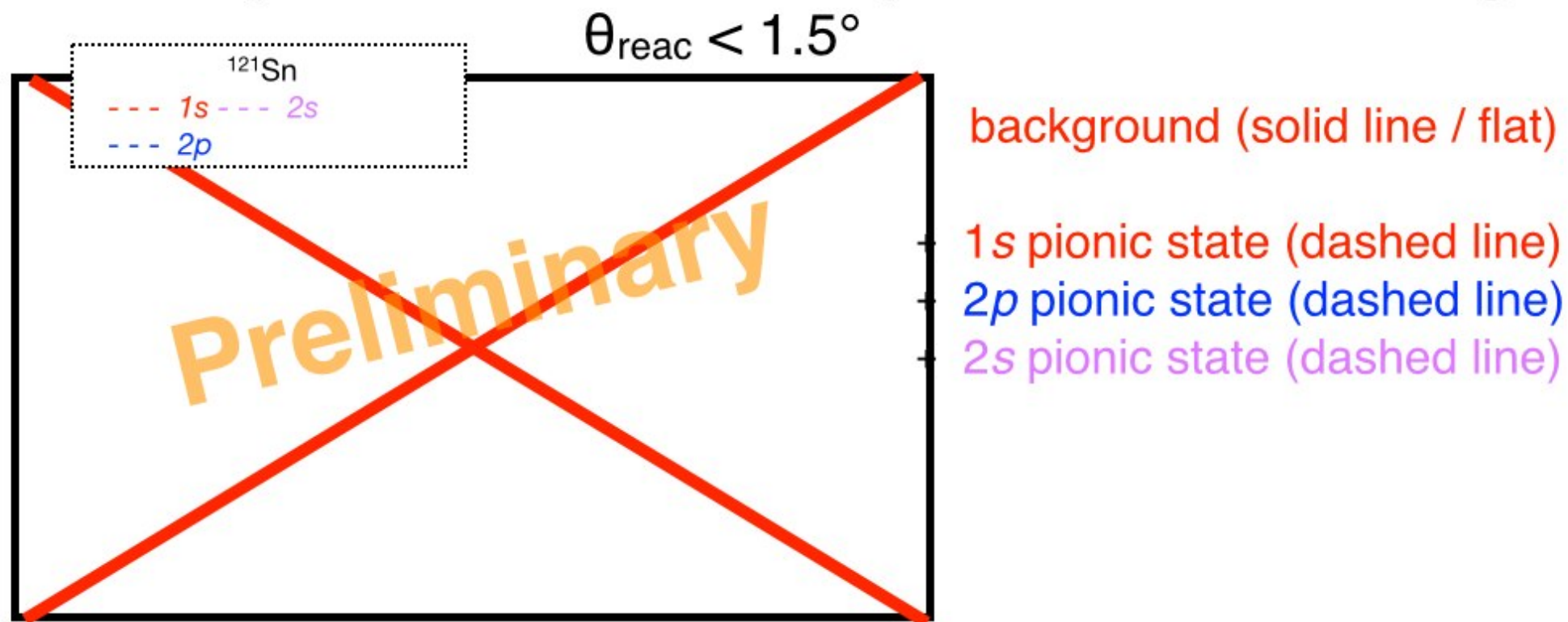
neutron hole	$E_{ex}$ [MeV]	relative strength for pionic 1s state
$2d_{3/2}$	0.000	0.09
$1h_{11/2}$	0.006	0.001
$3s_{1/2}$	0.060	<b>1</b>
$1g_{7/2}$	0.926	0.003
$2d_{5/2}$ (i)	1.121	0.12
$2d_{5/2}$ (ii)	1.403	0.06



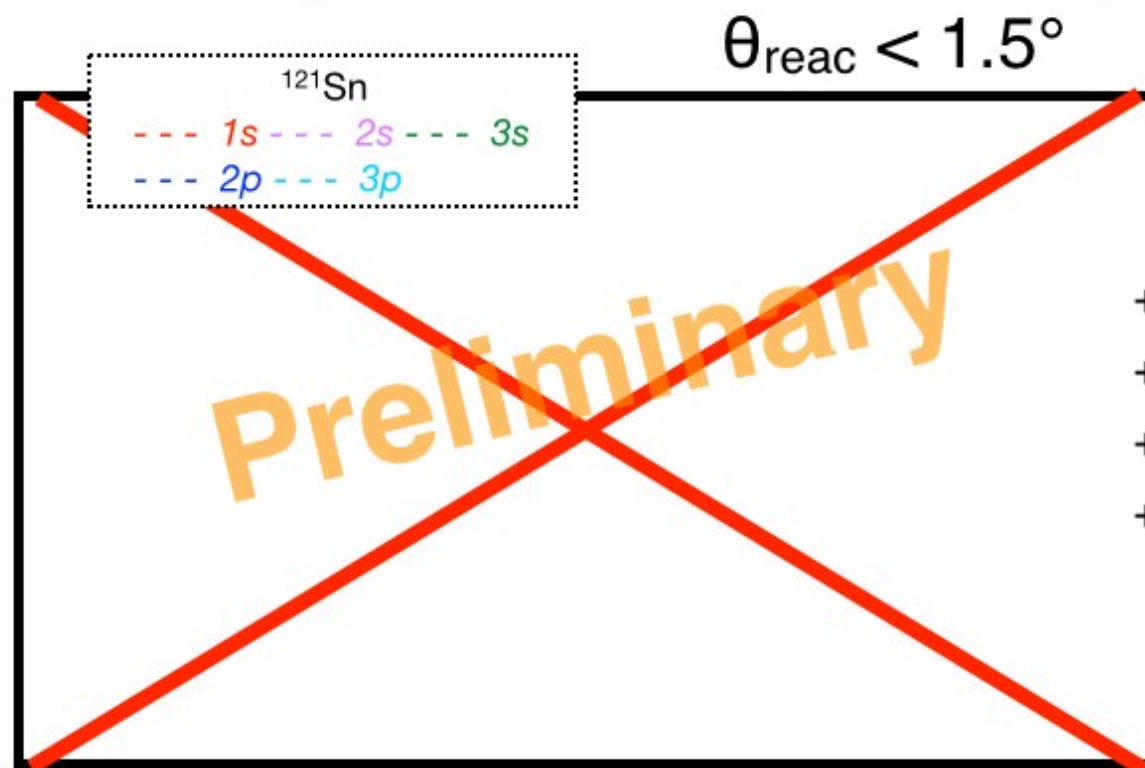
## Decomposition of the $E_{ex}$ spectrum : $^{122}\text{Sn}$ target



## Decomposition of the $E_{ex}$ spectrum : $^{122}\text{Sn}$ target



## Decomposition of the $E_{ex}$ spectrum : $^{122}\text{Sn}$ target



background (solid line / flat)

1s pionic state (dashed line)

2p pionic state (dashed line)

2s pionic state (dashed line)

3p, 3s state (dashed line)

### Fitting parameter

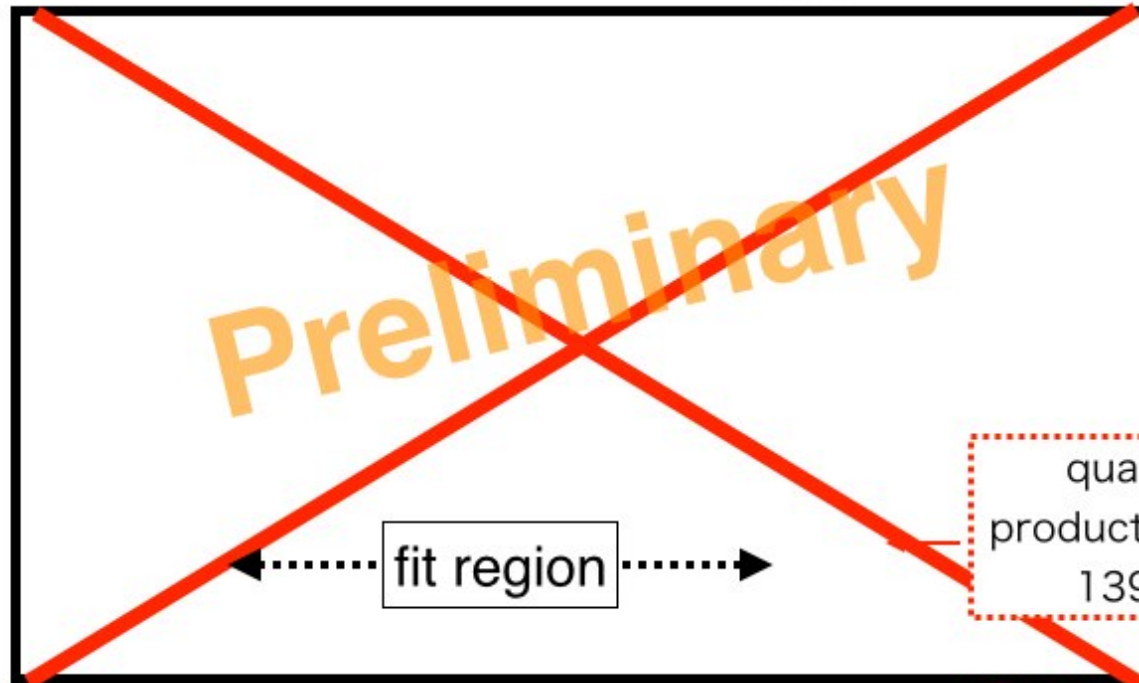
- relative strength of each state
- $BE_{1s}$ ,  $BE_{2p}$ ,  $BE_{2s}$
- $\Gamma_{1s}$ ,  $\Gamma_{2p}$

### Fixed parameter

- $BE_{3p}$ ,  $BE_{3s}$
- $\Gamma_{2s}$ ,  $\Gamma_{3p}$ ,  $\Gamma_{3s}$

# Decomposition of the $E_{ex}$ spectrum : $^{122}\text{Sn}$ target

$$\theta_{\text{reac}} < 1.5^\circ$$



2p state → good calibration peak

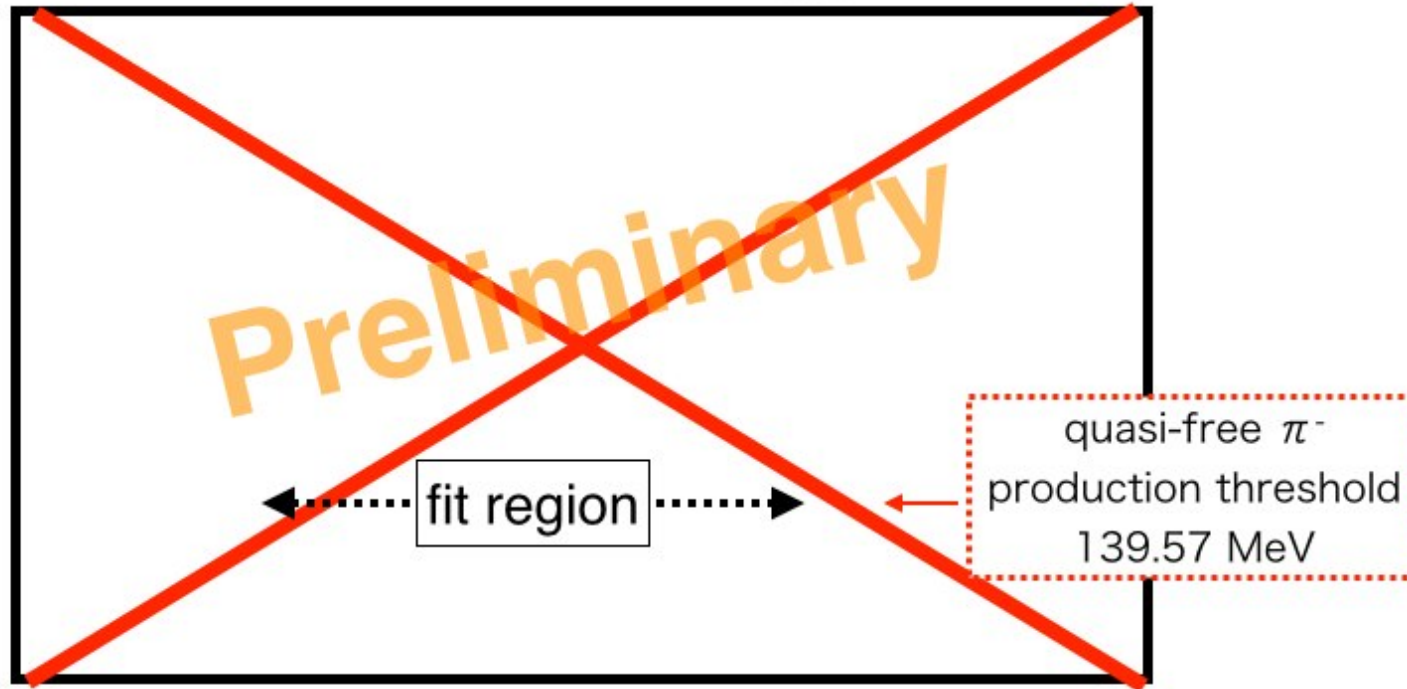
Deduced  $B_{1s}$ ,  $\Gamma_{1s}$ ,  $B_{2p}$  in  $^{121}\text{Sn}$

$(nl)$	$B_{nl}$ [MeV]	$\Gamma_{nl}$ [MeV]
1s	XXX ± XXX (stat.) ± XXX (sys.)	XXX ± XXX (stat.) ± XXX (sys.)
2p	XXX ± XXX (stat.) ± XXX (sys.)	XXX ± XXX (stat.) ± XXX (sys.)
1s - 2p	<b>XXX ± XXX (stat.) ± XXX (sys.)</b>	—



# Decomposition of the $E_{ex}$ spectrum : $^{117}\text{Sn}$ target

$$\theta_{\text{reac}} < 1.5^\circ$$



## Deduced $B_{1s}$ in $^{117}\text{Sn}$

$(nl)$	$B_{nl}$ [MeV]	$\Gamma_{nl}$ [MeV]
1s	$XXX \pm XXX$ (stat.) $\pm XXX$ (sys.)	—

# Optical potential deduced from $B_{1s}$ , $\Gamma_{1s}$ , $B_{2p}$ in $^{121}\text{Sn}$

$$\ast \Delta B = B_{1s} - B_{2p}$$



Calculate likelihood for each  $b_1$  and  $\text{Im}B_0$   
 $\rightarrow$  obtain best value

$$V_s(r) = -\frac{2\pi}{\mu} [\epsilon_1 \{b_0 \rho(r) + b_1 \delta \rho(r)\} + \epsilon_2 B_0 \rho(r)^2].$$

$\ast b_0$ ,  $\text{Re}B_0$  are deduced from data of light / symmetric pionic atoms

	$b_1$ [ $m_\pi^{-1}$ ]	$\text{Im}B_0$ [ $m_\pi^{-4}$ ]
This exp.	<b>XXX <math>\pm</math> XXX</b>	XXX $\pm$ XXX
GSI	- 0.115 $\pm$ 0.007	0.0472 $\pm$ 0.0013
In vacuum	- 0.0868 $\pm$ 0.0014	-

# Optical potential deduced from $B_{1s}$ , $\Gamma_{1s}$ , $B_{2p}$ in $^{121}\text{Sn}$

$$\ast \Delta B = B_{1s} - B_{2p}$$



Calculate likelihood for each  $b_1$  and  $\text{Im}B_0$   
 → obtain best value

$$V_s(r) = -\frac{2\pi}{\mu} [\epsilon_1 \{b_0 \rho(r) + b_1 \delta \rho(r)\} + \epsilon_2 B_0 \rho(r)^2].$$

$\ast b_0, \text{Re}B_0$  are deduced from

most precise evaluation the value of  $b_1$  in medium  
 further evidence that chiral symmetry breaking is partially restored at finite densities.

This exp.	<b>XXX ± XXX</b>	XXX ± XXX
GSI	- 0.115 ± 0.007	0.0472 ± 0.0013
In vacuum	- 0.0868 ± 0.0014	—

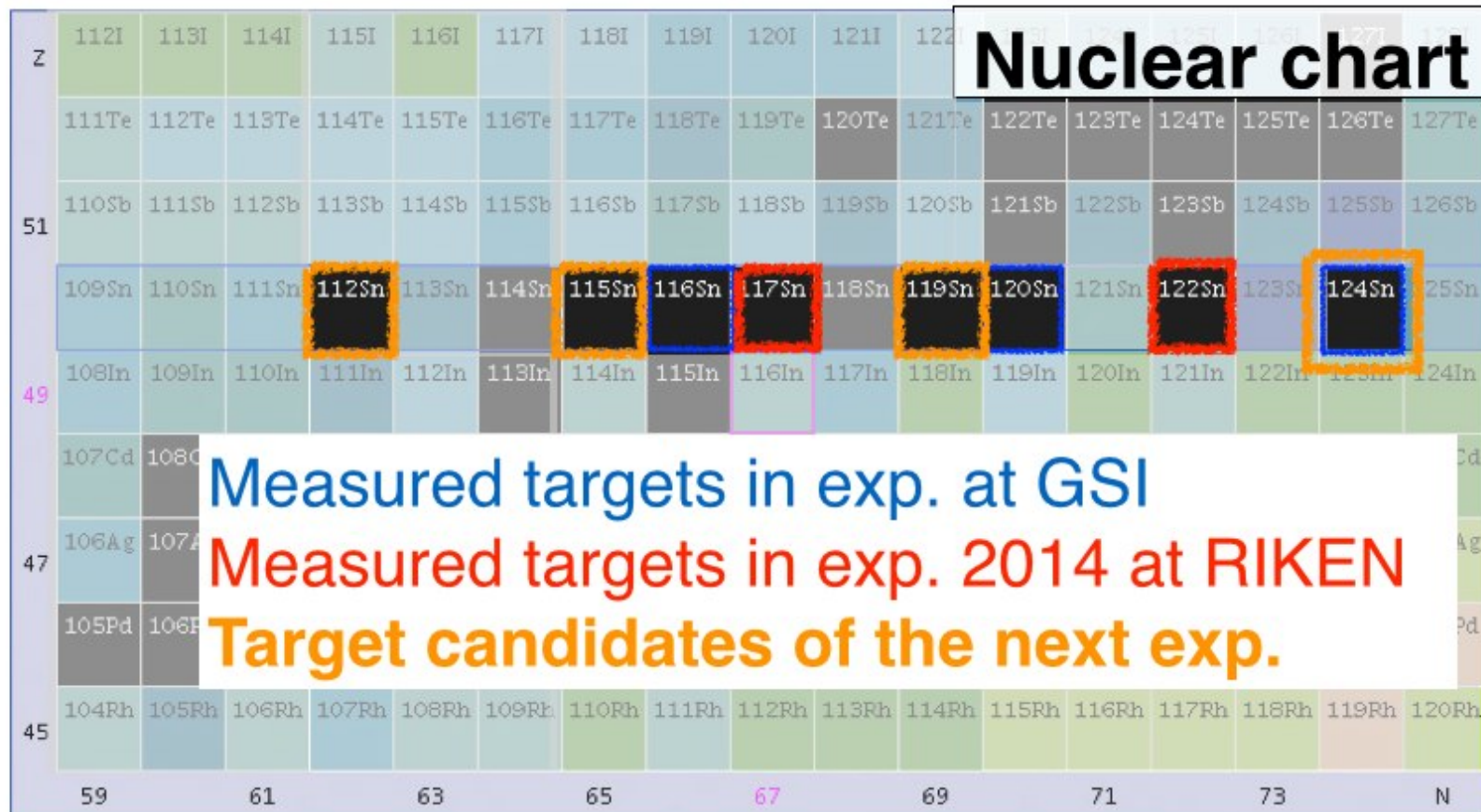


# Summary

- Deeply-bound pionic atom is good probe for QCD in finite density, especially for quark condensate via  $b_1$  parameter in  $\pi$  - A potential.
- To determine the  $b_1$  precisely, experiments of pionic Sn isotopes are on going at RIKEN.
- In the first exp. , we measured with the target of  $^{122,116}\text{Sn}$ , and succeed in
  - improvement of the resolution,
  - observation of the pionic 1s, 2p and 2s states in  $^{121, 116}\text{Sn}$ ,
- From the obtained binding energies and widths of pionic states,  $b_1$  in medium is evaluated most precisely, which is consistent with the partial restoration of chiral symmetry breaking.

## *(Near) future works*

NuDat



The next exp. are already approved in PAC at RIKEN with wider range of isotopes.

**The exp. will be performed in a few years.**