

CGC approach in p+A collisions

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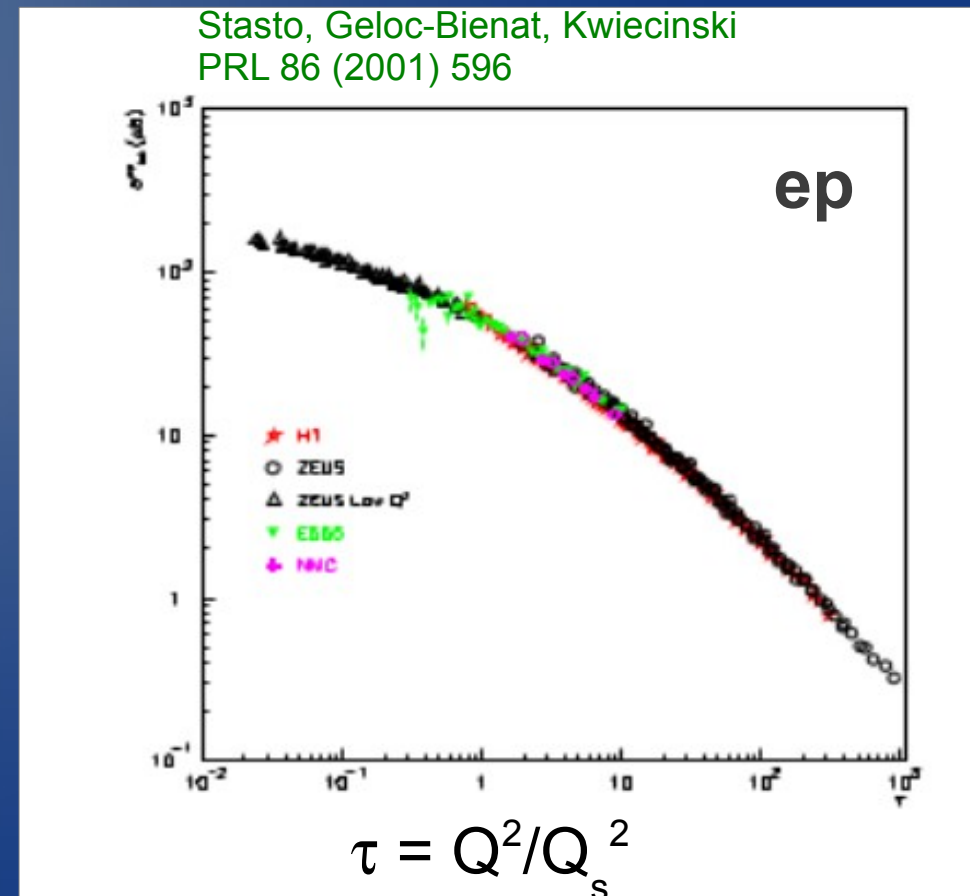
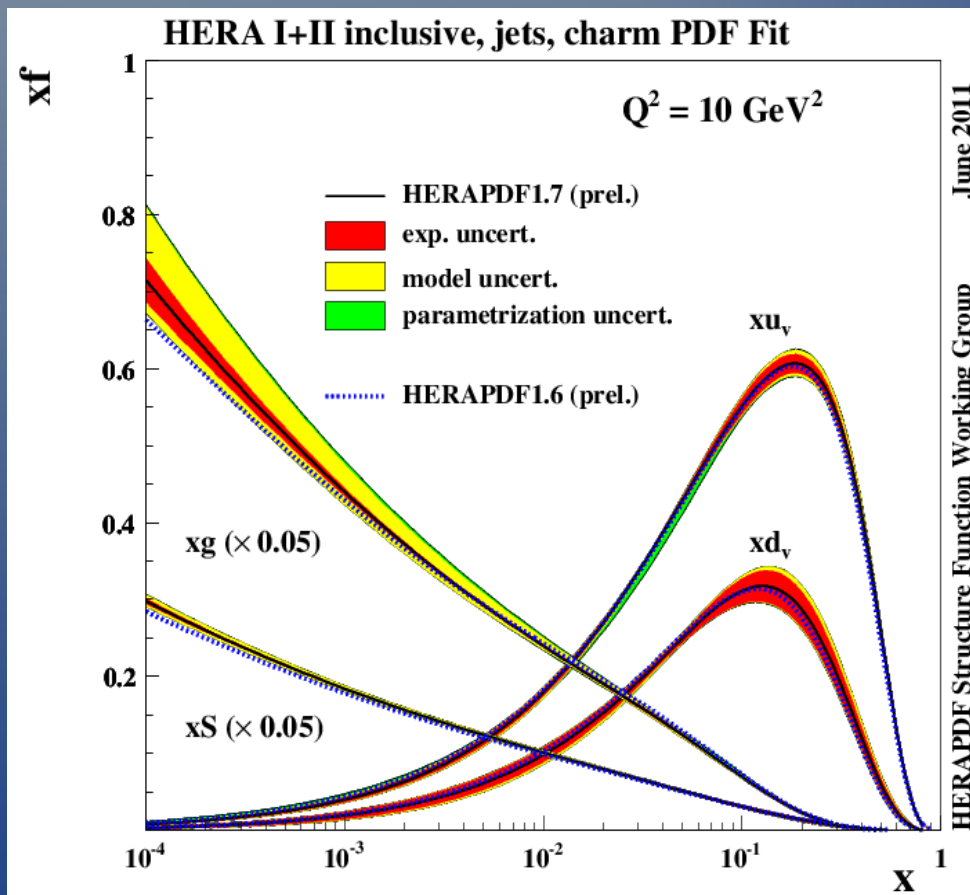
HF and K. Watanabe, arXiv:1304.2221[hep-ph]
J. Albacete, A. Dumitru, HF, Y. Nara, NPA897(2013)1
and others works

Outline

- Introduction
- rcBK phenomenology
 - single hadron production in forward region
 - heavy flavors
- Challenges
 - multi-point correlators
 - NLO extension
- Summary

Dense gluon and new scale

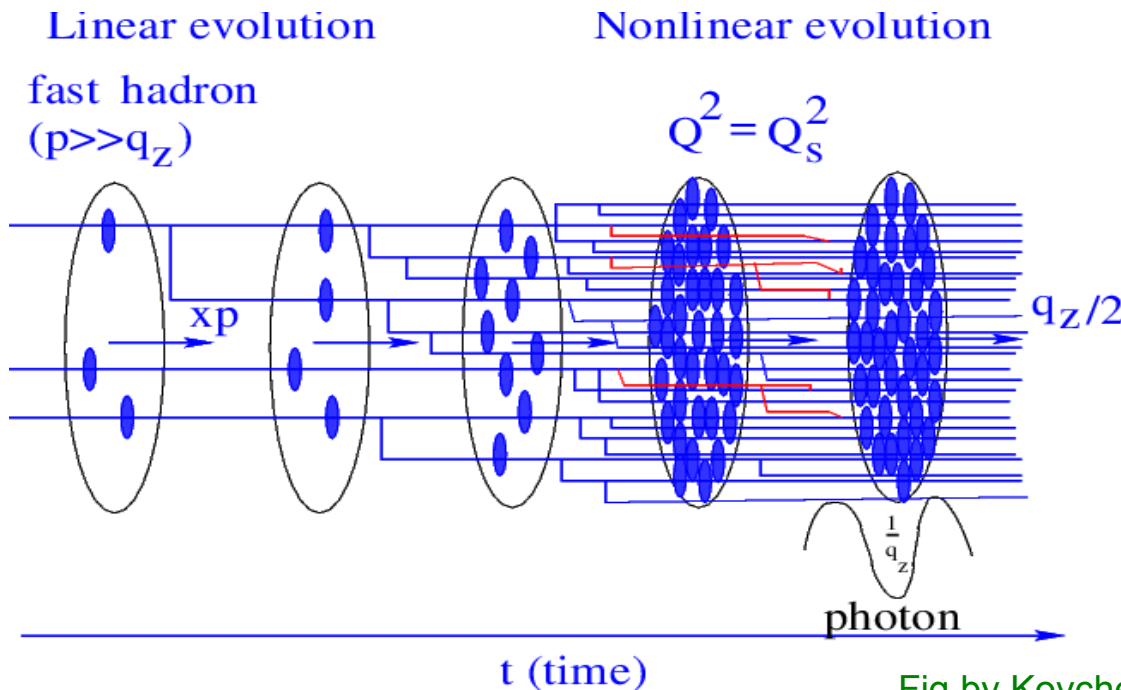
- Gluon number $xG(x, Q)$ increases at small x
- A new scaling suggests a new scale, Q_s^2



Intuitive picture of saturation

- The higher the energy is, the more parton fluctuations emerge
- Recombination becomes important when

$$Q_s^2(x) \sim \alpha_s \cdot \frac{xG(x, Q_s^2)}{\pi R_h^2}$$



Phase space density

$$\frac{dN}{d^2r d^2p} \sim \frac{xG(x, Q_s^2)}{\pi R_h^2 Q_s^2(x)} \sim \frac{1}{\alpha_s}$$

Fig by Kovchegov-Levin

Intuitive picture of saturation

- The higher the energy is, the more parton fluctuations appear
- Recombination becomes important when

$$Q^2(x) \sim \alpha_s \cdot \frac{xG(x, Q_s^2)}{\pi R_0^2}$$

$$Q_{s,A}^2(x) \sim \alpha_s \cdot \frac{AxG(x, Q_s^2)}{A^{2/3}\pi R_0^2} = A^{1/3}Q_{s,p}^2$$

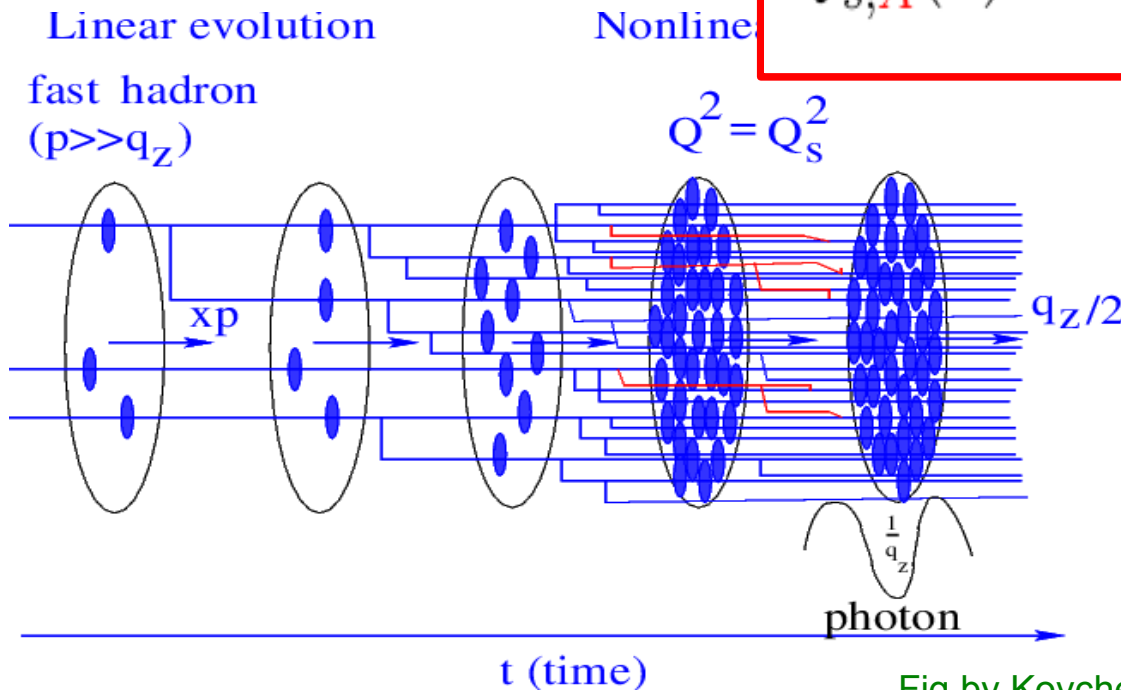
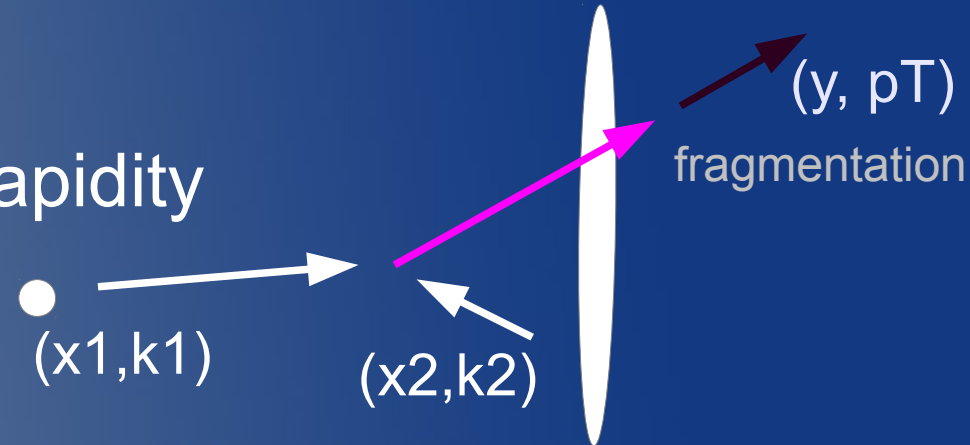


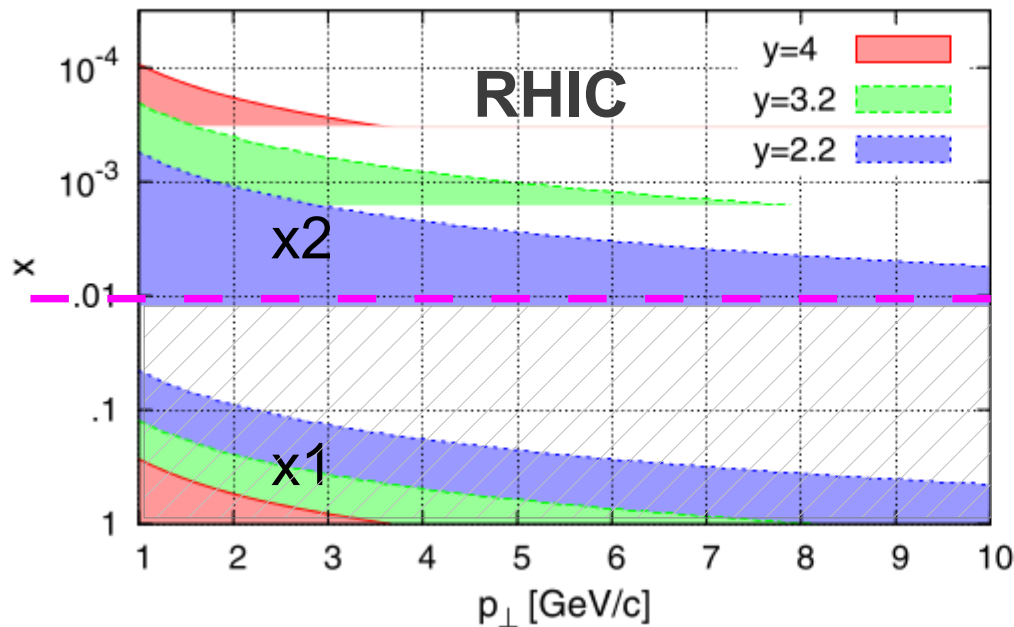
Fig by Kovchegov-Levin

Kinematic coverage at colliders

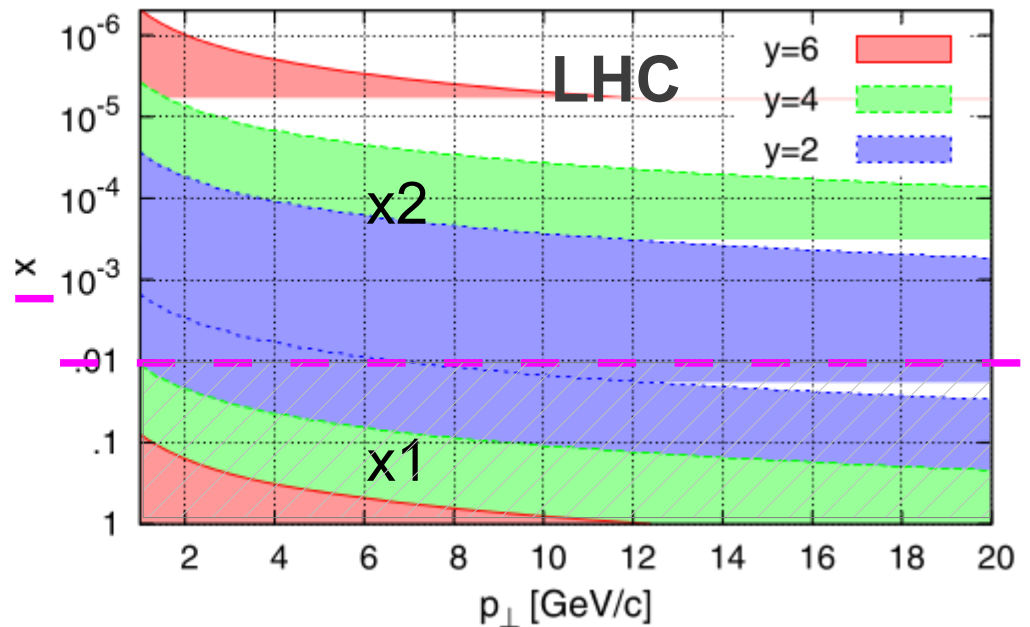
- $x_{1,2} \sim pT/\sqrt{s} \exp(\pm y)$
 - High energy, forward rapidity
- “small” $x < x_0 = 0.01$ for nuclei



Kinematic reach at $\sqrt{s}=200$ GeV



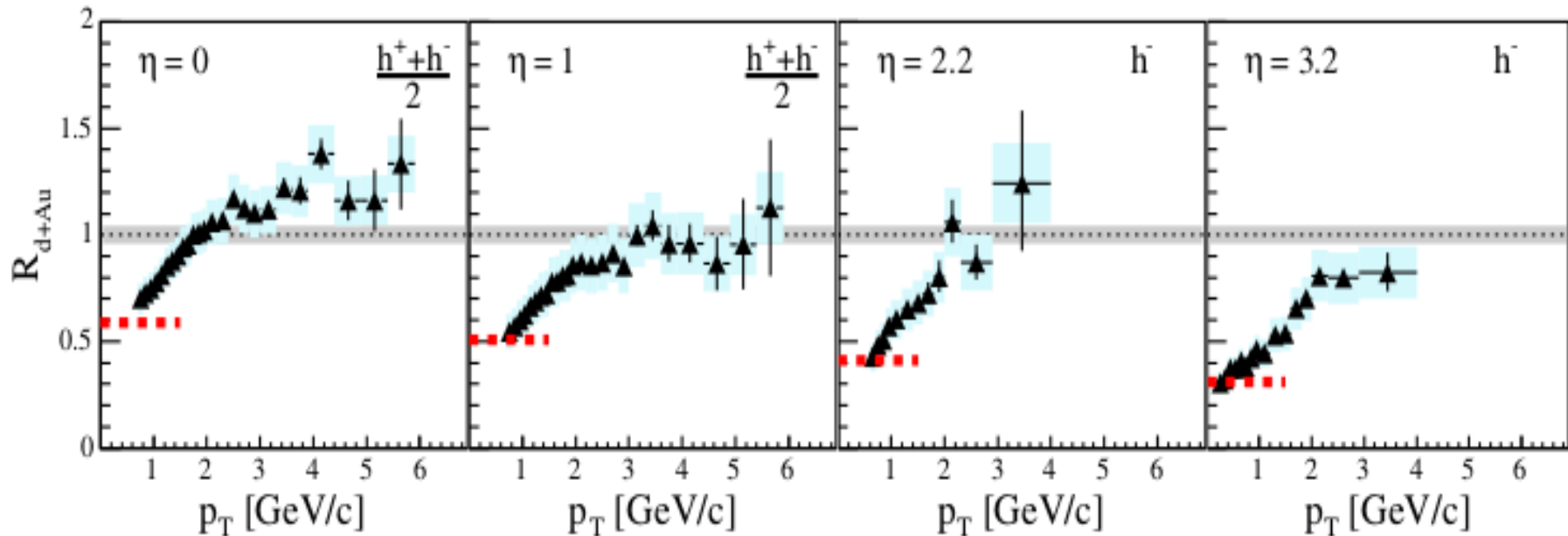
Kinematic reach at $\sqrt{s}=5$ TeV



Evidence in dA collisions

Forward suppression in d+Au at RHIC

Nuclear modification factor: $R_{dA} = \frac{1}{\langle N_{coll} \rangle} \frac{dN_{dA}/d^2p_{\perp}d\eta}{dN_{pp}/d^2p_{\perp}d\eta}$



- BRAHMS [arXiv: nucl-ex/0403005] shows
 $\eta = 0$: Cronin peak, $\eta \sim 3$: Suppression
- Qualitatively consistent with the CGC

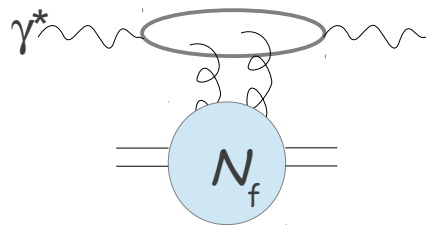
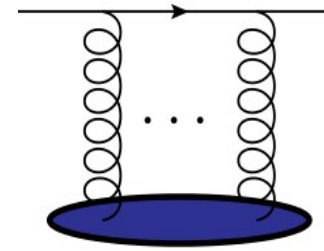
rcBK phenomenology

rcBK phenomenology

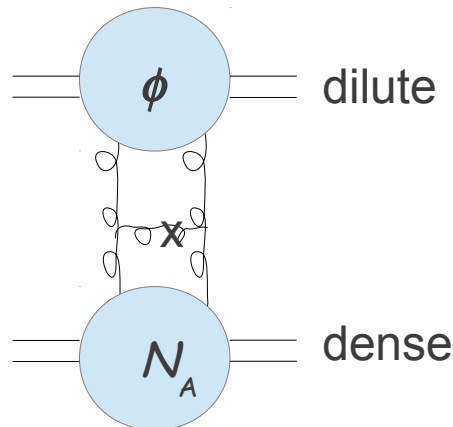
- Energetic parton coherently interact with target (eikonal multiple scatterings)

$$U(z_{\perp}) = \mathcal{P} \exp \left[ig \int dz^+ A^-(z^+, z) \right]$$

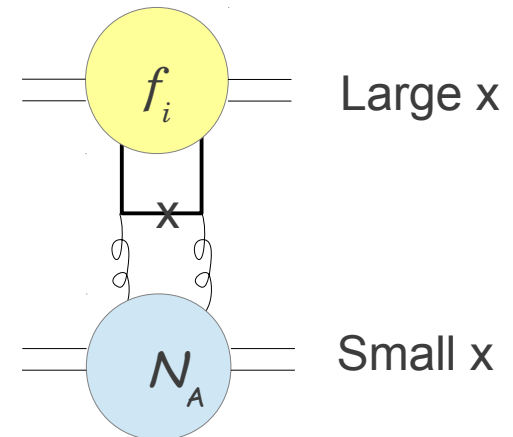
$$\mathcal{N}(\mathbf{x}_{\perp}, \mathbf{y}_{\perp}) = \int [D\rho] W_Y[\rho] \left[1 - \frac{1}{N_c} \text{tr}(U(\mathbf{x}_{\perp})U^{\dagger}(\mathbf{y}_{\perp})) \right]$$



dipole model



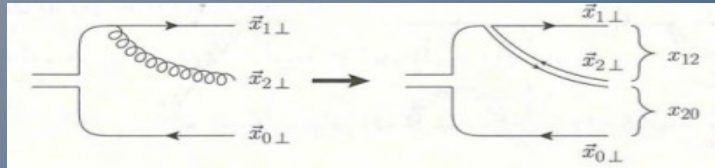
Kt factorization



DHJ factorization

rcBK phenomenology

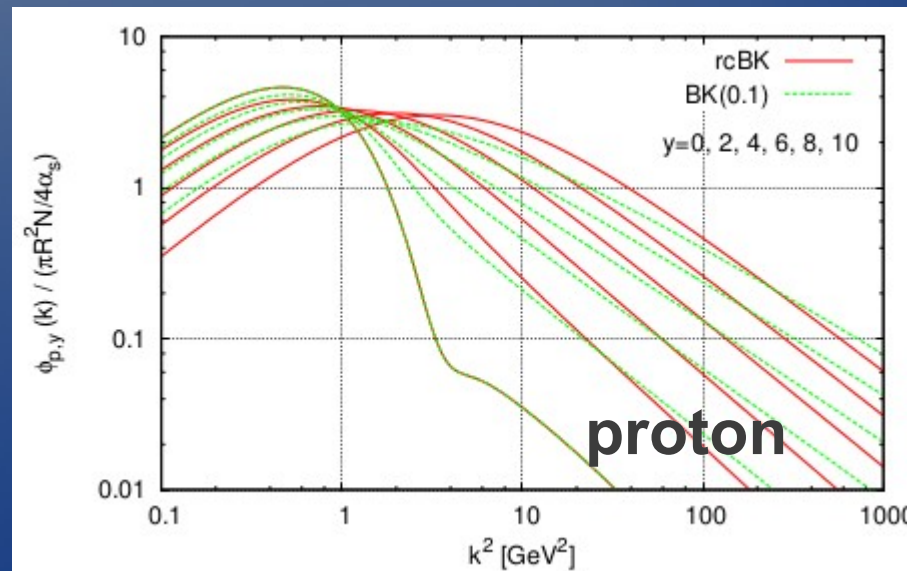
- Energy (x) dependence described by BK eqn with running coupling correction (in large N_c)



$$\frac{\partial \mathcal{N}(r, y)}{\partial y} = \int d^2 r_1 K^{\text{run}} [\mathcal{N}(r_1, y) + \mathcal{N}(r_2, y) - \mathcal{N}(r, y) - \mathcal{N}(r_1, y)\mathcal{N}(r_2, y)]$$

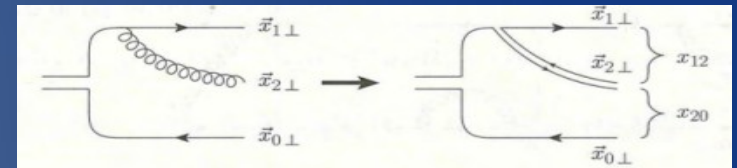
Example of evolution

Kovchegov-Weigert, Balitsky-Chirilli



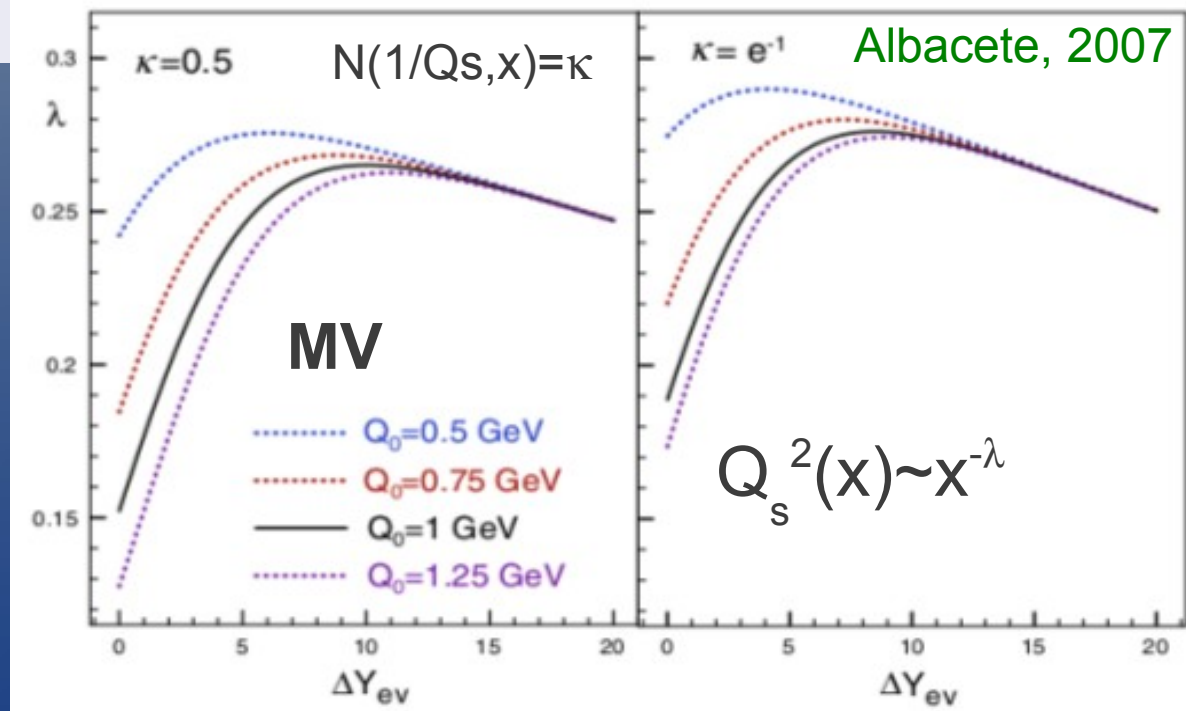
rcBK phenomenology

- In large N_c limit, evolution of dipole is closed



- Balitsky-Kovchegov eqn with running coupling

$$\frac{\partial \mathcal{N}(r, y)}{\partial y} = \int d^2 r_1 \kappa^{\text{run}} [\mathcal{N}(r_1, y) + \mathcal{N}(r_2, y) - \mathcal{N}(r, y) - \mathcal{N}(r_1, y) \mathcal{N}(r_2, y)]$$



gert, Balitsky-Chirilli

Constraining N with DIS data

Initial condition and evolution

Albacete et al. PRD80

$$\mathcal{N}_F(r, x=x_0) = 1 - \exp \left[-\frac{(r^2 Q_{s0, \text{proton}}^2)^\gamma}{4} \ln \left(\frac{1}{\Lambda r} + e \right) \right]$$

$$\alpha_s(r) = 1/[b_0 \ln(4C/r\Lambda + a)]$$

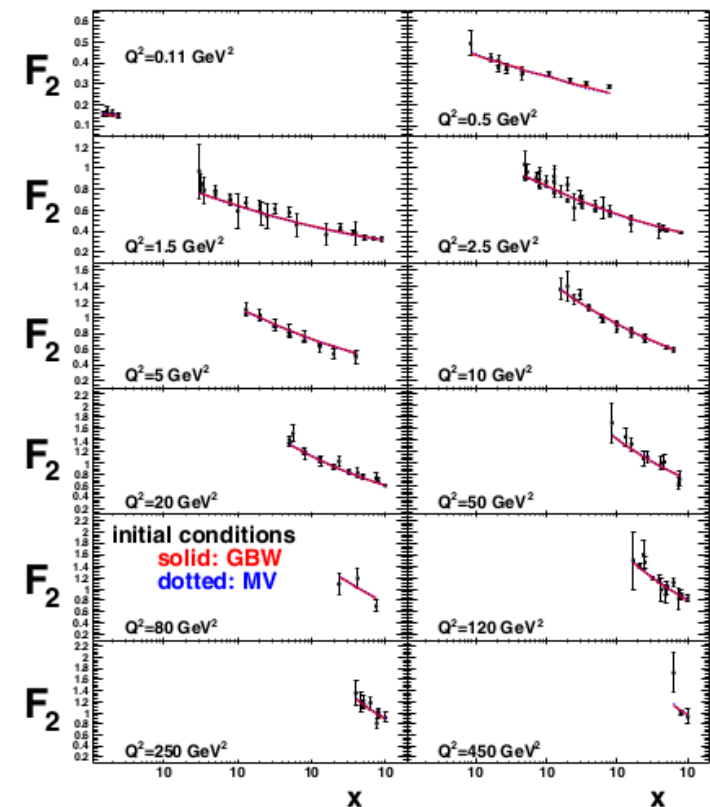
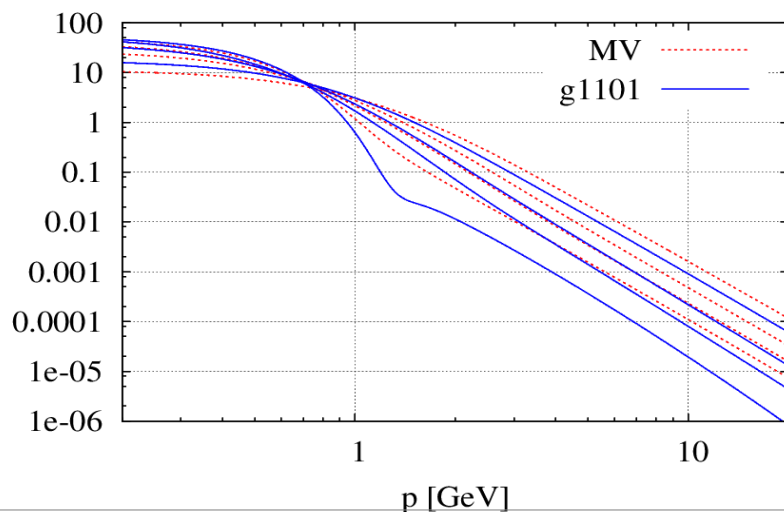
Fit to HERA-DIS data

$$\sigma_{T,L}(x, Q^2) = \sigma_0 \int_0^1 dz \int dr |\Psi_{T,L}(z, Q^2, \mathbf{r})|^2 \mathcal{N}(r, Y)$$

Parameter set

UGD Set	$Q_{s0, \text{proton}}^2$ (GeV ²)	γ	α_{fr}	C
MV	0.2	1	0.5	1
g1.119	0.168	1.119	1.0	2.47
g1.101	0.157	1.101	0.8	1

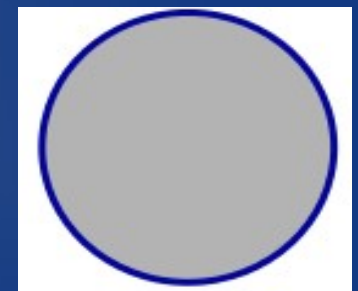
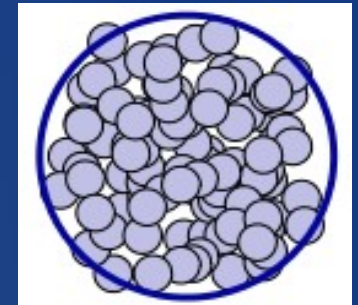
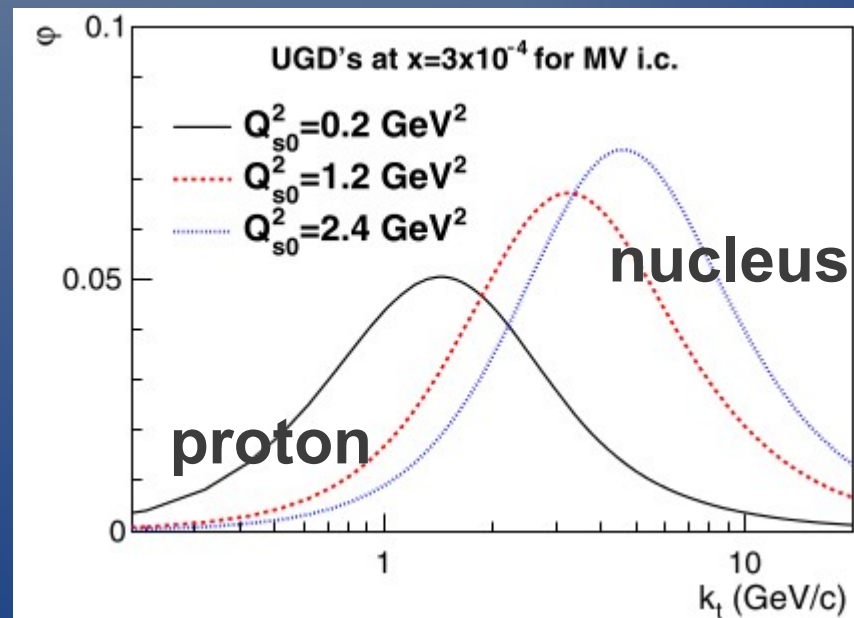
$N(k, y)$, $y=0, 1.5, 3, 6$



Modeling of nuclear target

$$Q_{s0A}^2 = A^{1/3} Q_{s0}^2$$

- Effective Q_s^2 w/o impact param dependence – simplest
- MC modeling with local $Q_s^2(b)$
- Low k_T gluon dist is more suppressed



Hadron production

- INPUT: gluon dist from rcBK in large N_c

$$\phi(k, y) \sim k^2 N_A(k, y), \quad 1 - N_A = (1 - N_F)^2$$

- kT-factorization for small $x_{1,2}$ ($y \sim 0$)

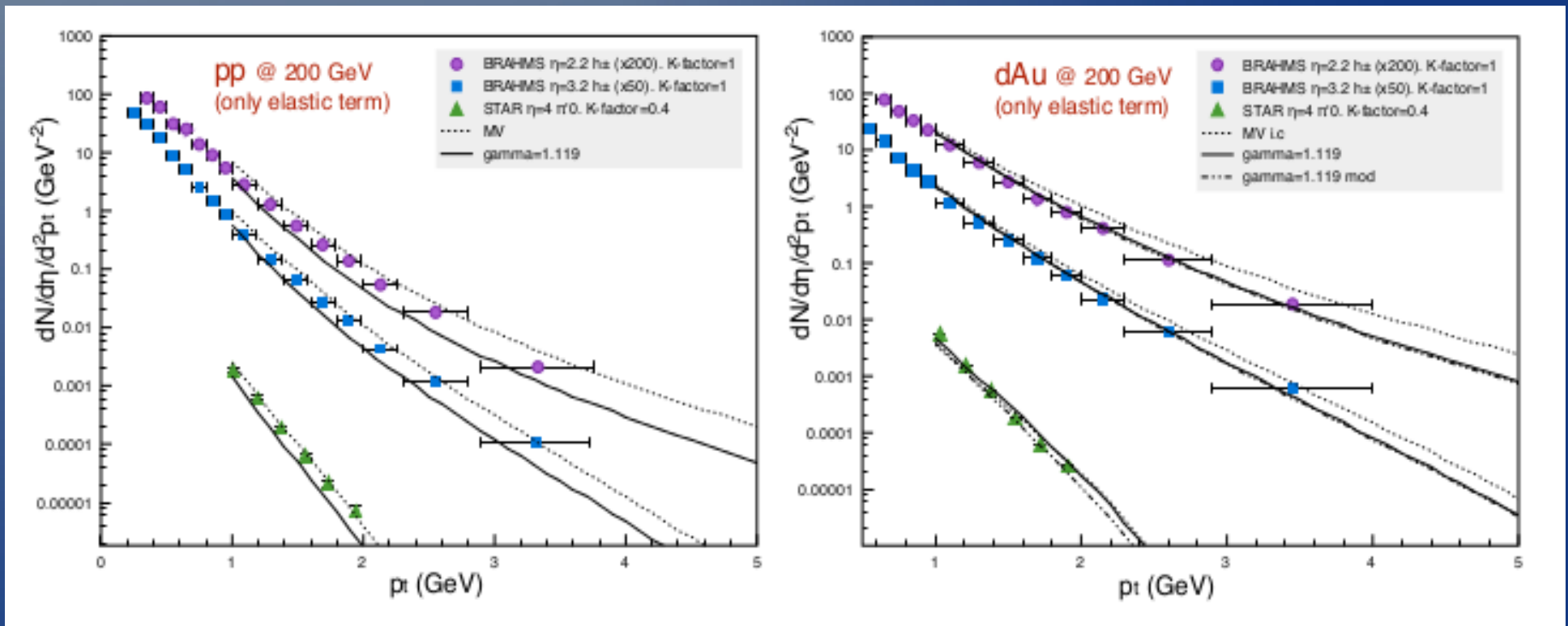
$$\frac{d\sigma^{A+B \rightarrow g}}{dy d^2 p_t d^2 R} = K^k \frac{2}{C_F} \frac{1}{p_t^2} \int \frac{d^2 k_t}{4} \times \int d^2 b \alpha_s(Q) \varphi_P\left(\frac{|p_t + k_t|}{2}, x_1; b\right) \varphi_T\left(\frac{|p_t - k_t|}{2}, x_2; R - b\right)$$

- DHJ *hybrid* formula for small x_2 for $y > 0$

$$\left[\frac{dN_h}{d\eta d^2 k} \right]_{\text{el}} = \frac{1}{(2\pi)^2} \int_{x_F}^1 \frac{dz}{z^2} \left[\sum_q x_1 f_{q/p}(x_1, Q^2) \tilde{N}_F\left(x_2, \frac{k}{z}\right) D_{h/q}(z, Q^2) + x_1 f_{g/p}(x_1, Q^2) \tilde{N}_A\left(x_2, \frac{k}{z}\right) D_{h/g}(z, Q^2) \right],$$

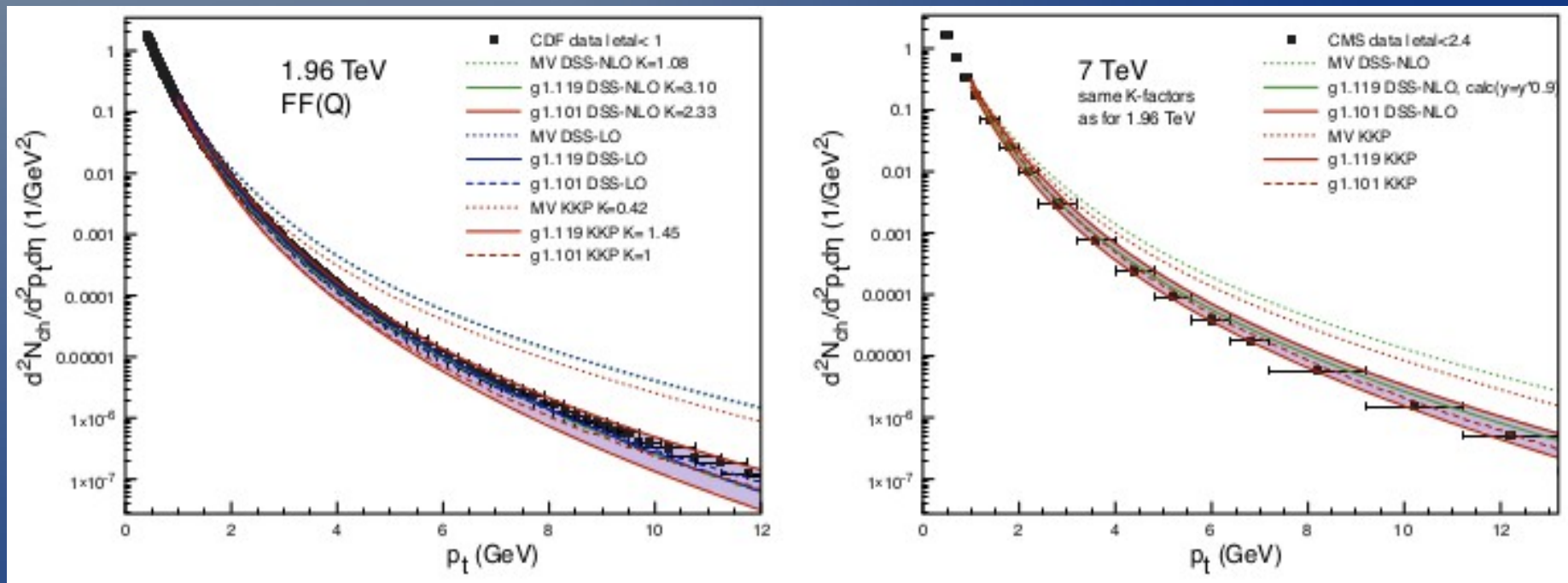
Hadron spectrum in pp, dA at RHIC

- DHJ formula (CTEQ6, DSS)
- The same normalization chosen for pp and dA
 - particle dependent

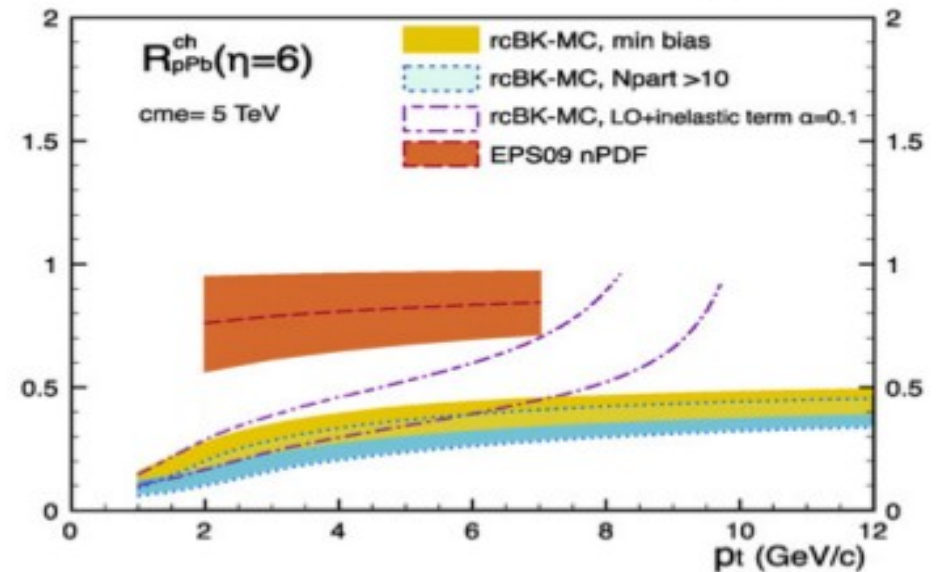
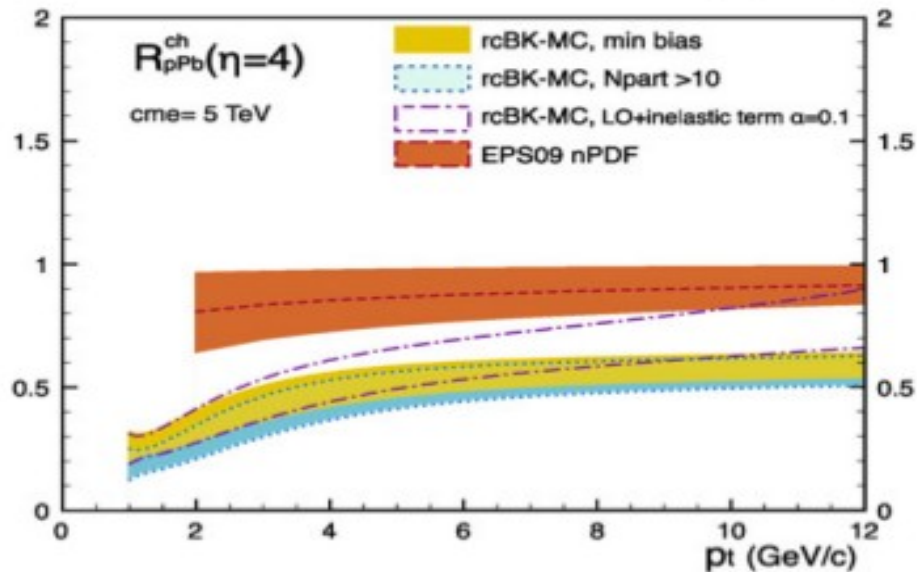
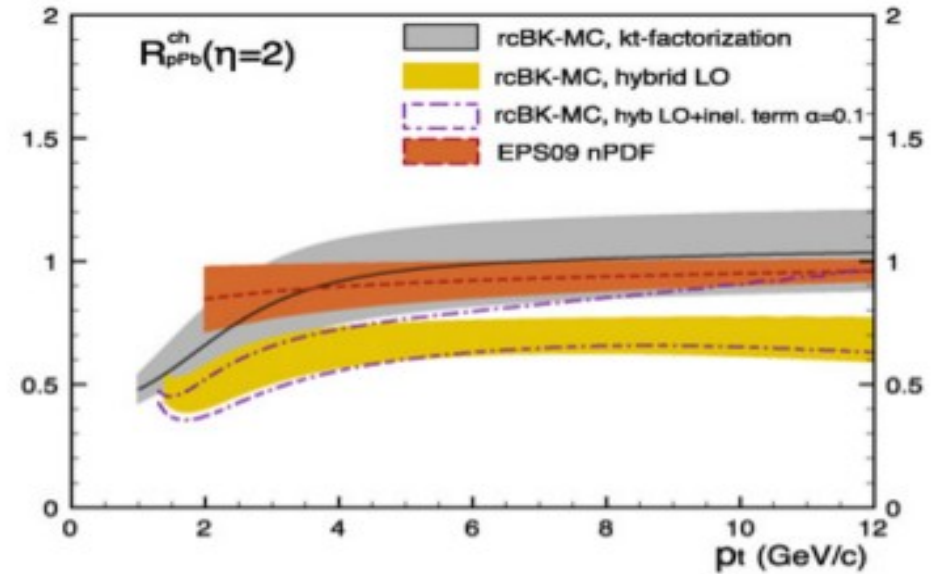
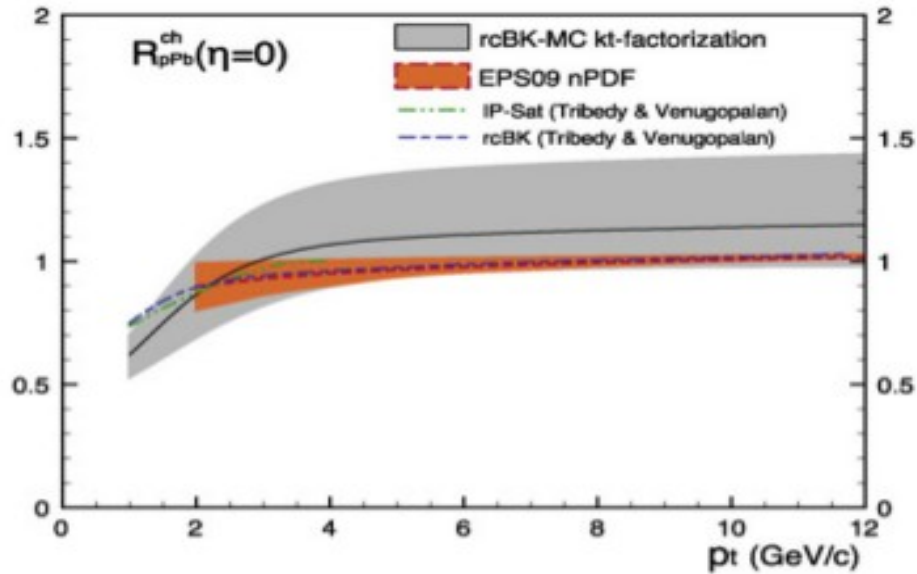


Hadron spectrum at pp colliders

- kT-factorized formula
 - Normalized at $p_T=1$ GeV for $\sqrt{s}=1.96$ TeV
 - uGD set ($\gamma \sim 1.1$) describes energy and p_T dependences



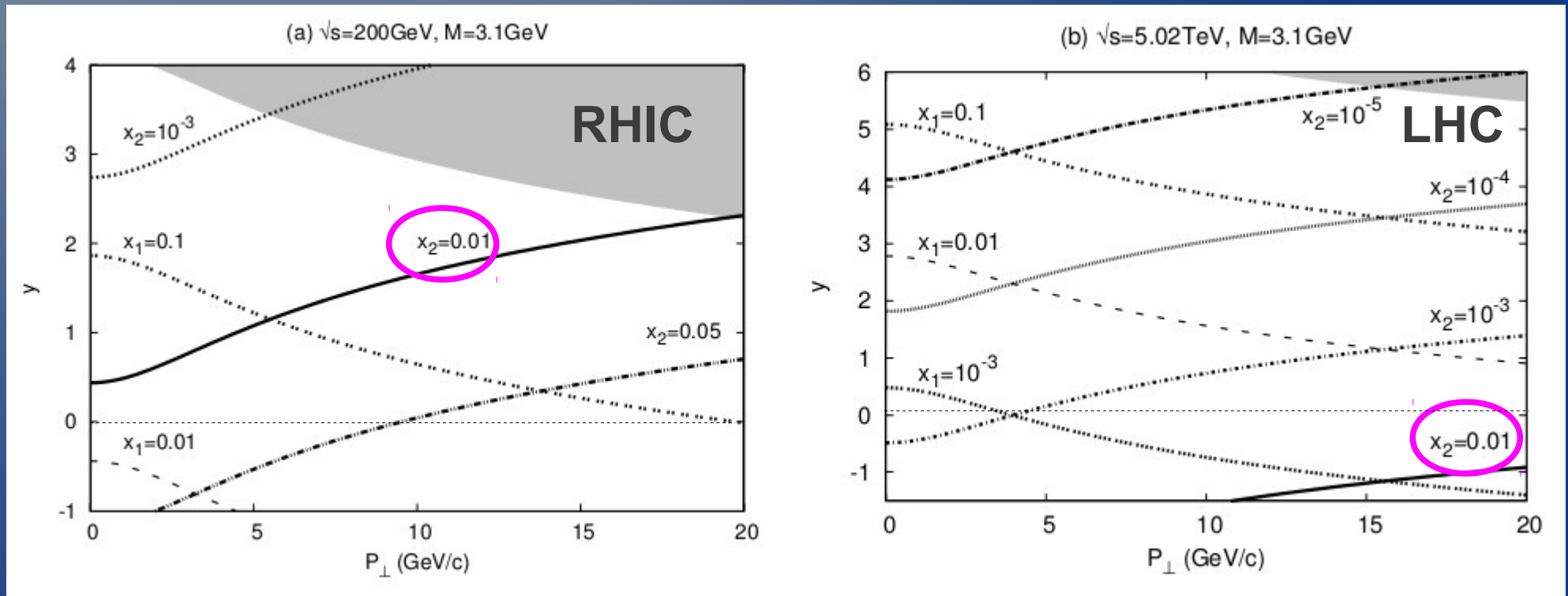
y-dependence of RpA at the LHC



Heavy quarks are sensitive?

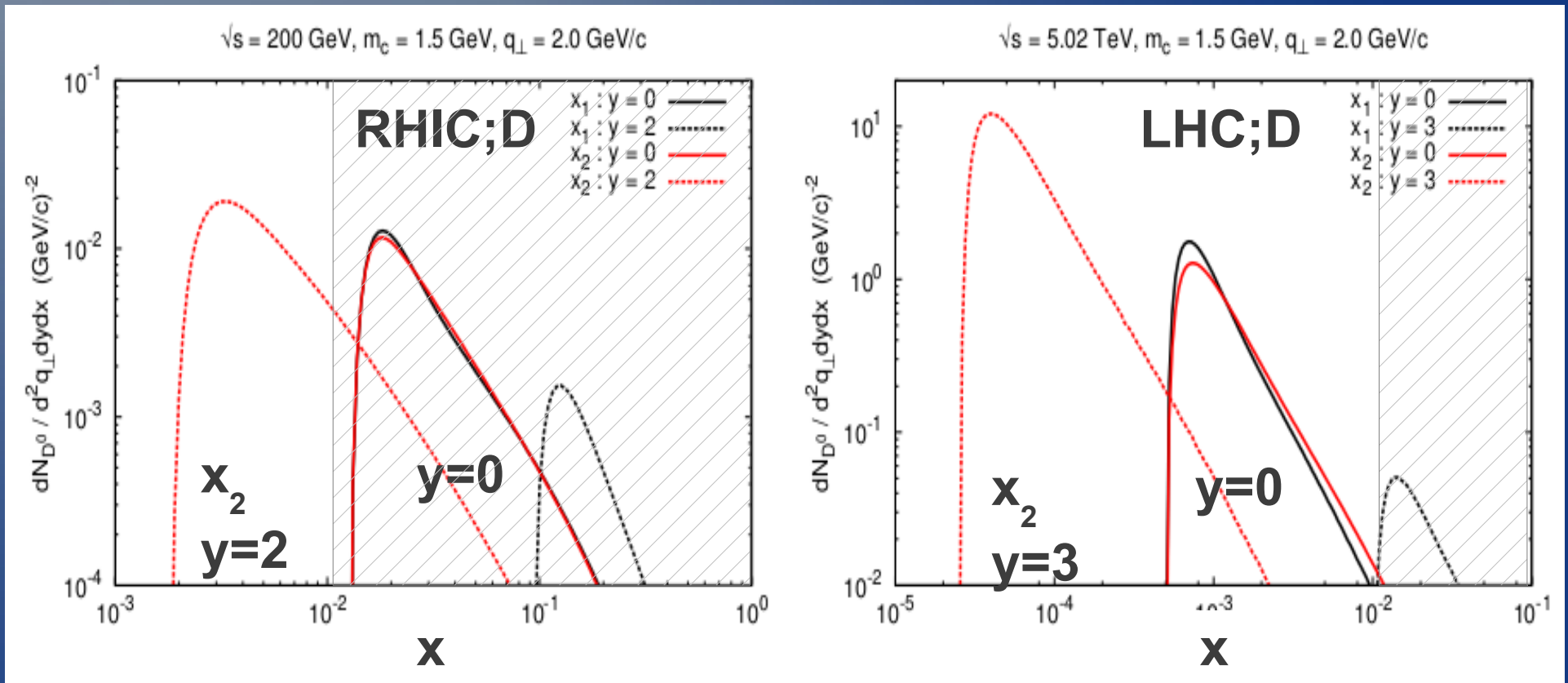
Heavy quarks are sensitive?

- Quarkonium; $gg \rightarrow J/\psi$
- maybe at RHIC, while must be at the LHC



How about open heavy flavors ?

- $c \rightarrow D$
- at $q_T = 2 \text{ GeV}$



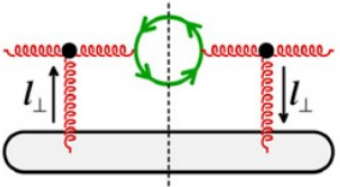
Quark pair production in pA

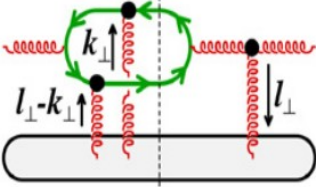
- Formula at $O(\rho_p \rho_A^{00})$ in large N_c limit

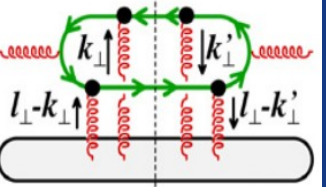
$$\frac{dN_{q\bar{q}}}{d^2\mathbf{p}_\perp d^2\mathbf{q}_\perp dy_p dy_q} = \frac{1}{\pi R_A^2} \frac{\alpha_s^2 N}{8\pi^4 d_A} \frac{1}{(2\pi)^2} \int_{\mathbf{k}_{2\perp}, \mathbf{k}_\perp} \frac{\Xi(\mathbf{k}_{1\perp}, \mathbf{k}_{2\perp}, \mathbf{k}_\perp)}{\mathbf{k}_{1\perp}^2 \mathbf{k}_{2\perp}^2} \phi_{A,y_2}^{q\bar{q},g}(\mathbf{k}_{2\perp}, \mathbf{k}_\perp) \varphi_{p,y_1}(\mathbf{k}_{1\perp})$$

Gelis-Blaizot-Venugopalan
HF-Gelis-Venugopalan

- Pair production --> multi-parton correlators

$$\phi_A^{g,g}(\mathbf{l}_\perp) =$$


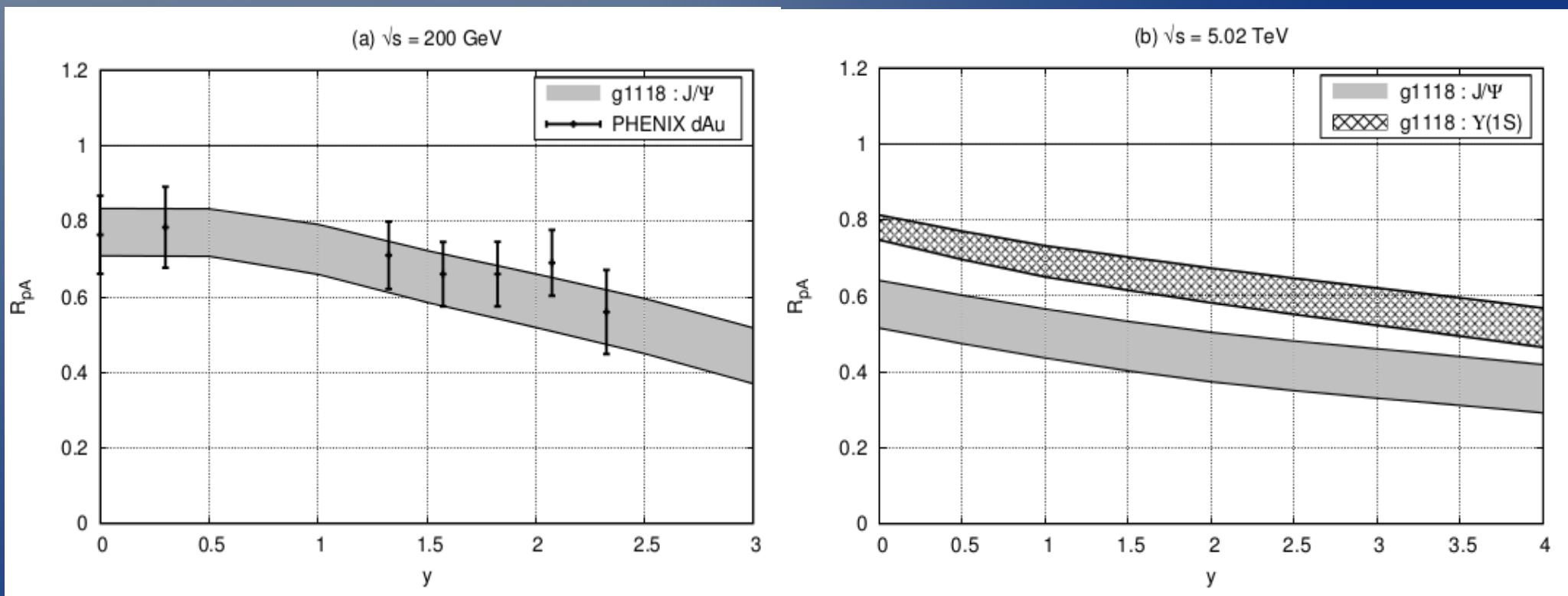
$$\phi_A^{q\bar{q},g}(\mathbf{l}_\perp; \mathbf{k}_\perp) =$$


$$\phi_A^{q\bar{q},q\bar{q}}(\mathbf{l}_\perp; \mathbf{k}_\perp, \mathbf{k}'_\perp) =$$


- 4-pt & 3-pt functions simplify to a product of fundamental 2-pt funs in large N_c limit

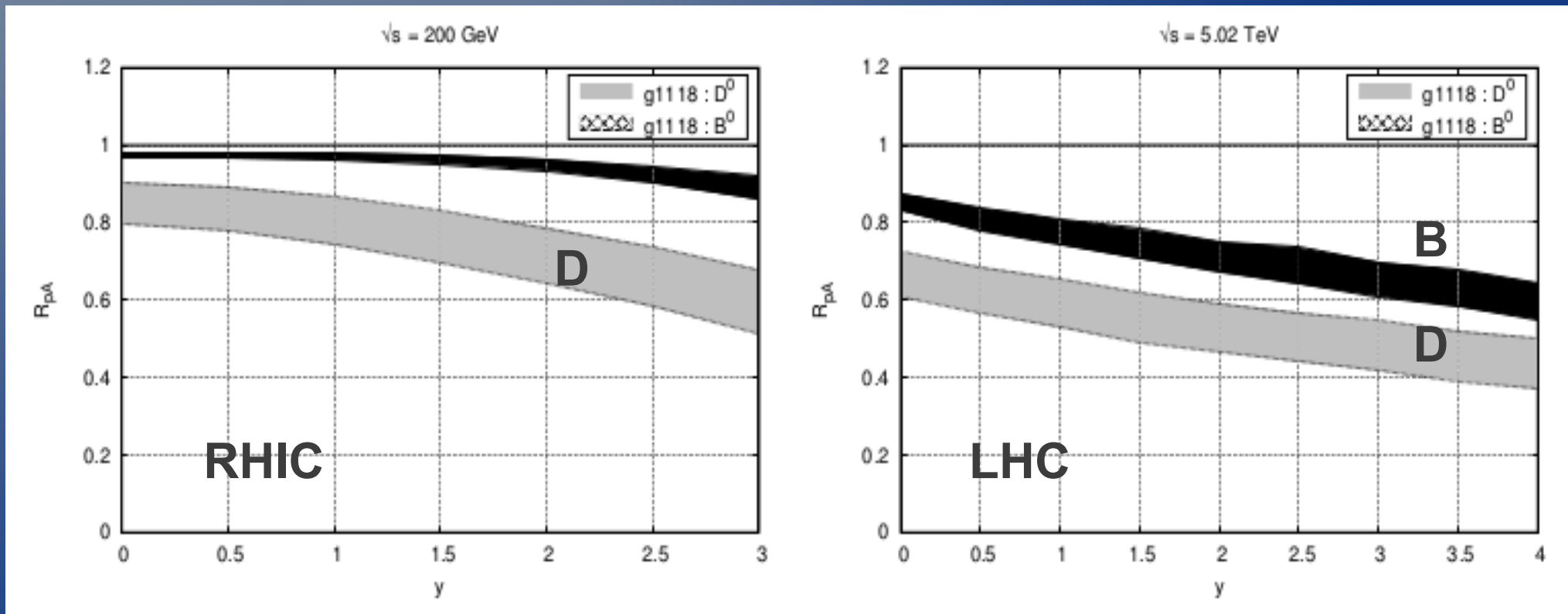
$R_{pA}(y)$ for J/psi at the LHC

- Color evaporation model
- More suppression at the LHC than at RHIC
- Upsion also suppressed significantly at LHC



RpA(y) for D,B (prelim)

- Suppression due to saturation is seen



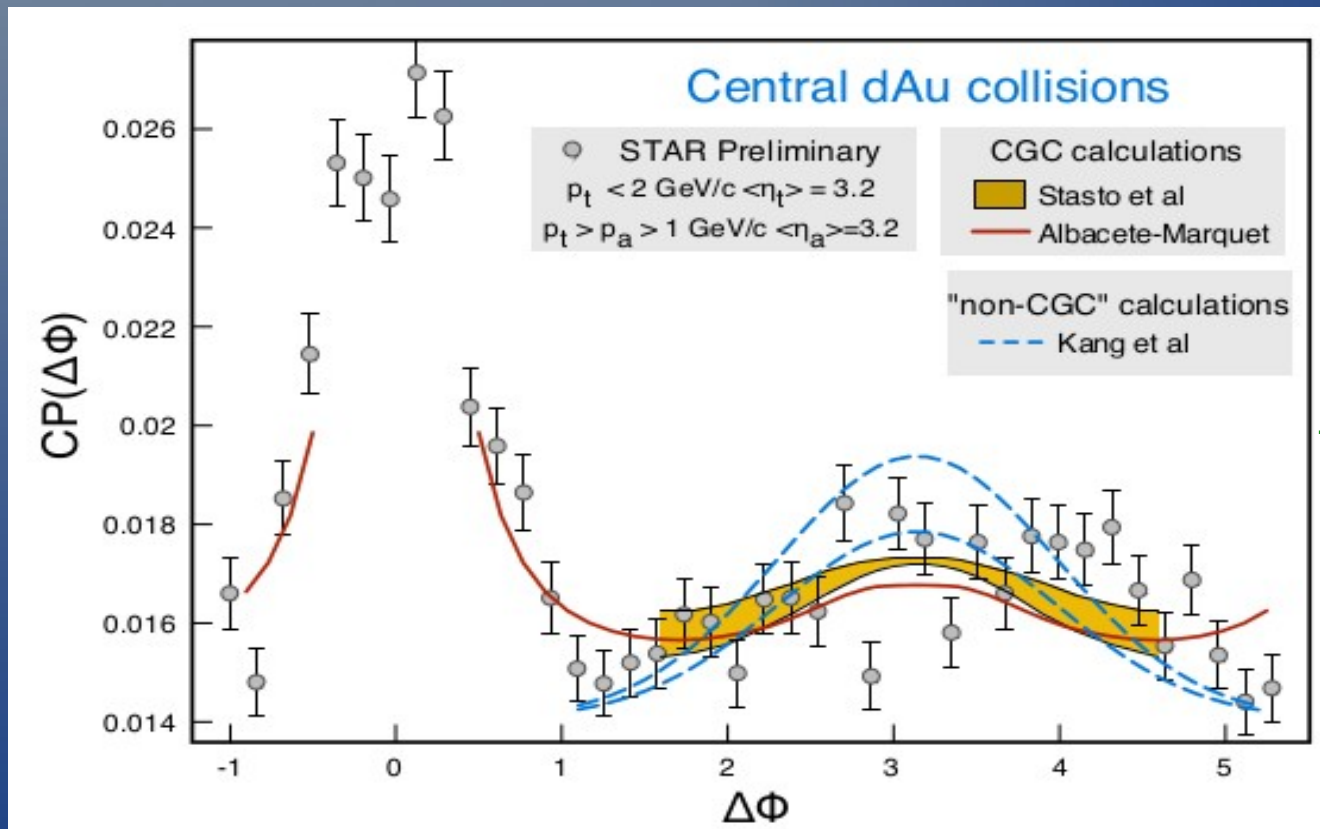
Challenges

- Multi parton correlators from JIMWLK
- Towards NLO calculations in CGC

di-hadron correlations at RHIC

- Extend hybrid formula to 2 particle production

$$N_{pair}(\Delta\phi) = \int_{y_i, |p_{i\perp}|} \frac{dN^{pA \rightarrow h_1 h_2 X}}{d^3p_1 d^3p_2}, \quad N_{trig} = \int_{y, p_\perp} \frac{dN^{pA \rightarrow hX}}{d^3p}$$

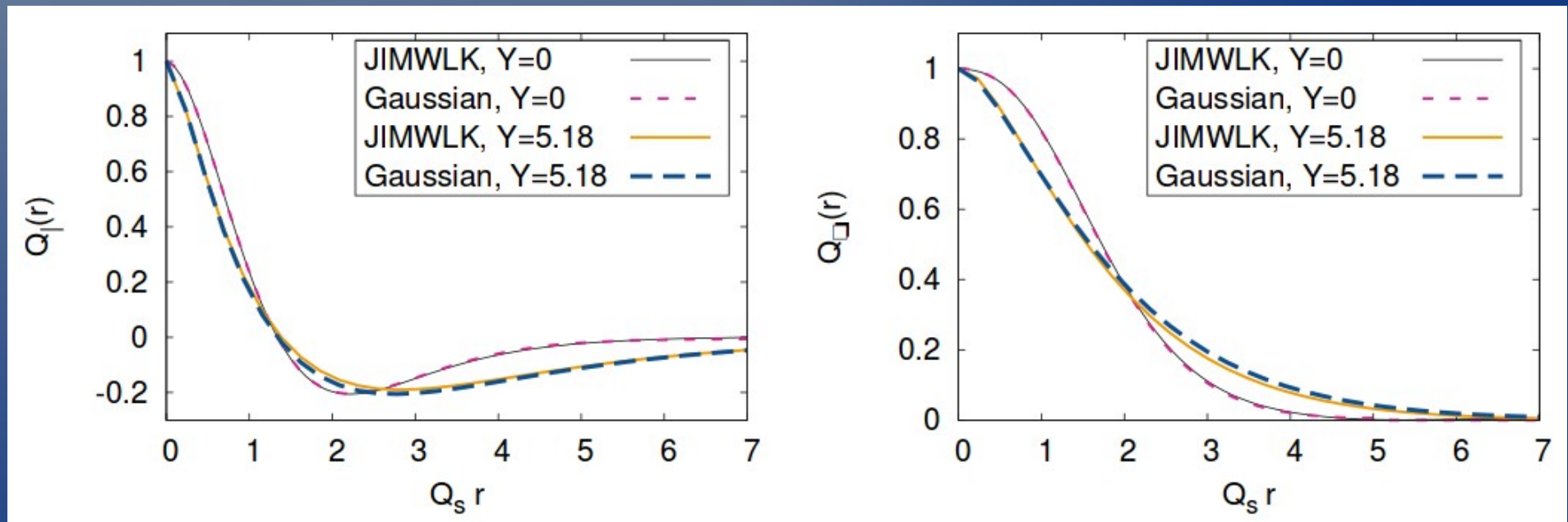


Albacete-Marquet, PRL 105, (2010)
 Stasto et al. arXiv:1109.1817 [hep-ph]
 Z.-B. Kang, et al. PRD85, (2012)

di-hadron correlations at RHIC

- Extend hybrid formula to 2 particle production
- involves 4-pt correlator, which needs JIMWLK

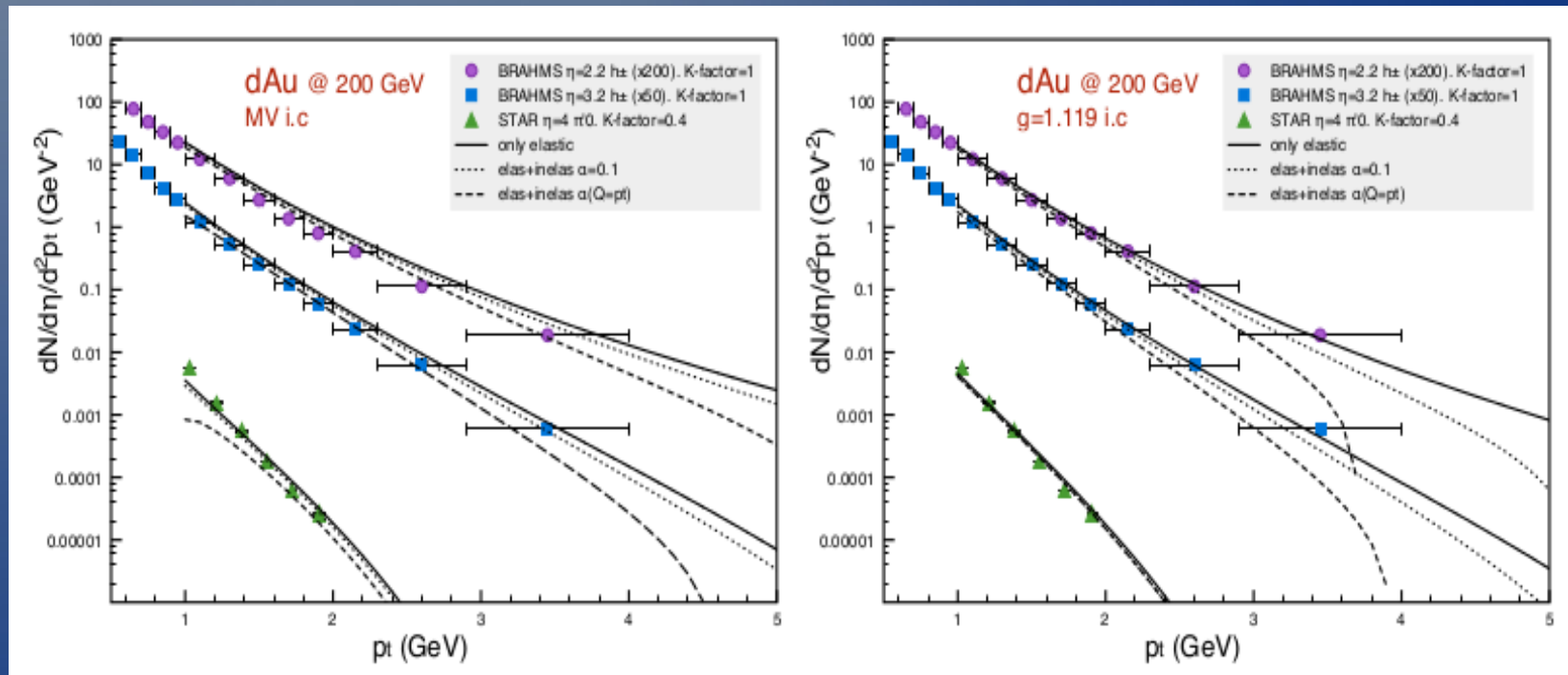
$$\hat{Q}(\mathbf{x}_T, \mathbf{y}_T, \mathbf{u}_T, \mathbf{v}_T) = \frac{1}{N_c} \text{Tr} U(\mathbf{x}_T) U^\dagger(\mathbf{y}_T) U(\mathbf{u}_T) U^\dagger(\mathbf{v}_T).$$



Dumitru et al., Phys.Lett.B706, 219 (2011)

Towards NLO calculation

- Part of NLO contributions is known to be large
- Full NLO calculation is desired



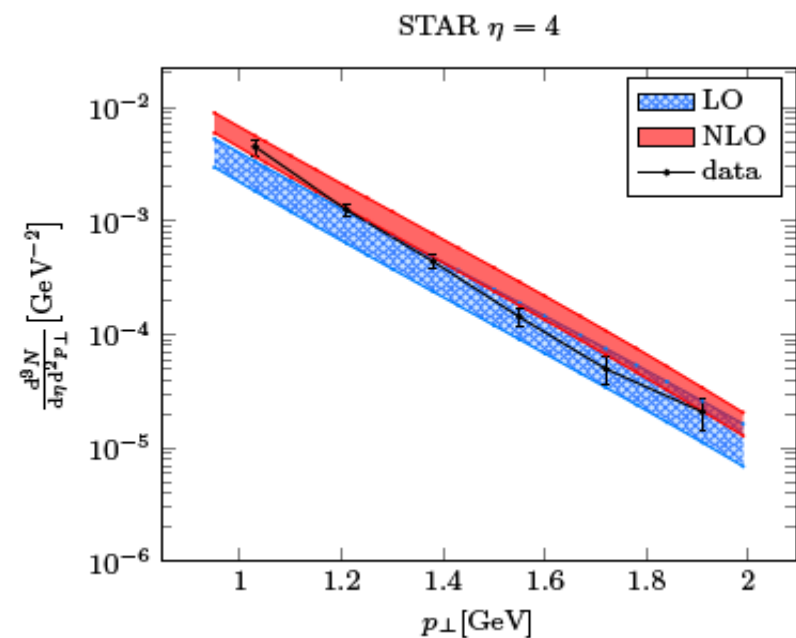
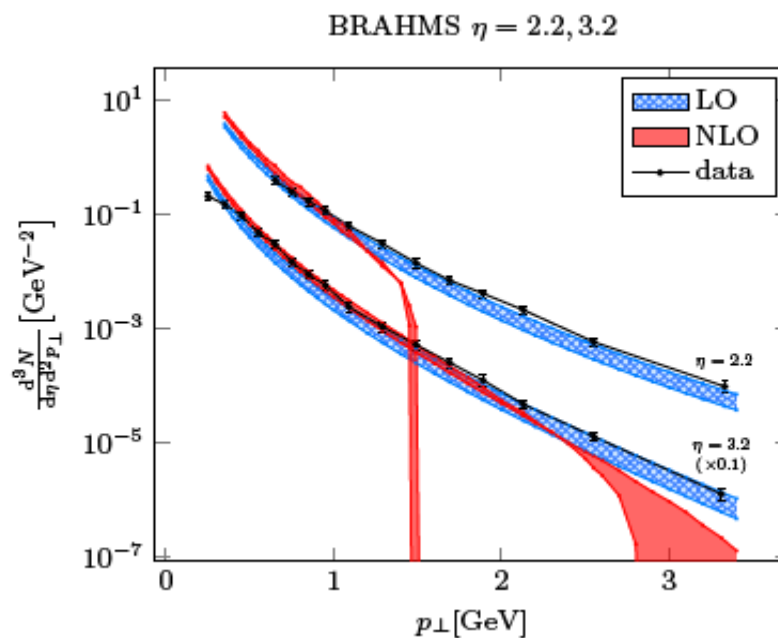
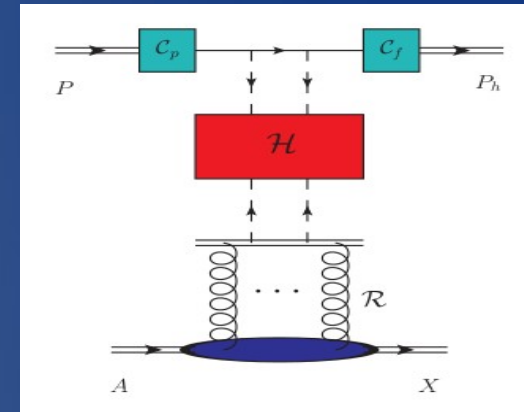
Albacete, Dumitru, HF, Nara, NPA897,1 2012

Towards NLO calculation

Chirilli, Xiao, Yuan (2012)
 Stasto, Xiao, Zaslavsky, arXiv:1307.4057

- Rapidity, collinear divergences identified and absorbed in uGD, pdf&FF, resp.

$$d\sigma = \int x f_a(x) \otimes D_a(z) \otimes \mathcal{F}_a^{xg}(k_\perp) \otimes \mathcal{H}^{(0)} + \frac{\alpha_s}{2\pi} \int x f_a(x) \otimes D_b(z) \otimes \mathcal{F}_{(N)ab}^{xg} \otimes \mathcal{H}_{ab}^{(1)},$$



MSTW NLO
 DSS NLO
 rcBK

Summary

- rcBK phenomenology fairly works for inclusive obs in pA from RHIC to LHC
- For more robust results,
 - multi-point correlations for more exclusive obs.
 - NLO extension from LO
- Important for providing the correct I.C. for AA

dN/dy in p-Pb collisions

ALICE:
Phys.Rev.Lett. 110 (2013) 032301

