



北京大学
PEKING UNIVERSITY

iTHES workshop on "Exploration of hidden symmetries in atomic nuclei"

Saturday, 27 July 2013 from 9:00 to 18:00
at RIKEN Wako Campus

Hidden symmetries in atomic nuclei: a brief introduction

Jie Meng 孟杰

北京大学物理学院

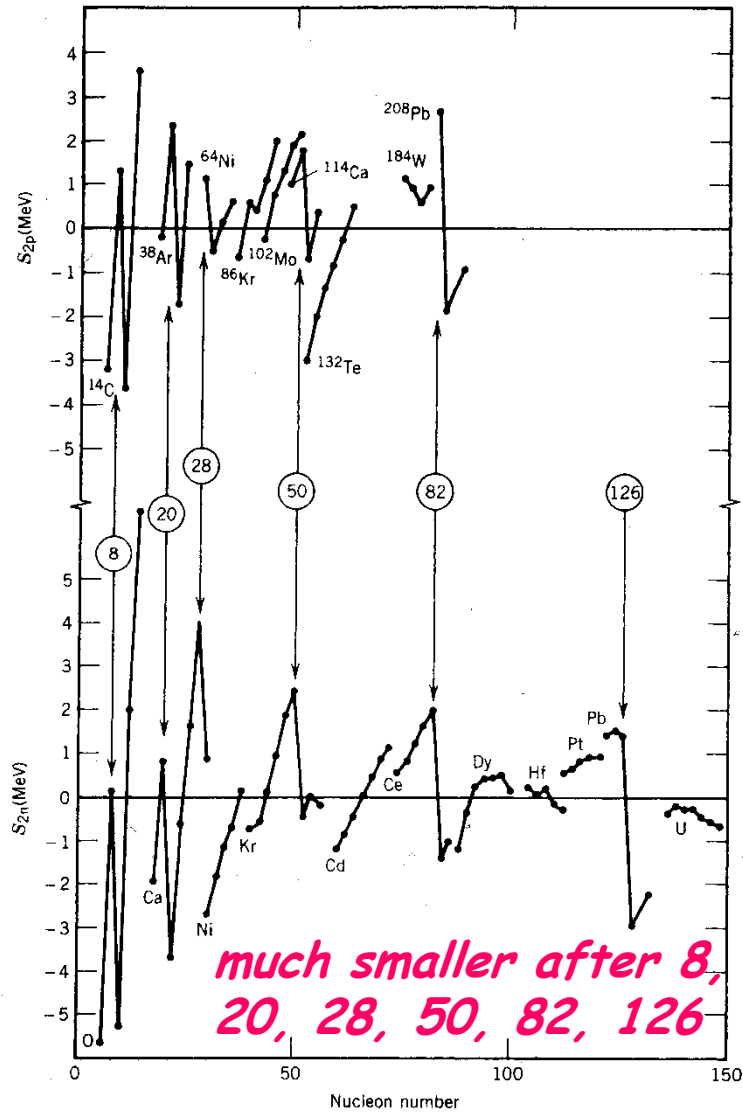
School of Physics, Peking University

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Nuclear magic number

two-nucleon separation energy



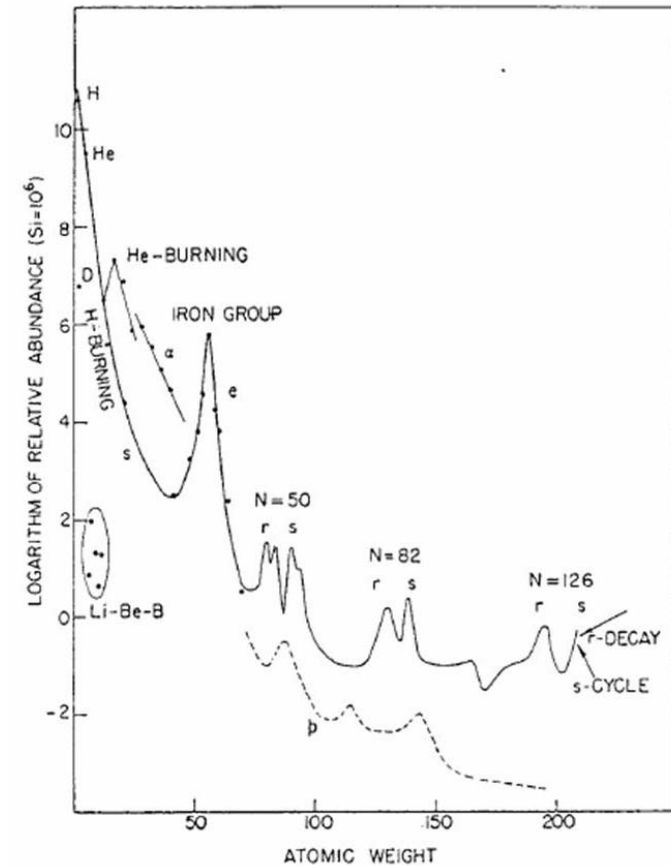
ΔS_{2p}

ΔS_{2n}

Nuclear properties change periodically with proton or neutron number

原子核性质随中子数或质子数周期性地变化

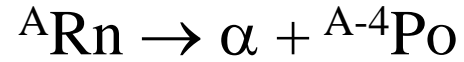
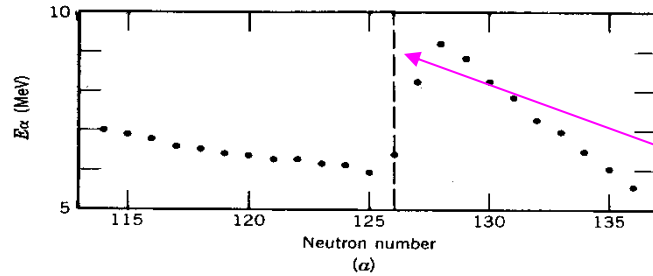
Abundance curve





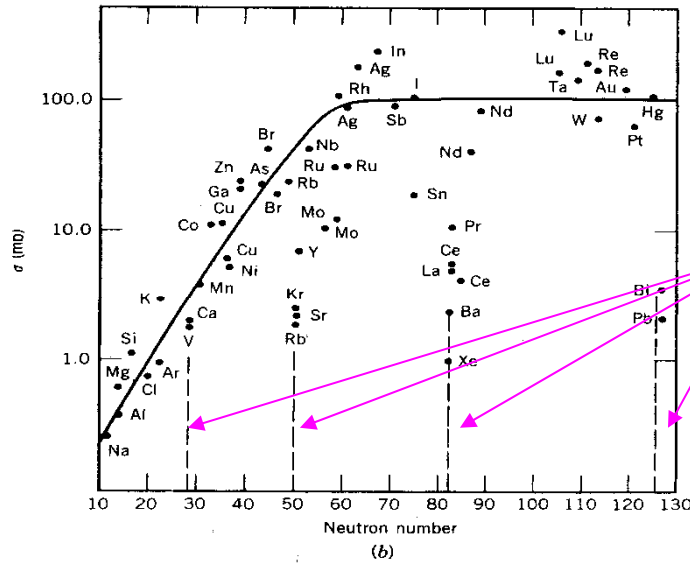
Nuclear magic number

T_α



Sudden rise at $N = 126$

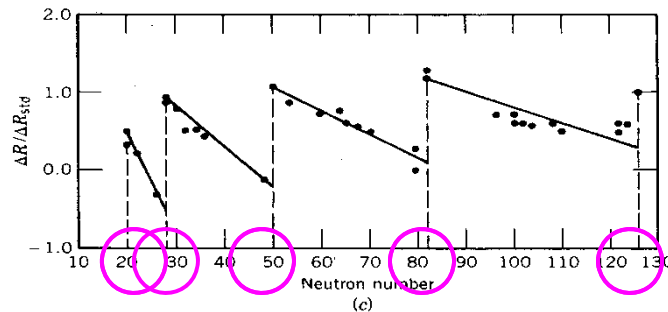
σ



Neutron capture cross section σ

Very small σ at $N = 28, 50, 82, 126$

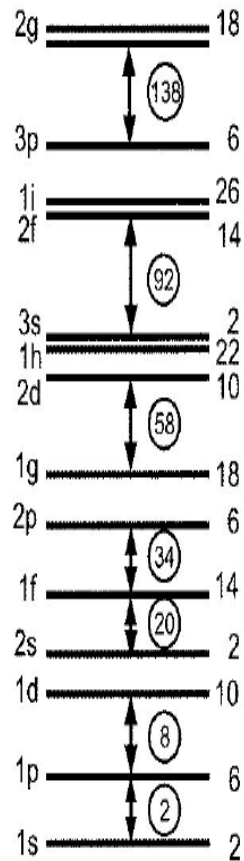
$\frac{\Delta R}{\Delta R_{avg}}$



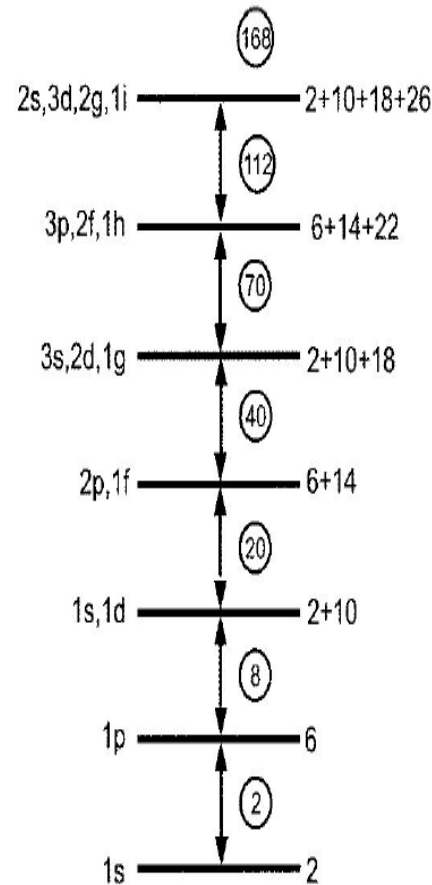
*Abrupt change in nuclear radius
at $N = 20, 28, 50, 82, 126$*



infinite square well potential



harmonic oscillator potential



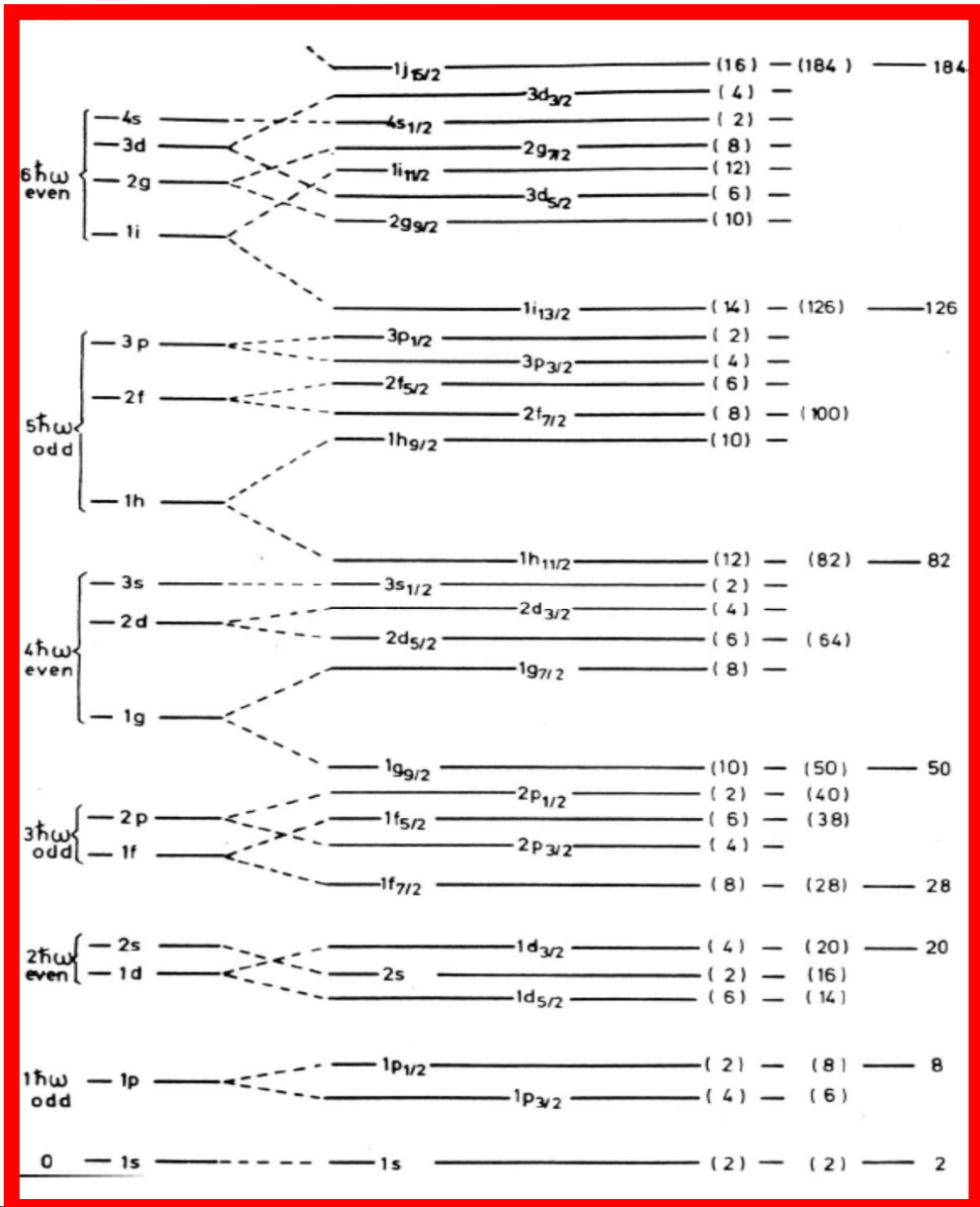
Only smallest magic numbers are reproduced. Why?

Spin-orbit term is necessary

Mayer and Jensen et al., 1949

$$V_{so} = -V_{ls} \frac{1}{r} \frac{\partial V(r)}{\partial r} \underline{l \cdot s}$$

M. G. Mayer, Phys. Rev. 75(1949)1969; 78(1950).
O. Haxel, J. H. D. Jensen, and H. E. Suess, Phys. Rev. 75(1949)1766; Z. Phys. 128(1950)295;



Great for:

- magic numbers
- ground state properties
- some low lying excited states

强自旋-轨道壳模型是原子核结构的所有微观理论的出发点

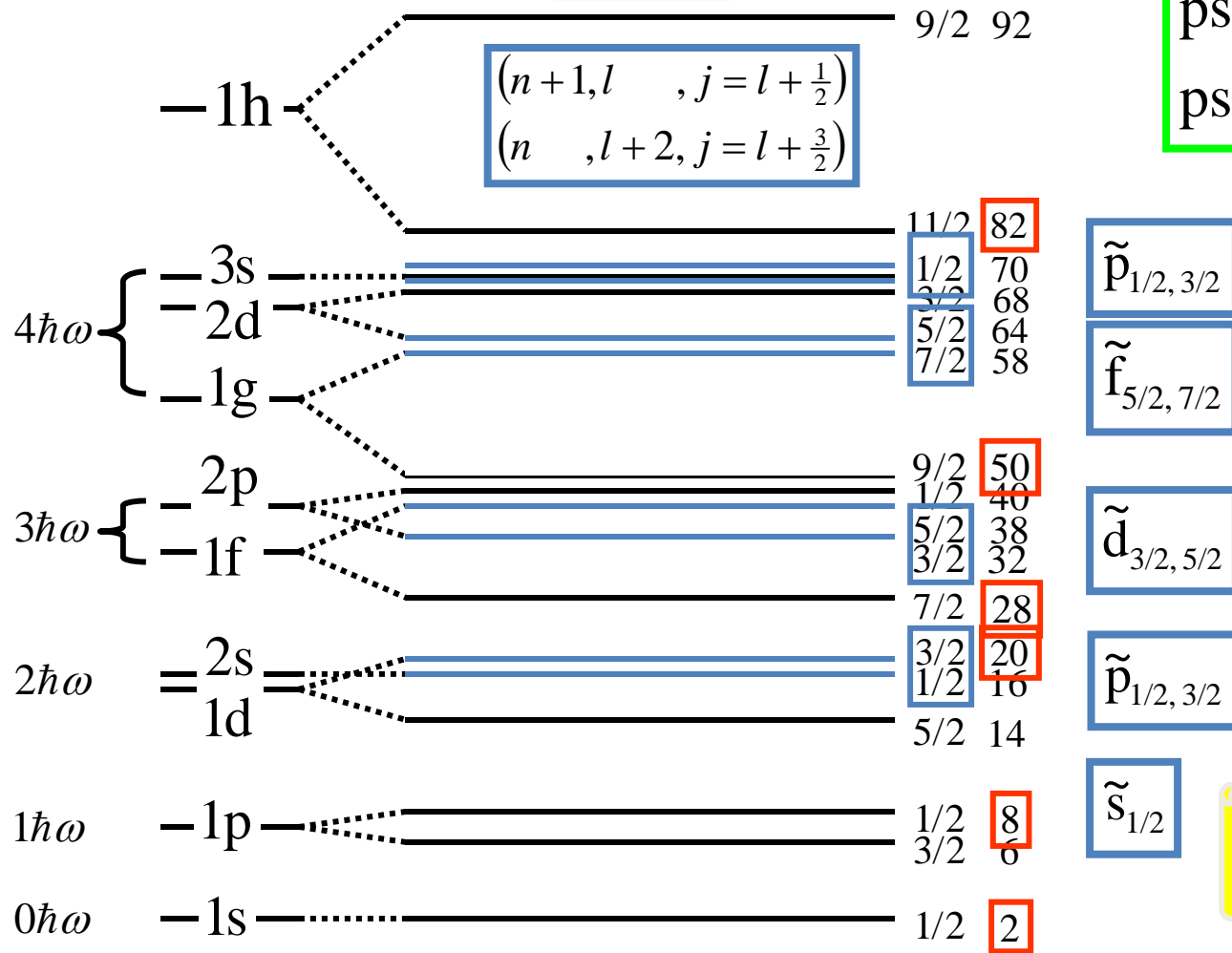
S. G. Nilsson的轴对称变形核的强自旋-轨道壳模型能级系

S. G. Nilsson, Mat. Fys. Medd. Dan. Vid. Selsk. 29, No.16(1955).
 S. G. Nilsson, et al., Nucl. Phys. A131(1969) 1.



Woods-Saxon

$$\kappa \vec{l} \cdot \vec{s}$$



pseudo-orbit : $\tilde{l} = l + 1$
pseudo-spin : $\tilde{s} = 1/2$

Hecht & Adler
NPA137(1969)129

Arima, Harvey & Shimizu
PLB30(1969)517



Physica Scripta. Vol. 26, 267–272, 1982

Pseudospin in Rotating Nuclear Potentials

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and

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Received July 7, 1982; accepted July 10, 1982

Abstract

It is found that the concept of pseudo-spin and pseudo-oscillator quantum numbers, introduced by previous authors, can be helpful in understanding the qualitative features of quasiparticle motion in rotating potentials, for the orbits with normal parity (excluding the high j orbits). In the first part of the paper, the use of pseudo-spin in static deformed potentials is briefly reviewed.



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Earlier attempt to understand PSS

VOLUME 68, NUMBER 14

PHYSICAL REVIEW LETTERS

6 APRIL 1992

Pseudospin Symmetry in Nuclear Physics

C. Bahri and J. P. Draayer

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Department of Physics, University of California, Los Angeles, Los Angeles, California 90024

(Received 12 November 1991)

The origin and consequences of pseudospin symmetry in nuclear physics, which is exact for an oscillator potential with one-body orbit-orbit (v_{ll}) and spin-orbit (v_{ls}) interaction strengths in the ratio $\mu \equiv 2v_{ll}/v_{ls} = 0.5$, are considered. Specifically, the $v_{ls} \approx 4v_{ll}$ condition is consistent with relativistic mean-field results and a pseudo LS coupling scheme. When deformation dominates, pseudospin extends to pseudo $SU(3)$, which is applicable to superdeformation.



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Since PSS was suggested, lots of nuclear phenomena have been interpreted in connection with the PSS, for example ...



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Nuclear superdeformed configurations

VOLUME 59, NUMBER 13

PHYSICAL REVIEW LETTERS

28 SEPTEMBER 1987

Abundance and Systematics of Nuclear Superdeformed States; Relation to the Pseudospin and Pseudo-SU(3) Symmetries

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(Received 31 December 1986)

Results of a multidimensional (β_2 , γ , β_4 , ω , Z , and N) search for nuclear superdeformed configurations are presented. Calculations based on a realistic deformed average field give a relatively strong dependence of the "super" elongation on the particle number. This dependence is shown to be a cyclic function of the particle number. It originates from the pseudospin and pseudo-SU(3) symmetries which are obeyed approximately in a realistic average field.

PACS numbers: 21.10.Gv, 21.60.Fw



Natural-Parity States in Superdeformed Bands and Pseudo SU(3) Symmetry at Extreme Conditions

W. Nazarewicz,⁽¹⁾ P. J. Twin,⁽²⁾ P. Fallon,⁽²⁾ and J. D. Garrett⁽³⁾

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(Received 18 December 1989)

RAPID COMMUNICATIONS

PHYSICAL REVIEW C

VOLUME 44, NUMBER 5

NOVEMBER 1991

Spin determination and quantized alignment in the superdeformed bands in ^{152}Dy , ^{151}Tb , and ^{150}Gd

J. Y. Zeng,^(1,2,3) J. Meng,^(2,3) C. S. Wu,^(1,2,3) E. G. Zhao,^(1,3)
Z. Xing,⁽⁴⁾ and X. Q. Chen⁽⁴⁾

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⁽⁴⁾*Department of Modern Physics, Lanzhou University, Lanzhou 730000, China*

(Received 13 May 1991)

The spin of the lowest level observed in the superdeformed band in ^{152}Dy is determined to be $I_0=25$ rather than the previously assigned $I_0=22$. As a result, we have $\mathcal{J}^{(1)} > \mathcal{J}^{(2)}$ and both $\mathcal{J}^{(1)}$ and $\mathcal{J}^{(2)}$ decrease very slowly with rotational frequency. From the spin determination of the two pairs of identical bands in ^{152}Dy and $^{151}\text{Tb}^*$, ^{151}Tb and $^{150}\text{Gd}^*$, evidence has been obtained to support the assumption of quantized alignments in units of $\hbar/2$.



PHYSICAL REVIEW C **80**, 034303 (2009)

Properties of the rotational bands in the transitional nucleus ^{189}Pt

W. Hua and X. H. Zhou*

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Graduate School of the Chinese Academy of Sciences, Beijing 100049, People's Republic of China*

Y. H. Zhang, Y. Zheng, M. L. Liu, F. Ma, S. Guo, L. Ma, S. T. Wang, N. T. Zhang, Y. D. Fang, X. G. Lei, and Y. X. Guo
Institute of

PHYSICAL REVIEW C **78**, 064301 (2008)

Identification of pseudospin partner bands in ^{108}Tc

Q. Xu (徐强),¹ S. J. Zhu (朱胜江),^{1,2,*} J. H. Hamilton,² A. V. Ramayya,² J. K. Hwang,² B. Qi (齐斌),³ J. Meng (孟杰),³
J. Peng (彭靖),⁴ Y. X. Luo,^{2,5} J. O. Rasmussen,⁵ I. Y. Lee,⁵ S. H. Liu,² K. Li,² J. G. Wang (王建国),¹ H. B. Ding (丁怀博),¹
L. Gu (顾龙),¹ E. Y. Yeoh (杨韵颐),¹ and W. C. Ma⁶

¹*Department of Physics, Tsinghua University, Beijing 100084, People's Republic of China*

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⁴*Department of Physics, Beijing Normal University, Beijing 100875, People's Republic of China*

⁵*Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA*

⁶*Department of Physics, Mississippi State University, Mississippi State, Mississippi 39762, USA*

(Received 31 July 2008; published 8 December 2008)

High-spin structures in the neutron-rich ^{108}Tc nucleus have been reinvestigated by measuring the prompt γ -rays from spontaneous fission of ^{252}Cf . A previously known collective band has been extended up to higher spin states and a new side band has been identified. These doublet bands are proposed as pseudospin partner bands with configurations $\pi 1/2^+[431] \otimes \nu[\widetilde{3}125/2^+, 3/2^+]$, which is a first identification in $A \sim 100$ region. The particle-rotor model (PRM) was applied to calculate levels and $B(M1)/B(E2)$ ratios of the bands in ^{108}Tc . The calculated results are in good agreement with the experimental values.

School of Physics, a

(Received

High-spin st
of 88 and 95 M
Rotational ban
shape and a ne
transition sequ
the $\nu i_{13/2}^{-2} \nu f_{5/2}$
as a pair of pse
theoretical calc



Even for Magnetic moments, transitions and γ -vibrational states

- ❑ On the validity of the pseudo-spin concept for axially symmetric deformed nuclei
D. Troltenier, W. Nazarewicz, Z. Szymański, J.P. Draayer, Nucl. Phys. A 567, 591 (1994).
- ❑ I-forbidden M1 transitions and pseudospin symmetry
P. von Neumann-Cosel and J. N. Ginocchio, Phys. Rev. C 62, 014308 (2000)
- ❑ Neutron number dependence of the energies of the γ -vibrational states in nuclei with $Z \sim 100$ and the manifestation of pseudospin symmetry
R. V. Jolos, N. Yi. Shirikova, and A. V. Sushkov, Phys. Rev. C 86, 044320 (2012)



Some themes in the study of very deformed rotating nuclei

Ben Mottelson

Nordita

Blegdamsvej 17

Very briefly the concept of pseudo-spin is built on the following observations

(Hecht and Adler, Nucl.Phys. A137, 129 (69); Arima, Harvey, and Shimizu, Phys.Lett. 30B, 517 (69)) :

- a) The observed single particle spectra in spherical nuclei exhibit a number of approximate degeneracies of orbits with the same parity and $\Delta j=1$:

$$\left. \begin{array}{l} (g_{7/2}, d_{5/2}) \\ (d_{3/2}, s_{1/2}) \end{array} \right\} \text{ for } N = 50 - 82$$

$$\left. \begin{array}{l} (h_{9/2}, f_{7/2}) \\ (f_{5/2}, p_{3/2}) \end{array} \right\} \text{ for } N = 82 - 126$$

$$\left. \begin{array}{l} (i_{11/2}, g_{9/2}) \\ (g_{7/2}, d_{5/2}) \\ (d_{3/2}, s_{1/2}) \end{array} \right\} \text{ for } N > 126$$

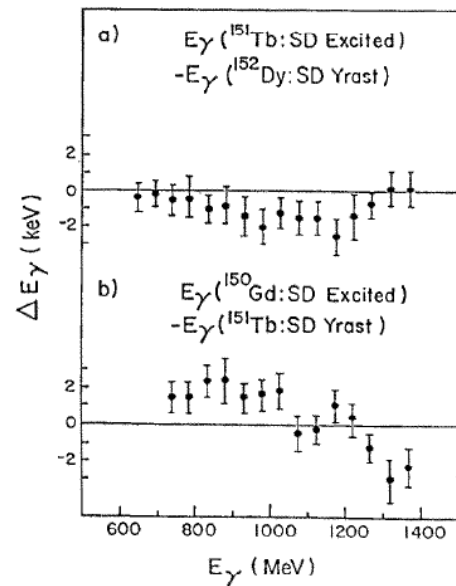


Fig.6: Comparison of transition energies for super-deformed rotational bands in ^{152}Dy , ^{151}Tb , and ^{150}Gd .



Relativistic mean field / Covariant density functional theory

$$\left[\boldsymbol{\alpha} \cdot \mathbf{p} + V(\mathbf{r}) + \beta (M + S(\mathbf{r})) \right] \psi_i = \varepsilon_i \psi_i$$

$$\begin{cases} V(\mathbf{r}) = g_\omega \omega(\mathbf{r}) + g_\rho \tau_3 \rho(\mathbf{r}) + e \frac{1 - \tau_3}{2} A(\mathbf{r}) \\ S(\mathbf{r}) = g_\sigma \sigma(\mathbf{r}) \end{cases}$$



$$\psi_{n\kappa m}^N(\mathbf{r}) = \frac{1}{r} \begin{pmatrix} i G_{n\kappa}(r) Y_{jm}^l(\Omega) \\ -F_{\tilde{n}\kappa}(r) Y_{jm}^{\tilde{l}}(\Omega) \end{pmatrix}$$

$n = \text{node number} + 1$

$$\begin{cases} j = l \pm 1/2 \\ \kappa = (-)^{j+l+1/2} (j + 1/2) \\ \tilde{l} = l + (-)^{j+l-1/2} \end{cases}$$

$$(2s_{1/2}, 1d_{3/2}) \Rightarrow (\tilde{p}_{1/2,3/2})$$

$$(\tilde{n} = 2) \tilde{p}_{1/2,3/2}$$

$$2s_{1/2} = \begin{pmatrix} n = 2, l = 0, j = l + \frac{1}{2} \\ \tilde{n} = 2, \tilde{l} = 1, j = \tilde{l} - \frac{1}{2} \end{pmatrix}$$

$$1d_{3/2} = \begin{pmatrix} n = 1, l = 2, j = l - \frac{1}{2} \\ \tilde{n} = 2, \tilde{l} = 1, j = \tilde{l} + \frac{1}{2} \end{pmatrix}$$

Pseudo quantum numbers are nothing
but the quantum numbers of the lower component.

Ginocchio PRL78(97)436



$$V_{\pm}(r) = V(r) \pm S(r)$$

$$M_{\pm}(\varepsilon_N, r) = M \pm \varepsilon_N \mp V_{\mp}(r)$$

$$\psi_{n\kappa m}^N(\mathbf{r}) = \frac{1}{r} \begin{pmatrix} i G_{n\kappa}(r) Y_{jm}^l(\Omega) \\ -F_{\tilde{n}\kappa}(r) Y_{jm}^{\tilde{l}}(\Omega) \end{pmatrix}$$

$$\left[-\frac{1}{2M_+} \left(\frac{d^2}{dr^2} + \frac{1}{2M_+} \frac{dV_-}{dr} \frac{d}{dr} - \frac{l(l+1)}{r^2} \right) - \frac{1}{4M_+^2} \frac{\kappa}{r} \frac{dV_-}{dr} + M - V_+ \right] G = +\varepsilon_N G$$

$$\left[-\frac{1}{2M_-} \left(\frac{d^2}{dr^2} - \frac{1}{2M_-} \frac{dV_+}{dr} \frac{d}{dr} + \frac{\tilde{l}(\tilde{l}+1)}{r^2} \right) + \frac{1}{4M_-^2} \frac{\tilde{\kappa}}{r} \frac{dV_+}{dr} + M - V_- \right] F = -\varepsilon_N F$$

For nucleons,

- ✘ $d[V(r)-S(r)]/dr=0 \Rightarrow$ spin symmetry
- ✓ $d[V(r)+S(r)]/dr=0 \Rightarrow$ pseudo-spin symmetry

Meng, Sugawara-Tanabe, Yamaji, Ring, Arima, PRC58 (1998) 628R
Meng, Sugawara-Tanabe, Yamaji, Arima, PRC59 (1999) 154



$$V_{\pm}(r) = V(r) \pm S(r)$$

$$M_{\pm}(\varepsilon_A, r) = M \mp \varepsilon_A \mp V_{\mp}(r)$$

$$\psi_{n\kappa m}^A(\mathbf{r}) = \frac{1}{r} \begin{pmatrix} -F_{n\kappa}(r) Y_{jm}^l(\Omega) \\ i G_{\tilde{n}\kappa}(r) Y_{jm}^{\tilde{l}}(\Omega) \end{pmatrix}$$

$$\left[-\frac{1}{2M_-} \left(\frac{d^2}{dr^2} - \frac{1}{2M_-} \frac{dV_+}{dr} \frac{d}{dr} + \frac{l(l+1)}{r^2} \right) + \frac{1}{4M_-^2} \frac{\kappa}{r} \frac{dV_+}{dr} + M - V_- \right] F = -\varepsilon F$$

$$\left[-\frac{1}{2M_+} \left(\frac{d^2}{dr^2} + \frac{1}{2M_+} \frac{dV_-}{dr} \frac{d}{dr} - \frac{\tilde{l}(\tilde{l}+1)}{r^2} \right) - \frac{1}{4M_+^2} \frac{\tilde{\kappa}}{r} \frac{dV_-}{dr} + M - V_+ \right] G = +\varepsilon G$$

For anti-nucleons,

- ✘ $d[V(r)-S(r)]/dr=0 \Rightarrow$ pseudo-spin symmetry
- ✓ $d[V(r)+S(r)]/dr=0 \Rightarrow$ spin symmetry

Zhou, Meng & Ring PRL92(03)262501



For nucleons, the smaller component F

Zhou, Meng & Ring
PRL92(03)262501

$$\left[-\frac{\vec{\nabla}^2}{2M_-^*} - [V(r) - S(r)] + \frac{1}{4M_-^2} \frac{1}{r} \frac{d}{dr} [V(r) + S(r)] \left(1 + \vec{l} \cdot \vec{\sigma} \right) + M \right] F = -\varepsilon_N F$$

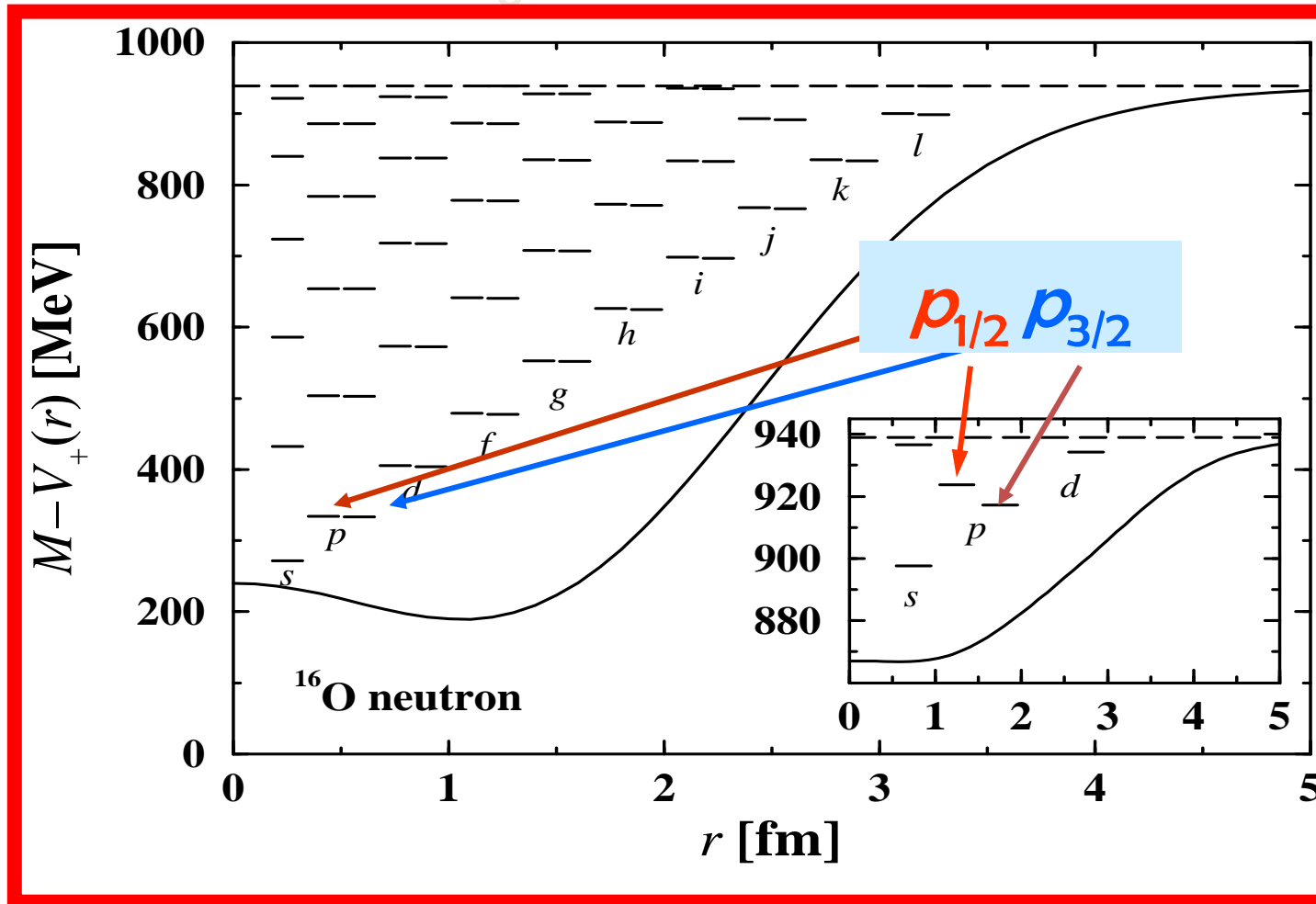
For anti-nucleons, the larger component F

$$\left[-\frac{\vec{\nabla}^2}{2M_-^*} - [V(r) - S(r)] - \frac{1}{4M_-^2} \frac{1}{r} \frac{d}{dr} [V(r) + S(r)] \left(1 + \vec{l} \cdot \vec{\sigma} \right) + M \right] F = +\varepsilon_A F$$

The factor $\frac{1}{4M_-^2}$ is ~400 times smaller for anti nucleons!

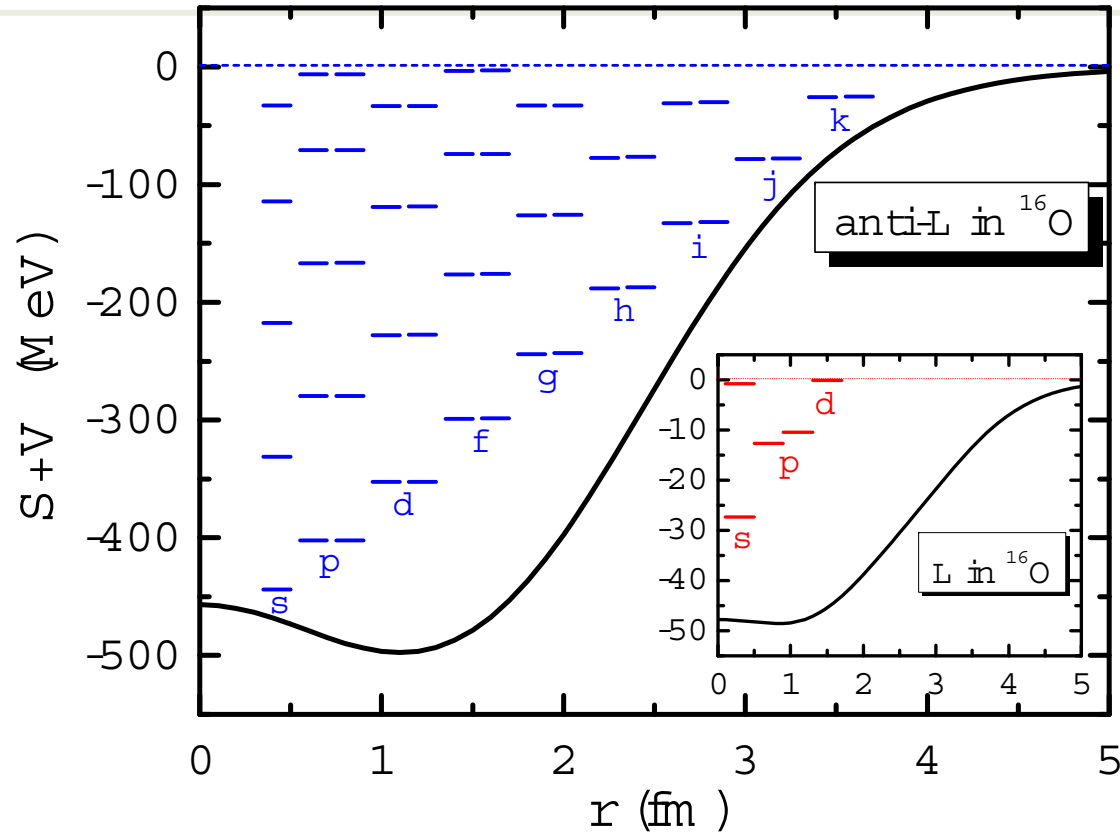


Zhou, Meng & Ring, PRL92 (03) 262501





Song, Yao, Meng, CHIN. PHYS. LETT. Vol. 26, No. 12 (2009) 122102



Experiment: exploration of the relativistic symmetry by hyperon



Kenichiro Arita (Nagoya Institute of Technology)

Haozhao Liang (RIKEN)

Shan-Gui Zhou (Institute of Theoretical Physics, CAS, China)

...

- Origins of nuclear shell structure
- Spin symmetry in Dirac negative energy spectrum
- Perturbative interpretation of relativistic symmetries possible
- Pseudospin symmetry in resonances
- Pseudospin symmetry in supersymmetric quantum mechanics
- Pseudospin for deformed nucleus
- ...

- Liang, Shen, Zhao, Meng, Phys. Rev. C 87, 014334 (2013) [13 pages]
- B.-N. Lu, E.-G. Zhao, and S.-G. Zhou, Phys. Rev. Lett. 109, 072501 (2012).
- Bing-Nan Lu, En-Guang Zhao, Shan-Gui Zhou arXiv:1305.1524

Thank you for your attention!