

# Ab-initio study of nuclear four-body reactions

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## Correlated Gaussian + Microscopic R-matrix Method

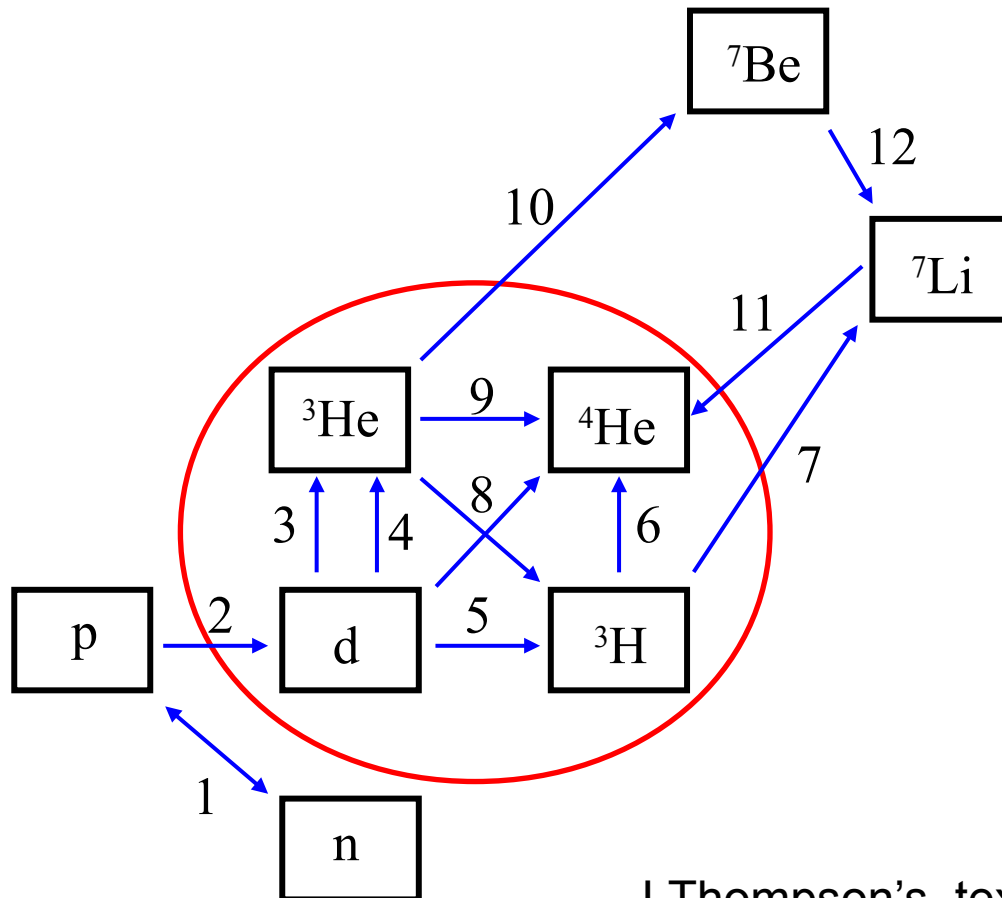
P. Descouvemont, D. Baye, Y. Suzuki, S. Aoyama, K. Arai, AIP ADVANCES, (2014)041011.

S. Aoyama, K. Arai, Y. Suzuki, P. Descouvemont, D. Baye, FBS52, (2012)97.

K. Arai, S. Aoyama, Y. Suzuki, P. Descouvemont, D. Baye, PRL107 (2011) 132502.

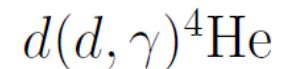
# Dominant reactions in primordial nucleosynthesis

Normally, the primordial nucleosynthesis is explained by the reaction chain calculation which is based on a simple nuclear model or a mere extrapolation from experiments.



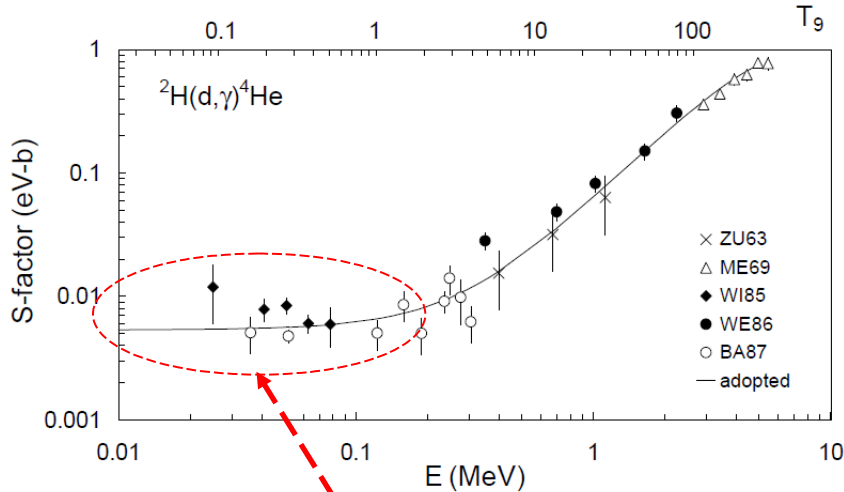
I.Thompson's textbook

- 1:  $n \leftrightarrow p$
- 2:  $p(n, \gamma)d$
- 3:  $d(p, \gamma)^3\text{He}$
- 4:  $d(d, n)^3\text{He}$
- 5:  $d(d, p)^3\text{H}$
- 6:  $^3\text{H}(d, n)^4\text{He}$
- 7:  $^3\text{H}(^4\text{He}, \gamma)^3\text{H}$
- 8:  $^3\text{He}(n, p)^3\text{H}$
- 9:  $^3\text{He}(d, p)^4\text{He}$
- 10:  $^3\text{He}(^4\text{He}, \gamma)^7\text{Be}$
- 11:  $^7\text{Li}(p, ^4\text{He})^4\text{He}$
- 12:  $^7\text{Be}(n, p)^7\text{Li}$



Can we understand these reactions in *ab-initio* way?  
Is there some effects of tensor interaction in Big-Bang?

# Effect of tensor interaction in low energy region



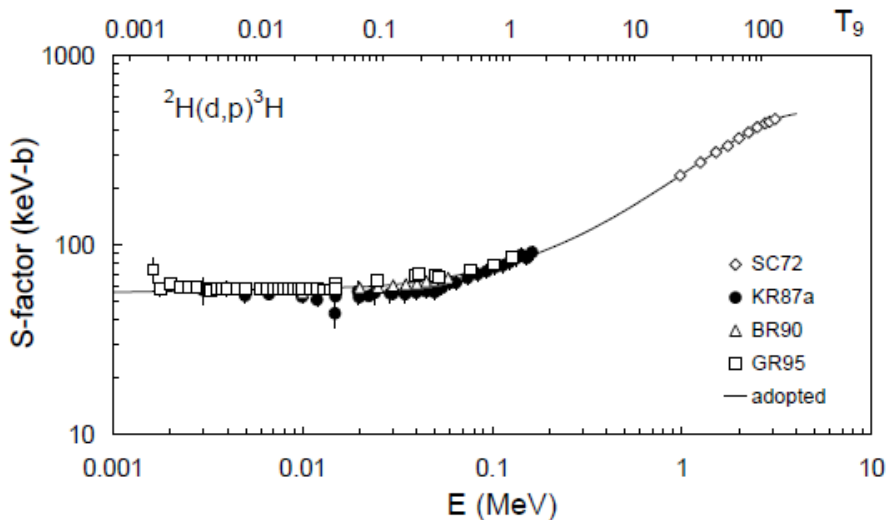
Nacre compilation

( C.Angulo et al.NPA656(1999)3)

${}^2\text{H}(d,\gamma){}^4\text{He}$  Astrophysical  
S-factor

$d + d$  S-wave  $\rightarrow$  D-state in the  $0^+$  g.s. of  ${}^4\text{He}$

H. J. Assenbaum and K. Langanke,  
PRC36(1987)17



${}^2\text{H}(d, p){}^3\text{H}$ ,  ${}^2\text{H}(d, n){}^3\text{He}$

Effect of Tensor Interaction?

# Hamiltonian(4-body case)

$$H = \sum_{i=1}^4 T_i - T_{\text{cm}} + \sum_{i<j}^4 V_{ij} + \sum_{i<j<k}^4 V_{ijk},$$

## Realistic Interaction: AV8' (+Coulomb+3NF)

$V_{ij}$ : Central+LS+Tensor+Coulomb

Pudliner, Pandharipande, Carlson, Pieper, Wiringa: PRC56(1997)1720

$V_{ijk}$ : Effective three nucleon force

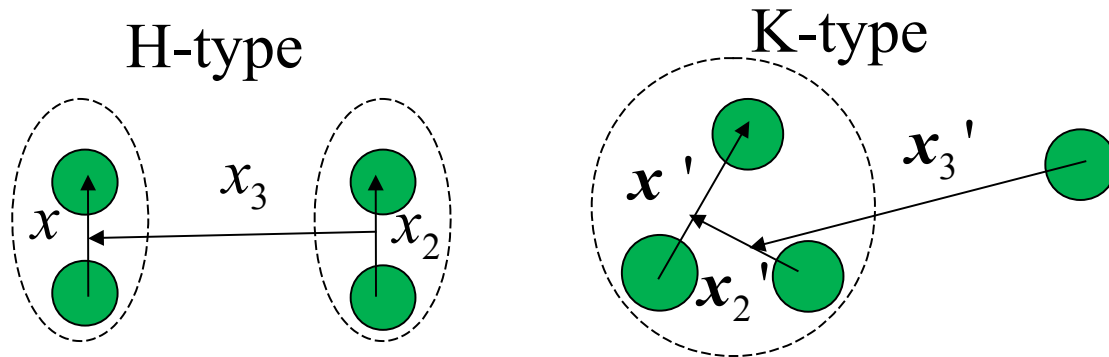
Hiyama, Gibson, Kamimura, PRC 70(2003)031001

## Effective Interaction: MN (+Coulomb)

$V_{ij}$ : Central+Coulomb

Thompson, LeMere, Tang, NPA(1977)286

Gaussian  
Basis  
Function



# Correlated Gaussian function with triple global vectors for four nucleon system

Unnatural parity 0-

$L_1=L_2=L_{12}=L_3=1$

$$F_{L_1 L_2 (L_{12}) L_3 L M}(u_1, u_2, u_3, A, \mathbf{x})$$

$$= \exp\left(-\frac{1}{2} \tilde{\mathbf{x}} A \mathbf{x}\right) \left[ \left[ \mathcal{Y}_{L_1}(\tilde{u}_1 \mathbf{x}) \mathcal{Y}_{L_2}(\tilde{u}_2 \mathbf{x}) \right]_{L_{12}} \mathcal{Y}_{L_3}(\tilde{u}_3 \mathbf{x}) \right]_{L M}$$

Single gloval vector

Double global vector

New extension(triple)

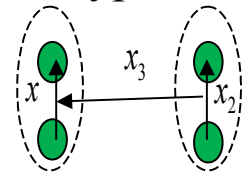
$$\mathcal{Y}_{L_i M_i}(\tilde{u}_i \mathbf{x}) = |\tilde{u}_i \mathbf{x}|^{L_i} Y_{L_i M_i}(\widehat{\tilde{u}_i \mathbf{x}}) \quad \tilde{u}_i \mathbf{x} = \sum_{j=1}^{N-1} (u_i)_j \mathbf{x}_j$$

For H-type, we can choose,  $\tilde{u}_1=(1,0,0)$ ,  $\tilde{u}_2=(0,1,0)$  and  $\tilde{u}_3=(0,0,1)$

We also write the K-type basis function in the same form.

$$\exp\left(-\frac{1}{2} \tilde{\mathbf{x}}' A_K \mathbf{x}'\right) \left[ \left[ \mathcal{Y}_{L_1}(\mathbf{x}'_1) \mathcal{Y}_{L_2}(\mathbf{x}'_2) \right]_{L_{12}} \mathcal{Y}_{L_3}(\mathbf{x}'_3) \right]_{L M}$$

H-type

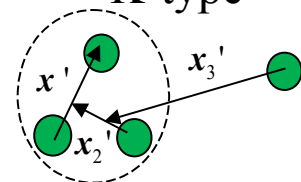


$$\mathbf{x}' = \tilde{U}_{KH} \mathbf{x}$$

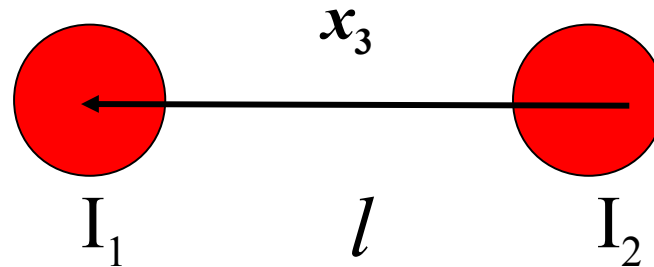
$$\tilde{u}_1=(1,0,0), \tilde{u}_2=(0, -\frac{1}{2}, 1) \text{ and } \tilde{u}_3=(0, \frac{2}{3}, \frac{2}{3})$$

$$A = (u_1 u_2 u_3) A_K \begin{pmatrix} \tilde{u}_1 \\ \tilde{u}_2 \\ \tilde{u}_3 \end{pmatrix} = \widetilde{U}_{KH} A_K U_{KH}$$

K-type



# Microscopic R-matrix Method



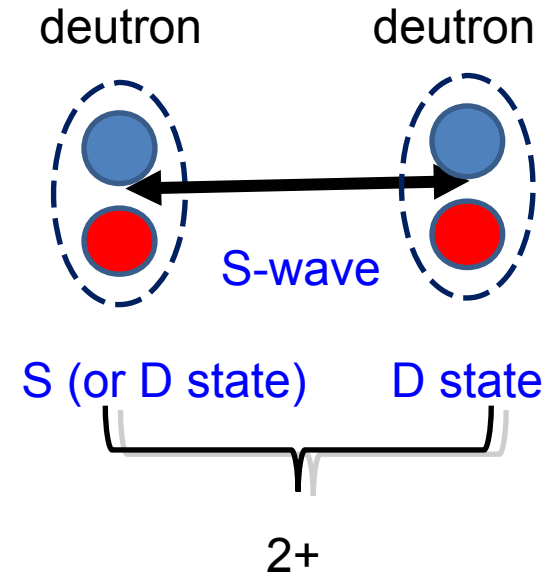
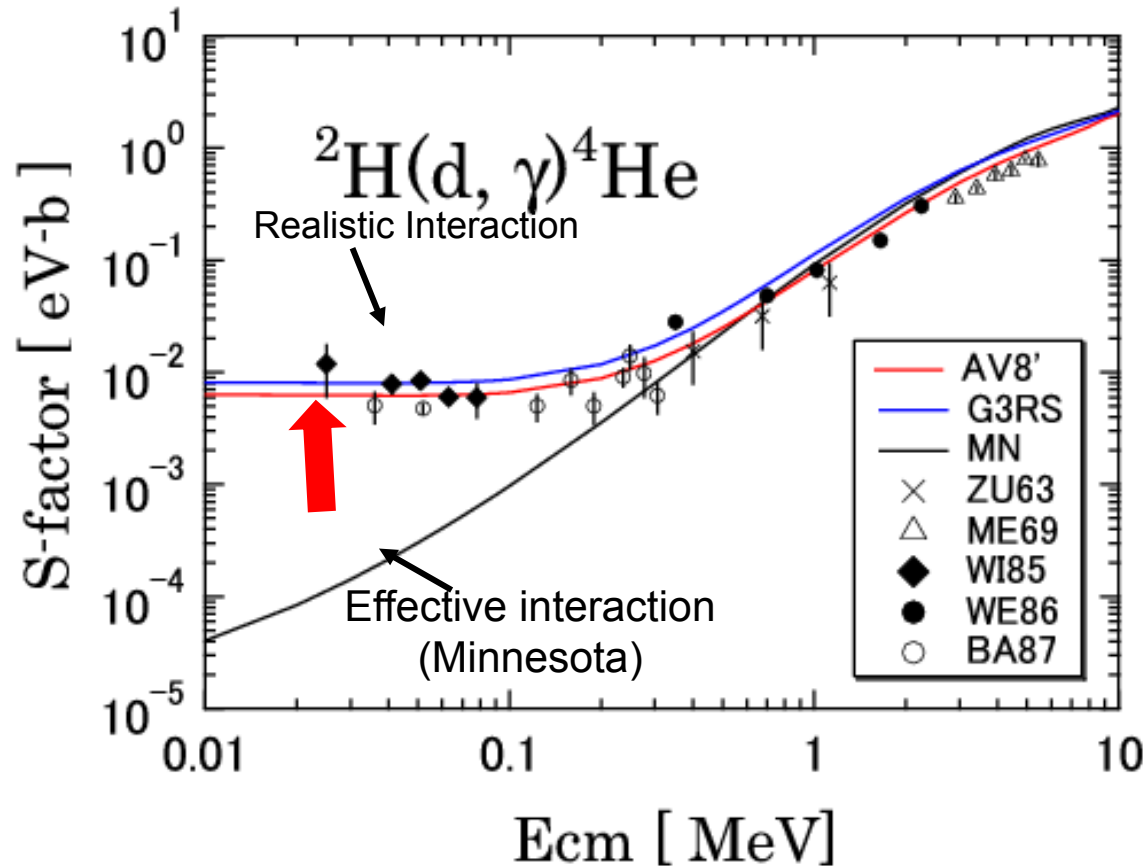
$$[[I_1 I_2]_I l]_{JM}$$

- $a$  : channel radius (13-15fm)
- $x_3 < a$  --- Gaussian expansion
- $x_3 > a$  ---  $I_1(ka) \delta_{\alpha\alpha'} - S_{\alpha\alpha'} O_1(ka)$  or  $W_{l+1/2,\eta}(2ka)$

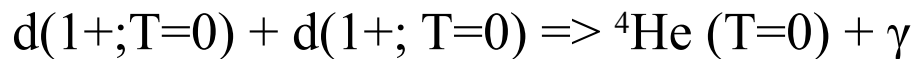
e.g. D. Baye, P. -H.Heenen, M. Libert-Heinemann, NPA291(1977).

# Radiative capture

K. Arai, S. Aoyama, Y. Suzuki, P. Descouvemont and D. Baye, PRL107 (2011) 132502.

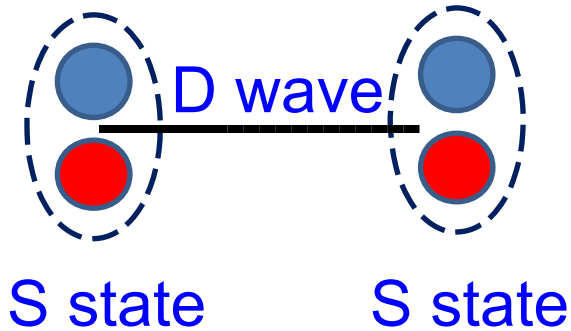


E2 transition is not reduced so much because of d-wave component.

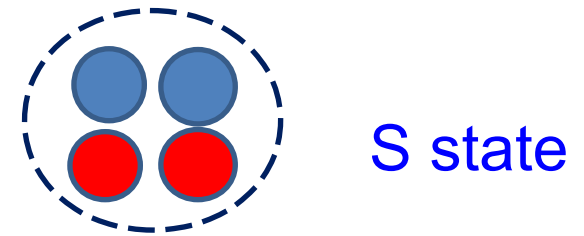


We can add a new evidence of D-wave components (tensor) of deuteron and  ${}^4\text{He}$ .

2+ continuum state  
Centrifugal Barrier



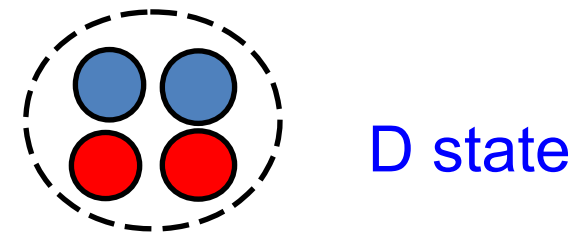
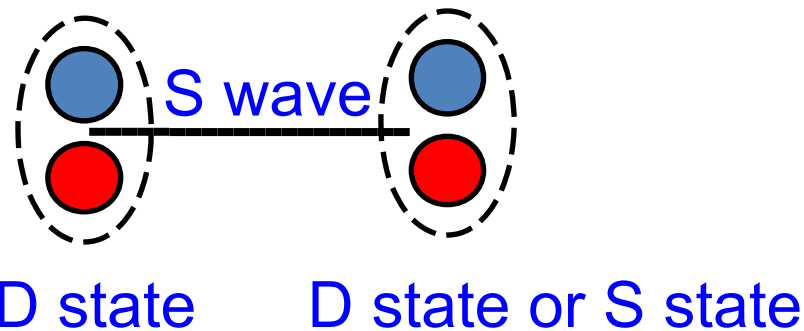
G.S. state of 4He (0+)



E2 transition



No centrifugal Barrier



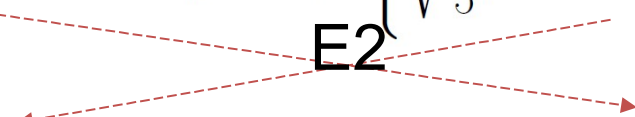
Initial dd state in S-wave

$$|^5S_2 : J = 2 \rangle \sim (1 - P_D(d))|L = 0, S = 2 \rangle + \sqrt{2P_D(d)} \left\{ \sqrt{\frac{1}{5}}|L = 2, S = 0 \rangle + \dots \right\}$$

Final  $^4\text{He}$  state

$$|^4\text{He} : J = 0 \rangle \sim \sqrt{1 - P_D(\alpha)}|L = 0, S = 0 \rangle + \sqrt{P_D(\alpha)}|L = 2, S = 2 \rangle$$

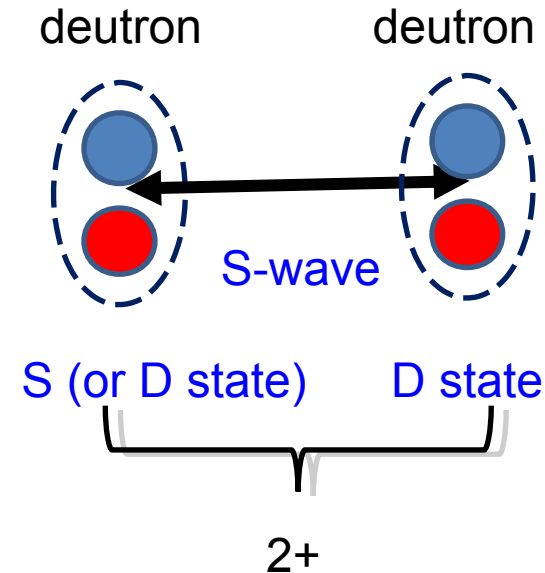
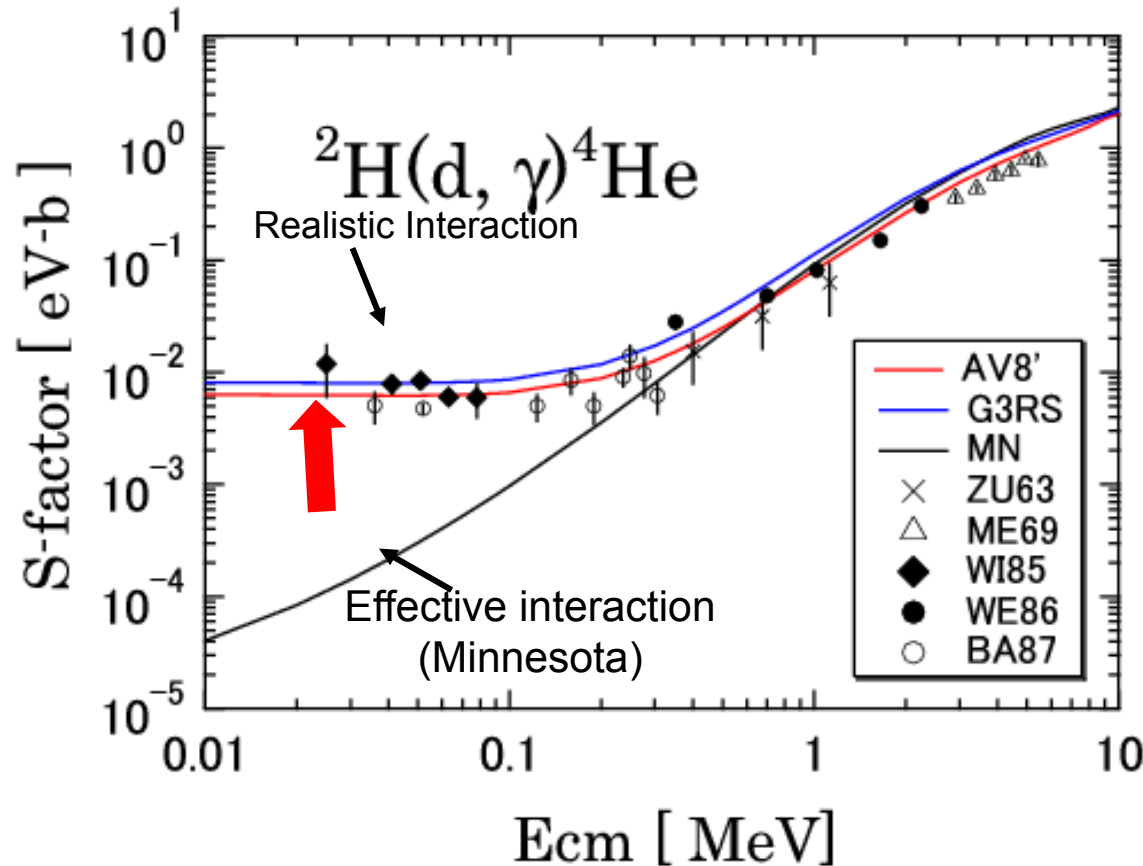
E2



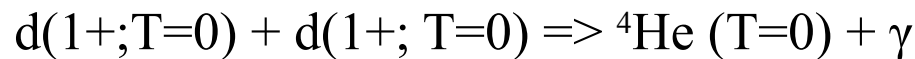


# Radiative capture

K. Arai, S. Aoyama, Y. Suzuki, P. Descouvemont and D. Baye, PRL107 (2011) 132502.



E2 transition is not reduced so much because of d-wave component.



# E1 transition

First term of E1 is an iso-vector operator

$$\begin{aligned}\mathcal{M}_{1\mu}^E &= e \sum_{i_1}^4 g_l^{(i)} r_i Y_{1\mu}(\hat{\mathbf{r}}_i) \\ &\approx e \sum_{i=1}^4 t_{i3} \mathbf{r}_i \\ &\propto \mathbf{R}_{c.m.}^n - \mathbf{R}_{c.m}^p\end{aligned}$$

Second term of E1 is an iso-scalar

$$\mathcal{M}_{1\mu}^E \approx -e \sum_i^A t_{i3} r_i Y_{1\mu}(\hat{\mathbf{r}}_i) - e \frac{k^2}{60} \sum_i^A r_i^3 Y_{1\mu}(\hat{\mathbf{r}}_i)$$

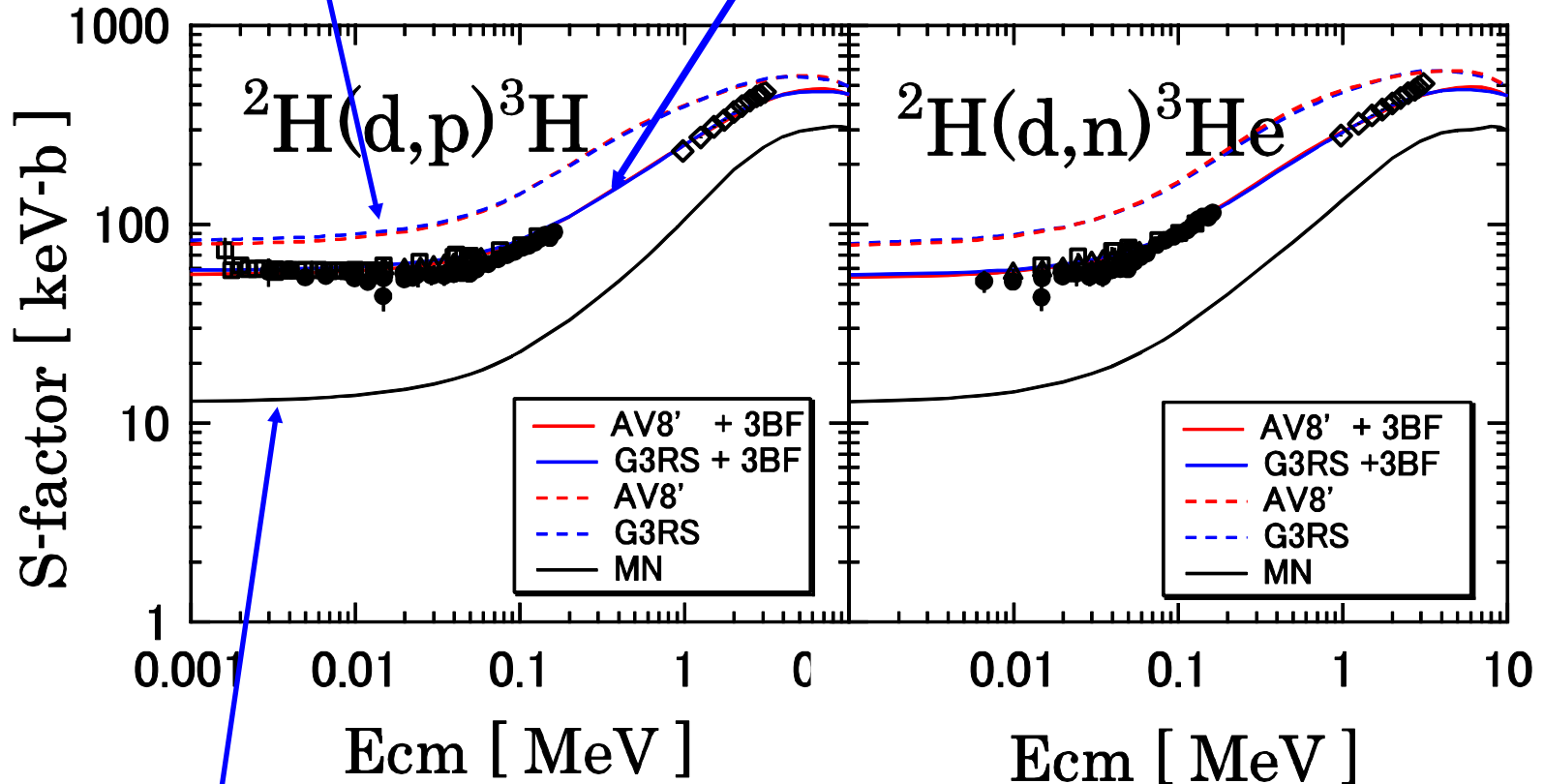
D.Baye, PRC86(2012)039306

order is  $1/30 \times (kr)^2 \times E1$

# Transfer reaction

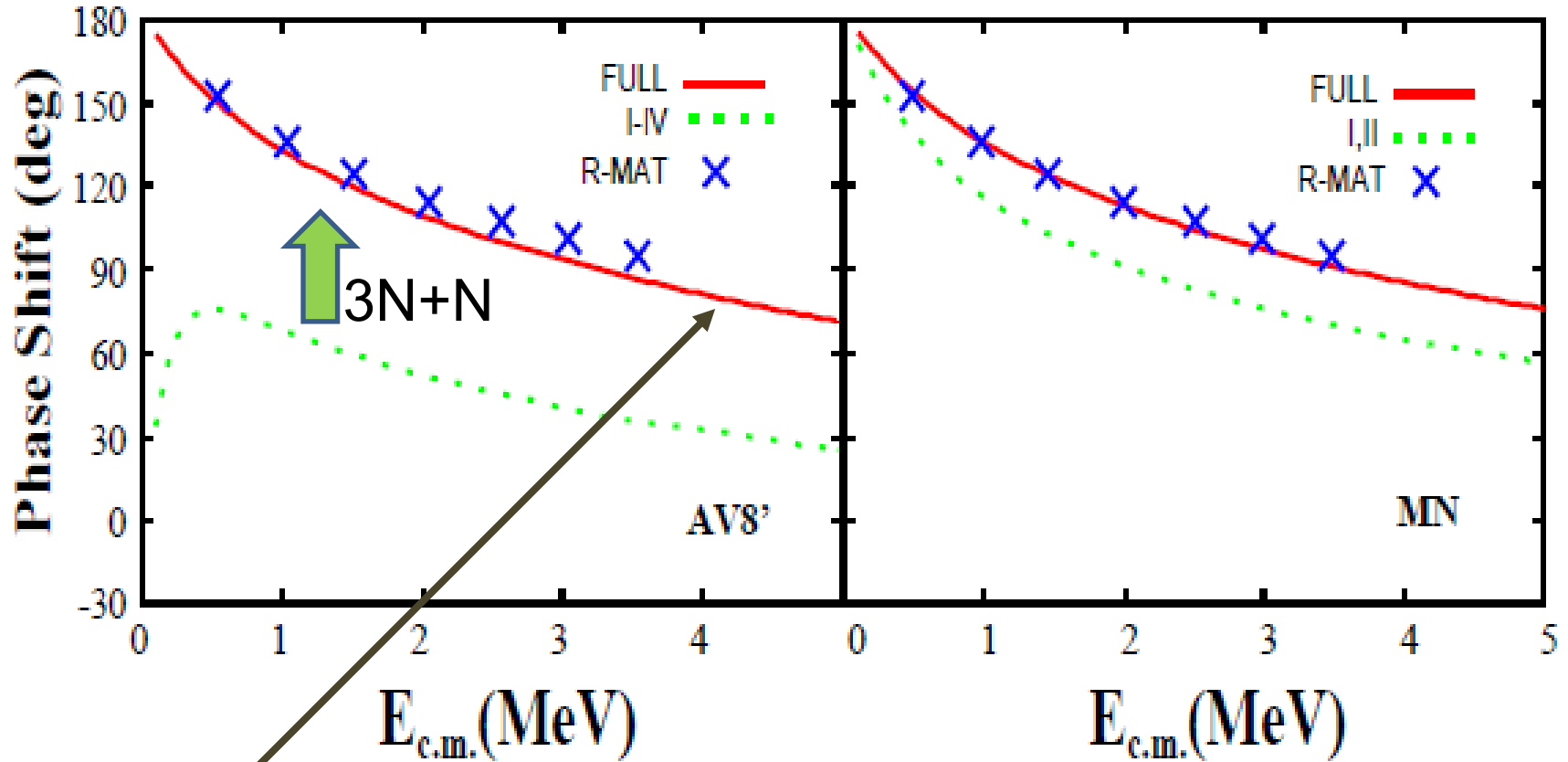
Realistic Interaction  
without 3NF

Realistic Interaction



Effective interaction

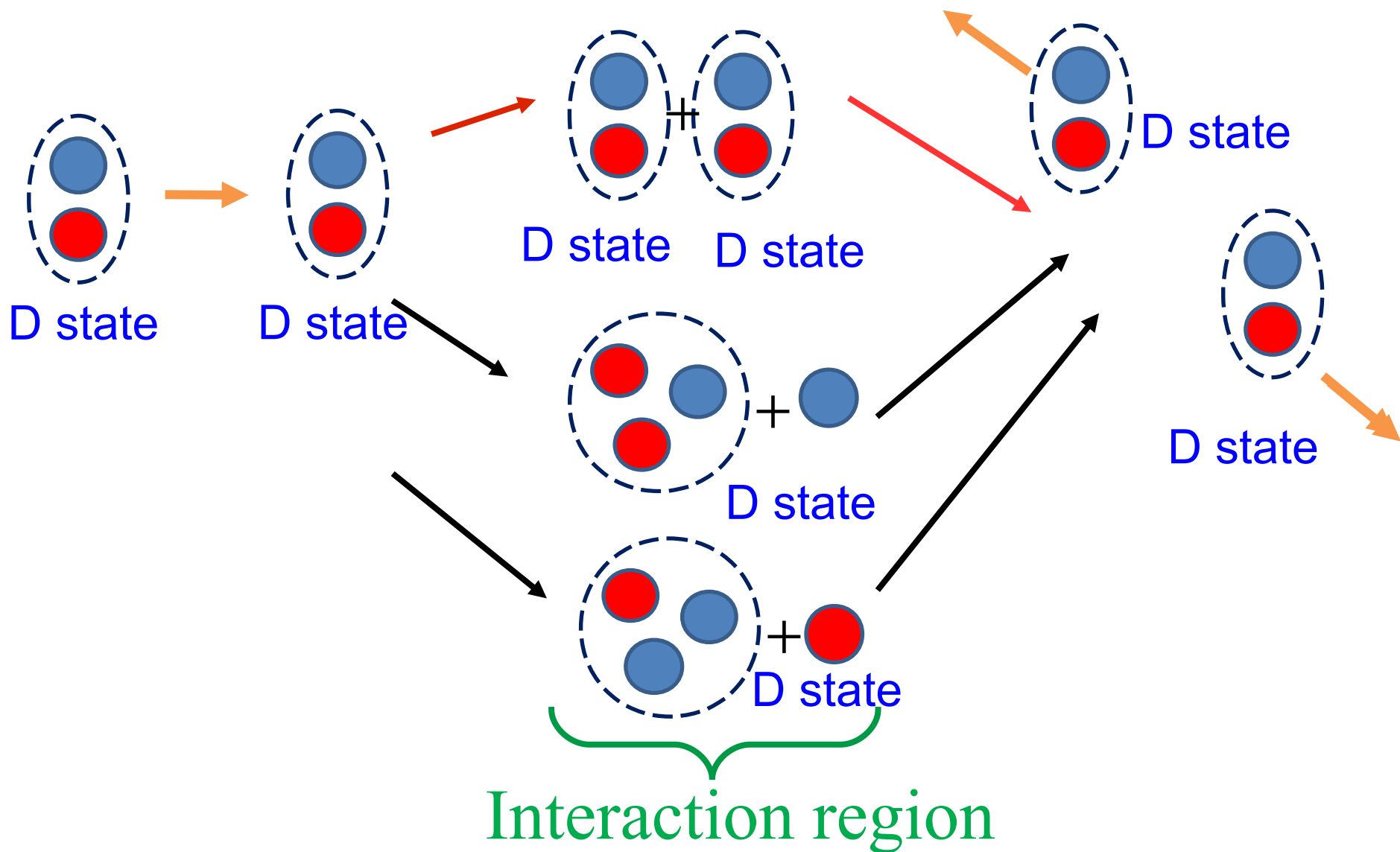
## $^1S_0$ d+d elastic phase shift (0+)



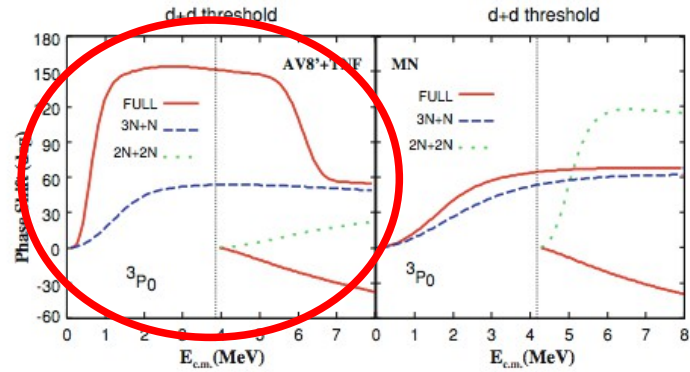
For effective interaction, d+d scattering picture is good!

# Coupling between d+d channel and 3N+N channels

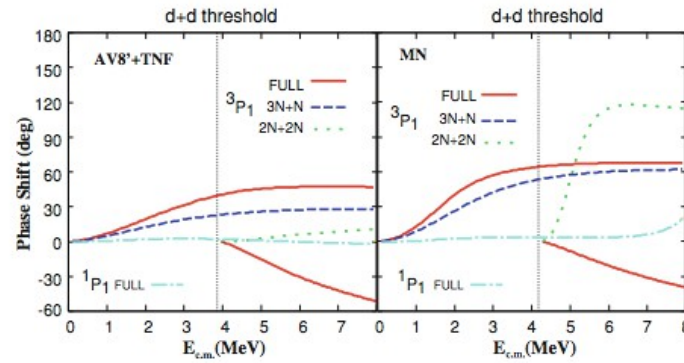
Tensor force makes the coupling in the scattering strong



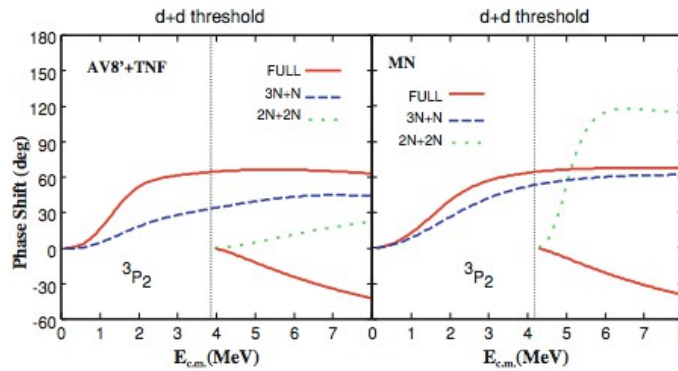
# Realistic int.    Effective int.



0-

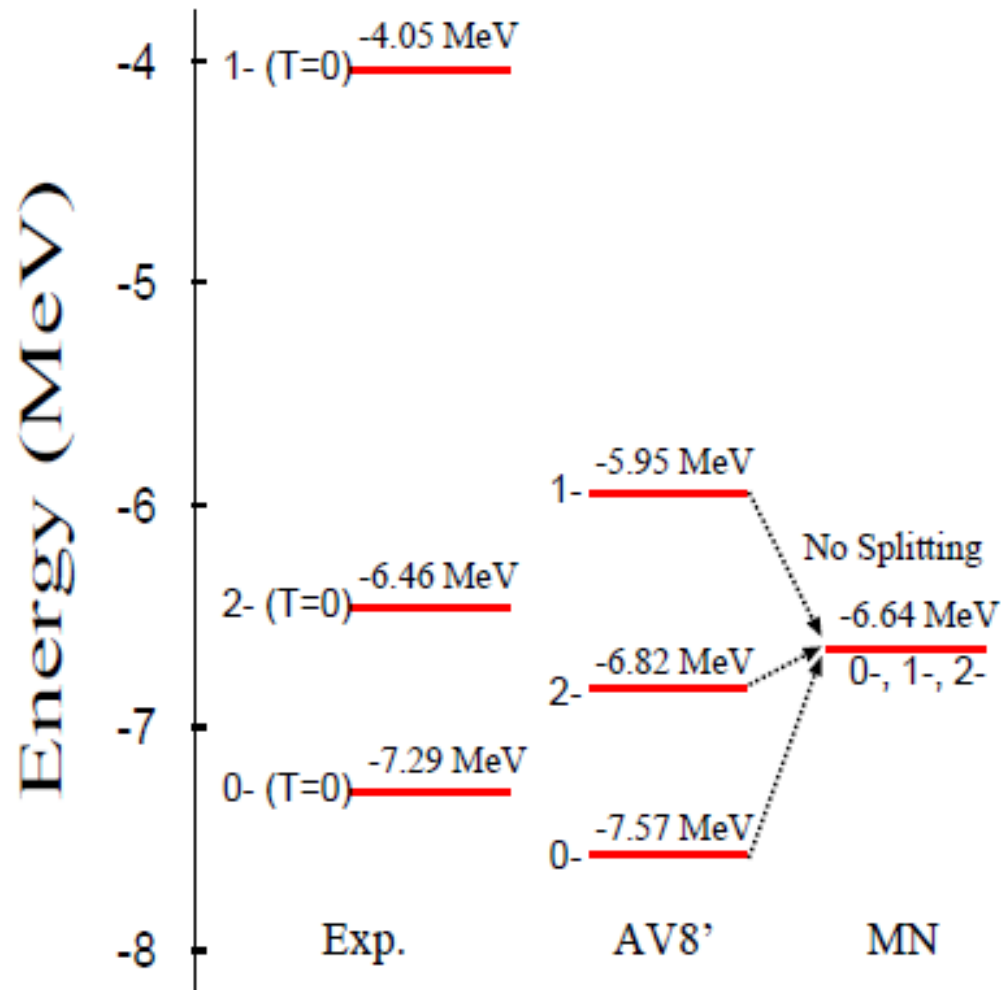


1-



2-

# Energy levels for negative parity states



Effective interaction (MN) gives same phase shift for 0-.1-.2- !

# Included channels in the present calculation

model		channel	
FULL	2N+2N	I	$d(1^+)+d(1^+)$
			$d(1^+)+d^*(1^+)$
			$d^*(1^+)+d^*(1^+)$
		II	$\bar{d}(0^+)+\bar{d}(0^+)$
			$\bar{d}(0^+)+d^*(0^+)$
			$d^*(0^+)+d^*(0^+)$
		III	$d^*(2^+)+d^*(1^+)$
			$d^*(2^+)+d^*(2^+)$
			$d^*(3^+)+d^*(1^+)$
		IV	$d^*(3^+)+d^*(2^+)$
			$d^*(3^+)+d^*(3^+)$
			$2n(0^+)+2p(0^+)$
		V	$2n(0^+)+2p^*(0^+)$
			$2n^*(0^+)+2p(0^+)$
			$2n^*(0^+)+2p^*(0^+)$
1	$t(\frac{1}{2}^+)+p(\frac{1}{2}^+)$		
	$t^*(\frac{1}{2}^+)+p(\frac{1}{2}^+)$		
2	$h(\frac{1}{2}^+)+n(\frac{1}{2}^+)$		
	$h^*(\frac{1}{2}^+)+n(\frac{1}{2}^+)$		

Thanks to the reduction of basis function by SVM for the sub-system. We can reduce the dimension of matrix elements very much!

Dimensions of matrix elements for FULL in the LS-coupled case

	N
0+	6660
1+	16680
2+	22230
0-	4200
1-	11670
2-	12480

The number of M.E. is  $N(N+1)/2$ .

For 2+, it takes about 200 days with 1CPU(1Core). And we need about 20Gbyte memory for the MRM calculation(half day).

All pseudo states (discretized continuum state) are employed in the MRM calculation.



# Summary

1. For astrophysical S-factor in  $d(d, \gamma)^4\text{He}$ ,  $d(d, p)^3\text{H}$ ,  $d(d, n)^3\text{He}$  reactions, tensor interaction is important to reproduce experiment.
2. In the  $d+d$  elastic reaction, the breaking of deuteron due to tensor interaction is large.

## Next

5-nucleon systems