

hadron spectroscopy from lattice QCD

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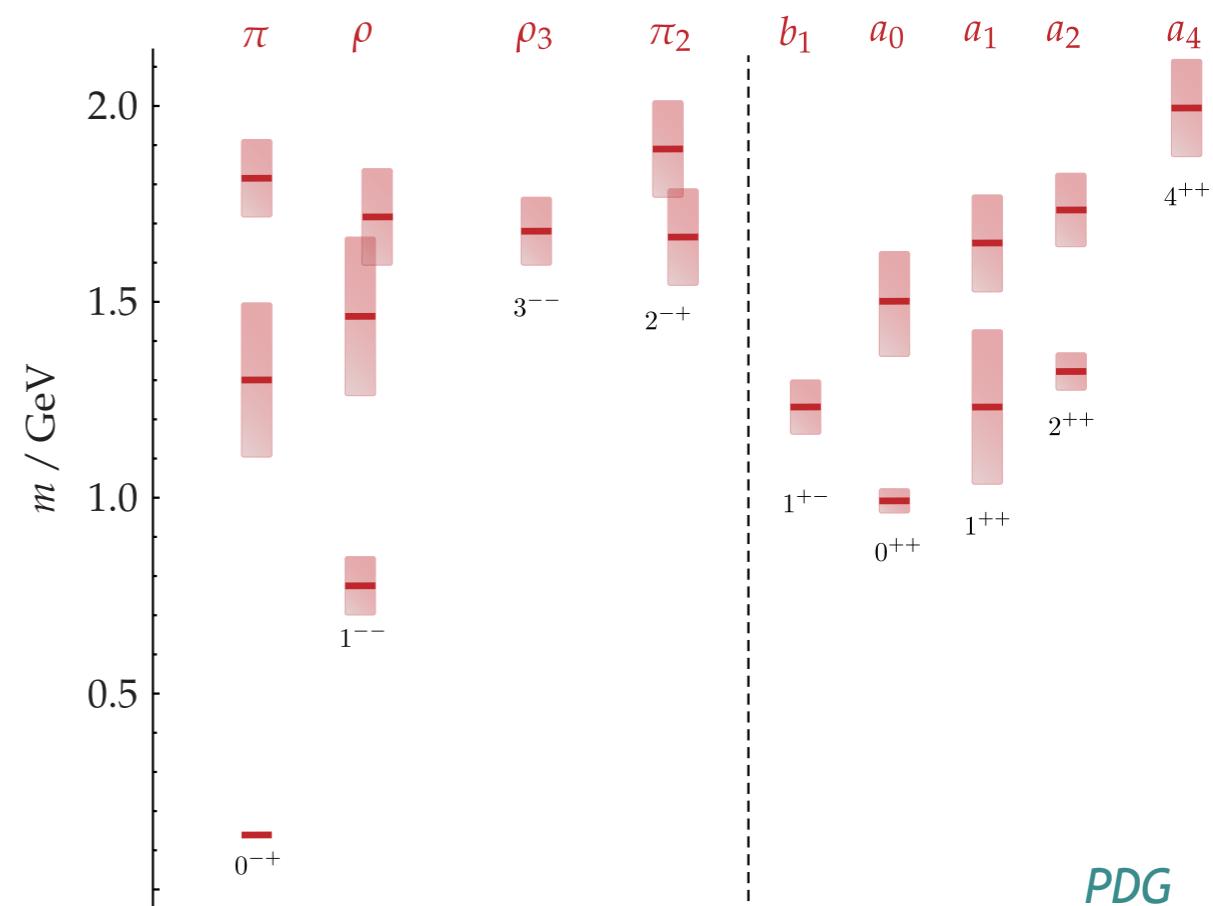
understanding the hadron spectrum

- how are mesons and baryons assembled from quarks and gluons ?
 - e.g. why only isospin ≤ 1 ?
 - do excited gluonic fields play a role?

QUANTUM CHROMODYNAMICS

$$\begin{aligned} \mathcal{L}_{\text{QCD}} = & \bar{\psi}(i\gamma^\mu \partial_\mu - m)\psi \\ & - g \bar{\psi}\gamma^\mu\psi A_\mu \\ & - \frac{1}{2} \text{tr } F_{\mu\nu}F^{\mu\nu} \end{aligned}$$

ISOSPIN=1 MESON SPECTRUM



hadrons from lattice QCD

- discretize fields on a finite Euclidean hypercubic grid
- integrate out the quark fields
- Monte Carlo sample 100s of configurations of gluon fields
- spectrum from two-point correlation functions

$$C(t) = \langle 0 | \mathcal{O}(t) \mathcal{O}^\dagger(0) | 0 \rangle$$

$$C(t) = \langle 0 | \mathcal{O}(0) e^{-Ht} \mathcal{O}^\dagger(0) | 0 \rangle$$

$$C(t) = \langle 0 | \mathcal{O}(0) e^{-Ht} \sum_{\mathfrak{n}} |\mathfrak{n}\rangle \langle \mathfrak{n}| \mathcal{O}^\dagger(0) | 0 \rangle$$

$$C(t) = \sum_{\mathfrak{n}} e^{-E_{\mathfrak{n}} t} \langle 0 | \mathcal{O}(0) | \mathfrak{n} \rangle \langle \mathfrak{n} | \mathcal{O}^\dagger(0) | 0 \rangle$$

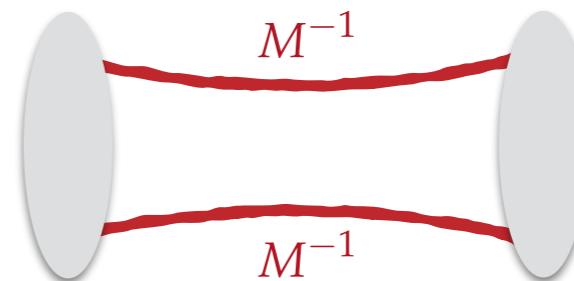


hadrons from lattice QCD

- spectrum from two-point correlation functions

$$C(t) = \langle 0 | \mathcal{O}(t) \mathcal{O}^\dagger(0) | 0 \rangle$$

$$\mathcal{O} \sim \bar{\psi} \Gamma \psi$$



$$\int d^4x \bar{\psi}(x) i\gamma^\mu (\partial_\mu + igA_\mu(x)) \psi(x) \xrightarrow{\text{discretize}} \sum_{xy} \bar{\psi}_x M_{xy} \psi_y$$

integrate out
quark fields

$$\psi_x \bar{\psi}_y \rightsquigarrow M_{xy}^{-1}$$

$$\langle \bar{\psi}_t \Gamma_t \psi_t \bar{\psi}_0 \Gamma_0 \psi_0 \rangle \rightsquigarrow M_{0t}^{-1} \Gamma_t M_{t0}^{-1} \Gamma_0$$

“Wick contract”

matrix multiply and trace

excited states from correlators

- how to get at excited QCD eigenstates ?

– optimal operator for state $|\mathfrak{n}\rangle$: $\Omega_{\mathfrak{n}}^\dagger \sim \sum_i v_i^{(\mathfrak{n})} \mathcal{O}_i^\dagger$

for a basis of
meson operators $\{\mathcal{O}_i\}$

- can be obtained (in a variational sense) from the matrix of correlators

$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j^\dagger(0) | 0 \rangle$$

- by solving a generalized eigenvalue problem

$$C(t)v^{(\mathfrak{n})} = C(t_0)v^{(\mathfrak{n})} \lambda_{\mathfrak{n}}(t)$$

eigenvalues

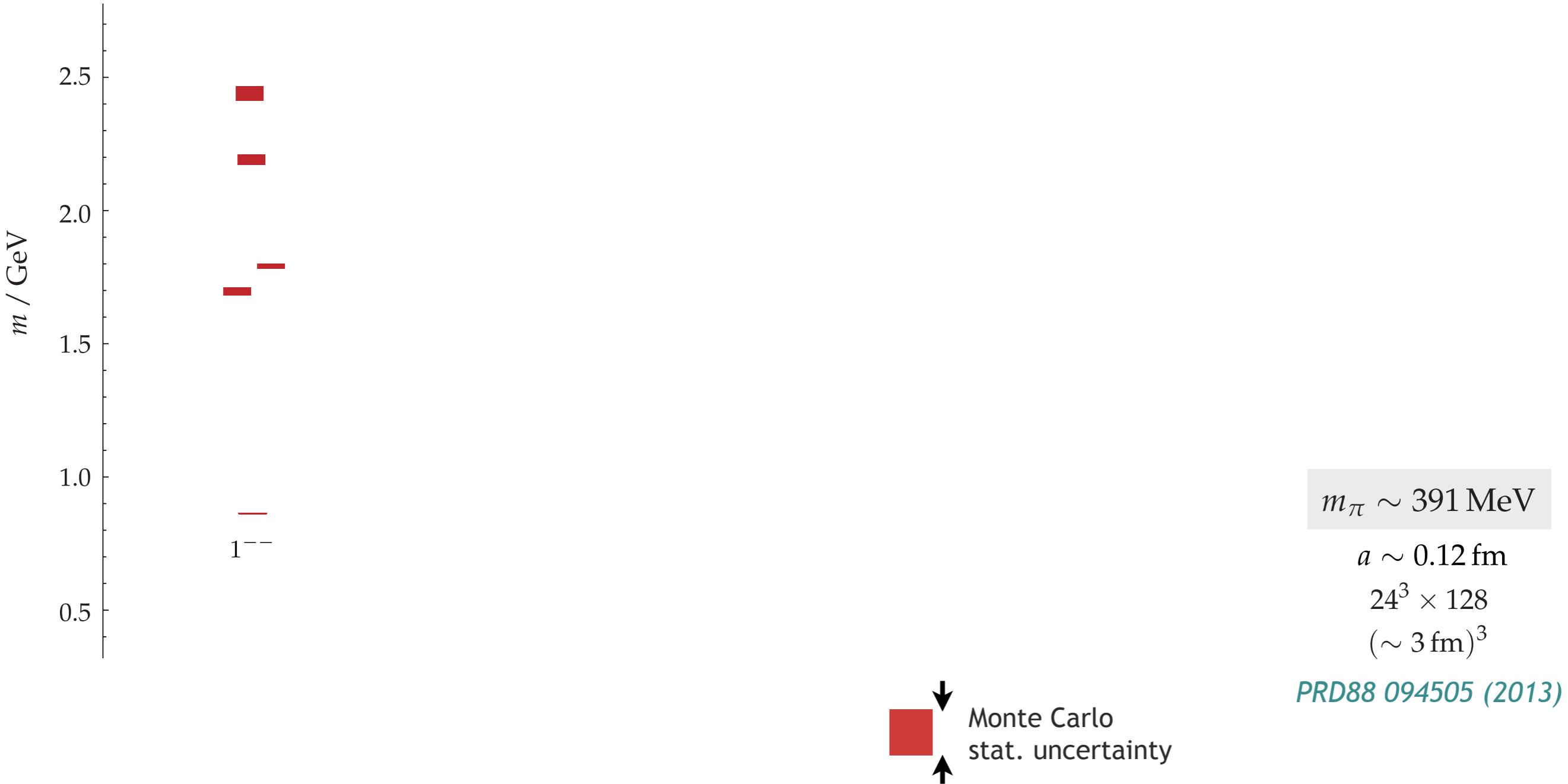
$$\lambda_{\mathfrak{n}}(t) \sim e^{-E_{\mathfrak{n}}(t-t_0)}$$

- a large basis can be constructed using covariant derivatives :

$$\mathcal{O} \sim \bar{\psi} \Gamma \overleftrightarrow{D} \dots \overleftrightarrow{D} \psi$$

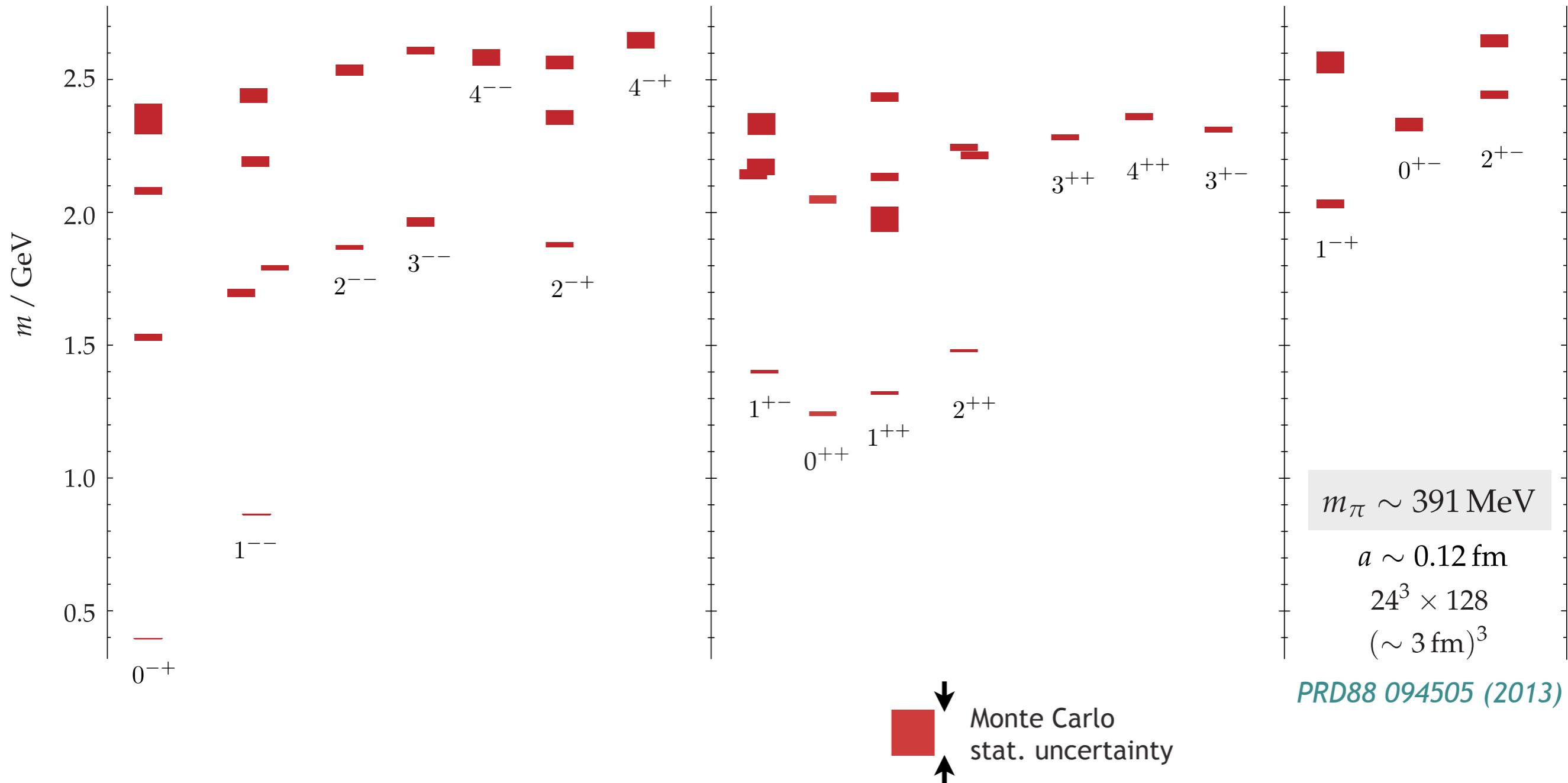
excited state spectra

- for example vector meson sector, using more than 20 operators



excited state spectra

- meson spectrum for a range of J^{PC}



excited hadrons are resonances

PHYSICAL REVIEW D

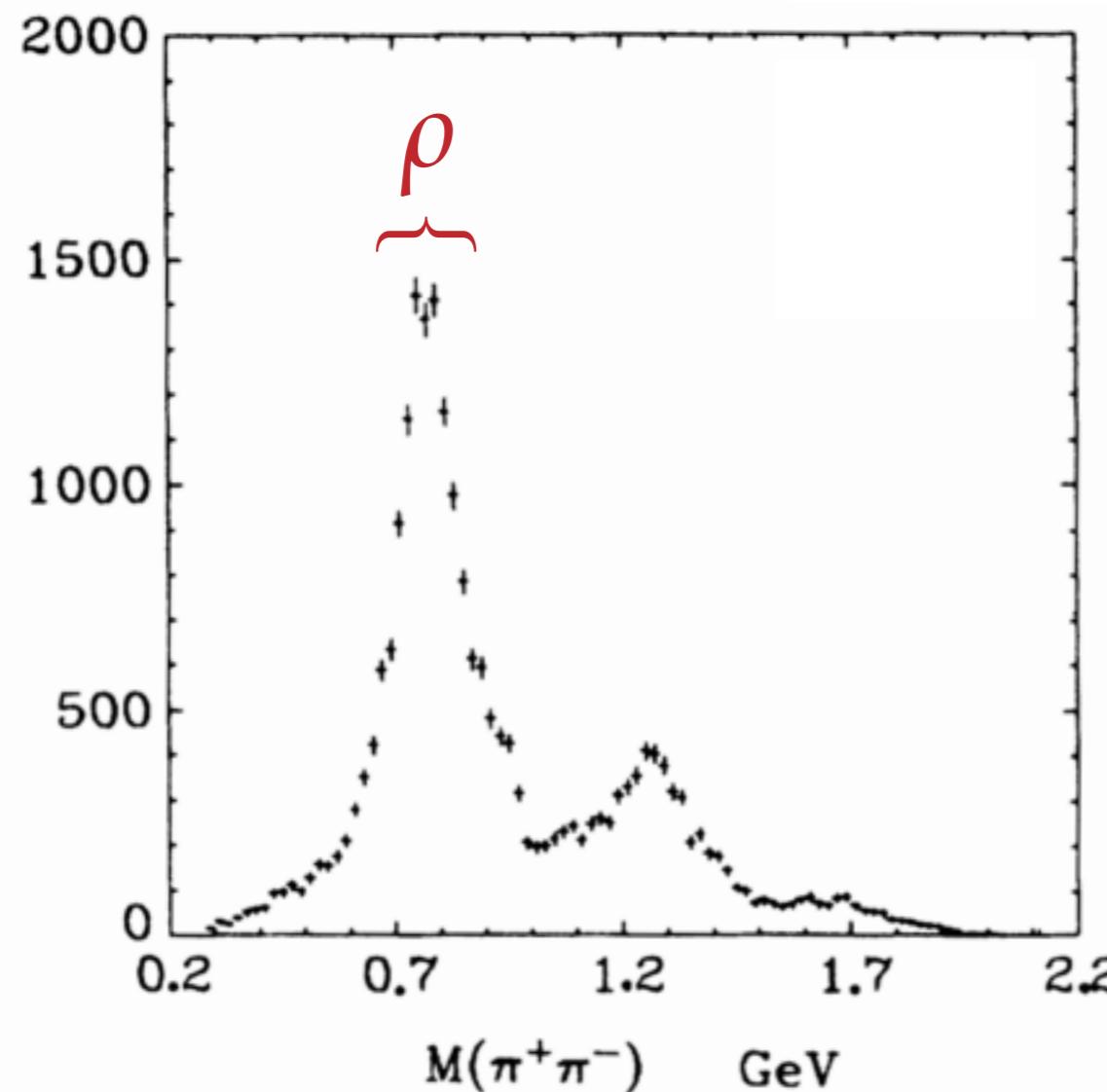
VOLUME 7, NUMBER 5

1 MARCH 1973

$\pi\pi$ Partial-Wave Analysis from Reactions $\pi^+p \rightarrow \pi^+\pi^-\Delta^{++}$ and $\pi^+p \rightarrow K^+K^-\Delta^{++}$ at 7.1 GeV/c†

S. D. Protopopescu,* M. Alston-Garnjost, A. Barbaro-Galtieri, S. M. Flatté,‡
 J. H. Friedman,§ T. A. Lasinski, G. R. Lynch, M. S. Rabin,|| and F. T. Solmitz
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(Received 25 September 1972)



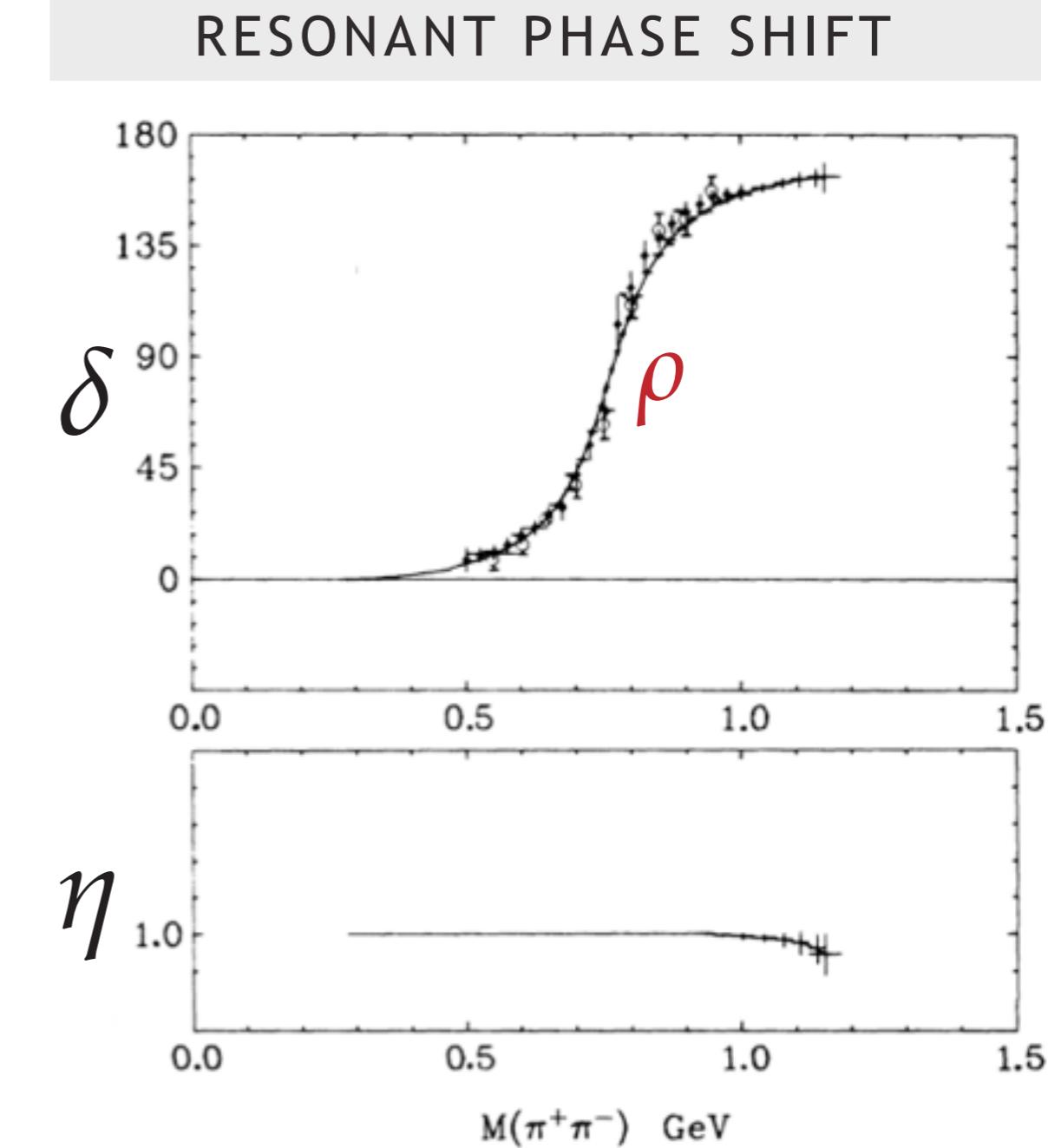
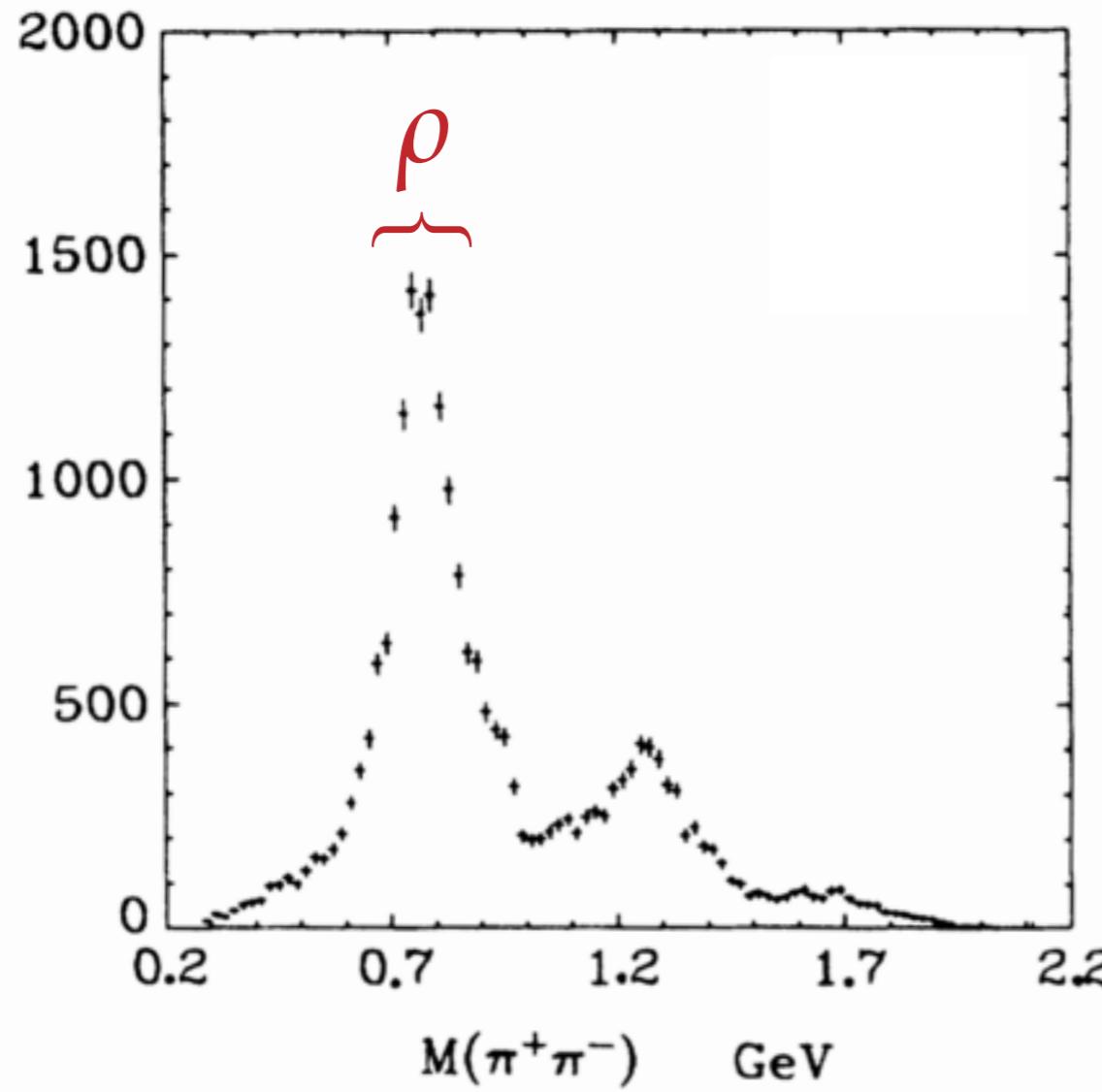
expand angular dependence
in *partial waves*

PARTIAL WAVE AMPLITUDE

$$f_\ell = \frac{1}{2i} (\eta_\ell e^{2i\delta_\ell} - 1)$$

$\eta = 1$ elastic
 $\eta \leq 1$ inelastic

excited hadrons are resonances



scattering in a finite cubic volume

- expect a discrete spectrum in a finite periodic volume

$$\psi(x + L) = \psi(x)$$

e.g. free particle $e^{ip(x+L)} = e^{ipx}$

quantized momentum $p = \frac{2\pi}{L}n$

- for an interacting theory

$$\cot \delta_\ell(E) = \mathcal{M}_\ell(E, L)$$

LÜSCHER ...

elastic scattering phase-shift

known function

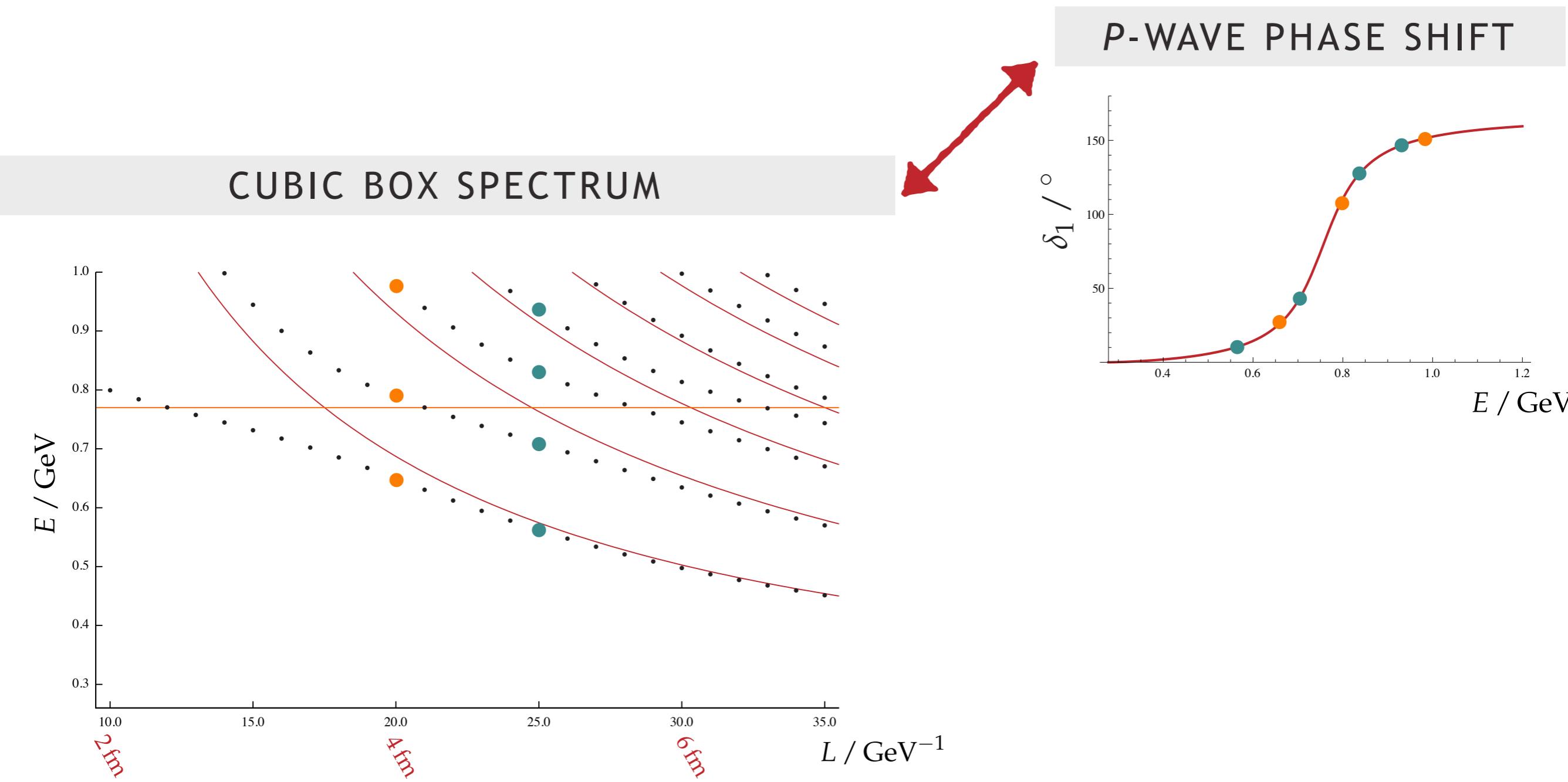
discrete energies in a finite-volume



discrete values of the phase-shift

scattering in a finite cubic volume

- e.g. experimental $\pi\pi$ $J=1$ P -wave scattering amplitude



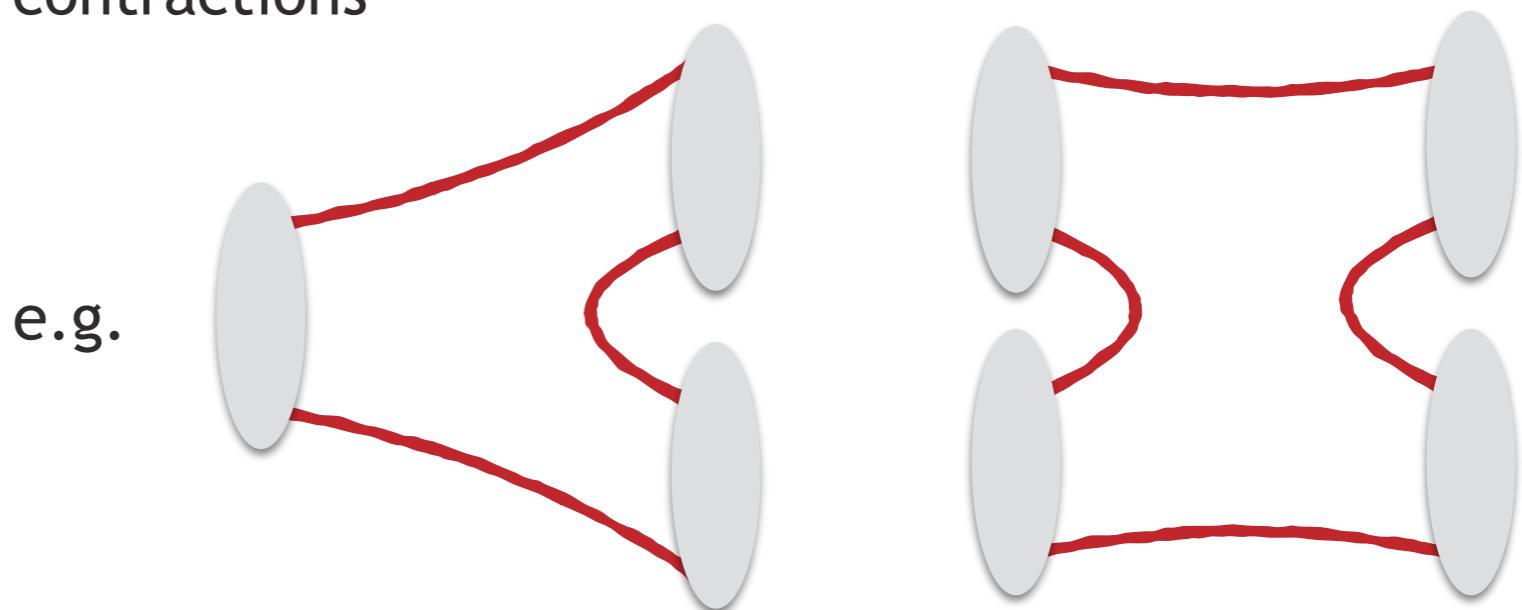
determining the finite-volume spectrum

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- include operators which resemble a pair of pions

$$\sum_{\hat{k}_1, \hat{k}_2} C(\Lambda, \vec{P}; \vec{k}_1, \vec{k}_2) \pi^\dagger(\vec{k}_1) \pi^\dagger(\vec{k}_2)$$
$$\pi^\dagger \sim \bar{\psi} \Gamma \psi$$

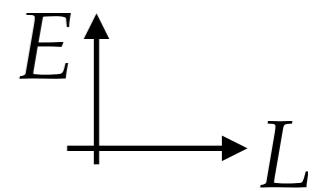
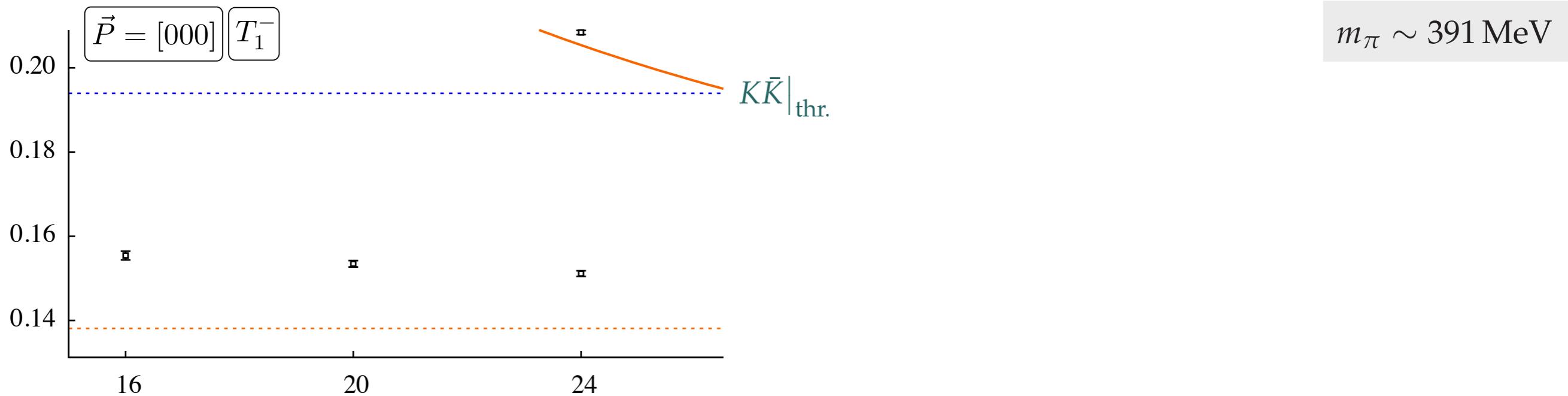
- form correlator matrix with both $\bar{\psi} \Gamma \psi$ and $\pi\pi$ -like
 - more complicated Wick contractions



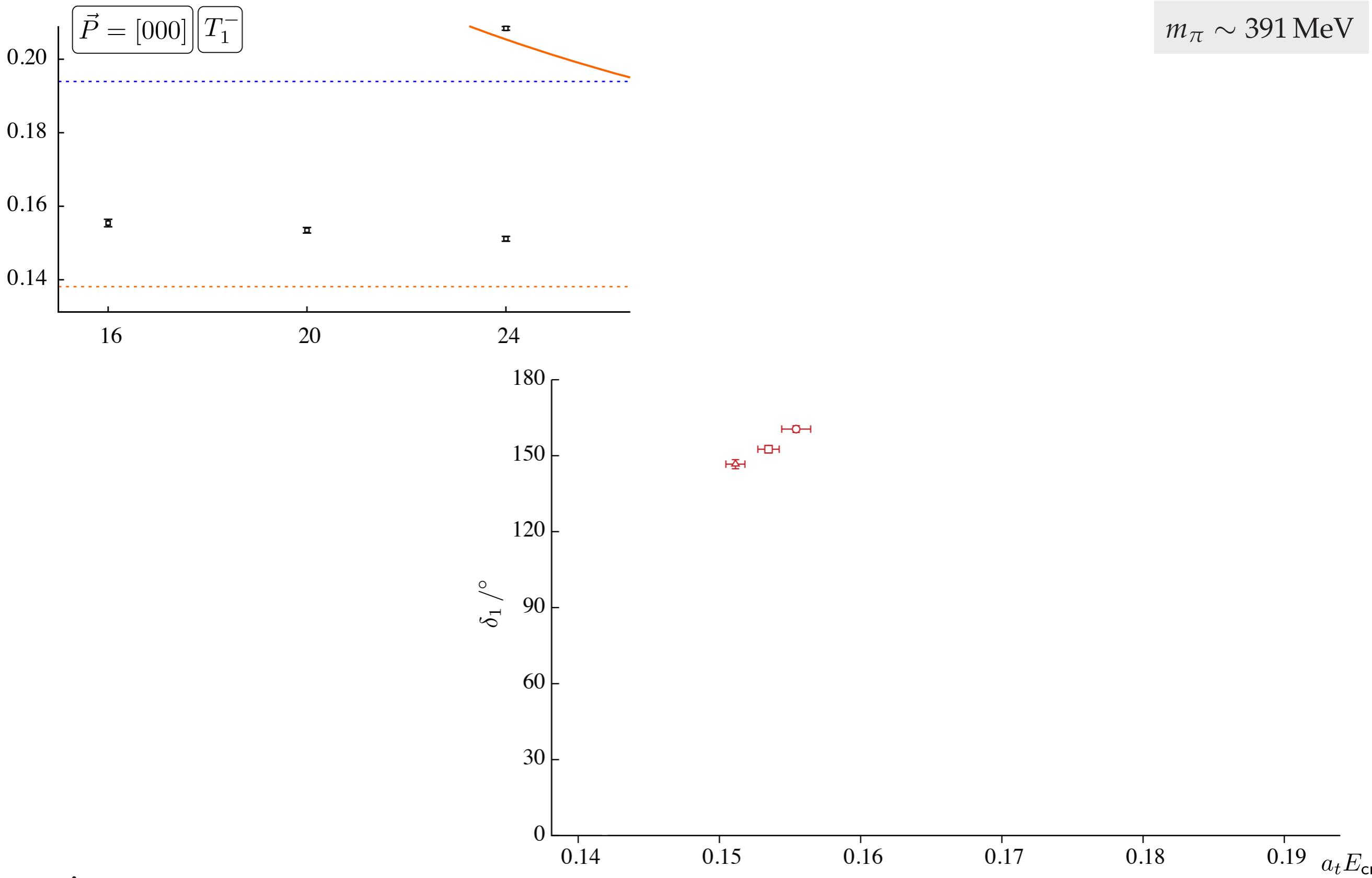
- quark ‘annihilation’ lines have in the past been a computational challenge

$$M_{tt}^{-1}$$

finite-volume spectrum

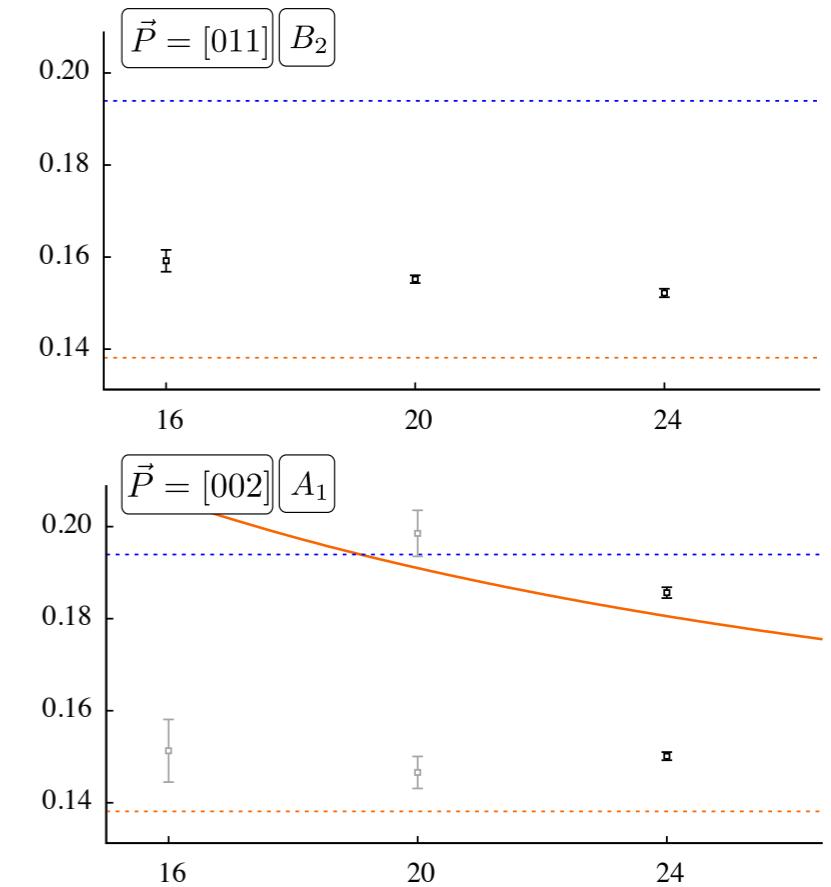
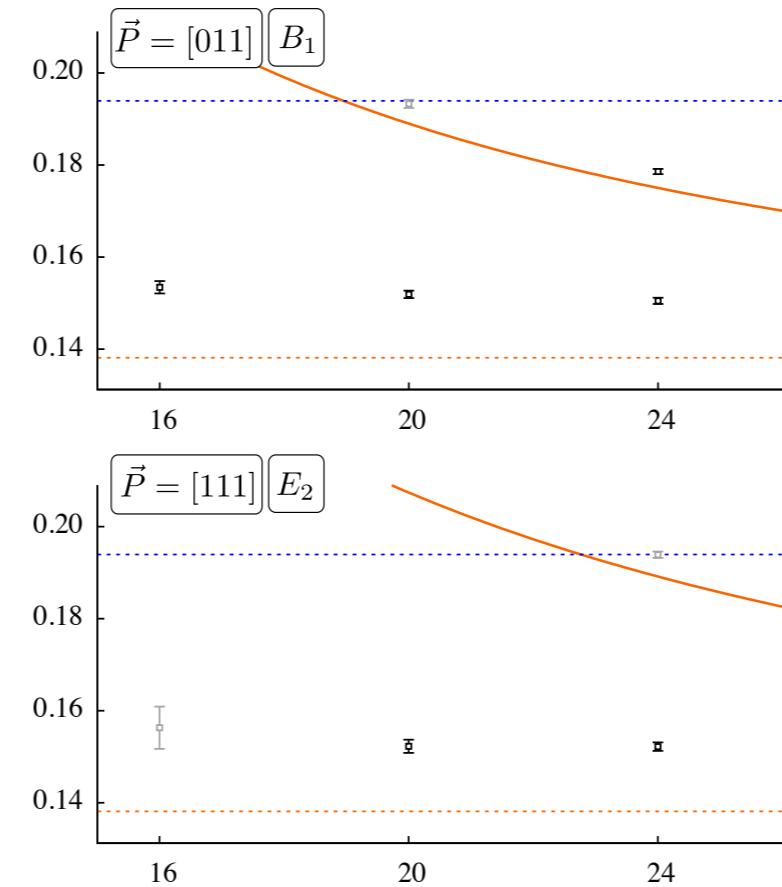
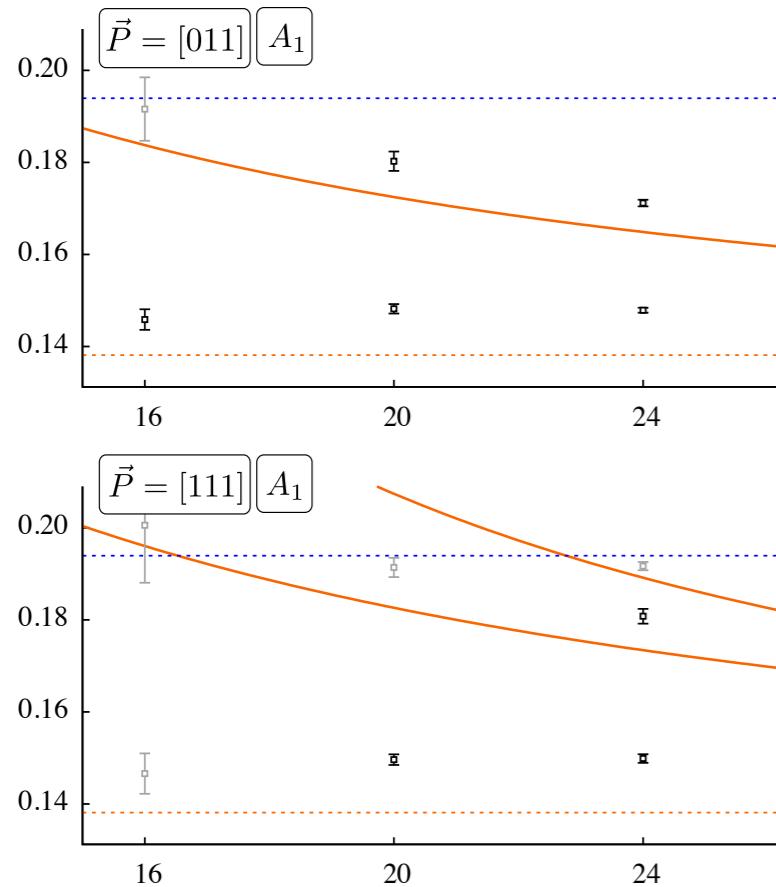
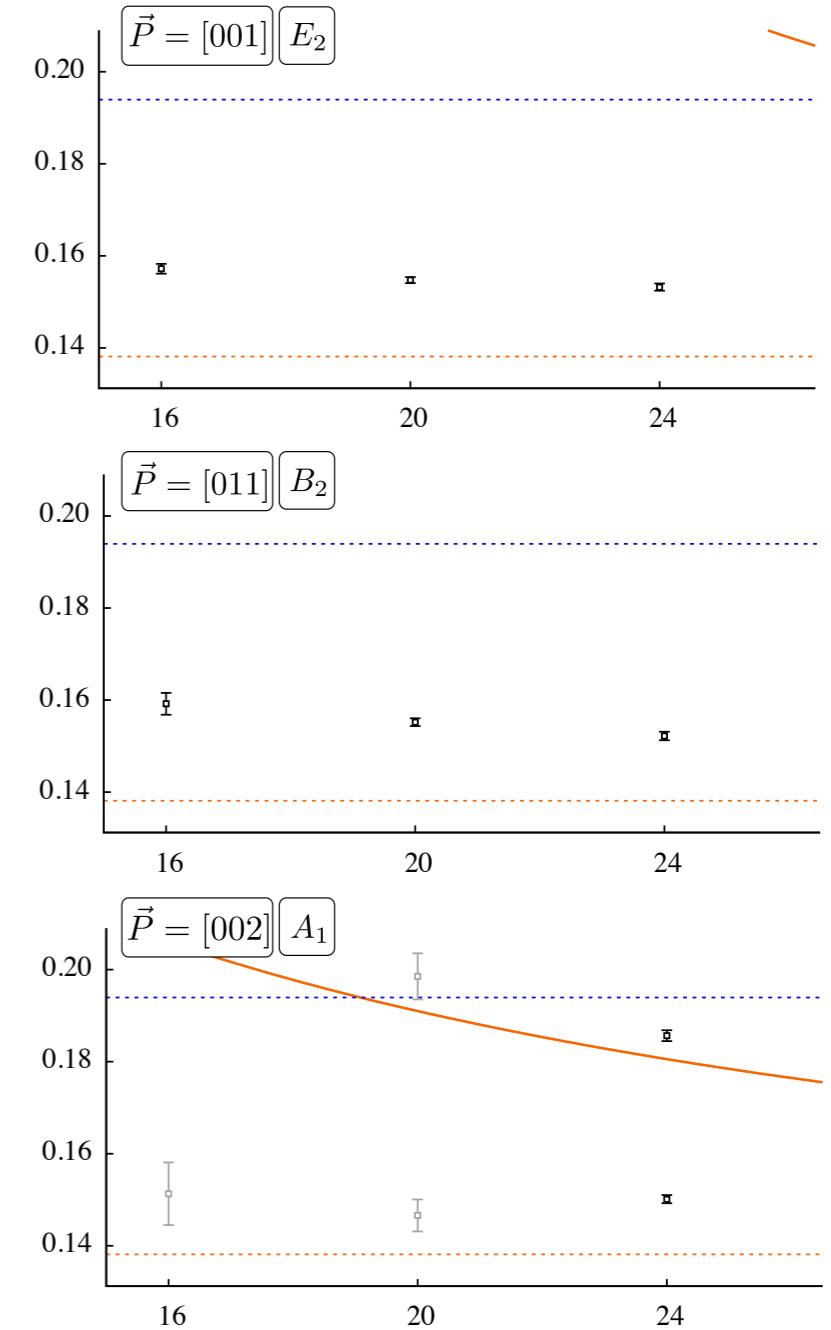
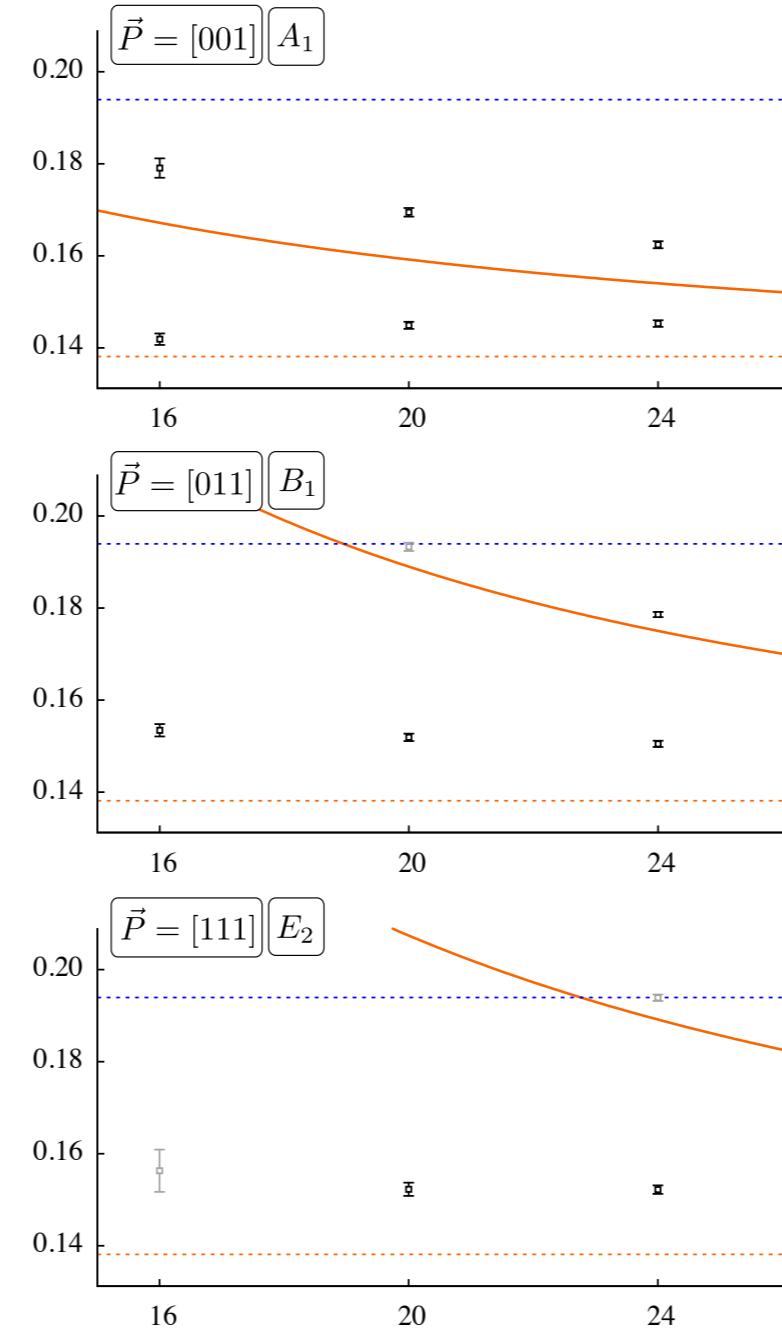
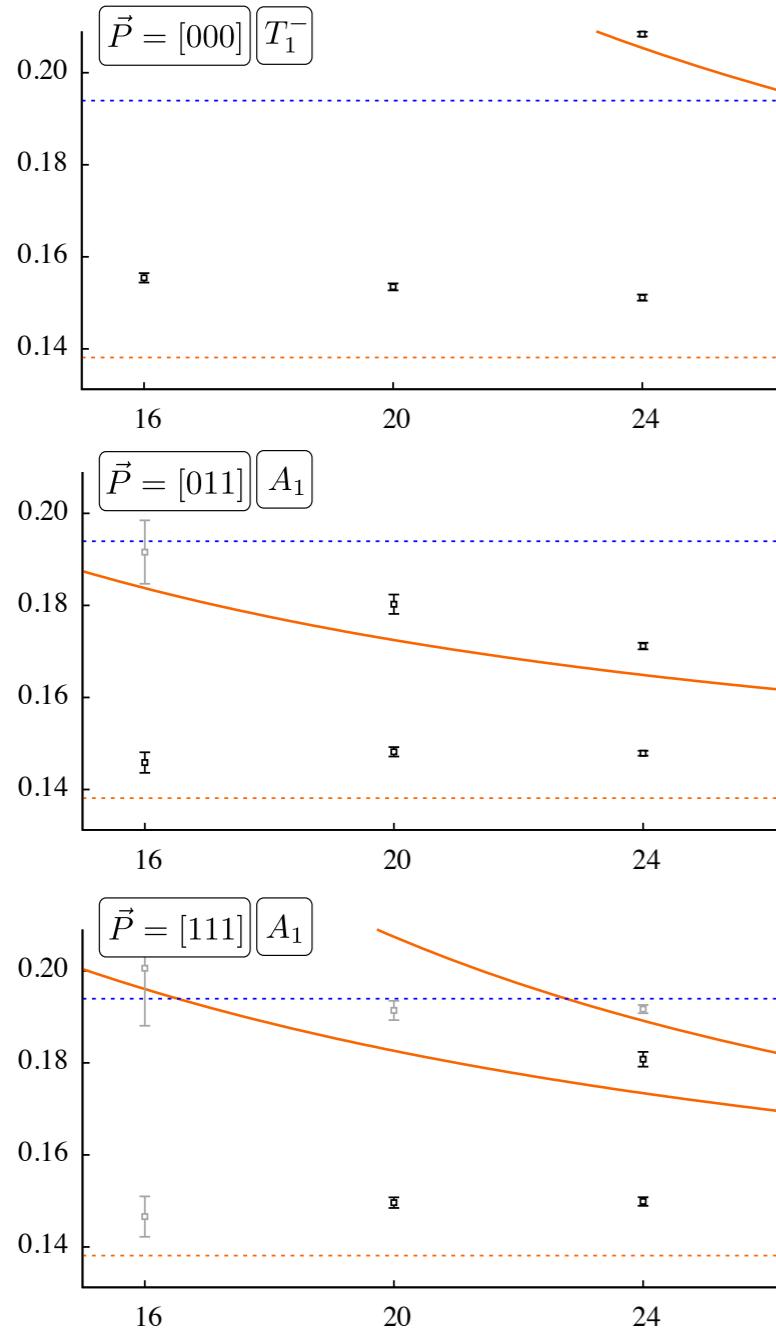


finite-volume spectrum

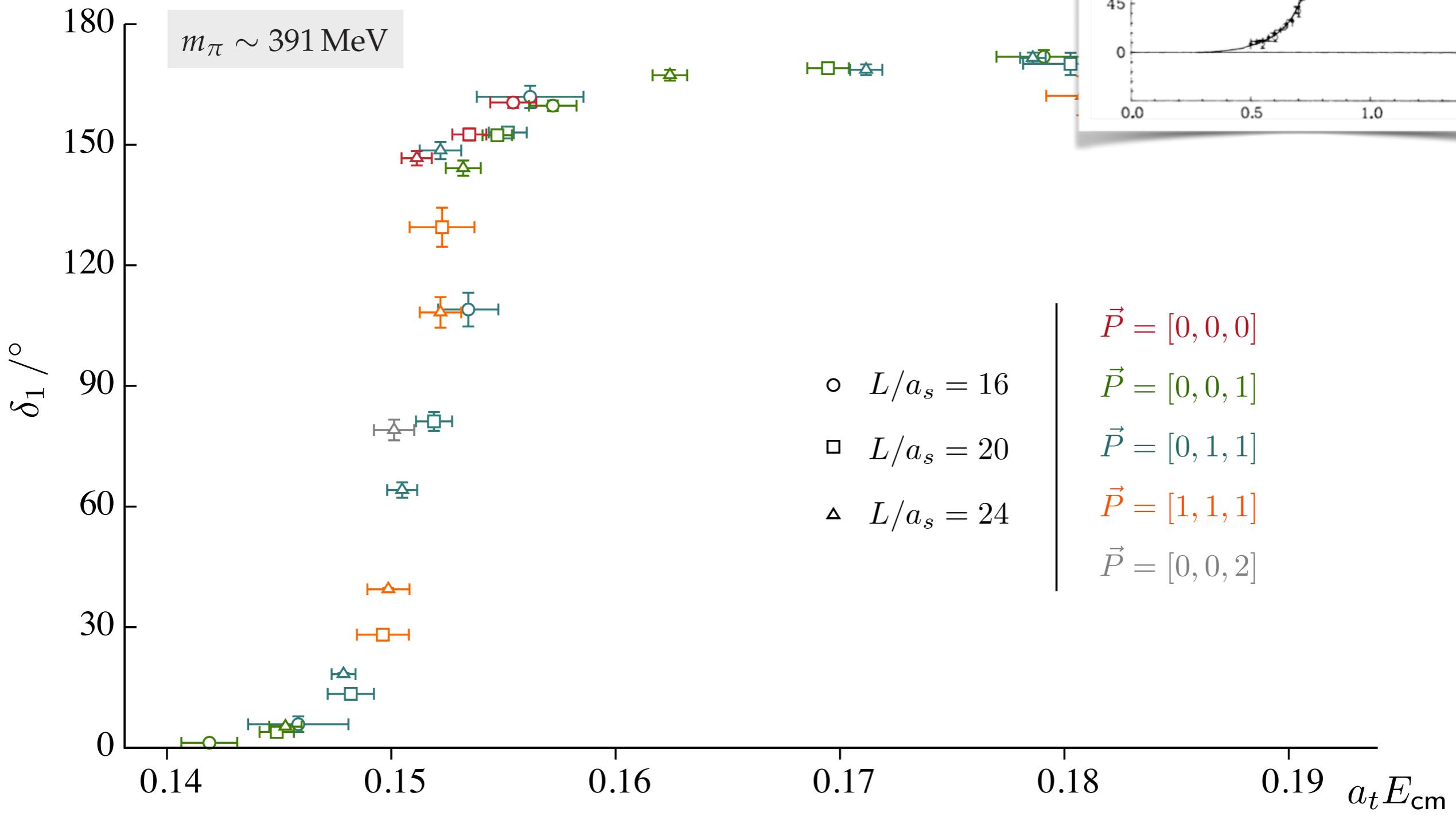


finite-volume spectrum - moving frames

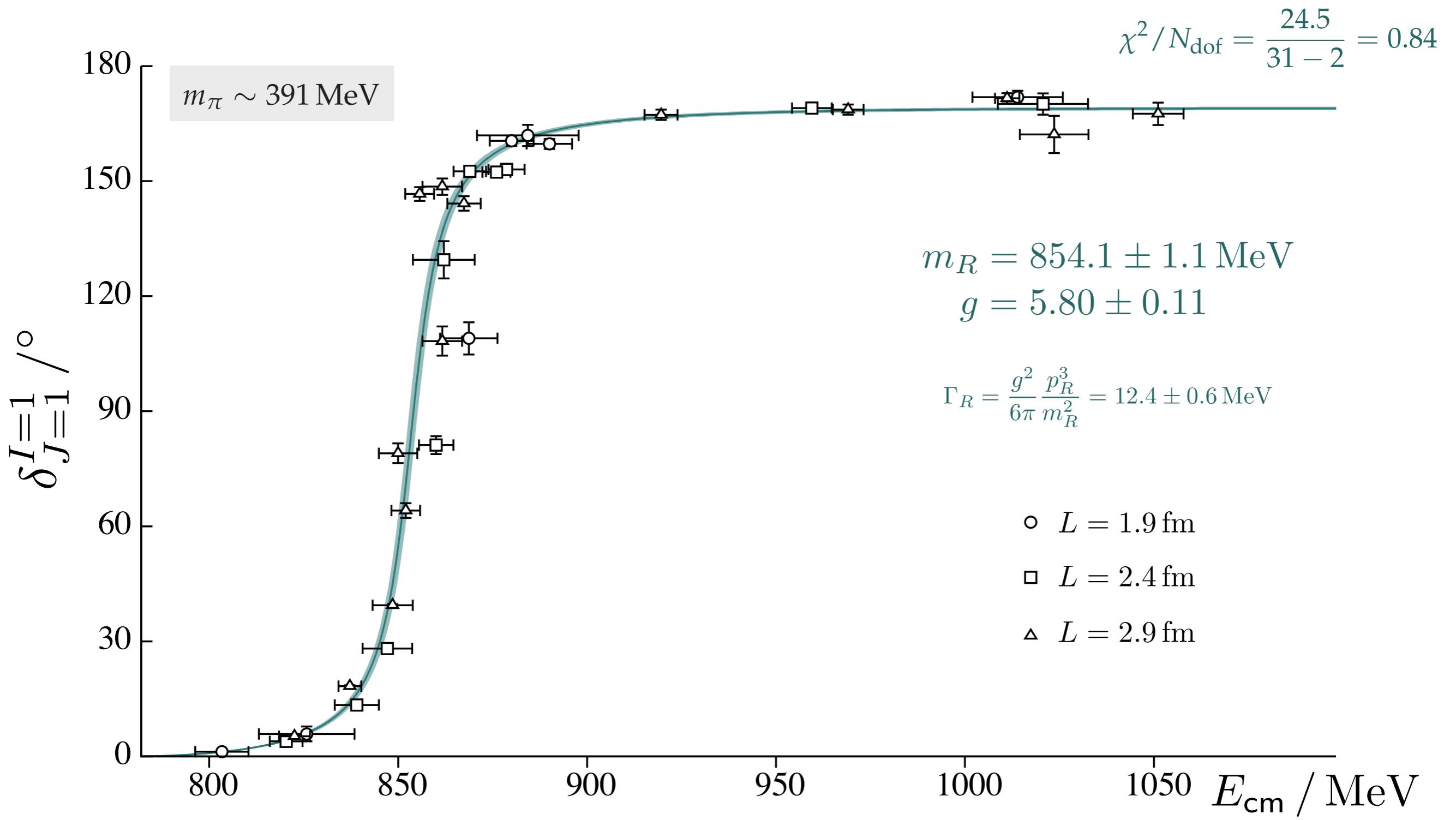
$m_\pi \sim 391 \text{ MeV}$



$\pi\pi$ P -wave phase-shift



ρ resonance



coupled-channel scattering

- finite-volume formalism recently derived (multiple methods)

*HE, JHEP 0507 011
HANSEN, PRD86 016007
BRICENO, PRD88 094507
GUO, PRD88 014051*

$$\det \left[([t^{(\ell)}(E)]_{ij}^{-1} + i\rho_i(E) \delta_{ij}) - \delta_{ij} \mathcal{M}_\ell(p_i(E)L) \right] = 0$$

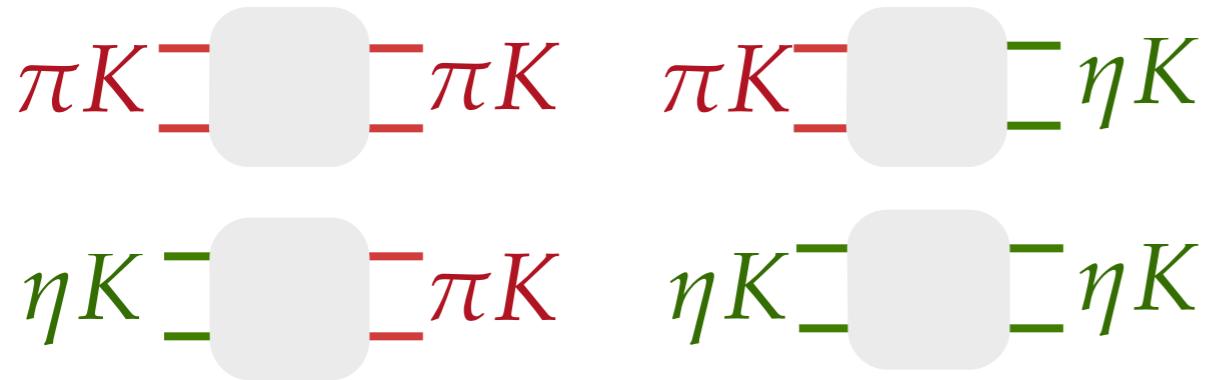
scattering matrix phase space known functions

- but this is one equation for multiple unknowns (per energy level)
 - parameterize the energy dependence of t
 - try to describe a spectrum globally

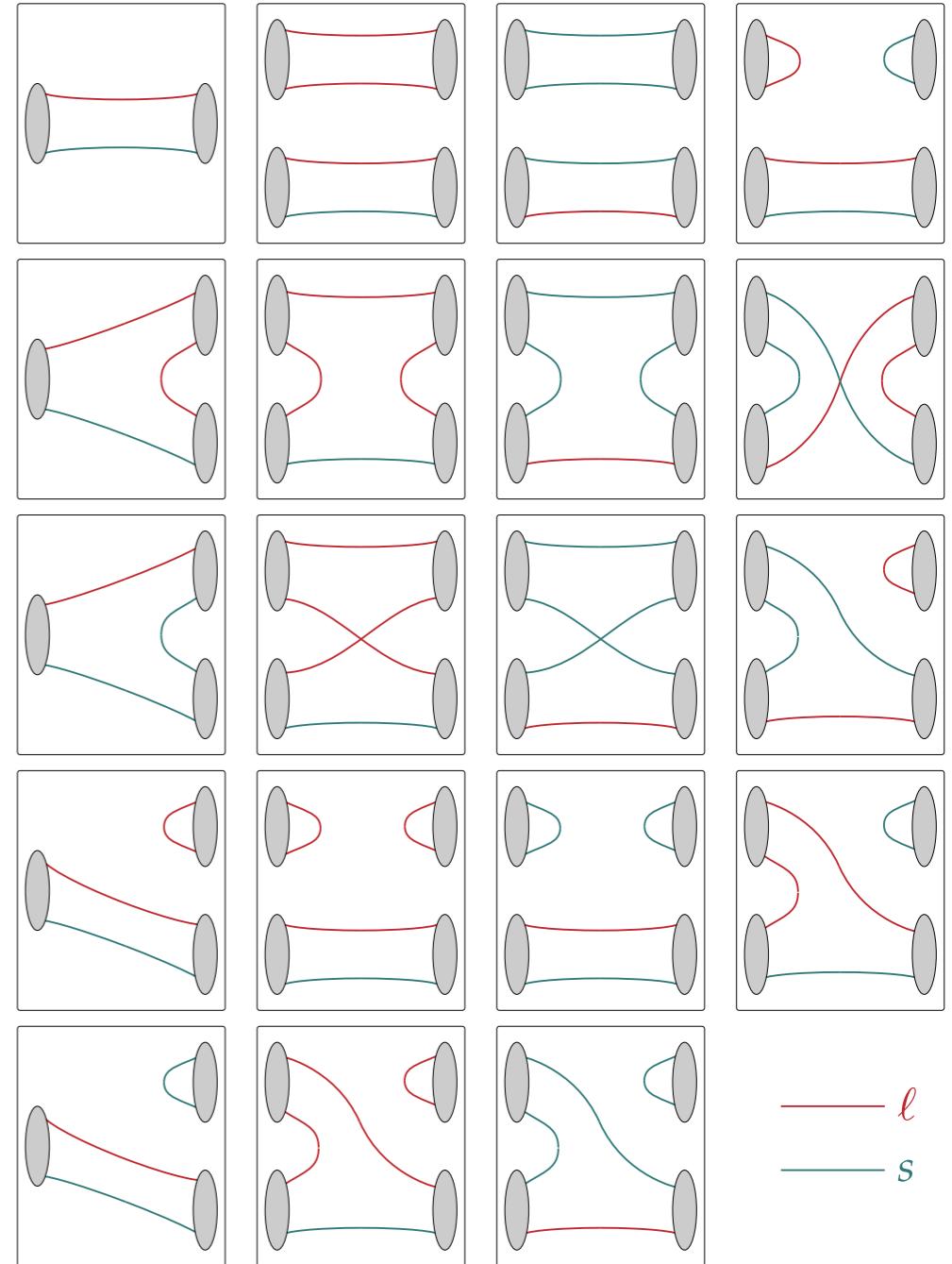
$\pi K/\eta K$ scattering & kaon resonances

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- an example of coupled-channel scattering



WICK CONTRACTIONS



- compute finite-volume spectrum

$$\bar{u} \Gamma s$$

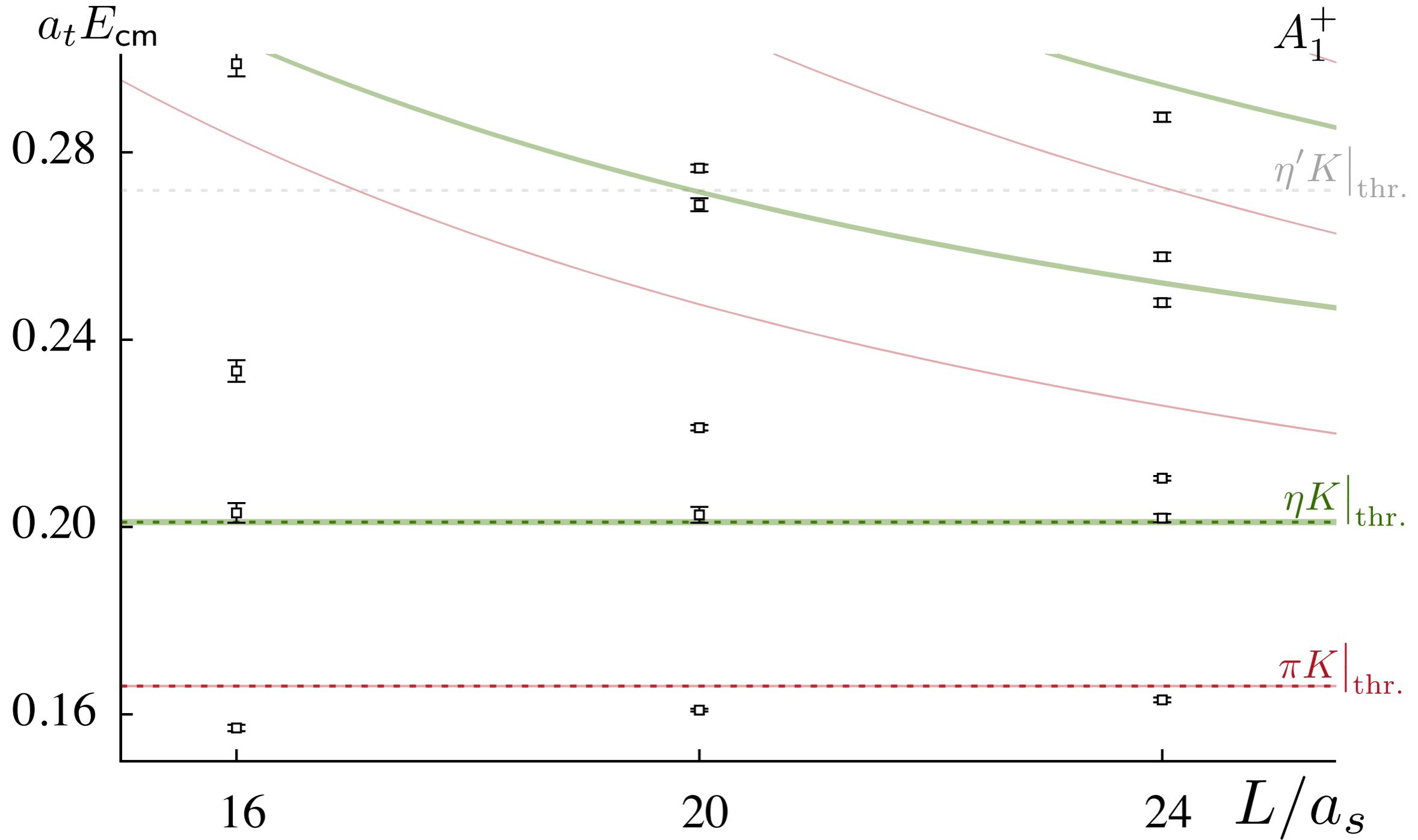
$$\sum_{\hat{k}_1, \hat{k}_2} C(\Lambda, \vec{P}; \vec{k}_1, \vec{k}_2) \pi^\dagger(\vec{k}_1) K^\dagger(\vec{k}_2)$$

$$\sum_{\hat{k}_1, \hat{k}_2} C(\Lambda, \vec{P}; \vec{k}_1, \vec{k}_2) \eta^\dagger(\vec{k}_1) K^\dagger(\vec{k}_2)$$

$\pi K/\eta K$ scattering & kaon resonances

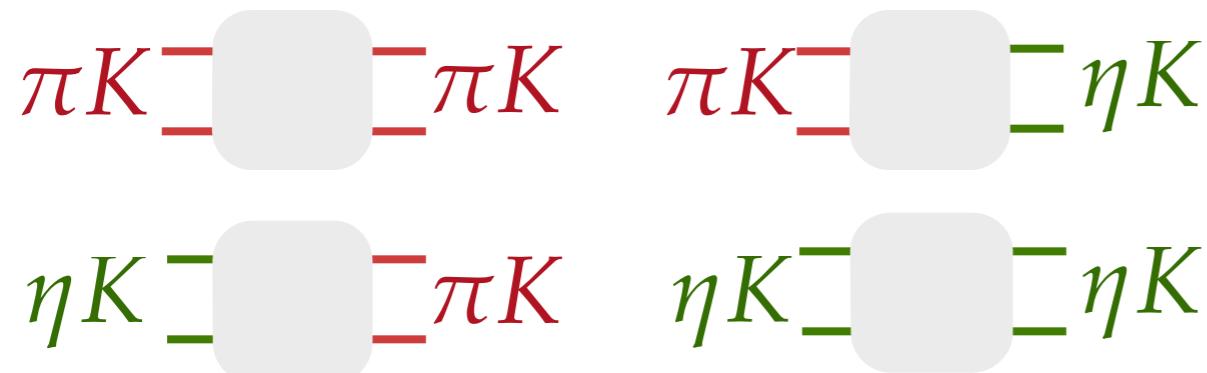
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- rest frame “S-wave” spectrum

 $m_\pi \sim 391 \text{ MeV}$ 

$\pi K/\eta K$ scattering & kaon resonances

- parameterize the t -matrix in a unitarity conserving way



$$t_{ij}^{-1}(E) = K_{ij}^{-1}(E) + \delta_{ij} I_i(E)$$

$$K_{ij}(E) = \frac{g_i g_j}{m^2 - E^2} + \gamma_{ij}$$

- vary the parameters, solving

$$\det \left[([t^{(\ell)}(E)]_{ij}^{-1} + i\rho_i(E) \delta_{ij}) - \delta_{ij} \mathcal{M}_\ell(E, L) \right] = 0$$

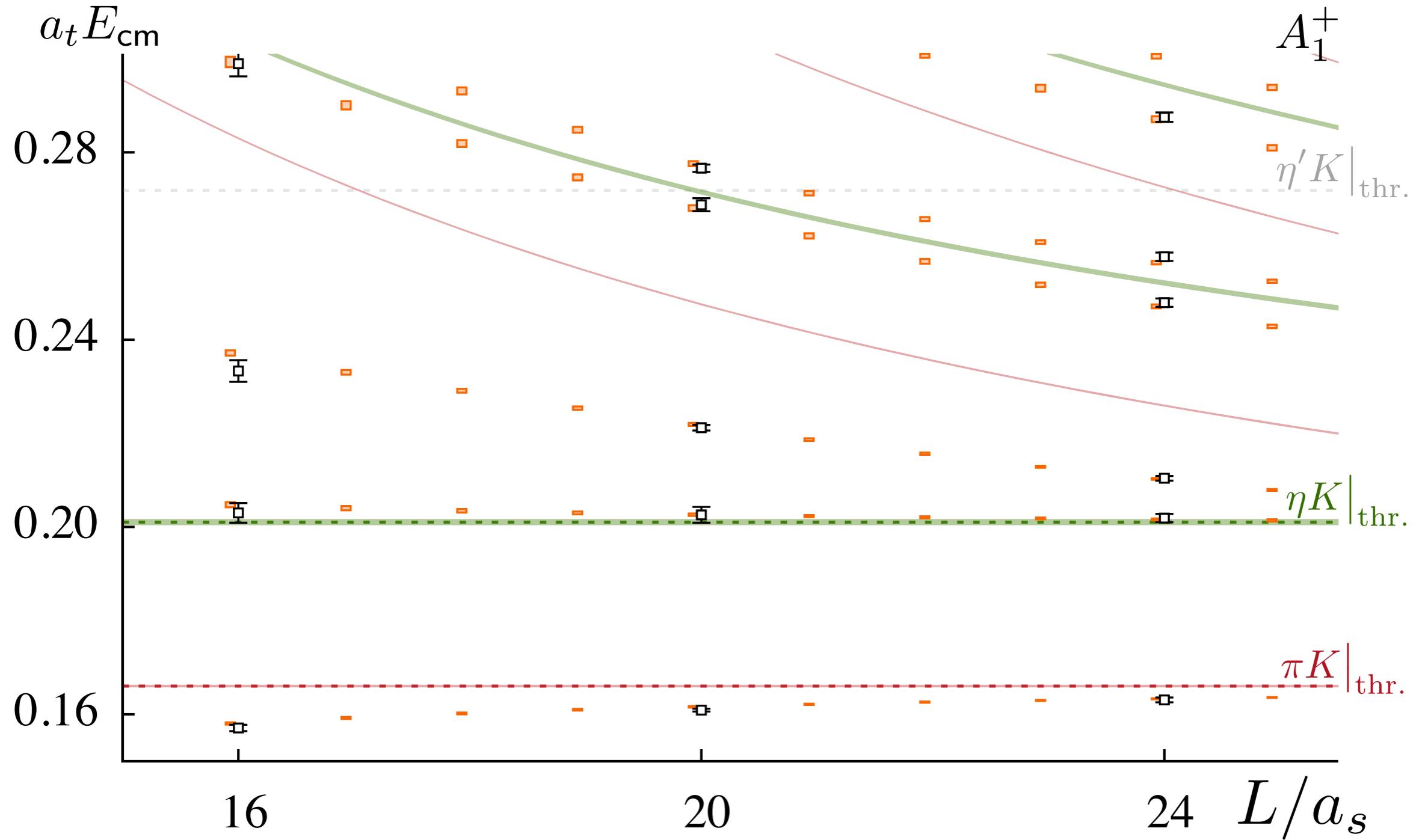
for the spectrum each time

$\pi K/\eta K$ scattering

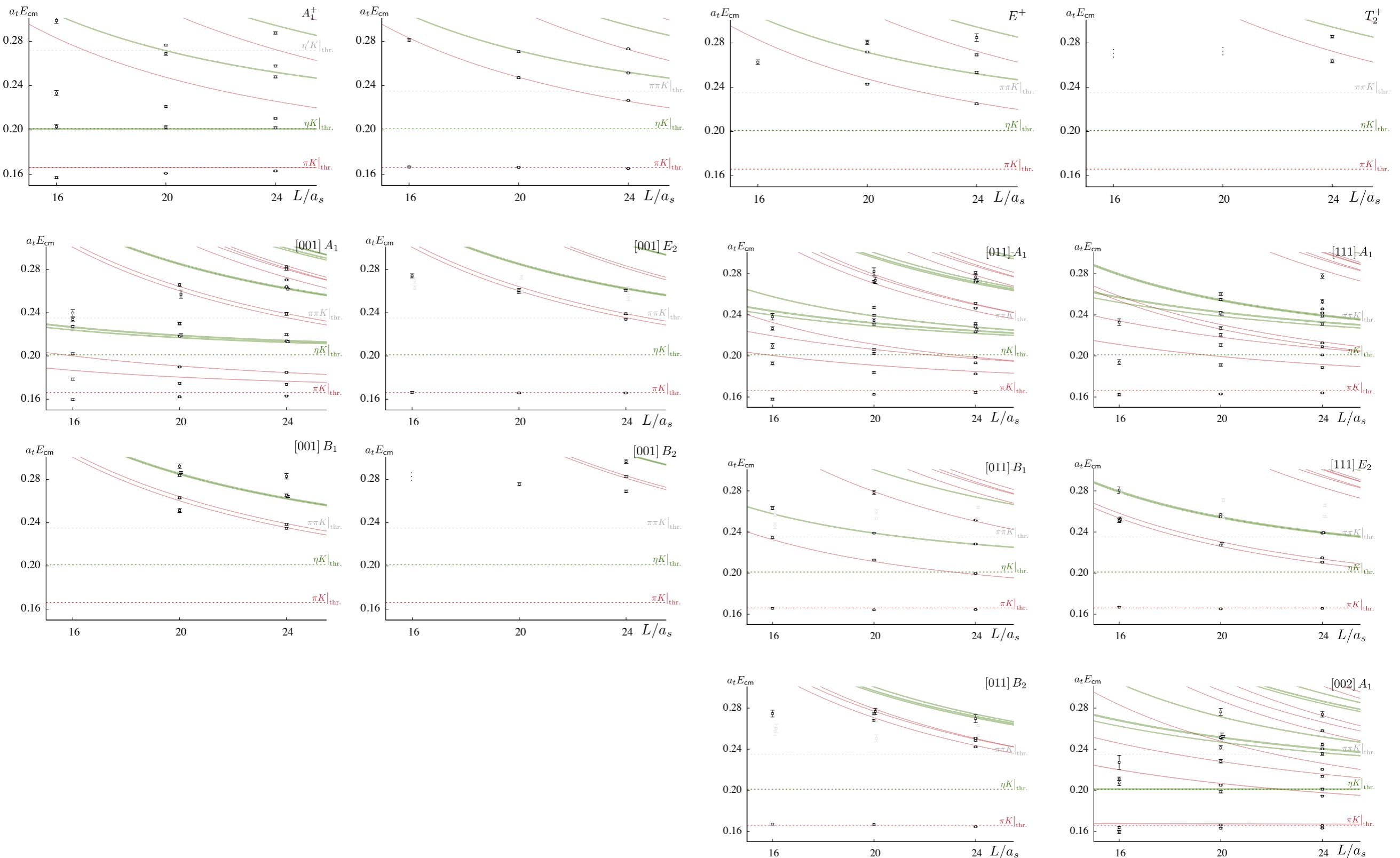
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model
description $\chi^2/N_{\text{dof}} = \frac{6.40}{15 - 6} = 0.71$

$m_\pi \sim 391 \text{ MeV}$



spectra in moving frames



$\pi K/\eta K$ scattering

$m_\pi \sim 391$ MeV

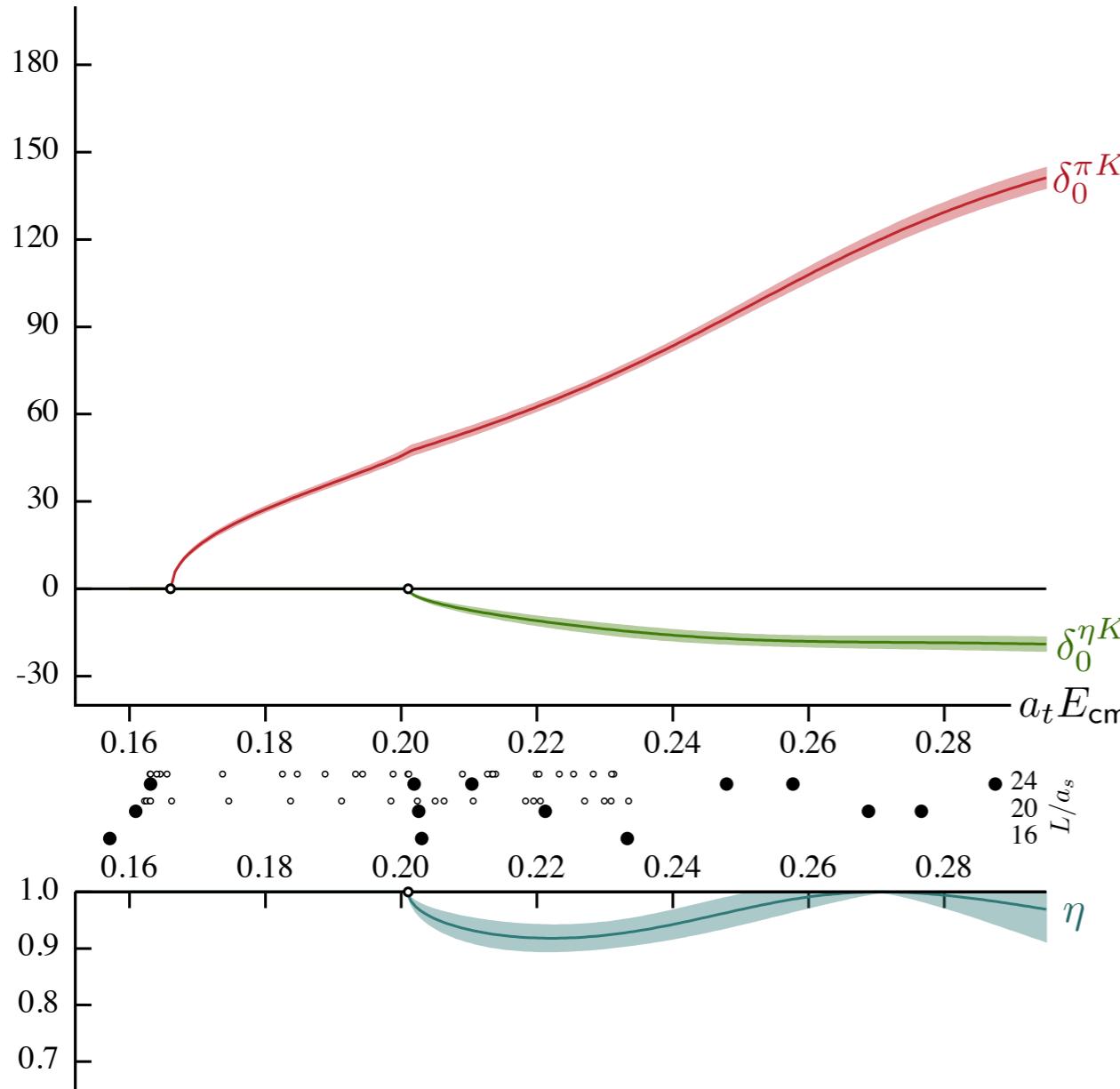
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- describe all the finite-volume spectra

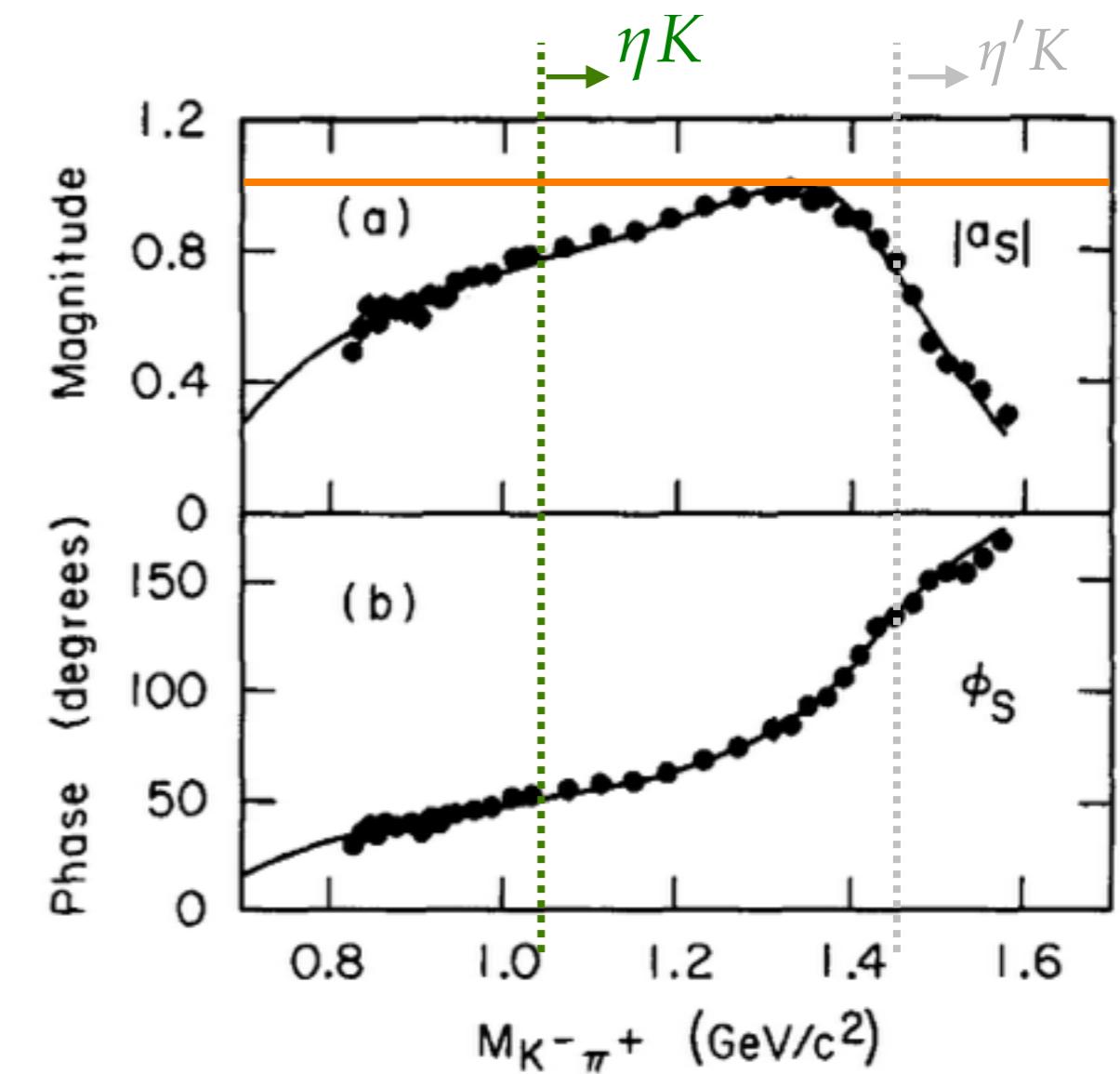
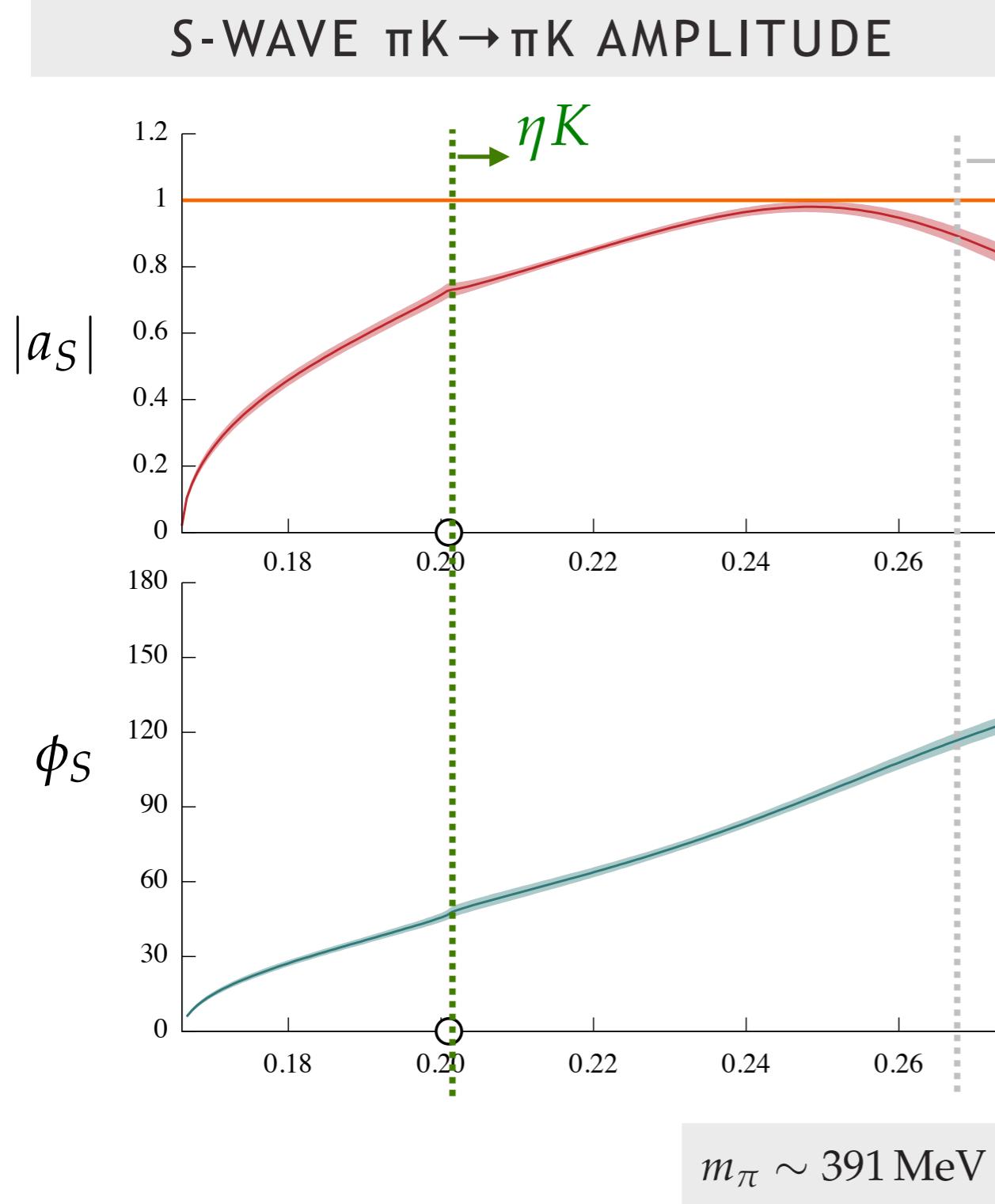
$$\chi^2/N_{\text{dof}} = \frac{49.1}{61 - 6} = 0.89$$

$$S_{\pi K, \pi K} = \eta e^{2i\delta^{\pi K}}$$
$$S_{\eta K, \eta K} = \eta e^{2i\delta^{\eta K}}$$

S-WAVE $\pi K/\eta K$ SCATTERING

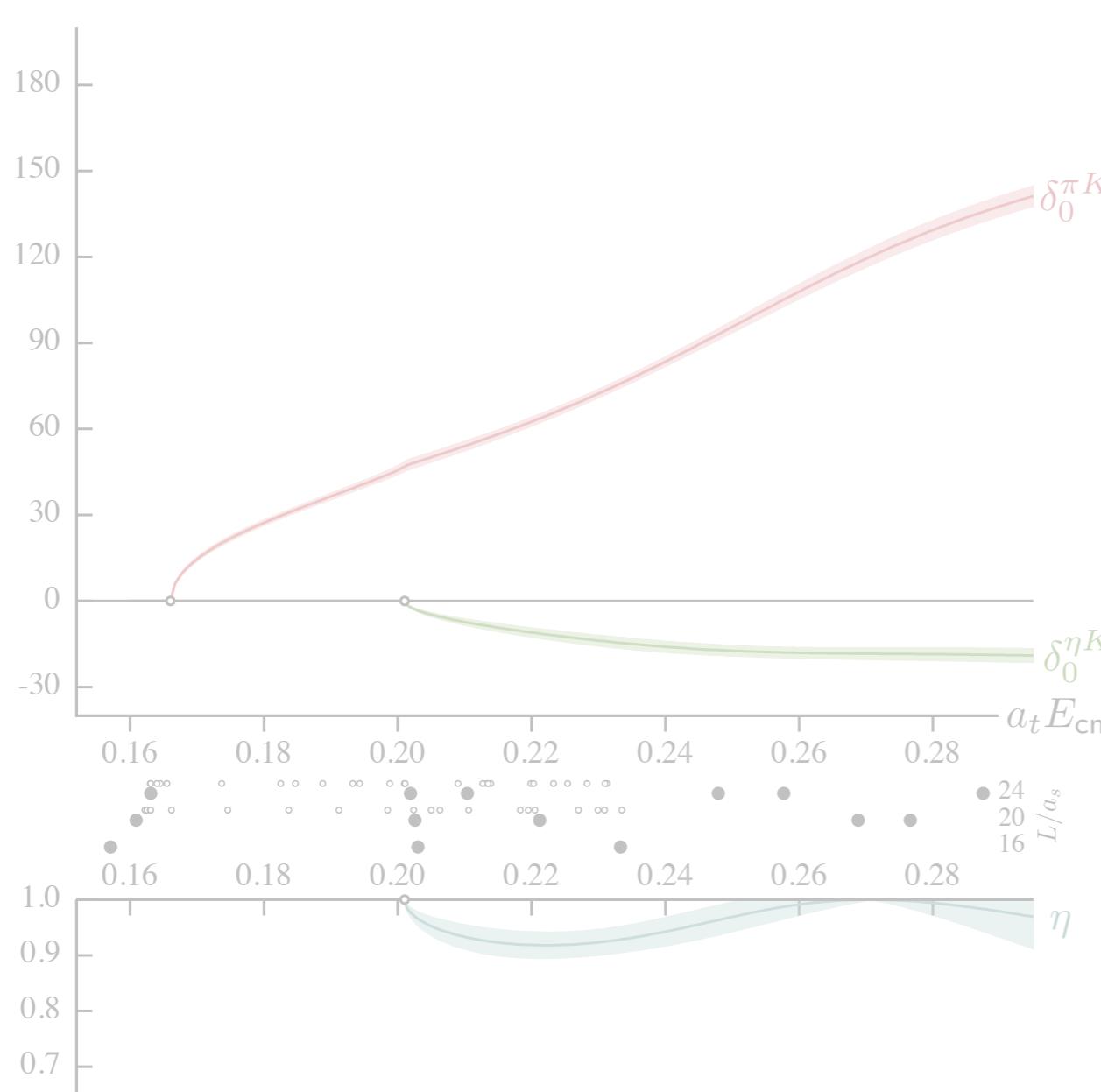


versus experimental scattering

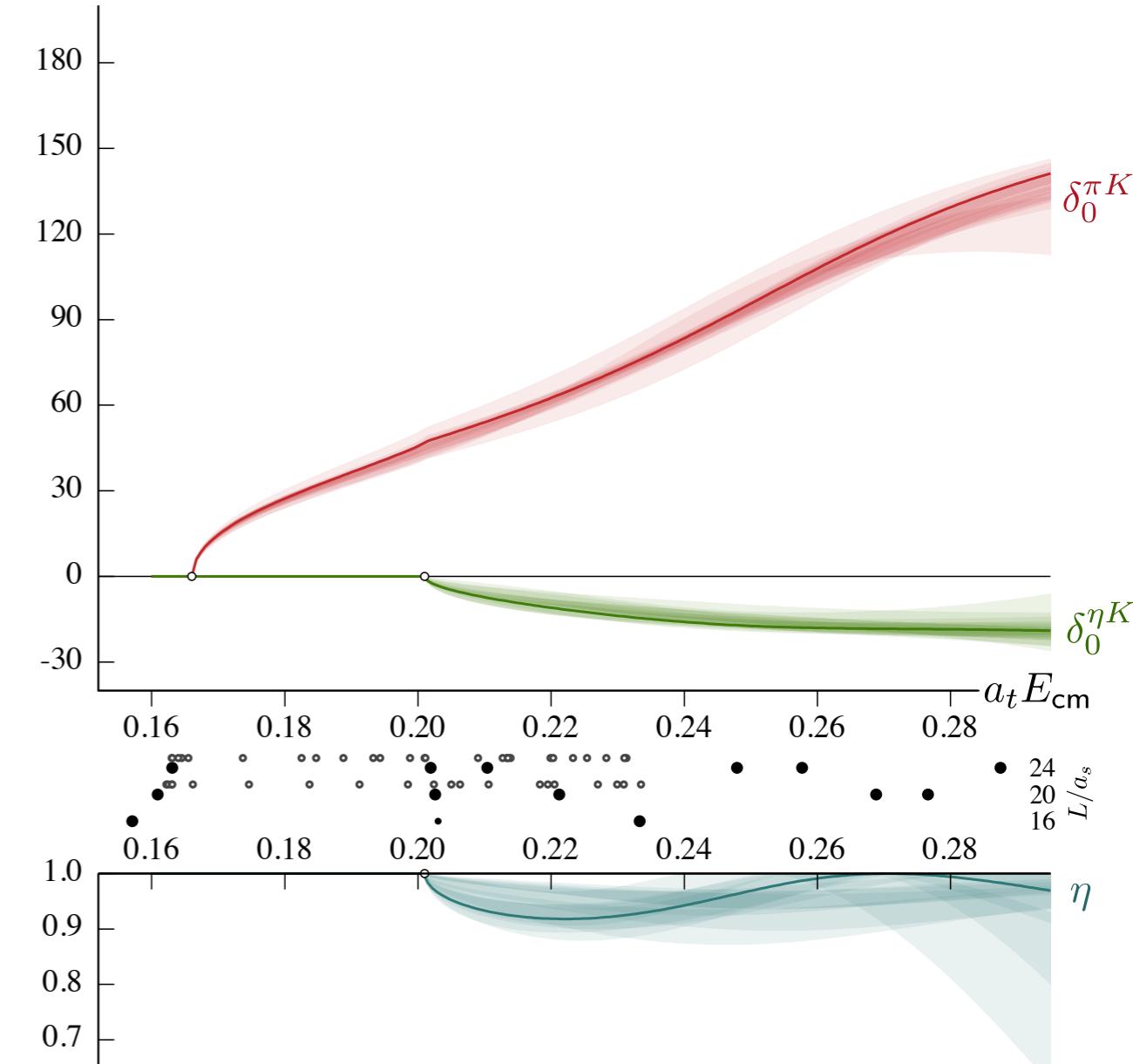


LASS, NPB296 493

- are the result parameterization dependent ?
 - try a range of parameterizations ...



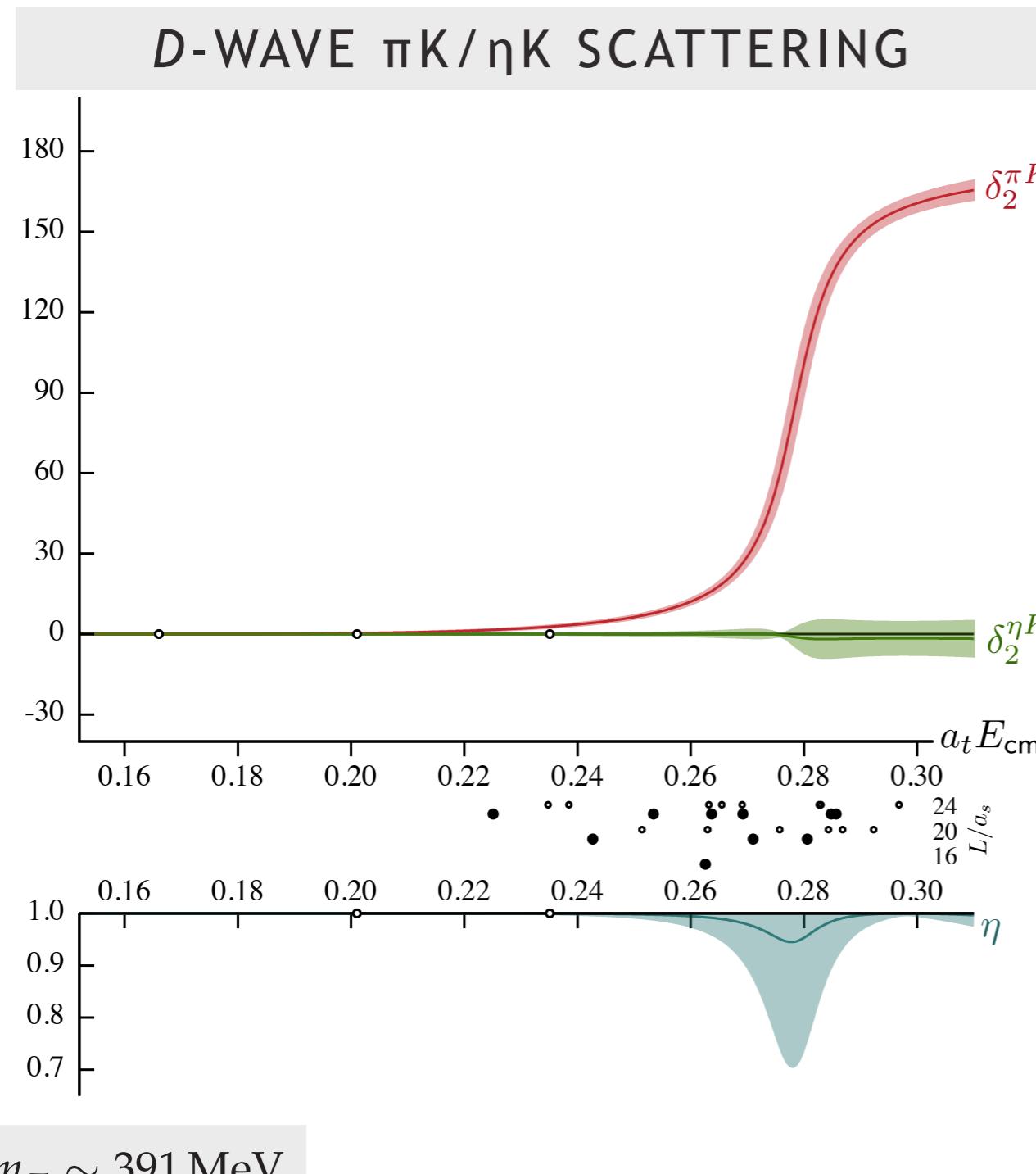
S-WAVE $\pi K/\eta K$ SCATTERING



– gross features are robust

$\pi K/\eta K$ scattering

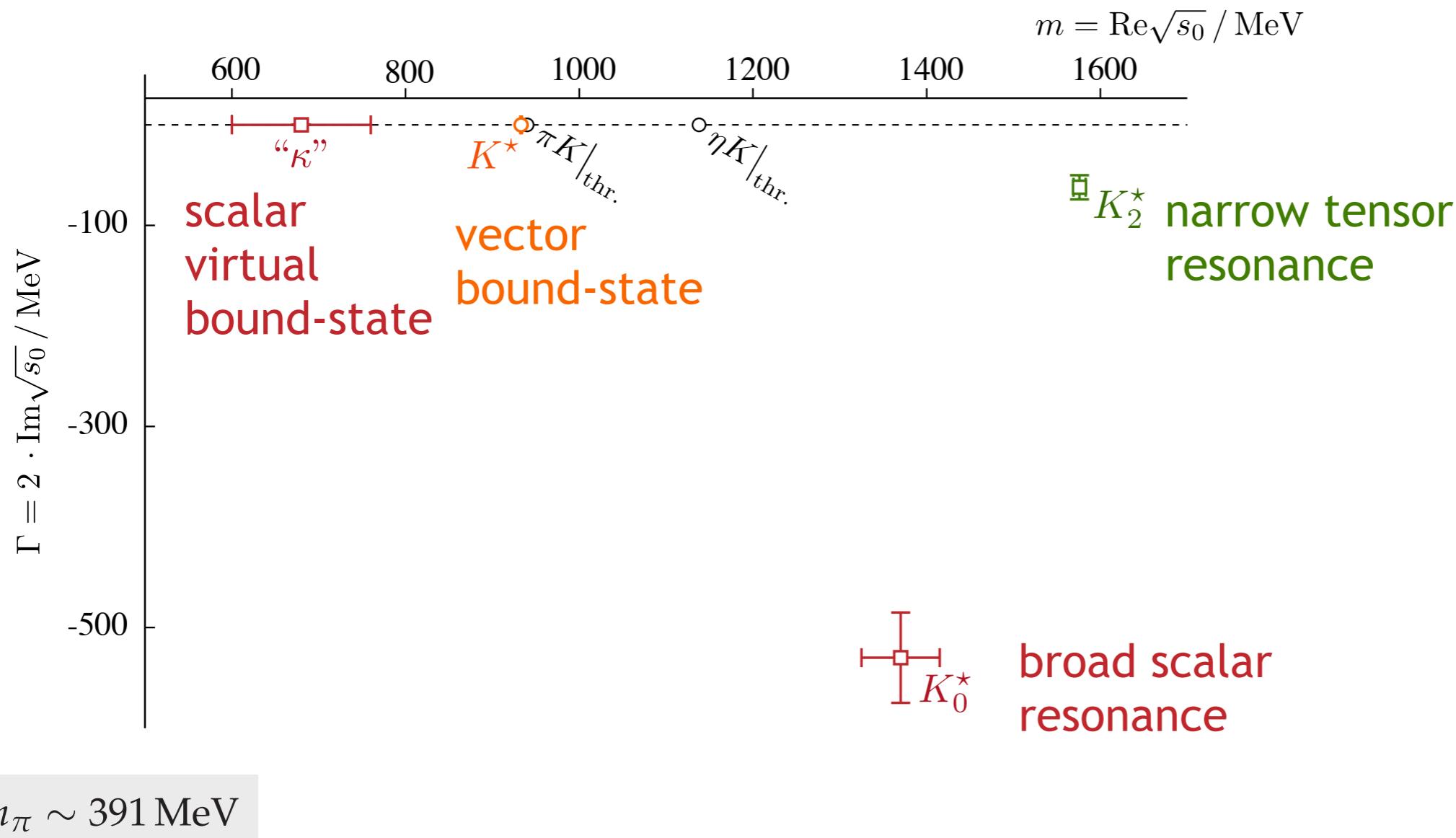
- clear narrow resonance in D -wave scattering



$m_\pi \sim 391 \text{ MeV}$

singularity content

- t -matrix poles as least model-dependent characterization of resonances



hadron spectroscopy from lattice QCD

- excited hadron spectra from lattice QCD
 - qualitative features of the hadron spectrum
 - hybrid hadrons seem to play a role
- hadrons as resonances in scattering amplitudes from lattice QCD
 - elastic & coupled-channel
 - but at large quark masses - challenges to go lower

computational - big volumes, expensive

physics - multi-hadron final states