

クォークと超弦理論

橋本幸士

橋本数理物理学研究室

2014.12.26 仁科センター コロキウム

クォークと異次元

橋本幸士

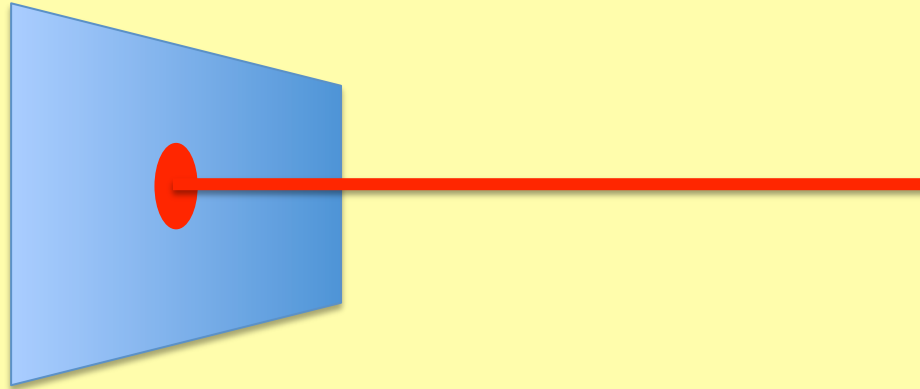
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Preprint archive system

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RIKEN-MP-94	On the stability of the vacuum in the Standard Model	Taro Kimura	2014.10
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RIKEN-MP-87	On the stability of the vacuum in the Standard Model	Masato Taki	2014.03
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0 クォーク = 異次元方向に伸びた弦



1 チャーモニウム と ブラックホール

2 重イオン衝突 と「とがった」異次元

3 原子核 と 隠れた対称性

HEP

3 のレコードが見つかりました。

1. The Large N limit of superconformal field theories and supergravity

(10389) [Juan Martin Maldacena](#) ([Harvard U.](#)). Nov 1997. 19 pp.
Published in [Int.J.Theor.Phys.](#) **38** (1999) 1113-1133, [Adv.Theor.Math.Phys.](#) **2** (1998) 231-252
HUTP-97-A097, HUTP-98-A097
DOI: [10.1023/A:1026654312961](#)
e-Print: [hep-th/9711200](#) | [PDF](#)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)
[ADS Abstract Service](#)

[レコードの詳細](#) - [Cited by 10389 records](#) 1000+

AdS/CFT対応

2. A Model of Leptons

(9196) [Steven Weinberg](#) ([MIT](#), [LNS](#)). Nov 1967. 3 pp.
Published in [Phys.Rev.Lett.](#) **19** (1967) 1264-1266
DOI: [10.1103/PhysRevLett.19.1264](#)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)
[ADS Abstract Service](#); [Phys. Rev. Lett. Server](#)

[レコードの詳細](#) - [Cited by 9196 records](#) 1000+

3. Measurements of Omega and Lambda from 42 high redshift supernovae

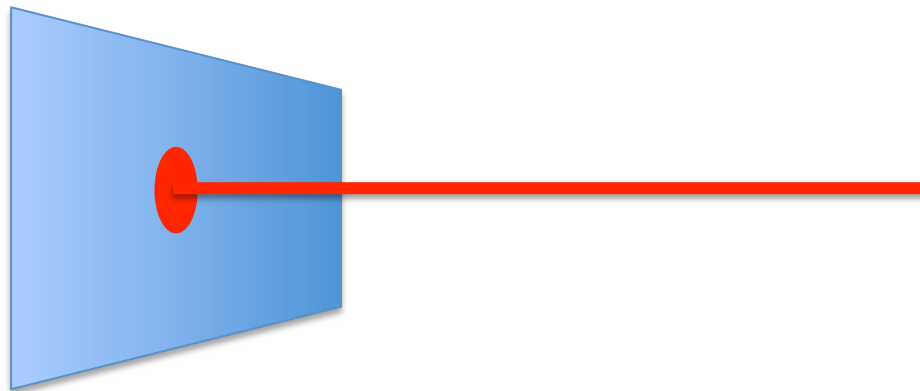
(8187) [Supernova Cosmology Project Collaboration](#) ([S. Perlmutter](#) ([UC](#), [Berkeley](#), [CfPA](#)) *et al.*). Dec 1998. 33 pp.
Published in [Astrophys.J.](#) **517** (1999) 565-586
LBNL-41801, LBL-41801
DOI: [10.1086/307221](#)
e-Print: [astro-ph/9812133](#) | [PDF](#)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)
[CERN Document Server](#); [ADS Abstract Service](#); [OSTI Information Bridge Server](#)

[レコードの詳細](#) - [Cited by 8187 records](#) 1000+

0

クォーク = 異次元方向に伸びた弦



0-1

クォーク間ポテンシャル = 弦の影

0-2

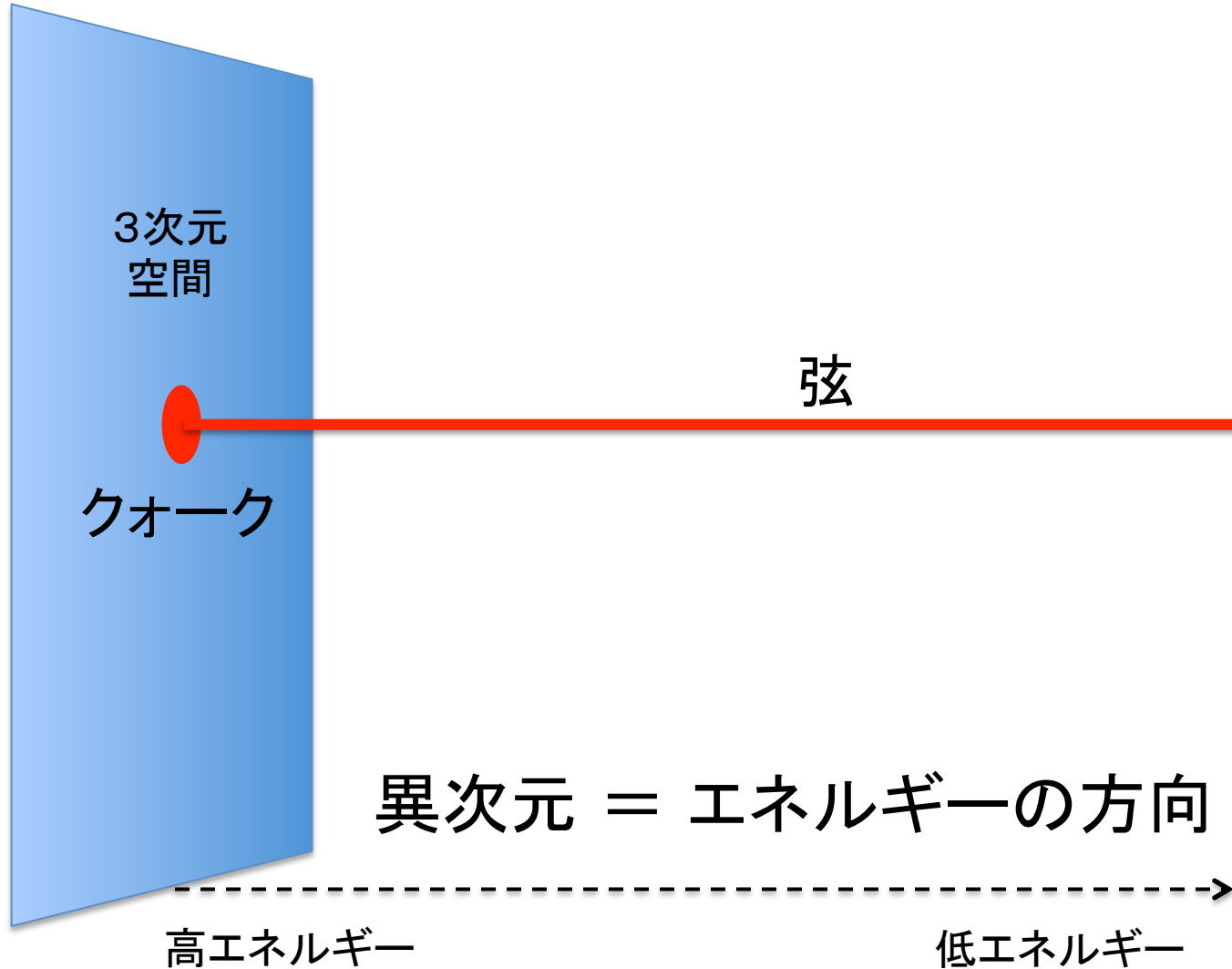
QCDスケール = 異次元の終わり

0-3

閉じ込め力 = 終わりを這う弦

0-0

クォーク = 異次元方向に伸びた弦



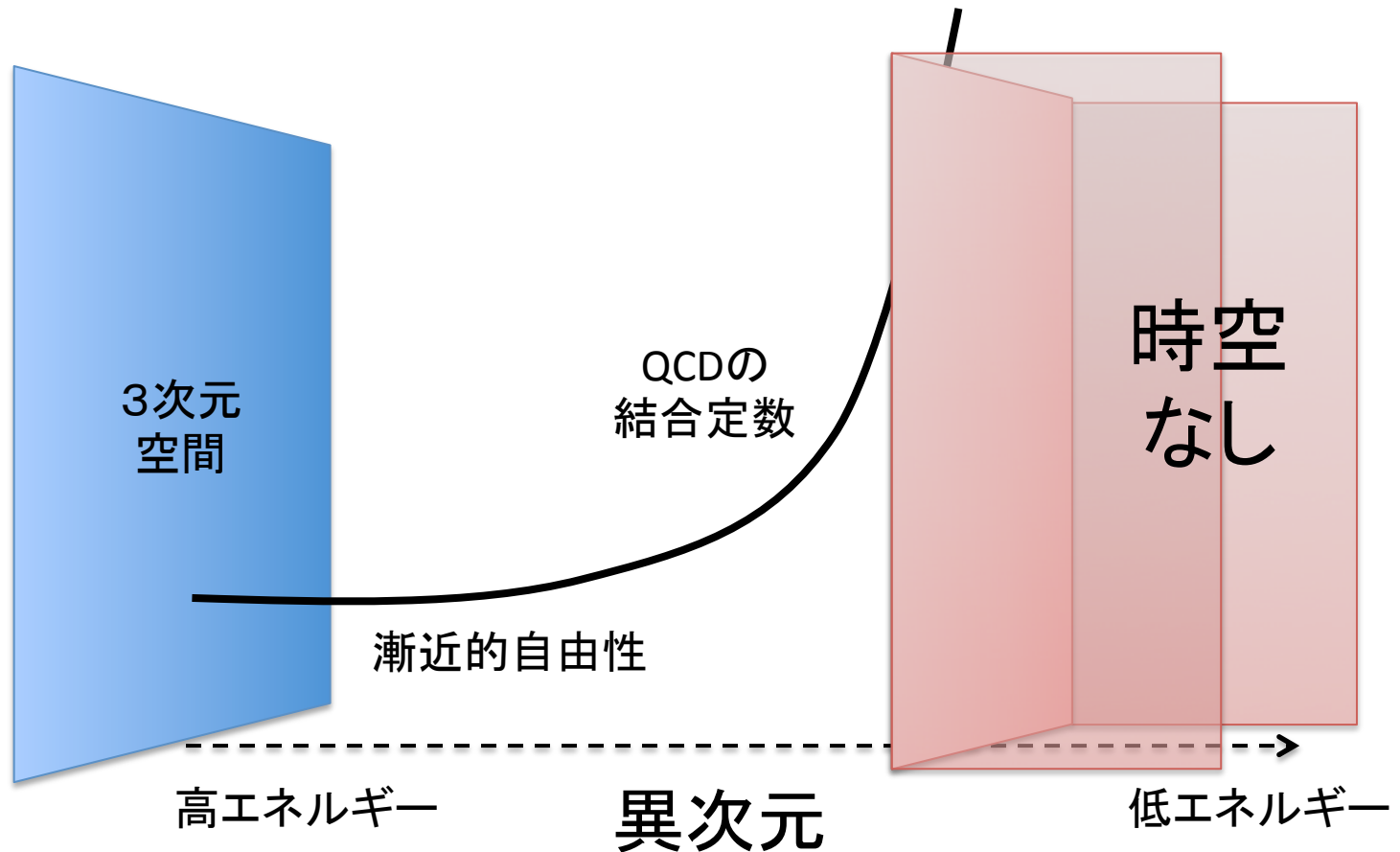
0-1

クォーク間ポテンシャル = 弦の影



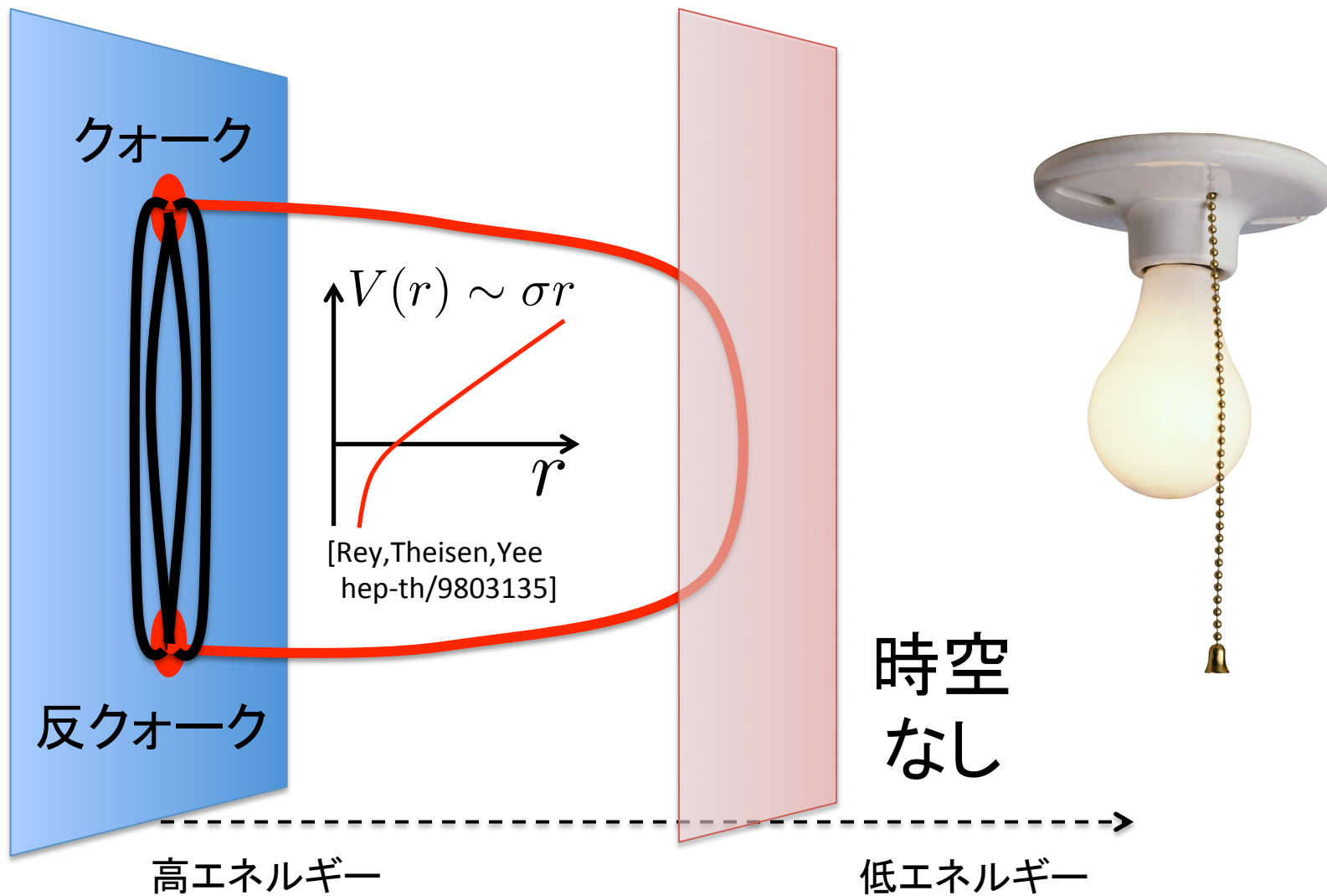
0-2

QCDスケール = 異次元の終わり

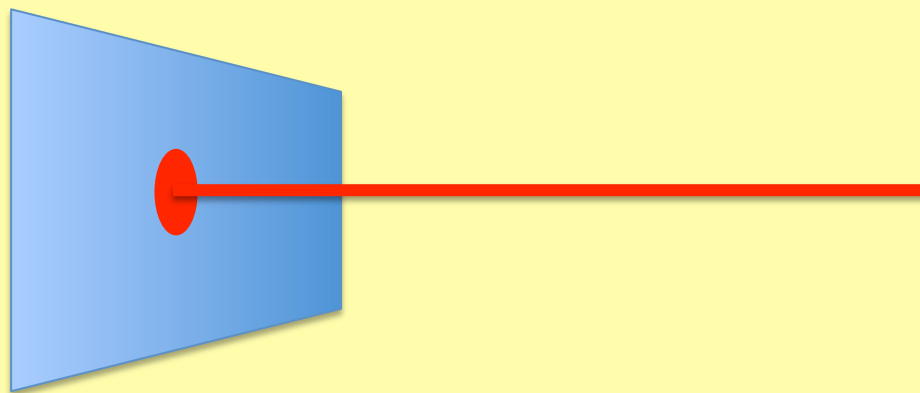


0-3

閉じ込め力 = 終わりを這う弦



0 クォーク = 異次元方向に伸びた弦



1 チャーモニウム と ブラックホール

2 重イオン衝突 と「とがった」異次元

3 原子核 と 隠れた対称性

1

チャーモニウム と ブラックホール

ArXiv:1411.0618

OU-HET-837, RIKEN-MP-95

Entropic destruction of heavy quarkonium in non-Abelian plasma from holography

Koji Hashimoto^{1,2} and Dmitri E. Kharzeev^{3,4}

¹Department of Physics, Osaka University, Toyonaka, Osaka 560-0043, Japan

²Mathematical Physics Lab., RIKEN Nishina Center, Saitama 351-0198, Japan

1-1

チャーモニウムの謎

Lattice QCD indicates a large amount of entropy associated with the heavy quark-antiquark pair immersed in the quark-gluon plasma. This entropy grows as a function of the inter-quark distance giving rise to an entropic force that can be very effective in dissociating the bound quarkonium

1-2

クォーク対のエントロピー: 格子VS超弦

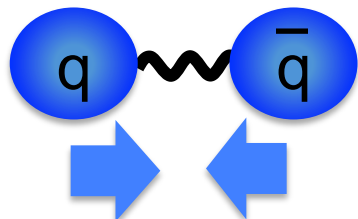
with the heavy quark-antiquark pair in two theories: i) $\mathcal{N} = 4$ supersymmetric Yang-Mills and ii) a confining Yang-Mills theory obtained by compactification on a Kaluza-Klein circle. In both cases we find the entropy growing with the inter-quark distance and evaluate the effect of the corresponding

1-3

ブラックホールを見ている

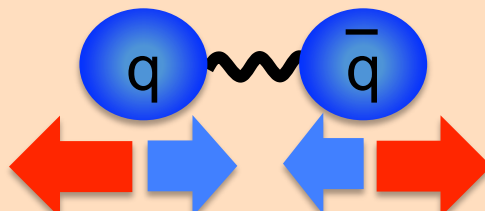
that describes the deconfinement transition. In terms of the boundary theory, this entropy likely emerges from the entanglement of a “long string” connecting the quark and antiquark with the rest of the system.

真空

 $T=0$ 

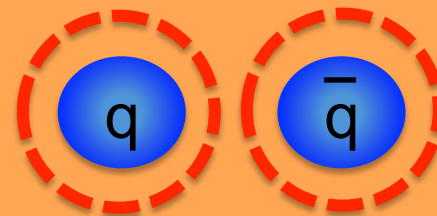
閉じ込め力

RHIC

 $T=300 \text{ MeV}$ 

エントロピー力

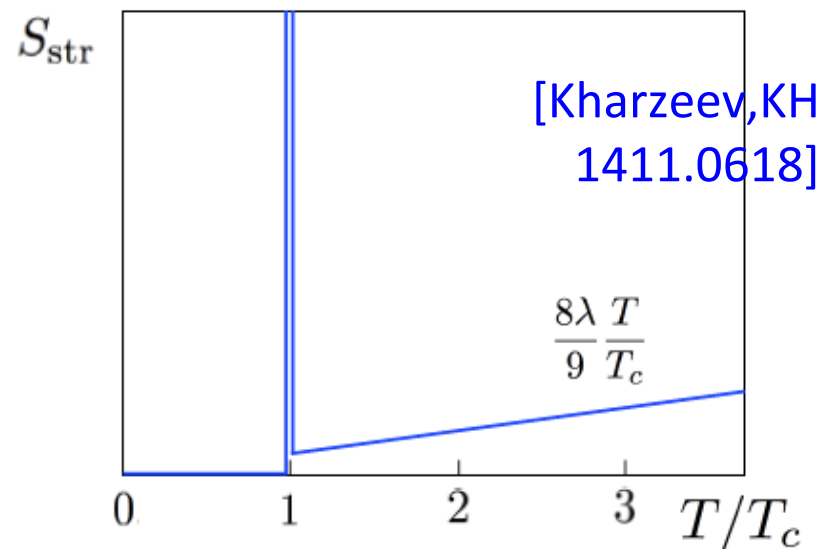
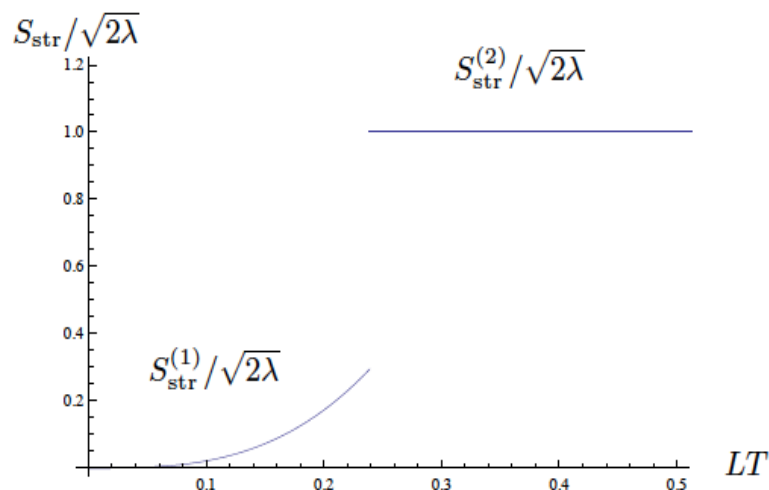
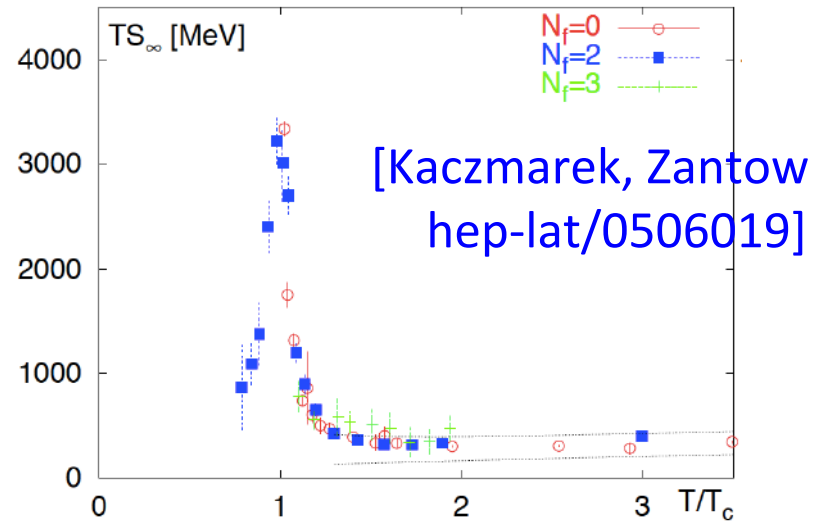
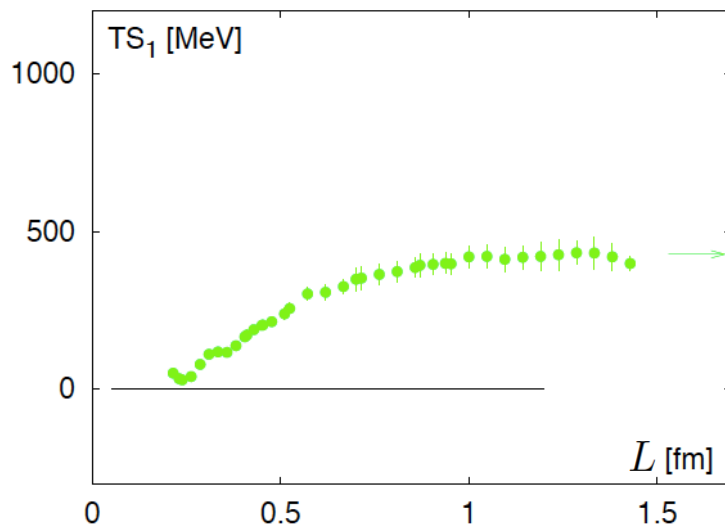
LHC

 $T=400 \text{ MeV}$ 

デバイ遮蔽

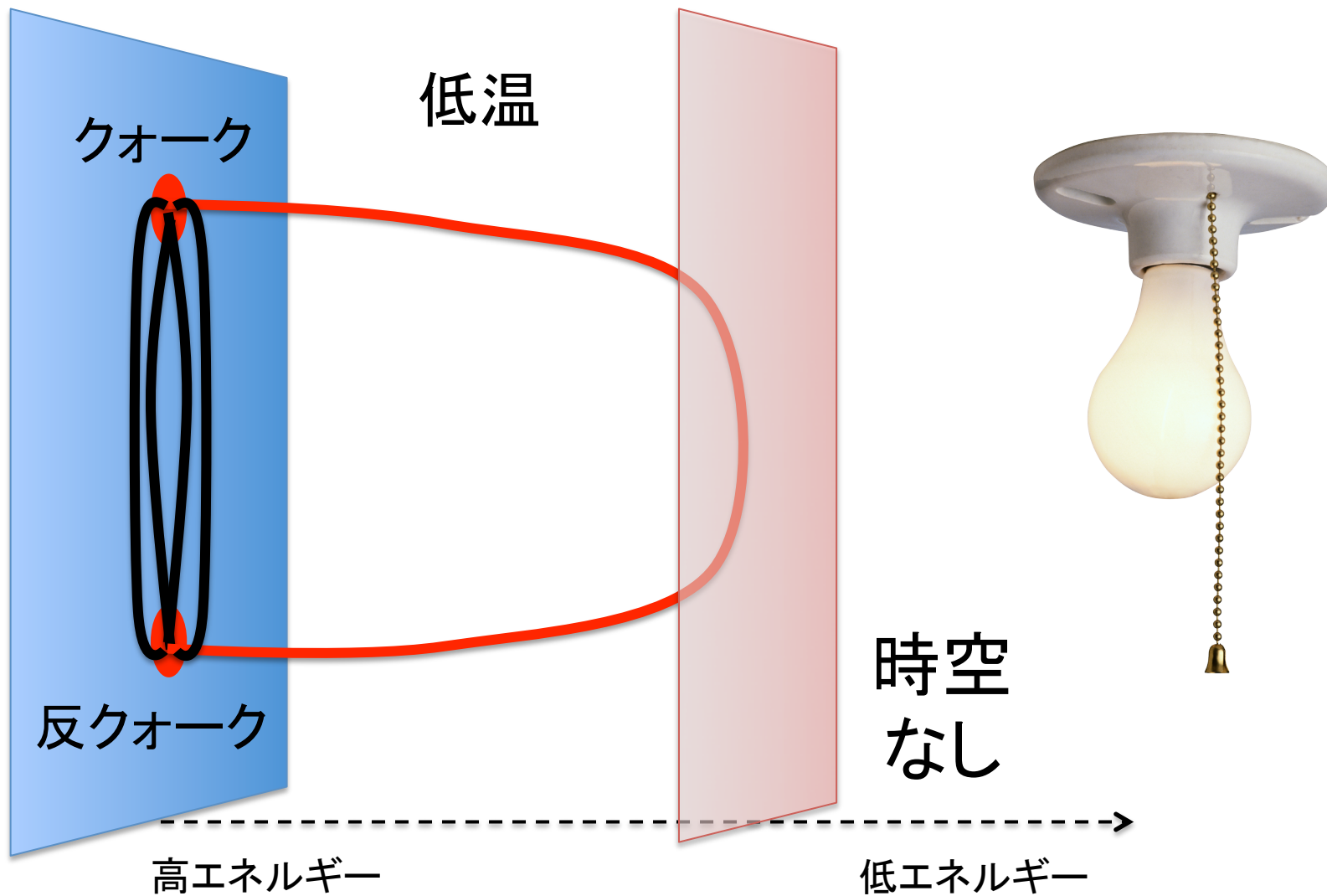
$$F = T \frac{\partial S}{\partial L}$$

1-2 クォーク対のエントロピー: 格子 VS 超弦



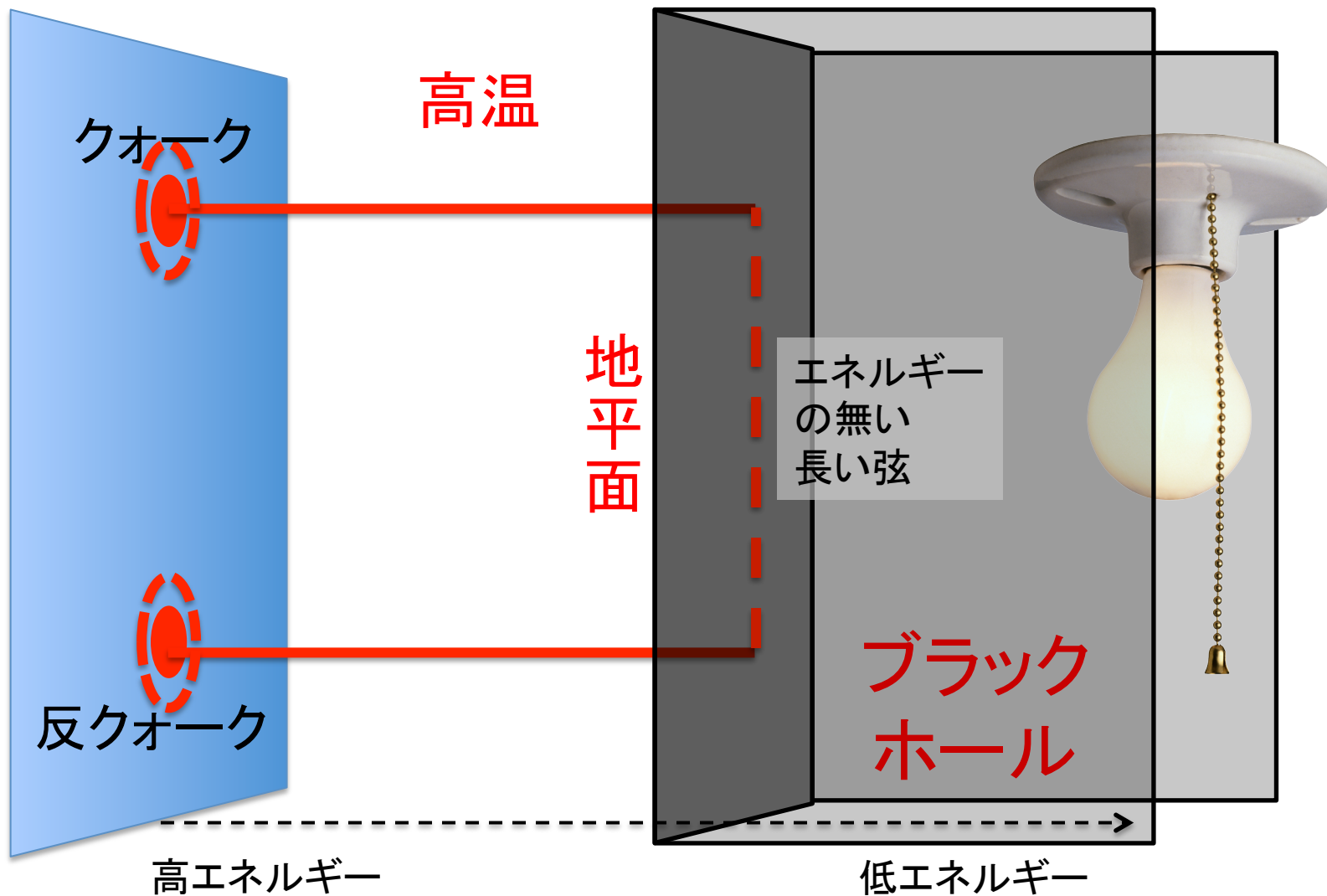
1-3

ブラックホールを見ている

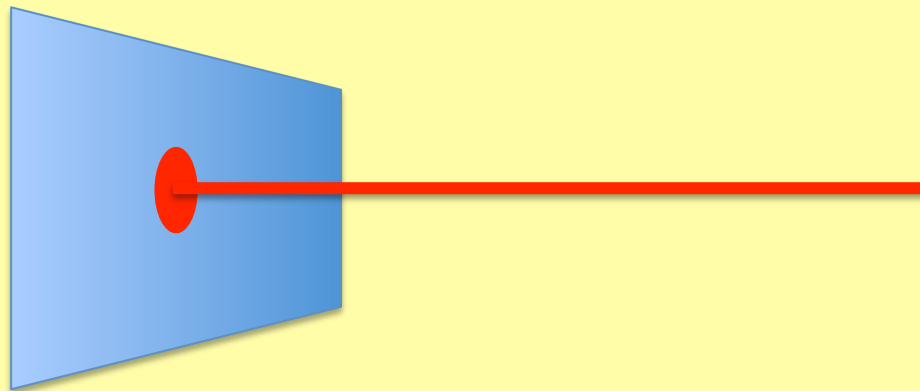


1-3

ブラックホールを見ている



0 クォーク = 異次元方向に伸びた弦



1 チャーモニウム と ブラックホール

2 重イオン衝突 と「とがった」異次元

3 原子核 と 隠れた対称性

2

重イオン衝突 と「とがった」異次元

ArXiv:1412.4964

Meson turbulence at quark deconfinement from AdS/CFT

Koji Hashimoto^{1,2}, Shunichiro Kinoshita³, Keiju Murata⁴, and Takashi Oka⁵¹Department of Physics, Osaka University, Toyonaka, Osaka 560-0043, Japan²Mathematical Physics Lab., RIKEN Nishina Center, Saitama 351-0198, Japan³Osaka City University Advanced Mathematical Institute, Osaka 558-8585, Japan⁴Keio University, 4-1-1 Hiyoshi, Yokohama 223-8521, Japan and⁵Department of Applied Physics, University of Tokyo, Tokyo 113-8656, Japan

1412.4254

1408.6293

1407.0798

1403.6336

1307.7423

Based on the QCD string picture at confining phase, we conjecture that the deconfinement tran-

2-1

クォーク質量はDブレーンで

distribution to each meson level n scales as n^α with the universal scaling $\alpha = -5$. The universality is checked for various ways to attain the quark deconfinement: a static electric field below/around the critical value, a time-dependent electric field quench, and a time-dependent quark mass quench, all result in the turbulent meson condensation with the universal power $\alpha = -5$ around the decon-

2-2

電場でとがるDブレーン

I. TURBULENCE AND QUARK DECONFINEMENT

How the quarks are confined at the vacuum of quantum chromodynamics (QCD) is one of the most fundamental

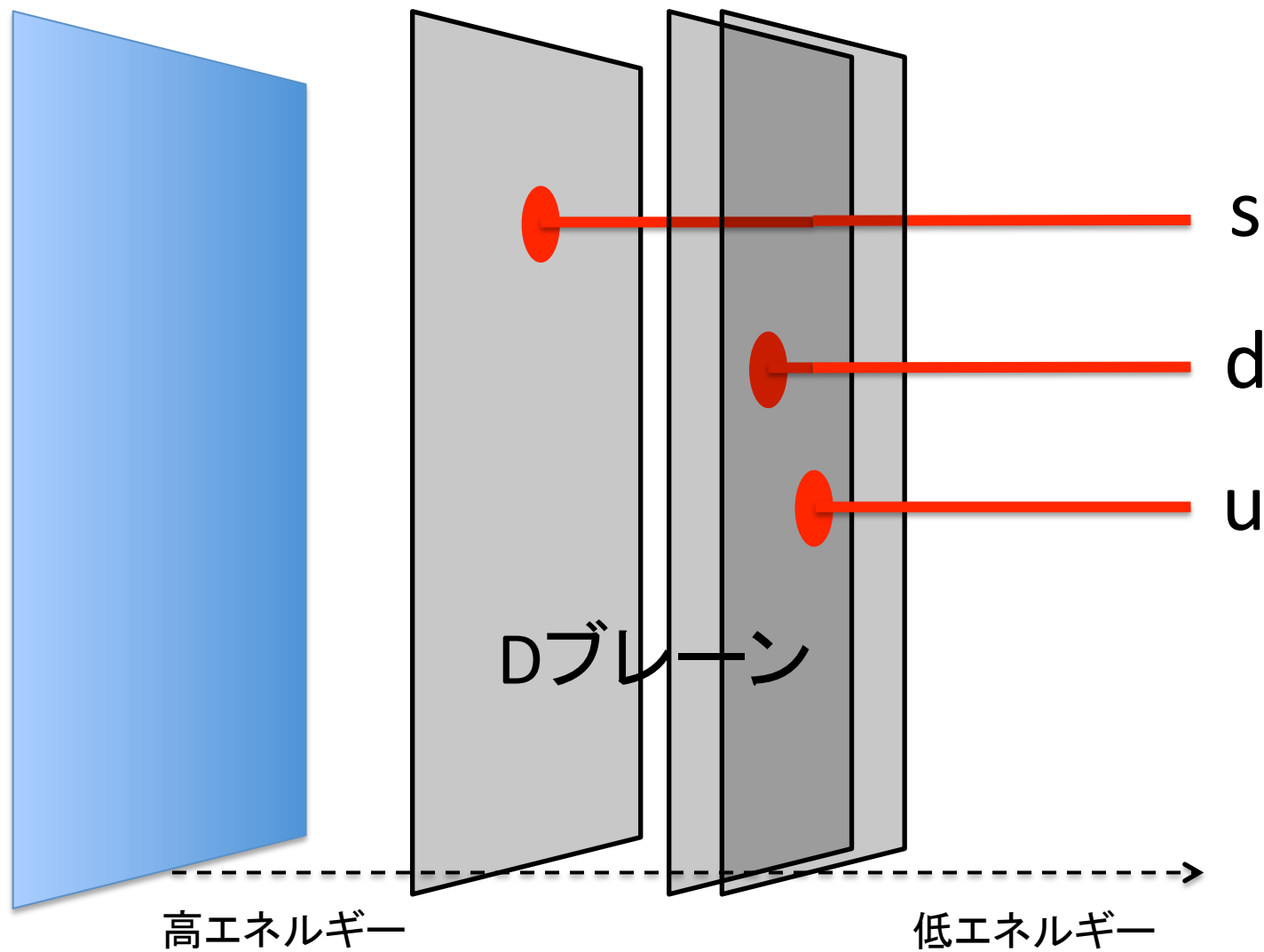
2-3

重イオン衝突: 変化する弱電場で非閉じ込め

in various manner as one departs from the vacuum with the help of some external forces. The forces include for example a finite temperature, a finite quark density and electric fields. Depending on how you break the vacuum confinement, the resultant deconfined phases show various aspects with various global symmetries. This variety makes

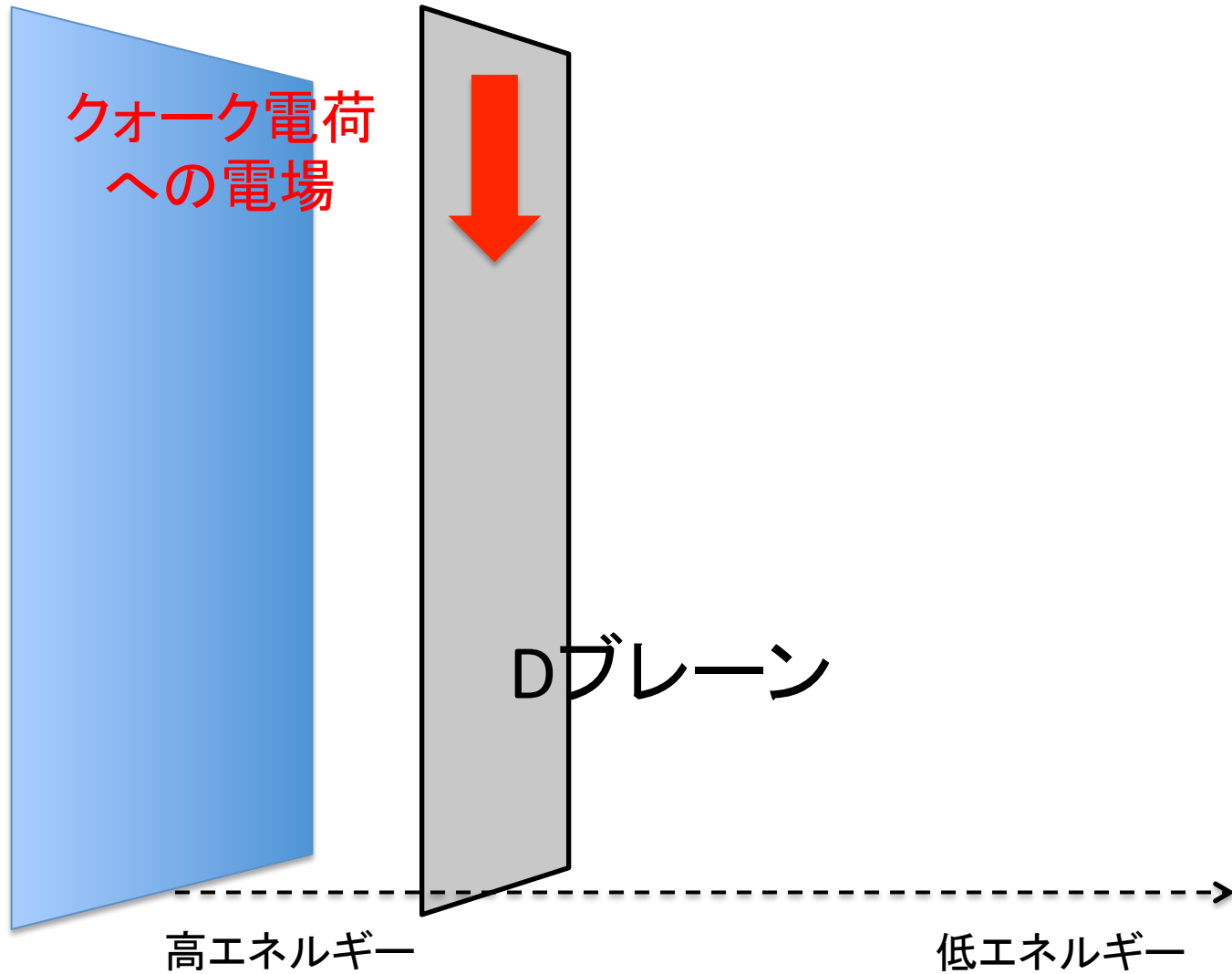
2-1

クォーク質量はDブレーンで



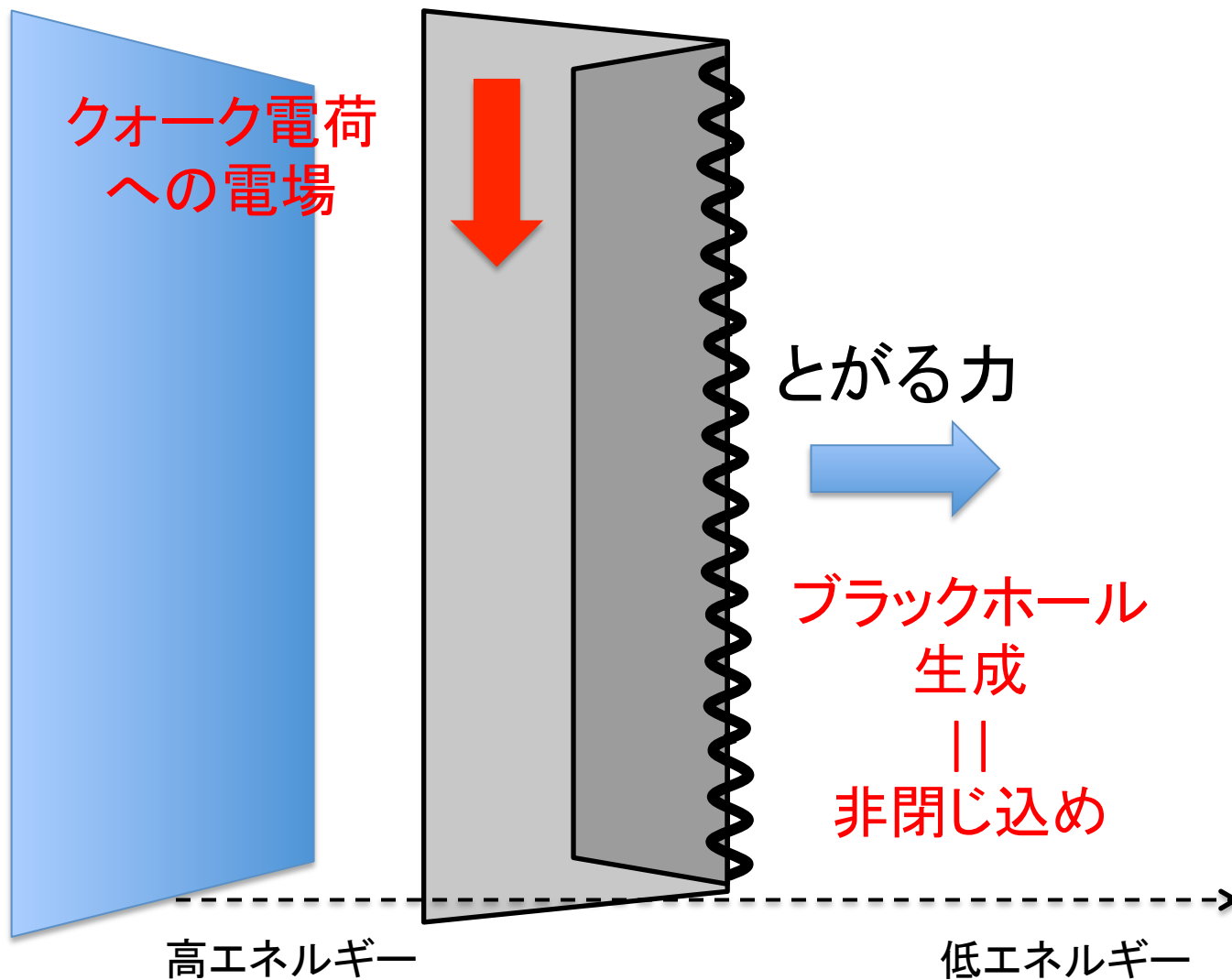
2-2

電場でとがるDブレーン



2-2

電場でとがるDブレーン



衝突では時間変化する弱い電場が存在

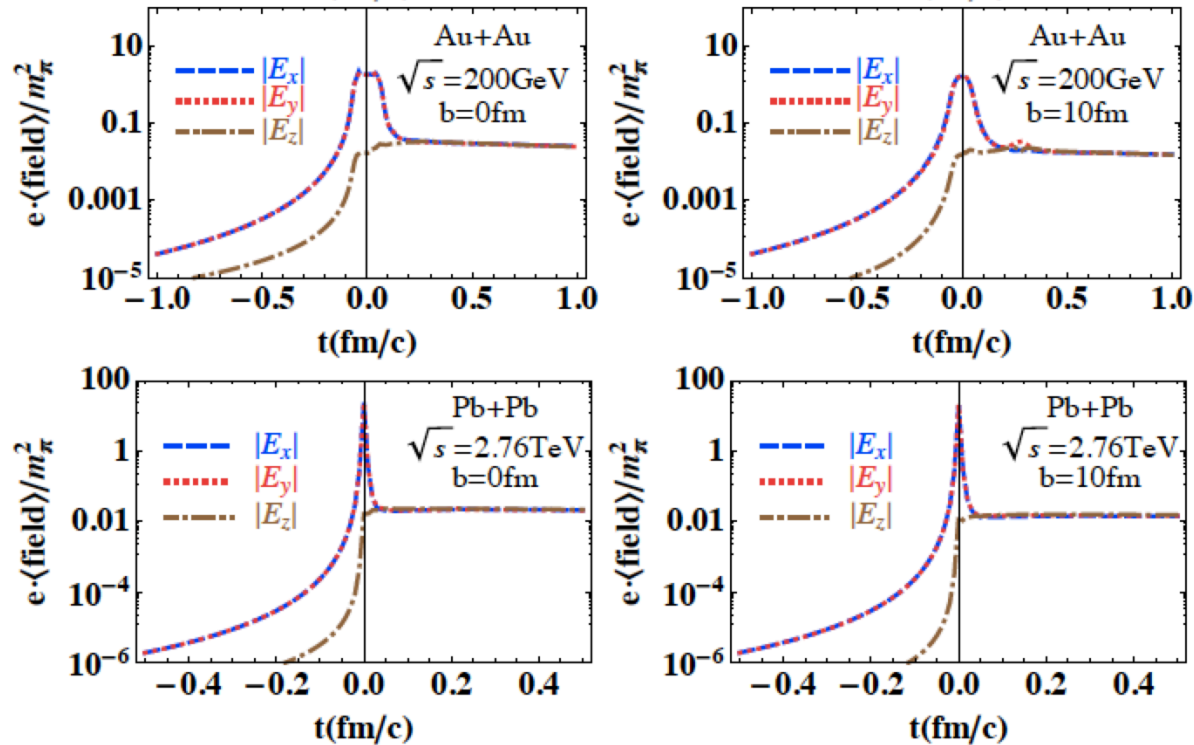


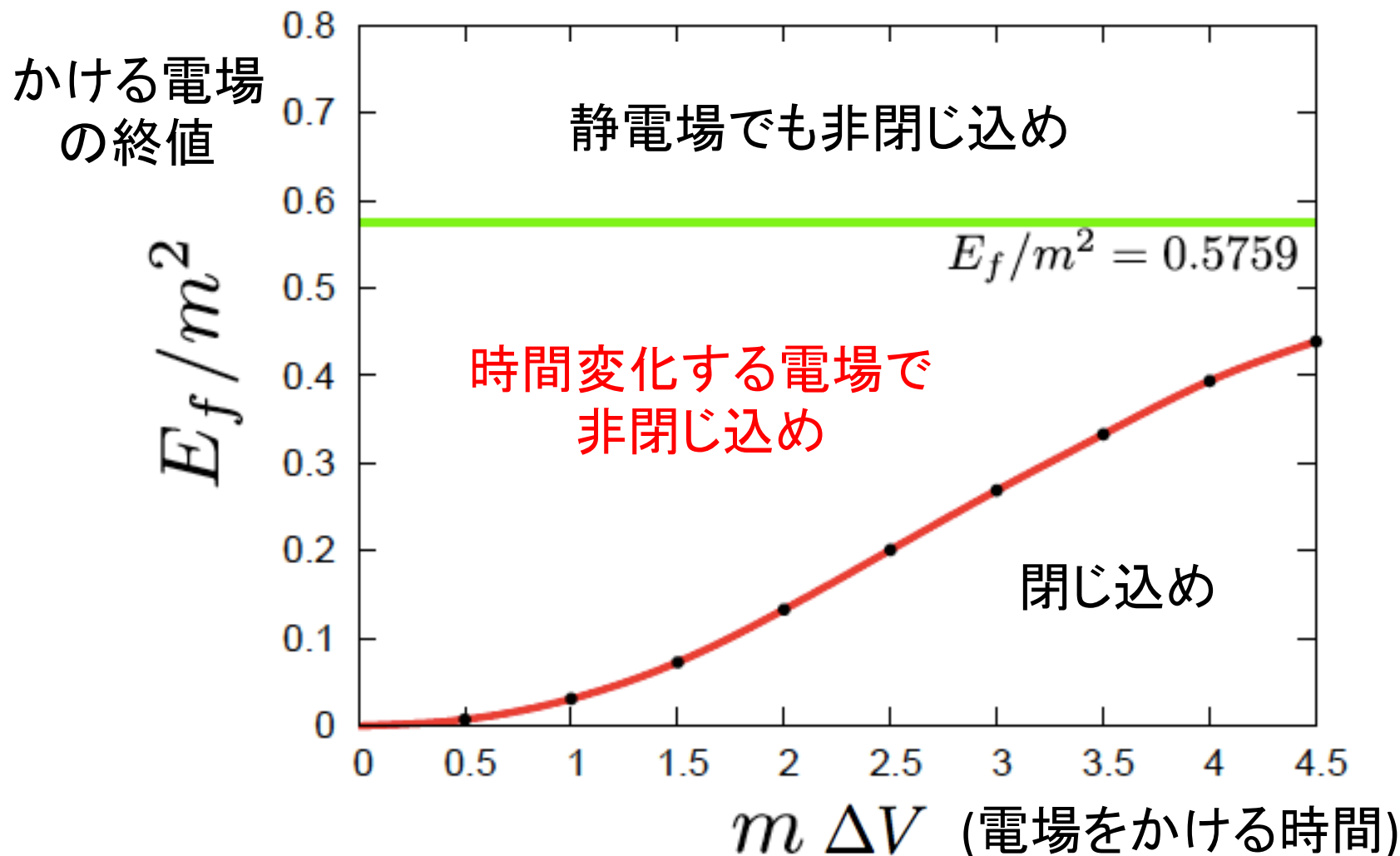
FIG. 7: (Color online) The time evolution of electromagnetic fields at $r = 0$ with impact parameter $b = 0$, and $b = 10$ for Au + Au collisions at $\sqrt{s} = 200 \text{ GeV}$ and Pb + Pb collisions at $\sqrt{s} = 2.76 \text{ TeV}$. After collision, the remnants can essentially slow down the decay of the transverse fields, and enhance the longitudinal fields.

[Deng, Huang, 1201.5108]

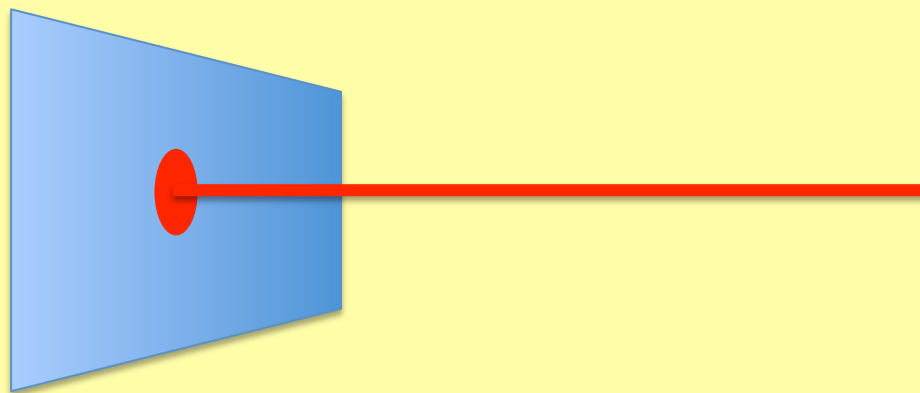
2-3

重イオン衝突: 変化する弱電場で非閉じ込め

N=2超対称QCD の ダイナミカル相図



0 クォーク = 異次元方向に伸びた弦



1 チャーモニウム と ブラックホール

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3 原子核 と 隠れたゲージ対称性

3

原子核 と 隠れた対称性

A Matrix Model for Baryons and Nuclear Forces

ArXiv:1003.4988

0809.3141

0910.2303

0911.1035

1005.4412

1006.3612

1103.5688

1203.5386

3-1

核子は小さなDブレーン

[†] Theory Division, CERN, CH-1211 Geneva 23, Switzerland

[‡] School of Physics, Korea Institute for Advanced Study, Seoul 130-722, Korea

3-2

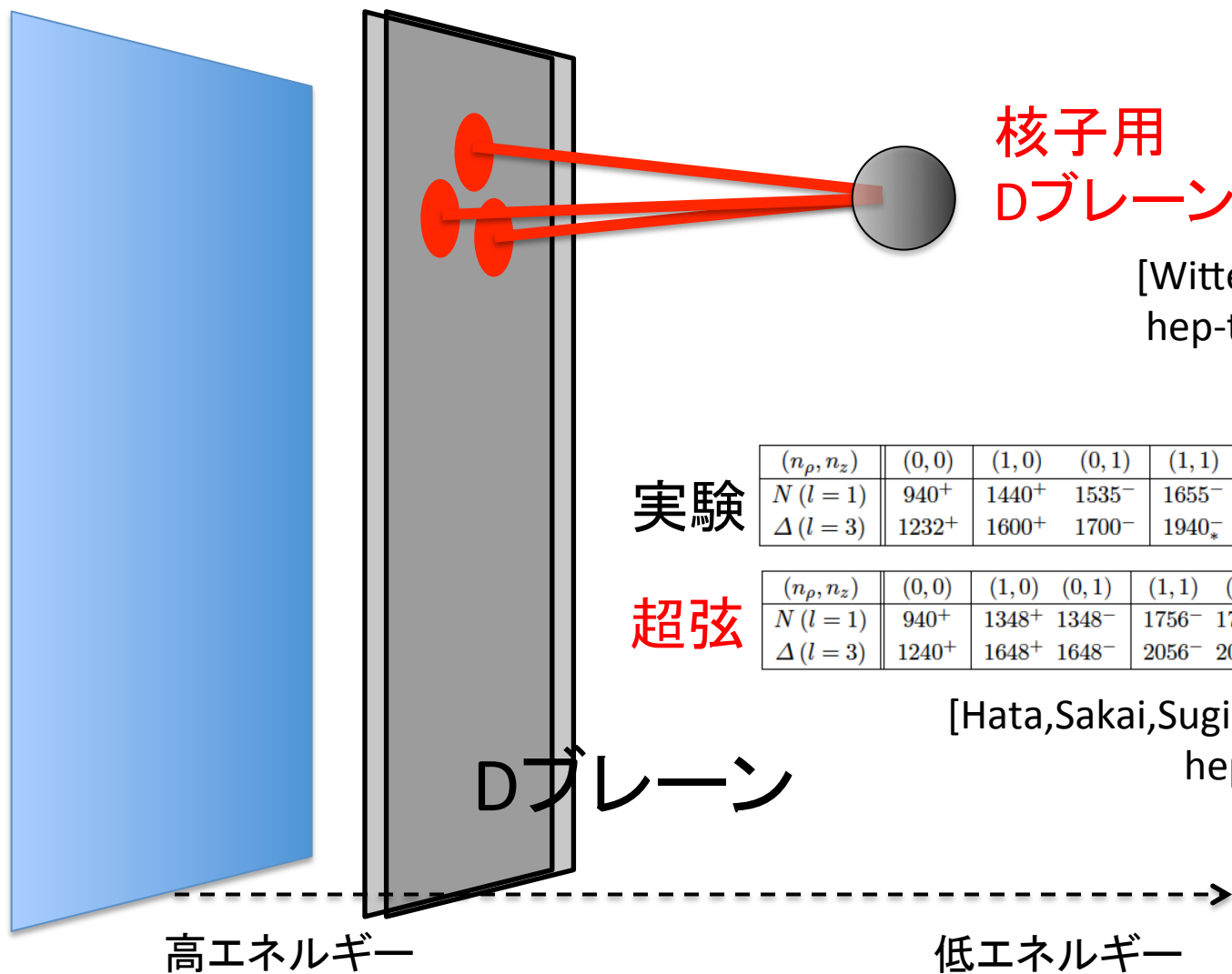
核力は弦の励起

3-3

原子核は行列

embedded in the D4-D8 model of large N_c holographic QCD. The positions of k baryons are unified into $k \times k$ matrices, with spin/isospin of the baryons encoded in

核子は小さなDブレーン



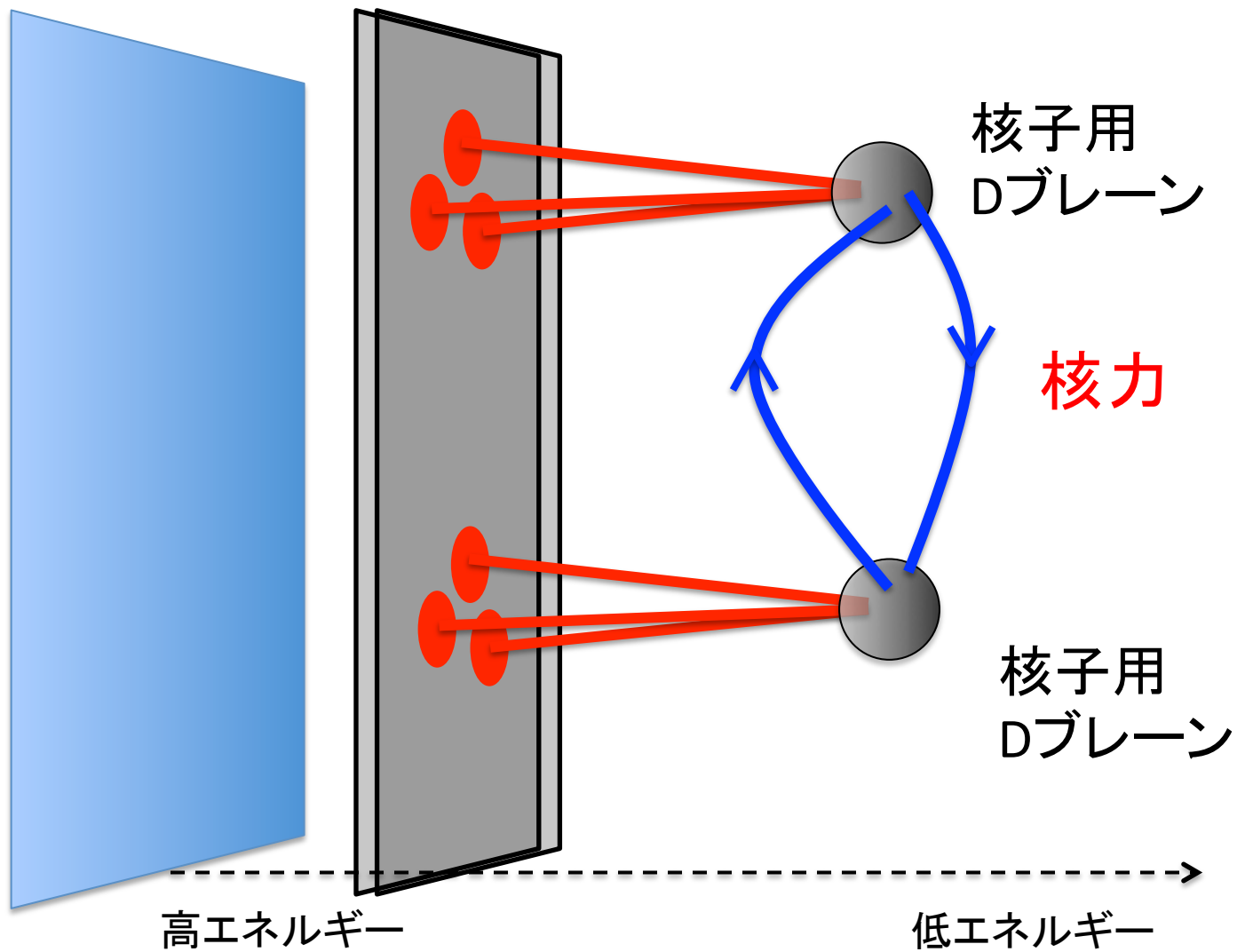
(n_ρ, n_z)	(0, 0)	(1, 0)	(0, 1)	(1, 1)	(2, 0)/(0, 2)	(2, 1)/(0, 3)
$N(l=1)$	940 ⁺	1440 ⁺	1535 ⁻	1655 ⁻	1710 ⁺ , ?	2090 ⁻ , ?
$\Delta(l=3)$	1232 ⁺	1600 ⁺	1700 ⁻	1940 [*]	1920 ⁺ , ?	?, ?

(n_ρ, n_z)	(0, 0)	(1, 0)	(0, 1)	(1, 1)	(2, 0)/(0, 2)	(2, 1)/(0, 3)
$N(l=1)$	940 ⁺	1348 ⁺	1348 ⁻	1756 ⁻	1756 ⁺ , 1756 ⁺	2164 ⁻ , 2164 ⁻
$\Delta(l=3)$	1240 ⁺	1648 ⁺	1648 ⁻	2056 ⁻	2056 ⁺ , 2056 ⁺	2464 ⁻ , 2464 ⁻

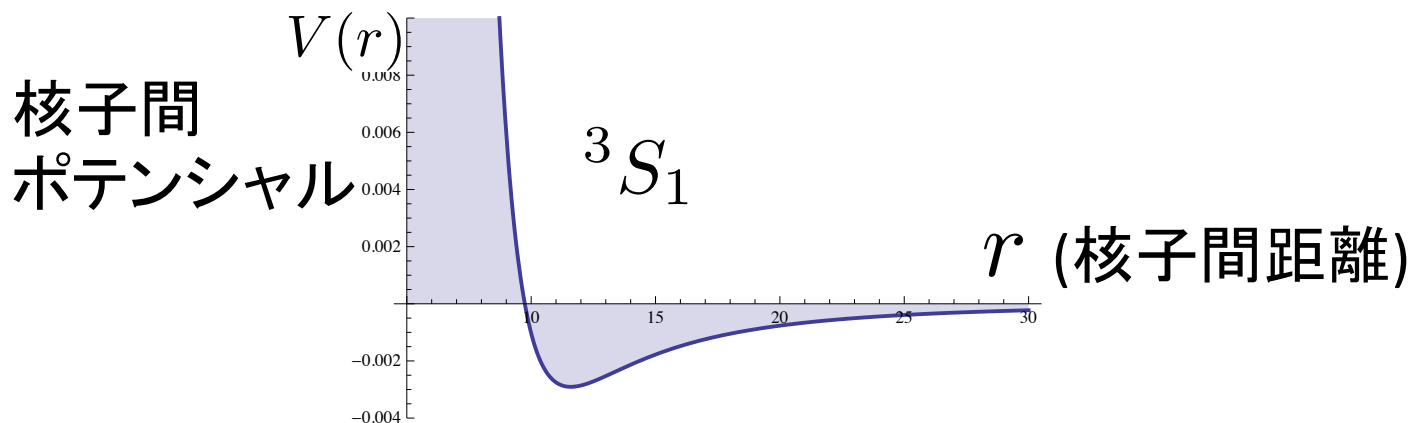
[Hata, Sakai, Sugimoto, Yamato
hep-th/0701280]

3-2

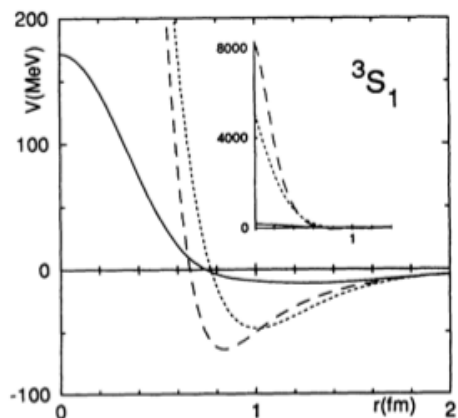
核力は弦の励起



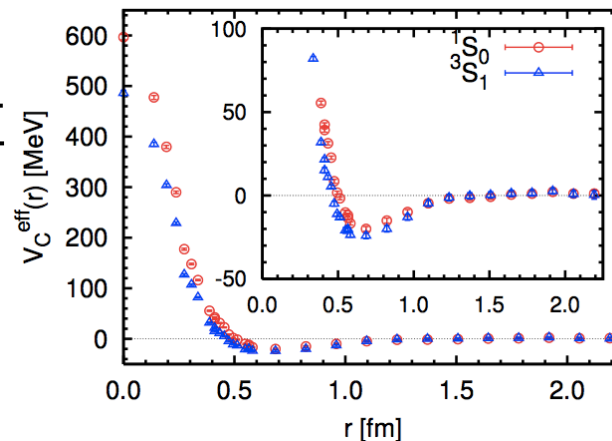
超弦理論:



実験

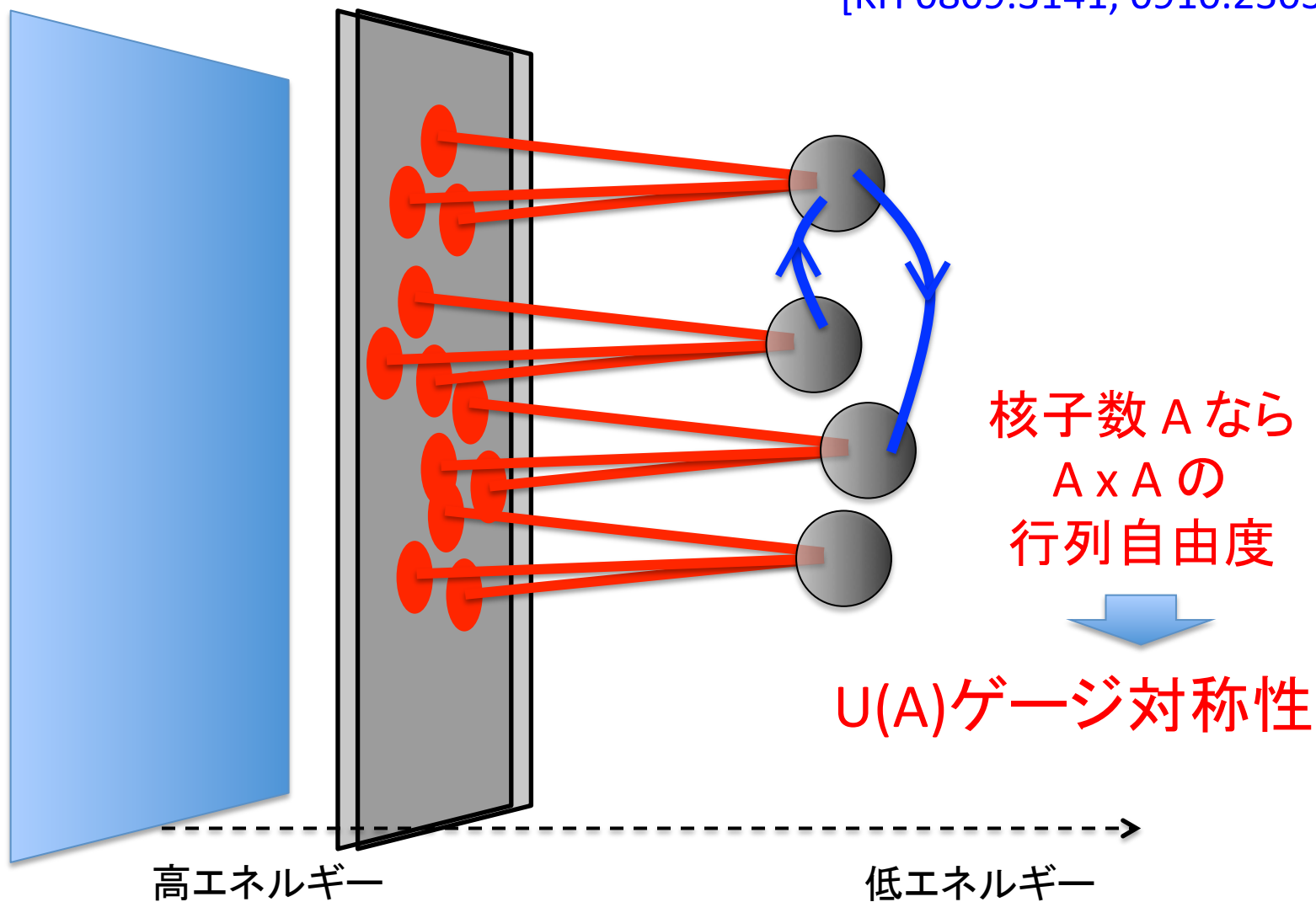


[Stoks,Klomp,Terheggen,deSwart ('94)]

格子QCD
数値計算

[Aoki,Ishii,Hatsuda ('07)]

[KH 0809.3141, 0910.2303]

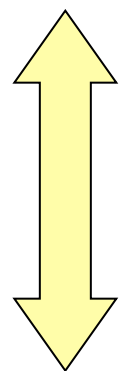


3-3

原子核 は 行列

[Iizuka,Yi, KH 1003.4988]

$$S = c \int dt \operatorname{tr} \left[\frac{1}{2} (D_t X^I)^2 - \frac{g^2}{4} [X^I, X^J]^2 \right] + \dots$$



X^M ($M = 1, 2, 3, 4$) : $A \times A$ エルミート行列

$X^{M=1,2,3}$ の固有値 = **A個の核子の位置**

パラメータは二つだけ: $c, g \leftrightarrow \lambda, M_{\text{KK}} \leftrightarrow m_\rho, f_\pi$

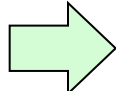
原子核は
量子力学

$$S \sim \int dt \left[\sum_{i=1}^A \frac{1}{2} m_N |\partial_t \vec{x}^{(i)}|^2 + \sum_{i \neq j} V(x^{(i)} - x^{(j)}, \dots) + \dots \right]$$

A=1 : 励起スペクトル

バリオンはフェルミオンで、スペクトルは： [Iizuka,Yi, KH 1003.4988]
[Iizuka, KH 1006.3612]

$$M = M_0 + \frac{1}{\sqrt{6}} \left[\sqrt{(I/2 + 1)^2 + N_c^2} + 2n_\rho + 2n_{X^4} + 2 \right]$$

 $I = J = 1/2$ $940^+, 1359^+, 1359^-, 1778^+, 1778^-, \dots$
 M_N, M_Δ 実験： $940^+, 1440^+, 1535^-, 1710^+, 1655^-, \dots$

A=2 : 斥力芯

$$V_T(\vec{r}) = 2\pi I_1^i I_2^i \frac{N_c}{\lambda M_{KK}} \frac{1}{r^2}$$

$$V_C(\vec{r}) = \pi \left(\frac{3^3}{2} - 8 I_1^i I_2^i J_1^j J_2^j \right) \frac{N_c}{\lambda M_{KK}} \frac{1}{r^2}$$

斥力芯を再現、 $1/r^2$ の依存性、3フレーバーでユニバーサル

A=3 : 三体力

[Iizuka, KH.1005.4412]

問題1： 原子核の形成

[T.Morita, KH 1103.5688]

問題： 次の量子力学を A が大きい時に解きなさい。

$$S = c \int dt \operatorname{tr} \left[\frac{1}{2} (D_t X^I)^2 - \frac{g^2}{4} [X^I, X^J]^2 \right]$$

答： 量子力学の真空：固有値が中心に集まる [Luscher, 1983]

→ 必然的に原子核が形成する

問題2： 原子核半径 $\propto A^{1/3}$

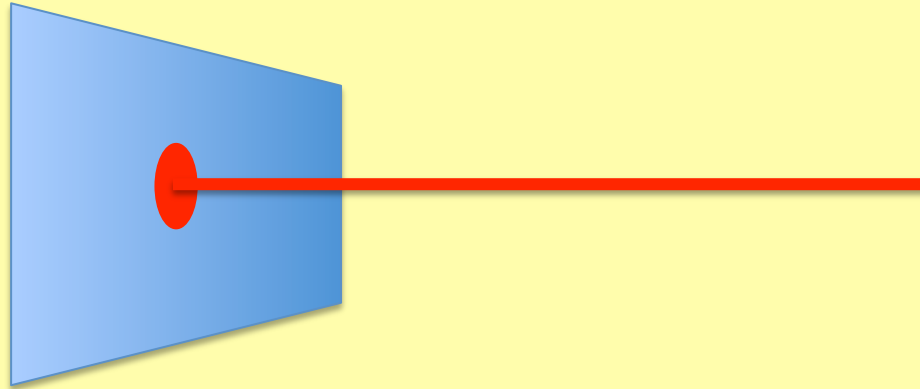
[T.Morita, KH 1103.5688]

$$\begin{aligned} \sqrt{r_{\text{mean}}^2} &\equiv \sqrt{\frac{1}{A} \operatorname{tr} \langle X^I X^I \rangle} \\ &= \frac{3^{5/2} \pi^{2/3}}{2^{5/6} 5^{1/6}} \frac{A^{1/3}}{M_{\text{KK}} (N_c \lambda^2)^{1/3}} \end{aligned}$$

$$\Rightarrow \sqrt{r_{\text{mean}}^2} \sim 0.7 A^{1/3} \text{ [fm]}$$

現実的な原子核半径
の解析的導出

0 クォーク = 異次元方向に伸びた弦



1 チャーモニウムとブラックホール

2 重力衝突と「とがった」素粒子

3 原子核と隠れた対称性

異次元