



### Hadron Structure from Lattice QCD

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### Acknowledgements

- Thanks to everyone who sent data, plots, references
  - Special thanks to .....
    - Alex Chambers
    - Jack Dragos
    - Tim Harris
    - Huey-Wen Lin
    - Rudolf Rödl
    - Shoichi Sasaki

Sorry if I don't cover your work, there are way too many topics to cover in 40mins



- Motivation Experimental challenges & theoretical issues
- Nucleon charges (Axial, tensor, scalar)
- Momentum fraction
- Spin content, including  $\Delta s \ \Delta G$
- EM form factors (including strangeness, light nuclei, charge symmetry)
- Pion,  $ho 
  ightarrow \pi\gamma$  form factors
- Transverse Momentum Distributions & "quasi"-PDFs
- Summary & outlook

### Motivation

- Major goal of nuclear physics community
  - understand the structure and behaviour of strongly interacting matter in terms of its basic constituents - quarks and gluons
- An important step towards this goal is the characterisation of the internal structure of the nucleon
  - Driving force behind several experimental programs, e.g. JLab 12 GeV upgrade



### Deep inelastic scattering (DIS) experiments



nucleon terms of its quark and gluon (parton) constituents

 The same cannot be said of our understanding of the nucleon spin, mass, magnetic moment, ....

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No, only a "spin puzzle"





(EIC white paper [1212.1701])

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(EIC white paper [1212.1701])

But how is the other 70% of the proton spin distributed between these components?

How is the spin of the proton distributed between its constituents?

Jaffe & Manohar (1990):

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + \mathcal{L}_q + \Delta G + \mathcal{L}_g$$

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + L_q + J_g$$

$$\Delta \Sigma = \Delta u + \Delta d + \Delta s = 0.33(3)(5)$$

Why so small?

Due to large negative  $\Delta s$  ?

Much effort to determine  $\Delta s$  experimentally  $\int_0^1 g_1^p(x) dx = \frac{1}{36}(4a_0 + 3a_3 + a_8)$ 

e.g. COMPASS, HERMES x > 0.004  $x \ge 0.02$ 

Also  $g_A$  and semileptonic hyperon decays assuming SU(3) symmetry





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Why so small?

Large  $\Delta G$  ?

#### **Recent global analysis of RHIC (PHENIX & STAR) and COMPASS data**

0.0

[de Florian et al., PRL113 012001((2014)]

$$\Delta G \approx \int_{0.005}^{0.2} \Delta g(x) \, dx = 0.1^{+0.06}_{-0.07}$$

(negligible contribution x > 0.2, poorly constrained 0.001 < x < 0.05)





### **Electromagnetic Form Factors**

$$\langle p', s' | J^{\mu}(\vec{q}) | p, s \rangle = \bar{u}(p', s') \left[ \gamma^{\mu} F_1(q^2) + i\sigma^{\mu\nu} \frac{q_{\nu}}{2m} F_2(q^2) \right] u(p, s)$$

$$G_E(Q^2) = F_1(Q^2) - \tau F_2(Q^2)$$
$$G_M(Q^2) = F_1(Q^2) + F_2(Q^2)$$

### Elastic Scattering - Polarisation Transfer

Polarisation transfer experiments at JLab revealed a surprising behaviour for  $G_E/G_M$ 

Precise results now available up to 8-9 GeV<sup>2</sup>

What is the origin of the linear fall-off?

Does  $G^p_E$  change sign?



## Size of the Proton

Charge radius of the nucleon

$$G_E^p(Q^2) = 1 - \frac{Q^2}{6} \langle r^2 \rangle_E + p + \dots$$

+  $\sim 7\sigma$  discrepancy between µH and H / e-p scattering





- An understanding of nucleon spin is a major goal of many current and future experimental programs
- 12 GeV JLab upgrade

- Origins of quark confinement
- Spin and flavour structure of the proton and neutron (PDF's, GPD's, TMD's...)
- Quark structure of nuclei
- Probe potential new physics through high precision tests of the Standard Model
- Flavour separation of EM form factors

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Proton spin (GPDs, Gluon)

· COMPASS (II)

- TMDs
- Pi and K polarisabilities

- An understanding of nucleon spin is a major goal of many current and future experimental programs
- 12 GeV JLab upgrade
- COMPASS (II)
- RHIC

• (PHENIX, STAR)

- spin and orbital angular momentum of quarks and gluons
- transverse spin structure of the proton

- An understanding of nucleon spin is a major goal of many current and future experimental programs
- 12 GeV JLab upgrade
- COMPASS (II)
- RHIC
- MAMI (Mainz):

- EM form factors, polarisabilities
- Structure of nuclei
- Parity violating strange EM, axial form factor

- An understanding of nucleon spin is a major goal of many current and future experimental programs
- 12 GeV JLab upgrade
- COMPASS (II)
- RHIC
- MAMI (Mainz)
- Electron-Ion Collider (JLab or BNL)

- Dramatic improvement in understanding role of sea quarks and glue
- "Missing spin" provided by gluons
- High-energy probes of partons' transverse momenta - contribution from orbital motion

### Lattice Hadron Structure

### Lattice 3pt function

Most common method for determining matrix elements relevant for hadron structure calculations - 3pt function



 $G(t,\tau,\vec{p},\vec{p}') = \sum_{s,s'} e^{-E_{\vec{p}'}(t-\tau)} e^{-E_{\vec{p}'}} \Gamma_{\beta\alpha} \langle \Omega | \chi_{\alpha}(0) | N(p',s') \rangle \langle N(p',s') | \mathcal{O}(\vec{q}) | N(p,s) \rangle \langle Np,s) | \overline{\chi}_{\beta}(0) | \Omega \rangle$ 

#### For large times $1 \ll \tau$ $1 \ll t - \tau$





#### Relatively simple to compute on the lattice (p=q=0, isovector)

Good benchmark for hadron structure (understanding systematic errors)

### Determination of $g_A$ on the Lattice

#### Summary plot: RQCD PRD91 (2015) 054501



## Determination of g<sub>A</sub> on the Lattice

Thorough investigations of

Lattice spacing

**Quark mass** 

**Finite volume** 

**Contamination from excited states** 

**Improved axial current** 

[Dragos, Tue 17:30] [Gupta, Thu 10:40] [Harris, Sat 10am] [Ohta, Wed 18:10] [von Hippel, Sat 10:40] [Yamanaka, Wed 17:10] [Yamazaki, Wed 14:20 (Had Spec & Int)]

[Schiller, Sat 10:20]

g<sub>A</sub> appears to be very sensitive to Lattice systematics

Lots of effort in reducing systematic errors

flow on for other quantities

## Determination of g<sub>A</sub> on the Lattice



## Determination of $g_A$ on the Lattice

#### **Renormalisation free ratio**



QCDSF: PLB 732 (2014) 41

 $g_A/f_{\pi}$  Ratio extrapolation agrees with exp

**FV** effects cancel in HBChPT

How is the spin of the proton distributed between its constituents?

**Recall Ji's sum rule:** 

$$\frac{1}{2} = \sum_{a} J_q(\mu^2) + J_g(\mu^2)$$

Express in terms of moments of Generalised Parton Distributions

$$J_{q/g} = \frac{1}{2} \left[ A_{20}^{q/g}(q^2 = 0) + B_{20}^{q/g}(q^2 = 0) \right]$$



which are obtained from the matrix elements of the energy momentum tensor

$$\langle P'|T^{\mu\nu}|P\rangle = \overline{U}(P') \left\{ \gamma^{\mu} \overline{P}^{\nu} A_{20}(q^2) + \frac{i\sigma^{\mu\rho} q_{\rho} \overline{P}^{\nu}}{2m_N} B_{20}(q^2) + \frac{q^{\mu} q^{\nu}}{m_N} C_{20}(q^2) \right\} U(P)$$

Determine the quark orbital angular momentum

$$L_q = \frac{1}{2} \left[ A_{20}^q (q^2 = 0) + B_{20}^q (q^2 = 0) - \Delta q \right]$$

But  $\langle x \rangle^{q/g} = A_{20}^{q/g} (q^2 = 0)$ 

# Quark Momentum Fraction $\langle x \rangle_q = \int_0^1 dx \, x (q(x) + \bar{q}(x))$

- First moment of the (isovector) nucleon parton distribution function
- Notorious for producing lattice results  $\approx$  2x too large for isovector nucleon
- Known to be sensitive to excited state contamination
- Results near physical mass inconclusive



### Spin of the Proton Ji's sum rule

**Recall Ji's sum rule requires moments of Generalised Parton Distributions** 



### Spin of the Proton Ji's sum rule

**Recall Ji's sum rule requires moments of Generalised Parton Distributions** 

For reliable results for  $L_q$ 

> need to control systematic errors for  $\, g_A \,$  and  $\langle x 
angle^{u-d}$ 

disconnected contribitions for:

Spin of the Proton  $\Delta s = \int_0^1 dx \left[ \Delta s(x) + \Delta \bar{s}(x) \right]$ 

 $\Delta s$  is a purely quark-line disconnected contribution

#### A challenge on the lattice

Standard procedure: Use stochastic (random noise) sources

200000

e.g.[QCDSF/RQCD PRL108 (2012) 222001] $m_{\pi} = 285 \,\mathrm{MeV}$  $\Delta s = -0.020(10)(4)$  $\overline{\mathrm{MS}}$  $\mu = \sqrt{7.4} \,\mathrm{GeV}$ 

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An alternative: Feynman-Hellmann method [CSSM/QCDSF PRD90 (2014) 014510]

$$\frac{\partial E_H(\lambda)}{\partial \lambda} = \frac{1}{2E_H(\lambda)} \left\langle H \left| \frac{\partial S(\lambda)}{\partial \lambda} \right| H \right\rangle$$



2000

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Also allows for full nonperturbative determination of singlet renormalisation constants

$$g_A = Z_A^{\rm NS} g_A^{\rm latt}$$
$$\Delta \Sigma = Z_A^{\rm S} \Delta \Sigma^{\rm latt}$$

[CSSM/QCDSF, PLB740 (2015) 30]





### Disconnected Spin Contributions $\Delta s$



favours a small and negative  $\Delta s$ 

### **Disconnected Contributions**



] q

### Glue contribution to proton spin Yang Thu 9:50am

Ji et al. [PRL 111 (2013) 112002] proposed a glue spin density operator



Tensor Charge, 
$$g_T = \int dx \left[ \delta u(x) - \delta d(x) \right]$$

**Recent interest driven by** 

Phenomenological determinations now available using SIDIS data (HERMES, COMPASS)

Implications for new physics

Precision neutron  $\beta$  decay studies sensitive to possible BSM scalar and tensor interactions [PNDME PRD85 (2012) 054512]

> Require neutron matrix elements of appropriate low-energy operators

i.e. 
$$g_T^{u-d}$$
,  $g_S^{u-d}$  (to ~10%)

Quantify the contribution of the quark EDM to the nEDM and set bounds on new sources of CP violation [PNDME 1506.04196, 1506.06411]

$$d_n = d_u g_T^u + d_d g_T^d + d_s g_T^s$$

(including disconnected)

See plenary talk by T. Izubuchi (Tue 11:45)

Tensor Charge, 
$$g_T = \int dx \left[ \delta u(x) - \delta d(x) \right]$$

[Gupta, Thu 10:40]



**Comprehensive study of systematics by PNDME** 

Introduced 'FLAG-like' colour-coding system for each systematic error

10% uncertainty target achieved

### Scalar Charge, $g_S$

$$\langle N | \bar{q}q | N \rangle$$

Commonly used for 
$$\sigma$$
 terms  $\sigma_l^H = m_l \langle H | (\bar{u}u + \bar{d}d) | H \rangle$   $\sigma_s^H = m_s \langle H | \bar{s}s | H \rangle$   
[R.Young, Latt'12 review]

Recent interest: new physics contributions to  $\beta$  decay:  $g_S^{u-d}$  to 10%

[PNDME PRD85 (2012) 054512]

Severe excited state contamination - take care!



#### CSSM/QCDSF (Dragos, Tue 17:30)

### Scalar Charge, $g_S$

 $\langle N | \bar{q}q | N \rangle$ 

Large scatter in lattice determinations

#### Only determined to ~30%



### **Electromagnetic Form Factors**

Several new simulations with near-physical quark masses

LHPC [PRD90 (2014) 074507]	$m_{\pi} \approx 149 \mathrm{MeV}$	
Mainz [arXiv:1504.04628]	$m_{\pi} \approx 193 \mathrm{MeV}$	Shintani, Wed 17:30
PACS-CS	$m_{\pi} \approx 145 \mathrm{MeV}$	Yamazaki, Wed 14:20 (Had Spec & Int)
PNDME	$m_{\pi} \approx 130 \mathrm{MeV}$	Gupta, Thu 10:40

**Progress in accounting for excited state contamination** 



Shintani, Wed 17:30

Mainz [arXiv:1504.04628]

#### **Evidence of excited state** contamination in **EM FFs**

Effect becomes more problematic at light quark masses

#### Effect increases $G_E$ at each $Q^2$

Improved approach to physical point



### **Electromagnetic Form Factors**

Comparison of LHPC ( $m_{\pi} \approx 149 \,\mathrm{MeV}$ )

PRD90 (2014) 074507

PNDME (  $m_\pi pprox 130\,{
m MeV}$  )

Gupta, Thu 10:40



Good agreement with parameterisation of experimental data

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Lattice form factors now available with physical masses

(Direct or small chiral extrapolation)

Determine accurate radii @  $m_{\pi} \approx 140 \,\mathrm{MeV}$  :

Use finite-volume and excited-state corrected results

Need small  $Q^2 < 4m_\pi^2 \sim 0.08 \,\mathrm{GeV}^2$  (2 $\pi$  threshold)

dipole fit to large Q<sup>2</sup> data introduces unwanted model dependence

"z-fit" [Hill & Paz, PRD82 (2010) 113005 ] more reliable



Lattice form factors difficult to determine at large  $Q^2 > 2-3$  GeV<sup>2</sup> (even at large m<sub> $\pi$ </sub>)

**Standard method:** 

fix 
$$\vec{p}' = 0$$
 vary  $\vec{q} = -\vec{p}$ 



Breit frame  $\vec{p}' = -\vec{p}$  to maximise Q<sup>2</sup> at fixed  $\vec{p}$ ?

**New methods?** 









### Neutron Form Factors

**Experimental determination of neutron EM form factors difficult since** 



### Magnetic Structure of Light Nuclei [NPLQCD PRL 113 (2014) 252001 & 1506.05518]

Magnetic moments and polarizabilities of nucleons and light nuclei with  $A \le 4$ 



shell-model configuration captures their dominant structures

#### **Deviations similar to experiment**

Implications for experiments using  $^{2}\mathrm{H},~^{3}\mathrm{H},~^{3}\mathrm{He}\,$  for neutron properties

### Strangeness Form Factors

- Understanding hidden flavour A fundamental challenge of hadronic physics
- Contributions arise entirely through interactions with QCD vacuum
- Extensive experimental searches
  - JLAB (G0, HAPPEX), MIT-Bates (SAMPLE), Mainz (A4)



## Strangeness Form Factors

### **Qweak experiment (JLab)**

- ullet Implications for experimental determination of the proton's weak charge  $Q^p_W$
- $Q_W^p$  : neutral current analog to the proton's electric charge
- Precisely predicted to be small in the SM

Constrain new parity-violating (PV) physics between electrons and light quarks



### Strangeness Form Factor



## Strangeness Form Factor

#### **Comparison with experimental contraints**



Presents a challenge for the next generation of experiments

Reached a precision where violations of charge symmetry will become important

### Charge Symmetry

*u* quarks in the proton  $\equiv d$  quarks in the neutron Experimental determinations of  $G^s_{E,M}(Q^2)$  rely on charge symmetry

(As does the inclusion of nuclear data for Qweak)

EM and weak interactions give access to different combinations of  $G^{p,(u/d/s)}$ 

$$G^{p,\gamma} = \frac{2}{3}G^{p,u} - \frac{1}{3}(G^{p,d} + G^{p,s})$$
$$G^{p,Z} = \left(1 - \frac{8}{3}\sin^2\theta_W\right)G^{p,u} - \left(1 - \frac{4}{3}\sin^2\theta_W\right)(G^{p,d} + G^{p,s})$$

Assume charge symmetry  $(G^{p,u} = G^{n,d}, G^{p,d} = G^{n,u}, G^{p,s} = G^{n,s})$ 

$$= \bigvee \left( G_{E/M}^{p,s} = \left( 1 - 4\sin^2 \theta_W \right) G_{E/M}^{p,\gamma} - G_{E/M}^{n,\gamma} - G_{E/M}^{p,Z} \right)$$

### Charge Symmetry Violation [CSSM/QCDSF PRD91 (2015) 113006]

Determine the degree to which charge symmetry is violated in EM form factors by

**Combining chiral perturbation theory fits to isospin-averaged hyperon FFs** 

Input  $m_u/m_d$  from experiment (or FLAG)

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Charge symmetry satisfied to better than 0.2% (QCD)

Violations due to **QED**? Challenge for future QCD+QED simulations

### Pion Form Factor

- Asymptotic normalisation known from  $\pi \to \mu + \nu {\rm decay}$ 

$$F_{\pi}(Q^2 \to \infty) = \frac{16\pi\alpha_s(Q^2)f_{\pi}^2}{Q^2}$$

- Allows to study the transition from the soft to hard regimes
- Low Q<sup>2</sup>: measured directly by scattering high energy pions from atomic electrons



## Pion Form Factor

[Bijnens-Ecker, 2014]

[T. Kaneko, Tue 16:30 (WD & ME)]

- JLQCD: overlap quarks,  $290 \le m_\pi \le 540 \,\mathrm{MeV}\,$  (all-2-all propagators)
- NNLO SU(3) ChPT + N3LO analytic [Bijnens-Ecker, 2014 ] combined fit to

 $F_{\pi^+}(q^2), \ F_{K^+}(q^2), \ F_{K^0}(q^2)$   $\langle r_{\pi^+}^2 \rangle = 0.458(15)(38) \text{ fm}^2, \ \langle r_{K^+}^2 \rangle = 0.380(12)(32) \text{ fm}^2, \ \langle r_{K^0}^2 \rangle = -0.055(11)(45) \text{ fm}^2$   $\textbf{PDG'14:} \ \langle r_{\pi^+}^2 \rangle = 0.452(11) \text{ fm}^2, \ \langle r_{K^+}^2 \rangle = 0.314(35) \text{ fm}^2, \ \langle r_{K^0}^2 \rangle = -0.077(10) \text{ fm}^2$ 



#### Access large Q<sup>2</sup> via Feynman-Hellmann





### **Quark Distribution Functions**

Recent proposal by Ji: $\tilde{q}_{lat}(x,\Lambda,P_z) = \int \frac{dz}{4\pi} e^{-izk} h(z,\Lambda,P_z),$ "Quasi"-PDFs $h(z,\Lambda,P_z) = \langle \vec{P} | \bar{\psi}(z) \gamma_z \left( \prod_n U_z(n\hat{z}) \right) \psi(0) | \vec{P} \rangle,$ [PRL 110 (2013) 262002] $\psi(0)$  $\Psi_z$  $\psi(0)$  $\psi(0)$  $\bar{\psi}(z)$  $\Psi_z$  $\psi(0)$ 

**Connection to continuum through perturbative matching** 

#### 2 lattice calculations [H-W.Lin et al. PRD91 (2015) 054510] 1.5 $P_z = \frac{4\pi}{L}$ 1.0 $\sigma$ $6\pi$ 0.5 $P_z =$ 0 -0.5-0.50 0.51.0-1.01.5х





spin-orbit correlations in hadrons

Accessed via Semi-Inclusive Deep Inelastic Scattering or Drell-Yan [JLab, HERMES, COMPASS, RHIC]

**Relevant matrix elements:** 

$$\langle P | \bar{q}(0) \Gamma \mathcal{U}[0, \eta v, \eta v + b, b] q(b) | P \rangle$$
  
Gauge link "along the light cone"

#### [LHPC: PRD85 (2012) 094510]



### Transverse Momentum Distributions [Engelhardt, Fri 15:40 (WD & ME)]

Nucleon update with  $m_{\pi} \approx 170 \,\mathrm{MeV}$  (2+1 DWF)



#### Dependence of SIDIS/DY limit on $|b_T|$

#### Recent investigated pion TMD [LHPC 1506.07826]

understand approach to "light cone"

### Summary & Outlook

**Excellent progress in understanding and controlling systematic errors** 

> Precise results at the physical point now achievable for

isovector charges

electromagnetic form factors

With additional progress in computing disconnected contributions

Improved understanding of long term questions, e.g.

decomposition of proton spin

role of hidden flavour

Right time to tackle more "exotic" quantities relevant for upcoming experiments

quasi-PDFS, TMDs, nuclear effects, ...

### International Nuclear Physics Conference (INPC2016)



2004: Göteborg, Sweden 2007: Tokyo, Japan 2010: Vancouver, Canada 2013: Florence, Italy **2016: Adelaide, Australia** 

with dedicated lattice sessions on nuclear physics, hadron structure, finite density, ...

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### Adelaide, Australia

### September 11-16, 2016

with dedicated lattice sessions on nuclear physics, hadron structure, finite density, ...