

Flavor Filtered Fermions

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Naive fermions

- Naive Dirac operator (T_μ covariant shift/parallel transporter)

$$D_N = \frac{1}{2} \gamma_\mu (T_\mu - T_\mu^\dagger)$$

- Naive propagator

$$\frac{1}{D_N + m}$$

- Has doubler poles at momenta (in the free, massless case)

$$p_\mu \in \{0, \pi\}$$

Basic idea

- Naive propagator

$$\frac{1}{D_N + m}$$

- Filter out doublers

$$F \frac{1}{D_N + m} F^\dagger$$

- F: “flavor filter”, simple example (free case, momentum space)

$$F = \prod_{\mu} (1 + \cos p_{\mu}) / 2$$
$$F \rightarrow \begin{cases} 1, p = \{0, 0, 0, 0\} \\ 0, \text{any } p_{\mu} = \pi \end{cases}$$

Flavor Filtered Fermions

- Zeros of $F \rightarrow$ poles in $D \rightarrow$ nonlocal D in real space
- Use Ginsparg-Wilson relation to provide way out

$$D^{-1}\gamma_5 + \gamma_5 D^{-1} = 2\rho\gamma_5$$

- (massless) Flavor Filtered Fermion

$$D_F^{-1} = \rho + F \frac{1}{D_N} F^\dagger$$

$$D_F = \frac{1}{\rho} \left[1 - F \frac{1}{\rho D_N + F^\dagger F} F^\dagger \right]$$

Mass term

- Simplest version

$$D_F^{-1} = \rho + F \frac{1}{D_N + m} F^\dagger$$

- Can spin diagonalize inverse \rightarrow 4 staggered solves
- Could use other operators for mass term $m \rightarrow mG$
 - Give heavier mass to doublers
 - Make operator in denominator more expensive, but improve condition number

Action

- Rewrite determinant

$$\begin{aligned} |D_F| &= \left| \frac{1}{\rho} \left[1 - F \frac{1}{\rho(D_N + m) + F^\dagger F} F^\dagger \right] \right| \\ &= \frac{|D_N + m|}{|\rho(D_N + m) + F^\dagger F|} \end{aligned}$$

- Can view as partial quenching of 15 flavors
 - $F^\dagger F$ mass term \rightarrow 15 light, 1 heavy state

Tuning filter

- Free case: simplest version

$$\begin{aligned} F &= \prod_{\mu} (1 + \cos p_{\mu}) / 2 \\ &= \prod_{\mu} (2 + T_{\mu} + T_{\mu}^{\dagger}) / 4 \end{aligned}$$

- Interacting case: link renormalization requires tuning filter
Symmetrize product over permutations

$$F = \text{Sym} \prod_{\mu} [(1 - \alpha) + \alpha(T_{\mu} + T_{\mu}^{\dagger})/2]$$

- $\alpha=0$: naive fermions
 $\alpha=0.5$: untuned filter
 $\alpha>0.5$: tuned filter

Tests

- Testing on quenched lattices
- Plaquette action
- Volume: $32^3 \times 64$
- 3 couplings
- Scale from Wilson flow

beta	w0 *	"a" (fm, Nf=2+1) *
6.0	1.766(2)	0.099(1)
6.2	2.442(4)	0.072(1)
6.4	3.273(13)	0.054(1)

* BMW [S. Borsányi, et al., JHEP09(2012)010]



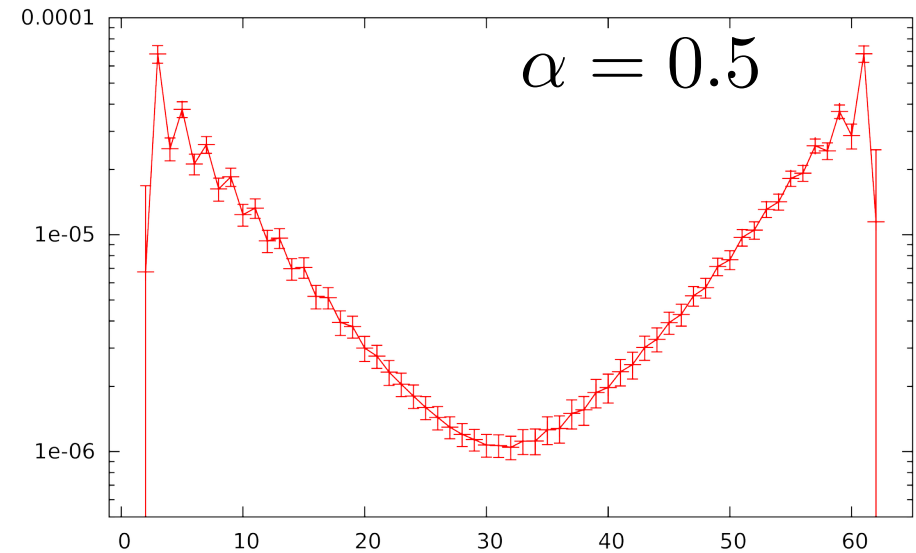
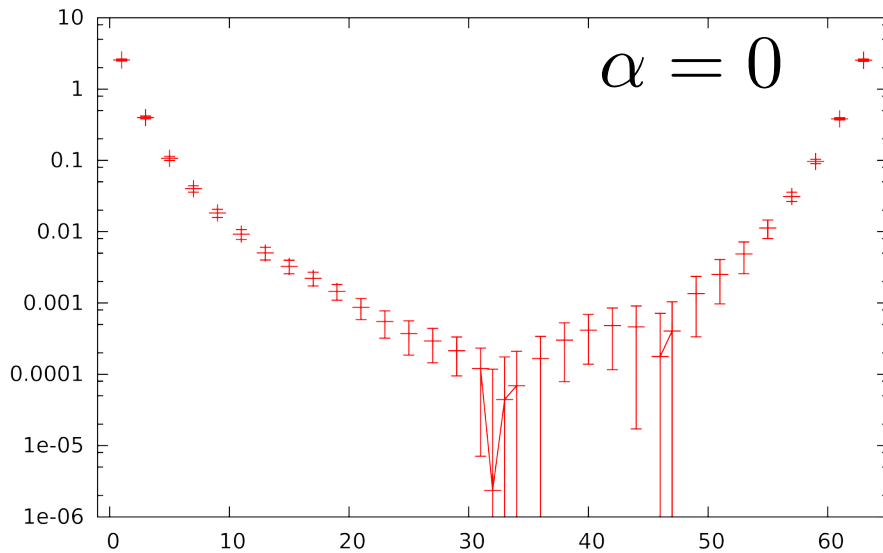
Meson spectrum

- Local mesons:
 - Treat like Wilson/Overlap fermions: insert appropriate gamma
 - Watch oscillating “time doubler” disappear as alpha tuned from naive fermions to filtered case
- Non-local mesons:
 - Gives complete information on suppression of doublers
 - Useful to think in terms of conversion from naive to staggered basis

$$\begin{aligned}\gamma_n &\rightarrow \gamma_n (\gamma_n \otimes \xi_n) \\ T_n &\rightarrow \gamma_n (1 \otimes \xi_n) \\ (-1)^{x+x_n} &\rightarrow 1 (\gamma_n \otimes \xi_n)\end{aligned}$$

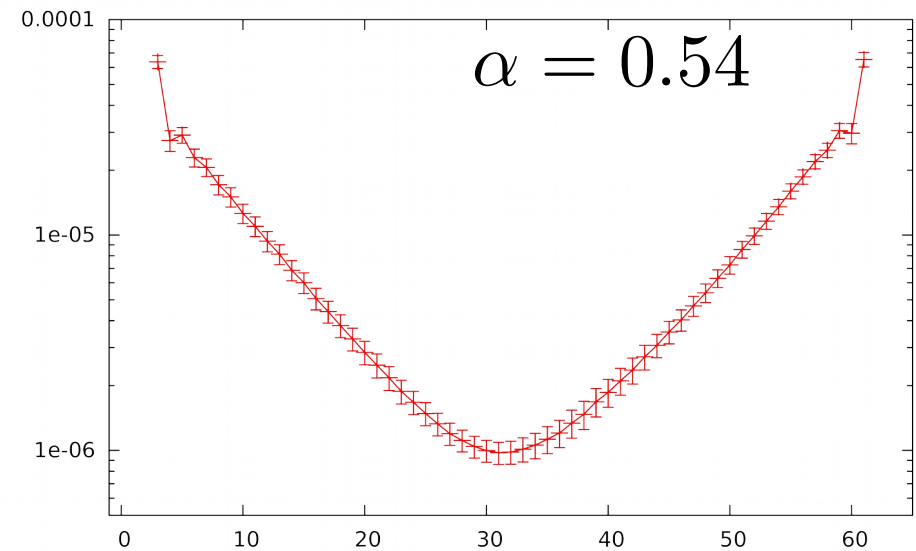
staggered flavor \nearrow \nwarrow staggered spin, taste

Local meson correlator



▲
Oscillating state goes
negative, not shown

- $\beta=6.4, m=0.0069$
- Local meson operator γ_{45}



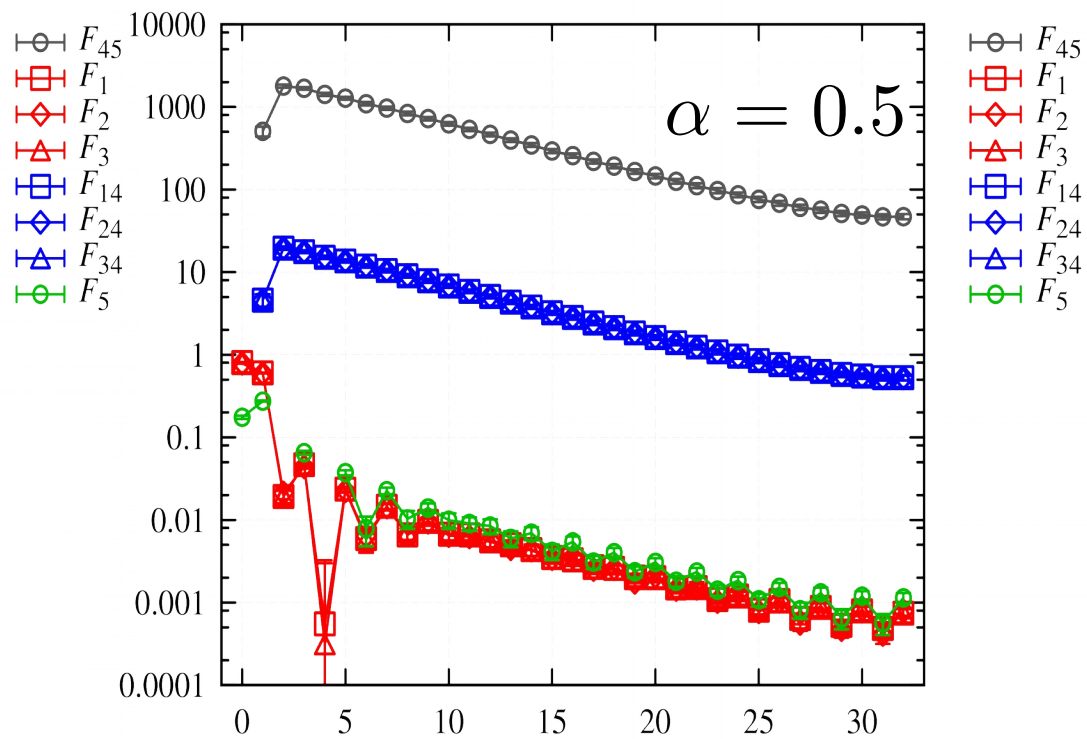
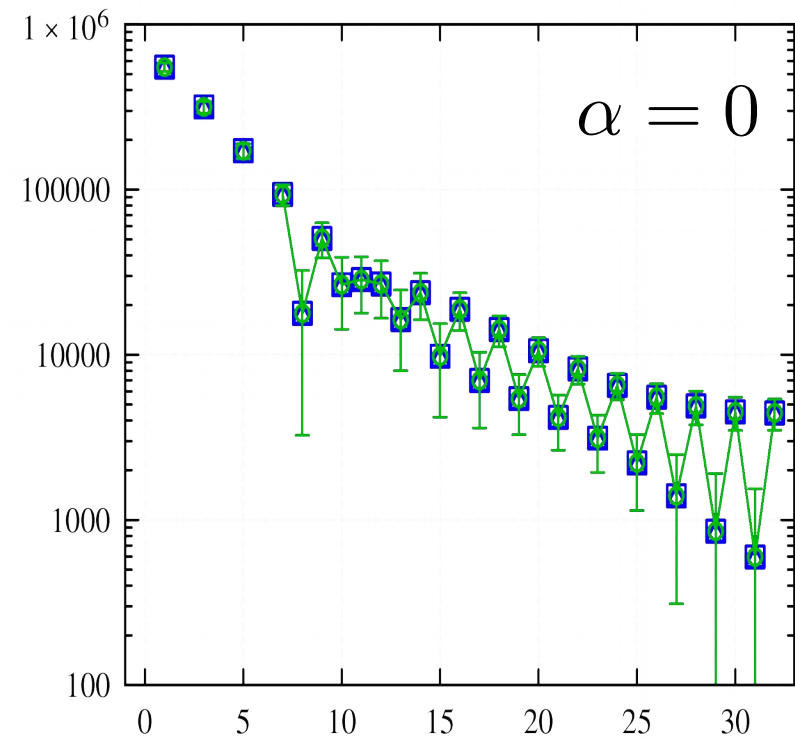
Meson spectrum

- Local mesons:
 - Treat like Wilson/Overlap fermions: insert appropriate gamma
 - Watch oscillating “time doubler” disappear as alpha tuned from naive fermions to optimally filtered case
- Non-local mesons:
 - Gives complete information on suppression of doublers
 - Useful to think in terms of conversion from naive to staggered basis

$$\begin{aligned}\gamma_n &\rightarrow \gamma_n (\gamma_n \otimes \xi_n) \\ T_n &\rightarrow \gamma_n (1 \otimes \xi_n) \\ (-1)^{x+x_n} &\rightarrow 1 (\gamma_n \otimes \xi_n)\end{aligned}$$

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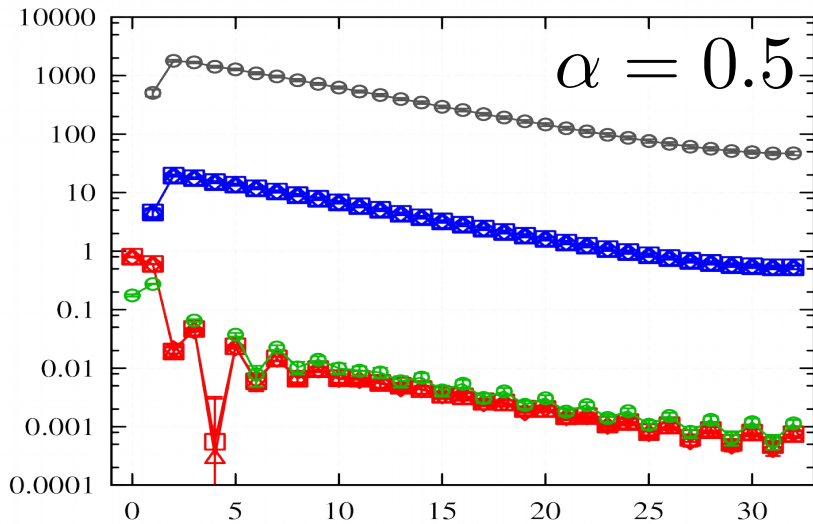
Non-local meson correlator



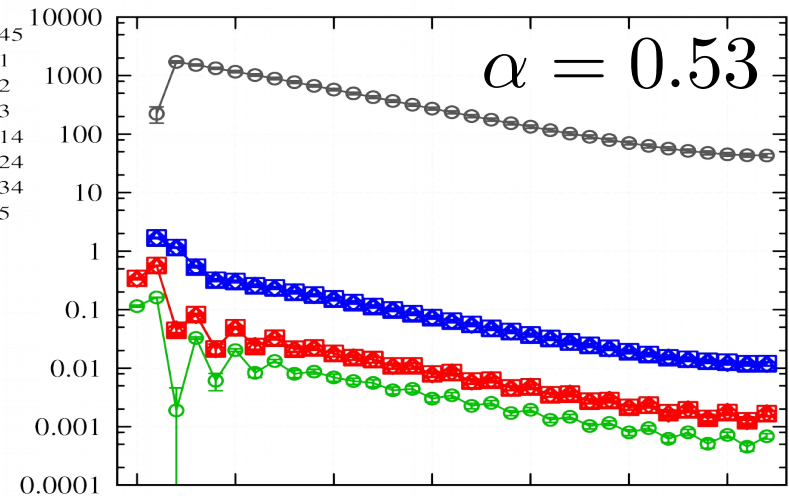
- $\beta=6.4, m=0.0069$

- Correlator of staggered states: $\gamma_F(\gamma_{45} \otimes \xi_{45})$

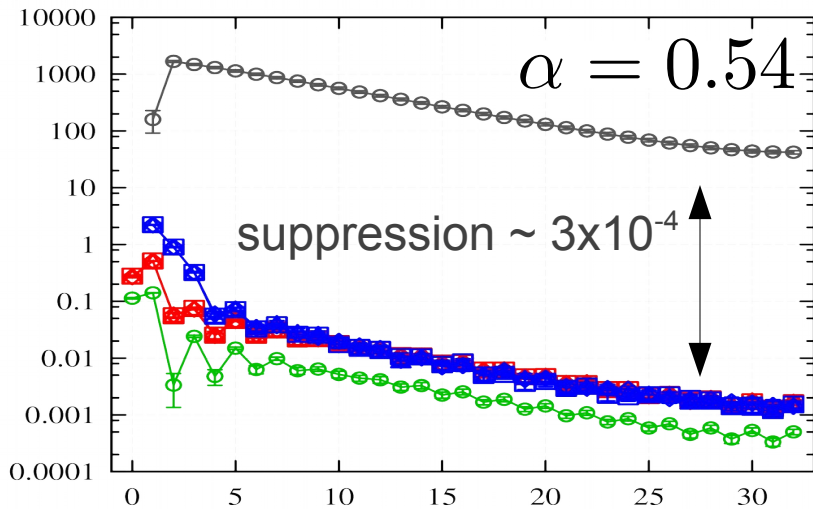
Non-local meson correlator



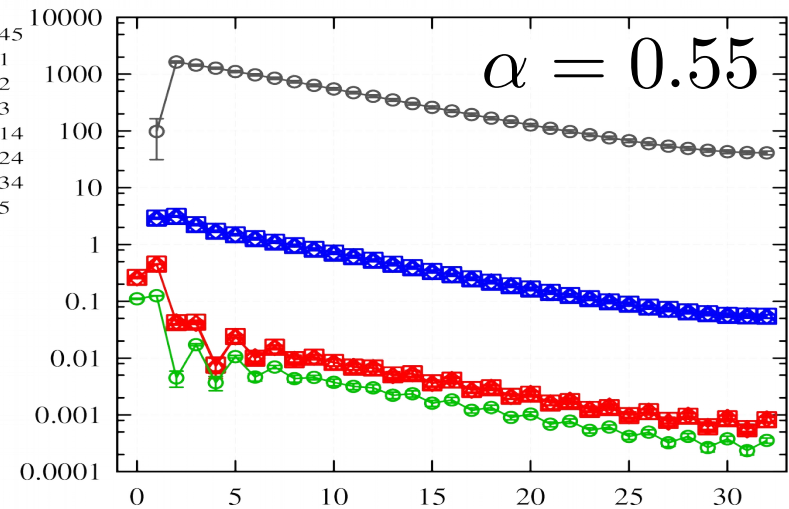
- F_{45}
- F_1
- F_2
- F_3
- F_{14}
- F_{24}
- F_{34}
- F_5



- F_{45}
- F_1
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- F_{34}
- F_5

Tuning filter

- Untuned filter ($\alpha = 0.5$)
 - Suppression $\sim 1e-2$
- Coarse tuning ($\alpha = 0.54$)
 - Suppression $\sim 3e-4$
- Finer tuning ($\alpha = 0.5352$)
 - Suppression $\sim 1e-4$

- Stout smearing (2 steps)
 - Reduces need for tuning
 - Untuned suppression ($\alpha = 0.5$) $\sim 3e-4$
 - Fine tuning suppression ($\alpha = 0.50615$) $\sim 1e-6$

- Still exploring optimal tuning



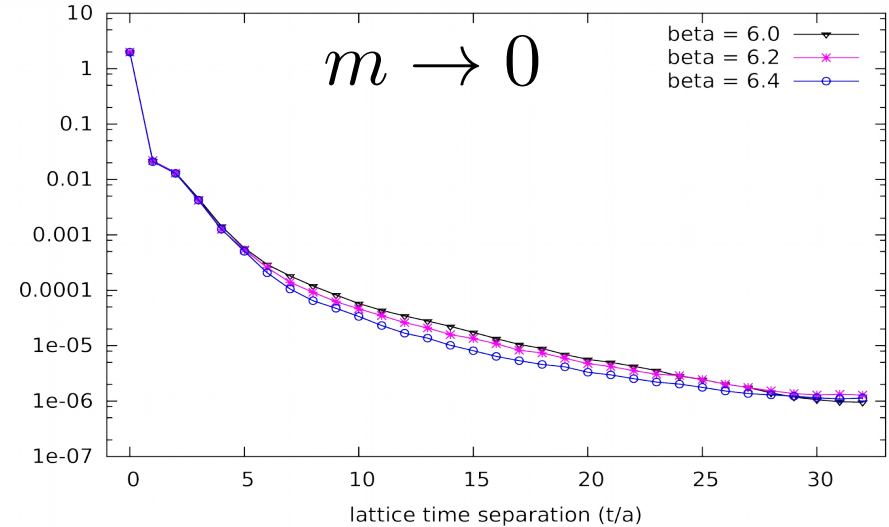
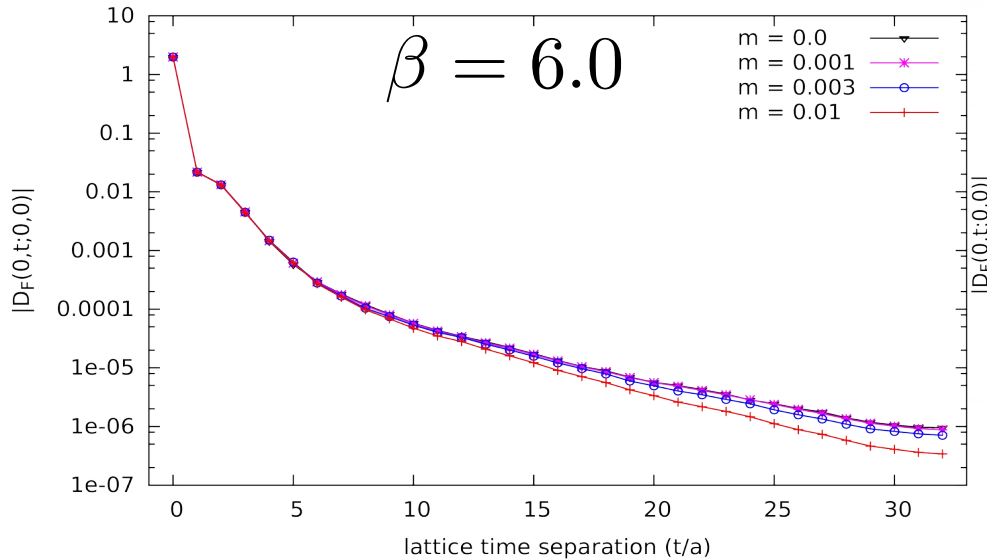
Fine tuning

- Simple filter with one parameter

$$F = \text{Sym} \prod_{\mu} \left[(1 - \alpha) + \alpha(T_{\mu} + T_{\mu}^{\dagger})/2 \right]$$

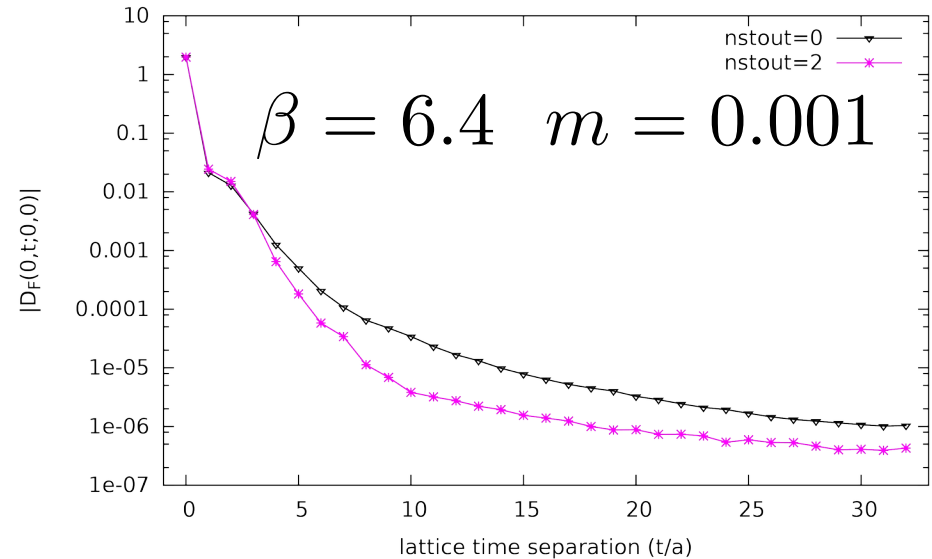
- Expands into 5 distinct operators, one for each number of hops (0,1,2,3,4)
- Could tune each coefficient separately (only ratio matters, so really 4 parameters to tune)
- Use non-local meson spectrum as guide
- Tuning 1 fairly easy, maybe 2 ok, 4 probably difficult

Dirac operator locality ($\alpha=0.5$)



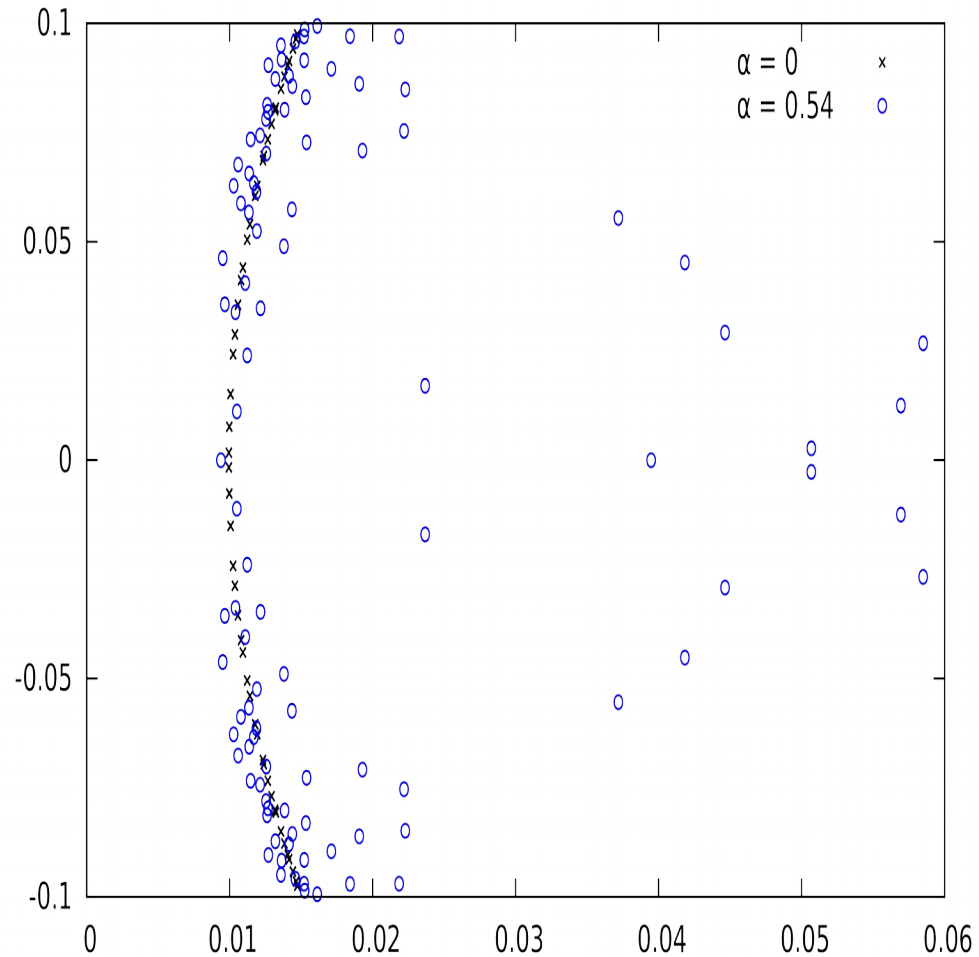
- Tail grows with m (effective mass gets lighter); reaches finite limit as $m \rightarrow 0$
- More local at intermediate distances as 'a' decreased
- Larger distances requires better filtering to reduce contribution

$$D_F = \frac{1}{\rho} \left[1 - F \frac{1}{\rho(D_N + m) + F^\dagger F} F^\dagger \right]$$



Dirac operator topology

- Eigenvalues of D_F on 8^4 lattice
- $\beta=5.8$, $m=0.01$
- $\alpha=0.54$ eigenvalues scaled to roughly match $\alpha=0$
- 15 eigenvalues move to larger real values, leaving 1 behind
- Effect goes away for smaller quark mass
- Need to check limit $a \rightarrow 0$ first, then $m \rightarrow 0$

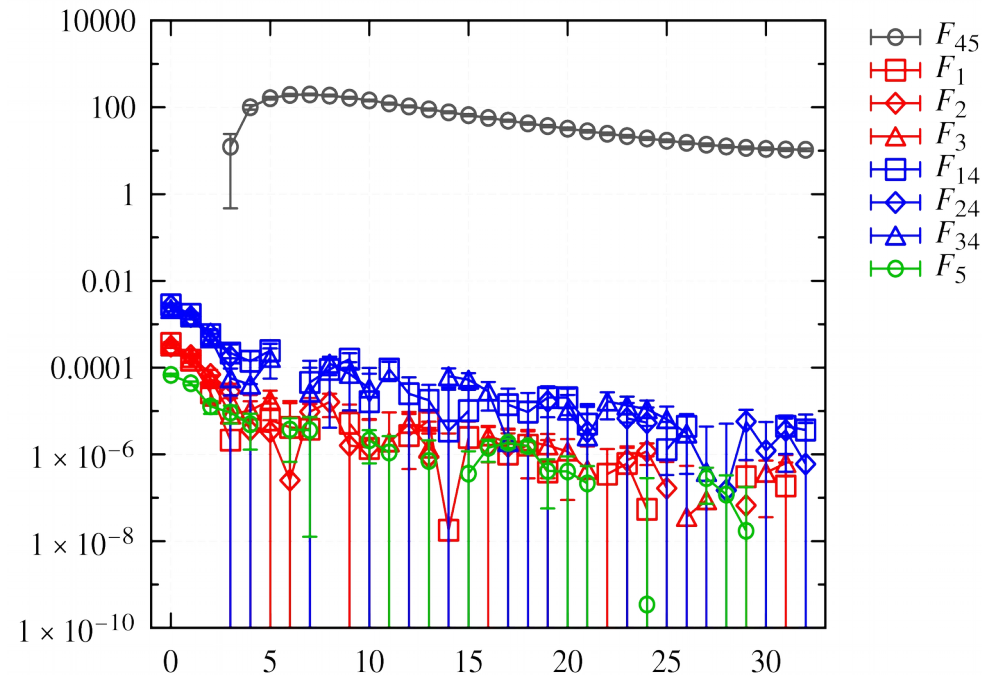
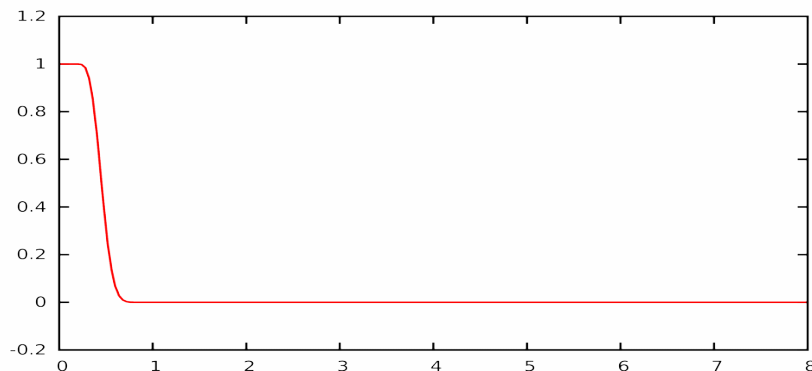


Alternative filters

- Use Wilson term to form filter

$$W = 4 - \frac{1}{2} \sum_{\mu} (T_{\mu} + T_{\mu}^{\dagger})$$

- Keep small eigenvalues of W , remove large
- Step function would introduce nonlocality
- Try smooth function
- Polynomial (order 64)



- $\beta=6.4, m=0.0069$

$$\gamma_F(\gamma_{45} \otimes \xi_{45})$$

- suppression $< 10^{-6}$
(too noisy to measure accurately)

Alternative filters

- Is there an optimal filter (at any cost)?
- What about something like:

$$\sim \frac{D_N}{D_{overlap}}$$

- Ratio preserves physical mode, cancels doublers
- Expensive (no better than just using overlap)
- Maybe a good enough approximation is cheaper?

Flavor (taste) Filtered Staggered?

- Start with staggered, filter out bad tastes

$$F \frac{1}{D_S + m} F^\dagger$$

- 2 flavor (taste)

$$F = (1 - \alpha) + \alpha(1 \otimes \xi_5)$$

- 1 flavor possible, but breaks rotational symmetry (see staggered-wilson literature)
 - Is there a variation that makes this feasible?

Summary

- Constructed new lattice Dirac operator
 - Satisfies Ginsparg-Wilson relation
 - Approximately 1 flavor (needs more study)
 - Depending on choice of filter, can probably satisfy basic locality requirement
 - Signs of emerging topological zero modes
- Cheaper than overlap/DW (in simplest form)
- Comparable expense to Wilson
 - Wilson has effective MG
- Staggered still cheaper
 - Can do $N_f=2$ flavor filtered staggered

Extra slides



