

The Nature of the Roberge-Weiss Transition in $N_f = 2$ QCD with Wilson Fermions on $N_t = 6$ Lattices

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in collaboration with
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Kobe, Japan

Outline

QCD Phase Diagram

Imaginary Chemical Potential and Roberge Weiss Symmetry

The QCD Phase Structure for Imaginary μ

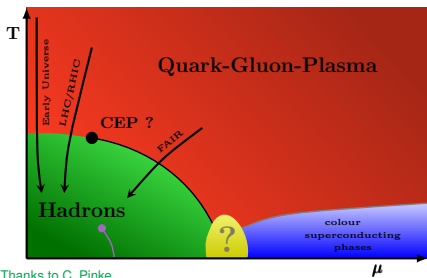
Previous and Ongoing Studies

Summary and Perspectives

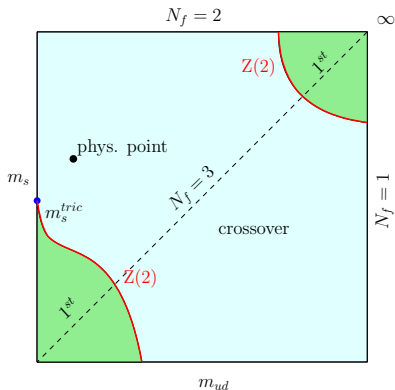
QCD Phase Diagram

QCD Phase Diagram

Columbia plot, $\mu = 0$



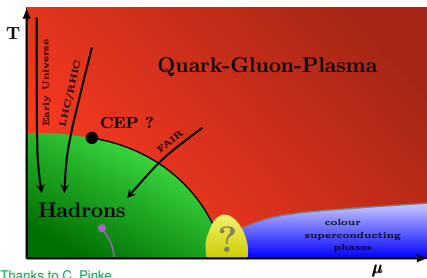
Thanks to C. Pinke



- Sign problem spoils Hybrid Monte Carlo simulations for $\mu > 0$.
- Simulate at $\mu = 0$ and apply e.g. extrapolation techniques.
- Choose purely imaginary chemical potential $\mu = i\mu_i$.

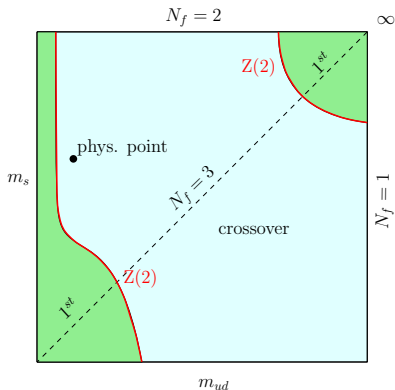
QCD Phase Diagram

QCD Phase Diagram



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Columbia plot, $\mu = 0$



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Imaginary Chemical Potential and Roberge Weiss Symmetry

- **No sign problem** → Hybrid Monte Carlo applicable.
- QCD partition function symmetries (Roberge-Weiss symmetry)

$$Z(\mu) = Z(-\mu), \quad Z\left(\frac{\mu}{T}\right) = Z\left(\frac{\mu}{T} + i\frac{2\pi n}{3}\right).$$

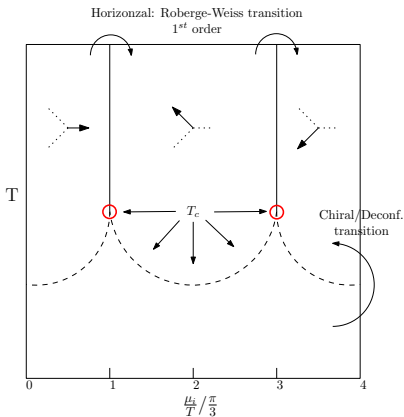
Roberge & Weiss, Nucl. Phys. B275 (1986)

- Adding $\mu_i/T = 2\pi n/3$, $n \in \mathbb{Z}$ is equivalent to $Z(3)$ transformation.
- Describes completely equivalent physics.
- Centre symmetry is a good symmetry again.
- Centre sectors separated by

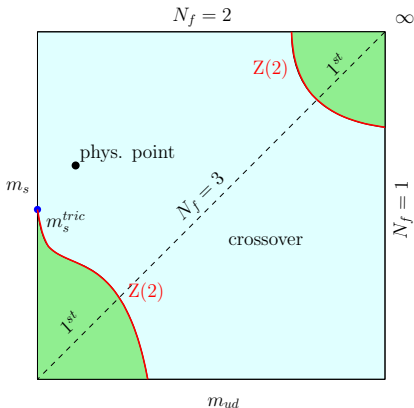
$$\mu_i^c = \frac{\pi T}{3}(2n + 1), \quad n \in \mathbb{Z}.$$

Imaginary Chemical Potential and Roberge Weiss Symmetry

Roberge-Weiss phase diagram



Columbia plot at $\mu = 0$

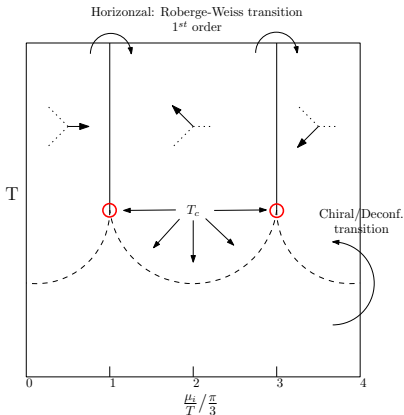


- Polyakov loop on the lattice: $L(n) = \frac{1}{V} \prod_{\tau_0}^{N_\tau} U_0(\tau, n)$
- Phase of the Polyakov is sensitive to centre sector transitions

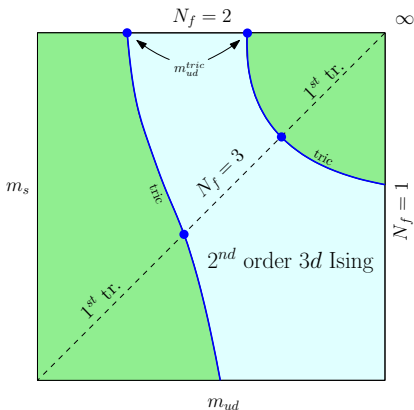
$$\text{Tr} L^g = e^{-i \frac{2\pi n}{3}} \text{Tr} L$$

Imaginary Chemical Potential and Roberge Weiss Symmetry

Roberge-Weiss phase diagram



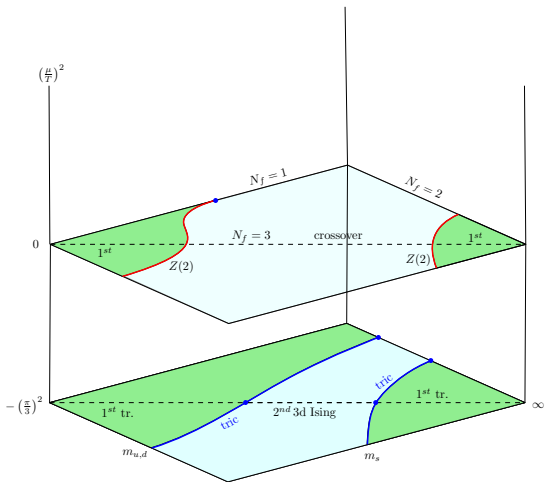
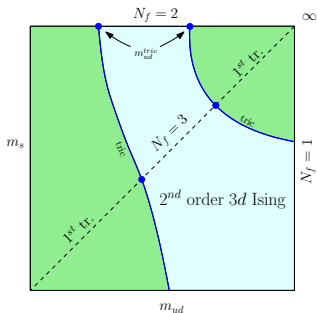
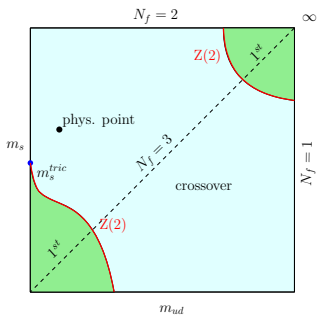
Columbia plot at $\mu_c = i\pi T/3$



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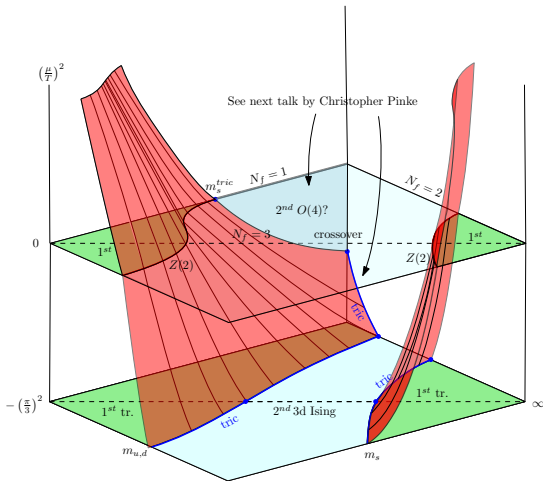
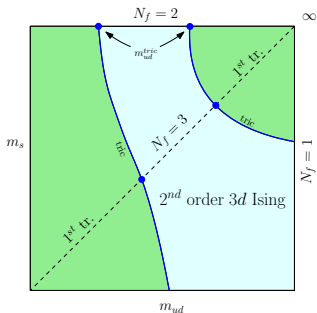
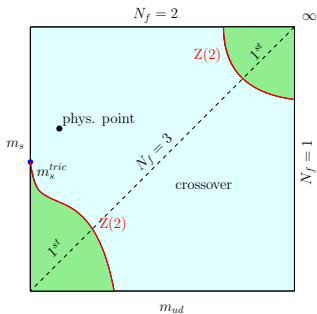
$$\text{Tr}L^g = e^{-i\frac{2\pi n}{3}} \text{Tr}L$$

The QCD Phase Structure for Imaginary μ



P. de Forcrand and O. Philipsen, JHEP (2007)
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 M. Fromm, J. Langelage, S. Lottini and O. Philipsen, JHEP (2012)
 H. Saito et al., arXiv:1309.2445 (hep-lat)

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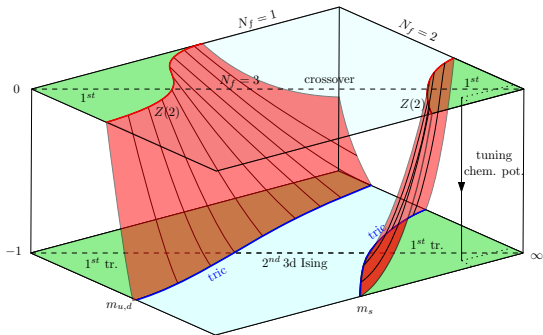
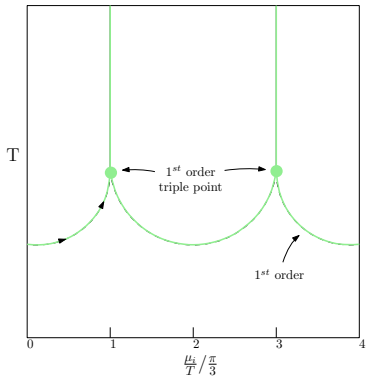
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The QCD Phase Structure for Imaginary μ

1st order



1st order triple

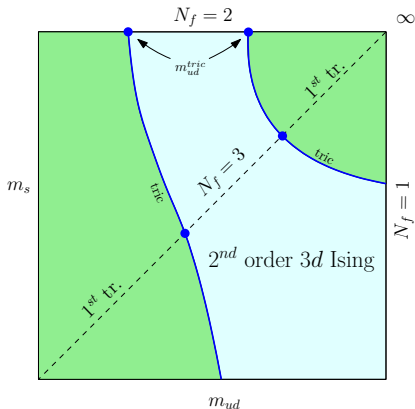


D'Elia, M. and Sanfilippo, F. Phys. Rev. D80 (2009)
 C. Pinke, O. Philipsen, Phys.Rev. D89 (2014)

Previous and Ongoing Studies

Studies done on $N_\tau = 4$:

- Staggered, $N_f = 2, 2 + 1, 3$
- Wilson, $N_f = 2$



- P. de Forcrand and O. Philipsen, JHEP (2007)
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M. D'Elia and F. Sanfilippo, Phys. Rev. D 80 (2009)
P. de Forcrand and O. Philipsen, Phys. Rev. Lett. 105 (2010)
G. Cossu, M. D'Elia and F. Sanfilippo, Phys. Rev. D 83(2011)
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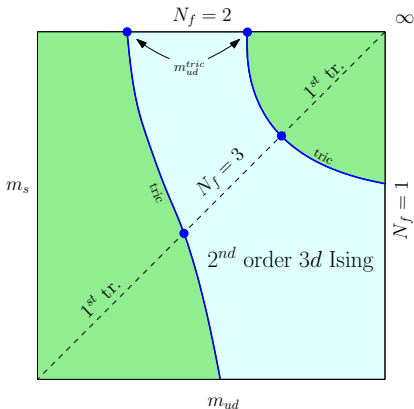
- Staggered, $N_f = 2, 2 + 1, 3$
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Ongoing for $N_\tau = 6$:

- Wilson, $N_f = 2$
- $\kappa \in (0.1000, 0.1650)$
- $\mathcal{O}(15)\beta$ for 3-4 $N_\sigma \in (16, 36)$
- $\mathcal{O}(80k - 200k)$ trajectories
- OpenCL based **CL²QCD** code

Bach, Philipsen, Pinke (arxiv:1411.5219v2)(2014)

<https://github.com/CL2QCD/cl2qcd.git>



P. de Forcrand and O. Philipsen, JHEP (2007)

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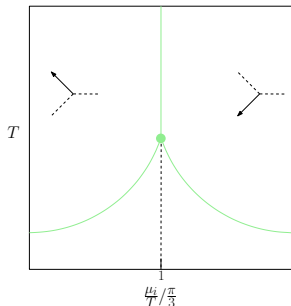
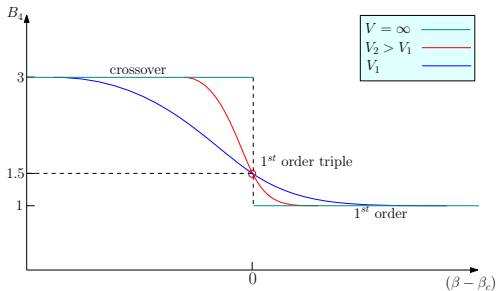
C. Pinke, O. Philipsen, Phys.Rev. D89 (2014)

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Extracting the Order of a Phase Transition

$$B_4(\beta) = \frac{\langle (L_{\text{Im}} - \langle L_{\text{Im}} \rangle)^4 \rangle}{\langle (L_{\text{Im}} - \langle L_{\text{Im}} \rangle)^2 \rangle^2}$$

$$\lim_{V \rightarrow \infty} B_4(\beta_c) = \begin{cases} 1 & , 1^{\text{st}} \text{ order} \\ 1.5 & , 1^{\text{st}} \text{ order triple} \\ 1.604 & , 2^{\text{nd}} \text{ order Z(2)} \\ 3 & , \text{crossover} \end{cases}$$



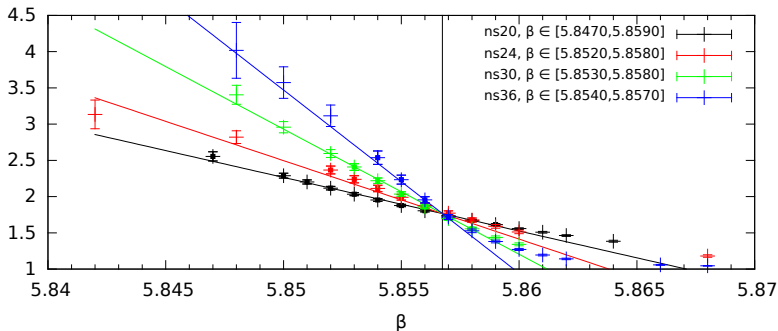
$$B_4(\beta, N_\sigma) = B_4(\beta_c, \infty) + a(\beta - \beta_c) N_\sigma^{1/\nu} + \dots$$

Finite Size Scaling

- Smoothing/interpolating Data points with Ferrenberg-Swensen reweighting.
- Fitting of

$$B_4(\beta_c, N_\sigma) = B_4(\beta_c, \infty) + a(\beta - \beta_c)N_\sigma^{1/\nu}$$

to reweighted points.



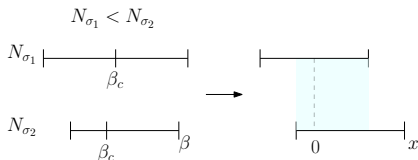
$$\kappa = 0.11, \quad N_\tau = 6$$

Finite Size Scaling

$$\begin{aligned} B_4(\beta_c, N_\sigma) &= B_4(\beta_c, \infty) + a(\beta - \beta_c)N_\sigma^{1/\nu} + \dots \\ &= B_4(\beta_c, \infty) + a \quad x \quad + \dots \end{aligned}$$

- Fit criteria:

- ✓ $Q \approx 50\%$
- ✓ $\chi^2 \approx 1$
- ✓ Overlap in $x \geq 80\%$
- ✓ Symmetry of x around 0

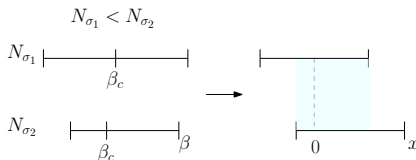


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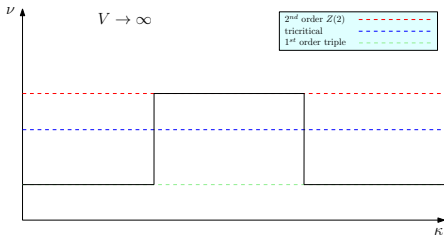
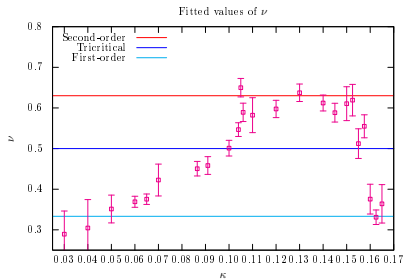


- ν less prone to finite size effects
→ better suited to extract order of phase transition
- Determine order of phase transition according to

$$\nu = \begin{cases} 1/3 & 1^{st} \text{ order triple} \\ 1/2 & \text{tricritical} \\ 0.63 & 2^{nd} \text{ order } 3D \text{ Ising} \end{cases}$$

Results

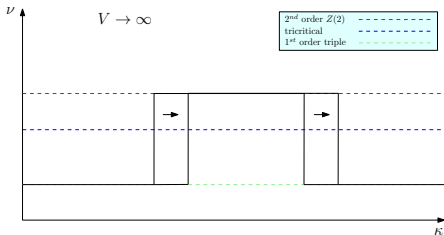
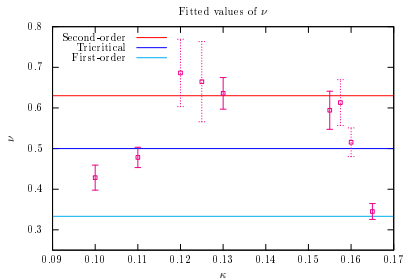
Results for studies with $N_f = 2$ flavors of Wilson fermions on $N_\tau = 4$ lattices



- κ as function of the bare quark mass: $\kappa = (2(am + 4))^{-1}$

Results

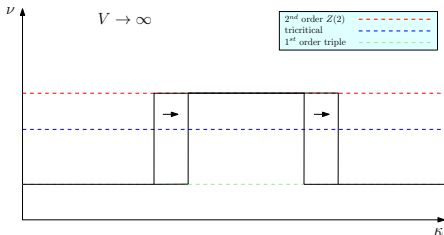
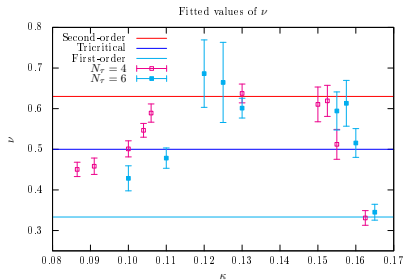
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- Shift in 1st order region to smaller kappa: $T = 1/(aN_\tau)$
 \Rightarrow Tricritical mass: $m_\pi^{tric} \approx 730 \text{ MeV} \rightarrow m_\pi^{tric} \approx 660 \text{ MeV}$

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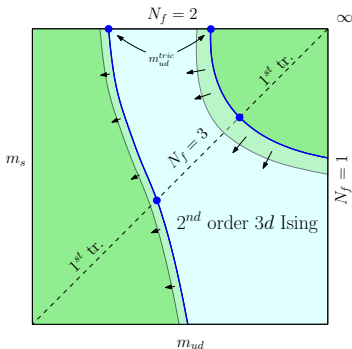
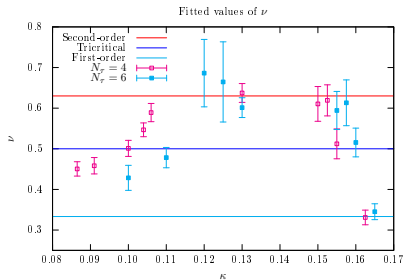
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Summary and Perspectives

Summary

- No sign problem for imaginary chemical potential.
- QCD phase diagram constrained by imaginary chemical potential region.
- $N_\tau = 6$ studies with Wilson fermions ongoing.
- Shift in 1st/2nd order region.

Perspectives

- Compare Wilson results to $N_\tau = 6$ staggered results (ongoing).
- Start studies for $N_\tau = 8$.
- Extend Wilson fermion studies to $N_f = 3$.