

Nucleon generalized form factors from lattice QCD near the physical quark mass

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July 15, 2015, Lattice 2015



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An outline

- 1 GPD motivation
- 2 Lattice setup
- 3 Extracting GFFs
- 4 Excited state analysis
- 5 Results
- 6 Conclusion and outlook

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Generalized Parton Distributions (GPDs)

- introduced late '90s
- comprehensive description of the hadron structure

Contain information about

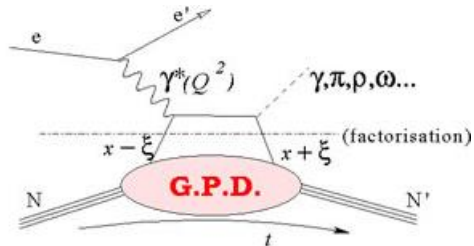
- traditional form factors and parton distributions (limiting cases)
- quark total angular momentum

In particular important for

- spin structure of the proton
- deeply virtual Compton scattering (DVCS)

Kinematic definitions

- $\Delta = P' - P$
- $\bar{P} = (P' + P)/2$
- $n \cdot n = 0$
- $n \cdot \bar{P} = 1$



Kinematic variables

- x : longitudinal momentum fraction
- $\xi = -n \cdot \Delta/2$: longitudinal momentum transfer
- $Q^2 = -t = \Delta^2 = P' - P$: momentum transfer

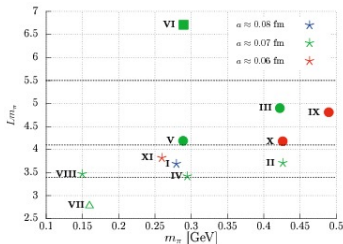
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General setup

- $N_f = 2$ mass-degenerated NP improved Wilson-clover Fermions with Wilson Gauge action (RQCD, QCDSF)
- Apply Wuppertal smearing to suppress excited states
- Kinematic setup: $\vec{p}_f \cdot \vec{p}_f = 0$ and $\vec{p}_i \cdot \vec{p}_i \neq 0$

Today's setup

- Show results for ensemble VIII
- 150 MeV m_π
- $L \cdot m_\pi = 3.47$
- $t_f/a = 9, 12, 15$ with $t_{15} \sim 1$ fm



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Consider Vector GFF

$$\langle N(P') | \mathcal{O}_V^{\mu\nu} | N(P) \rangle = S(\mu, \nu) \bar{U}(P') \left[i\gamma^\mu \bar{P}^\nu A_{2,0}(t) + \frac{i\sigma^{\mu\alpha} \Delta^\alpha \bar{P}^\nu}{2m_N} B_{2,0}(t) + \frac{\Delta^\mu \Delta^\nu}{m_N} C_2(t) \right] U(P)$$

Remarks:

- Traces are already subtracted
- Euclidean signature (after Wick-Rotation)
- Compute the (kinematic) prefactors of the GFFs They define the coefficient matrix M

Extract the GFF by solving an overdetermined system of equations:

$$\epsilon = \left[M \cdot \begin{bmatrix} A_{20} \\ B_{20} \\ C_2 \end{bmatrix} - c_* \right]^T \cdot \text{cov}^{-1}(c_*) \cdot \begin{bmatrix} M \cdot \begin{bmatrix} A_{20} \\ B_{20} \\ C_2 \end{bmatrix} - c_* \end{bmatrix}$$

Remarks:

- Find A_{20} , B_{20} , C_2 such that ϵ is minimized
- c_* can be fitted from lattice data

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Imporant (first) step:

- Suppress excited states (Wuppertal smearing)

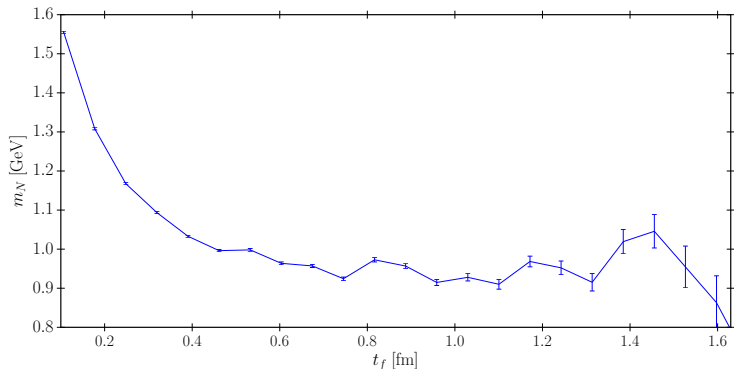


Figure : Effective mass plot of a nucleon (Ens. VIII).

Strategy of the analysis

- Perform excited state analysis for GFFs with simultaneous fits to 2-pt and 3-pt functions (extract c_*)
- Carefully analyze introduced systematics (e.g. fit ranges)

This is challenging

- More noise due to derivative
- Even more noise if $Q^2 \neq 0$
- Fit ansatz more complicated if $Q^2 \neq 0$

Key ingredients

- Improve 3-pt signal by averaging data of equal rows of M
- Simultaneous fit of all c_* (stabilizing noisy channels)

Nucleon 2-pt function with one exponent:

$$C_2^{1e}(p, t) = z_N(\vec{p}) \frac{E(\vec{p}, m_N) + m_N}{E(\vec{p}, m_N)} e^{-E(\vec{p}, m_N)t}$$

Nucleon 2-pt function with two exponents:

$$C_2^{2e}(p, t) = z_N(\vec{p}) \frac{E(\vec{p}, m_N) + m_N}{E(\vec{p}, m_N)} e^{-E(\vec{p}, m_N)t} \\ + z_{N'}(\vec{p}) \frac{E' + m_{N'}}{E'} e^{-E't}$$

Remarks

- Assume continuum dispersion relation for N but not for N'

Nucleon 3pt function:

$$\begin{aligned}
 C_3(p, t) = & \sqrt{z_N^0 z_N^{pi}} \mathbf{c} \mathcal{E}(t, m_N, E(\vec{p}, m_N)) + \\
 & \sqrt{z_N^0 z_{N'}^{pi}} x \mathcal{E}(t, m_N, E') + \\
 & \sqrt{z_{N'}^0 z_N^{pi}} x \mathcal{E}(t, m'_N, E(\vec{p}, m_N)) +
 \end{aligned}$$

with

$$\mathcal{E}(t, M, E) = e^{-M(t_{\text{sink}} - t)} e^{-E(t - t_{\text{src}})}$$

Remarks

- Need to get \mathbf{c}

Scan fit range dependency

- Fit range 2-pt 1 exp.: (9,25), (10,25), (11,25), (12,25), (13,25)
- Fit range 2-pt 2 exp.: (1,25), (2,25), (3,25), (4,25), (5,25)
- Fit range 3-pt: (2,14), (3,13), (4,12), (4,25), (5,25)

Do resampling

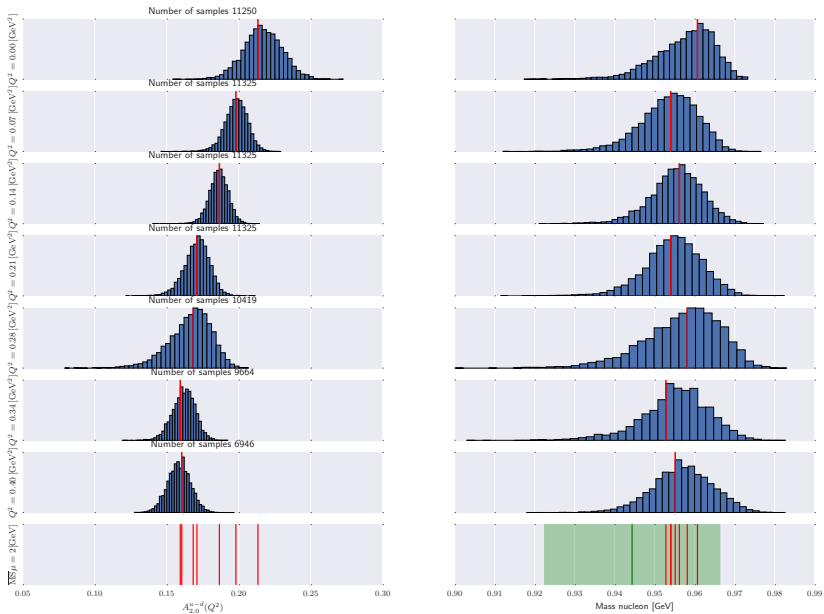
- Create bootstrap ensemble ($ss = 150$) for each fit range combination

Remarks

- Solve the $5 \cdot 5 \cdot 3 \cdot 151 = 11325$ overdetermined systems
- Create a histogram
- Weight the fits properly
- Mix statistical and systematic distributions

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$A_{2,0}^{u-d}(Q^2)$ for Ens. VIII with ($t_f = 15$)



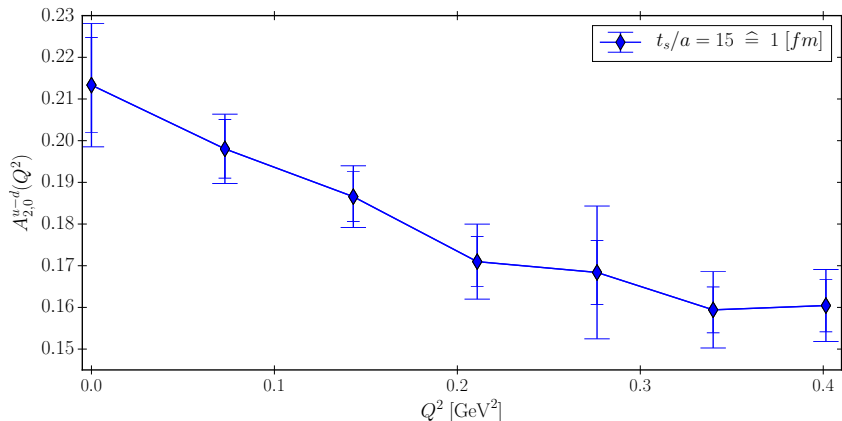


Figure : Inner error is statistical, outer error is total

Remarks:

- Linear behavior for small Q^2 as expected by LO χ -PT
- Finite volume effects of $\langle x \rangle_{u-d} = A_{2,0}^{u-d}(Q^2 = 0)$ are small (χ -PT result)

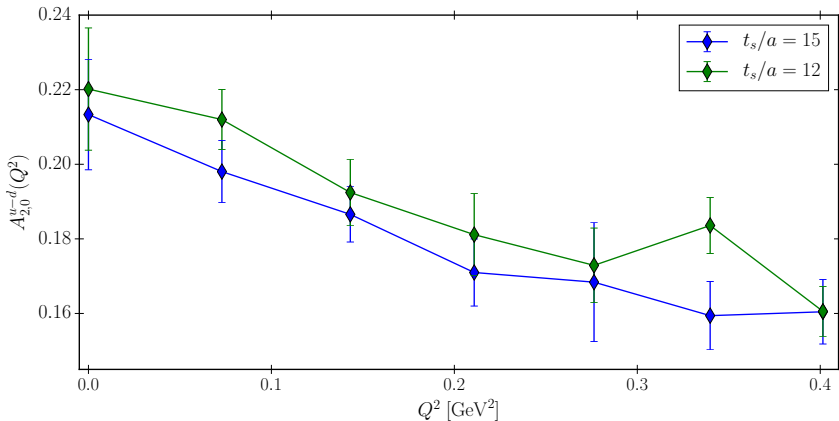


Figure : Comparison of two different sink locations.

Remarks:

- $A_{2,0}^{u-d}(Q^2 = 0) = \langle x \rangle_{u-d} \neq 0.155$ is no x-state effect

Ji's sumrule

$$J^{u-d} = \frac{A_{2,0}^{u-d}(Q^2 = 0) + B_{2,0}^{u-d}(Q^2 = 0)}{2}$$

Remarks

- $B_{2,0}^{u-d}(Q^2 = 0)$ not directly accessible.
- Calculate $(A_{2,0}^{u-d}(Q^2 \neq 0) + B_{2,0}^{u-d}(\neq 0))/2$.
- Perform linear fit and draw histogram (preliminary)

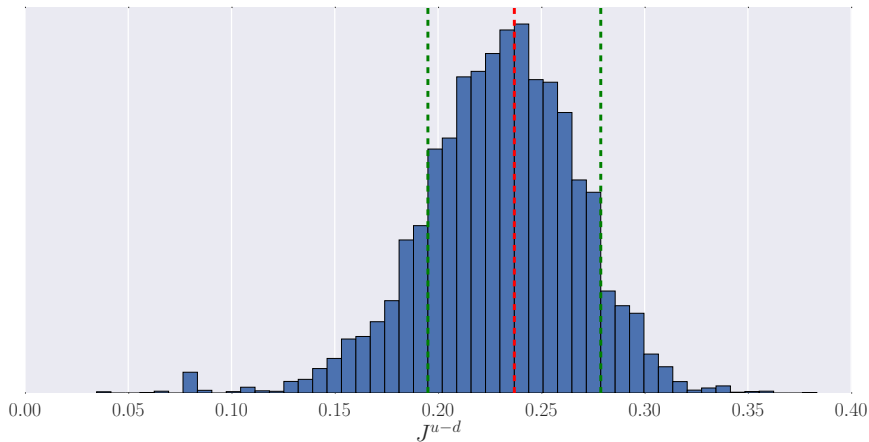


Figure : $J^{u-d} = 0.237 \pm 0.042$ for Ens. VIII with $(t_s/a = 15)$.

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Conclusion

We have demonstrated our excited states analysis for some GFFs on our lightest ensemble. At least in this case excited states are under control at a source sink separation of 1 [fm].

To do

- Apply analysis to all ensembles.
- Apply different weights to the histograms.
- Perform combined χ -PT fits to all ensembles simultaneously.