

Complex Langevin simulation in condensed matter physics

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Complex Langevin Method

Euclidean path integral:

$$Z = \int D\Phi e^{-S[\Phi]}$$

$$S[\Phi] \neq S^*[\Phi] \longrightarrow \text{"sign problem"}$$

ex) chemical potential, electric field, theta term, real time,
matrix model, non-relativistic system, spin system

Complex Langevin Method

complex Langevin equation:

$$\Phi \in \mathbf{R} \rightarrow \tilde{\Phi} \in \mathbf{C}$$

$$\frac{\partial}{\partial \theta} \tilde{\Phi} = -\frac{\delta}{\delta \tilde{\Phi}} S[\tilde{\Phi}] + \eta$$

fictitious time

noise field

expectation value:

$$\langle \hat{O}[\tilde{\Phi}] \rangle = \lim_{\theta \rightarrow \infty} \langle \hat{O}[\tilde{\Phi}(\theta)] \rangle_{\eta}$$

Complex Langevin Method

complex Langevin simulations of

Bose system & Fermi system

in condensed matter physics

- ✓ for the ab-initio analysis of condensed matter systems
- ✓ for a test of the complex Langevin method

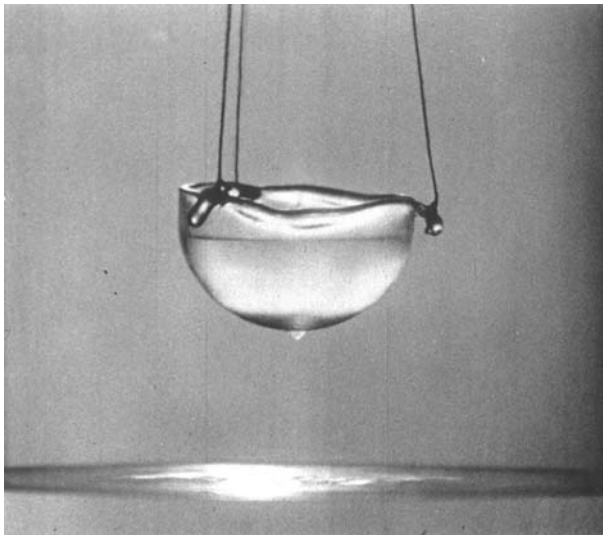
Bose System

non-relativistic Bose gas:

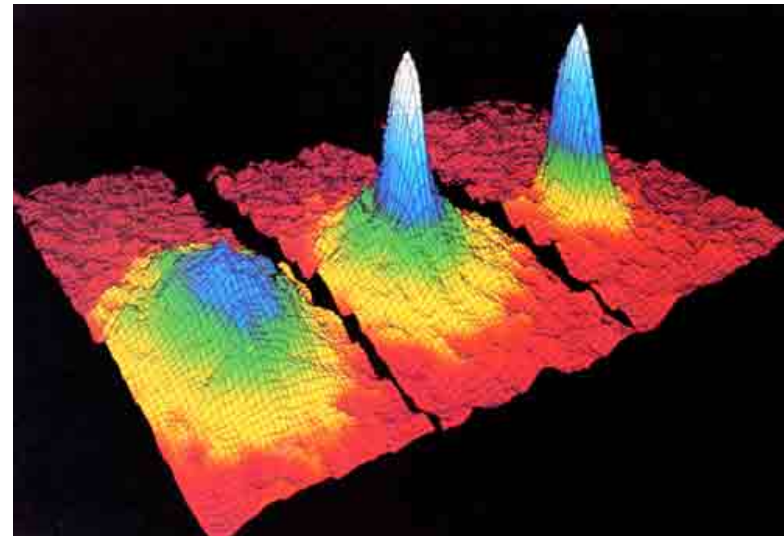
$$S_0[\Phi_1, \Phi_2] = \int d\tau d^3x \left[\Phi^*(x) \left(\frac{\partial}{\partial \tau} - \mu - \frac{1}{2m} \Delta \right) \Phi(x) + \frac{1}{4} \lambda |\Phi(x)|^4 \right]$$

anti-Hermitian

superfluid ^4He
from wikipedia

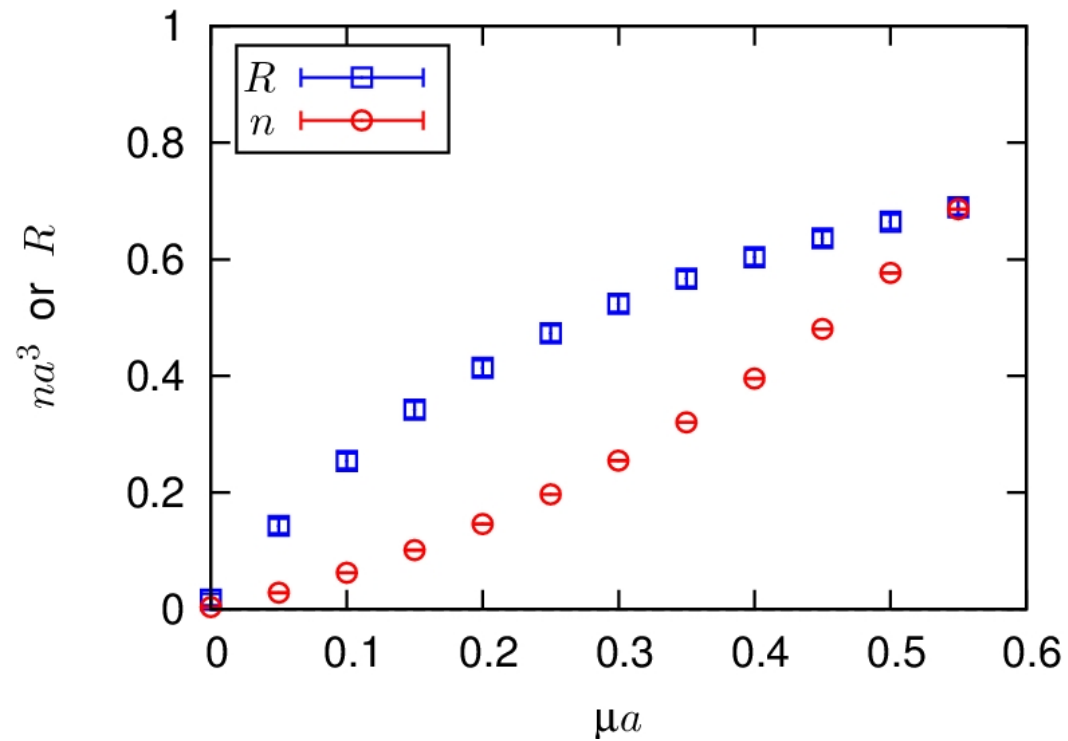


atomic BEC
JILA group experiment



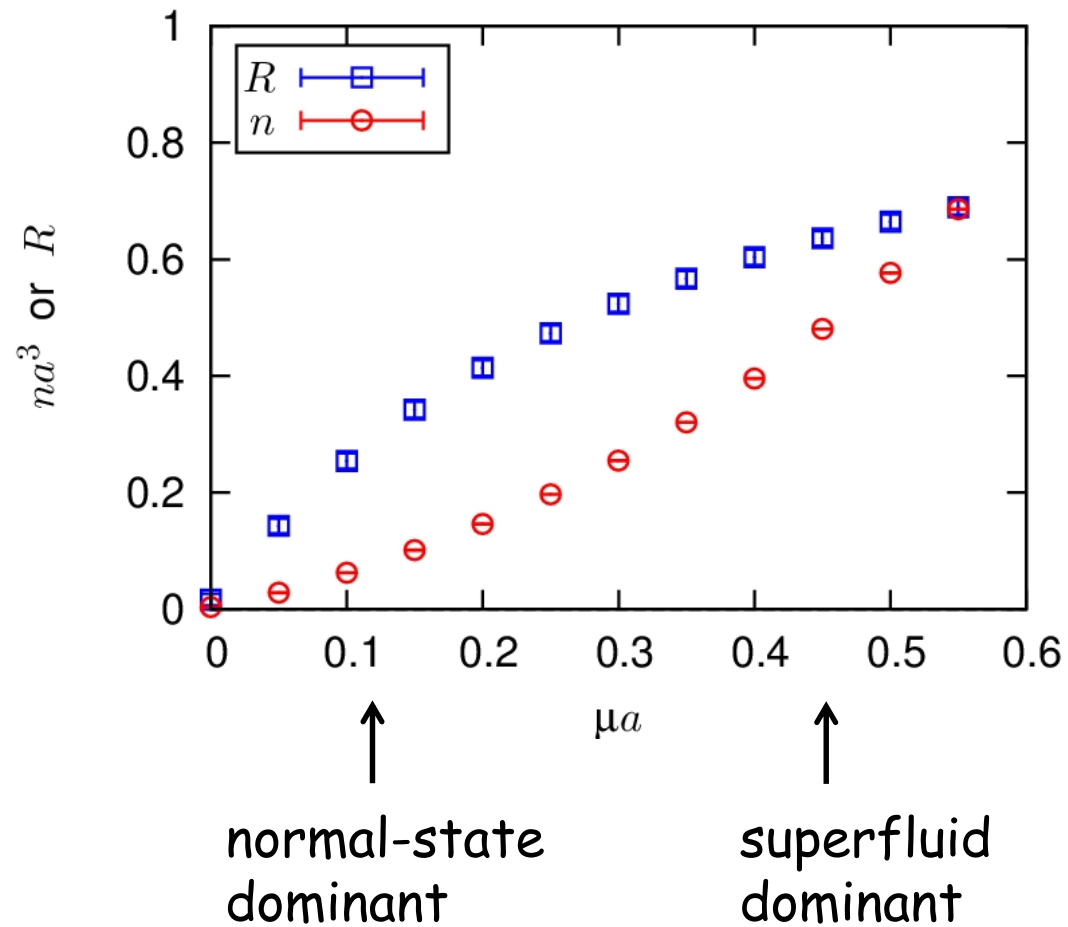
Bose System

condensate fraction:
$$R = \lim_{|x-y| \rightarrow \infty} \frac{\langle \Phi^*(x)\Phi(y) \rangle}{\langle \Phi^*(x)\Phi(x) \rangle}$$



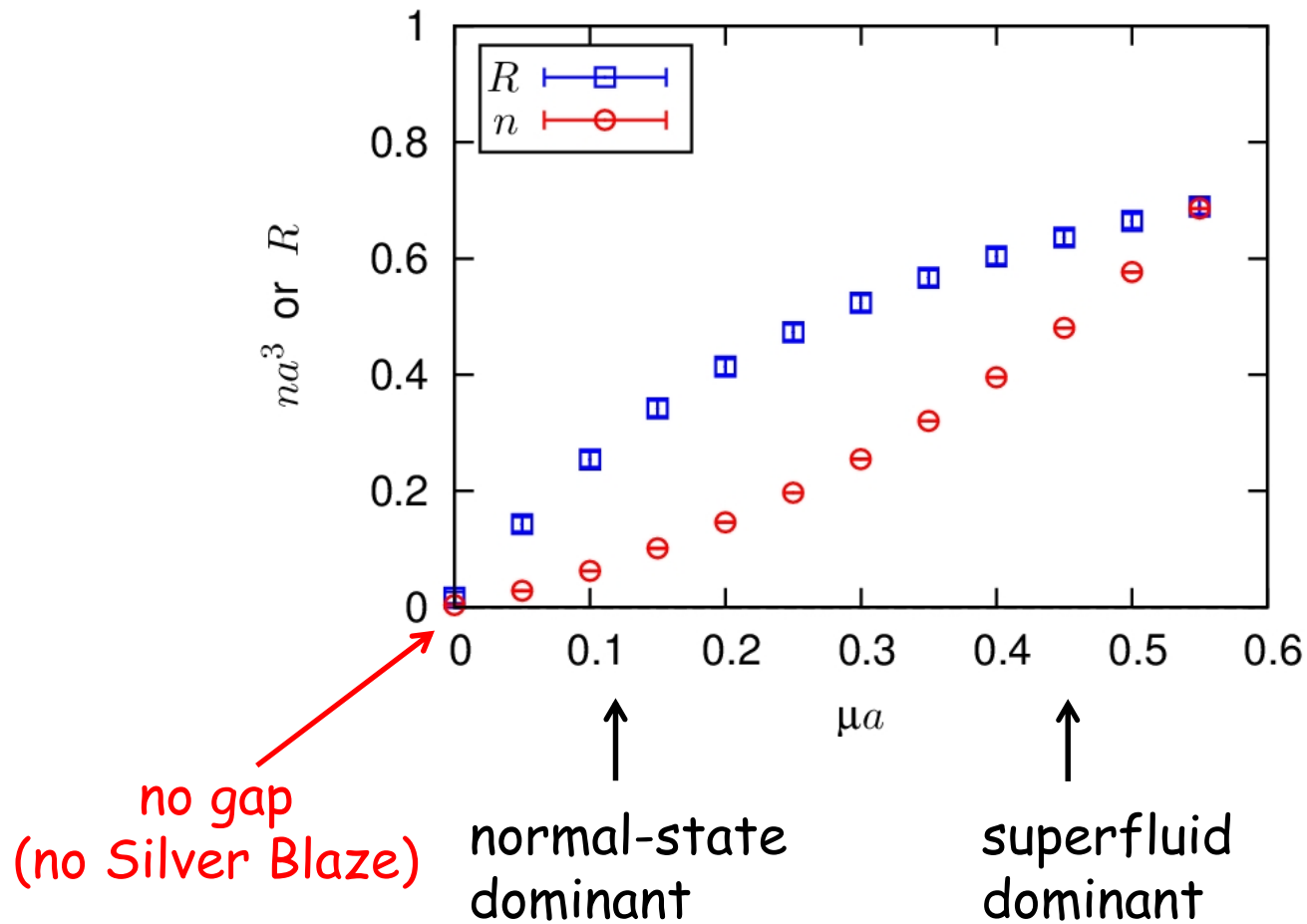
Bose System

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Bose System

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Rotation

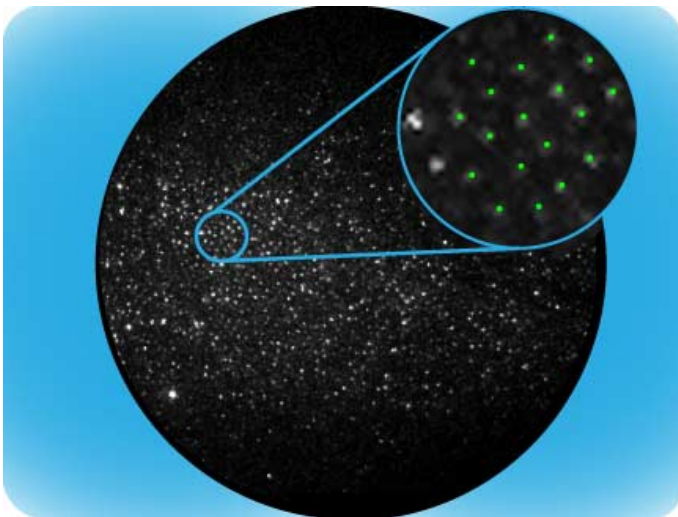
rotating Bose gas:

$$S_{\Omega} = S_0 - \Omega L_z$$

$$L_z = -i \int d\tau d^3x \Phi^*(x) \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right) \Phi(x)$$

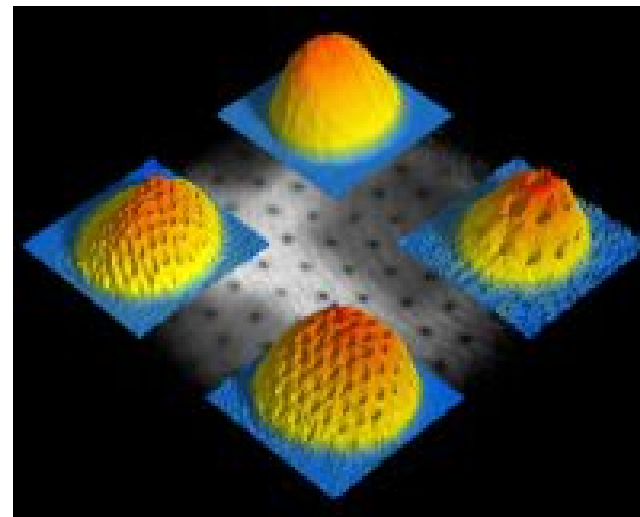
vortex in ^4He

Maryland group experiment

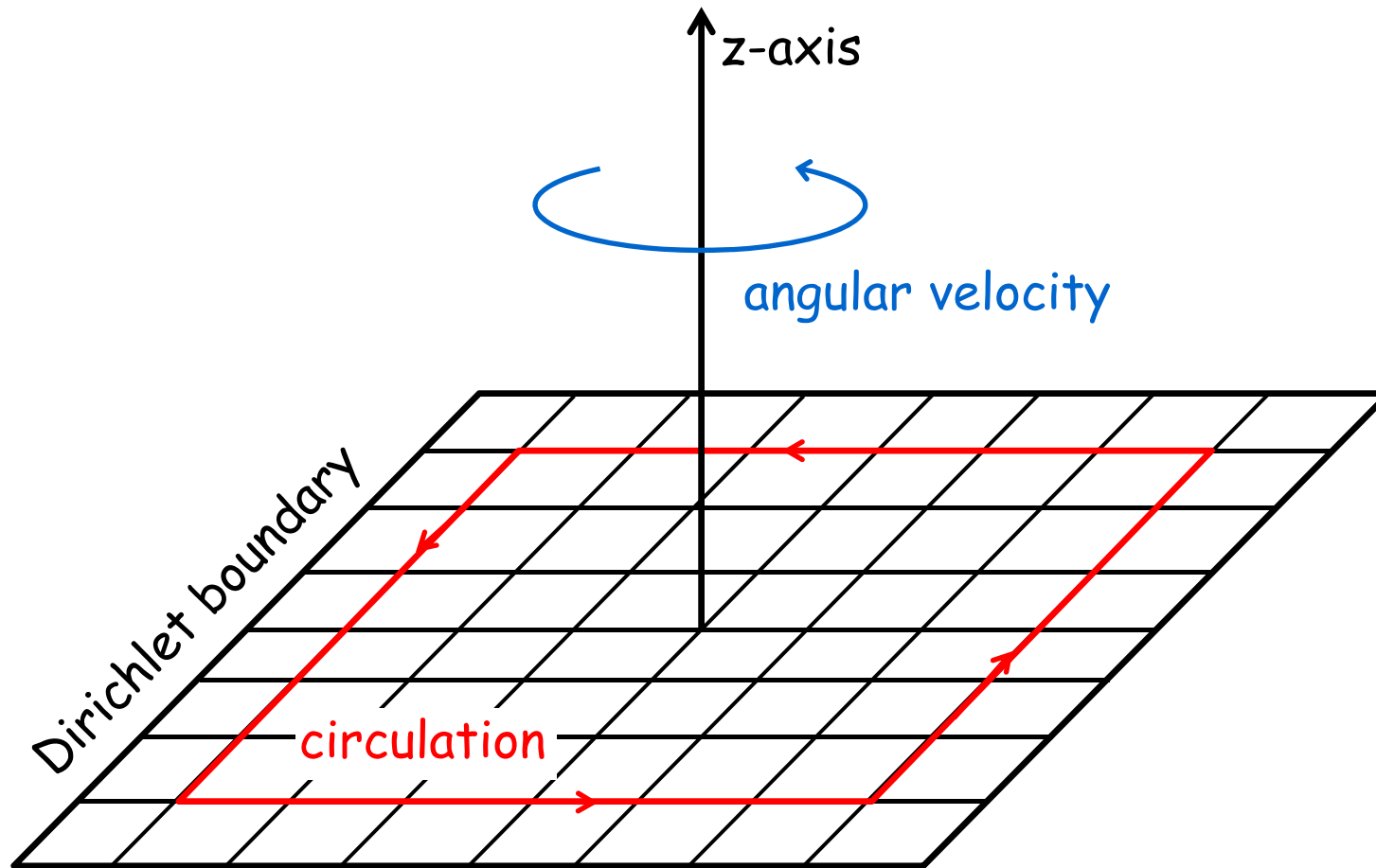


vortex in atomic BEC

MIT group experiment



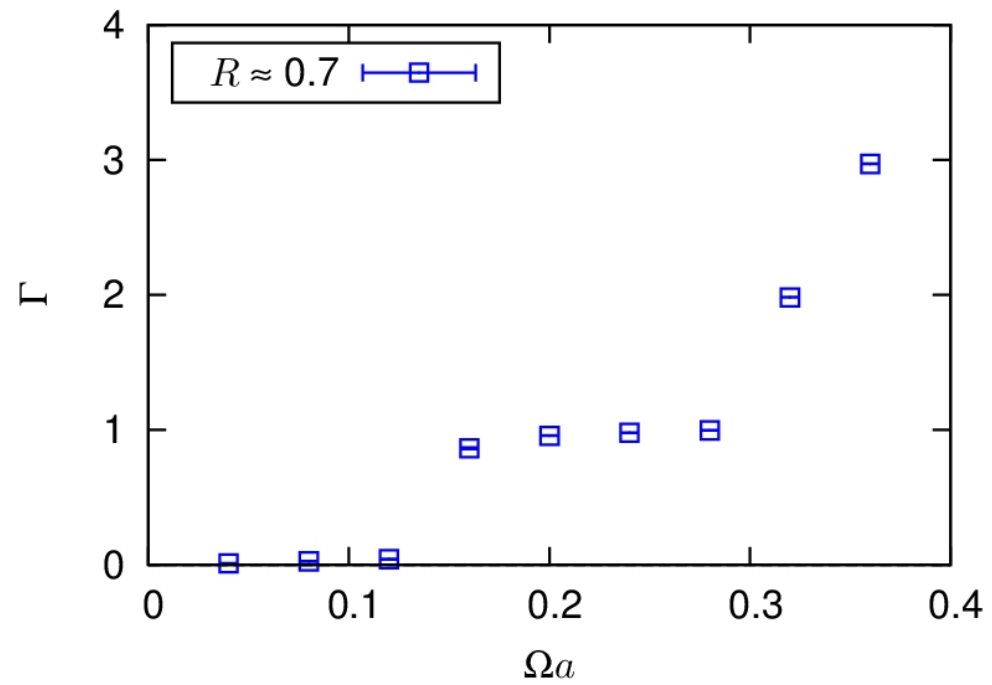
Rotation



circulation : $\hat{\Gamma} = \frac{1}{2\pi} \oint \delta\Theta = N_{\text{vortex}}$

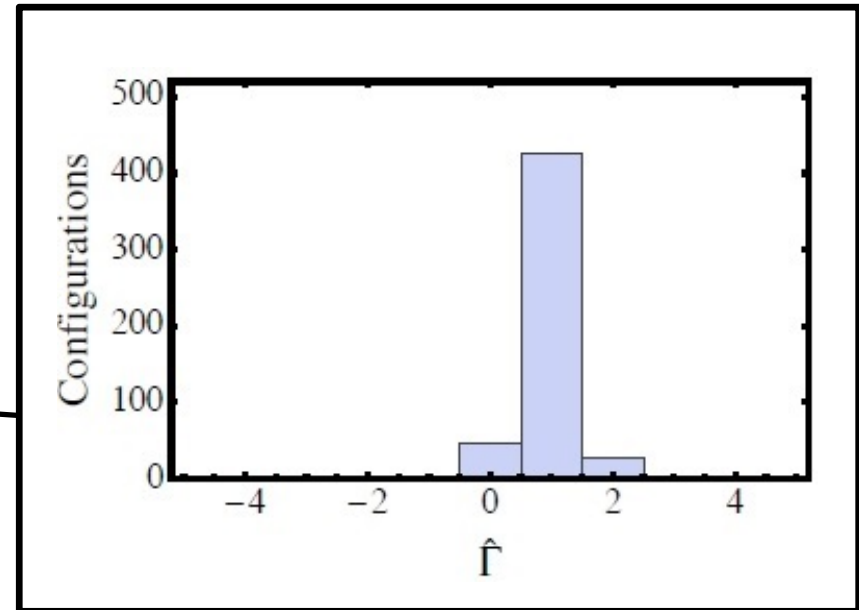
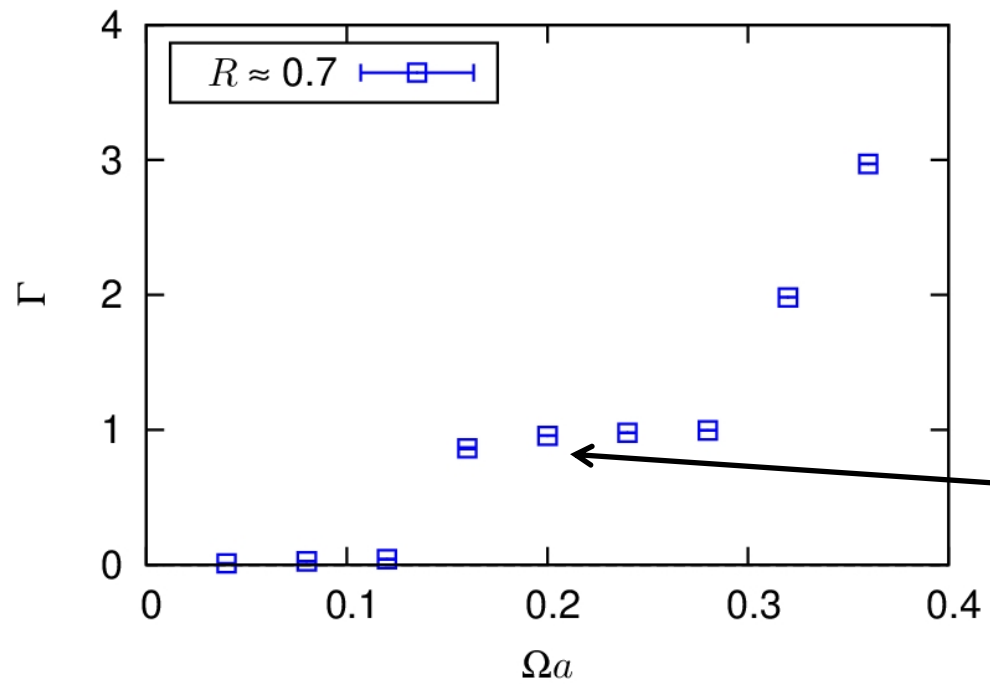
Rotation

circulation (superfluid dominant)



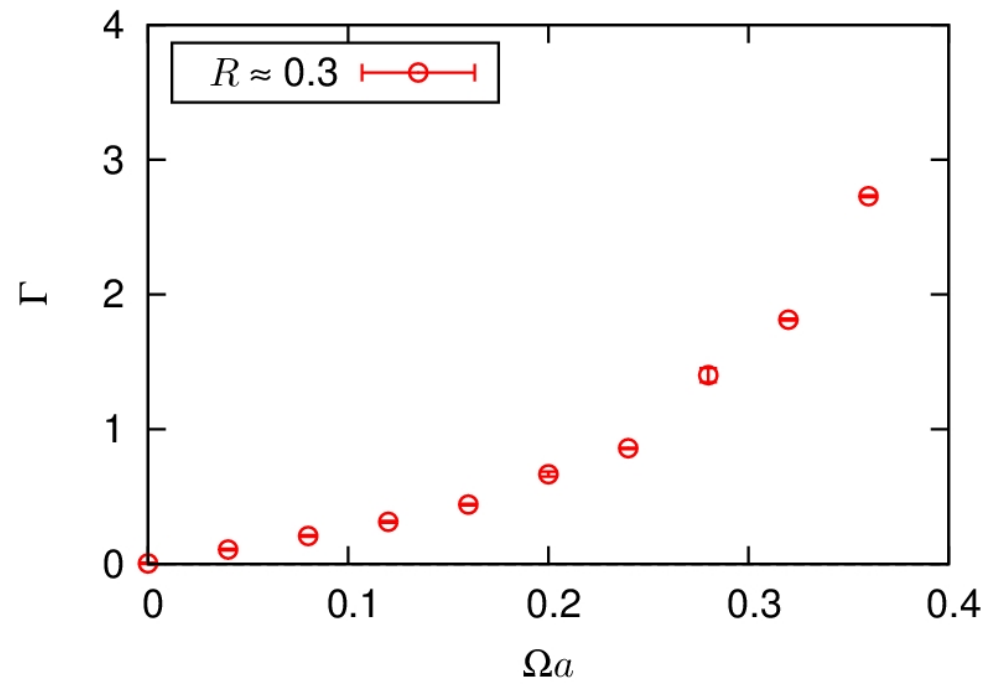
Rotation

circulation (superfluid dominant)



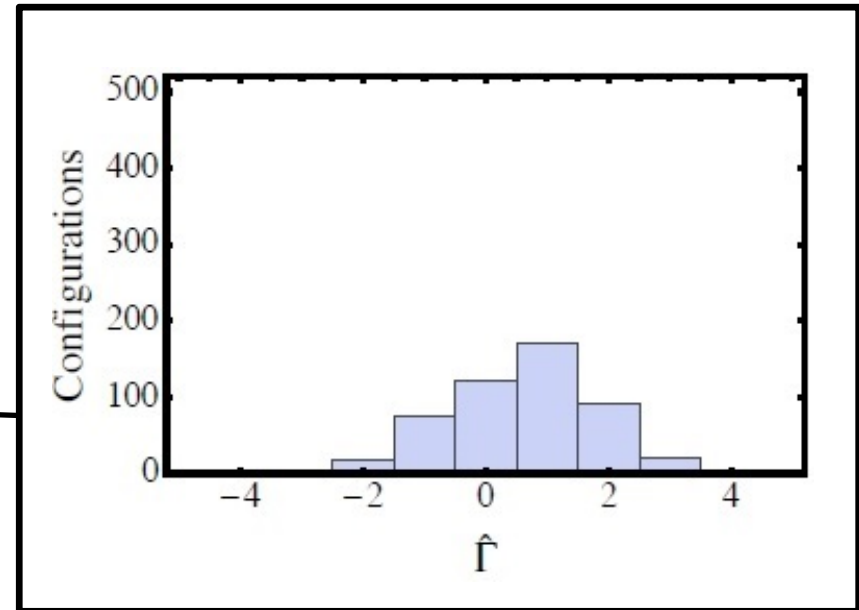
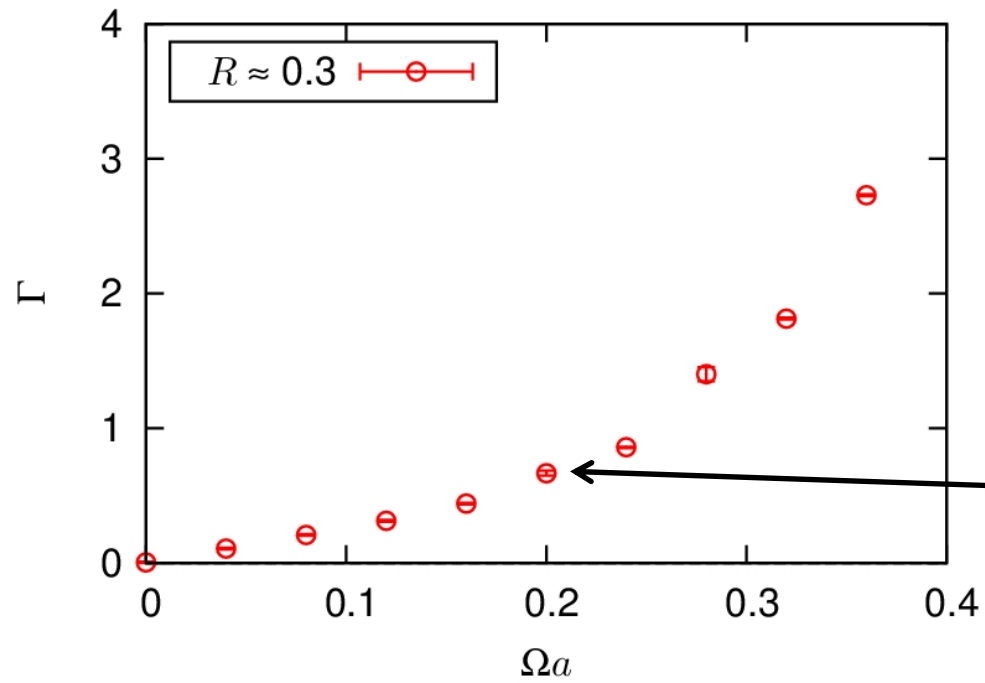
Rotation

circulation (normal-state dominant)



Rotation

circulation (normal-state dominant)

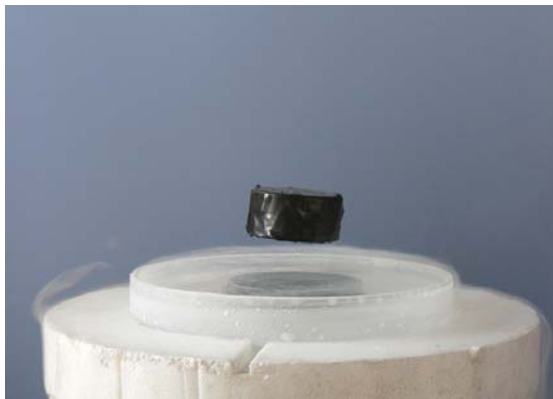


Fermi System

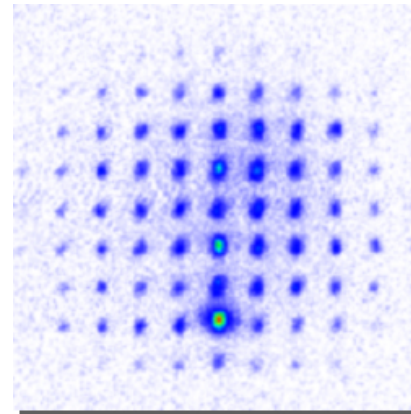
Hubbard model:

$$S[\Psi_{\uparrow}^*, \Psi_{\uparrow}, \Psi_{\downarrow}^*, \Psi_{\downarrow}] = \int d\tau \sum_x \left[\sum_{i=\uparrow, \downarrow} \Psi_i^*(x) \left(\frac{\partial}{\partial \tau} - \mu_i \right) \Psi_i(x) \right. \\ \left. - \sum_{i=\uparrow, \downarrow} \sum_j t_i (\Psi_i^*(x) \Psi_i(x + \hat{j}) + \Psi_i^*(x + \hat{j}) \Psi_i(x)) \right. \\ \left. + U \Psi_{\uparrow}^*(x) \Psi_{\uparrow}(x) \Psi_{\downarrow}^*(x) \Psi_{\downarrow}(x) \right]$$

high-Tc superconductor
from wikipedia



optical lattice of atom gas
LMU-MPQ group experiment



Fermi System

Hubbard model:

$$Z = \int D\Phi \det K_{\uparrow}[\Phi] \det K_{\downarrow}[\Phi] e^{-S_A[\Phi]}$$

$$K_i[\Phi] = \frac{\partial}{\partial \tau} - \mu_i + \Phi(x) - \sum_j t_j (T_{+j} + T_{-j})$$

$$S_A[\Phi] = \int d\tau \sum_x \frac{1}{2|U|} \Phi^2(x)$$

Fermi System

$$\det K[t, \mu] \in \mathbf{R}$$

balanced:

$$\det K[t, \mu] \det K[t, \mu] \geq 0$$

no problem

imbalanced:

$$\det K[t_{\uparrow}, \mu_{\uparrow}] \det K[t_{\downarrow}, \mu_{\downarrow}] \geq 0$$

sign problem

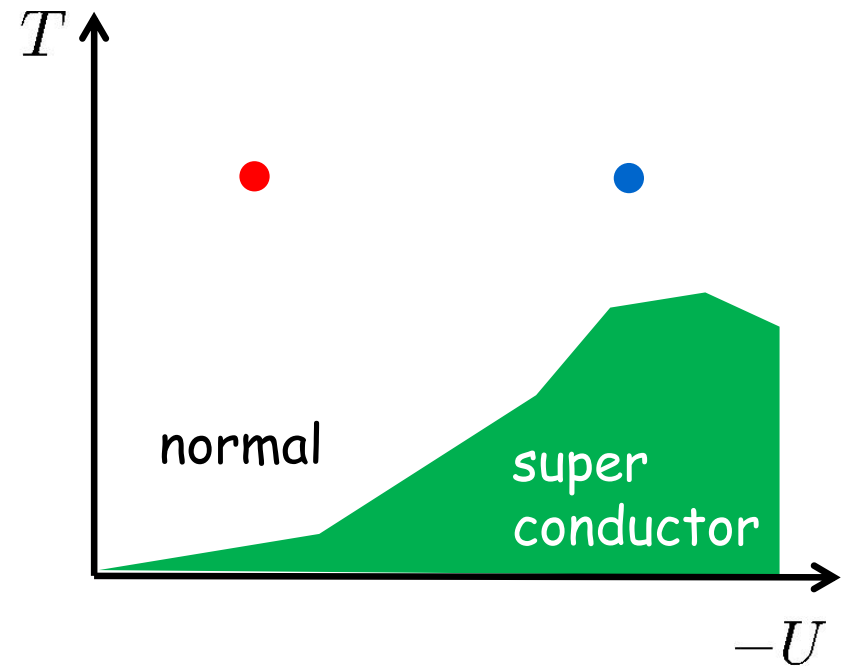
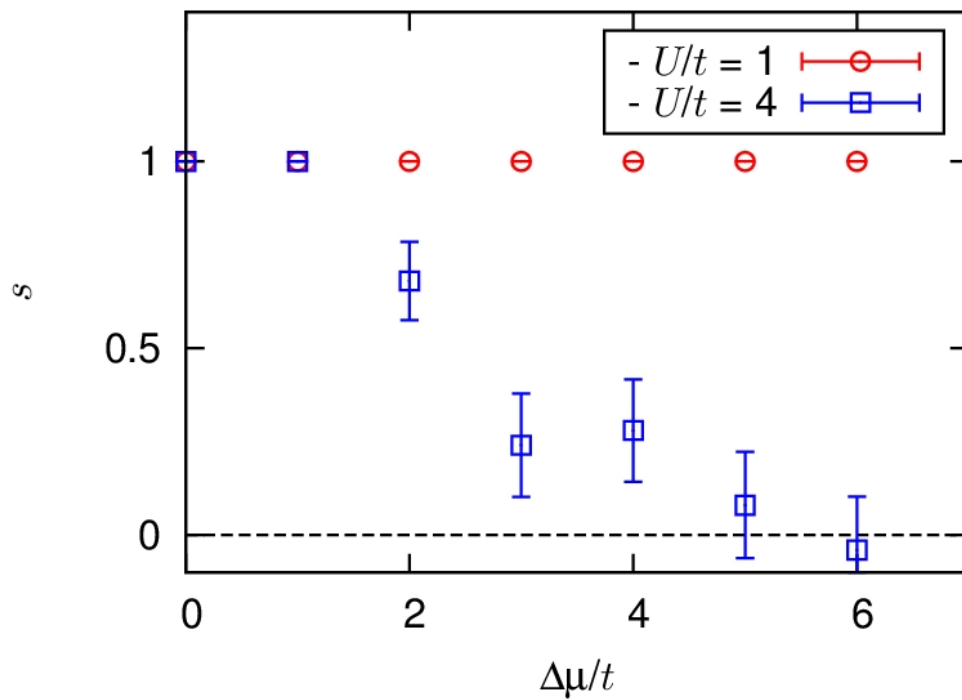
cf) QCD

$$\det D[m, \mu] \det D[m, \mu] \in \mathbf{C}$$

sign problem

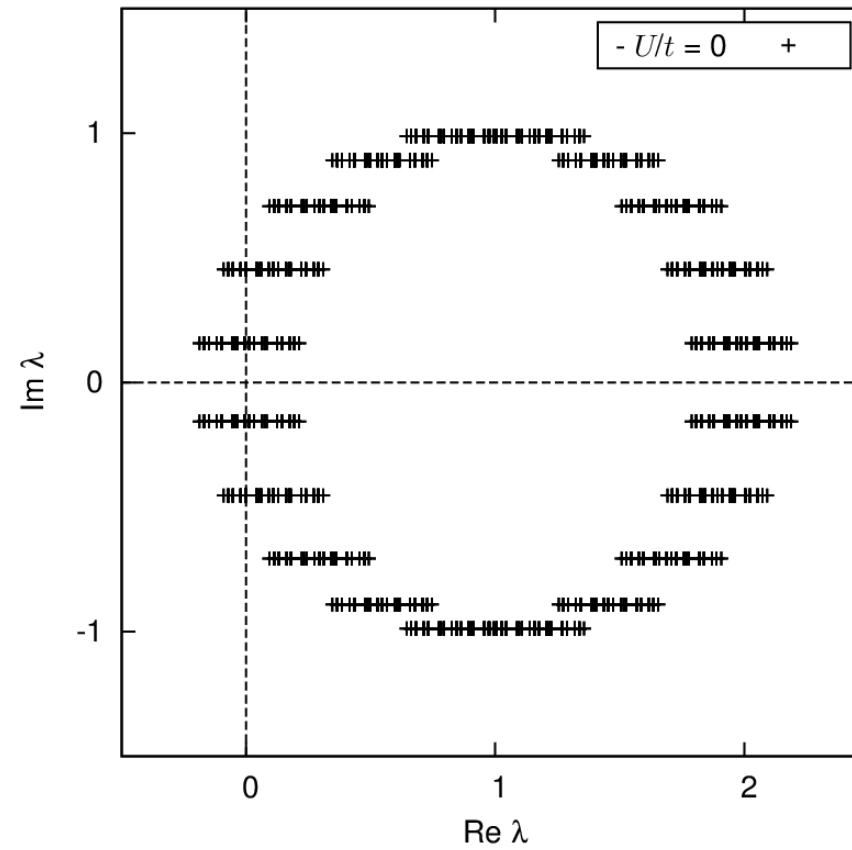
Fermi System

average sign in "sign-quenched" Monte Carlo



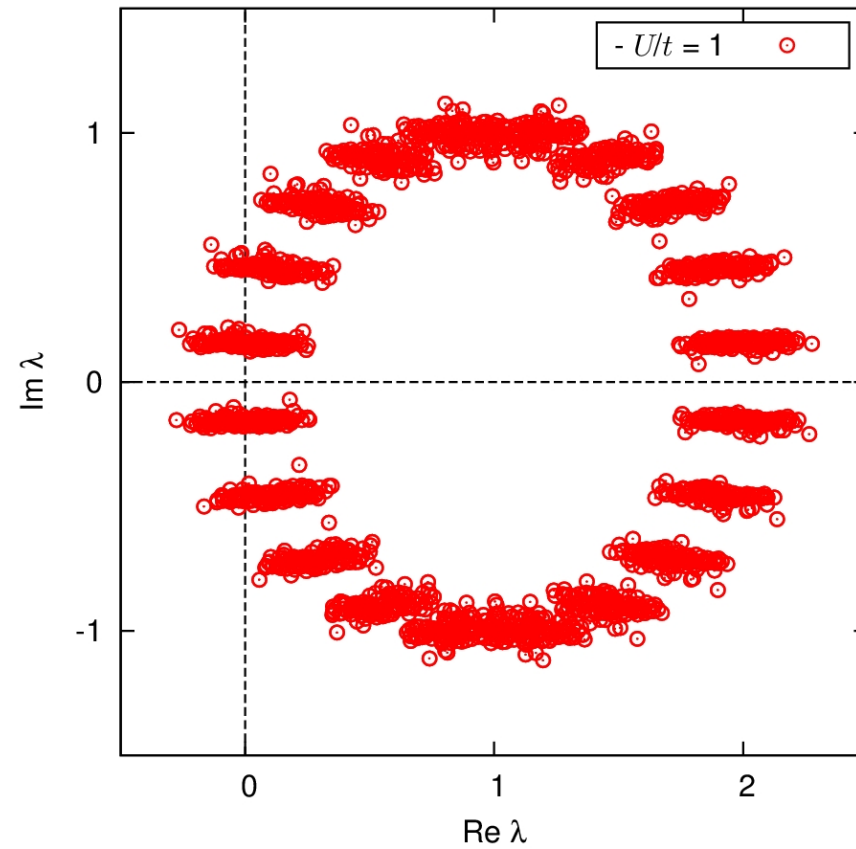
Fermi System

eigenvalues (free)



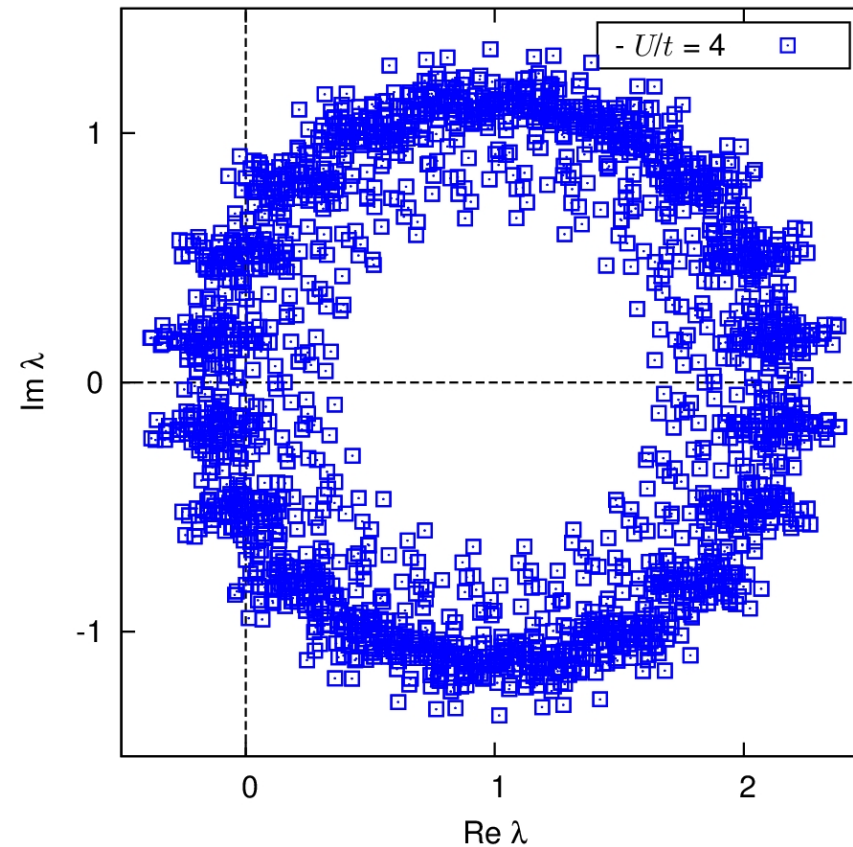
Fermi System

eigenvalues (weak coupling)



Fermi System

eigenvalues (strong coupling)

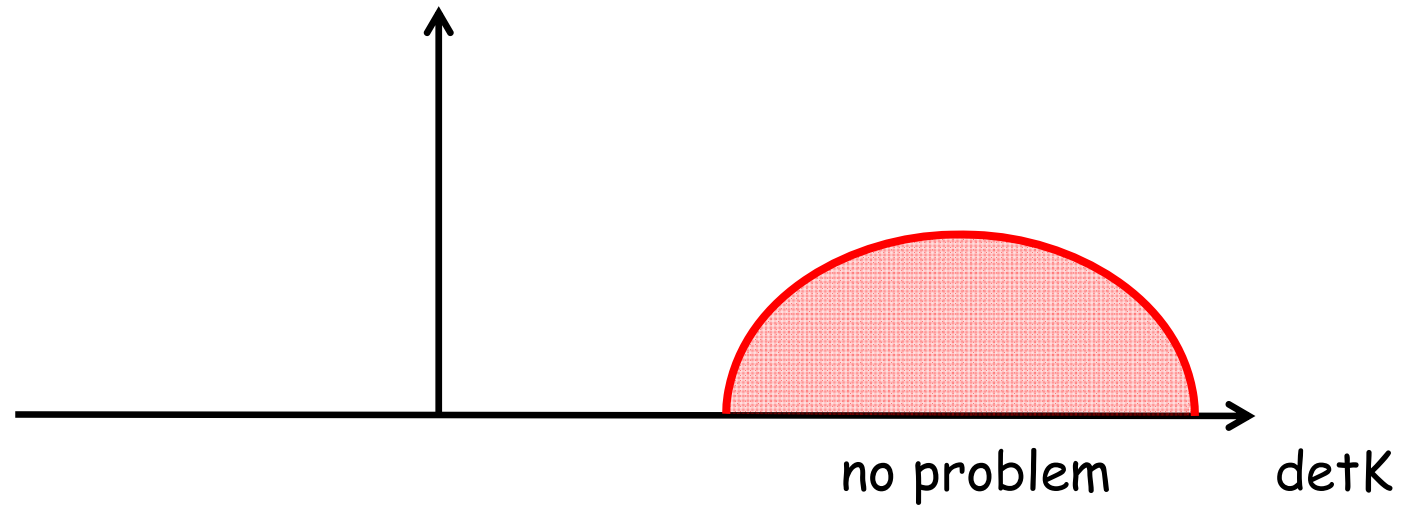


logarithmic singularity at zero eigenvalues

Mollgaard Splittorff '13, Greensite '14, Nishimura Shimasaki '15

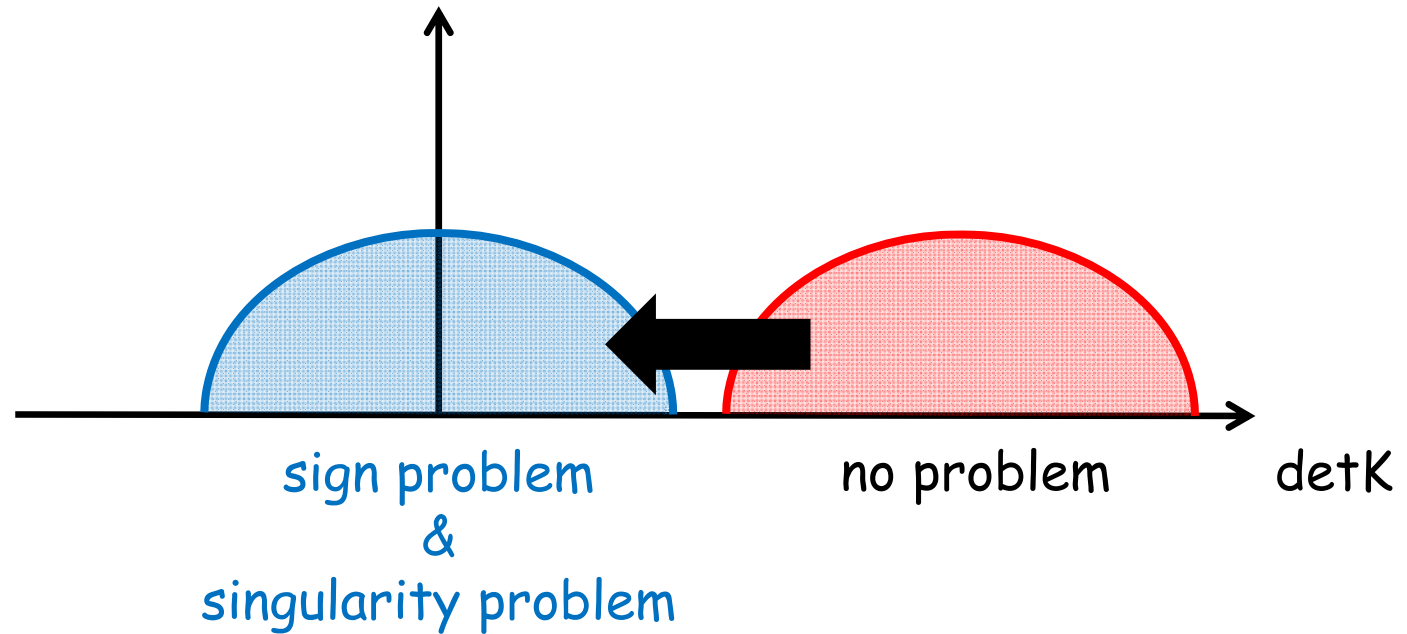
Fermi System

probability distribution of $\det K$



Fermi System

probability distribution of $\det K$



Summary

Boson is good.

Fermion is bad.