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# Analysis of short-distance current correlators using OPE

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for JLQCD Collaboration

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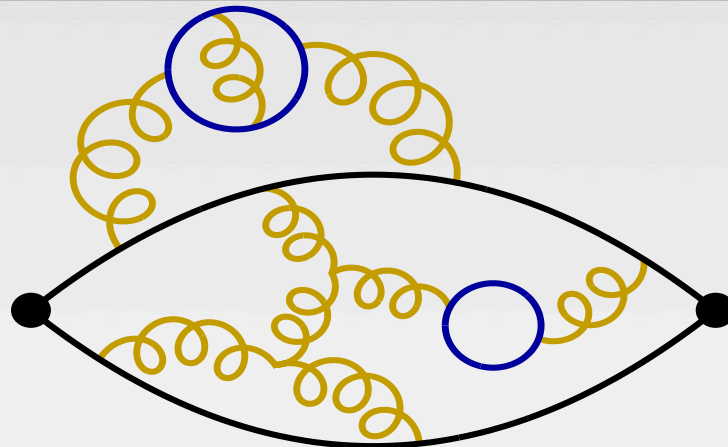
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# Current Correlators



High-energy



- Perturbative at short-distance  
(Continuum PT well-developed, 4-loop)
  - ▲ OPE in the middle
  - Non-perturbative at long-distance ( $\gtrsim 1/\Lambda_{\text{QCD}}$ )  
ChPT
- 0.1fm
- 1fm
- Test & use  
This region

# Operator Product Expansion (OPE)

$$\Pi(x) = \frac{c_0(\alpha_s)}{x^6} + \frac{c_2 m_q^2}{x^4} + \frac{c_{4,\bar{q}q} m_q \langle \bar{q}q \rangle + c_{4,G} \langle GG \rangle + \dots}{x^2} + \dots$$

Massless Perturbation  
(up to 4-loop)

tiny mass correction

First non-perturbative  
effect through condensates

- Test this theoretical formulation against the “experiment” on the lattice
- Use as an input for NPR

# Two-point Functions

- Definitions

$$\Pi_S(x) = \langle S(x)S(0)^\dagger \rangle, \quad \Pi_P(x) = \langle P(x)P(0)^\dagger \rangle,$$

$$\Pi_{V,\mu\nu}(x) = \langle V_\mu(x)V_\nu(0)^\dagger \rangle, \quad \Pi_{A,\mu\nu}(x) = \langle A_\mu(x)A_\nu(0)^\dagger \rangle,$$

$$\Pi_{V/A}(x) = \sum_\mu \Pi_{V/A,\mu\mu}(x)$$

- Non-singlet local operators are used

$$S(x) = \bar{u}d(x), \quad P(x) = \bar{u}i\gamma_5d(x),$$

$$V_\mu(x) = \bar{u}\gamma_\mu d(x), \quad A_\mu(x) = \bar{u}\gamma_\mu\gamma_5d(x)$$

- Degeneracy

- Massless perturbation

$$\Pi_S = \Pi_P, \quad \Pi_{V,\mu\nu} = \Pi_{A,\mu\nu}$$

- Non-perturbative effects breaks these degeneracies

# Outline

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## 0. Set up

- Ensembles: 2+1-flavor configs with Möbius domain-wall fermion

## 1. Test of OPE

- Subtraction/Elimination of discretization effects at short distances
- Convergence of perturbation theory
- Effect of condensates

## 2. Renormalization

- Discrimination of discretization effects and non-perturbative
- $Z_V, Z_A$
- $Z_S, Z_P$

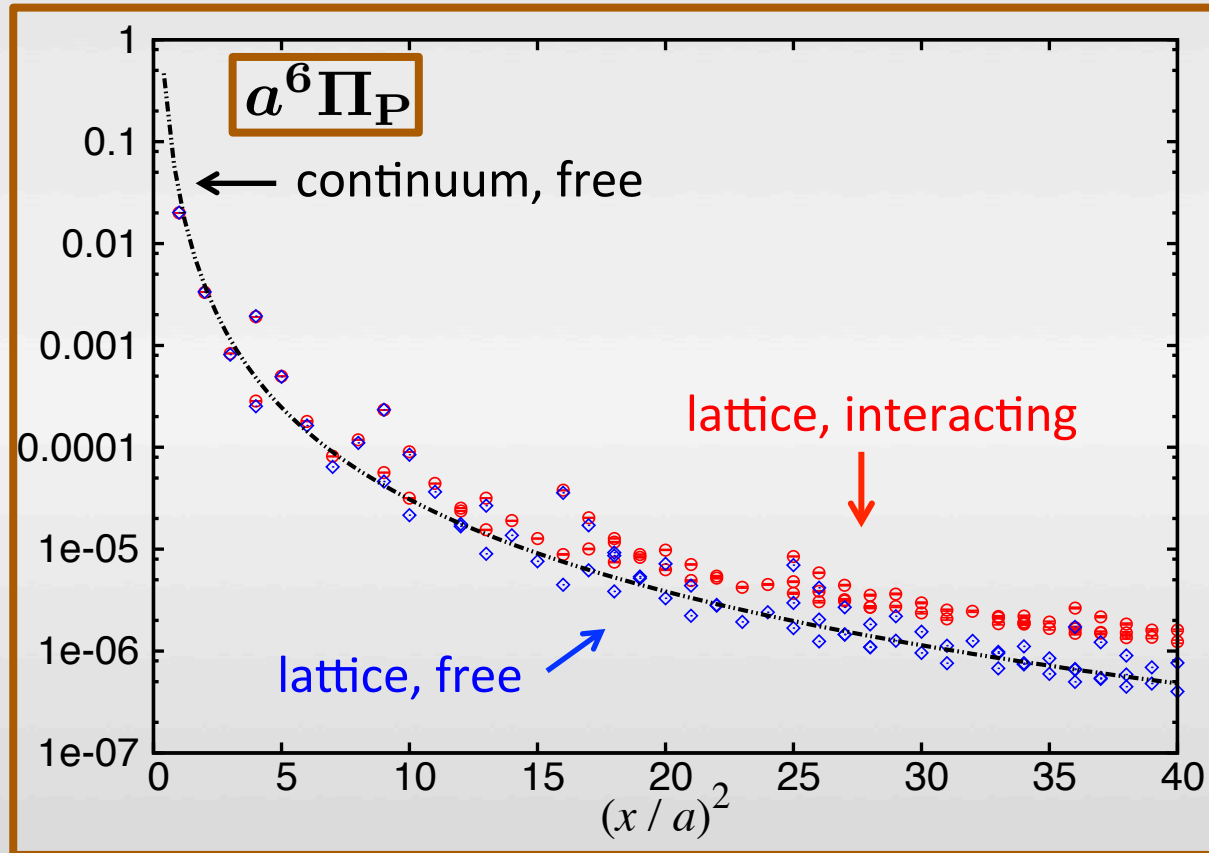
# Lattice Setup

- 2+1 Möbius DW fermions w/ 3 times stout smeared links
- Symanzik improved gauge action
- 3 kinds of different cutoff w/ nearly same physical volume

$a^{-1}$ [GeV]	Volume	$a m_s$	$a m_{ud}$ ( $M_\pi$ [MeV])
2.453(4)	$32^3 \times 64 \times 12$	0.0300	0.0070 (300), 0.0120 (400), 0.0190 (500)
		0.0400	0.0035(230), 0.0070(300), 0.0120 (400), 0.0190 (500)
	$48^3 \times 96 \times 12$	0.0400	0.0035 (230)
3.610(9)	$48^3 \times 96 \times 8$	0.0180	0.0042 (300), 0.0080 (400), 0.0120 (500)
		0.0250	0.0042 (300), 0.0080 (400), 0.0120 (500)
4.496(9)	$64^3 \times 128 \times 8$	0.0150	0.0030 (300)

# Discretization effect

- Subtraction of tree-level discretization effect

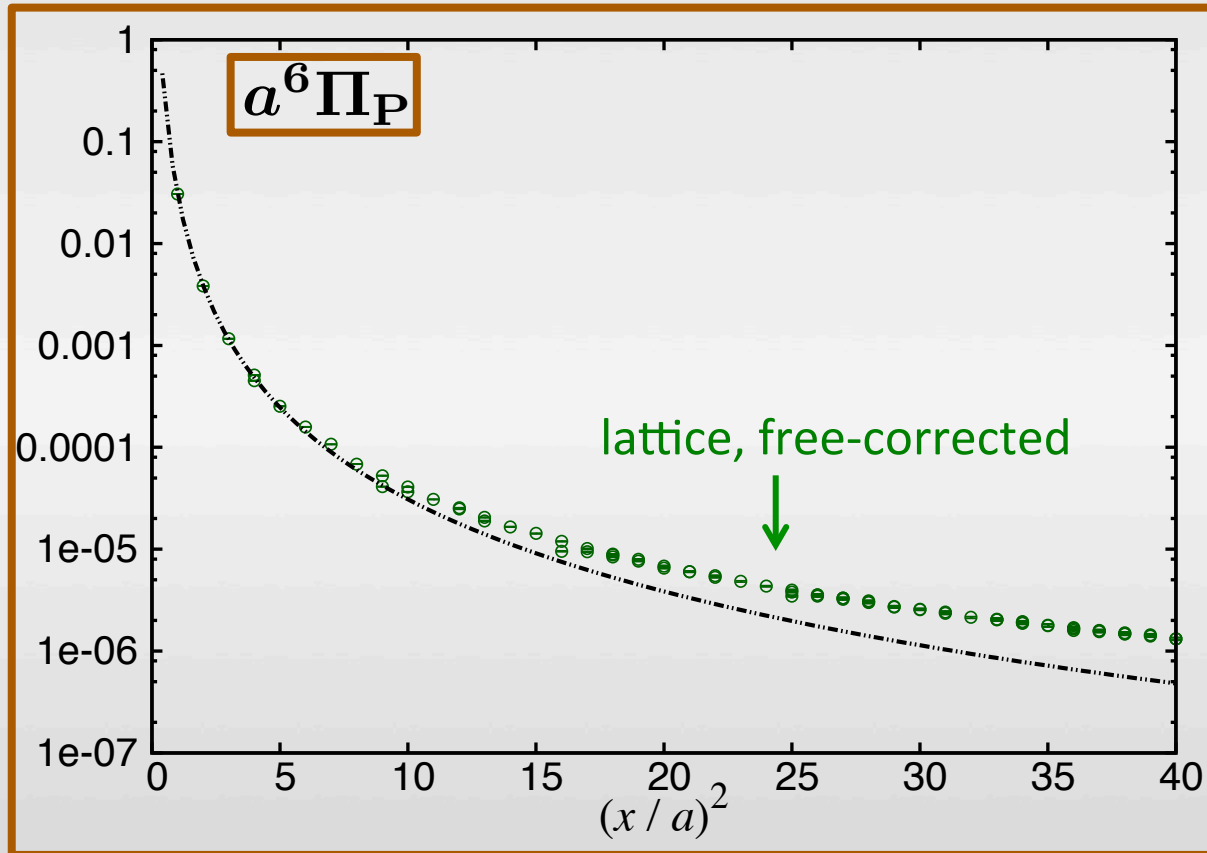


✦ Replace

$$\Pi_{\Gamma}^{lat}(x) \longrightarrow \Pi_{\Gamma}^{lat}(x) - \left( \Pi_{\Gamma}^{lat,free}(x) - \Pi_{\Gamma}^{cont,free}(x) \right)$$

# Discretization effect

- Subtraction of tree-level discretization effect



✦ Replace

$$\Pi_{\Gamma}^{lat}(x) \longrightarrow \Pi_{\Gamma}^{lat}(x) - \left( \Pi_{\Gamma}^{lat,free}(x) - \Pi_{\Gamma}^{cont,free}(x) \right)$$



# Continuum Perturbation

- Perturbative series

$$\Pi = c_0 + c_1 a_s(\mu) + c_2 a_s(\mu)^2 + \dots$$

$$a_s(\mu) = a_s(\mu^*) + O(a_s(\mu^*)^2)$$

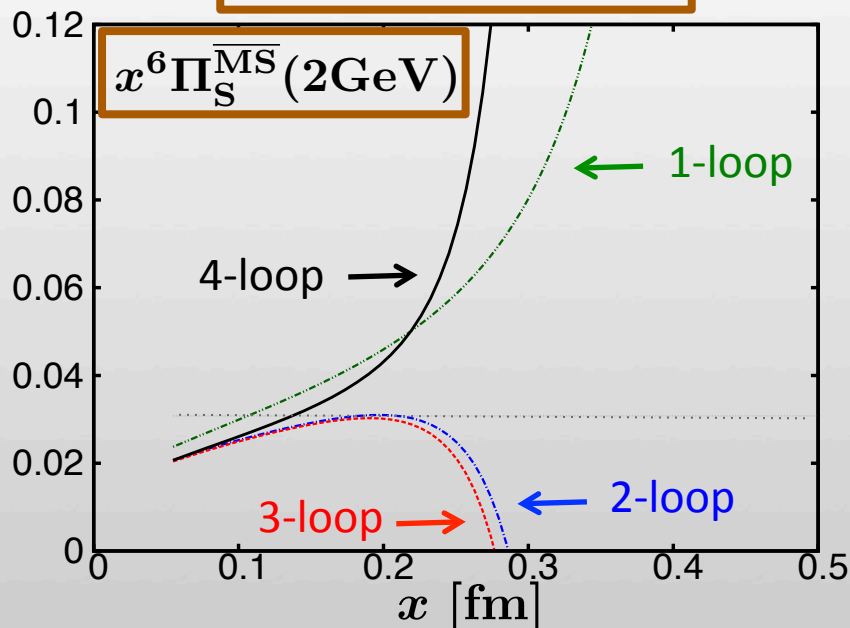
$$= c_0 + c_1^* a_s(\mu^*) + c_2^* a_s(\mu^*)^2 + \dots \leftarrow \text{good convergence}$$

Ref: Brodsky et al, 1983

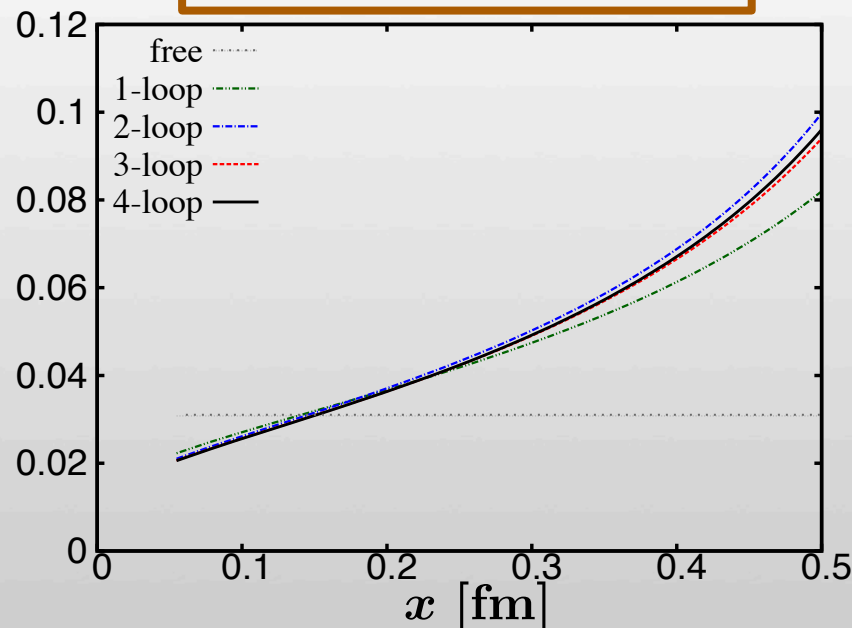
if  $c_i^*$  : small

- $\mu^*$  : arbitrary

Coupling scale  $\mu = x^{-1}$



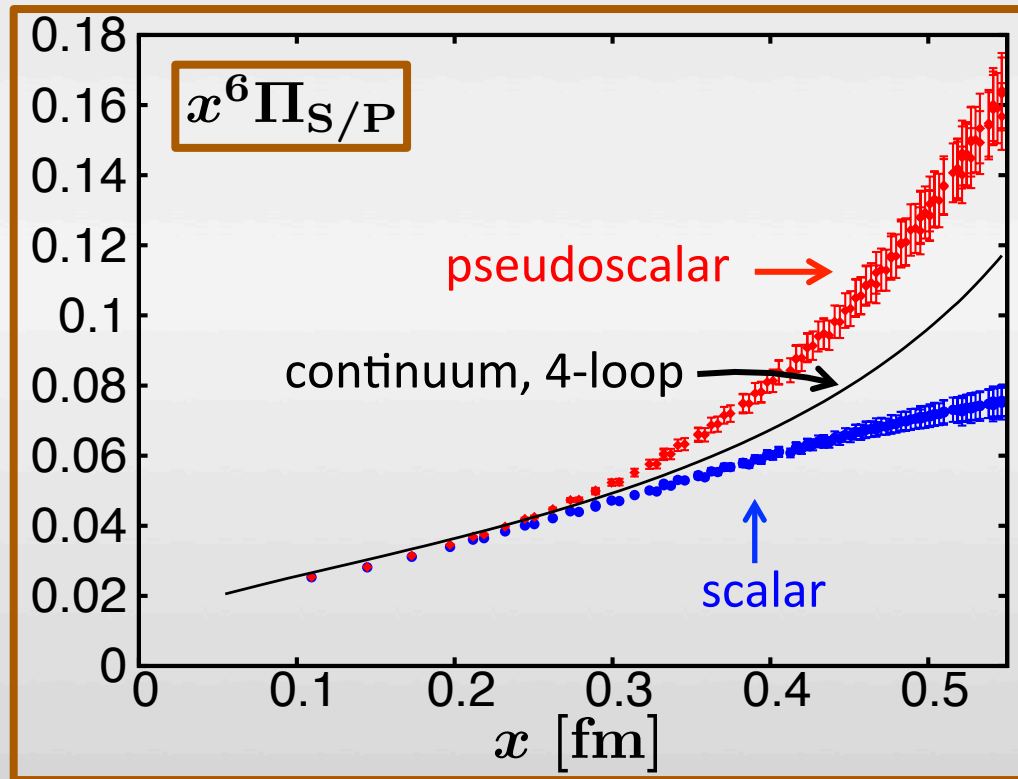
Optimal coupling scale  $\mu^*$



# Lattice V.S. Continuum

- Comparison

- ✦ Lattice :  $a^{-1} = 3.61 \text{ GeV}$ ,  $M\pi = 300 \text{ MeV}$ , RC is multiplied
- ✦ Continuum :  $\overline{\text{MS}}$  2GeV, 4-loop



- Non-perturbative effect  $\rightarrow$  Gap between 2 channels  
 $\Rightarrow$  Check consistency with OPE

# OPE — (Axial-)Vector Channel

- Momentum space

$$\begin{aligned}\Pi_{\Gamma,\mu\nu}(q) &= \int dx e^{-iqx} \Pi_{\Gamma,\mu\nu}(x) \\ &= (\delta_{\mu\nu}q^2 - q_\mu q_\nu) \Pi_{\Gamma}^{(1)}(q^2) - q_\mu q_\nu \Pi_{\Gamma}^{(0)}(q^2) \\ &= \underbrace{(\delta_{\mu\nu}q^2 - q_\mu q_\nu) \Pi_{\Gamma}^{(1+0)}(q^2)}_{\text{Conserved}} - \underbrace{\delta_{\mu\nu}q^2 \Pi_{\Gamma}^{(0)}(q^2)}_{\text{Non-conserved}}\end{aligned}$$

$$\Pi_{\Gamma}^{(1+0)}(q^2) \equiv \Pi_{\Gamma}^{(1)}(q^2) + \Pi_{\Gamma}^{(0)}(q^2)$$

- Non-conserved part

$$q_\mu \Pi_{\Gamma,\mu\nu}(q) = q_\nu q^2 \Pi_{\Gamma}^{(0)}(q^2)$$

$$\Pi_{\mathbf{V}}^{(0)}(q^2) = 0$$

$$\Pi_{\mathbf{A}}^{(0)}(q^2) = \frac{4m_q \langle \bar{q}q \rangle}{q^4} + O(m_q^2)$$

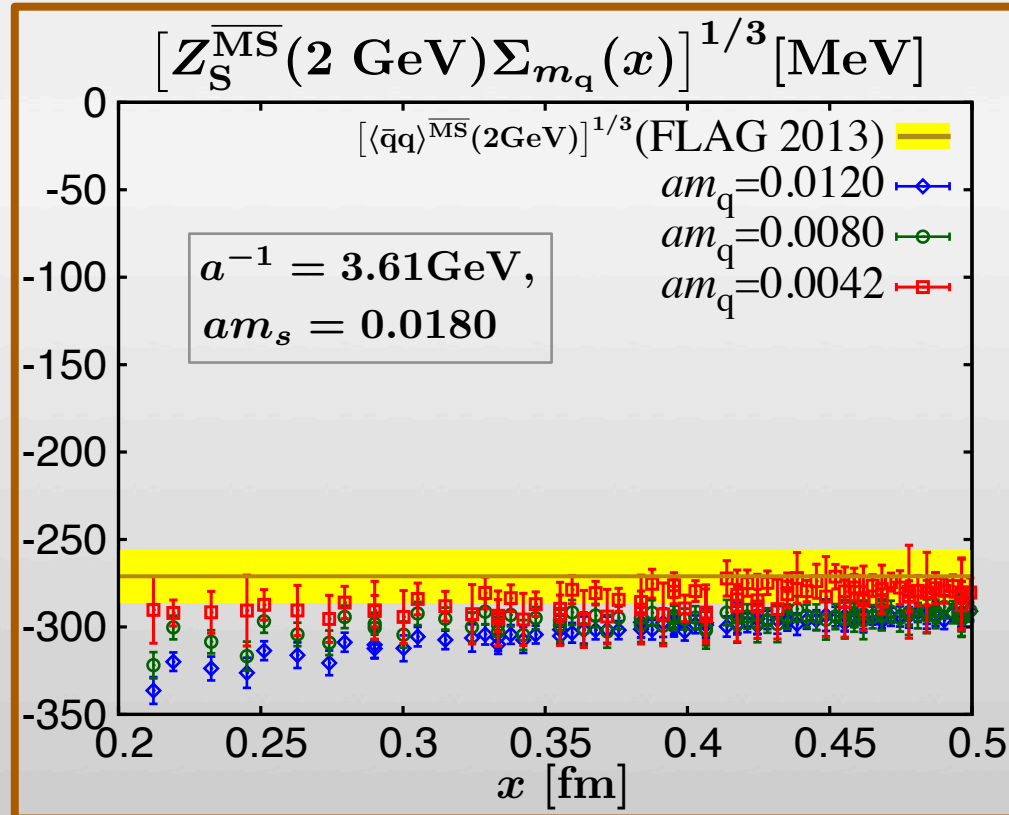
# OPE — (Axial-)Vector Channel

- Non-conserved part in X-space

$$\Sigma_{m_q}(x) \equiv -\frac{\pi^2}{2m_q} x^2 x_\nu \partial_\mu \Pi_{A-V, \mu\nu}(x) = \langle \bar{q}q \rangle + O(m_q) \cdot O(x^{-2})$$

$$\Pi_{A-V, \mu\nu}(x) \equiv \Pi_{A, \mu\nu}(x) - \Pi_{V, \mu\nu}(x)$$

- Very close to FLAG's value of  $\langle \bar{q}q \rangle = [-271(15) \text{ MeV}]^3$
- **Vector & Axial-vector :**  
**consistent with OPE**
- How about scalar & pseudo-scalar?



# OPE — (Pseudo-)Scalar Channel

● Gap :  $\Pi_{\Gamma-\Gamma'}(x) \equiv \Pi_{\Gamma}(x) - \Pi_{\Gamma'}(x)$

● OPE :  $\Pi_{P-S} \sim 0.5 \Pi_{V-A}$

● Lattice :  $\Pi_{P-S} \gg \Pi_{V-A}$

P - S : too large

● Known on quenched lattices

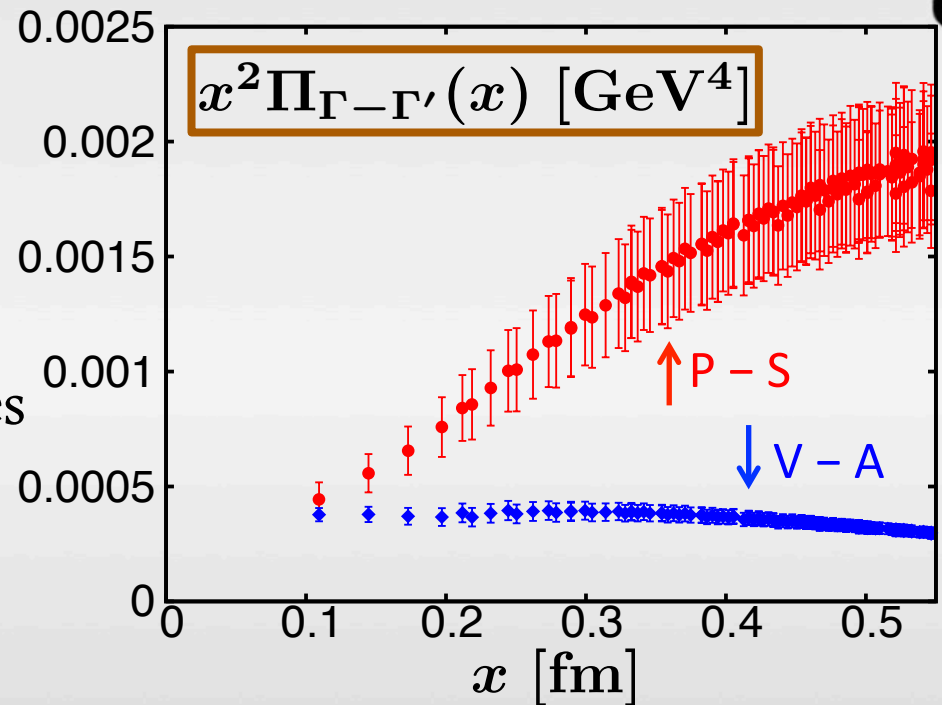
Ref: Chu et al, 1993;

Faccioli, DeGrand, 2003

● Not understood

● Maybe related to  $U(1)_A$  anomaly

● Detailed study underway



# Renormalization in X-space

- Renormalization to  $\overline{\text{MS}}$  scheme

$$\mathcal{O}^{lat}(1/a) \rightarrow \mathcal{O}^{\overline{\text{MS}}}(\mu) = Z^{\overline{\text{MS}}/lat}(1/a \rightarrow \mu) \mathcal{O}^{lat}(1/a)$$

- Renormalization condition in X-space method  
(Martinelli et al '97, Giménez et al 2004, Cichy et al 2012)

$$Z_{\Gamma}^{\overline{\text{MS}}/lat}(\mu)^2 \Pi_{\Gamma}^{lat}(x) = \Pi_{\Gamma}^{\overline{\text{MS}}}(\mu; x) \Rightarrow Z_{\Gamma}^{\overline{\text{MS}}/lat}(\mu) = \sqrt{\frac{\Pi_{\Gamma}^{\overline{\text{MS}}}(\mu; x)}{\Pi_{\Gamma}^{lat}(x)}}$$

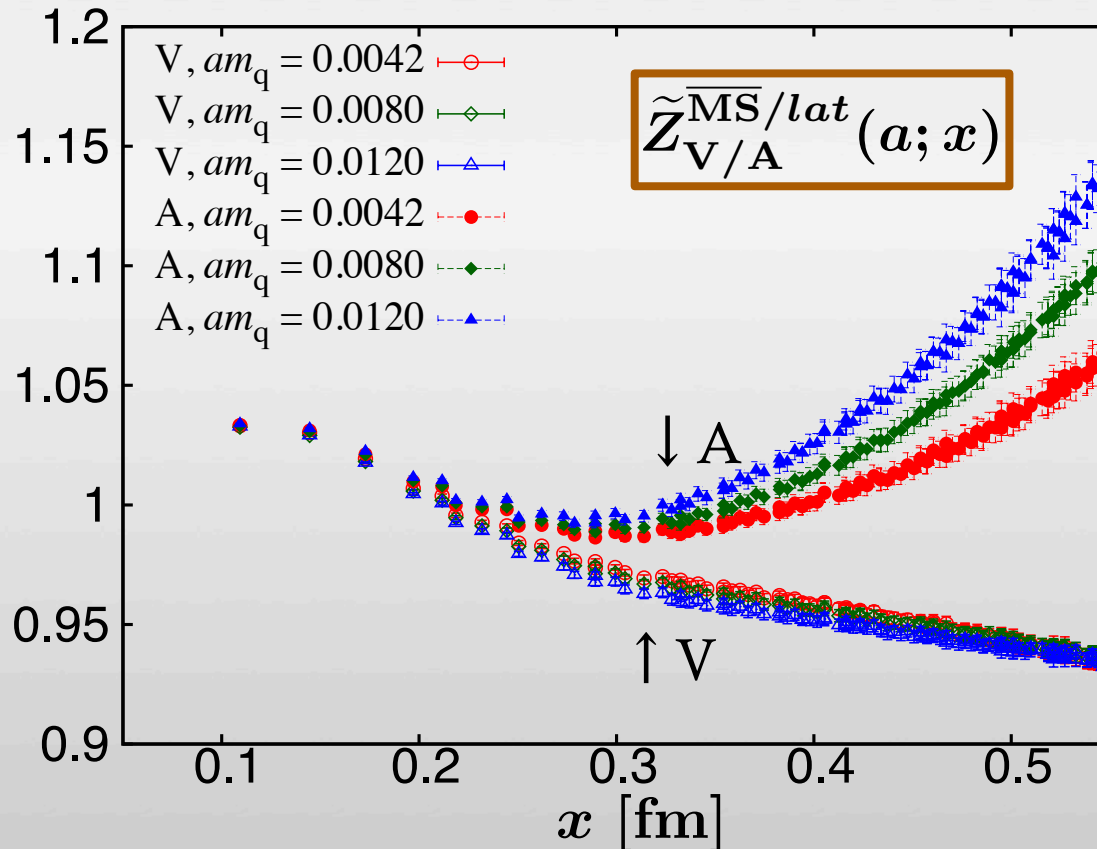
- Advantage

- Renormalization by **gauge invariant** quantities
- We know  $\Pi_{\Gamma}^{\overline{\text{MS}}}(\mu; x)$  to **4-loop** level (Chetyrkin, Maier, 2011)
- Window problem : we need to extract  $Z_{\Gamma}$  from  $a \ll x \ll \Lambda_{\text{QCD}}^{-1}$   
in order to avoid ① discretization effect, ② non-perturbative effect

# Renormalization of Vector Current

- $Z_V^{\overline{\text{MS}}} = Z_A^{\overline{\text{MS}}}$  : scale independent

- $x$ -dependence of  $\sqrt{\frac{\Pi_V^{\overline{\text{MS}}}(x)}{\Pi_{V/A}^{\text{lat}}(x)}} \equiv \tilde{Z}_{V/A}^{\overline{\text{MS}}/\text{lat}}(a; x)$  : not a constant



# Renormalization of Vector Current

- OPE

$$\tilde{Z}_{V/A}^{\overline{\text{MS}}/lat}(a; x) = Z_V^{\overline{\text{MS}}/lat}(a) + \cancel{c_2^{V/A} m_q^2 x^2} + (c_{4,\bar{q}q}^{V/A} m_q \langle \bar{q}q \rangle + c_{4,G}^{V/A} \langle GG \rangle + \cancel{c_{4,m}^{V/A} m_q^4}) x^4 + \dots$$

Too small  $\Rightarrow$  neglected

$$\underline{c_{4,\bar{q}q}^V / c_{4,\bar{q}q}^A = -3/5}$$

$$\begin{aligned} \Rightarrow \tilde{Z}_{(5V+3A)/8}^{\overline{\text{MS}}/lat}(a; x) &\equiv \sqrt{\frac{\Pi_V^{\overline{\text{MS}}}(x)}{(5\Pi_V^{lat}(x) + 3\Pi_A^{lat}(x))/8}} \\ &= Z_V^{\overline{\text{MS}}/lat}(a) + 0 \cdot m_q \langle \bar{q}q \rangle x^4 + c_{4,G} \langle GG \rangle x^4 + \dots \end{aligned}$$

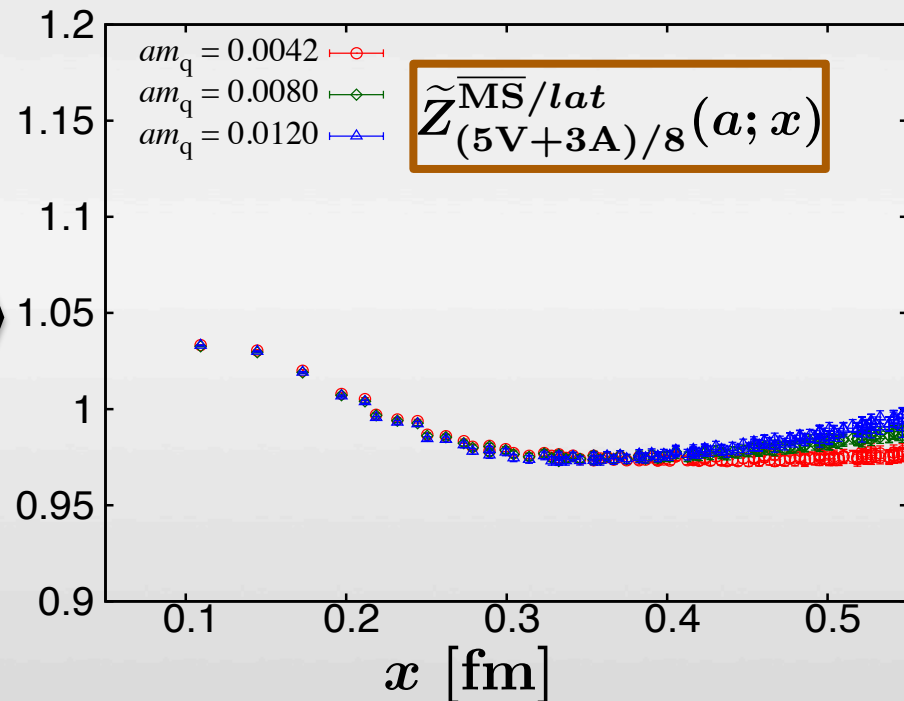
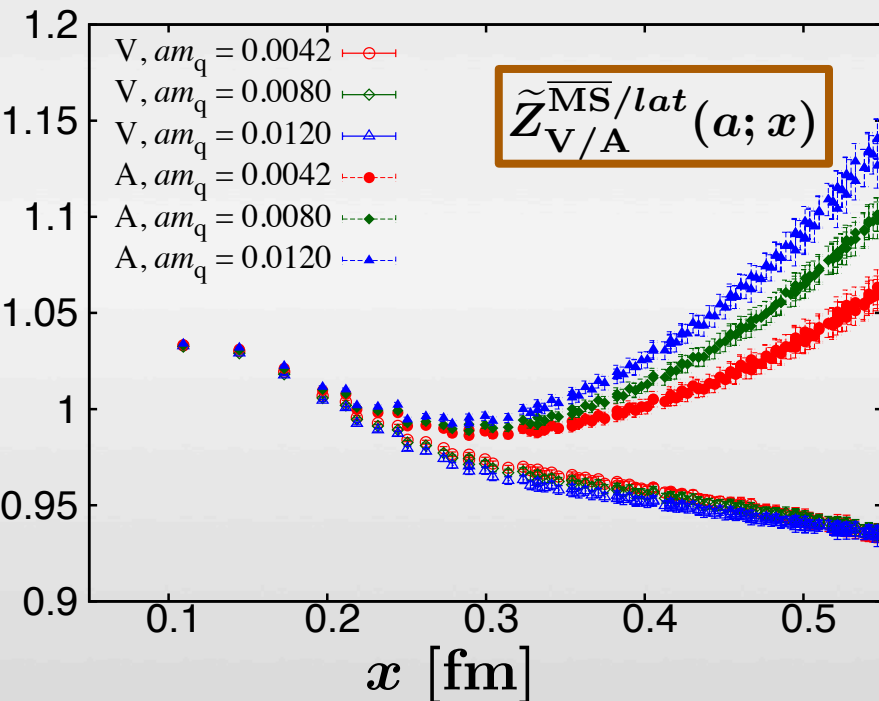
- Mass dependence in short-range can be eliminated



# Renormalization of Vector Current

- New combination

$$\tilde{Z}_{(5V+3A)/8}^{\overline{\text{MS}}/lat}(a; x) = Z_V^{\overline{\text{MS}}/lat}(a) + c_{4,G} \langle GG \rangle x^4 + \dots$$



- Mass dependence become much smaller

# Renormalization of Vector Current

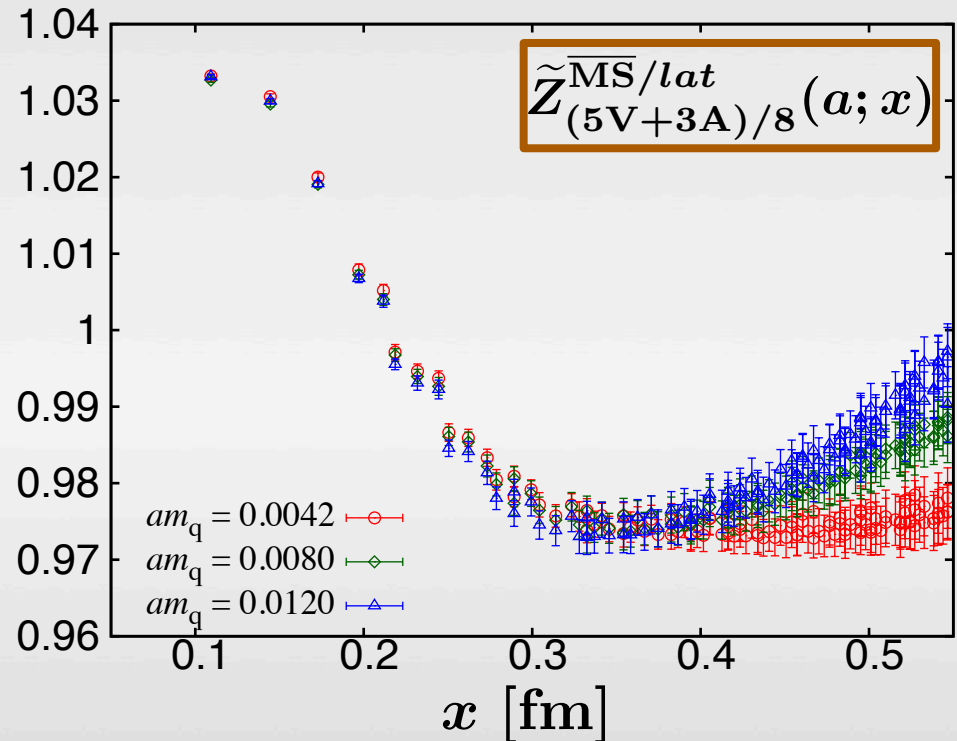
- Up to  $x^4$ , data are independent of mass

$\Rightarrow O(x^6)$  : necessary

at  $x > 0.4$  fm

$$c_{6,\bar{q}q^2} \langle \bar{q}q \rangle x^6 + c_{6,m_G} m_q^2 \langle GG \rangle x^6 + \dots$$

Zoomed version



# Renormalization of Vector Current

- Up to  $x^4$ , data are independent of mass

$\Rightarrow O(x^6)$  : necessary

at  $x > 0.4$  fm

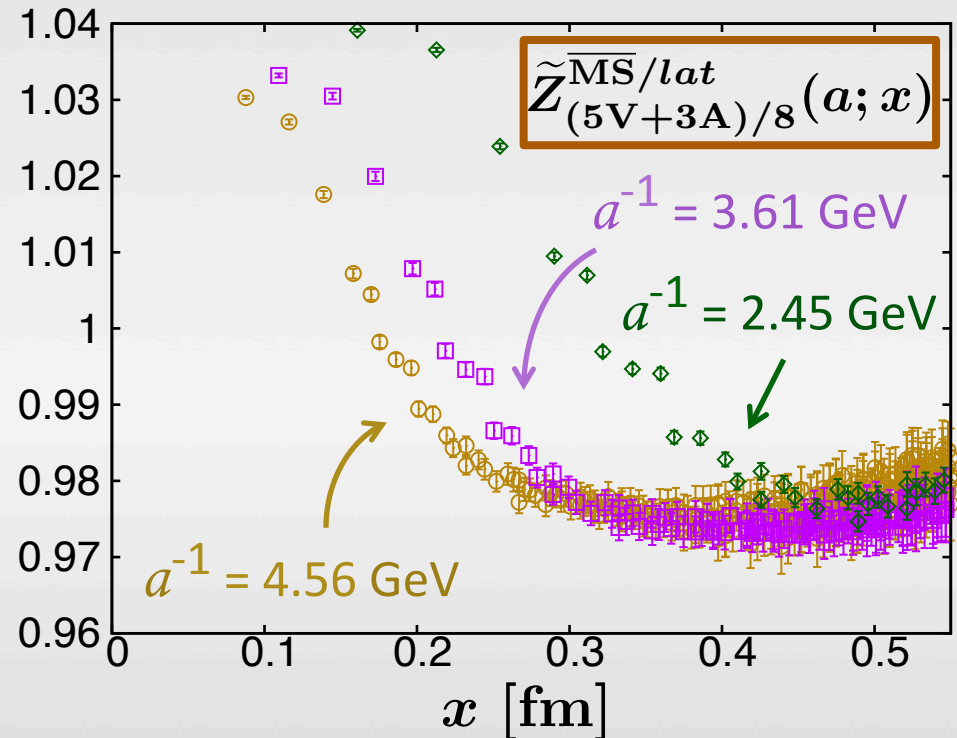
$$c_{6,\bar{q}q^2} \langle \bar{q}q \rangle x^6 + c_{6,m_G} m_q^2 \langle GG \rangle x^6 + \dots$$

- Increasing behavior in short range depends on  $a$

$\rightarrow$  Discretization effect  $(a/x)^2, (am_q)^2, (a/x)^4, \dots$

Significant

Result for each  $a$  at  $M_\pi = 300$  MeV



# Renormalization of Vector Current

- Fit function:

$$\tilde{Z}_{(5V+3A)/8}^{\overline{\text{MS}}/lat}(a; x) = Z_V^{\overline{\text{MS}}/lat}(a) + c_{-2}(a/x)^2 + c_4 x^4 + (c_6 + c'_6 m_q^2) x^6$$

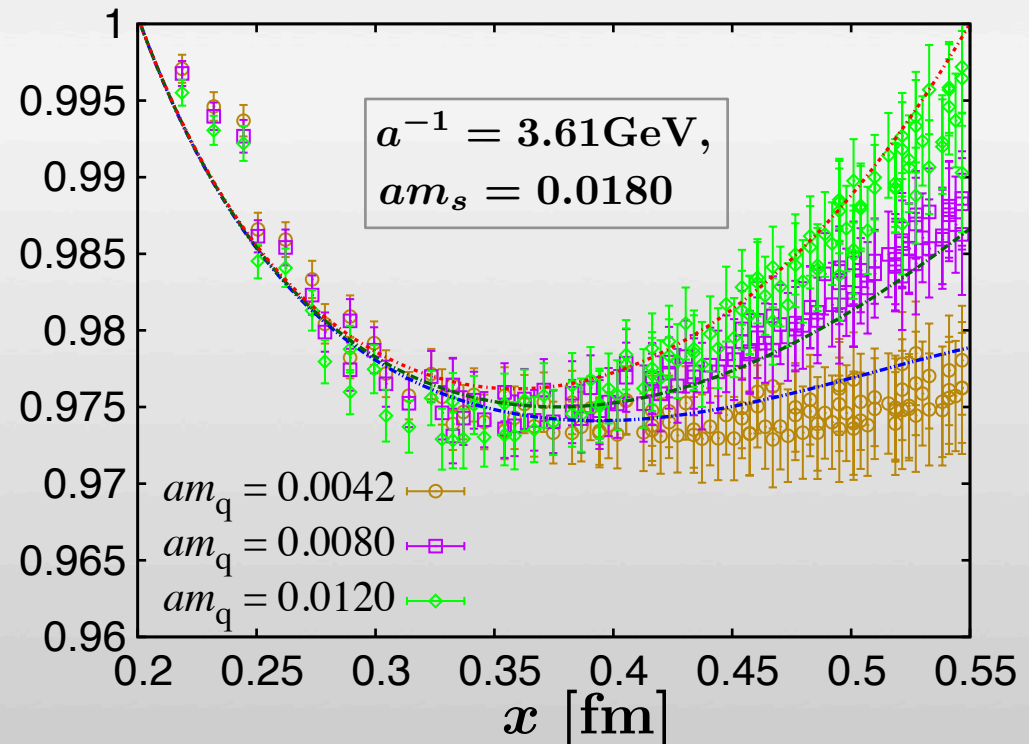
- $\chi^2$  fitting simultaneously among each ensembles

- Preliminary results

$$Z_V^{\overline{\text{MS}}/lat}(\beta = 4.17) = 0.951(4),$$

$$Z_V^{\overline{\text{MS}}/lat}(\beta = 4.35) = 0.956(3),$$

$$Z_V^{\overline{\text{MS}}/lat}(\beta = 4.47) = 0.961(3)$$



# Renormalization of Scalar Density

- Good combination of  $\Pi_S$  &  $\Pi_P$

Naïve average :  $\Pi_{(S+P)/2} = \frac{1}{2}(\Pi_S + \Pi_P)$  ←  $U(1)_A$ -singlet

- Subtracting remaining non-perturbative effect

$$\Pi_{(S+P)/2} = C_{pert}x^{-6} + (C_{4,\bar{q}q}m_q\langle\bar{q}q\rangle + C_{4,G}\langle GG\rangle)x^{-2} + \dots$$

Reduce number of fit parameters  
using

$$\Pi_{V-A} = -16C_{4,\bar{q}q}m_q\langle\bar{q}q\rangle x^{-2} + \dots$$

- ✦  $\Pi_{(S+P)/2+(V-A)/16}$  : independent of  $m_q \langle\bar{q}q\rangle x^{-2}$

# Renormalization of Scalar Density

- Fit function:

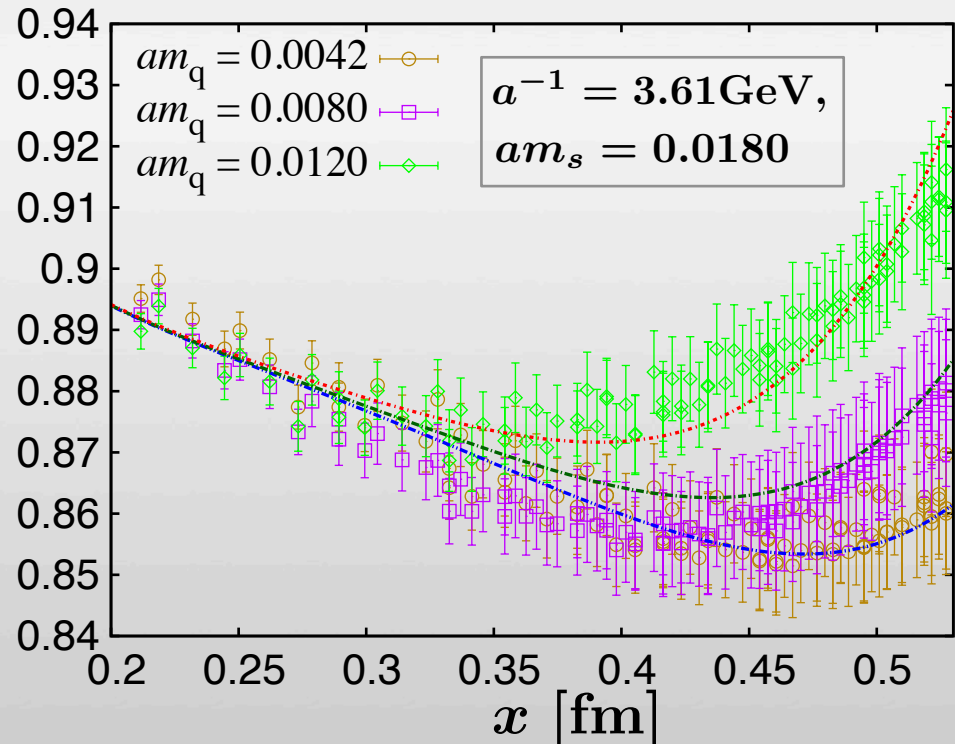
$$\begin{aligned} \tilde{Z}_{(S+P)/2+(V-A)/16}^{\overline{\text{MS}}/lat}(2 \text{ GeV}; a; x) &= Z_S^{\overline{\text{MS}}/lat}(2 \text{ GeV}; a) \\ &+ c_{-2}(a/x)^2 + c_4 x^4 + (c_6 + c'_6 m_q^2) x^6 \end{aligned}$$

- Preliminary results for  $Z_S^{\overline{\text{MS}}/lat}(2 \text{ GeV}; a)$

$$Z_S^{\overline{\text{MS}}/lat}(\beta = 4.17) = 0.975(9),$$

$$Z_S^{\overline{\text{MS}}/lat}(\beta = 4.35) = 0.881(7),$$

$$Z_S^{\overline{\text{MS}}/lat}(\beta = 4.47) = 0.842(5)$$



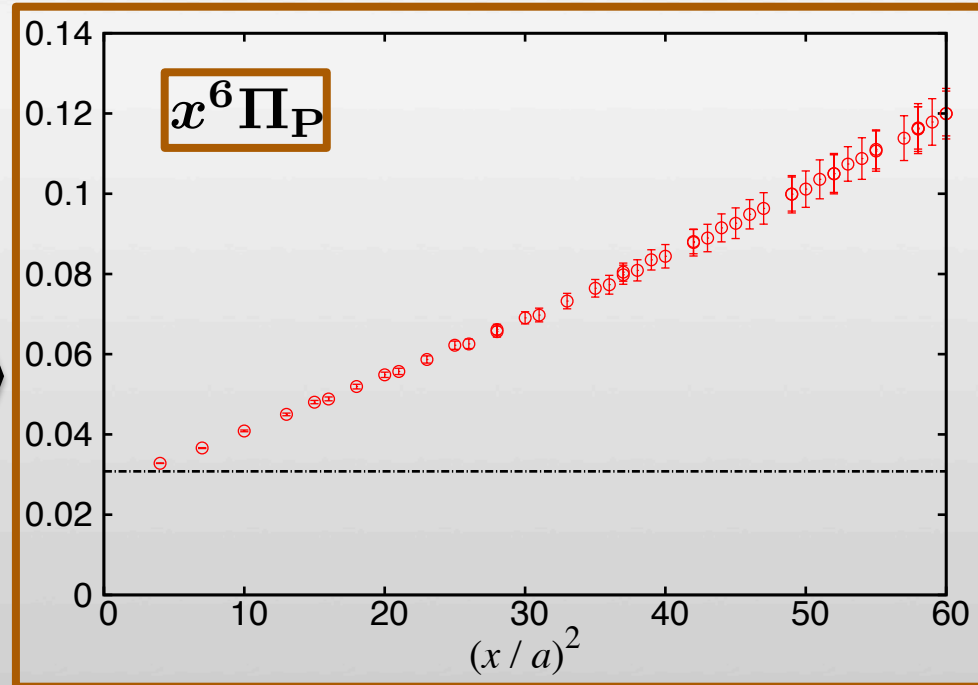
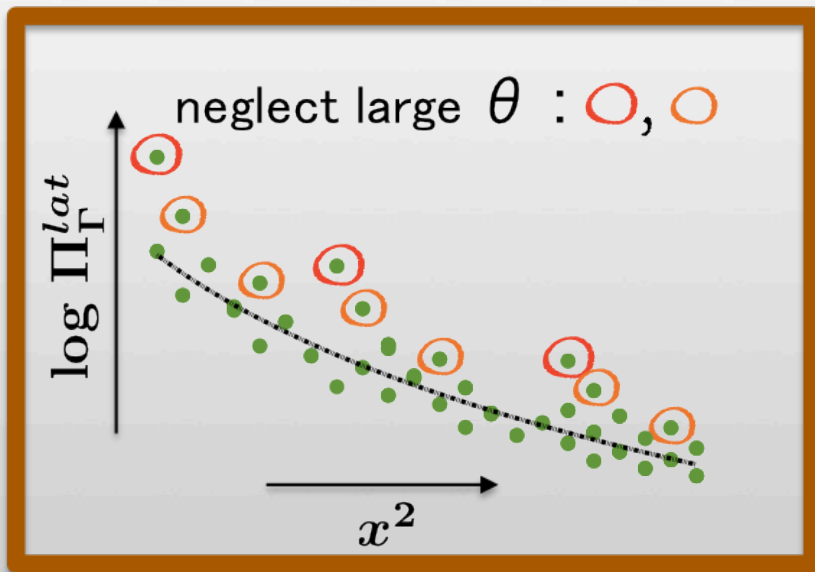
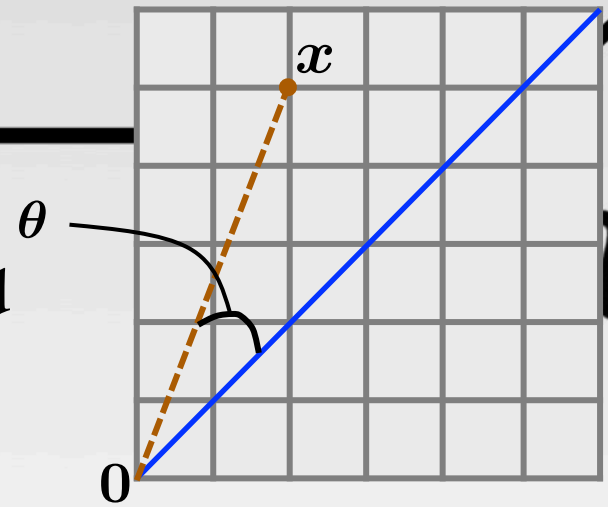
# Summary

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- We investigate two-point correlation functions for scalar and pseudoscalar densities, vector and axial-vector currents.
- Convergence of perturbative correlators is well improved by changing scale of coupling.
- Results of vector and axial-vector currents have good agreement with OPE.
- Chiral condensate is extracted from non-conserved part of the correlator of axial-vector current.
- Using the context of OPE, renormalization constants are determined with good precision ( $\lesssim 1\%$ ).

# Diagonal Cut

- $\theta$  : angle between  $x$  &  $(1,1,1,1)$
- Correlators at large  $\theta$  are more distorted
- Free correlators at small  $\theta$  are closer to those in continuum theory
- Data with  $\theta < 30^\circ$  are used





# Renormalization of Vector Current

- Fit function:

$$\tilde{Z}_{(5V+3A)/8}^{\overline{\text{MS}}/lat}(a; x) = Z_V^{\overline{\text{MS}}/lat}(a) + c_{-2}(a/x)^2 + c_4 x^4 + (c_6 + c'_6 m_q^2) x^6$$

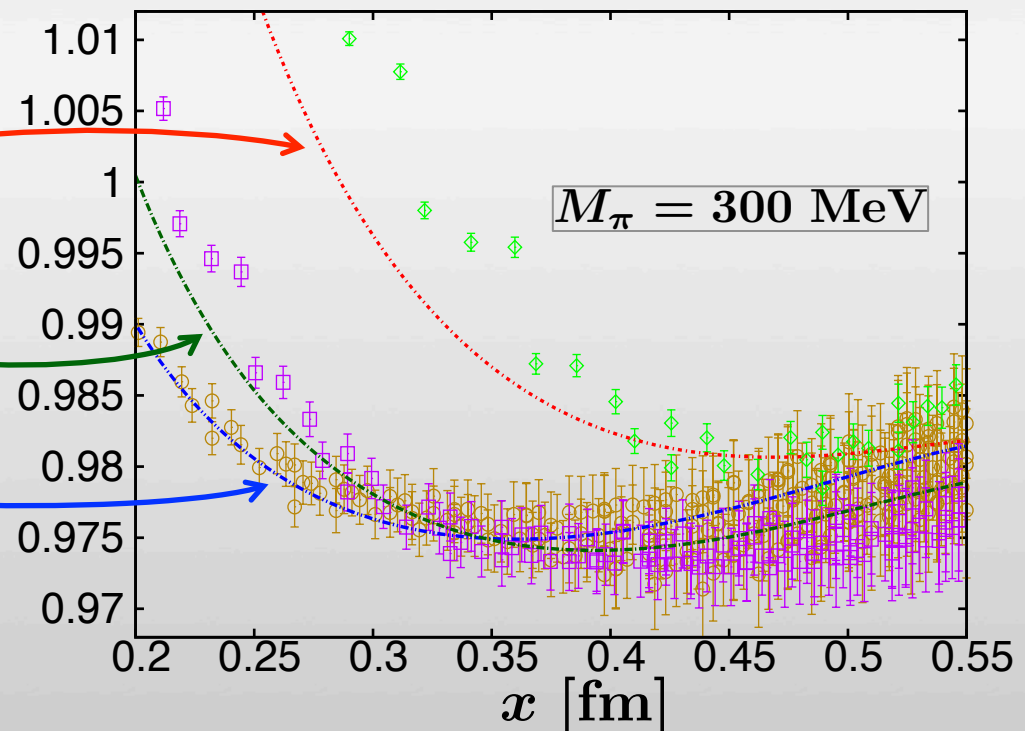
- Perform chi 2 fitting simultaneously for all ensembles

- Preliminary results

$$Z_V^{\overline{\text{MS}}/lat}(2.45 \text{ GeV}) = 0.951(4),$$

$$Z_V^{\overline{\text{MS}}/lat}(3.61 \text{ GeV}) = 0.956(3),$$

$$Z_V^{\overline{\text{MS}}/lat}(4.50 \text{ GeV}) = 0.961(3)$$



# Renormalization of Scalar Density

- Fit function:

$$\tilde{Z}_{(S+P)/2+(V-A)/16}^{\overline{\text{MS}}/lat}(2 \text{ GeV}; a; x) = Z_S^{\overline{\text{MS}}/lat}(2 \text{ GeV}; a) + c_{-2}(a/x)^2 + c_4 x^4 + (c_6 + c'_6 m_q^2) x^6$$

- Preliminary results for  $Z_S^{\overline{\text{MS}}/lat}(2 \text{ GeV}; a)$

