

$N_f=2+1+1$ renormalisation of four-quark operators

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Motivations

The story so far

- The Rome-Southampton method with momentum sources has proven to be very efficient
- We have presented last year a $N_f = 2 + 1 + 1$ strategy to treat the charm threshold
- ... and preliminary results for B_K

What more can we do?

- All systematics were not fully under control
- We need generalisation to SUSY B_K and $K \rightarrow \pi\pi$

Operators and chiral sectors

sector	dim	B_K	SUSY B_K	$K \rightarrow \pi\pi$ $I = 2$	$K \rightarrow \pi\pi$ $I = 0$
(27, 1)	1	✓	✓	✓	✓
(8, 8)	2		✓	✓	✓
(8, 1)	4				✓
(6, 6)	2		✓		

With our chiral action, we can consider each chiral sector separately.

$$Q^{(27,1)} = [\bar{s}\gamma_\mu(1 - \gamma_5)d] [\bar{s}\gamma^\mu(1 - \gamma_5)d] \quad (1)$$

$$Q_1^{(8,8)} = [\bar{s}\gamma_\mu(1 - \gamma_5)d] [\bar{s}\gamma^\mu(1 + \gamma_5)d] \quad (2)$$

$$Q_2^{(8,8)} = [\bar{s}(1 - \gamma_5)d] [\bar{s}(1 + \gamma_5)d] \quad (3)$$

$$Q_1^{(6,6)} = [\bar{s}\sigma_{\mu\nu}(1 - \gamma_5)d] [\bar{s}\sigma^{\mu\nu}(1 - \gamma_5)d] \quad (4)$$

$$Q_2^{(6,6)} = [\bar{s}(1 - \gamma_5)d] [\bar{s}(1 - \gamma_5)d] \quad (5)$$

Step-scaling

If one chooses a scaling trajectory and compute at constant physical masses the continuum limit

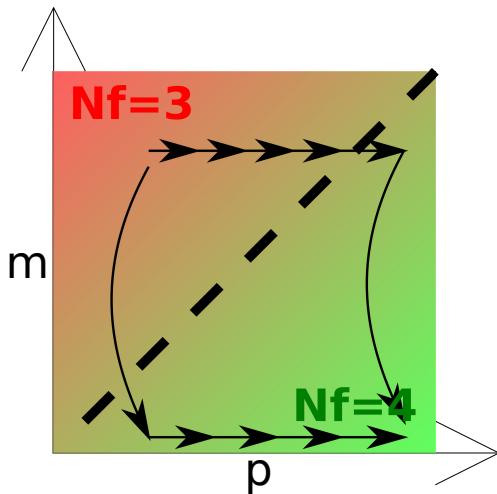
$$\sigma(p, p_0 | m) = \lim_{a \rightarrow 0} \frac{Z(p | m)}{Z(p_0 | m)}, \quad (6)$$

this is universal, i.e. only dependent on the continuum RGEs from one scale to another.

Does not depend on your scaling trajectory, your lattice action, the $O(4)$ orbit of momenta, and so on.

Therefore those ratios are convenient building-blocks that we can combine with as many ratios as we want into a telescopic product.

Charm threshold



Ensembles

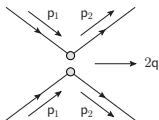
$N_f = 2 + 1$ Ensembles

Matrix elements have been computed on a wide set of (M)DWF ensembles, including two ensembles at the physical quark masses, and lattice spacing going up to 3 GeV.

$N_f = 2 + 1 + 1$ Ensembles

β	$L^3 \times T \times L_5$	m_l	m_c	a^{-1}
5.70	$32^3 \times 64 \times 12$	0.0047	0.243, 0.1, 0.0186	3.0 GeV
5.70	$32^3 \times 64 \times 12$	0.002	0.243	3.0 GeV
5.77	$32^3 \times 64 \times 12$	0.0044	0.213	3.6 GeV
5.84	$32^3 \times 64 \times 12$	0.0041	0.183, 0.0146	4.3 GeV
5.84	$32^3 \times 64 \times 12$	0.002	0.183	4.3 GeV

RI-SMOM scheme



Kinematics

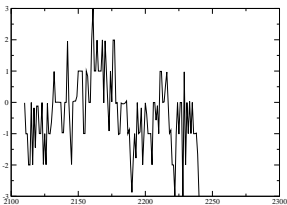
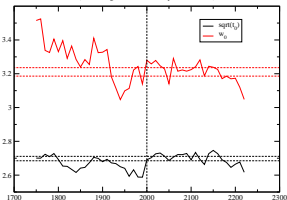
- Non-exceptional schemes avoid π pole
- $p_1^2 = p_2^2 = (p_1 - p_2)^2$
- no $\sum p_i$ combination cancels out
- many orientations satisfy this condition but cont. limit is universal

Renormalisation condition

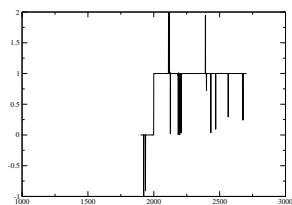
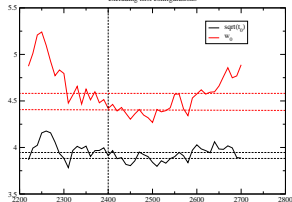
- $Z \text{Tr} [P_{ijkl} G_{ijkl}] = \text{Tr} [P_{ijkl} G_{ijkl}] |_{\text{tree}}$
- $P_{ijkl} = \gamma_i \delta_{ij} \gamma_k \delta_{kl}$ or $P_{ijkl} = \not{q}_{ij} \not{q}_{kl}$
different schemes allow us to evaluate the truncation error
- Very versatile method, with many knobs to turn
- Cheap once we have configs, excellent signal/noise

Wilson flow and topological charge

b5.70_2p2_mob2 flow history
(first configuration is b5.70_2p2, with b-c=5)



b5.84_2p2 flow history
excluding first configurations



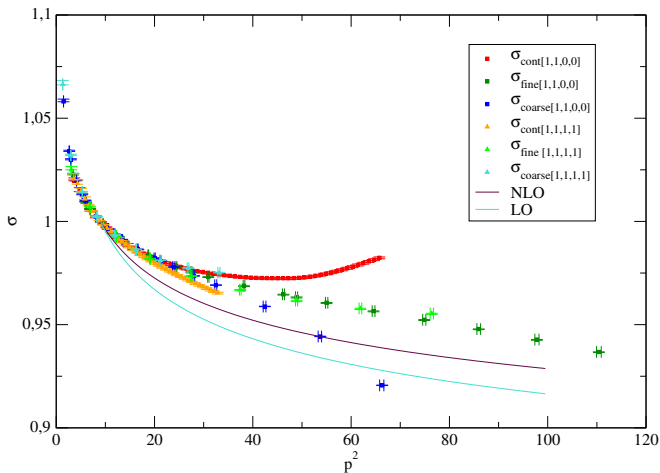
Continuum limit (intro)

The following results are produced through this process:

- 1 Compute step-scaling from 3 GeV for each ensemble (and orientation)
- 2 Interpolate them to a common set of p^2 values
- 3 For each interpolated p^2 and each orientation, extrapolate independently in a^2
- 4 The difference between two orientations is an estimate of the systematics from higher order discretisation terms

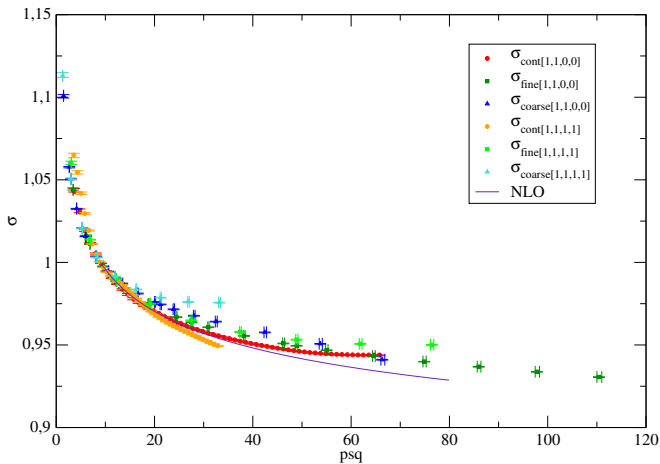
Continuum limit I

$(27,1) \gamma$



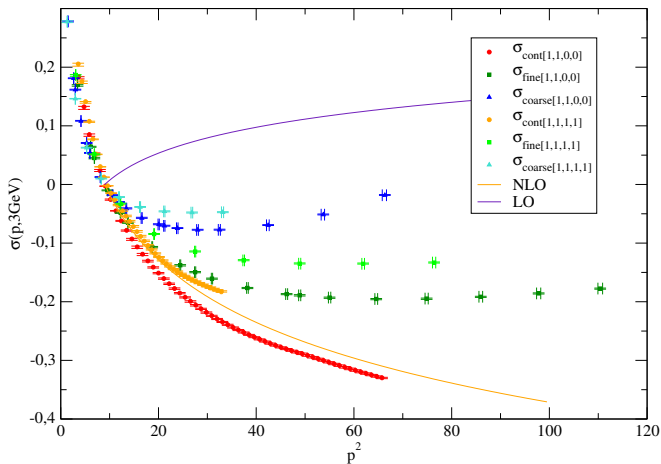
Continuum limit II

(27,1) qq

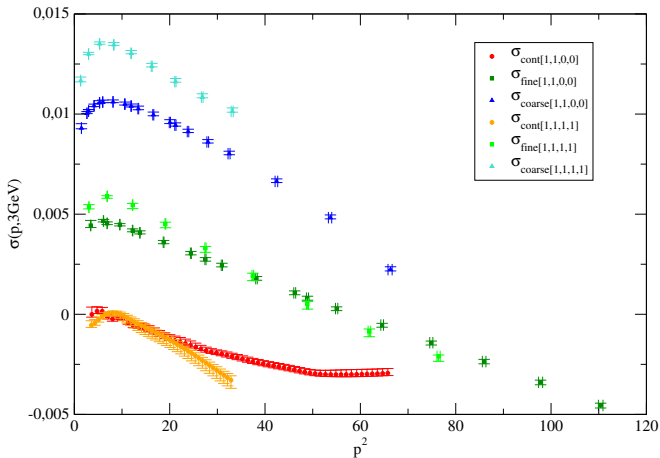


Continuum limit III

$(8,8)_{12} gg$

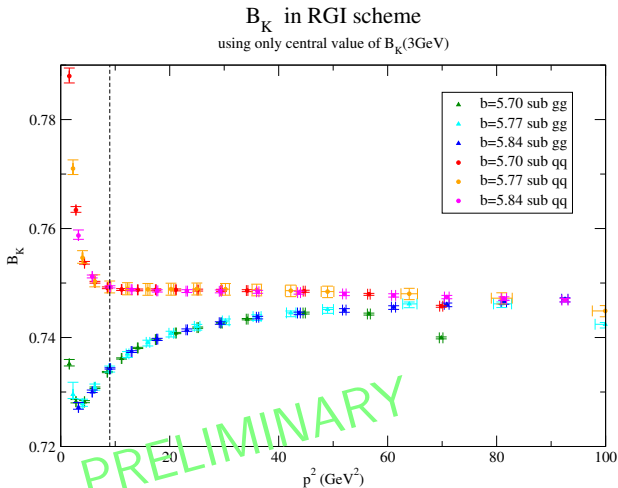


Continuum limit IV

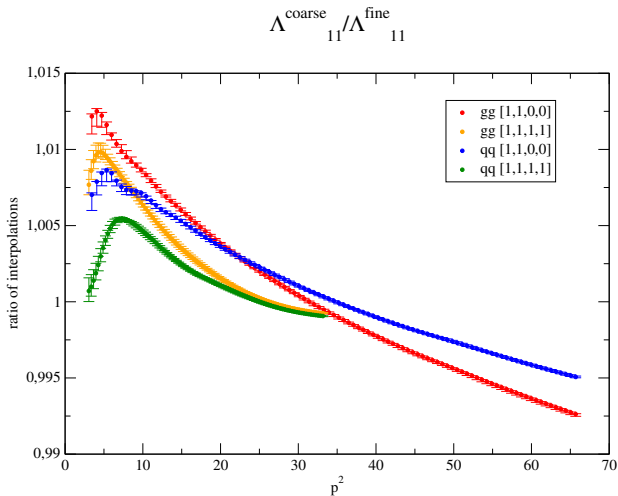
 $(6,6)_{12} \text{ gg}$


Method

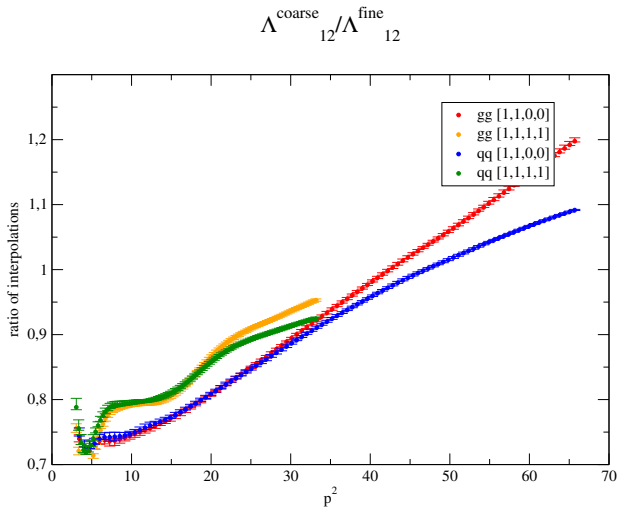
- Because most of the running comes from low-order PT, discretisation effects might be well-described by PT
- However we do not need any explicit perturbative formula
- At high- p^2 , Λ_{QCD} becomes irrelevant, dimensional analysis suggest discr is mostly $(ap)^n$ -dependent.
More precisely: the boosted coupling constant only grows logarithmically (see hep-lat/1412.0834)
- When combined with this all-order theoretical argument, the p^2 -dependence gives additional information
- Of course it doesn't replace a third ensemble, but the latter would be much more expensive than this project

B_K 

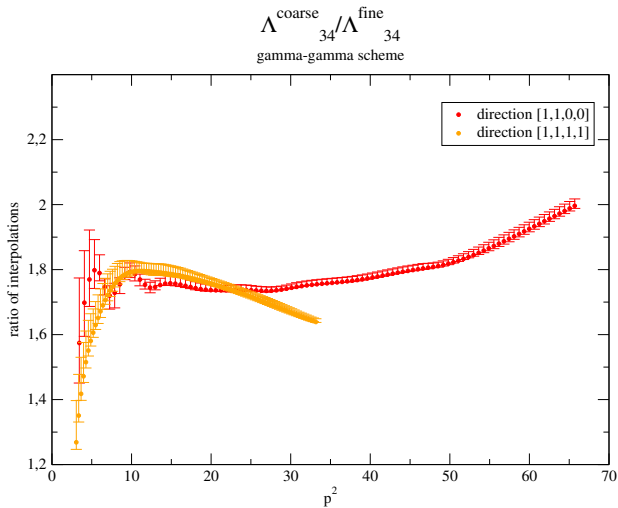
Perspectives for other sectors I



Perspectives for other sectors II

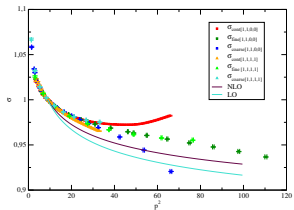
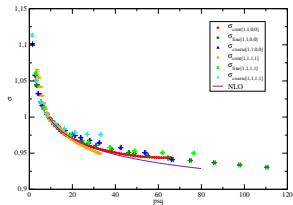


Perspectives for other sectors III

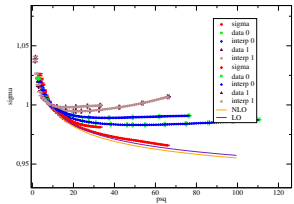


- We have presented preliminary results for the 5×5 SUSY B_K step-scaling
- A tiny “ $N_f = 3/N_f = 4$ ” has not been included here
- With $O(a^2)$ extrapolations the discretisation errors are huge
- A third ensemble would be most useful
- Generating new ensembles, with stats for the low-energy scale-setting, is painful
- It is important to extract as much info as possible from ensembles we already have
- discr. effects in SUSY B_K have a very smooth p^2 -dependence, this looks like a relatively operator-independent observation.

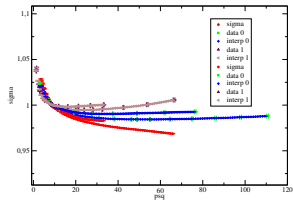
Thanks for your attention!

(27.1) $\gamma\gamma$ (27.1) $q\bar{q}$ 

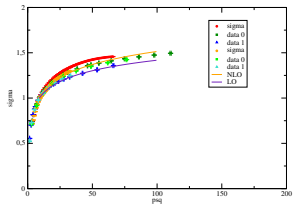
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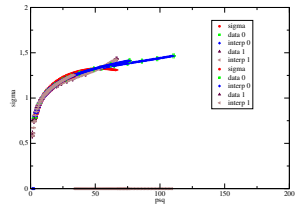
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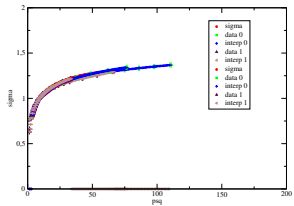
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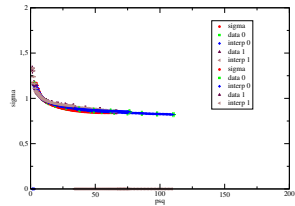
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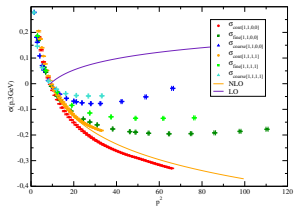
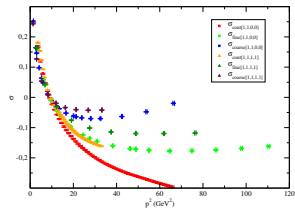


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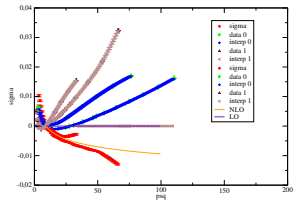


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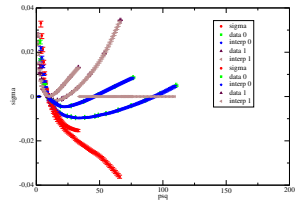


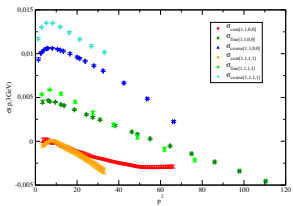
$(8,8)_{1,2} gg$  $(8,8)_{1,2} qq$ 

results/StepCont_GGdir1_2_1_agr



results/StepCont_QQdir1_2_1_agr



$(6.6)_{12}$ gg

results/StepCont_GGdir1_4_3.agr

