

Thermodynamics and reference scale of SU(3) gauge theory from gradient flow on fine lattices

Masakiyo Kitazawa

for FlowQCD Collaboration

(Asakawa, Hatsuda, Iritani, Itou, MK, Suzuki)

FlowQCD, arXiv:1503.06516

LATTICE2015, 2015/Jul./15, Kobe, Japan

Outline

1. Yang-Mills Gradient flow

2. Reference scale and Lattice spacing

FlowQCD, arXiv:1503.06516

3. Thermodynamics

Yang-Mills Gradient Flow

$$\partial_t A_\mu = D_\nu G_{\mu\nu}$$

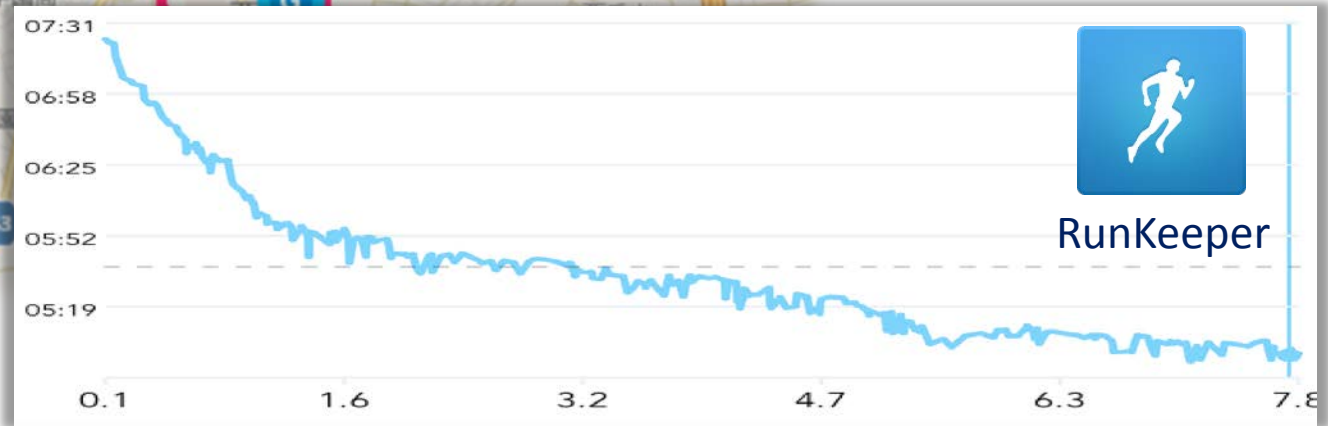
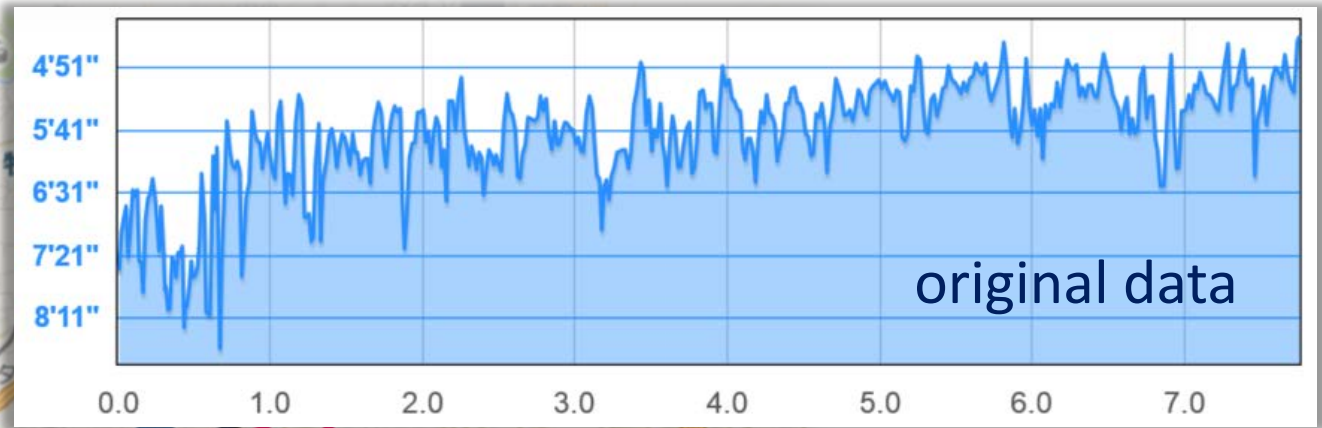
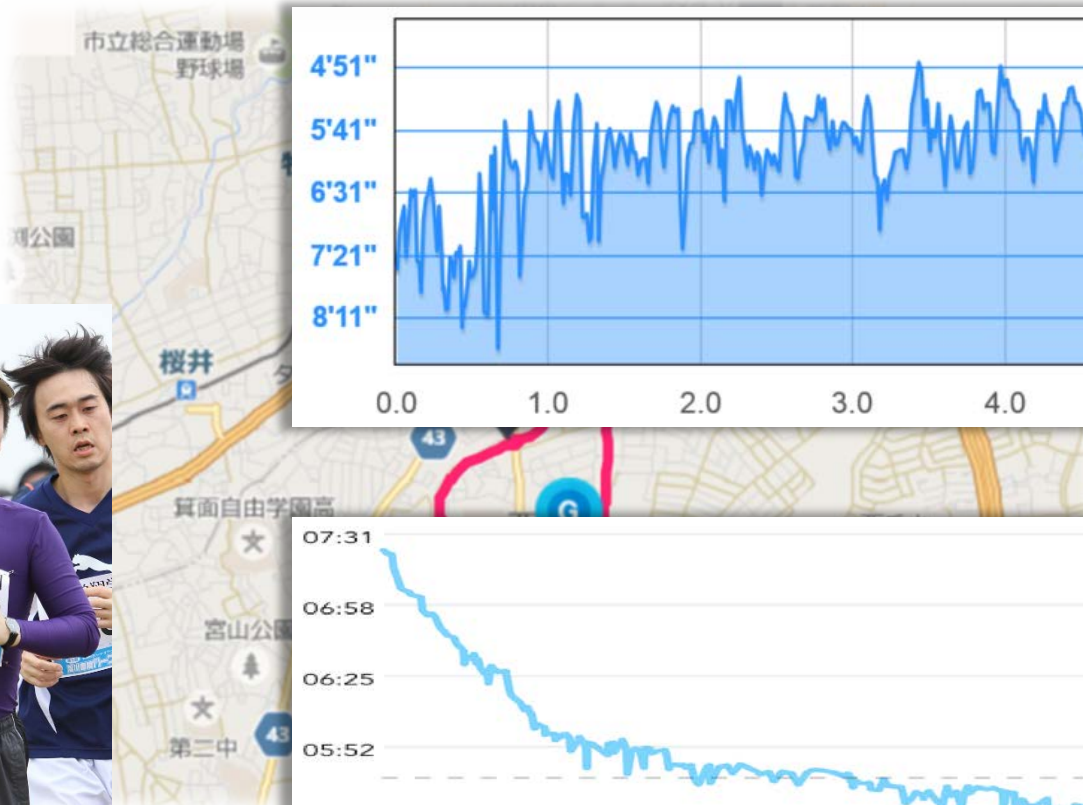
Gradient Flow and Jogging

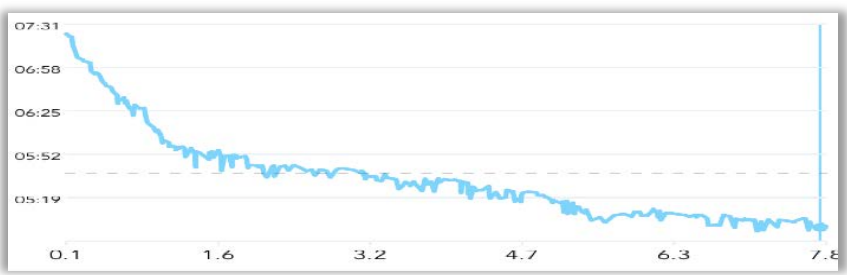
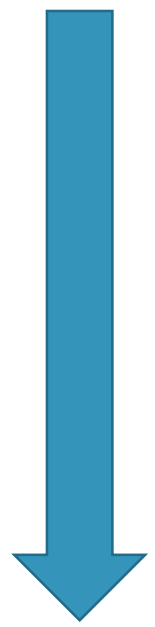
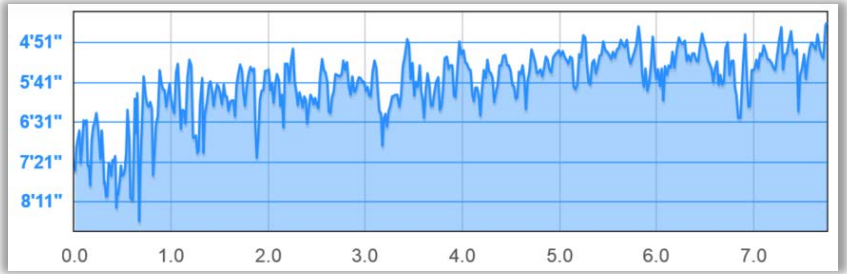


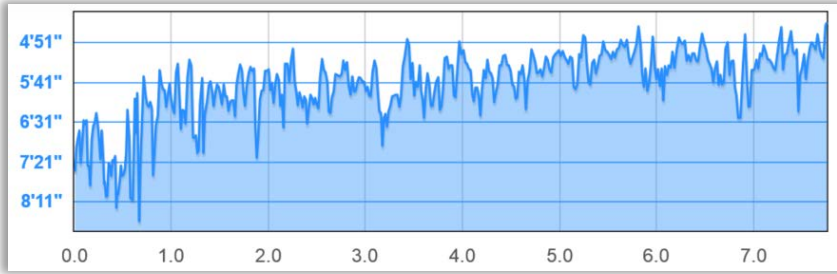
Gradient Flow and Jogging



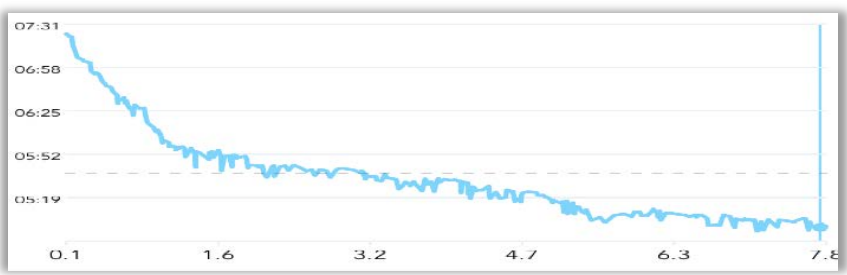
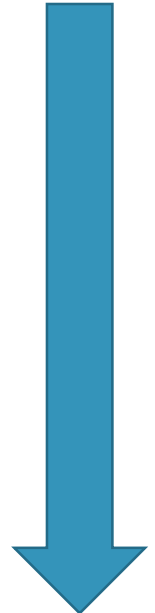
Gradient Flow and Jogging

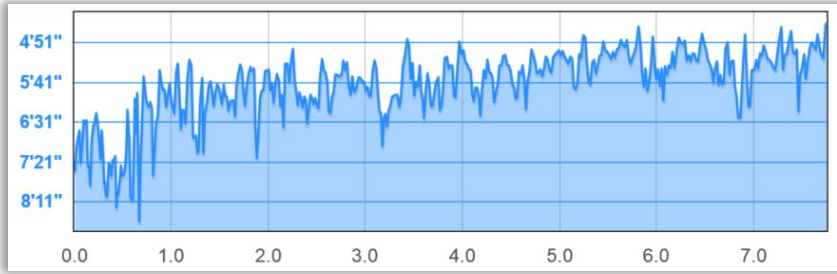






① $x(t) \rightarrow x'(t) \sim \int dt' \exp \left[-\frac{(t-t')^2}{2\sigma^2} \right] x(t')$

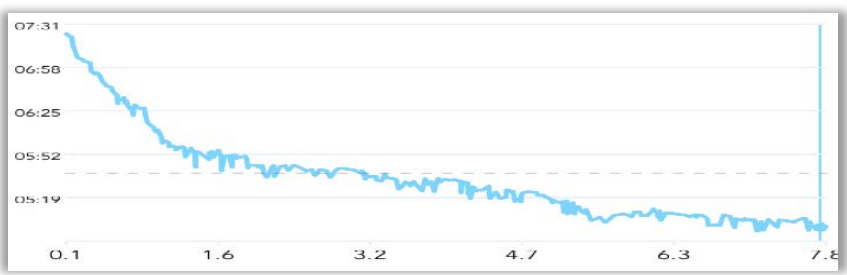


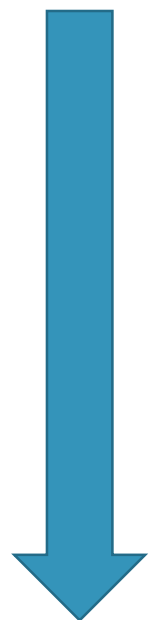
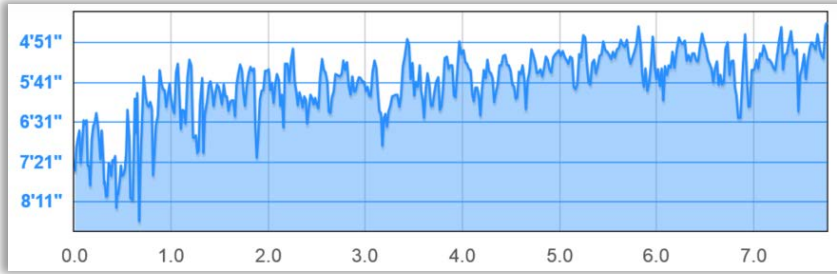


① $x(t) \rightarrow x'(t) \sim \int dt' \exp \left[-\frac{(t-t')^2}{2\sigma^2} \right] x(t')$

$\sigma = \sqrt{2s}$

② $\frac{d}{ds} x(t; s) = \frac{d^2}{dt^2} x(t, s) \quad x(t; 0) = x(t)$

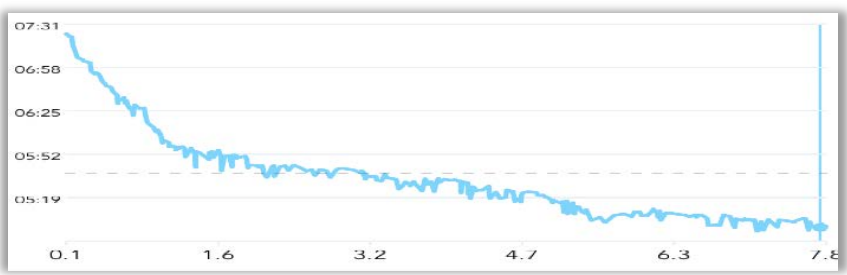




① $x(t) \rightarrow x'(t) \sim \int dt' \exp \left[-\frac{(t-t')^2}{2\sigma^2} \right] x(t')$

$\sigma = \sqrt{2s}$

② $\frac{d}{ds} x(t; s) = \frac{d^2}{dt^2} x(t, s) \quad x(t; 0) = x(t)$



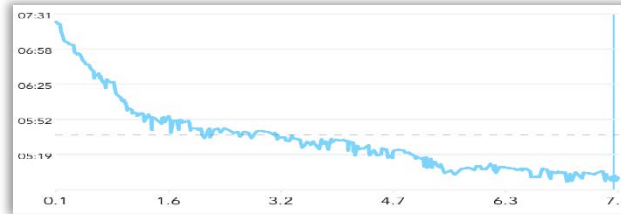
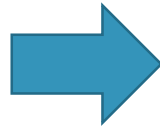
YM Gradient Flow

$$\partial_t A_\mu = D_\nu G_{\mu\nu} = \partial_\nu \partial_\nu A_\mu + \dots$$

- t: "flow time" dim: [E⁻²]
- smearing length $r = \sqrt{8t}$

Use of Gradient Flow 1: Observables at Nonzero Flow Time

$\langle \mathcal{O} \rangle_t$



- function of t
- independent of renormalization

Luescher, Weisz, 2011

independent of lattice spacing
(up to $O(a^2)$ corrections)

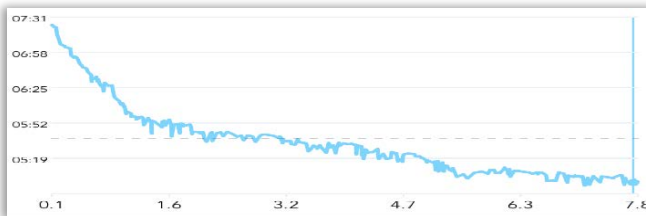
The t dep. can be used for the **scale setting** of the lattice.

Use of Gradient Flow 2: Small Flow Time Expansion

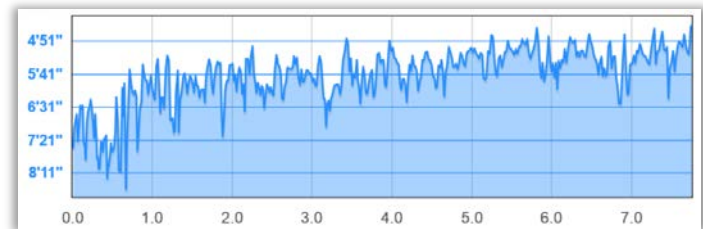
Luescher, Weisz, 2011
Suzuki, 2013

$$\tilde{\mathcal{O}}(t, x) \xrightarrow[t \rightarrow 0]{} \sum_i c_i(t) \mathcal{O}_i^R(x)$$

An operator
in blurred world



Renormalized operators
in original theory



less noisy, continuous

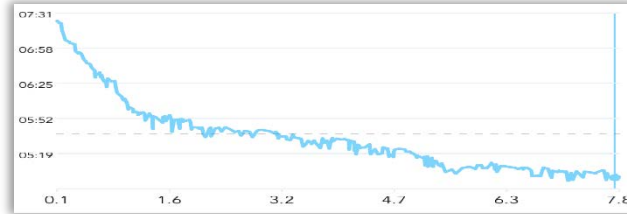
noisy, lattice discretized

Ex: energy momentum tensor, Suzuki, 2013; FlowQCD, 2014

Reference Scale and Lattice Spacing


Basic Idea

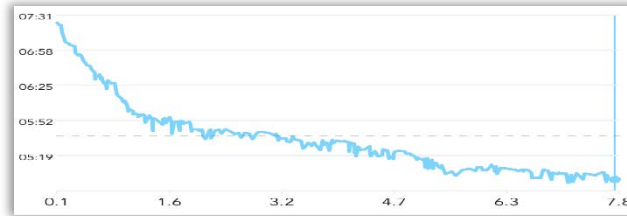
$$\langle \mathcal{O} \rangle_t \rightarrow$$



- function of t
- independent of renormalization

Basic Idea

$\langle \mathcal{O} \rangle_t$ 



- function of t
- independent of renormalization

Choice 1: $\mathcal{O} = t^2 E(t)$

Luescher, 2010

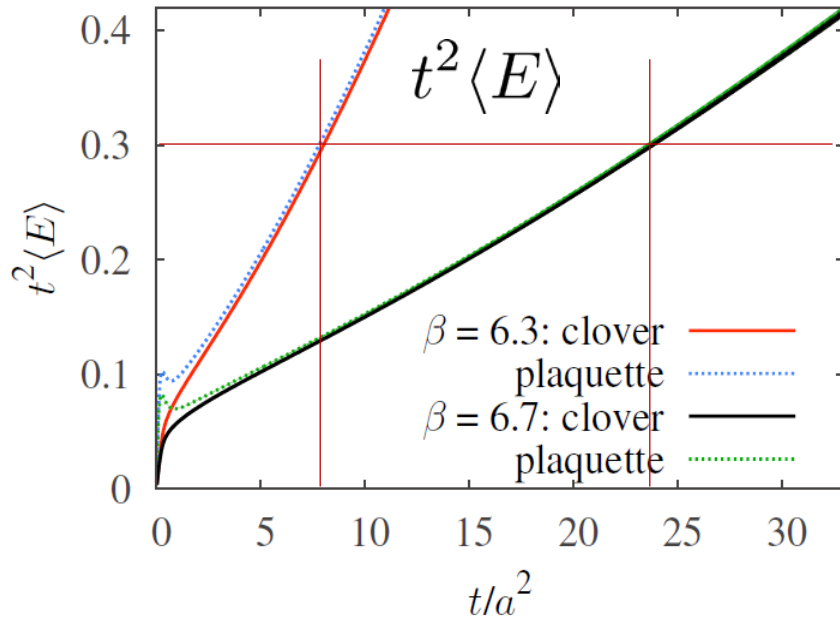
Choice 2: $\mathcal{O} = t \frac{d}{dt} t^2 E(t)$

Budapest-Wuppertal, 2012

$$E = \frac{1}{4} G_{\mu\nu}^a G_{\mu\nu}^a$$

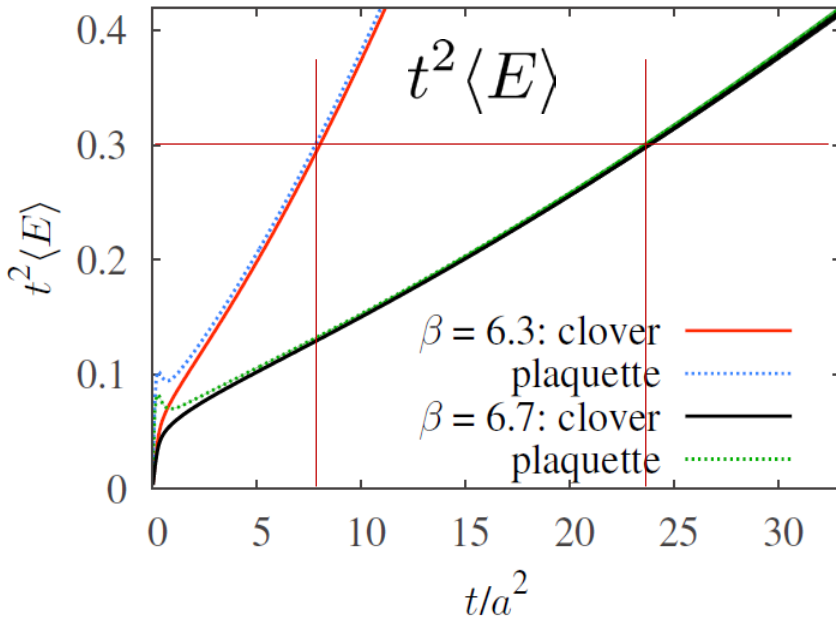
discretization effect suppressed

Behavior of $t^2 \langle E \rangle$

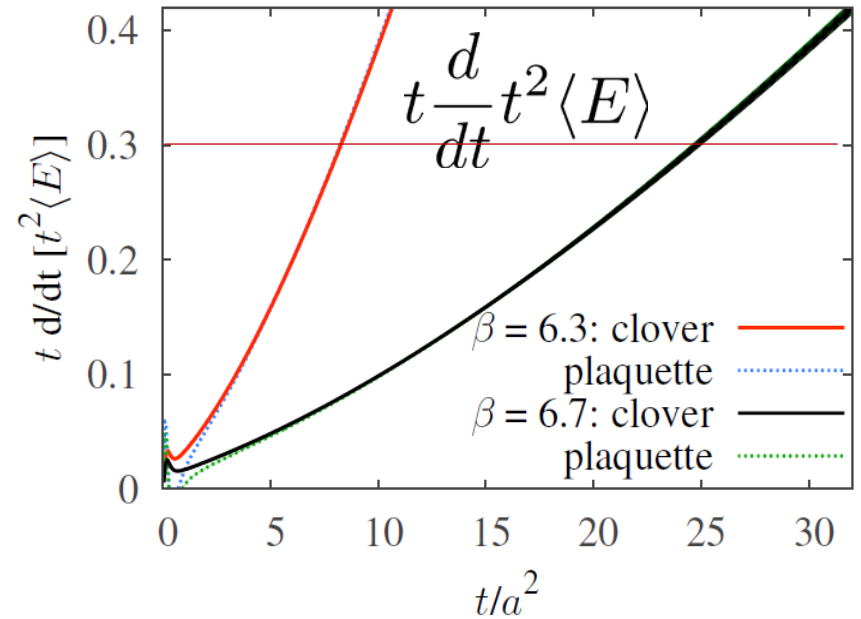


$$t_0 : t^2 \langle E \rangle|_{t=t_0} = 0.3$$

Behavior of $t^2 \langle E \rangle$



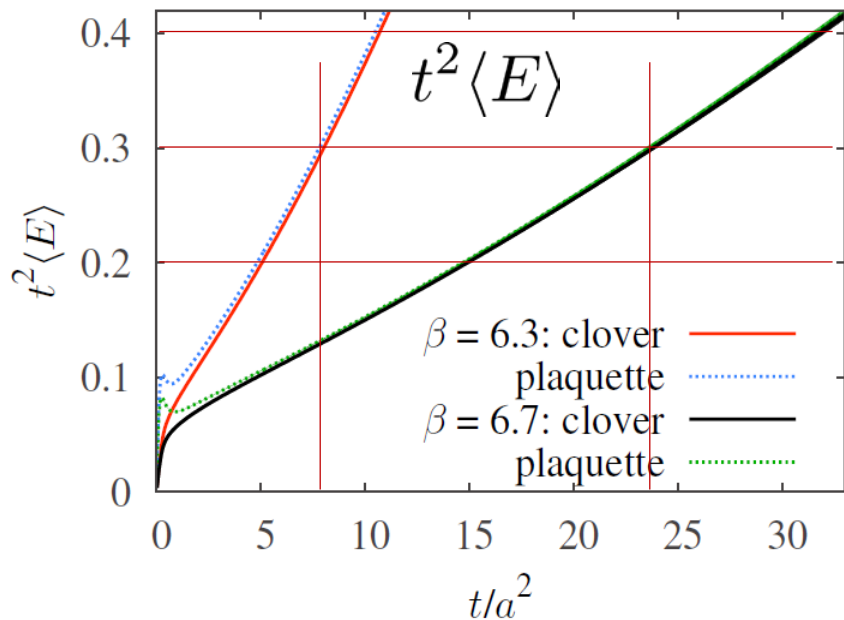
$$t_0 : t^2 \langle E \rangle |_{t=t_0} = 0.3$$



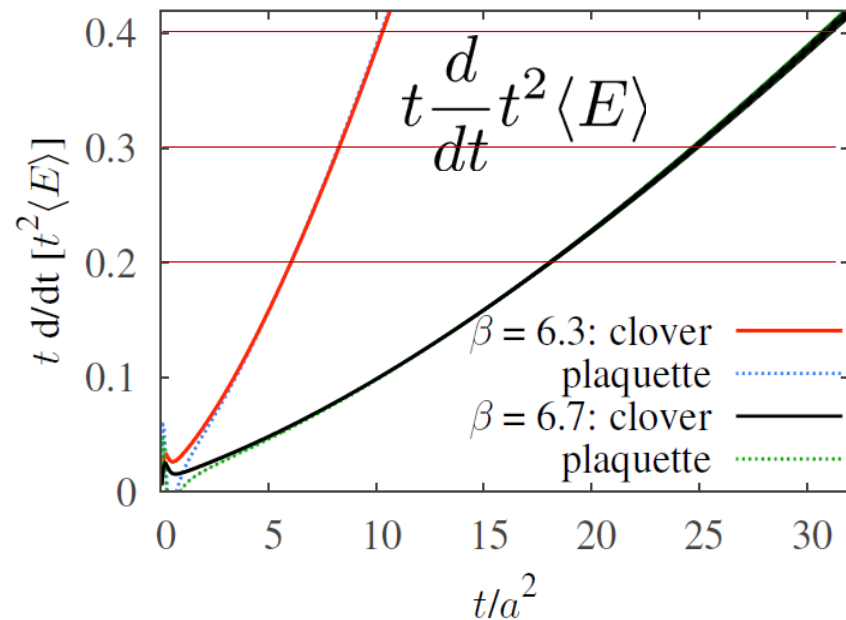
$$w_0 : t \frac{d}{dt} t^2 \langle E \rangle |_{t=w_0^2} = 0.3$$

- Advantages: {
- Less statistical errors
 - No fit analyses
 - No measurement of non-local obs.

Behavior of $t^2 \langle E \rangle$



$$t_0 : t^2 \langle E \rangle |_{t=t_0} = 0.3$$



$$w_0 : t \frac{d}{dt} t^2 \langle E \rangle |_{t=w_0^2} = 0.3$$

value in rhs is arbitrary

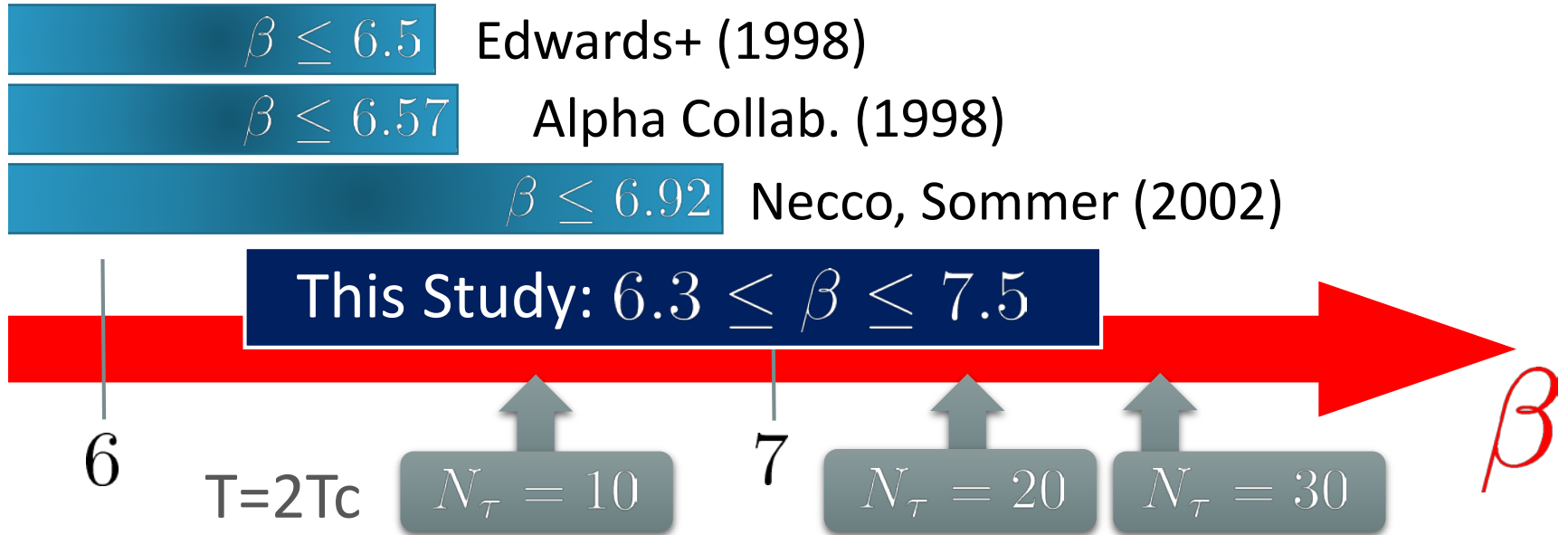
$t_{0.2}, t_{0.3}, t_{0.4}$

Our Choice



$w_{0.2}, w_{0.3}, w_{0.4}$

Lattice Spacing of SU(3) Wilson Action



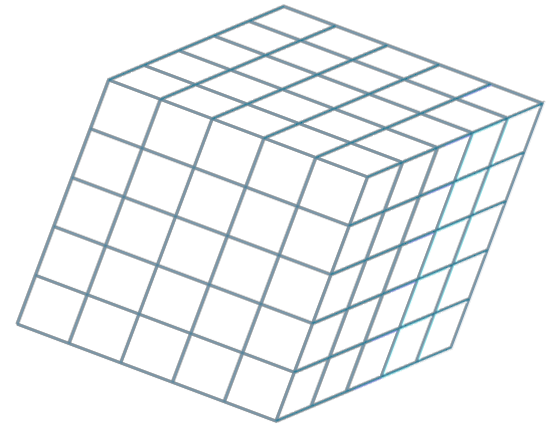
our parametrization for a

$$\log \left(\frac{w_{0.4}}{a} \right) (\beta) = \frac{4\pi^2}{33} \beta - 8.6853 + \frac{37.422}{\beta} - \frac{143.84}{\beta^2}$$

stat. err. < 0.4% / sys. err. < 0.7%

Numerical Setting

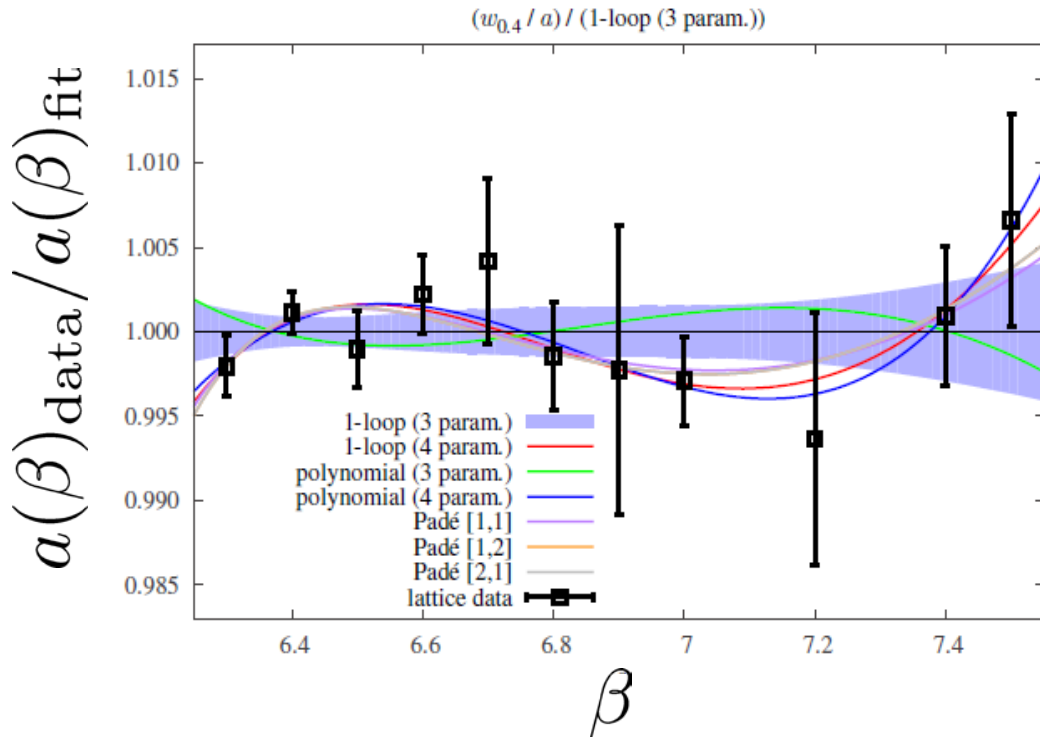
- SU(3) YM theory
- Wilson gauge action
- $w_{0.4} / w_{0.2}$ scaling



b	size	N_{conf}	b	size	N_{conf}
6.3	64 ⁴	30	6.9	64 ⁴	30
6.4	64 ⁴	100	7.0	96 ⁴	60
6.5	64 ⁴	49	7.2	96 ⁴	53
6.6	64 ⁴	100	7.4	128 ⁴	40
6.7	64 ⁴	30	7.5	128 ⁴	60
6.8	64 ⁴	100			

Each configuration is separated by 1000 updates (HB+OR⁵)
BlueGene/Q @ KEK

Parametrization with $w_{0.4}$



□ Use various fitting funcs

➔ estimate sys. err.

□ For $\beta=7.4$ and 7.5 , $w_{0.4}$ is obtained from $w_{0.2}$.

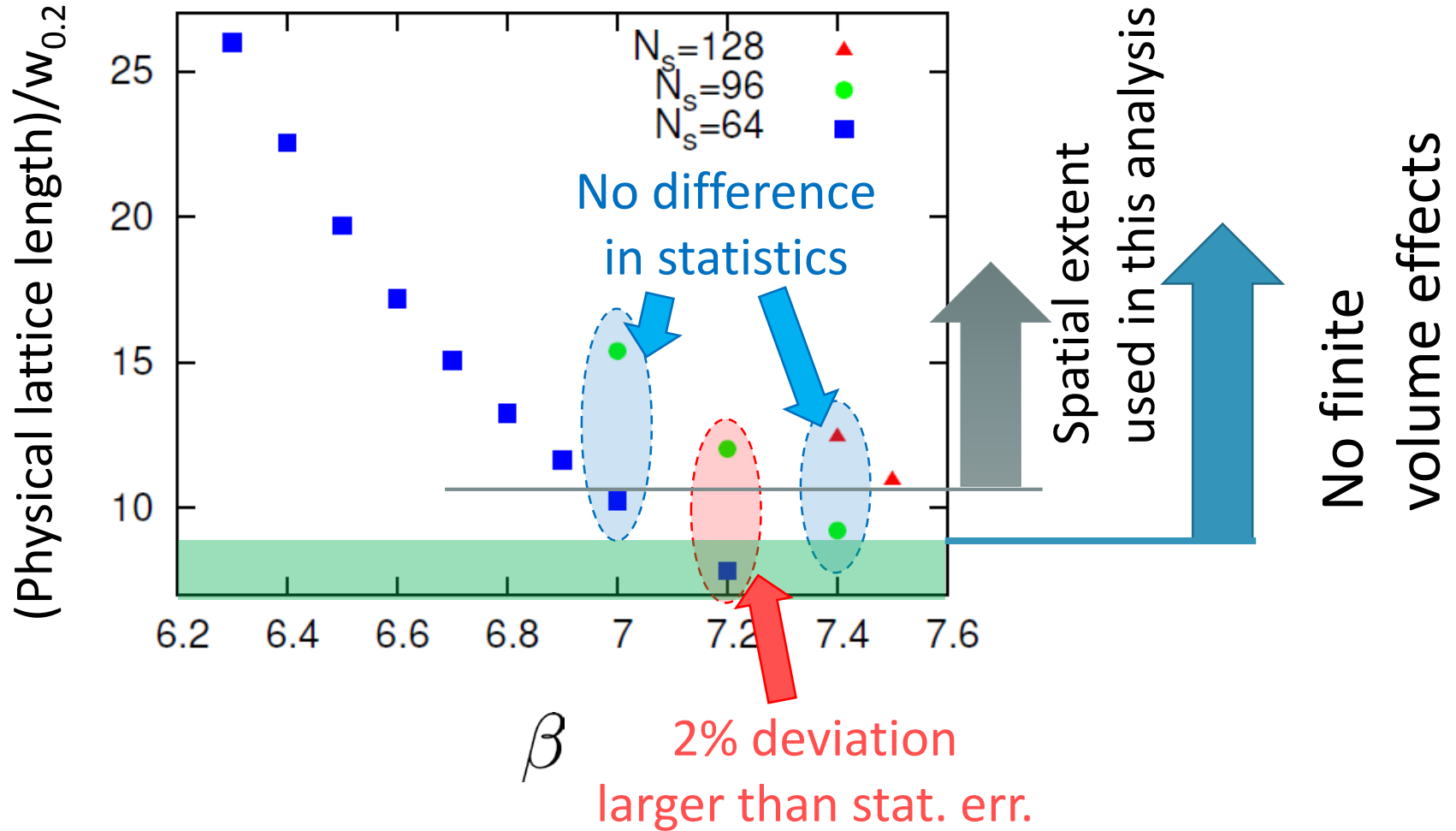
$$\log\left(\frac{w_{0.4}}{a}\right)(\beta) = \frac{4\pi^2}{33}\beta - 8.6853 + \frac{37.422}{\beta} - \frac{143.84}{\beta^2}$$

stat. err. < 0.4% / sys. err. < 0.7%

3 parameter fit

$\chi^2/\text{dof} = 0.92$

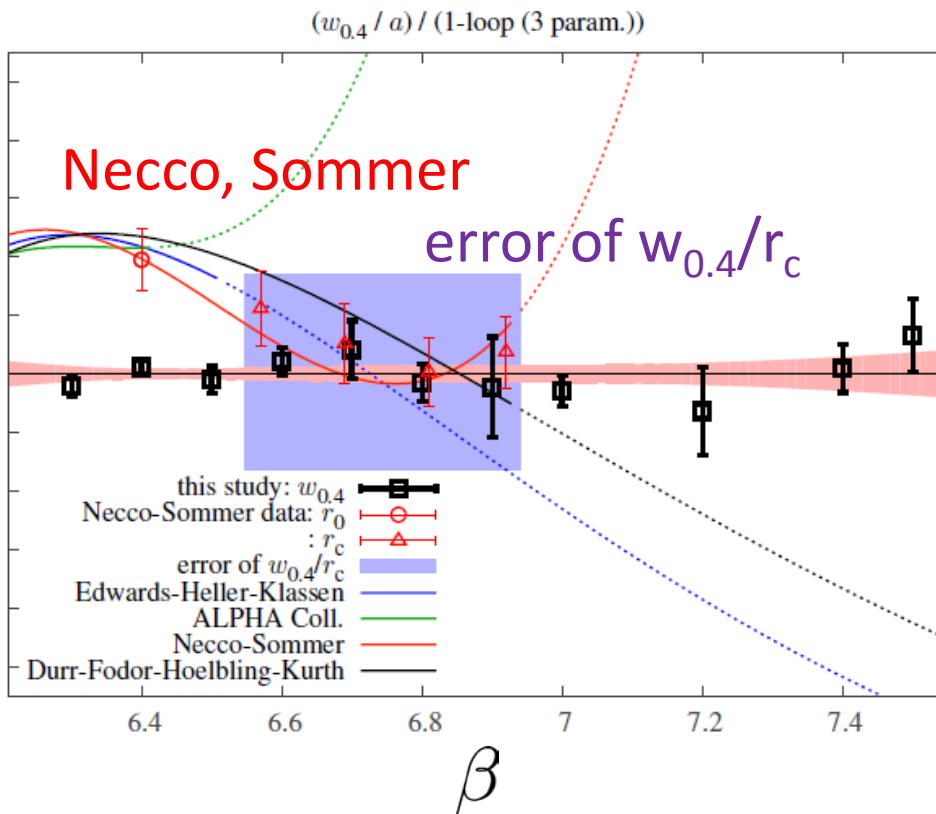
Finite Volume Effects



No finite volume effects within statistics

Comparison with Previous Studies

$a(\beta)_{\text{ours}} / a(\beta)_{\text{data, prev.}}$



Edwards, Heller, Klassen, 1998
 Alpha-Collaboration, 1998
 Necco, Sommer, 2002
 Durr, Fodor, Hoelbling, 2007

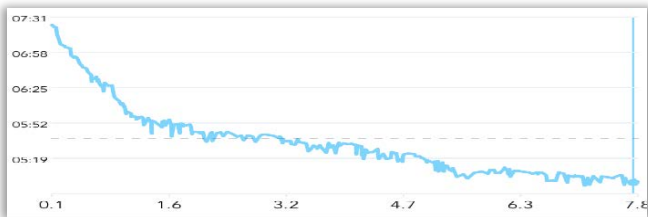
- Parametrizations agree with each other in available ranges.
- Errorbars of our data is smaller than previous ones.

Thermodynamics

Basic Idea: Small Flow Time Expansion

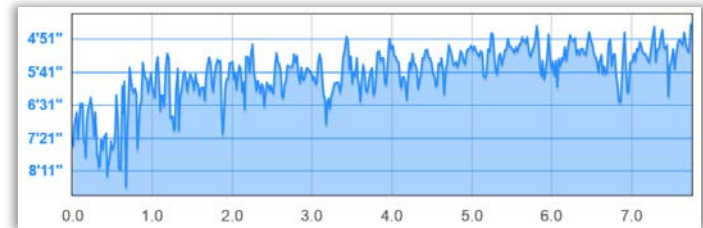
$$\tilde{\mathcal{O}}(t, x) \xrightarrow[t \rightarrow 0]{} \sum_i c_i(t) \mathcal{O}_i^R(x)$$

An operator
in blurred world



less noisy, continuous

Renormalized operators
in original theory



noisy, lattice discretized

SFTE of Energy-Momentum Tensor

Suzuki, 2013

- gauge-invariant dimension 4 operators

$$\left\{ \begin{array}{l} U_{\mu\nu}(t, x) = G_{\mu\rho}(t, x)G_{\nu\rho}(t, x) - \frac{1}{4}\delta_{\mu\nu}G_{\mu\nu}(t, x)G_{\mu\nu}(t, x) \\ E(t, x) = \frac{1}{4}\delta_{\mu\nu}G_{\mu\nu}(t, x)G_{\mu\nu}(t, x) \end{array} \right.$$

Remormalized EMT

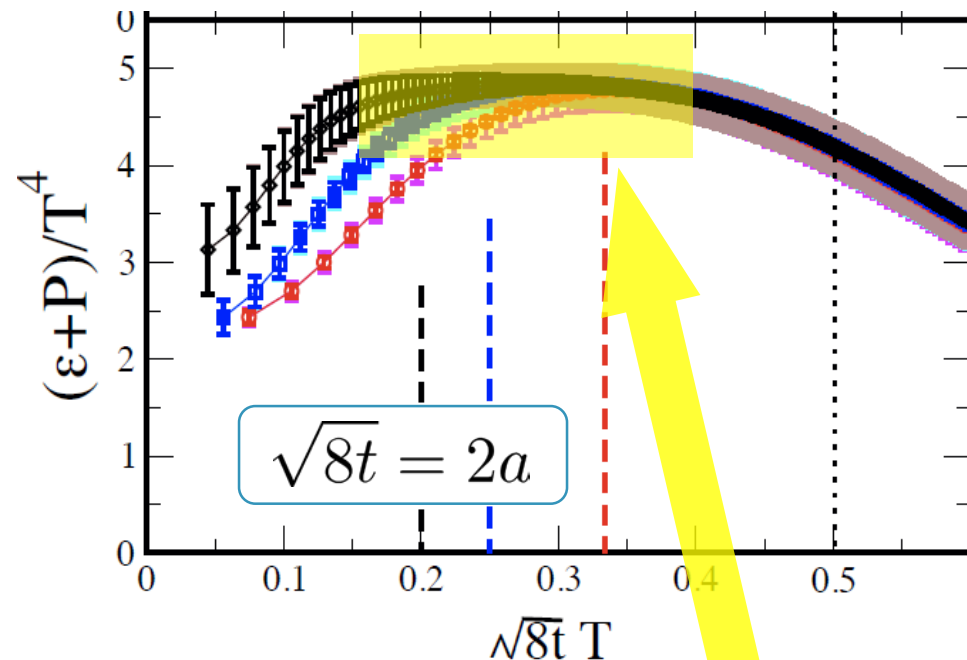
$$T_{\mu\nu}^R(x) = \lim_{t \rightarrow 0} \left[\frac{1}{\alpha_U(t)} U_{\mu\nu}(t, x) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} E(t, x)_{\text{subt.}} \right]$$

Suzuki coeffs. $\left\{ \begin{array}{l} \alpha_U(t) = g^2 [1 + 2b_0 s_1 g^2 + O(g^4)] \\ \alpha_E(t) = \frac{1}{2b_0} [1 + 2b_0 s_2 g^2 + O(g^4)] \end{array} \right. \quad g = g(\sqrt{8t})$

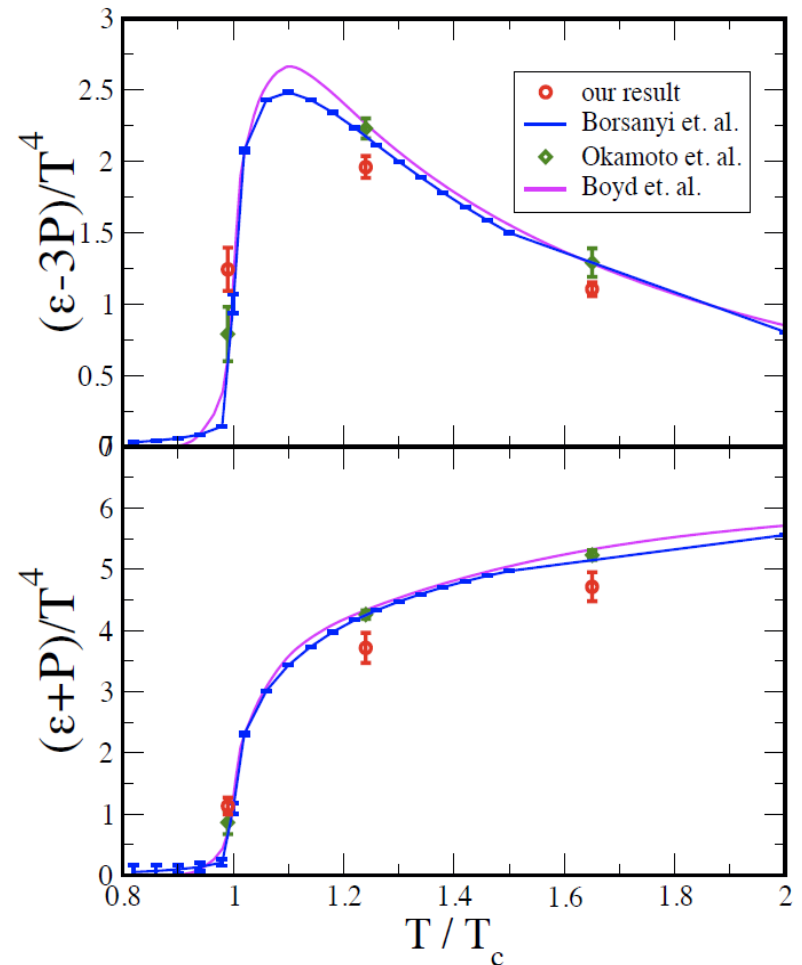
SU(3) Thermodynamics with $N_t=6, 8, 10$

FlowQCD, PRD90,011501 (2014)

Entropy density $T=1.65T_c$



Stable region for
 $2a < \sqrt{8t} < 0.4T^{-1}$

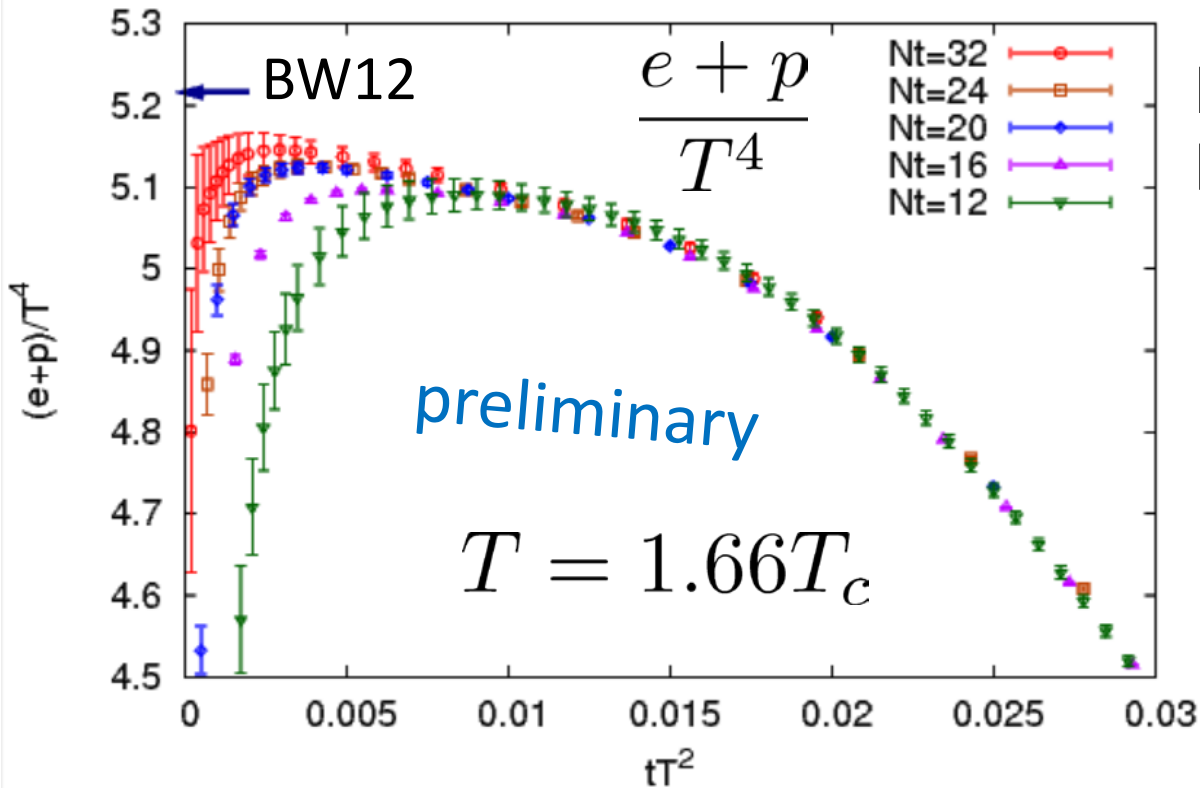


New Results: Thermodynamics (e+p)

$$\tilde{T}_{\mu\nu}(t) = \frac{1}{\alpha_U(t)} U_{\mu\nu}(t) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} E(t)_{\text{subt.}}$$

FlowQCD, in prep.

$$T_{\mu\nu}^R = \tilde{T}_{\mu\nu}(t) + O(t)$$



□ Existence of $O(t)$ effect

□ Linear behavior for

$$tT^2 < 0.015$$

$$(\sqrt{8t} < 0.35T^{-1})$$

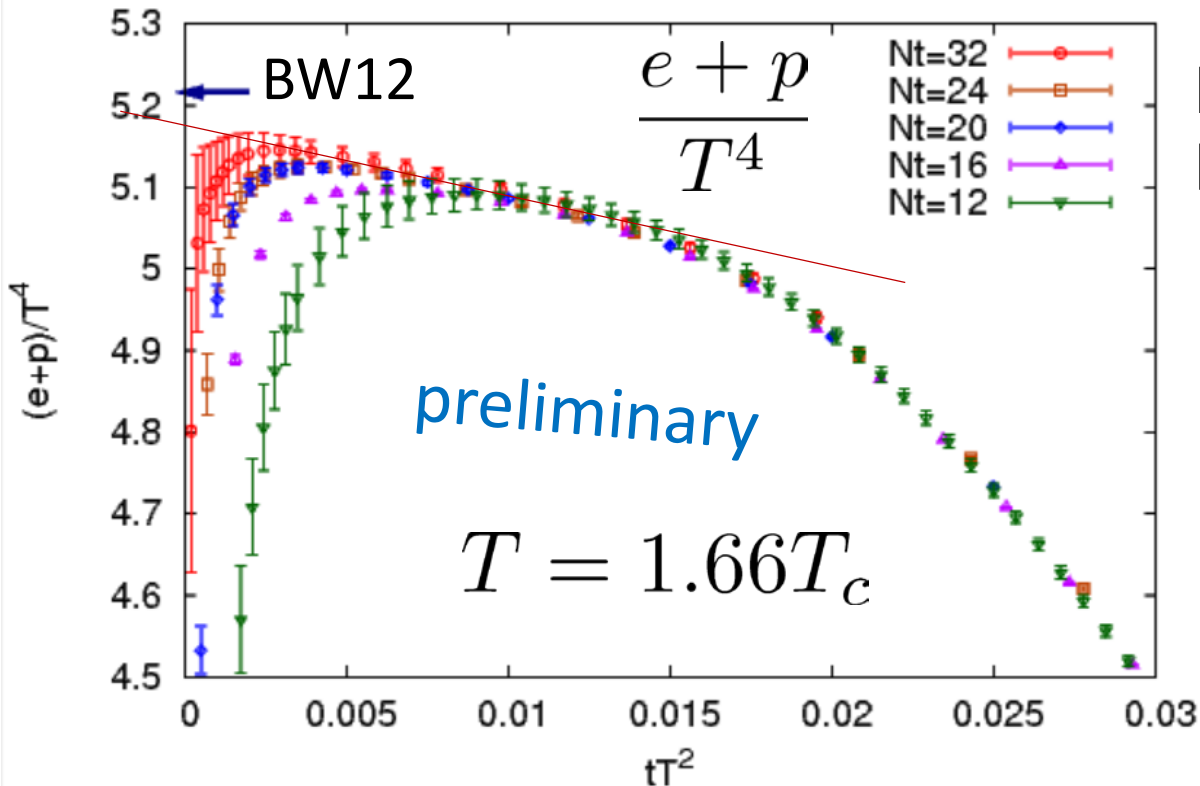
- $t \rightarrow 0$ limit is necessary

New Results: Thermodynamics (e+p)

$$\tilde{T}_{\mu\nu}(t) = \frac{1}{\alpha_U(t)} U_{\mu\nu}(t) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} E(t)_{\text{subt.}}$$

FlowQCD, in prep.

$$T_{\mu\nu}^R = \tilde{T}_{\mu\nu}(t) + O(t)$$



□ Existence of $O(t)$ effect

□ Linear behavior for

$$tT^2 < 0.015$$

$$(\sqrt{8t} < 0.35T^{-1})$$

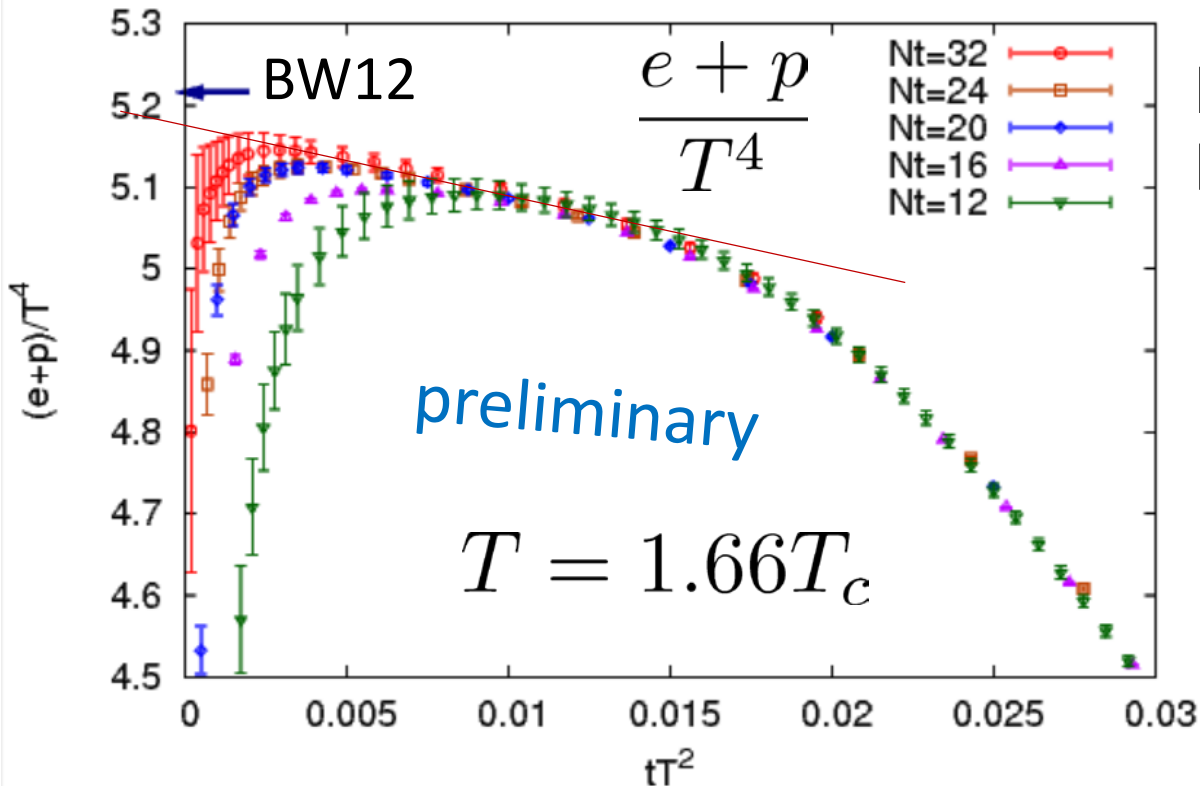
- $t \rightarrow 0$ limit is necessary

New Results: Thermodynamics (e+p)

$$\tilde{T}_{\mu\nu}(t) = \frac{1}{\alpha_U(t)} U_{\mu\nu}(t) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} E(t)_{\text{subt.}}$$

FlowQCD, in prep.

$$T_{\mu\nu}^R = \tilde{T}_{\mu\nu}(t) + O(t)$$



□ Existence of $O(t)$ effect

□ Linear behavior for

$$tT^2 < 0.015$$

$$(\sqrt{8t} < 0.35T^{-1})$$

- $t \rightarrow 0$ limit is necessary

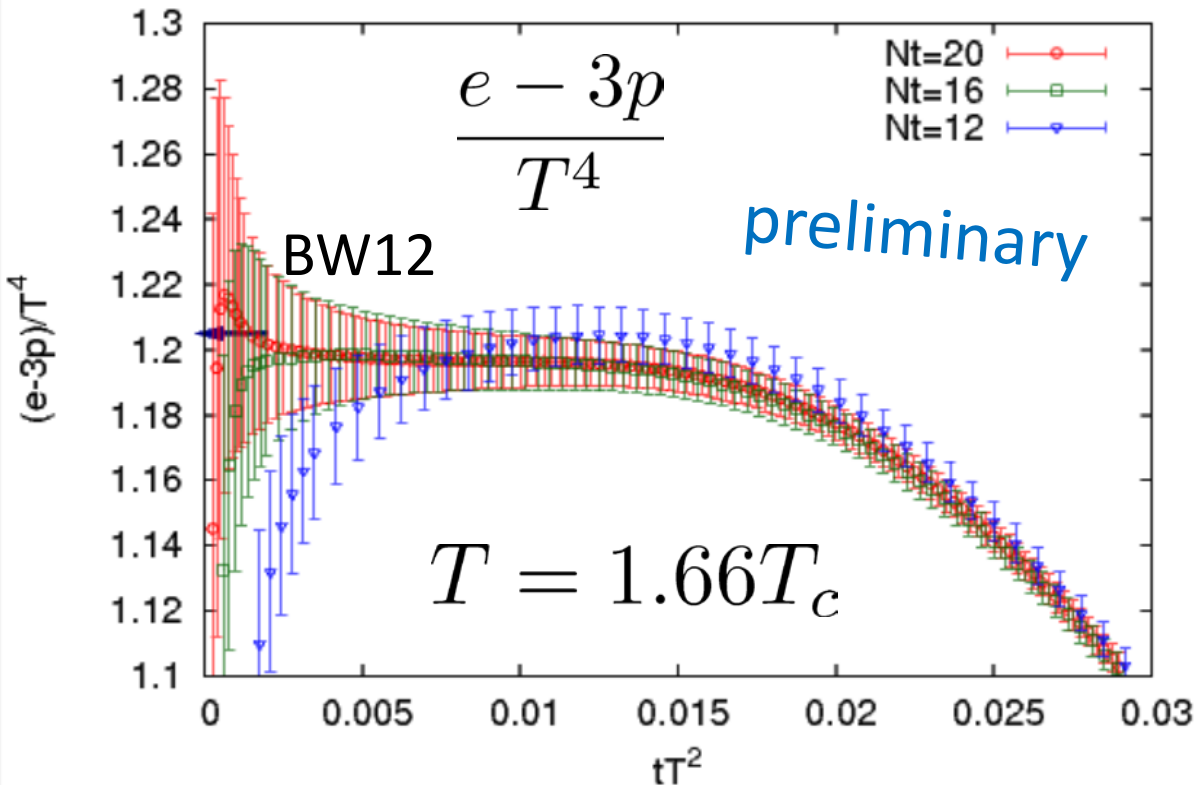
BW12: Budapest-Wuppertal, 2012



New Results: Thermodynamics (e-3p)

$$\tilde{T}_{\mu\nu}(t) = \frac{1}{\alpha_U(t)} U_{\mu\nu}(t) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} E(t)_{\text{subt.}}$$

FlowQCD, in prep.



□ Good agreement with previous study

□ Stable plateau for
 $tT^2 < 0.015$
($\sqrt{8t} < 0.35T^{-1}$)

Summary

Accurate lattice spacing of Wilson gauge action for $6.3 < \beta < 7.5$ is now available.

$$\log \left(\frac{w_{0.4}}{a} \right) (\beta) = \frac{4\pi^2}{33} \beta - 8.6853 + \frac{37.704}{\beta} - \frac{144.77}{\beta^2}$$

The EMT defined with gradient flow is nicely applied to the measurement of thermodynamics.

The $a \rightarrow 0$ behavior is checked on fine lattices.

- taking $t \rightarrow 0$ limit is needed.
- application to full QCD \rightarrow E. Itou, 18th, room 403

Summary

Accurate lattice spacing of Wilson gauge action for $6.3 < \beta < 7.5$ is now available.

$$\log \left(\frac{w_{0.4}}{a} \right) (\beta) = \frac{4\pi^2}{33} \beta - 8.6853 + \frac{37.704}{\beta} - \frac{144.77}{\beta^2}$$



The EMT defined with gradient flow is nicely applied to the measurement of thermodynamics.

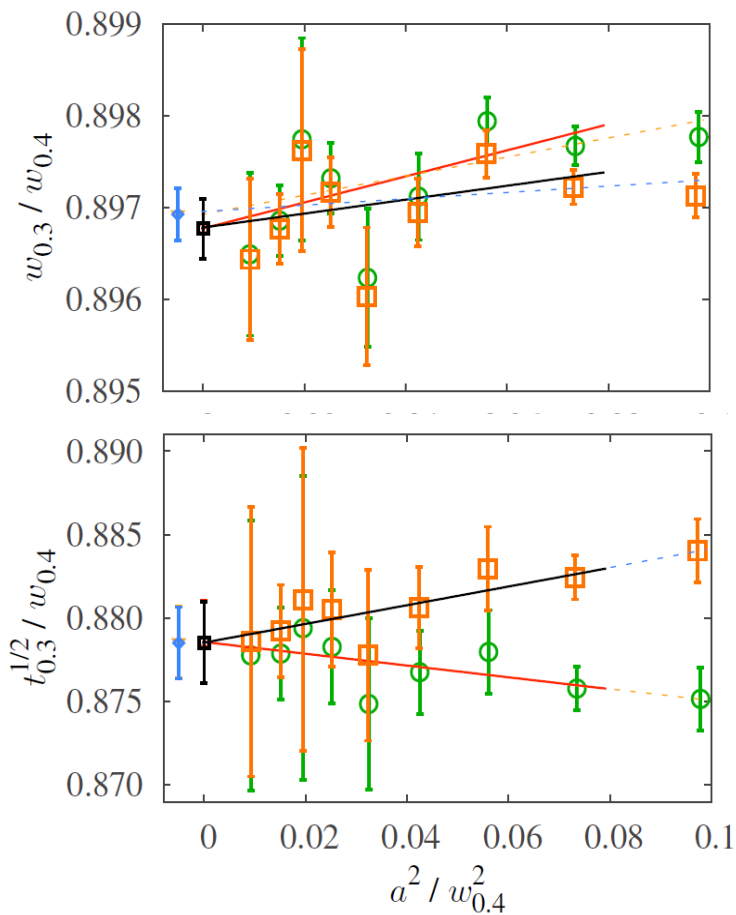
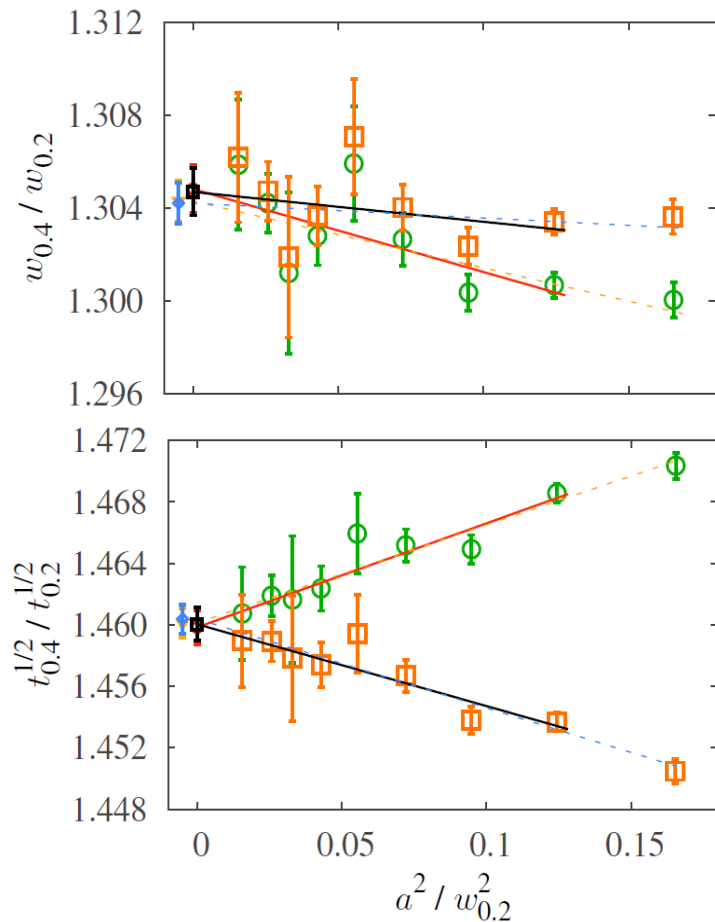
The $a \rightarrow 0$ behavior is checked on fine lattices.

- taking $t \rightarrow 0$ limit is needed.
- application to full QCD \rightarrow E. Itou, 18th, room 403



Backup

Various Reference Scales



plaq.
clover

$\sqrt{t_{0.4}}/w_{0.4}$	$\sqrt{t_{0.3}}/w_{0.4}$	$\sqrt{t_{0.2}}/w_{0.4}$	$w_{0.3}/w_{0.4}$	$w_{0.2}/w_{0.4}$
1.0164(32)(3)	0.8785(24)(0)	0.6952(18)(2)	0.8968(3)(2)	0.7665(6)(2)

$r_c/w_{0.4}$	$r_0/w_{0.4}$	$\sqrt{\sigma}w_{0.4}$	$T_c w_{0.4}$	$w_{0.4}\Lambda_{\overline{\text{MS}}}$
1.328(21)(7)	2.587(45)	0.455(8)	0.285(5)	0.233(19)