

Lattice study for conformal windows of
SU(2) and SU(3) gauge theories
with fundamental (unimproved staggered) fermions

C.-J. David Lin
National Chiao-Tung University, Taiwan

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Kobe, Japan

Motivation

- Controversy over the $SU(3)$, 12-flavour theory.
 - ➔ Control of the systematics.
- The $SU(2)$, 6-flavour theory is chirally-broken.
 - ➔ Next natural thing to study is the 8-flavour theory.

Running coupling and linearised beta function in SU(3) gauge theory with 12 flavours:

Gradient flow scheme with twisted BC and massless fermions

C.-J.D. Lin (NCTU)

K. Ogawa (NCTU  industry)

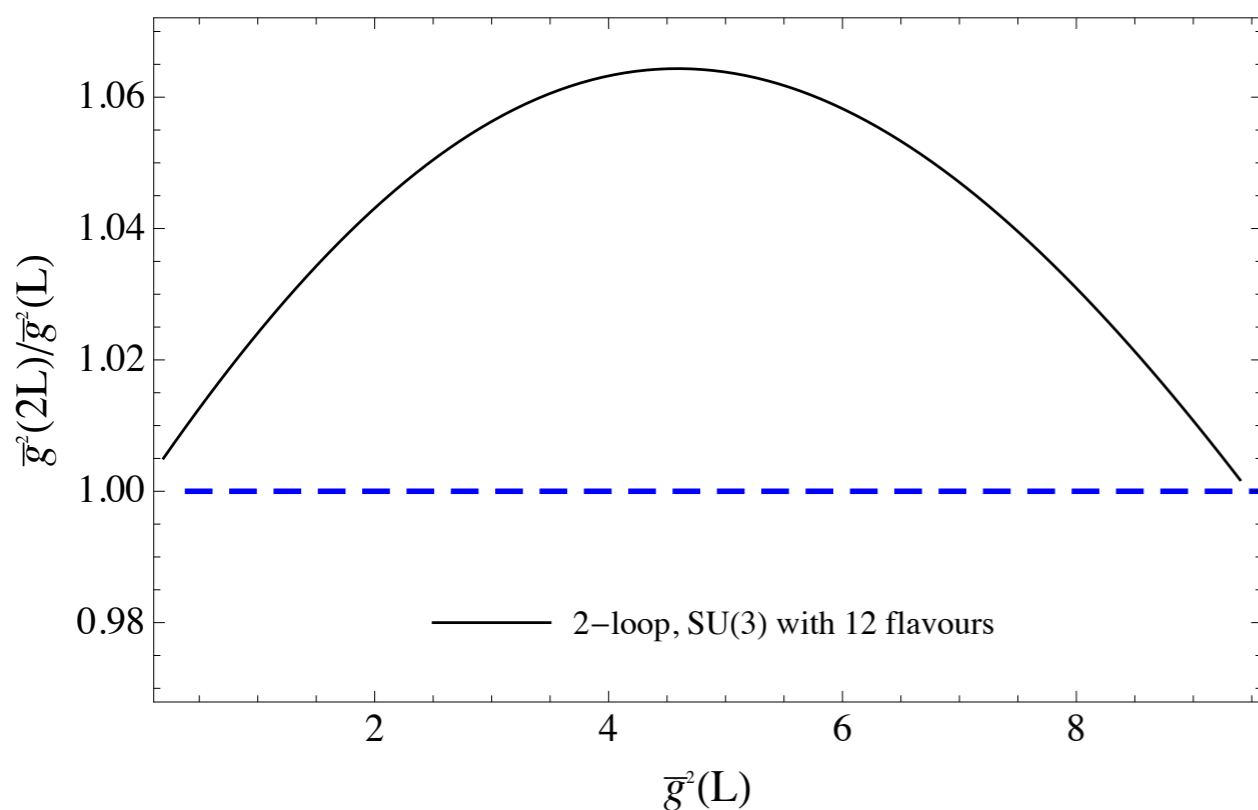
H. Ohki (KMI  RIKEN)

A. Ramos (CERN)

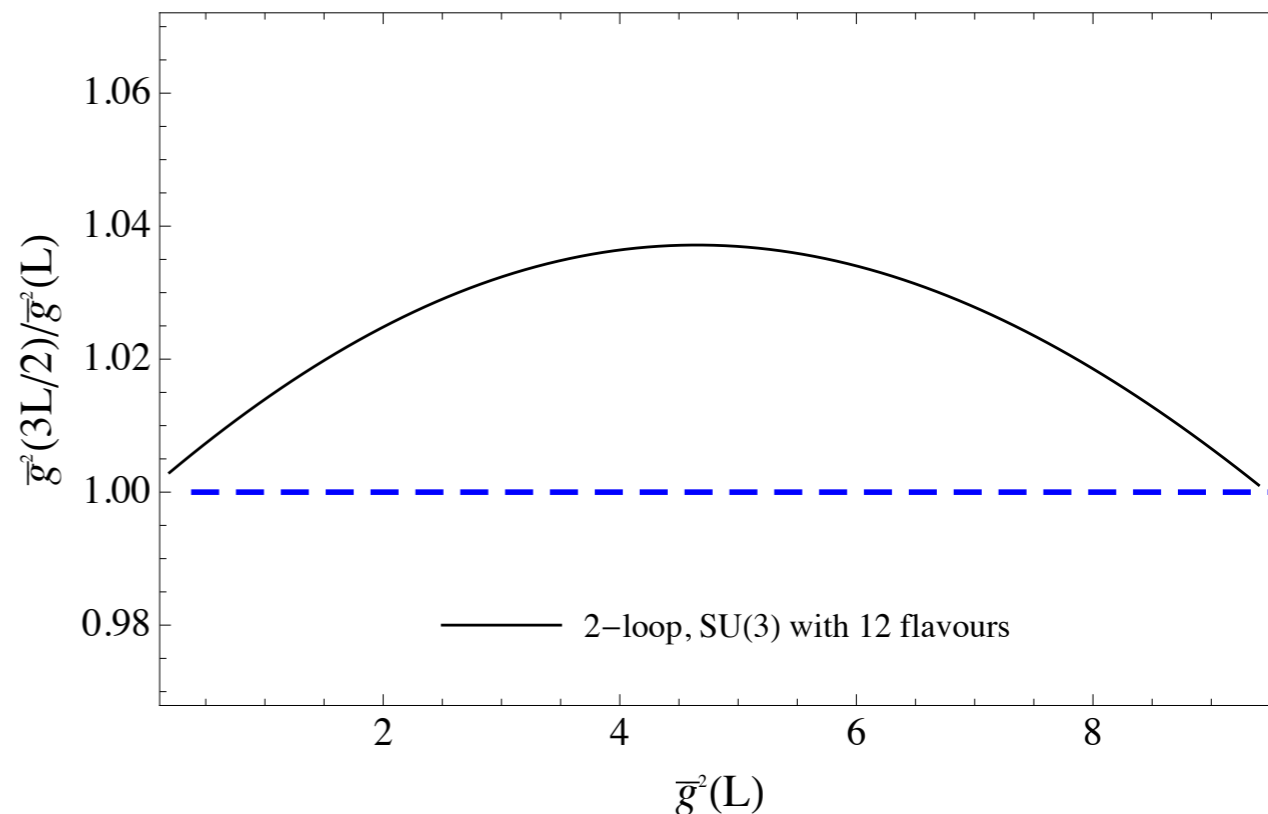
E. Shintani (Mainz U.  RIKEN)

How well do we have to control the errors?

In order to demonstrate the flow from the UV to the IR fixed points...



~1%



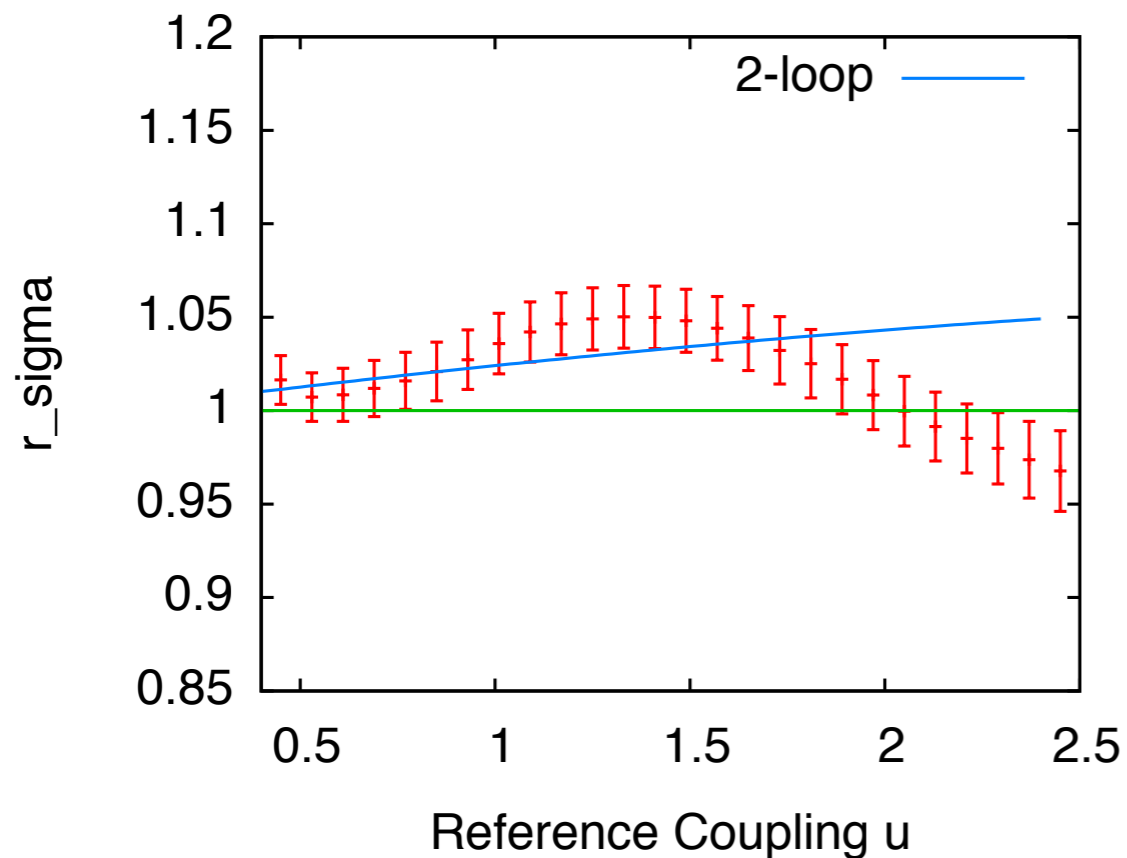
~0.6%

➡ It is important to choose good observables and strategies.

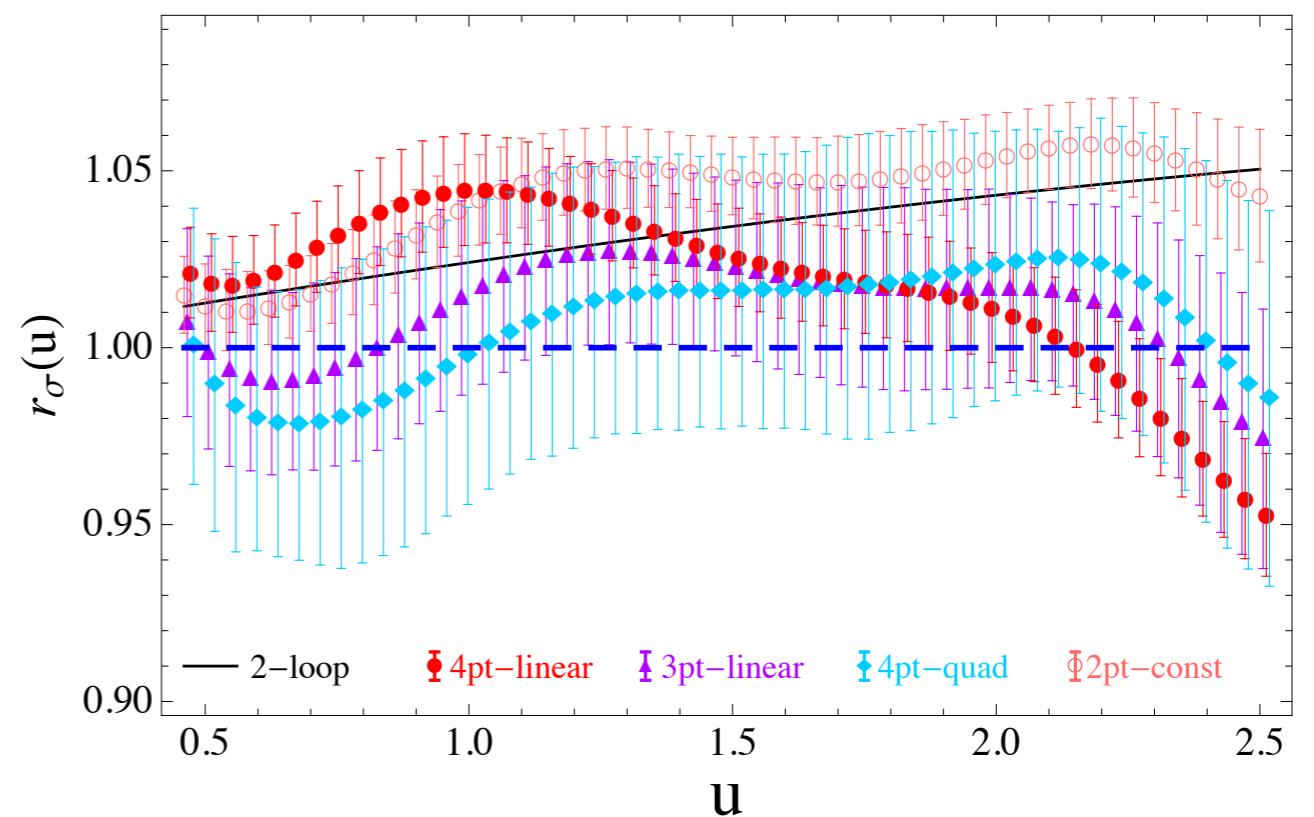
Example: The use of the Twisted Polyakov Loop scheme

C.-J.D.L., K.Ogawa, H.Ohki, E.Shintani, 2012

K.Ogawa, lattice 2013



without ($L/a=12 \rightarrow L/a = 24$)
systematics severely underestimated...



with ($L/a=12 \rightarrow L/a=24$)

It is challenging to draw any conclusion from such a “noisy scheme”.

The Gradient Flow scheme

- The quantity, $\langle E(t) \rangle = \frac{1}{4} \langle G_{\mu\nu}(t) G_{\mu\nu}(t) \rangle$, is finite when expressed in terms of renormalised coupling at positive flow time.

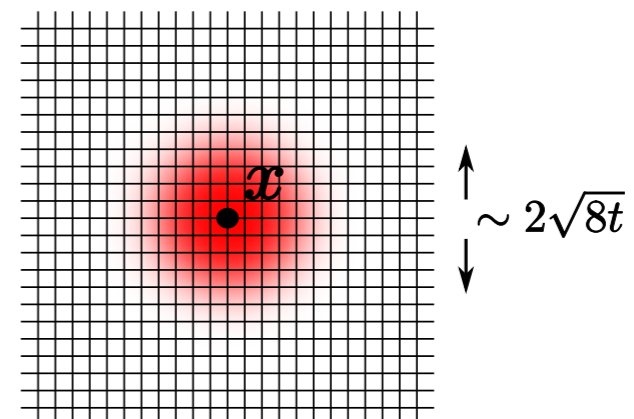
- In a colour-twisted box, can define,

$$\bar{g}_{\text{GF}}^2(L) = \mathcal{N}^{-1} t^2 \langle E(t) \rangle = \bar{g}_{\text{MS}}^2 + \mathcal{O}(\bar{g}_{\text{MS}}^4),$$

with tree-level improvement.

- Use the **clover operator**, as well as the **plaquette**, to extract $\langle E(t) \rangle$.

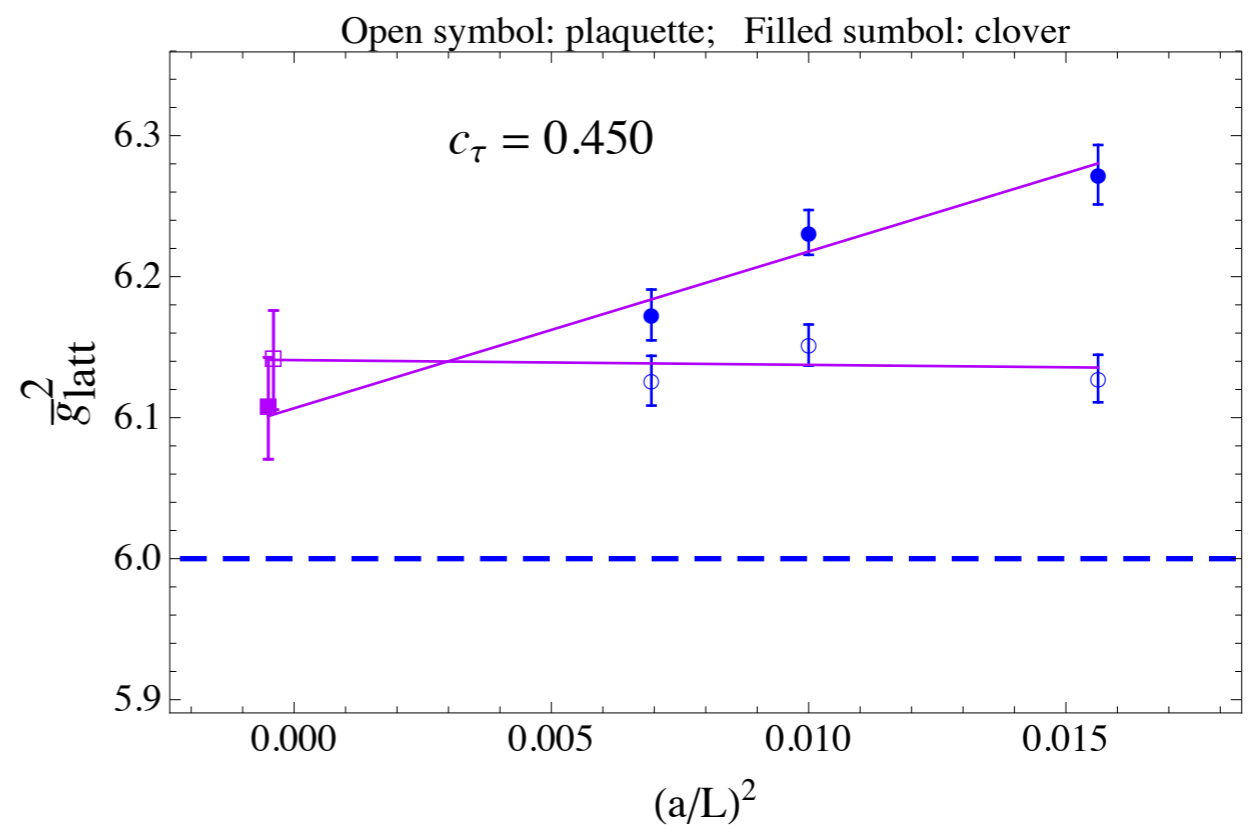
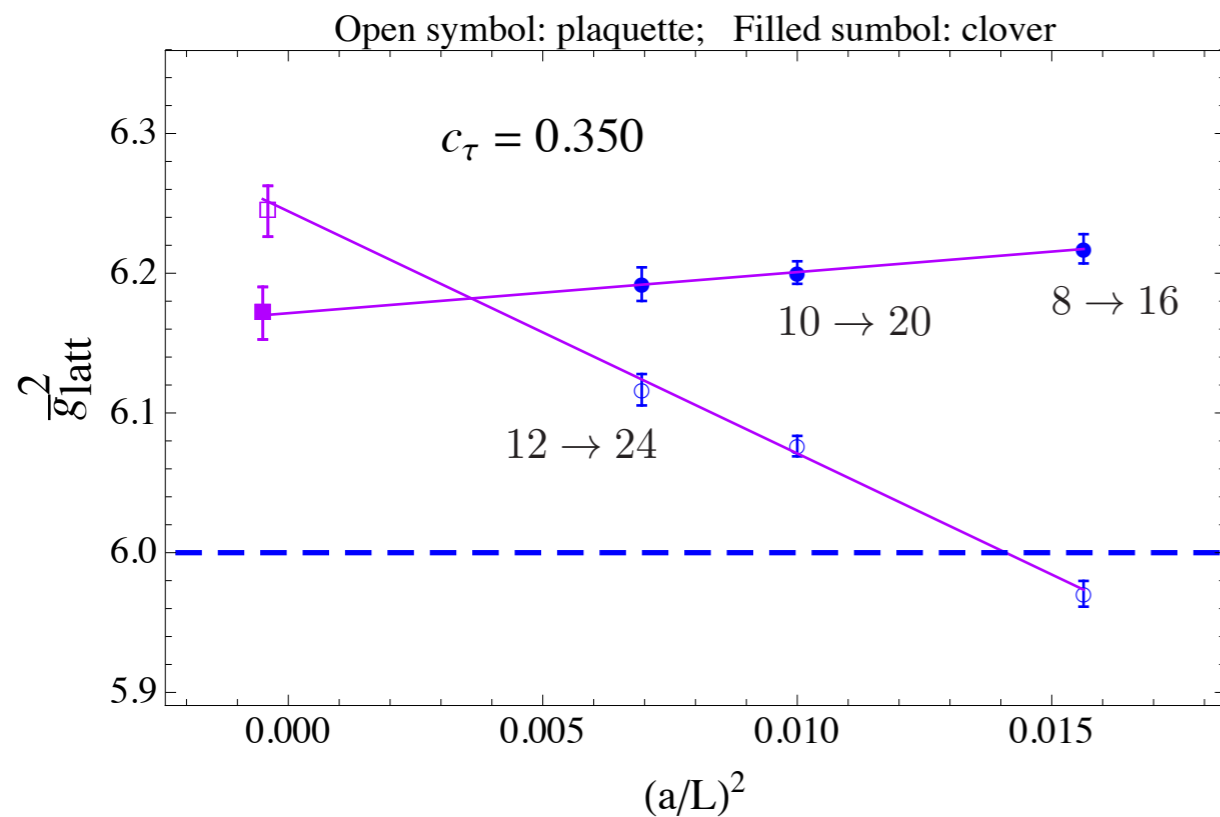
- Step scaling at fixed $c_\tau = \frac{\sqrt{8t}}{L}$.



Caution for the continuum extrapolation

Lattice artefacts arise from the action, the observable, and the flow.

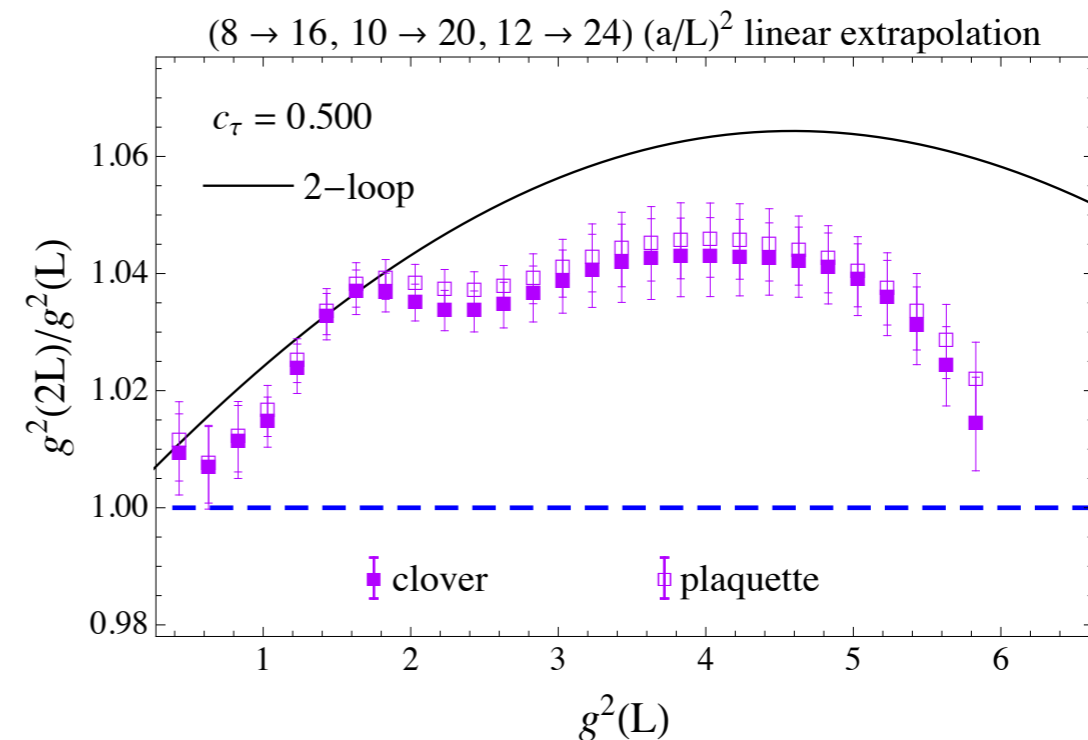
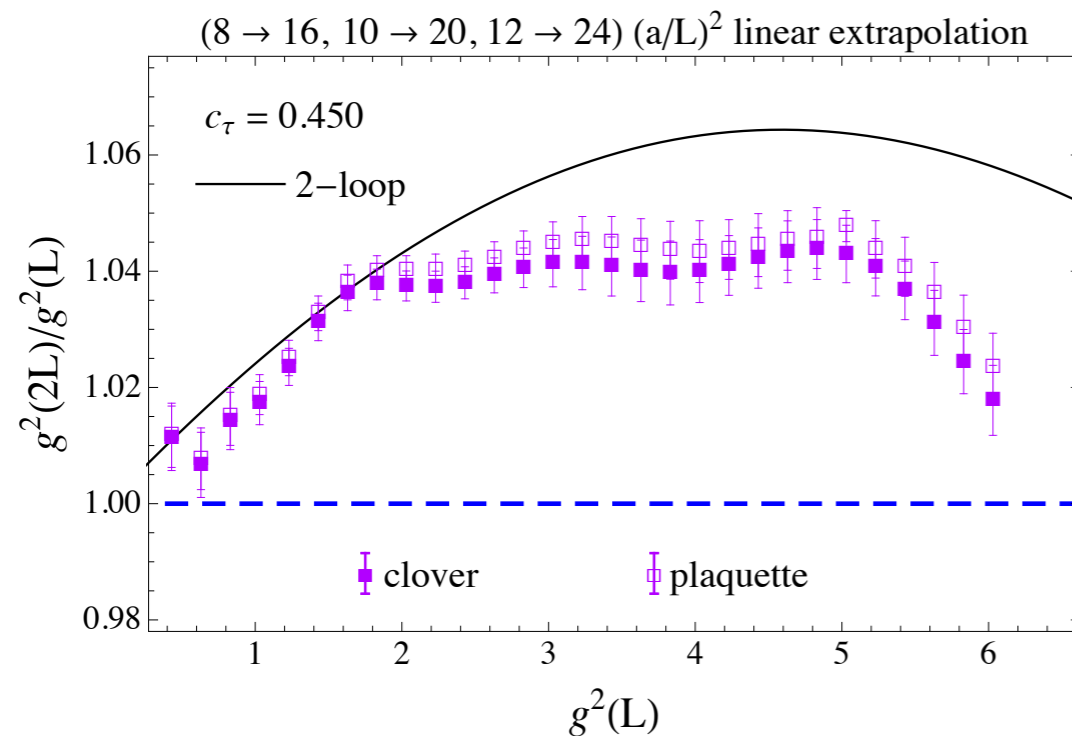
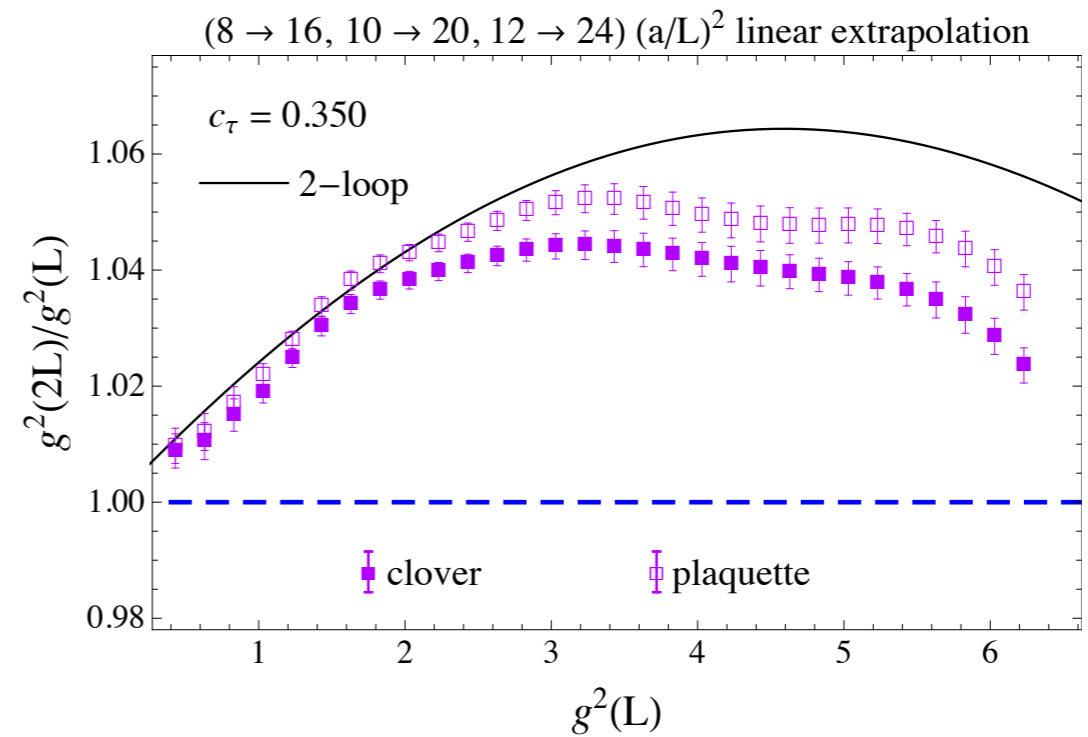
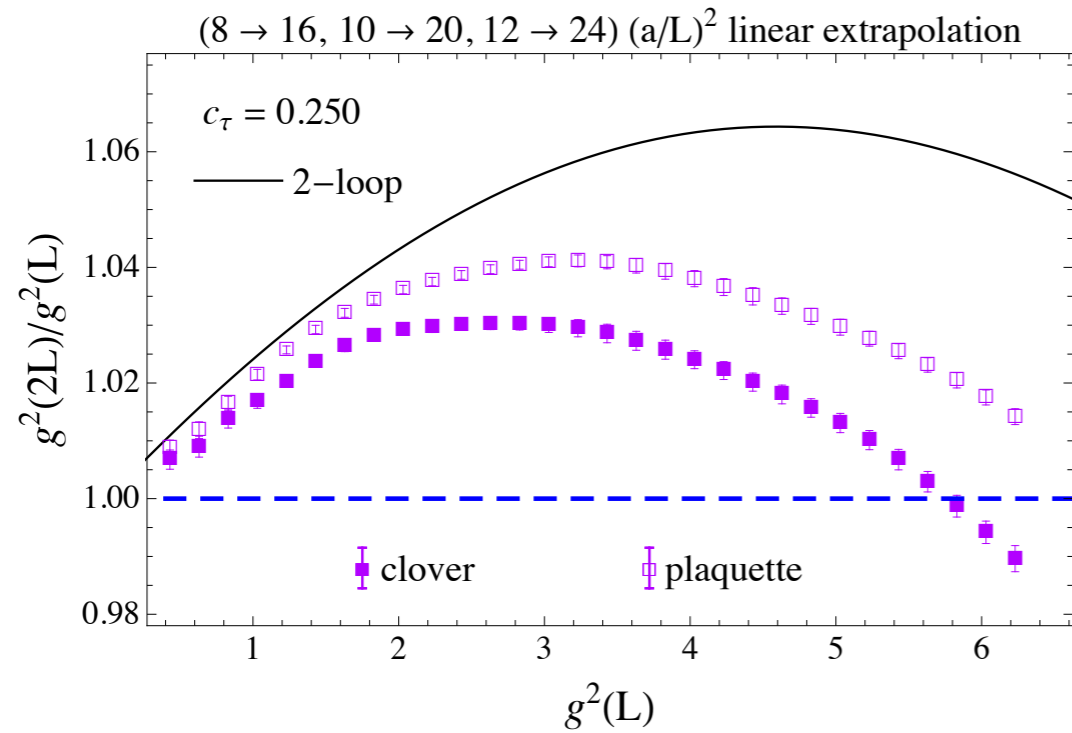
➡ Should improve all of these in principle.



1-5% extrapolation, but...

➡ It is important to use two discretisations and check the lattice artefacts.

The “conventional” continuum extrapolation



Require c_τ to be at least 0.45 to control the extrapolation.

Assuming we are close to an IRFP..

- The finite-size scaling behaviour is governed by the linearised beta function,

$$\beta(g^2) \equiv -\hat{L} \frac{dg^2}{d\hat{L}} = \gamma(g^2 - g_*^2)$$

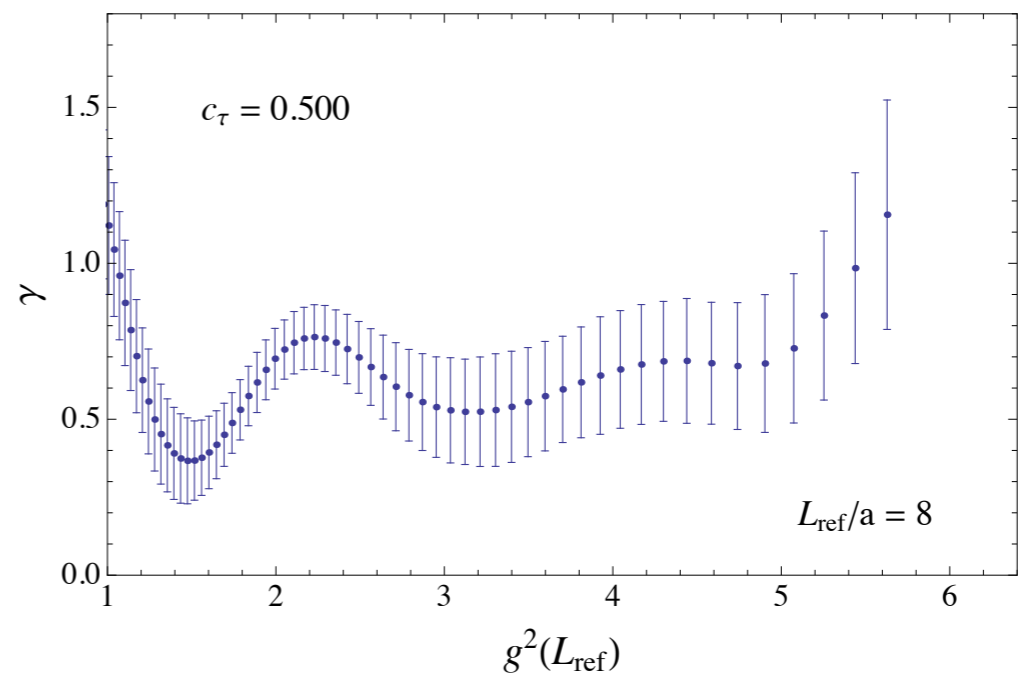
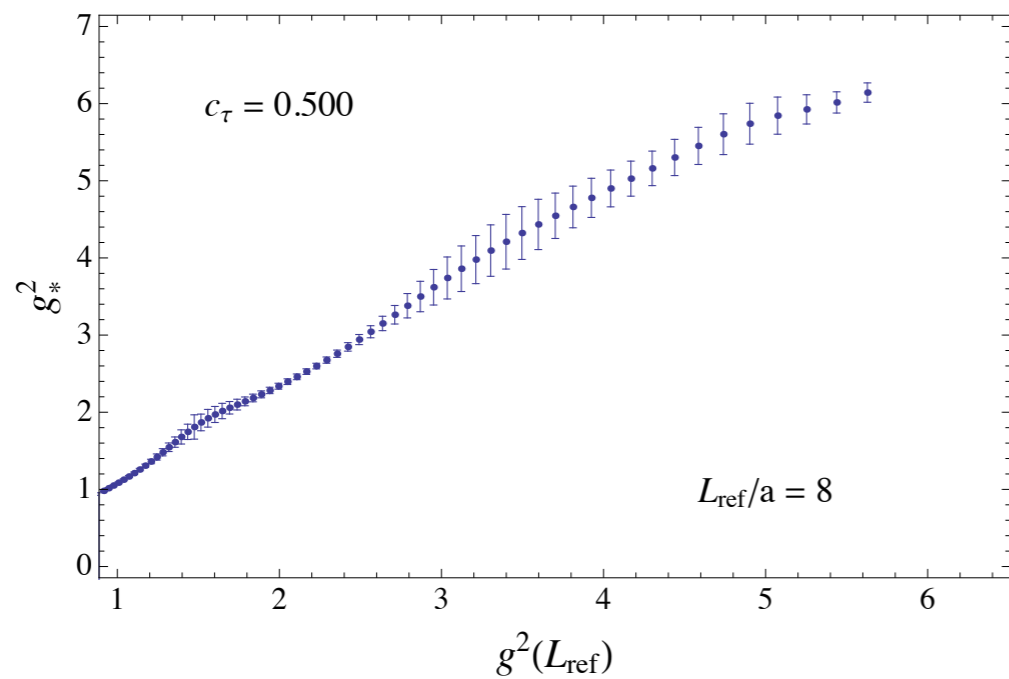
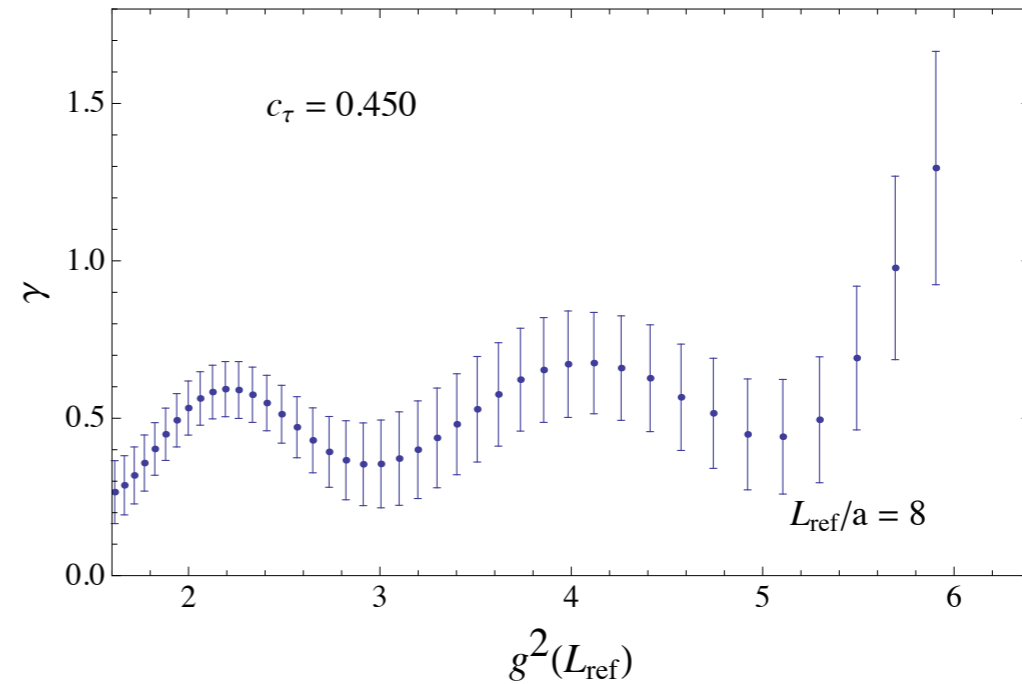
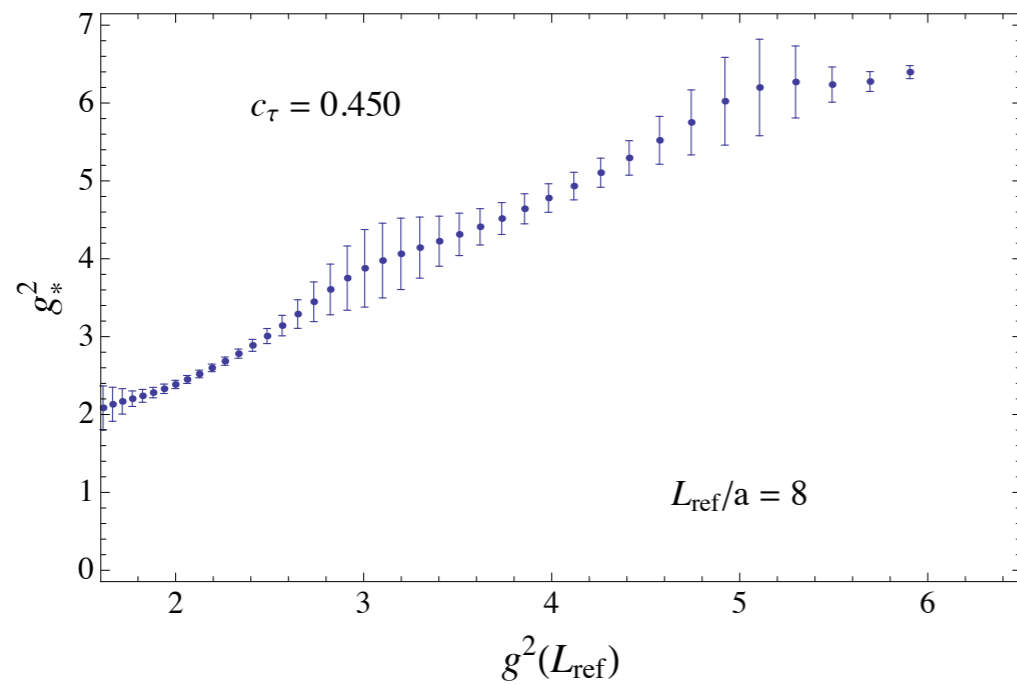
- Integrating from a reference volume, \hat{L}_{ref} , at fixed lattice spacing,

$$g^2(\hat{L}) = g_*^2 + \left[g^2(\hat{L}_{\text{ref}}) - g_*^2 \right] \left(\frac{\hat{L}}{\hat{L}_{\text{ref}}} \right)^{-\gamma}.$$

- At fixed lattice spacing, fit γ and g_* , then examine whether plateaus against $g^2(\hat{L}_{\text{ref}})$ are observed.

Results at two flow times

From the clover discretisation



Weak (if any) evidence for plateaus...

Conclusion for this theory


For the study of the running coupling...

- *Necessary* to control error to the sub-percentage level.
- *Necessary* to use two discretisations.
To check the control of the continuum extrapolation.
- The coupling runs slower than the two-loop prediction.
- If there is an IRFP, $g_*^2 > 6$ in the gradient-flow scheme.
- However, we do not see an IRFP in our data.

Non-thermal phase structure of SU(2) gauge theory with 8 flavours

The chiral condensate via RMT
Fermion masses 0.005, 0.010, and 0.015

C.Y-H. Huang (NCTU)

I. Kanamori (NCTU  Hiroshima U.)

C.-J.D. Lin (NCTU)

K. Ogawa (NCTU  industry)

H. Ohki (NCTU  RIKEN)

E. Rinaldi (LLNL)

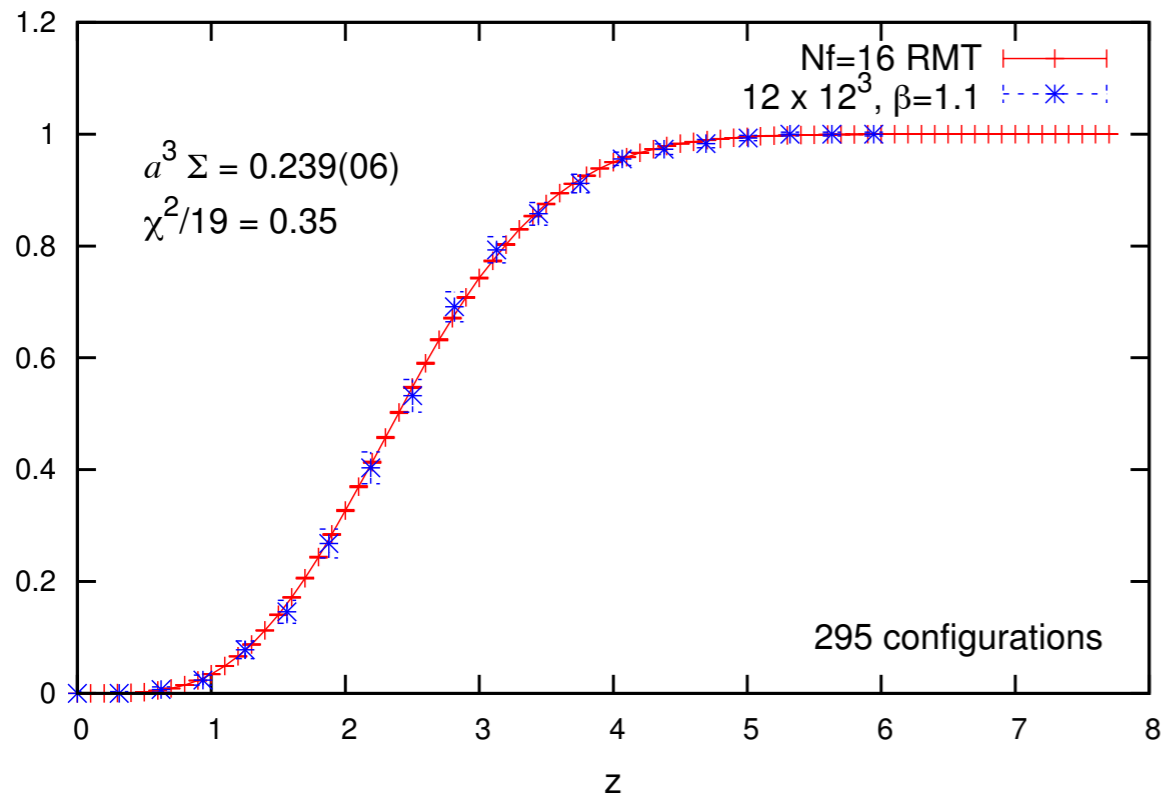
Analysis details

- Study the distribution of the lowest eigenvalue of the Dirac operator at zero topological charge.
- Compare to Random Matrix Theory.
 - ➔ Dyson index = 4.
- Two staggered flavours.
 - ➔ Eight flavours.
- Two-fold degeneracy from the $SU(2)$ gauge group.
 - ➔ Compare to $N_f=16$ RMT.
- Strong taste breaking effects.
 - ➔ Compare to $N_f=4$ and 8 RMT.

Fits to the RMT predictions

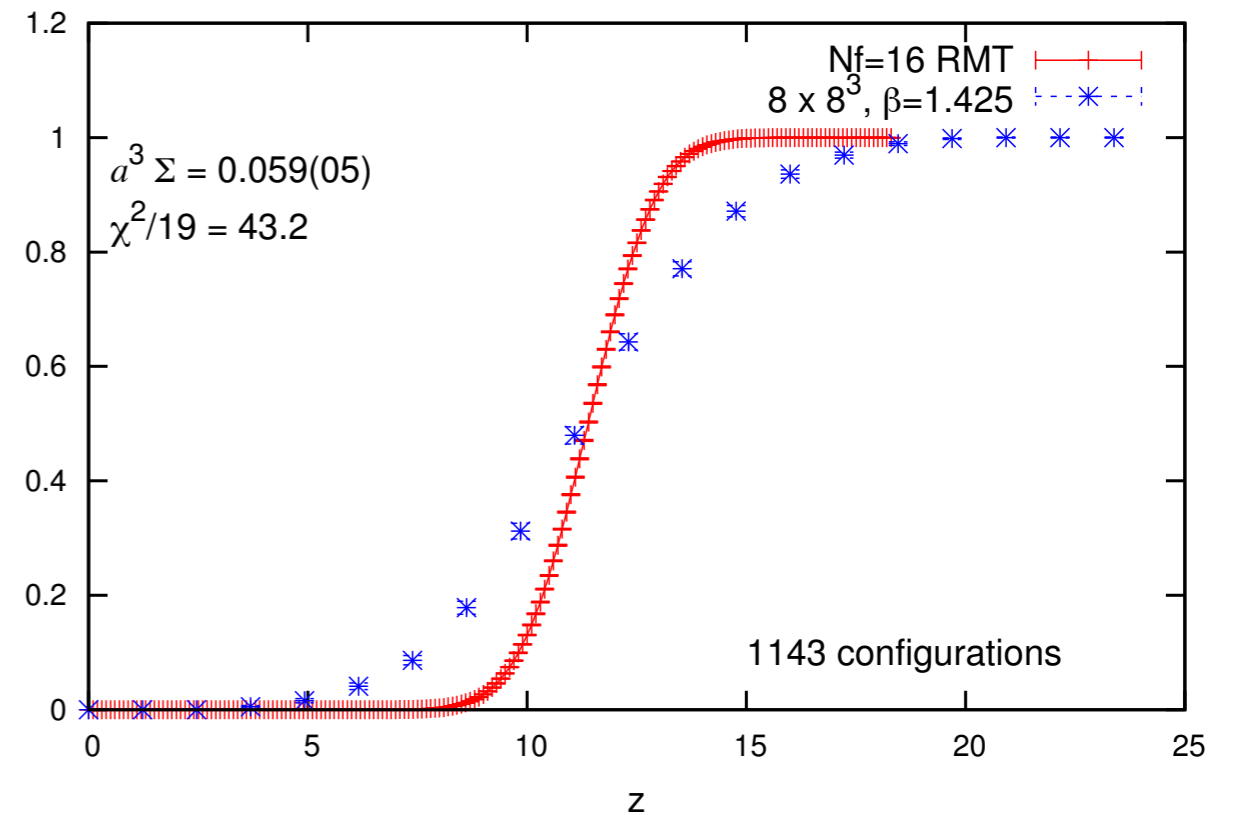
$$z = V \Sigma \sqrt{\lambda_1}$$

integrated eigenvalue distribution: $am=0.010$

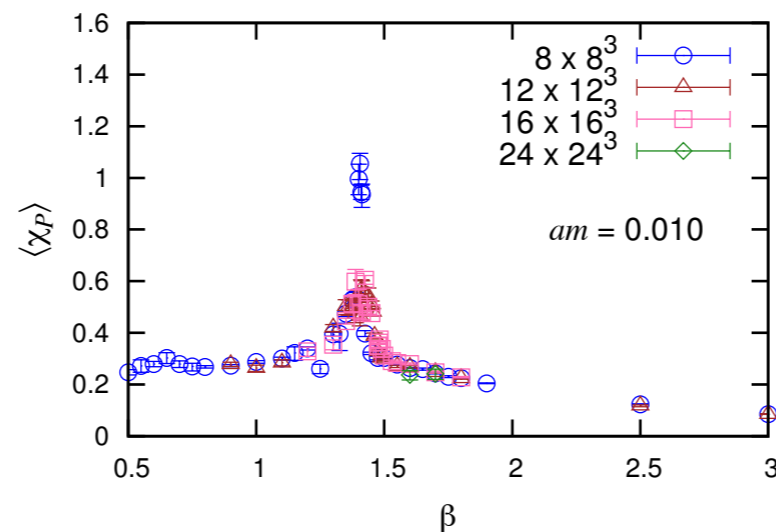


a good fit

integrated eigenvalue distribution: $am=0.010$

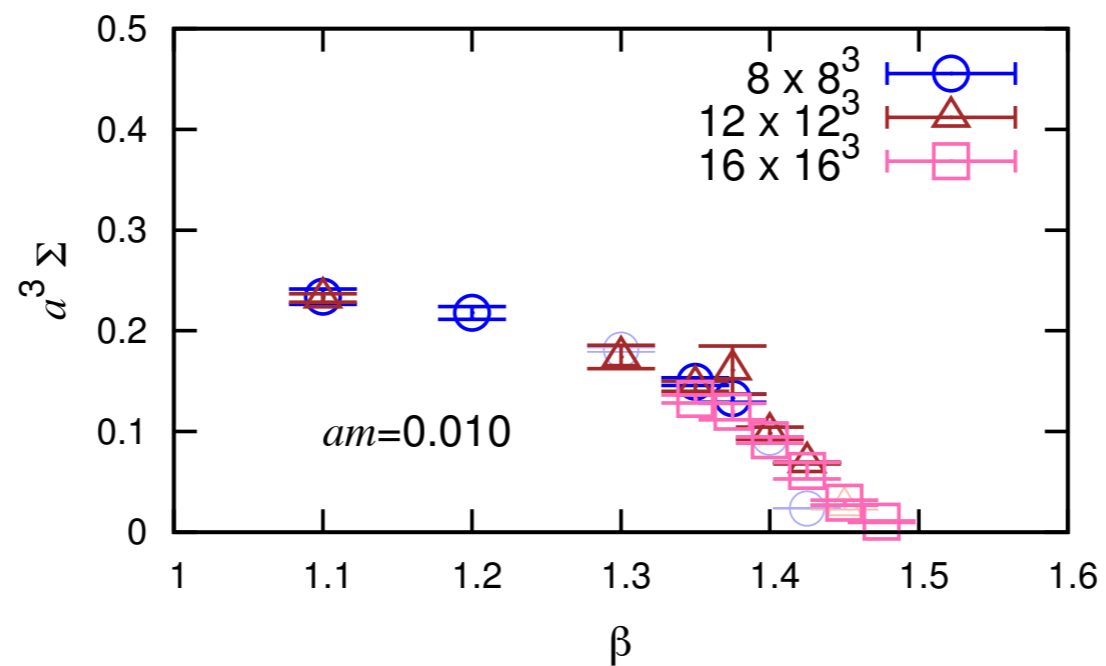


a bad fit
(chiral symmetry restoration?)

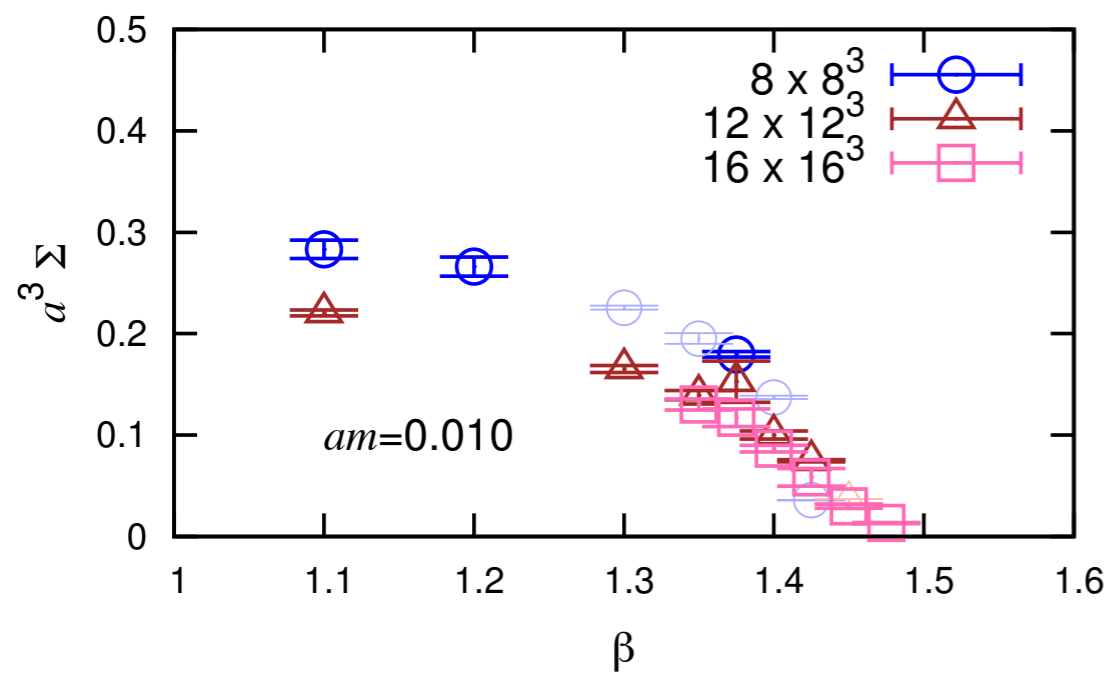


Results: volume dependence

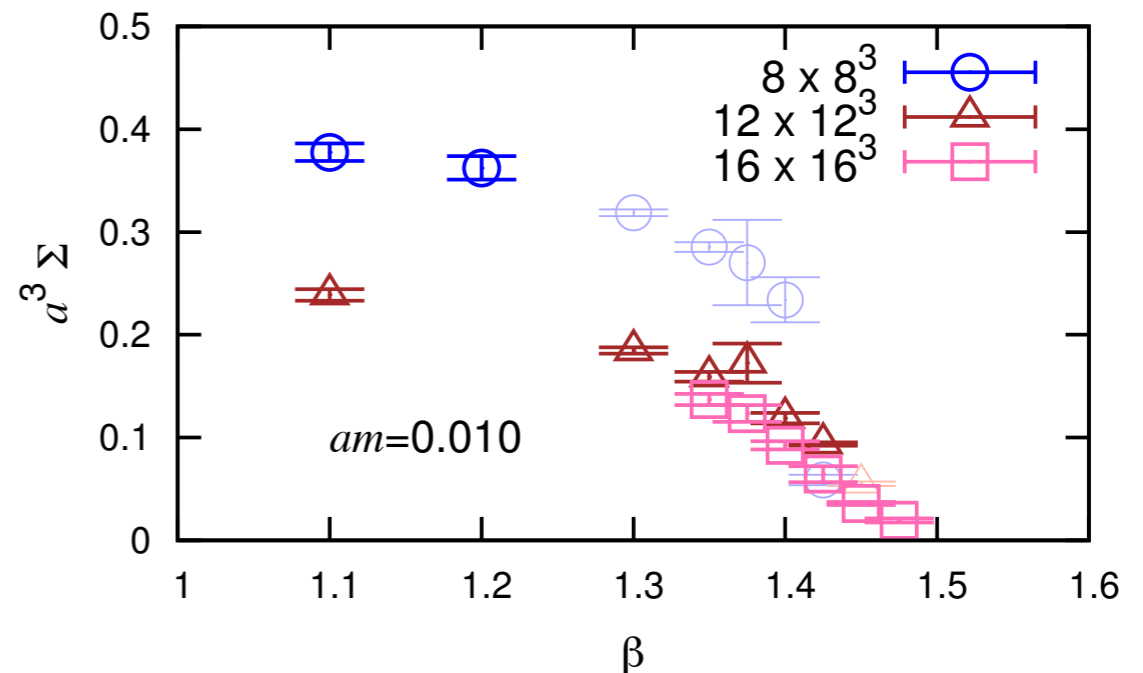
$a^3 \Sigma$, fit with Nf=4 RMT using lowest 1 eigenvalue



$a^3 \Sigma$, fit with Nf=8 RMT using lowest 1 eigenvalue



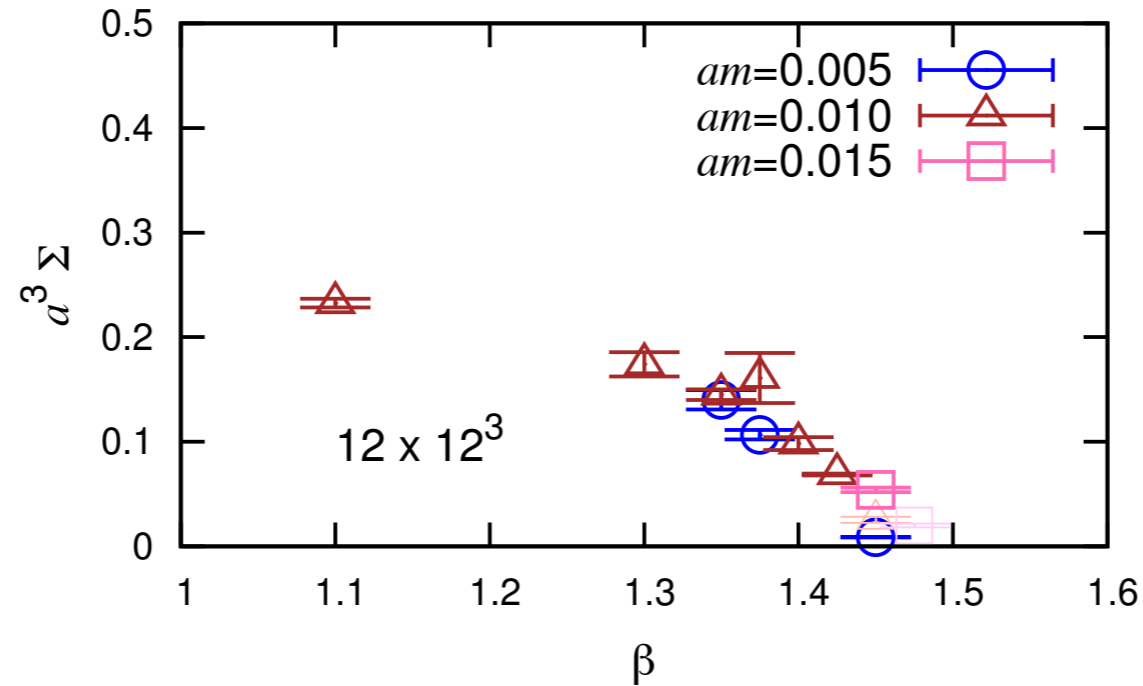
$a^3 \Sigma$, fit with Nf=16 RMT using lowest 1 eigenvalue



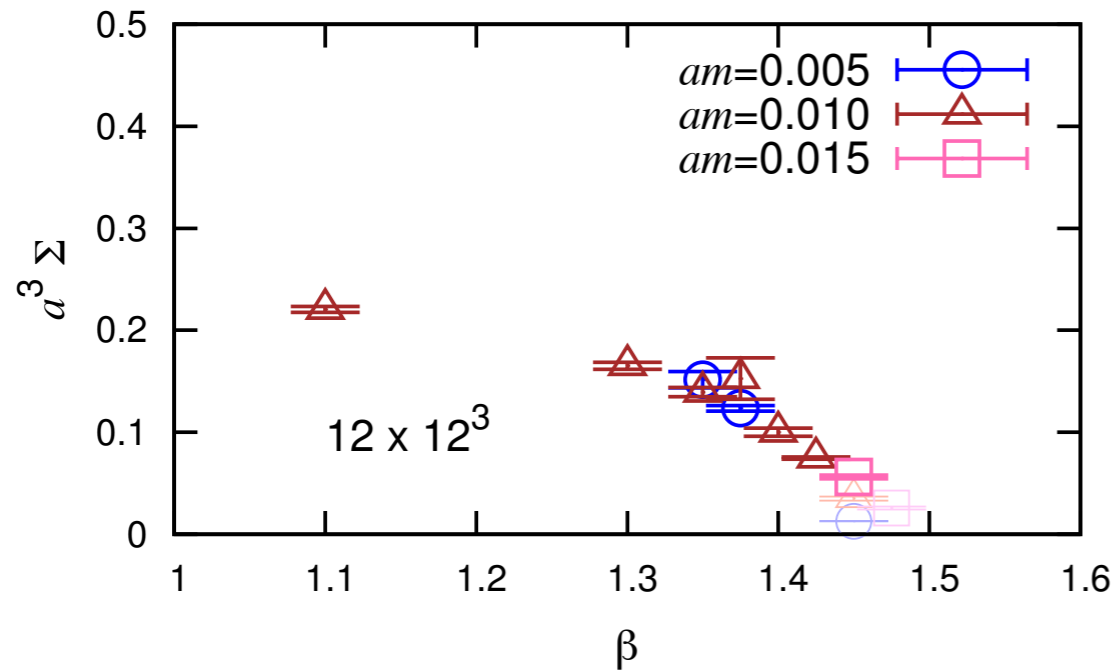
“Thin” symbols indicate bad RMT fits.

Results: mass dependence

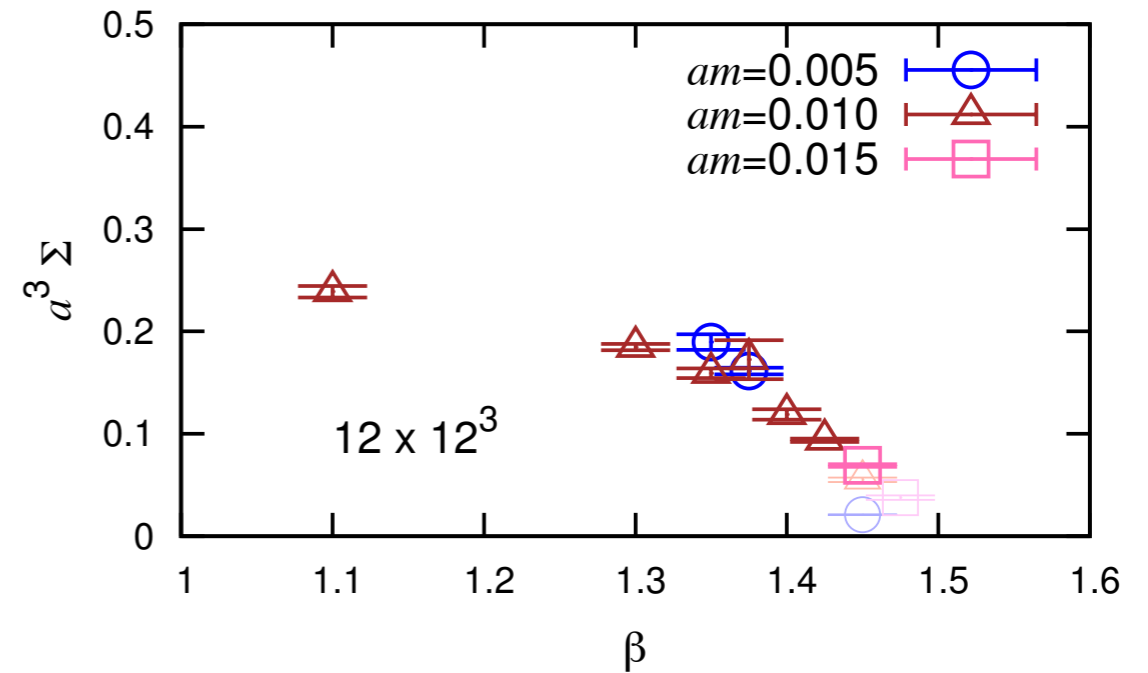
$a^3 \Sigma$, fit with Nf=4 RMT using lowest 1 eigenvalue



$a^3 \Sigma$, fit with Nf=8 RMT using lowest 1 eigenvalue



$a^3 \Sigma$, fit with Nf=16 RMT using lowest 1 eigenvalue



“Thin” symbols indicate bad RMT fits.

Discussion

- For $L = 8$, RMT analysis indicates taste breaking effects seem to be large.
- For $L \geq 12$, RMT analyses at $N_f=4, 8, 12, 16$ lead to consistent results.
- The chiral phase transition occurs at $\beta \sim 1.47$.
 - ➔ No visible mass and volume dependences.
 - ➔ Bulk phase transition.
- Strong taste breaking effects at β as large as 2.
 - ➔ Seems to be irrelevant for the chiral transition.
- The chiral transition does not look like a first order phase transition.