

High statistics analysis of nucleon form factor and charged in lattice QCD

Eigo Shintani (Mainz/AICS)

Hatmut Wittig(Mainz) and collaboration with Mainz-CLS group

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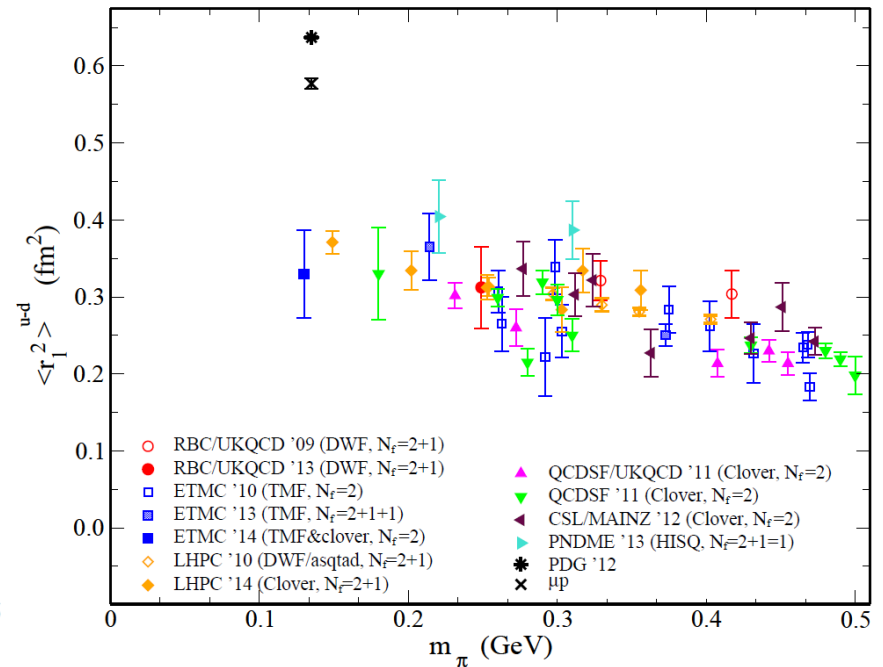
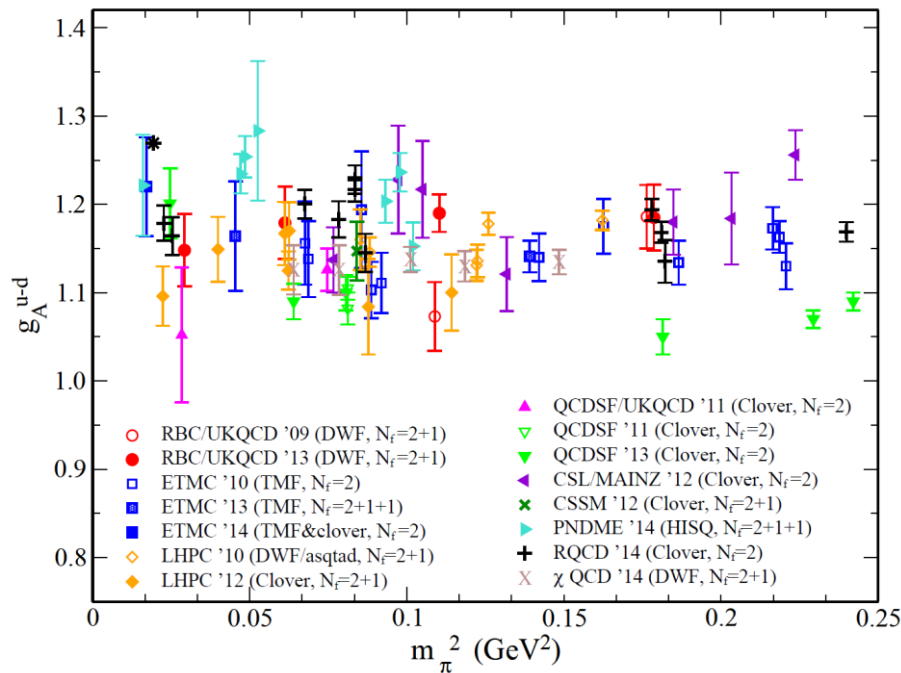
OUTLINE

- ▶ Introduction
- ▶ Error reduction technique
- ▶ Lattice result
 - ▶ Plateau and summation method for axial charge
 - ▶ Scalar and tensor charge
 - ▶ Isovector form factor
 - ▶ Charge radius
- ▶ Summary

1. Introduction

“Puzzle” of nucleon form factor in LQCD

Constantinou, lattice2014



- Many lattice efforts, $N_f=2, 2+1$ (also $2+1+1$) with Wilson, Twisted Wilson, DW, ...
- There is slight tension from experiment, even in different group
 $\Delta g_A \sim 5 - 10\%$, $\Delta r_E^2 \sim 10 - 20\%$
- Careful estimate of systematic uncertainty should be carried out.

1. Introduction

Computation of matrix element

► 2pt, 3pt function

$$\langle 0|\mathcal{N}(t)\mathcal{N}^\dagger(0)|0\rangle = |\langle 0|\mathcal{N}|N\rangle|^2 e^{-m_N t} + \underbrace{|\langle 0|\mathcal{N}|N'\rangle|^2 e^{-m'_N t}}_{\text{First excited state contamination}} + \dots$$

First excited state contamination

$$\begin{aligned} & \langle 0|T\{\mathcal{N}(t_s, 0)J_\mu(t, q)\mathcal{N}^\dagger(0, p)|0\rangle \\ &= \langle 0|\mathcal{N}|N\rangle\langle N|J_\mu|N\rangle\langle N|\mathcal{N}^\dagger|0\rangle e^{-E_N t - m_N(t_s - t)} + \langle 0|\mathcal{N}|N'\rangle\langle N'|J_\mu|N'\rangle\langle N'|\mathcal{N}^\dagger|0\rangle e^{-E'_N t - m'_N(t_s - t)} + \dots \\ &\simeq Z_N(0)Z_N(p)e^{-E_N t - m_N(t_{\text{sep}} - t)} \times \underbrace{[\{G_X, g_A\}]}_{\text{Matrix element of ground state}} + \underbrace{c_1 e^{-\Delta(t_{\text{sep}} - t)} + c_2 e^{-\Delta' t}}_{\text{First excited state contamination}} \end{aligned}$$

Matrix element
of ground state

First excited state contamination
 $\Delta = m'_N - m_N > 0, \Delta' = E'_N - E_N > 0$

- Ground state matrix element is able to be extracted from ratio of 3pt and 2pt function after removing excited state contamination.

Our strategy:

- To reduce statistical error, the all-mode-averaging (AMA) is applied.
- Systematic study of excited state contamination is performed in light pion mass and large volume, $m_\pi L > 4$.

2. Error reduction technique

AMA

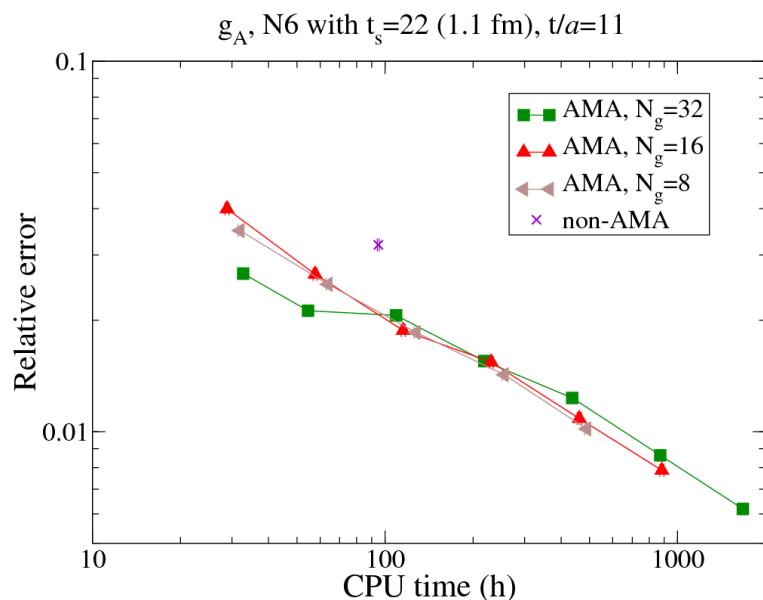
Blum, Izubuchi, ES (2013)

► Reduction of computational cost by using approximation

$$O^{(\text{imp})} = O^{(\text{rest})} + \frac{1}{N_G} \sum_{g \in G} O^{(\text{appx}),g}, \quad O^{(\text{rest})} = O - O^{(\text{appx})}$$

- O : high precision (10^{-10} residue) \Rightarrow expensive but small number of computation
- $O^{(\text{appx})}$: low precision ($\sim 10^{-2}$ residue) \Rightarrow cheap but large number of computation

Error reduction of AMA estimator $O^{(\text{imp})}$ is depending on quality of $O^{(\text{appx})}$.



- Parameter tuning of deflation field N_s which is related to performance of iteration algorithm.
- Cost of computing quark propagator is reduced to 1/5 and less.
- Total speed-up is about factor 2 and more. (depending on lattice size and pion mass)

3. Lattice results (preliminary)

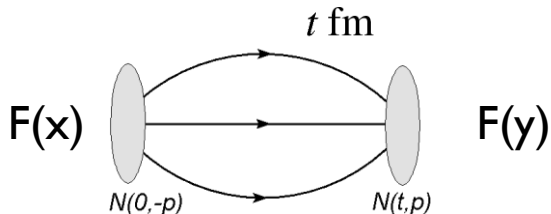
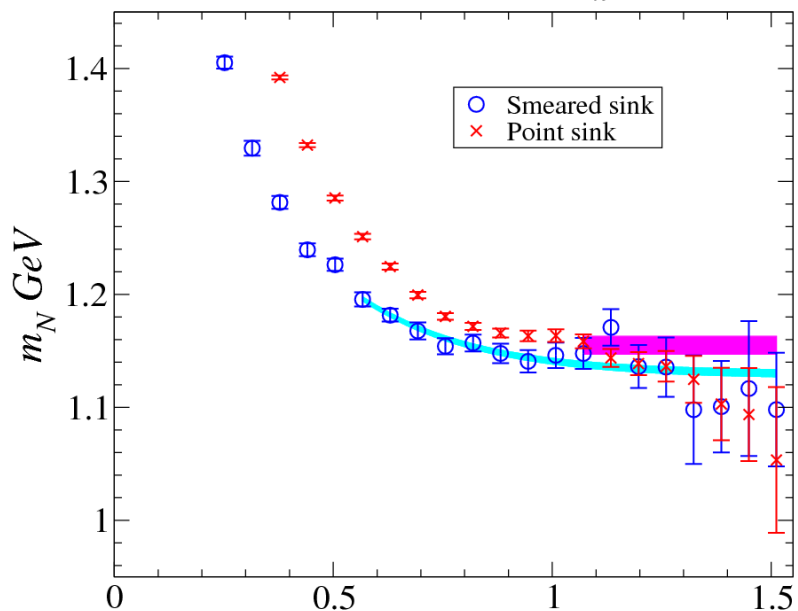
CLS config, $N_f = 2$ Wilson-clover fermion

	Lattice	a (fm)	m_π (GeV)	N_G	t_s (fm)	#conf	#meas(*)
E5	64×32^3 (2.0 fm) ³	0.063	0.456 ($m_\pi L=4.7$)	64	0.82, 0.95, 1.13	~480	~30,000
					1.32	994	63,616
					1.51	1605	102,720
F7	96×48^3 (3.0 fm) ³	0.063	0.277 ($m_\pi L=4.2$)	64	0.82, 0.95, 1.07	250	16,000
				128	1.20, 1.32	250	32,000
				192	1.51	250	64,000
N6	96×48^3 (2.4 fm) ³	0.05	0.332 ($m_\pi L=4.1$)	32	0.9	110	3,520
				32	1.1, 1.3	888	28,416
				32	1.5, 1.7	936	30,272
G8	128×64^3 (4.0 fm) ³	0.063	0.193 ($m_\pi L=4.0$)	80	0.88	184	14,720
				112	1.07	170	19,040
				160	1.26	178	28,480
				160	1.51	179	28,640

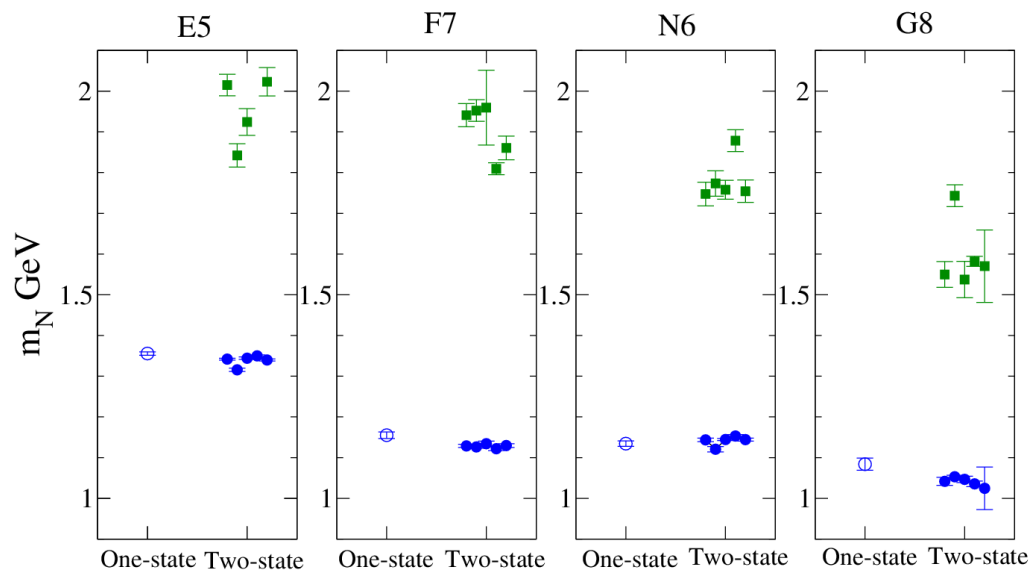
3. Lattice results (preliminary)

Nucleon mass and its excited state

F7: $(3.0 \text{ fm})^3$, $a^{-1}=3.13 \text{ GeV}$, $m_\pi=0.277 \text{ GeV}$



$F(x)$: Jacobian function with APE smearing link.



- The ground-state dominant, $t = 1 \text{--} 1.5 \text{ fm}$.
- Including the excited state, $t = 0.5 \text{--} 1.5 \text{ fm}$
- Fitting function
 - One-state : $Z e^{-m t}$,
 - Two-state : $Z e^{-m t} + Z' e^{-m' t}$
- almost comparable with two fitting results

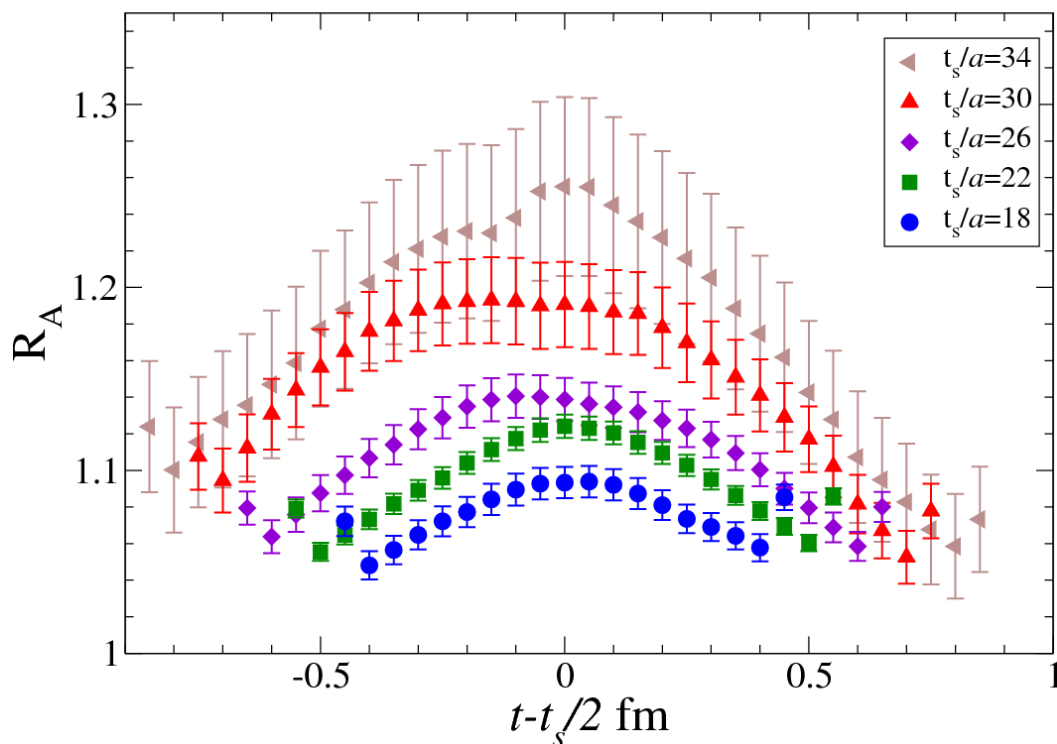
3. Lattice results (preliminary)

Axial charge

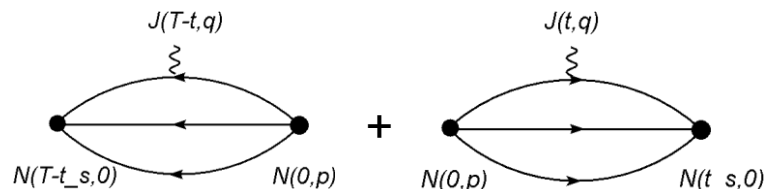
▶ Single ratio of 2pt and 3pt with fixed t_s

$$R_A(t, t_s) = Z \frac{\mathcal{P}\langle 0 | \mathcal{N}(t_s, 0) J_3(t, q) \mathcal{N}^\dagger(0, 0) | 0 \rangle}{\mathcal{P}\langle 0 | \mathcal{N}(t_s, 0) \mathcal{N}^\dagger(0, 0) | 0 \rangle} \simeq g_A + c_1 e^{-\Delta t_s} + c_2 e^{-\Delta'(t_s - t)}$$

N6: $(2.4 \text{ fm})^3$, $a^{-1} = 3.95 \text{ GeV}$, $m_\pi = 0.332 \text{ GeV}$



- Computation of 3pt and 2pt function at zero momentum with spin projection P.
- Signal is regarded as plateau.
- There is significant **size of excited state (2nd and 3rd terms) → fitting including 1st excited state**
- Forward and backward averaging



3. Lattice results (preliminary)

Extraction of g_A

▶ Ground and excited state ansatz

▶ Ground state dominance (plateau method)

$$R_A(t, t_s) = Z \frac{\mathcal{P}\langle 0 | \mathcal{N}(t_s, 0) J_3(t, q) \mathcal{N}^\dagger(0, 0) | 0 \rangle}{\mathcal{P}\langle 0 | \mathcal{N}(t_s, 0) \mathcal{N}^\dagger(0, 0) | 0 \rangle} \simeq g_A, \quad (t_s, t_s - t \gg 1)$$

- Evaluation from constant fitting for t with fixed t_s .
- To suppress the excited state contamination, measurement at large t_s is needed.

▶ First excited state (two-state)

PNDME(2014), RQCD(2014), ...

$$R_A(t, t_s) \simeq g_A + c \left(e^{-\Delta t_s} + e^{-\Delta(t_s - t)} \right)$$

- Δ is mass difference between ground and 1st excited state.

▶ Summation method

Capitani et al. PRD86 (2012)

$$R_A^{\text{sum}}(t_s) = \sum_{t=0}^{t_s} R_A(t, t_s) \simeq a_0 + t_s (g_A + O(e^{-\Delta t_s}))$$

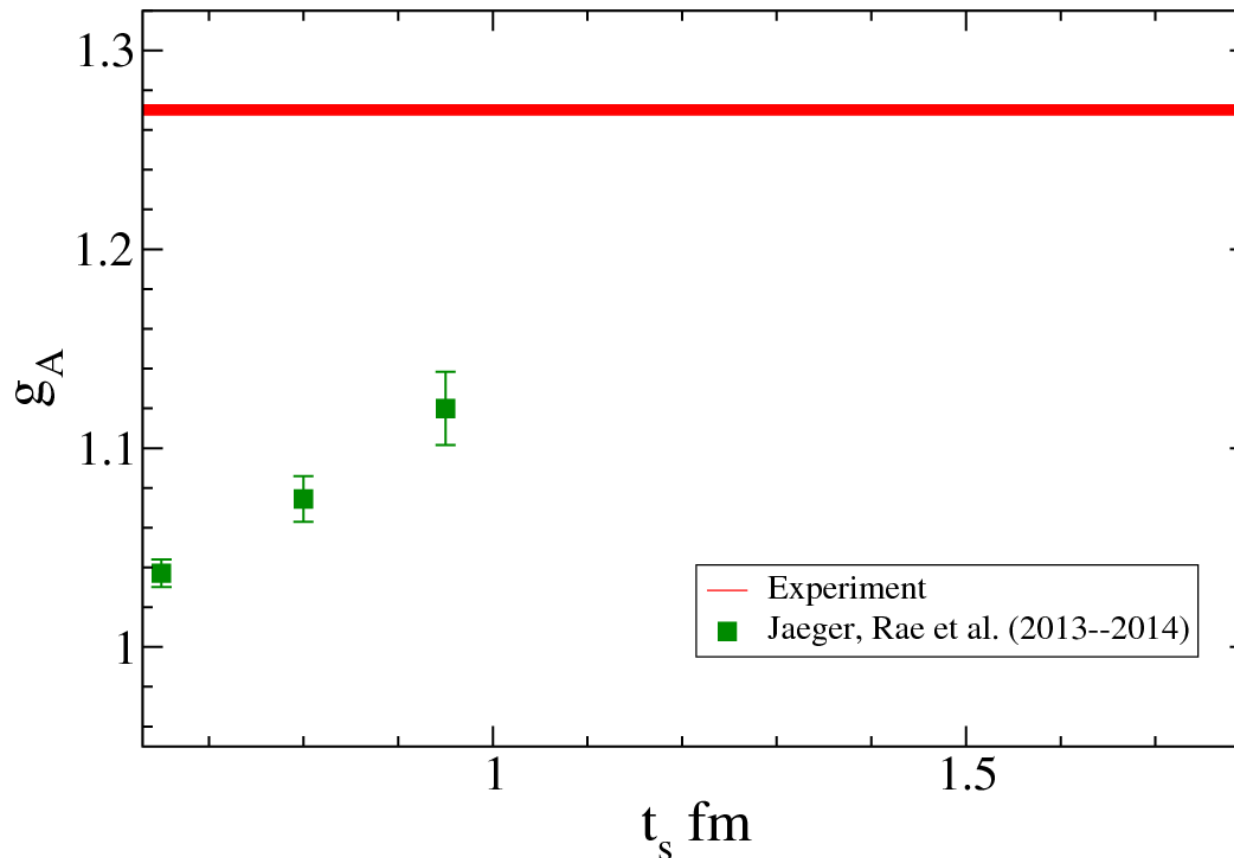
- Using summation in $[0, t_s]$ at fixed t_s , the excited state effect is $\sim O(e^{-\Delta t_s})$
- g_A is given from t_s linear part at $t_s \gg 1$.

3. Lattice results (preliminary): axial charge

Plateau method

► Non-AMA results at $t_s < 1$ fm

N6: $(2.4 \text{ fm})^3$, $a^{-1}=3.95 \text{ GeV}$, $m_\pi=0.332 \text{ GeV}$

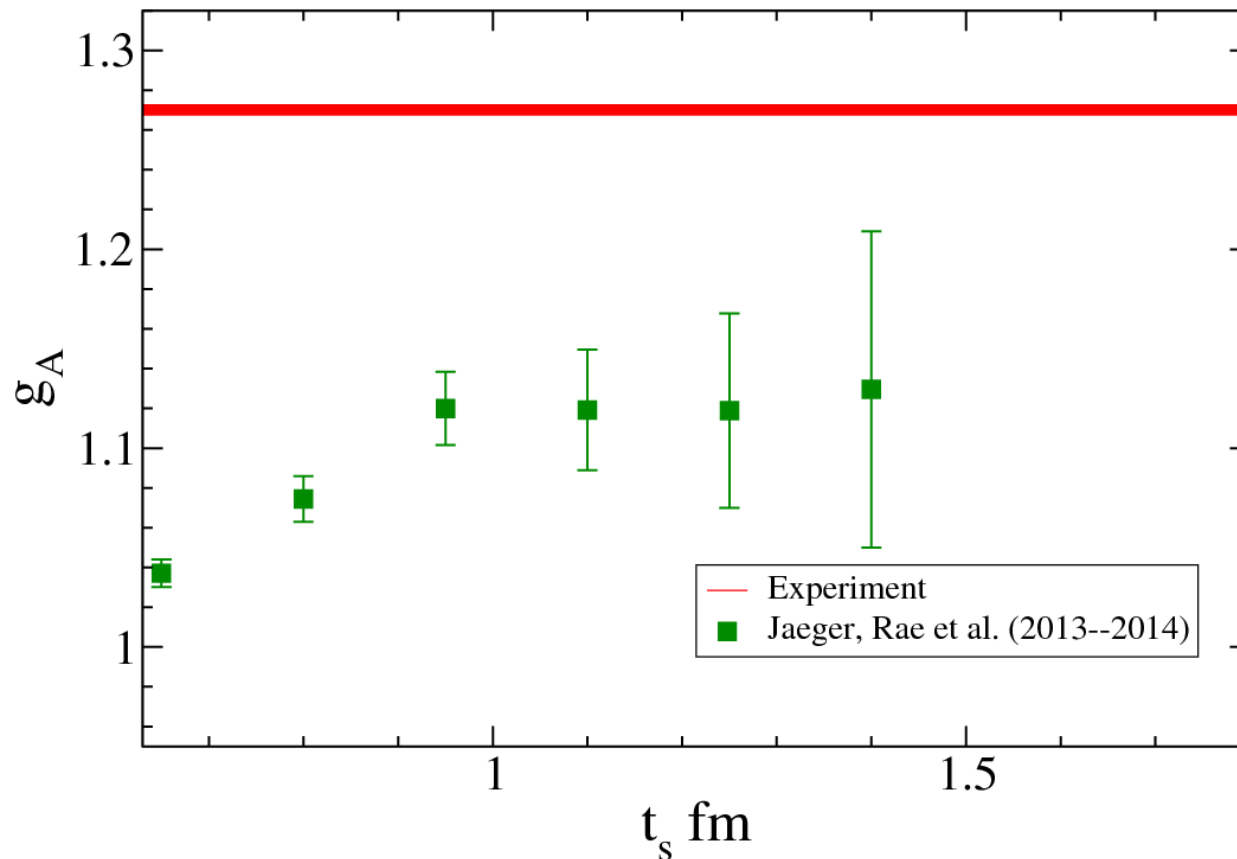


3. Lattice results (preliminary): axial charge

Plateau method

► Non-AMA results at $t_s < 1.5$ fm

N6: $(2.4 \text{ fm})^3$, $a^{-1}=3.95 \text{ GeV}$, $m_\pi=0.332 \text{ GeV}$

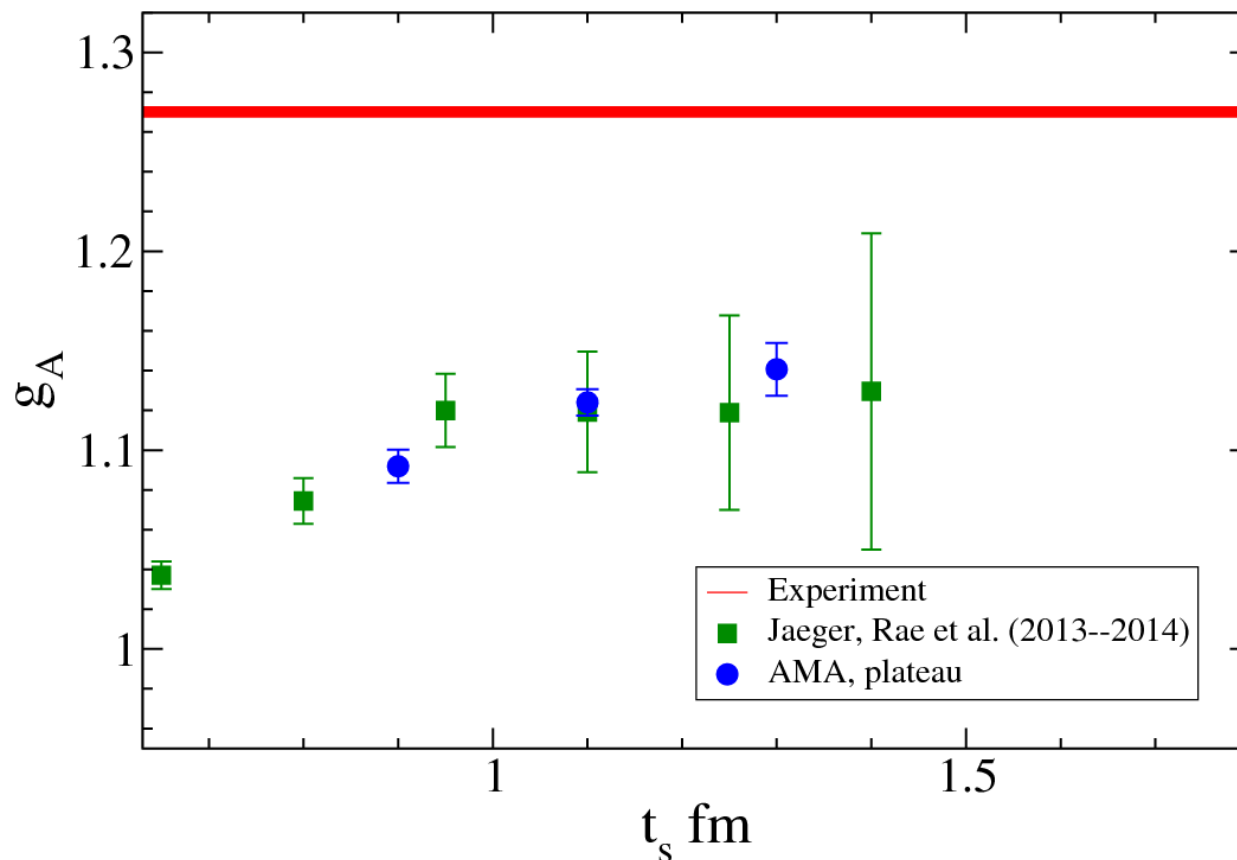


3. Lattice results (preliminary): axial charge

Plateau method

▶ AMA results at $t_s < 1.5$ fm

N6: $(2.4 \text{ fm})^3$, $a^{-1}=3.95 \text{ GeV}$, $m_\pi=0.332 \text{ GeV}$

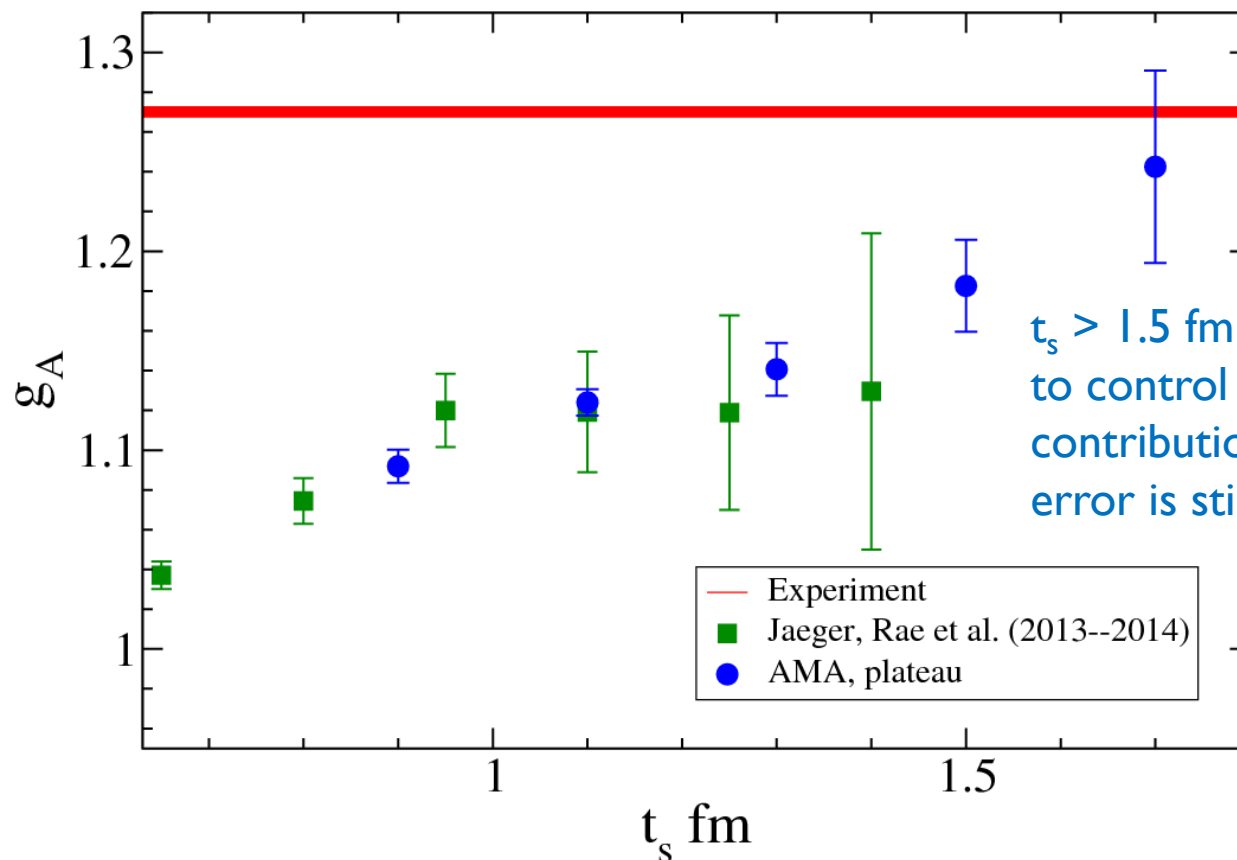


3. Lattice results (preliminary): axial charge

Plateau method

▶ AMA results at $t_s > 1.5$ fm

N6: $(2.4 \text{ fm})^3$, $a^{-1}=3.95 \text{ GeV}$, $m_\pi=0.332 \text{ GeV}$

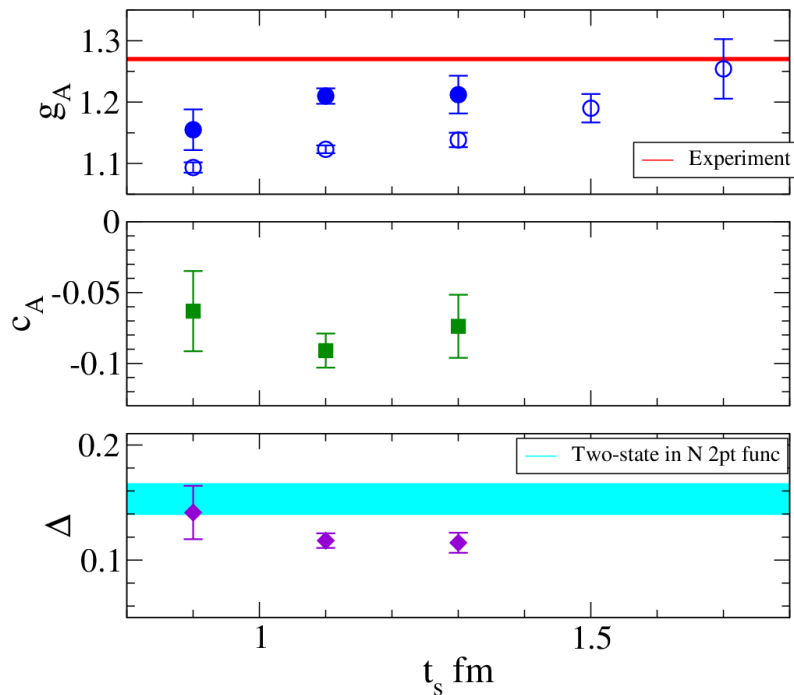


$t_s > 1.5$ fm region is much better to control the excited state contribution, although statistical error is still large.

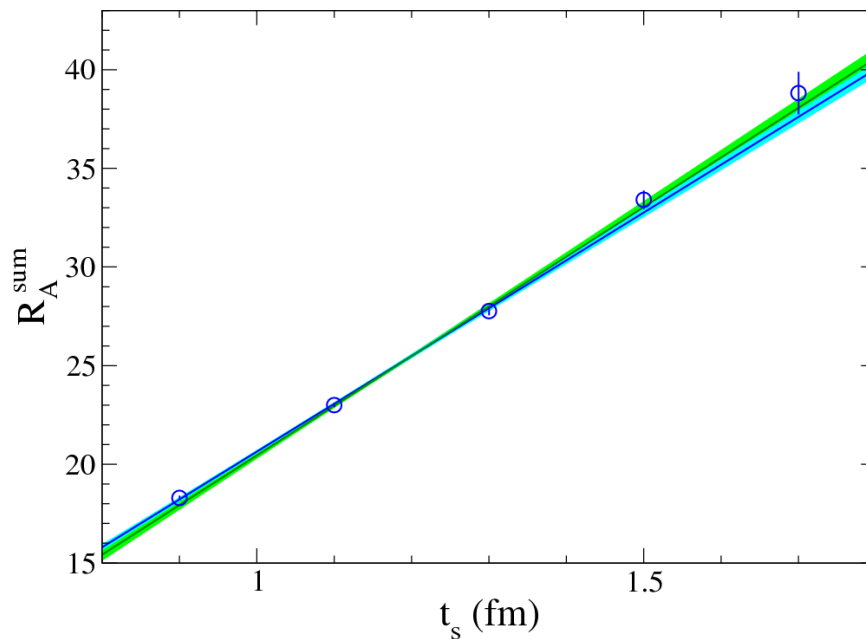
3. Lattice results (preliminary): axial charge

Two state and summation method

N6: $(2.4 \text{ fm})^3$, $a^{-1}=3.95 \text{ GeV}$, $m_\pi=0.332 \text{ GeV}$



N6: $(2.4 \text{ fm})^3$, $a^{-1}=3.95 \text{ GeV}$, $m_\pi=0.332 \text{ GeV}$

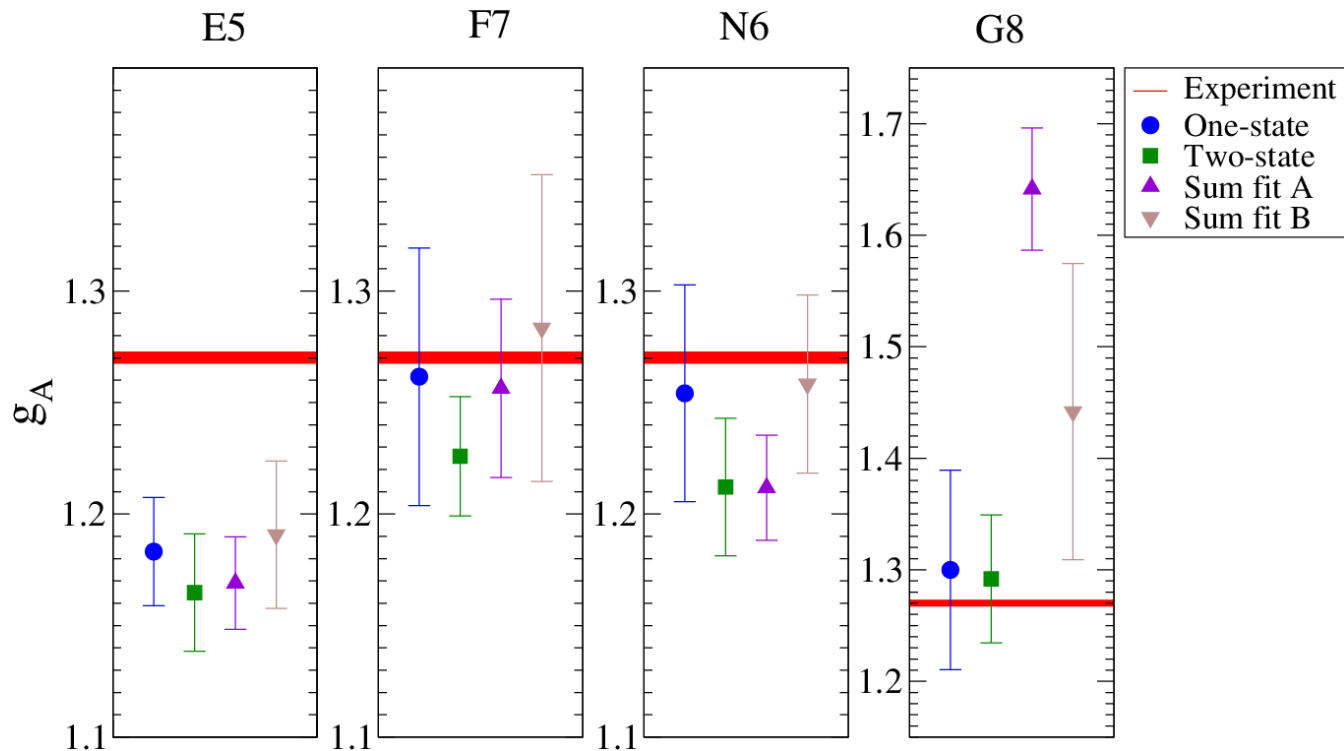


- After correction to excited state, g_A increases, and in agreement with plateau method in $t_s > 1.5 \text{ fm}$.
- Mass difference Δ is compatible with two state fit of 2pt function.

- Linear behavior which is consistent with linear ansatz as expected.
- Comparison between two fitting range:
 $t_s = (\text{fit A})[0.9, 1.7], (\text{fit B})[1.1, 1.7]$
 \Rightarrow estimate of systematic uncertainty



3. Lattice results (preliminary): axial charge Comparison

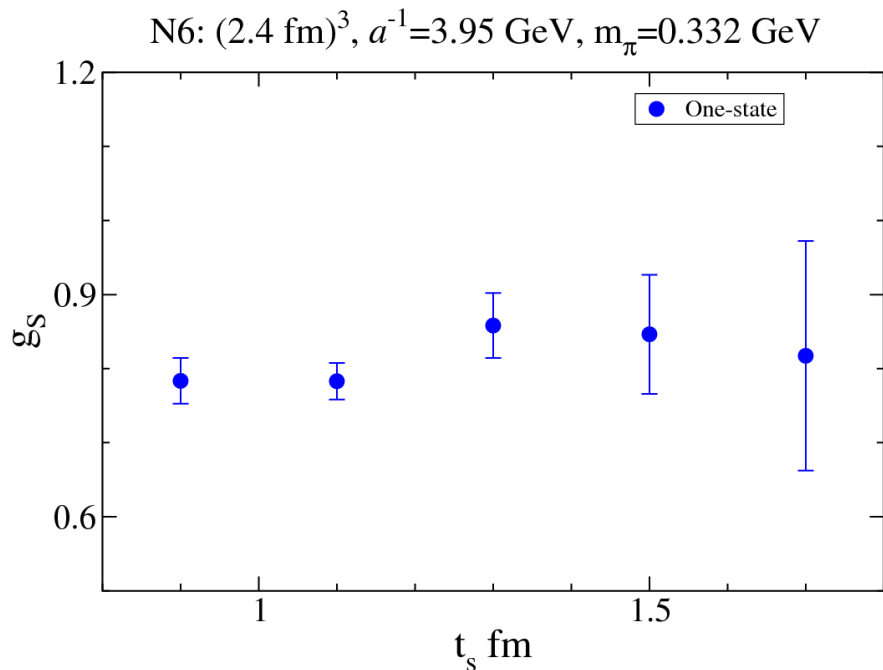


- Four methods provide comparable result except for G8 ensemble at $m_\pi = 0.19$ GeV .
- On G8 summation method with fit A (including short t_s) is discrepancy from others
→ expect systematic uncertainty in linear fit function.
- Finite pion mass effect of g_A is rather mild.

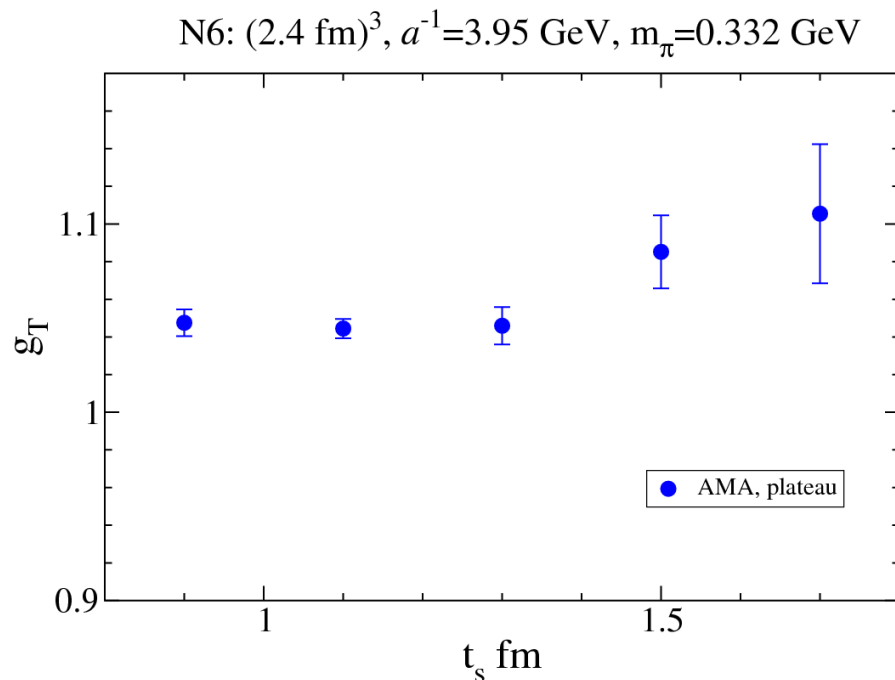
3. Lattice results (preliminary)

Scalar and tensor charge

Scalar (lattice)



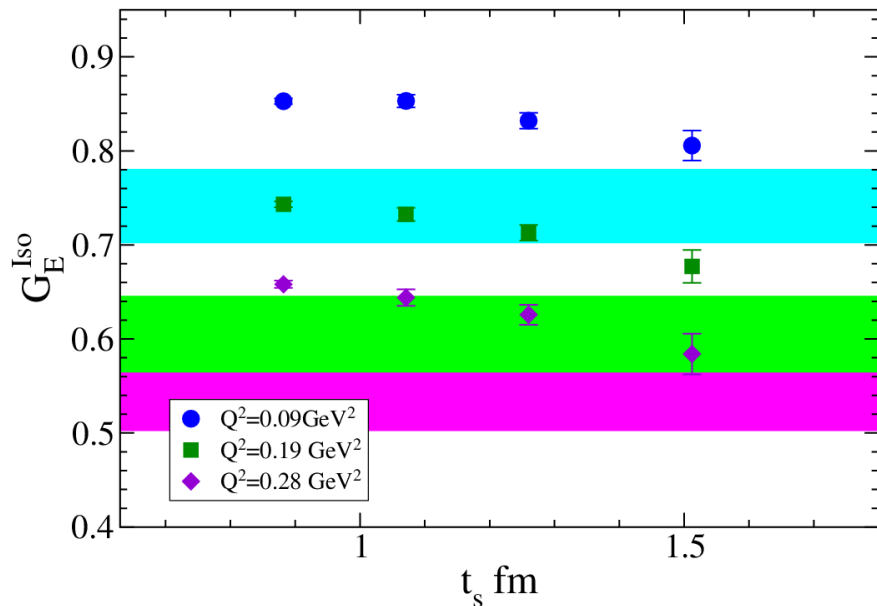
Tensor (lattice)



- There does not appear significant effect of excited state.

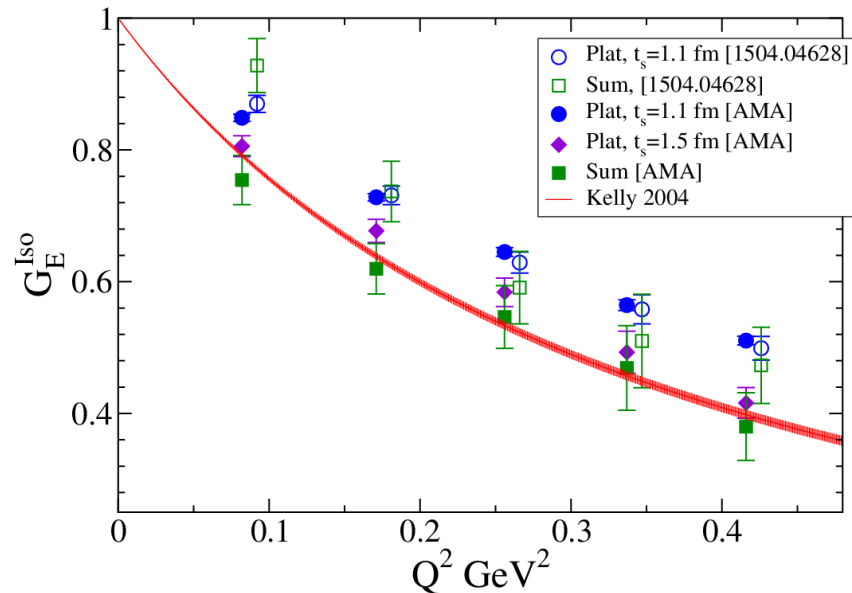
3. Preliminary results: Isovector form factor Analysis at large t_s

G8: $(4.0 \text{ fm})^3$, $a^{-1}=3.13 \text{ GeV}$, $m_\pi=0.193 \text{ GeV}$



- From $t_s > 1 \text{ fm}$, there is still tendency to decrease by $\sim 10\%$.
- Summation method and plateau method at $t_s > 1.5 \text{ fm}$ are compatible.

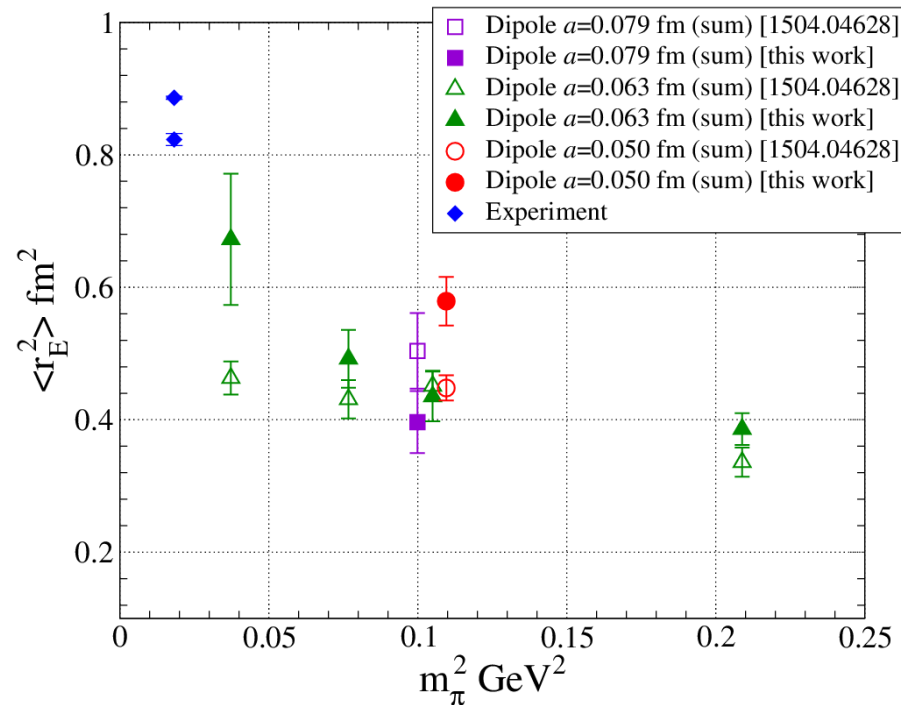
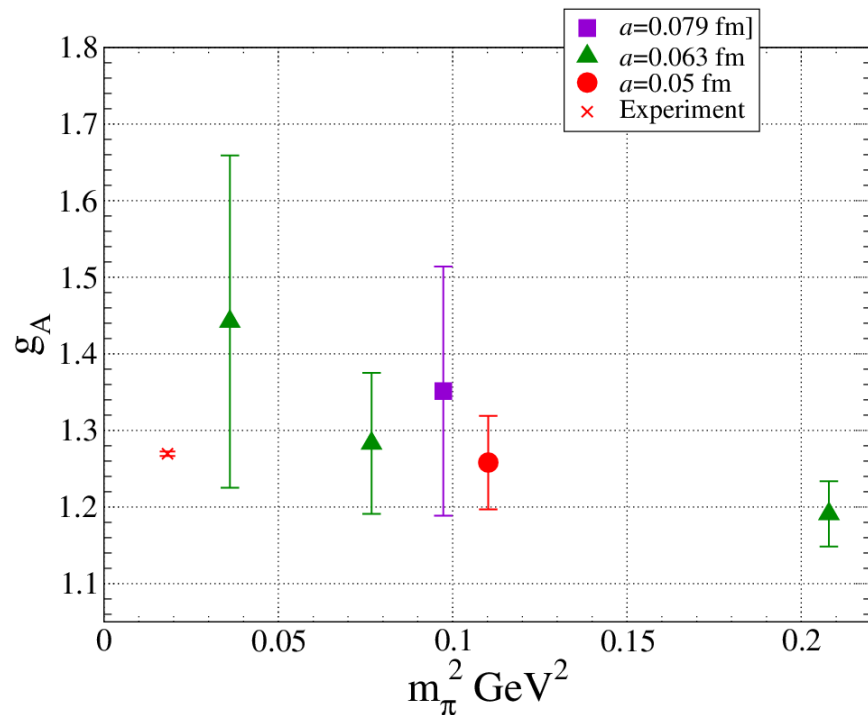
G8: $(4.0 \text{ fm})^3$, $a^{-1}=3.13 \text{ GeV}$, $m_\pi=0.19 \text{ GeV}$



- Comparison with previous work on the same ensemble.
- Large discrepancy between plateau method at $t_s = 1.1 \text{ fm}$ and 1.5 fm , due to excited state contamination.
- Approaching to experimental value.

3. Preliminary results

Axial charge and charge radius



- Analysis of axial charge and charge radius with large t_s up to 1.7 fm.
- Result has still large statistical error, even though statistics $O(10^5)$ is used.
- In $t_s = 1.1$ fm, there is still unsuppressed excited state effect, which may be one of the reason for large discrepancy from experiment \Rightarrow need more than 1.5 fm.
- Axial charge may not have strong m_π dependence, but $\langle r_E \rangle$ may have.

4. Summary

Summary

- ▶ High statistics calculation of nucleon form factor is performed in $N_f=2$ Wilson-clover at $Lm_\pi > 4$ with $m_\pi = 0.19--0.46$ GeV.
- ▶ All-mode-averaging technique is working well for reduction of statistical error.
- ▶ $t_s > 1.5$ fm is required for small contribution of excited state contamination in axial charge and (iso)vector form factor.
- ▶ Axial charge and charge radius are approaching to experimental value.
- ▶ Feasible study for application to $N_f = 2+1$ CLS configurations with open boundary condition. Tim Harris, talk on 18 July 10:00

Thank you for your attention.

Isvector form factor

▶ Ratio with momentum transition

$$R_G(t, t_s) = Z \frac{\mathcal{P}\langle 0 | \mathcal{N}(t_s, p_1) J_\mu(t, q) \mathcal{N}^\dagger(0, p_0) | 0 \rangle}{\mathcal{P}\langle 0 | \mathcal{N}(t_s, p_0) \mathcal{N}^\dagger(0, p_0) | 0 \rangle} K(p_1, p_0) \simeq G_X + d_1 e^{-\Delta t_s} + d_2 e^{-\Delta'(t_s - t)}$$

$$K(p_1, p_0) = \sqrt{\frac{C_{2\text{pt}}^{\text{lc}}(p_1, t_s - t) C_{2\text{pt}}^{\text{sm}}(p_0, t) C_{2\text{pt}}^{\text{lc}}(p_0, t_s)}{C_{2\text{pt}}^{\text{lc}}(p_0, t_s - t) C_{2\text{pt}}^{\text{sm}}(p_1, t) C_{2\text{pt}}^{\text{lc}}(p_1, t_s)}}$$

- The ratio consists of 3pt and 2pt, with combination of local “lc” and smeared “sm” sink.
- Matrix element with Sachs form factor

$$\langle N(\vec{p}_1) | J_\mu | N(\vec{p}_0) \rangle = \bar{u}(p_1) \left[F_1^v(q^2) \gamma_\mu + F_2 q_\nu \sigma_{\mu\nu} / 2m_N \right] u_N(p_0)$$

$$G_E = F_1 - \frac{q^2}{4m_N^2} F_2, \quad G_M = F_1 + F_2$$

- Form factor G_X as a function of q^2 , $q = p_1 - p_0$, in which $p_1 = (0, m_N)$ $p_0 = (p, E)$ are used.
- Systematic study of excited state contamination with plateau and summation method is necessary.

Improvement of standard deviation:

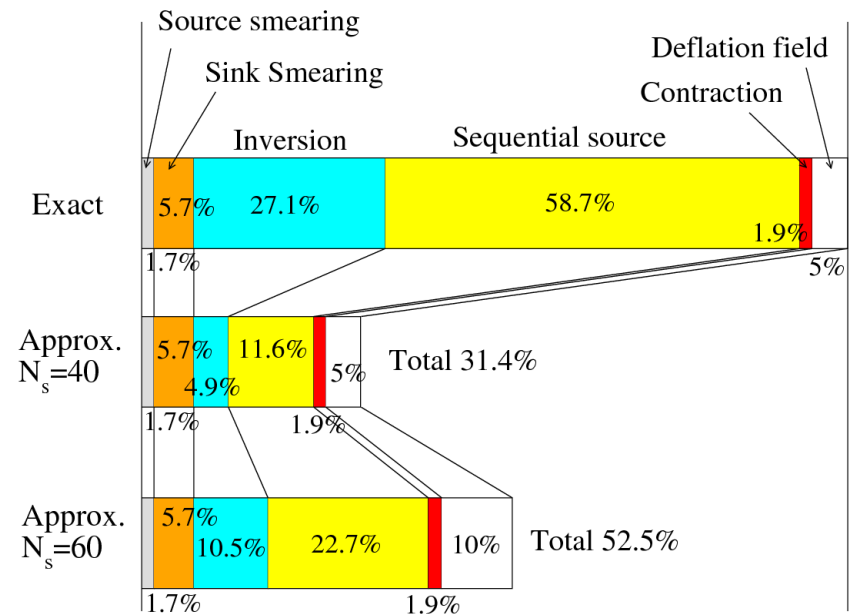
$$\frac{\sigma^{\text{imp}}}{\sigma} \simeq \sqrt{\frac{1}{N_G} + 2(1-r) + \frac{1}{N_g^2} \sum_{g \neq g'} r_{gg'}}$$

$$r = \frac{\langle \Delta O \Delta O^{(\text{appx})} \rangle}{\sigma \sigma^{(\text{appx})}} \quad r_{gg'} = \frac{\langle \Delta O^{(\text{appx}),g} \Delta O^{(\text{appx}),g'} \rangle}{\sigma^{(\text{appx}),g} \sigma^{(\text{appx}),g'}}$$

r : correlation between O and $O^{(\text{appx})}$

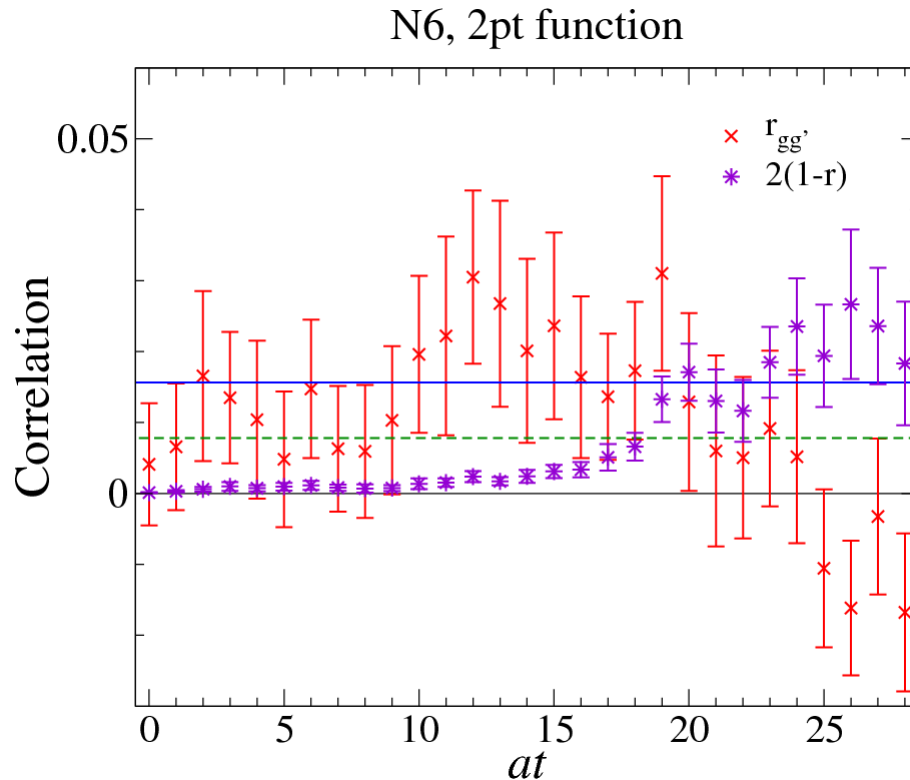
$r_{gg'}$: correlation between $O^{(\text{appx}),g}$ and $O^{(\text{appx}),g'}$

- $O^{(\text{appx})}$ has several tuning parameters to control of r and $r_{gg'}$
e.g. stopping condition, deflation field, source location



Performance test of AMA

► Correlation



Expected error reduction in AMA:

$$\frac{\sigma^{\text{imp}}}{\sigma} \simeq \sqrt{\frac{1}{N_G} + 2(1-r) + \frac{1}{N_g^2} \sum_{g \neq g'} r_{gg'}}$$

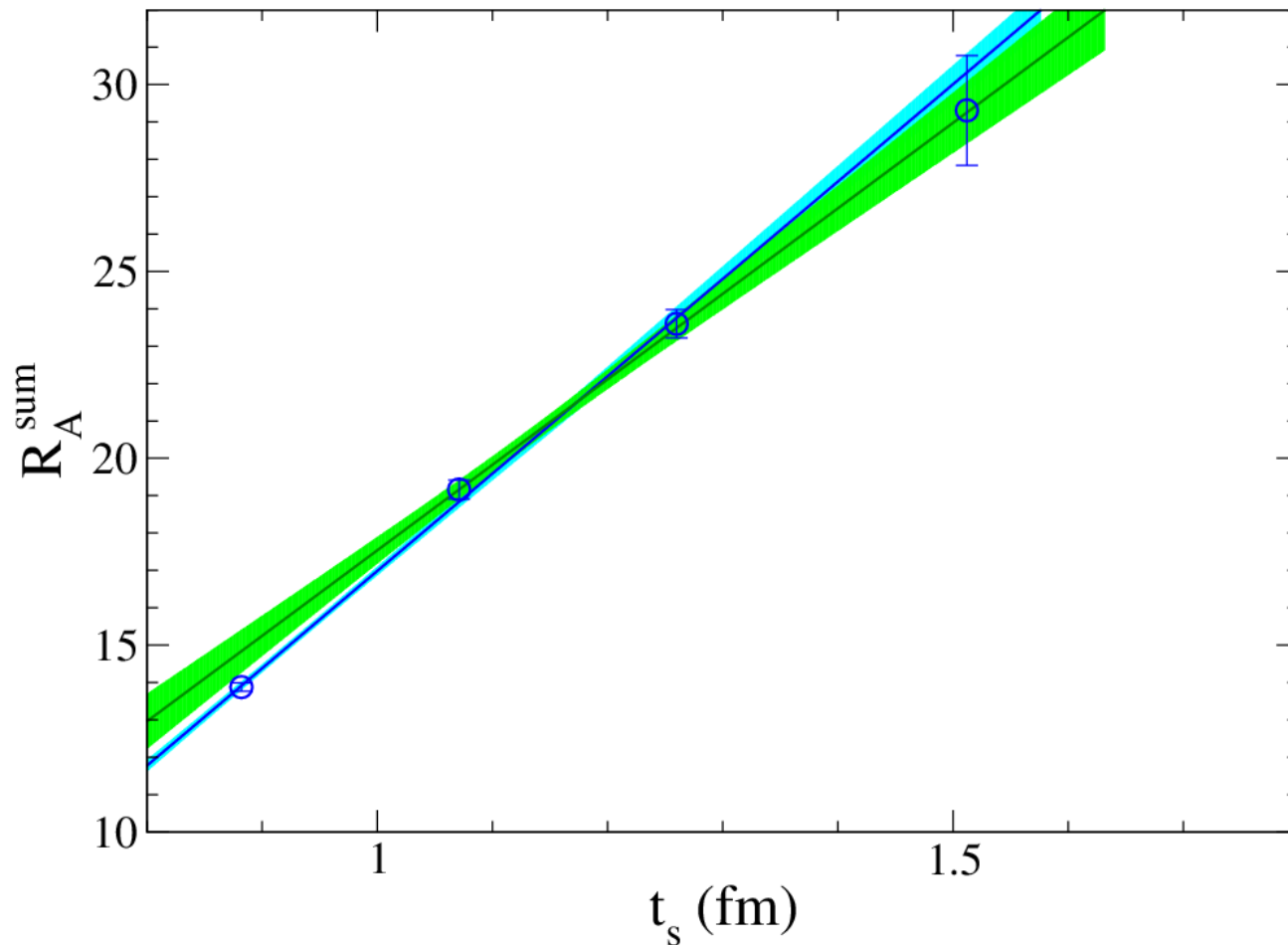
← $1/N_G = 1/64$

← $1/N_G = 1/128$

- $r_{gg'}$: correlation between $\mathcal{O}^{(\text{appx})}$ with g and g' transformation.
- $2(1-r)$: correlation between $\mathcal{O}^{(\text{appx})}$ and \mathcal{O} .
- At $t \sim 24$, size of correlation is similar to $1/N_G$, \Rightarrow maximum point to reduce error

Summation method on G8

G8: $(4.0 \text{ fm})^3$, $a^{-1}=3.13 \text{ GeV}$, $m_\pi=0.19 \text{ GeV}$



t dependence of G_E

