Understanding the problem with logarithmic singularities in the complex Langevin method

#### Shinji Shimasaki (KEK)

based on the work in collaboration with Jun Nishimura (KEK, SOKENDAI) Phys. Rev. D 92, 011501(R) (2015) (arXiv:1504.08359 [hep-lat])

## Complex Langevin Method (CLM) [Parisi 83][Klauder 84]

- A promising method for evading the sign problem in the path integral with a complex weight
- $\cdot\,$  e.g.) finite density QCD, real time dynamics,  $\cdots\,$
- sometimes works, sometimes fails
- recent progress : better understanding of the cause of the failure

## Recent understanding

CLM fails when

- the probability distribution has large excursions
  into the imaginary direction [Aarts, Seiler, Stamatescu 09]
  [Aarts, James, Seiler, Stamatescu 09 11]
- the action involves "log" :  $S = \log \Delta(\phi) + \cdots$ 
  - [Mollgaard, Splittorff 13][Greensite 14]
  - · CLM gives wrong results when the phase of  $\Delta(\phi)$  rotates frequently during the Langevin process
  - why? ambiguity of log branch ? … not clear

## Our understanding

#### [Nishimura-SS 15]

- $\cdot \,$  log branch is not the cause of the problem
- Rather, the singularity in the drift term causes the problem
- Indeed, the problem occurs even when the action involves non-logarithmic singularity

(ex) 
$$S = \frac{\beta}{(x+i\alpha)^2} + \frac{x^2}{2}$$



gauge cooling is useful here as well (Nagata's talk)

#### [Nishimura-SS 15]

$$Z = \int_{-\infty}^{\infty} dx (x + i\alpha)^p e^{-x^2/2}$$
 (  $\alpha \in \mathbf{R}$ ,  $p \in \mathbf{R}$ )

- action  $S = -\log(x + i\alpha)^p + x^2/2$  "log" or "Log"
- drift term  $-\frac{dS}{dx} = \frac{p}{x + i\alpha} x$
- Langevin eq.  $x(t) \rightarrow z(t) = x(t) + iy(t)$

$$\frac{dz(t)}{dt} = \frac{p}{z(t) + i\alpha} - z(t) + \eta(t)$$

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Defining the action is not necessary in formulating CLM

[Nishimura-SS 15]

$$Z = \int_{-\infty}^{\infty} dx \, w(x)$$

· drift term u

$$w(x)^{-1}\frac{dw(x)}{dx}$$

· Langevin eq.  $x(t) \rightarrow z(t) = x(t) + iy(t)$ 

$$\frac{dz(t)}{dt} = w(z)^{-1}\frac{dw(z)}{dz} + \eta(t)$$

• The problem is whether the ensemble average can reproduce the path integral with the complex weight w(x)

#### [Nishimura-SS 15]

$$Z = \int_{-\infty}^{\infty} dx (x + i\alpha)^p e^{-x^2/2}$$
 (  $\alpha \in \mathbf{R}$ ,  $p \in \mathbf{R}$ )

· Langevin eq.

$$\frac{dx(t)}{dt} = \operatorname{Re}\frac{p}{z(t) + i\alpha} - x(t) + \eta(t)$$
$$\frac{dy(t)}{dt} = \operatorname{Im}\frac{p}{z(t) + i\alpha} - y(t)$$

· compute the average of  $z^n$  in CLM

 $Z = \int_{-\infty}^{\infty} dx (x + i\alpha)^p e^{-x^2/2}$ 

p = 50

### Result (p=50)











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For CLM to work :

probability distribution associated with the Langevin process  $\int dx dy \mathcal{O}(x+iy) P(x,y) = \int dx \mathcal{O}(x) e^{-S}$ 

 We can show this equality if partial integration on the complex plane is allowed

[Aarts, Seiler, Stamatescu 09][Aarts, James, Seiler, Stamatescu 09 11]

• For partial integration to be justified :

•

#### The distribution should have a fast fall-off

[Aarts, Seiler, Stamatescu 09][Aarts, James, Seiler, Stamatescu 09 11]

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· The drift term should not have singularities

[Nishimura-SS 15]

$$\int_{-\infty}^{\infty} dx \frac{1}{x} \frac{d}{dx} e^{-x^2} \neq \int_{-\infty}^{\infty} dx \frac{1}{x^2} e^{-x^2}$$

• For partial integration to be justified :

· The distribution should have fast fall-off

OK, since the drift term =  $-z + \cdots$ 

· The drift term should not have singularities

The drift term is singular at  $z = -i\alpha$  ( $x = 0, y = -i\alpha$ ) ... OK, if the distribution is zero near the singularity.

## Result (p=50)



#### Radial distribution



## Non-logarithmic case

$$Z = \int_{-\infty}^{\infty} dx e^{-\frac{1}{(x+i\alpha)^2} - x^2/2}$$

• action 
$$S(z) = \frac{1}{(z+i\alpha)^2} + z^2/2$$

no log finite for real x

$$\cdot \text{ drift} = \frac{2}{(z+i\alpha)^3} - z$$

singular at  $z = -i\alpha$ 

## Result (non-log)

non-Log, p=2,  $\beta$ =1 1 exact — CLM ⊢→ 0.5 ×> ₩ 0  $\overset{'}{\Theta} \overset{\Theta}{\Theta} \overset{O}{\Theta} \overset{O}{O} \overset{O}$ -0.5 -1 0.5 2 1.5 0 1 α non-Log, p=2,  $\beta$ =1 exact — CLM ⊢→ 4  $\Theta \Theta \Theta \Theta \Theta \Theta$ 2 🖡  $Re < x^2 >$ 0 -2 -4 0.5 1.5 2 0 1

α



## Result (non-log)



## Summary

- Condition for CLM to work
- = probability distribution has a fast fall-off and is practically zero near the singularity
- Both are related to the justification of the partial integration
- "gauge cooling" should be useful in curing the singularity problem, too !!
- The new insights have great impact on applications to finite density QCD at low T with light quarks. (Nagata's talk)