# Beyond Standard Model: Lattice Calculations at the Energy Frontier

# Anna Hasenfratz University of Colorado Boulder

Lattice 2015, Kobe

I would like to thank the organizers for an amazing Lattice 2015 Conference, and for giving me the opportunity to review

"Lattice studies for BSM models, especially the strong dynamics"

I also want to thank everyone who sent me details of their work, answered my questions, and gave excellent talks, posters. I learned a lot.

I have to apologize to everyone whose work I cannot mention or review appropriately - and that is pretty much everyone.

## Strong Yukawa models : 1989 and 2014

#### Massive fermions without fermion bilinear condensates

Venkitesh Ayyar and Shailesh Chandrasekharan Department of Physics, Box 90305, Duke University, Durham, North Carolina 27708, USA

#### NON-PERTURBATIVE STUDY OF THE STRONGLY COUPLED SCALAR-FERMION MODEL \*

#### Anna HASENFRATZ

Supercomputer Computations Research Institute and Physics Department, The Florida State University, Tallahassee, FL 32306-4052, USA

and

Thomas NEUHAUS Institut für Theoretische Physik, Universität Bielefeld, D-4800 Bielefeld, Fed. Rep. Germany

Received 14 June 1988; revised manuscript received 18 January 1989

# In Memoriam



# Thomas Neuhaus 1956-2015

http://www.fz-juelich.de/SharedDocs/Meldungen/IAS/JSC/EN/2015/2015-06-obituary-thomas-neuhaus.html?nn=992970

# Strong Yukawa models : 1989

Generic phase diagram (scalar hopping vs Yukawa coupling):



Two paramagnetic phases  $\langle \Phi \rangle = 0 \quad \langle \bar{\psi}_n \psi \rangle = 0$ PMW : usual perturbative phase PMS : strong coupling phase with massive fermions :

PMS phase: mass generation without fermion bilinear condensate (but fermions decouple in any continuum limit)

# Strong Yukawa models : 2015

Use 2 copies of (reduced) staggered fermions

$$S = S_0 - \kappa \sum_{x,\hat{\alpha}} \sigma_x \sigma_{x+\hat{\alpha}} - Y \sum_{i=1,2} \sum_x \sigma_x \overline{\psi}_{x,i} \psi_{x,i}$$

Simulations at  $\kappa = 0$ :

$$S = S_0 - U \sum_{x} \left\{ \overline{\psi}_{x,1} \psi_{x,1} \overline{\psi}_{x,2} \psi_{x,2} \right\}$$

In 3D Scenario B is realized with 2nd order phase transition between PMW and PMS phases topological vs trivial insulator

#### **Dynamics?**

The system has topological "hedgehog" solutions that can condense similar to 2D XY model! S. Catterall, in prep

Preliminary results in 4D suggest the same structure

talk by V. Ayyar: arXiv:1410.6474



# More on Yukawa models

Talks by D. Chu, O. Akerlund

Yukawa models in the weak coupling region are "trivial"  $\rightarrow$  include a  $\Phi^6$  interaction and study this as an effective theory (mimics new dynamics)

– What is the effect of the  $\Phi^6$  on the IR dynamics ?

The model is not universal, but could show generic properties

The two groups use different (complementary ?) approaches

- Berlin-Taiwan : MC with overlap fermions
- Akerlund, DeForcrand, Steinbauer : effective MF (analytical)

# More on Yukawa models

Talks by D. Chu, O. Akerlund

Yukawa models in the weak coupling region are "trivial"  $\rightarrow$  include a  $\Phi^6$  interaction and study this as an effective theory (mimics new dynamics)

- What is the effect of the  $\Phi^6$  on the IR dynamics
  - vacuum stability problems are relaxed:
    - we do not sit on the edge
  - Finite T : the EW phase transition can become 1st order EW baryogenesis is possible
    - (but might need to consider gauge-scalar interactions)

# Gauge-Higgs unification on 5D orbifold

Idea: identify the Higgs field with the five-dimensional components of a gauge field



dimensional reduction via localization:

Alberti, Irges, Knechtli and Moir, arXiv:1506.06035; talk by G. Moir

 study of warped fifth dimension talk by E Lambrou Strong dynamics -Composite Higgs models

Challenges & progress

# Higgs era of particle physics

### 2012 : 125GeV Higgs boson



# Higgs era of particle physics

2012 : 125GeV Higgs boson



## 2015 : 3.4 σ excess in diboson



Do not get too excited....

# Higgs era of particle physics

### 2012 : 125GeV Higgs boson



What is the mechanism that keeps the Higgs light? Spontaneously broken symmetry -> massless Goldstone bosons

- Flavor symmetry: SSB leads to massless "pions"
- Scale symmetry: SSB leads to dilaton: near-conformal models

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What is the mechanism that generates SM fermion masses

- Bilinear coupling conformal TC :  $y_{\psi}\psi_L\psi_R\mathcal{O}_H$
- Linear coupling Partial composite models:  $\lambda_{\psi} \overline{\psi} \mathcal{O}_{\psi}$

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Golterman, Shamir :

calculate the couplings of the low energy effective theory in terms of baryonic correlation functions ; in preparation to lattice simulations

arXiv:1502.00390

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- Scale symmetry: SSB leads to dilaton: near-conformal models

Lattice simulations investigate strongly coupled but simple models Couple them to Standard Model  $\rightarrow$  things will change

# Challenge #1 : Is there a light dilaton?

Is the 0<sup>++</sup> scalar of a near-conformal system light\*?

\*How light is light enough?

# Light dilaton?

Is the 0<sup>++</sup> scalar of a near-conformal system light?

dilaton emerges from spontaneously conformal breaking dilaton mass ≈ value of the β function at breaking of conformal symmetry

(Csaki et al, arXiv:1305.3919)



# Light dilaton?

Is the 0<sup>++</sup> scalar of a near-conformal system light?

dilaton emerges from spontaneously conformal breaking dilaton mass ≈ value of the β function at breaking of conformal symmetry (Csaki et al, arXiv:1305.3919)



tuning is required for a light dilaton





![](_page_23_Figure_2.jpeg)

![](_page_24_Figure_2.jpeg)

![](_page_25_Figure_2.jpeg)

# Light dilaton?

![](_page_26_Figure_1.jpeg)

maybe we are lucky ....

# Challenge #2 : Chiral extrapolation

![](_page_27_Figure_1.jpeg)

Simulations are in a regime where the 0<sup>++</sup> state is degenerate/lighter than the pion

Chiral extrapolations need to take this into account:

some models, no solid approach yet

talk by J. Kuti

# Challenge #3: What is the (right?) model

Guiding principles

- chirally broken<sup>\*</sup>
- walking<sup>\*</sup> with large anomalous dimension
- light 0++

. . . .

\* or not??

# Challenge #3: What is the (right?) model

Guiding principles

- chirally broken<sup>\*</sup>
- walking<sup>\*</sup> with large anomalous dimension
- light 0++

Two approaches :

- Look at many models and find generic features, identify pitfalls, etc ( and then do large scale studies )
- Pick a likely model and do a large-scale study

![](_page_29_Figure_8.jpeg)

# Models & methods (this conference)

	Nf	Step Scaling	Finite Size S	Finite T	Eigen modes	Runnin coupli	Spectr. conn	Spectr. scalar	
	<b>1A</b>				Х		Х	X	
	2A	X			Х		Х	Х	
(2)	2F						Х		
SU	4F						Х		
	6F	X					Х		
	8F	X					Х		
	4F	X		Х	Х		Х	Х	
$\widehat{\mathbf{x}}$	8F	X	Х	Х	Х	Х	Х	Х	
SU(3	4+8				Х	Х	Х	Х	
	12F	X	Х	Х	Х		Х	Х	
	2S	X	Х		Х	Х	Х	Х	

# Large scale (+ exploratory) studies

	Nf	Step Scaling	Finite Size S	Finite T	Eigen modes	Runnin coupli	Spectr. conn	Spectr. scalar	
	<b>1A</b>				X		Х	Х	
	<b>2A</b>	X			Х		Х	Х	
(2)	2F						Х		
SU	4F						Х		
	6F	X					Х		
	8F	X					Х		
	4F	X		Х	X		Х	Х	
$\widehat{\mathbb{C}}$	8F	X	Х	Х	Х	Х	Х	Х	
SU(C	4+8				X	Х	Х	Х	
	12F	X	Х	Х	X		Х	Х	
	2S	X	Х		X	Х	Х	Х	

SU(3) 8F : LatKMI, LSD, LatHC;

talks by Ohki, Bennett, Rebbi, Weinberg, Nakayama

# Large scale (+ exploratory) studies

	N <sub>f</sub>	Step Scaling	Finite Size S	Finite T	Eigen modes	Runnin coupli	Spectr. conn	Spectr. scalar	
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	<b>2A</b>	X			X		Х	Х	
(2)	2F						Х		
SU	4F						Х		
	6F	Х					Х		
	8F	Х					Х		
	4F	Х		Х	Х		Х	Х	
3)	8F	Х	Х	Х	Х	Х	Х	Х	
SU(3	4+8				Х	Х	Х	Х	
S S	12F	Х	Х	Х	Х		Х	Х	
	<b>2S</b>	X	Х		Х	Х	Х	Х	

SU(3) 2S: LatHC ( + CP<sup>3</sup>, Boulder)

talks by Kuti, Wong, Santanu, Nogradi, Hansen posters by Holland, AH

# Exploratory studies

	Nf	Step Scaling	Finite Size S	Finite T	Eigen modes	Runnin coupli	Spectr. conn	Spectr. scalar	talks
	<b>1A</b>				Х		Х	Х	Bergner
	<b>2A</b>	Х			Х		Х	Х	Bergner
(2)	2F						Х		Tahtinen
SU	4F						Х		Tahtinen
	6F	Х					Х		Tahtinen
	8F	Х					Х		Leino
	4F	Х		Х	Х		Х	Х	Masafumi
3)	8F	Х	Х	Х	Х	Х	Х	X	
SU(3	4+8				Х	Х	Х	Х	Rebbi, Weinberg
U)	12F	Х	Х	Х	Х		Х	Х	D. Lin
	2S	Х	Х		Х	Х	Х	Х	

# In this talk

		Nf	Step Scaling	Finite Size S	Finite T	Eigen modes	Runnin coupli	Spectr. conn	Spectr. scalar	
		<b>1A</b>				X		Х	Х	
		2A	X			Х		Х	Х	
SU(2)	(2)	2F						Х		
	SU	4F						Х		
		6F	Х					Х		
		8F	Х					Х		
		4F	X		Х	Х		Х	Х	
	3)	8F	X	Х	Х	Х	Х	Х	Х	
	SU(S	<b>4+8</b>				Х	Х	Х	Х	
	رن ر	12F	X	Х	Х	Х		Х	Х	
		2S	X	Х		Х	Х	Х	Х	

An example: SU(3) N<sub>f</sub>=8 fundamental

Perturbative prediction:

2 loop : just below conformal window

3 & 4 loop MS-bar : strongly coupled IRFP

Questions for lattice studies:

- is the model chirally broken?
- is it walking?
- what is the anomalous dimension in the walking range?
- is there a light scalar?
- $-\pi$   $\pi$  scattering
- etc

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- etc

## Possibly the most frustrating BSM model

# SU(3) N<sub>f</sub>=8 fundamental - Step scaling

#### Two step scaling studies:

![](_page_37_Figure_2.jpeg)

- Step scaling function is monotonic up to  $g^2 = 15$ ,
- Consistent with 4-loop MS-bar
- Range limited by bulk transition

# SU(3) N<sub>f</sub>=8 fundamental - finite T

No spontaneous chiral symmetry breaking in the chiral limit yet

![](_page_38_Figure_2.jpeg)

D. Schaich et al (LSD), arXiv:1506.08791

## SU(3) N<sub>f</sub>=8 fundamental - finite size scaling

Finite size scaling: (in the chiral limit)

arXiv:1503.02359, talk by Nakayama

The propagator G(t,N) scales as (up to 1/N corrections)

$$G(t;g,m_q,N,\mu) = \left(\frac{N'}{N}\right)^{3-2\gamma} G(t';g',m'_q,N',\mu)$$

The quantity

$$\mathfrak{m}(t,N) = N \ln \frac{G(t,N)}{G(t+1,N)}. \qquad \tau = t/N_t$$

is scale invariant at the FP. Leading corrections from g<sup>2</sup>:

$$\mathfrak{m}(\tau, g, N) = \mathfrak{m}(\tau, g, N') + \mathcal{B}(g) \ln\left(\frac{N}{N'}\right) \partial_g \mathfrak{m}(\tau, g, N') + \cdots$$
$$= \mathfrak{m}(\tau, g, N')$$

 $\rightarrow$  tune  $\beta$  to find "curve collapse": IRFP

# SU(3) N<sub>f</sub>=8 fundamental - finite size scaling

![](_page_40_Figure_1.jpeg)

arXiv:1503.02359, talk by Nakayama

Numerical tests with Wilson fermions suggest that N<sub>f</sub>=8 is conformal De Silva et al, based on different arguments and staggered fermions reach the same conclusion arXiv:1506.06396

Caveats:

- do other operators predict the same  $\beta_{IRFP}$ ?
- does a volume squeezed chirally broken system look different?

An example: SU(3) N<sub>f</sub>=8 fundamental

Questions for lattice studies:

- is the model chirally broken?
- is it walking?
- what is the anomalous dimension ?
- is there a light scalar?

# SU(3) N<sub>f</sub>=8 fundamental - mass anomalous dimension

Anomalous dimension from Dirac operator mode number

![](_page_42_Figure_2.jpeg)

Anomalous dimension  $\gamma \approx 1$  makes this model exciting

## SU(3) N<sub>f</sub>=8 fundamental - Spectrum

Both LatKMI and LSD collaborations have extensive spectrum data

![](_page_43_Figure_2.jpeg)

## SU(3) N<sub>f</sub>=8 fundamental - 0 <sup>++</sup> Spectrum

0<sup>++</sup> is degenerate with the pion, at lightest pion  $m_{0^{++}} / m_{\rho} \lesssim 0.5$ 

![](_page_44_Figure_2.jpeg)

## SU(3) N<sub>f</sub>=8 fundamental - 0 <sup>++</sup> Spectrum

0<sup>++</sup> is degenerate with the pion, at lightest pion  $m_{0^{++}} / m_{\rho} \lesssim 0.5$ 

![](_page_45_Figure_2.jpeg)

So it is chirally broken with a light scalar, right?

# Challenge # 4: Conformal or chirally broken ?

	Step	Finite	Finite	Eigen	Runnin	Spectr.	Spectr.	Ref.
8F	?	?	?	?	?	?	?	

No definite signal of chiral symmetry breaking (or conformality)

Does it matter? YES

![](_page_46_Figure_4.jpeg)

# Challenge # 4: Conformal or chirally broken ?

		Step	Finite	Finite	Eigen	Runnin	Spectr.	Spectr.	Ref.
8	8F	?	?	?	?	?	?	?	

No definite signal of chiral symmetry breaking or conformality

Does conformality kill the model? NO

# Challenge # 4: Conformal or chirally broken ?

	Step	Finite	Finite	Eigen	Runnin	Spectr.	Spectr.	Ref.
8F	?	?	?	?	?	?	?	

No definite signal of chiral symmetry breaking or conformality

Does conformality kill the model? NO

We can move a conformal system into the chirally broken regime (add a 4-fermion operator)

Even more exciting!

# SU(3) with 4+8 flavors

A system with 4 massless (light) and 8 heavier flavors

- IR : 4 flavor, chirally broken
- UV : 12 flavor conformal

Tune the mass of the 8 "heavy" flavors to interpolate between conformal and chirally broken dynamics

Boston-Boulder coll. talks by Rebbi, Weinberg

## 4+8 flavors : "walking" running coupling

![](_page_50_Figure_1.jpeg)

 $g_{GF}^2(\mu)$  develops a "shoulder" as  $m_h \rightarrow 0$ : this is walking ! Walking range can be tuned arbitrarily with  $m_h$ 

## 4+8 flavors : "walking" running coupling

![](_page_51_Figure_1.jpeg)

 $g_{GF}^2(\mu)$  develops a "shoulder" as  $m_h \rightarrow 0$  : this is walking ! Walking range can be tuned arbitrarily with  $m_h$ 

## 4+8 flavors : "walking" running coupling

![](_page_52_Figure_1.jpeg)

 $g_{GF}^2(\mu)$  develops a "shoulder" as  $m_h \rightarrow 0$ : this is walking ! Walking range can be tuned arbitrarily with  $m_h$ 

# 4+8 flavors : spectrum (at m<sub>h</sub>=0.080)

# $M_{\rho}/M_{\pi}$ shows the approach to the chiral regime

# $0^{++}$ is just above $M_{\pi}$ very different from $N_f=4!$

![](_page_53_Figure_3.jpeg)

If tuning  $m_h \rightarrow 0$  does not give a light enough  $0^{++}$ , we can try 2+8 or a cascade of masses 2+ (6+1+1+1+1....)

# 4+8 flavors : spectrum (at m<sub>h</sub>=0.080)

# $M_{\rho}/M_{\pi}$ shows the approach to the chiral regime

# $0^{++}$ is just above $M_{\pi}$ very different from $N_f=4!$

![](_page_54_Figure_3.jpeg)

2-loop perturbation theory predicts an IRFP at  $g^2 \sim 10$ 3, 4 loop MS-bar  $g^2 \sim 6$ 

High statistics, extensive studies with staggered fermions favor chiral symmetry breaking with light scalar (LatHC, talks by Kuti, Wong, Santano, Nogradi, Holland)

Emerging Wilson fermion study is consistent both with conformal hyper scaling and chiral symm breaking

(CP<sup>3</sup>, Hansen)

Step scaling studies with Wilson fermions are in tension with staggered

(Boulder coll, AH et al)

![](_page_56_Figure_1.jpeg)

- the 3 lightest mass points might not be reliable
- the spectrum is not a good way to differentiate conformal vs chirally broken systems

#### Preliminary spectrum from CP<sup>3</sup>

#### talk by Hansen

![](_page_57_Figure_3.jpeg)

#### Preliminary spectrum from CP<sup>3</sup>

#### talk by Hansen

![](_page_58_Figure_3.jpeg)

The two collaborations probe quite different parameter spaces

Challenge #(n+1): Universality

Even in QCD we want to verify that Wilson / staggered/overlap fermions predict the same IR physics

In a walking theory that is much harder. In a conformal theory it might not even be true (Wilson Renormalization Group - universality requires identical symmetries!) :

![](_page_59_Figure_3.jpeg)

# Step scaling of N<sub>f</sub>=2 sextet

Step scaling

with staggered

Wilson fermions

![](_page_60_Figure_4.jpeg)

VS

# Step scaling of N<sub>f</sub>=2 sextet

Step scaling

with staggered

Wilson fermions

![](_page_61_Figure_4.jpeg)

VS

Is the tension due to universality? The question has to be settled

## SU(3) N<sub>f</sub>=2 sextet, Wilson fermions

#### Consistency checks:

![](_page_62_Figure_2.jpeg)

![](_page_62_Figure_3.jpeg)

## Conclusion

BSM lattice studies:

### Lattice calculations complement phenomenology : we could have real impact !

- These calculations are difficult, even when not high precision
- Plenty of challenges, also solutions

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Future :

I put my bet on systems that interpolate between conformality and chiral symmetry breaking with a tunable parameter

## Conclusion

BSM lattice studies:

Lattice calculations complement phenomenology : we could have real impact !

- These calculations are difficult, even when not high precision
- Plenty of challenges, also solutions

Future :

I put my bet on systems that interpolate between conformality and chiral symmetry breaking with a tunable parameter

Thank you

## Challenge #5: Control of walking

Can we control the amount of walking? need  $(F^2L^2) < N_f$  — usually not (or barely) satisfied

![](_page_66_Figure_2.jpeg)

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Can we control the amount of walking? need  $(F^2L^2) < N_f$  — usually not (or barely) satisfied

![](_page_67_Figure_2.jpeg)