

# Extracting the $\eta'$ meson mass from gluonic correlators in lattice QCD



 OSAKA UNIVERSITY  
Live Locally, Grow Globally

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[ JLQCD collaboration ]

# What I talked at Lattice 2014

## New method for topological susceptibility

### 1. Gluonic topology density after Wilson flow

$$q(x) = c \text{Tr} F_{\mu\nu} \tilde{F}^{\mu\nu}(x)$$

-> smooth gauge field, good noise reduction.

### 2. Subvolume topology $\chi_t^{\text{sub}} \equiv \int_{|x| < r_{\text{cut}}} d^4x \langle q(x)q(0) \rangle$

-> shorter auto-correlation than global Q.

### 3. Subtract bias from global topology

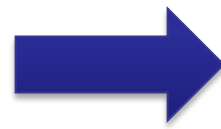
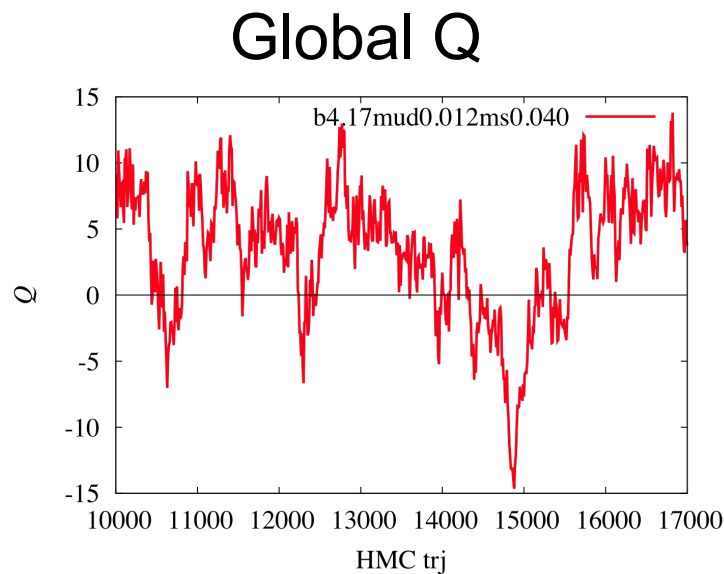
$$\langle q(x)q(0) \rangle_Q \xrightarrow{|x| \rightarrow \text{large}} \frac{1}{V} \left[ \frac{Q^2}{V} - \chi_t \right] + O(V^{-2}),$$

[Aoki, F, Hashimoto, Onogi 2007]

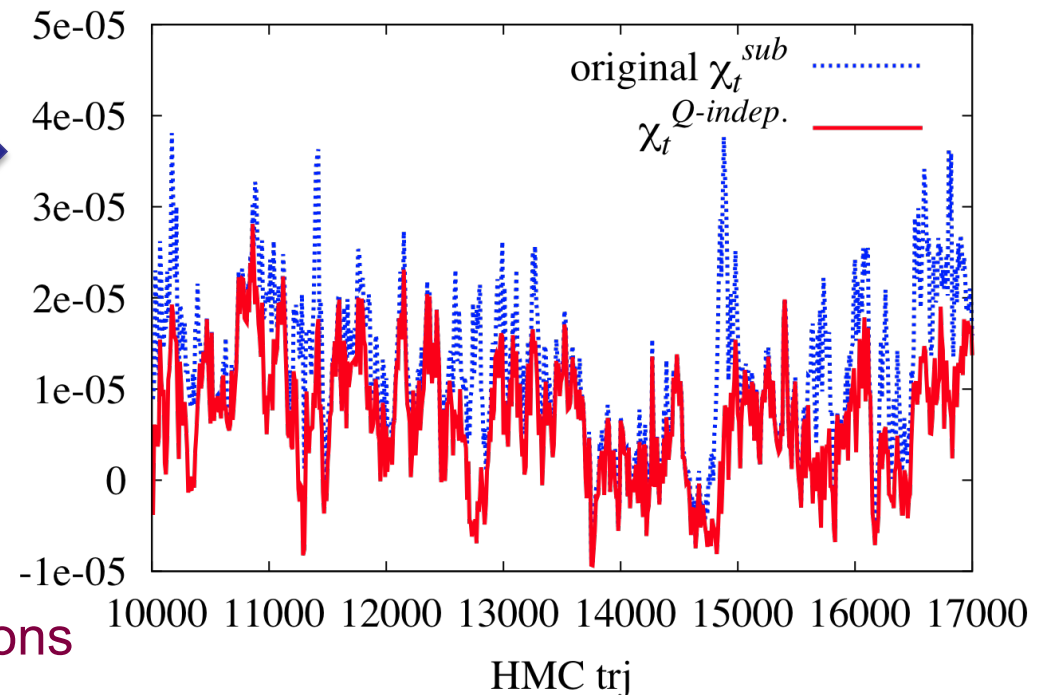
New  $\chi_t$  has shorter auto-correlation.

$$\langle \chi_t^{Q\text{-indep.}} \rangle = \frac{V}{V - V_{sub}} \left( \chi_t^{sub} - \frac{V_{sub}}{V^2} Q^2 \right)$$

L=32 b4.17mud0.012ms0.040



$\chi_t^{sub}$

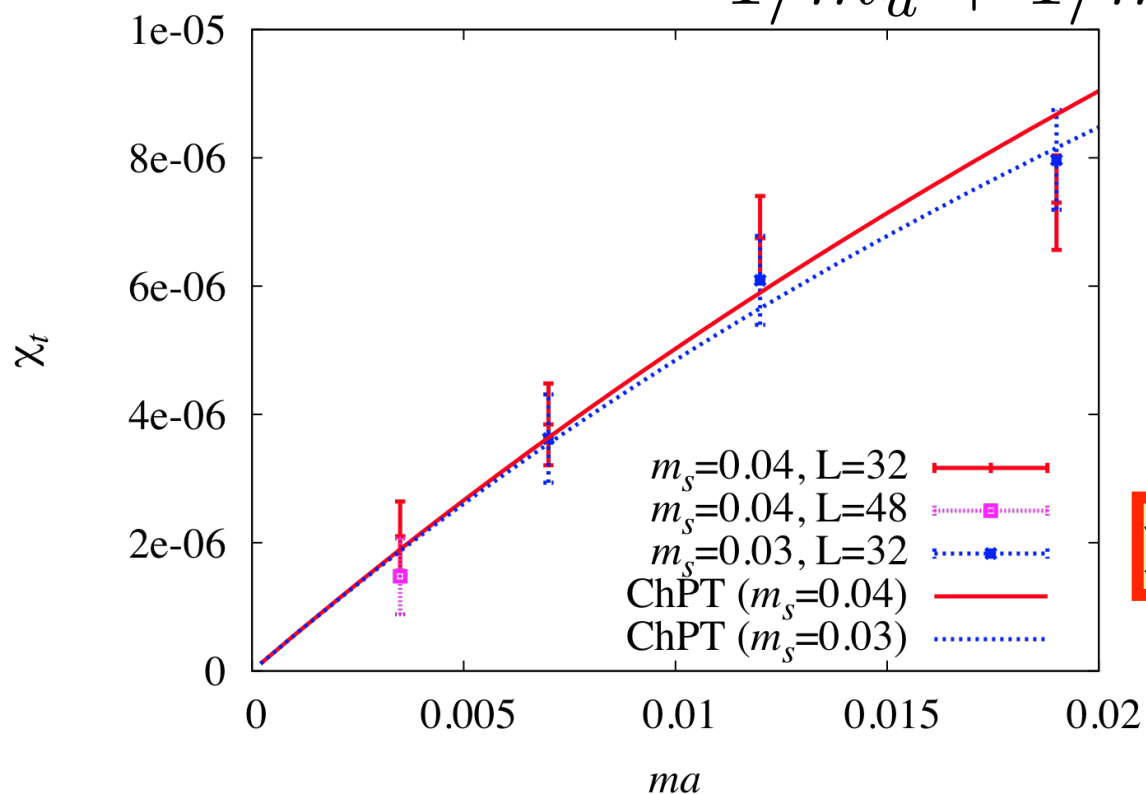


JLQCD 2014 with Mobius DW fermions

# Topological susceptibility

$\chi_t$   $Q$ -indep. agrees with ChPT prediction :

$$\chi_t^{\text{ChPT}} = \frac{\Sigma}{1/m_u + 1/m_d + 1/m_s}$$



JLQCD 2014  
with Mobius DW fermions

1-param fit

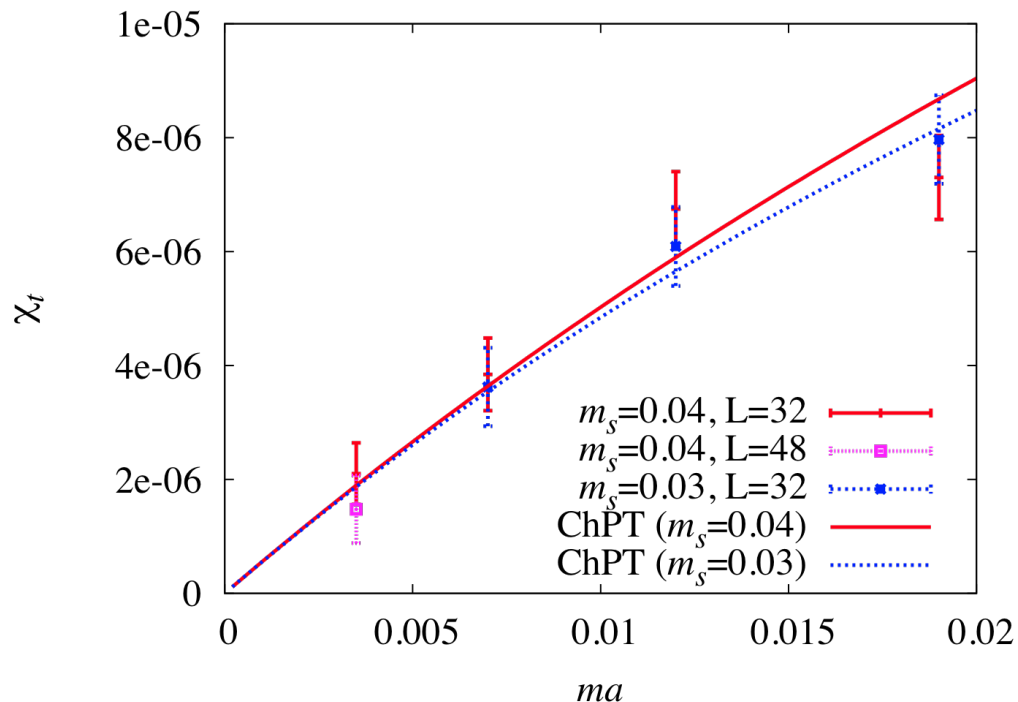
→

$$\Sigma^{\overline{\text{MS}}} 1/3 (2\text{GeV}) \sim 255\text{MeV.}$$

# Topological susceptibility

$\chi_t$   $Q$ -indep. agrees with ChPT prediction.

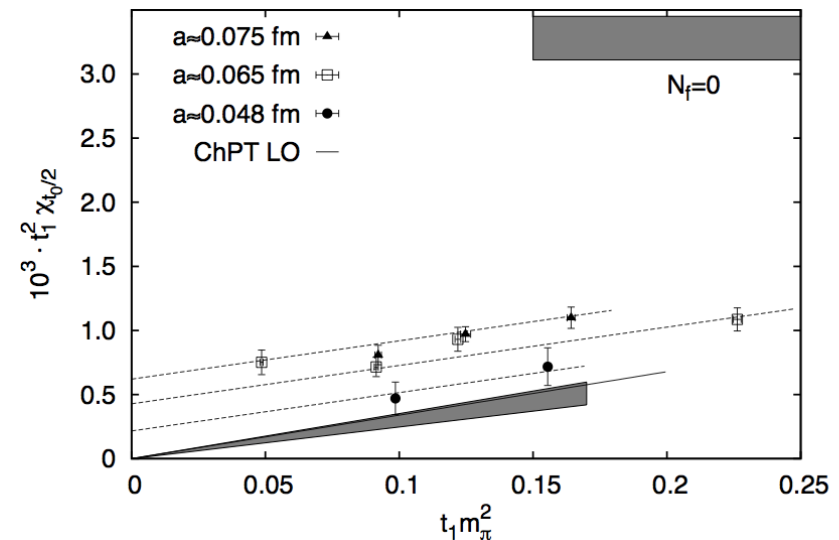
Our work w/ Mobius DW fermions



Cf. w/ Wilson fermions

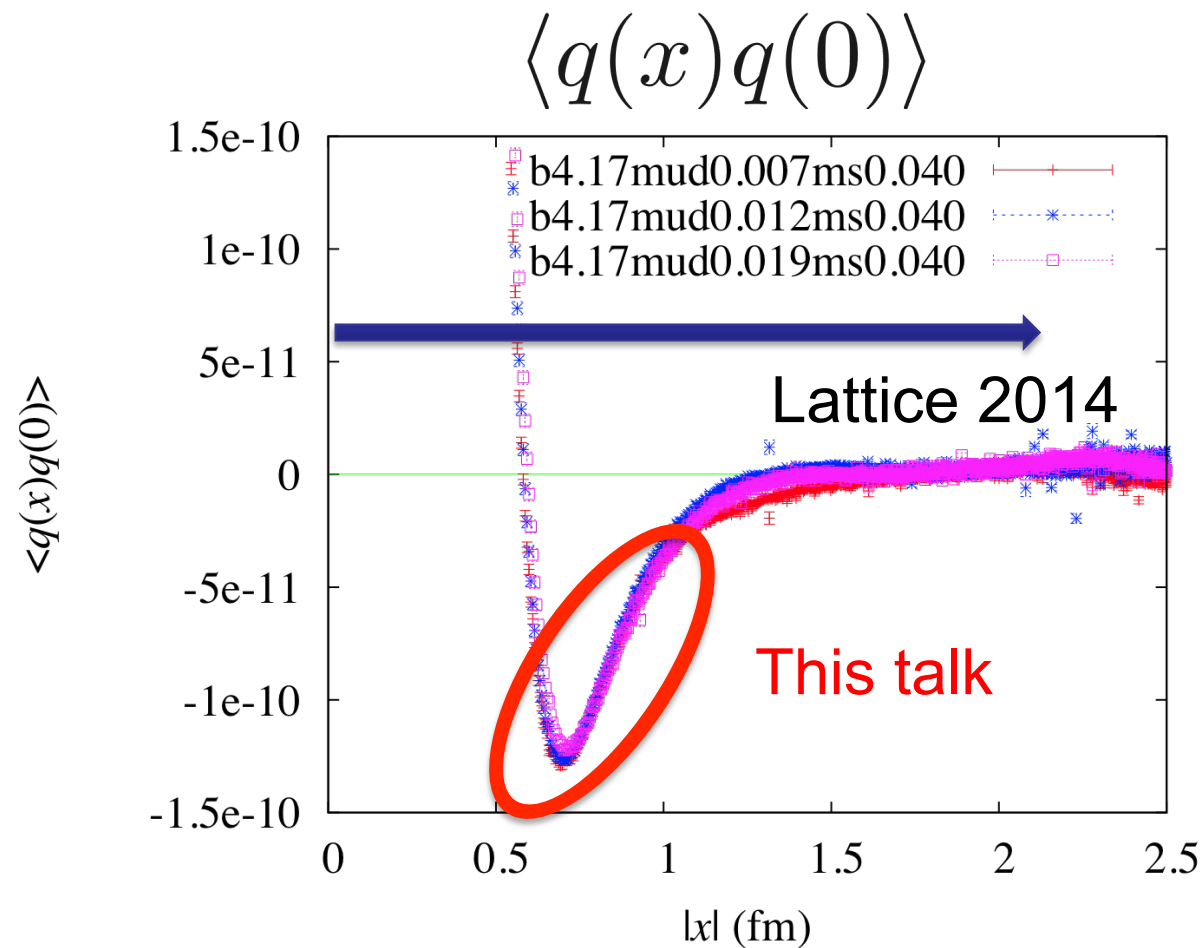
[Bruno, Schaefer, Sommer 14]

$N_f = 2$



# Today's target

= Same correlator but shorter region.





# 1. Introduction

# Eta-prime meson is interesting.

$$\eta' = \bar{\Psi} \gamma_5 \Psi \quad (\Psi = (u, d, s)^T)$$

would-be a Nambu-Goldstone boson  
for  $U(1)_A$  symmetry but heavy  
(due to anomaly) =  **$U(1)_A$  problem**

[related talks by Tomiya, Cossu, Sharma Sat. morning]




# Eta-prime meson is difficult.

$$\eta' = \bar{\Psi} \gamma_5 \Psi \quad (\Psi = (u, d, s)^T)$$

difficult quantity to treat on a lattice  
because of “disconnected” part  
in the correlators :

$$\langle \eta'(x) \eta'(y) \rangle \simeq \underbrace{\langle \bar{\Psi} \gamma_5 \Psi(x) \bar{\Psi} \gamma_5 \Psi(y) \rangle}_{\text{connected}} + \underbrace{\langle \bar{\Psi} \gamma_5 \Psi(x) \bar{\Psi} \gamma_5 \Psi(y) \rangle}_{\text{disconnected}}$$

# Disconnected part is expensive.



$$\langle \bar{\Psi} \gamma_5 \Psi(x) \bar{\Psi} \gamma_5 \Psi(y) \rangle = \langle \text{Tr} D^{-1}(x, y) D^{-1}(x, y)^* \rangle$$

➡ Enough to solve  $D(x, z)s(z) = v(x), \quad s(z) = \delta_{y,z}$



$$\langle \bar{\Psi} \gamma_5 \Psi(x) \bar{\Psi} \gamma_5 \Psi(y) \rangle = \langle \text{Tr} \gamma_5 D^{-1}(x, x) \text{Tr} \gamma_5 D^{-1}(y, y) \rangle$$

➡ Requires  $D(x, z)s(z) = v(x), \quad s(z) = \delta_{x,z}$

many times for different x and y.

# Disconnected part is noisy

due to pion's contamination.

$$\begin{aligned}
 \langle \eta'(x)\eta'(y) \rangle &\simeq \langle \bar{\Psi} \gamma_5 \Psi(x) \bar{\Psi} \gamma_5 \Psi(y) \rangle + \langle \bar{\Psi} \gamma_5 \Psi(x) \bar{\Psi} \gamma_5 \Psi(y) \rangle \\
 &= A e^{-m_\pi |x-y|} - A e^{-m_\pi |x-y|} + B e^{-m_{\eta'} |x-y|}
 \end{aligned}$$



noisy

noisy

our target

# In this work

We use

[Similar (quenched) works : Chowdhury et al. 2014,  
Ohki this conference 14<sup>th</sup> 14:00 BSM]

1. gluonic (topology density) operator :

$$\eta' = \frac{\bar{u}\gamma_5 u + \bar{d}\gamma_5 d + \bar{s}\gamma_5 s}{\sqrt{3}} \rightarrow \frac{1}{32\pi^2} \text{Tr} \epsilon_{\mu\nu\rho\sigma} F^{\mu\nu} F^{\rho\sigma}$$



Free from pion's fluctuation,  
Numerically less expensive.

2. smeared links via Wilson flow



Noise reduction.

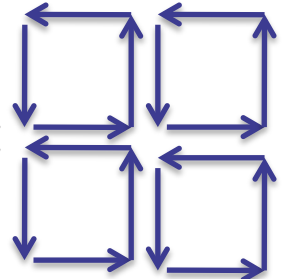
# Contents

- ✓ 1. Introduction
- 2. Topological charge density and (short) Wilson flow
- 3. Lattice simulations
- 4. Summary



## 2. Topological charge density and (short) Wilson flow

# Topological charge density

$$F_{\mu\nu}^{lat} = \text{Diagram} \Rightarrow q(x) = \frac{1}{32\pi^2} \text{Tr} \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu}^{lat} F_{\rho\sigma}^{lat}$$


1.  $\langle q(x)q(y) \rangle$  is **numerically cheap** with FFT (fast Fourier transformation).
2. **No** (direct) **pion's contamination**.
3. Global topological charge

$$\sum_x q(x) = Q + O(a^2)$$

# Wilson flow smearing

## Wilson flow equation

$$\partial_t A_\mu(t, x) = -\frac{\partial S_{YM}}{\partial A_\mu}, \quad A_\mu(0, x) = A_\mu(x)$$

smoothed length  $d \sim \sqrt{8t}$

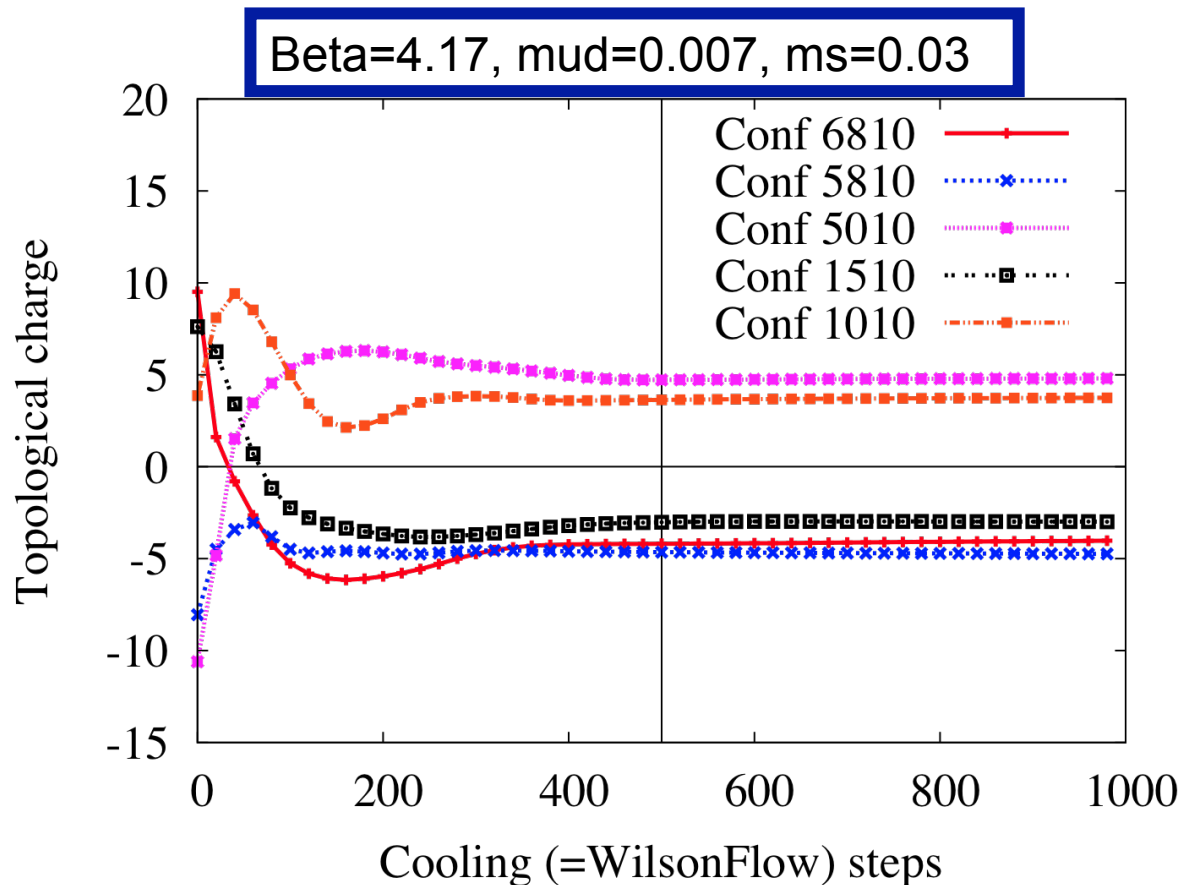
Gluonic correlators are UV finite.

[Luscher & Weisz 2011]



# Topological charge

Flow time history of  $\sum_x q(x) = Q + O(a^2)$

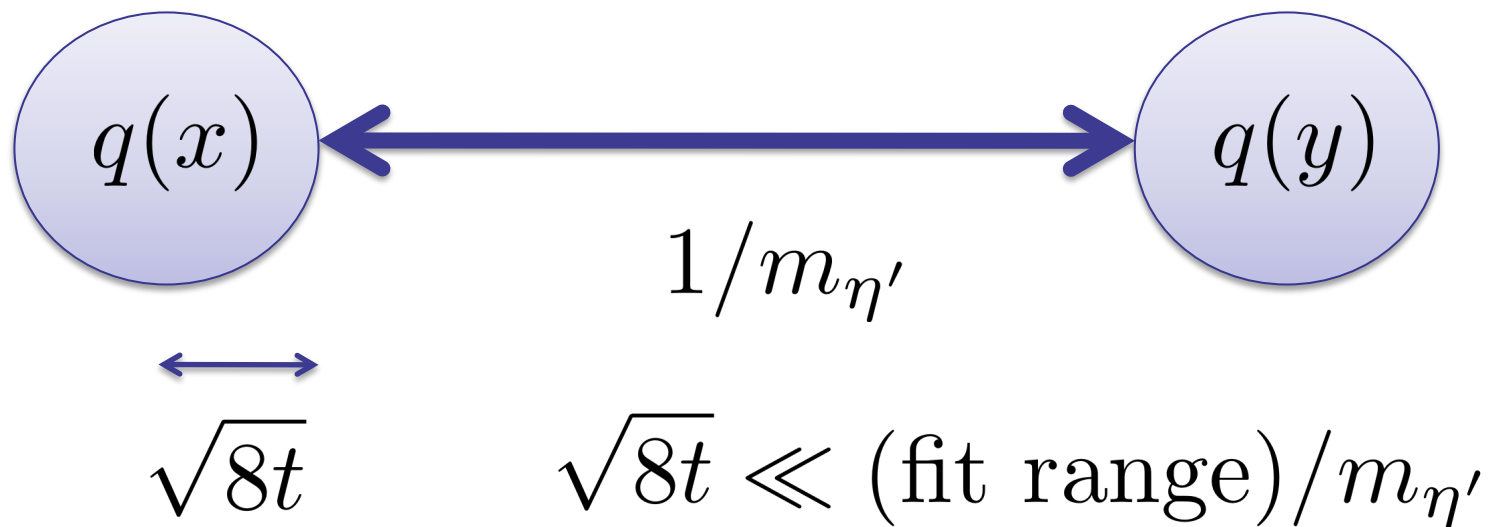


After Wilson  
flow around  
 $\sqrt{8t} \sim 0.5\text{fm}$

topological  
charge does  
not change.

# Wilson time flow has to be short...

Available window for flow time



Our choice :  $\sqrt{8t} < 0.25\text{fm}, \quad (\text{fit range}) > 0.7\text{fm}.$



# 3. Lattice simulations

# JLQCD's new project

Simulations on bigger & finer lattices started.

Computers @KEK: SR11000 ( 2 TFLOPS) + BG/L ( 57 TFLOPS)  
→ SR16000 (55 TFLOPS) + BG/Q (1.2 PFLOPS)

Lattice cut-off : 1.8 GeV → 2.5, 3.6, 4.5 GeV

Lattice size :  $16^3 \times 48$  →  $32^3 \times 64$ ,  $48^3 \times 96$ ,  $64^3 \times 128$

(Physical size : 1.8 fm → 2.6 fm ~ 4 fm )

Fermion action : overlap fermion → (improved) DomainWall

Pion mass : 200-400 MeV



Hitachi SR16000



IBM Blue Gene/Q

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Pion mass : 200-400 MeV

Our goal = 1% precision of (B)SM calculations  
(in particular, D & B mesons )



Hitachi SR16000



IBM Blue Gene/Q

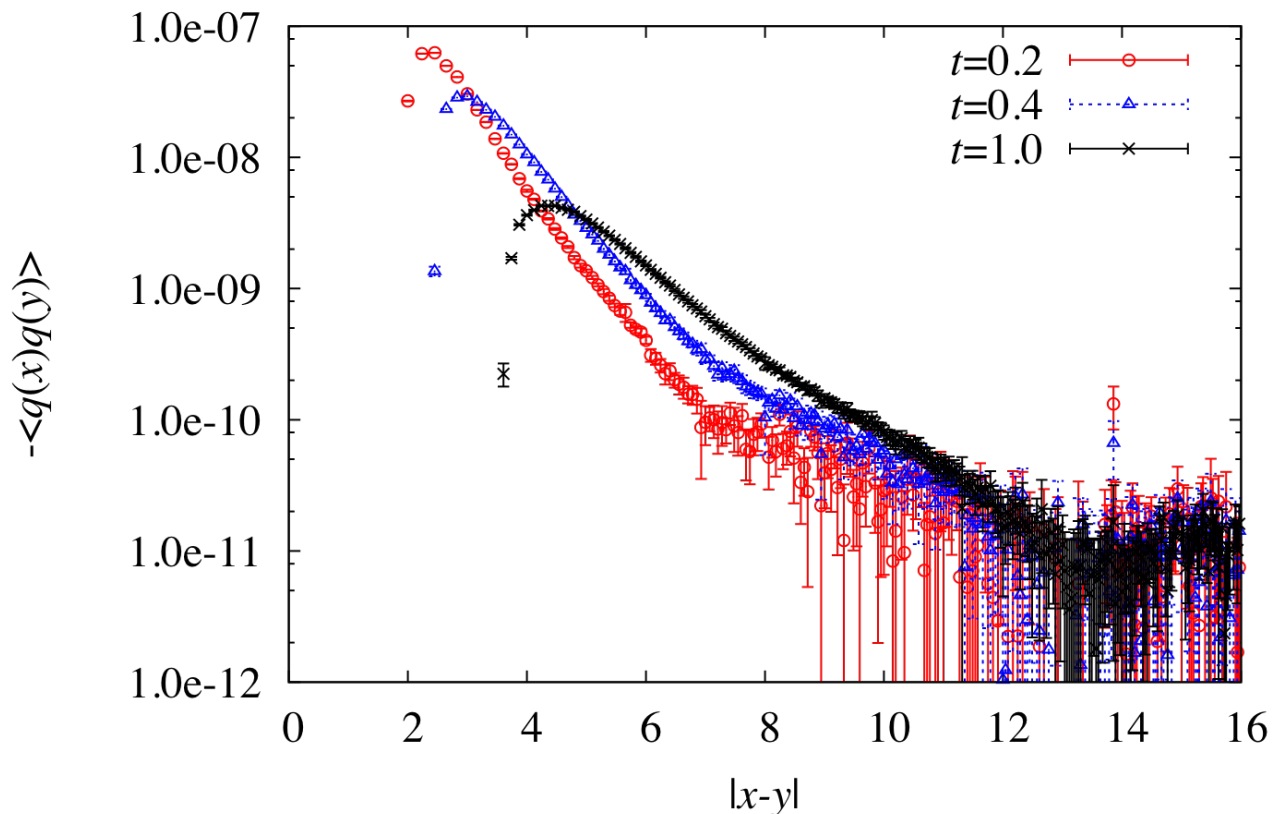
# Talks by JLQCD members

- B. Fahy, meson spectrum and decay const. 14<sup>th</sup>,  
16:30
- K. Nakayama, Charmonium, 15<sup>th</sup>, 16:50
- T. Suzuki, Dmeson decays, 16<sup>th</sup>, 11:20
- N. Yamanaka, nucleon axial and tensor charges, 15<sup>th</sup>,  
17:10
- G. Cossu, A. Tomiya, axial U(1) at finite T, 18<sup>th</sup> 9:00,  
9:20
- T. Kaneko light meson form factors, 14<sup>th</sup>, 16:30
- S. Hashimoto, Dirac spectrum, 15<sup>th</sup>, poster
- M. Tomii, renormalization & OPE, 15<sup>th</sup>, 16:30

# Topology density correlator

$$-\langle q(x)q(y) \rangle \sim AK_1(m_{\eta'}|x-y|)/|x-y|$$

b4.17 mud0.007 ms0.040



As the Wilson flow time  $t$  increases,

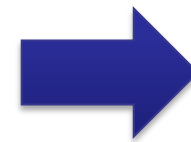
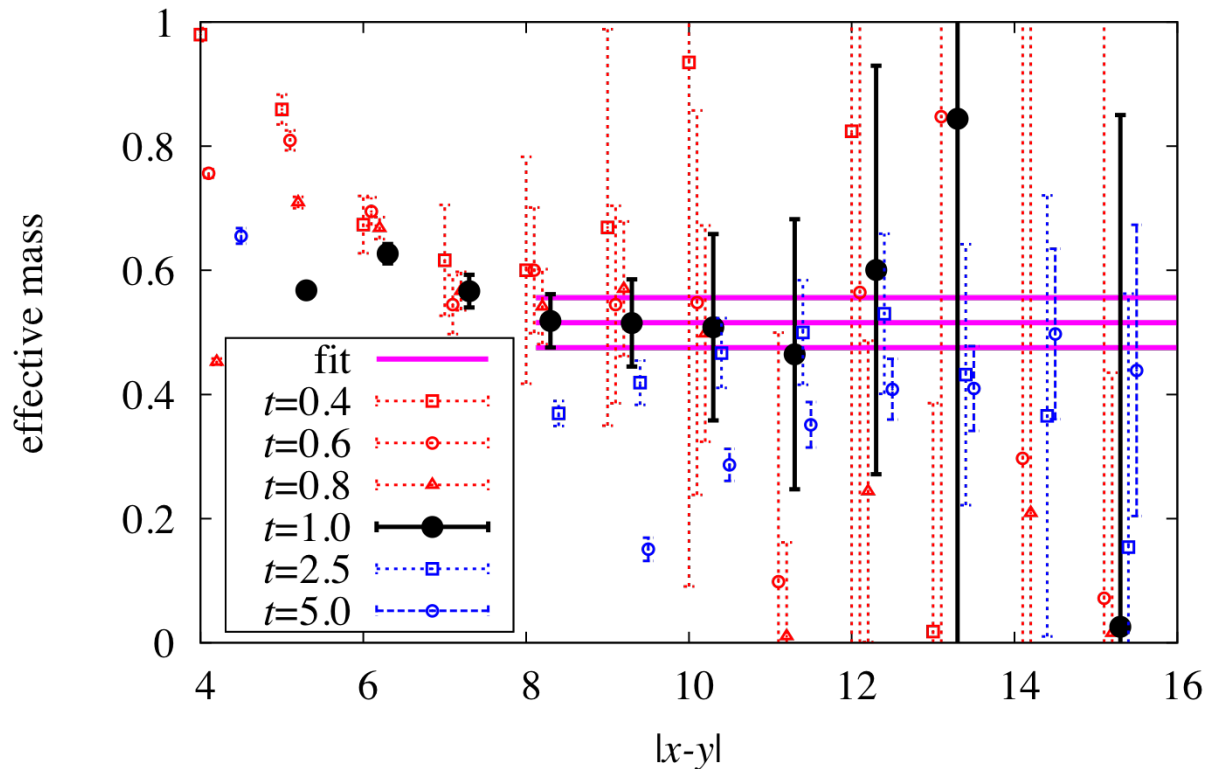
1. Better signals
2. Short ( $\sim \sqrt{8t}$ ) correlation destroyed.

[similar observations w/ different coolings : de Forcrand et al. 1997, Horvath et al. 2003, 2005, Koma et al. 2005... ]

# Effective mass plot

$$m_{\text{eff}}(|x - y|) \sim -\frac{1}{\Delta x} \ln \frac{\langle q(x + \Delta x)q(y) \rangle}{\langle q(x)q(y) \rangle}$$

b4.17 M1.00 mud0.0035 ms0.040

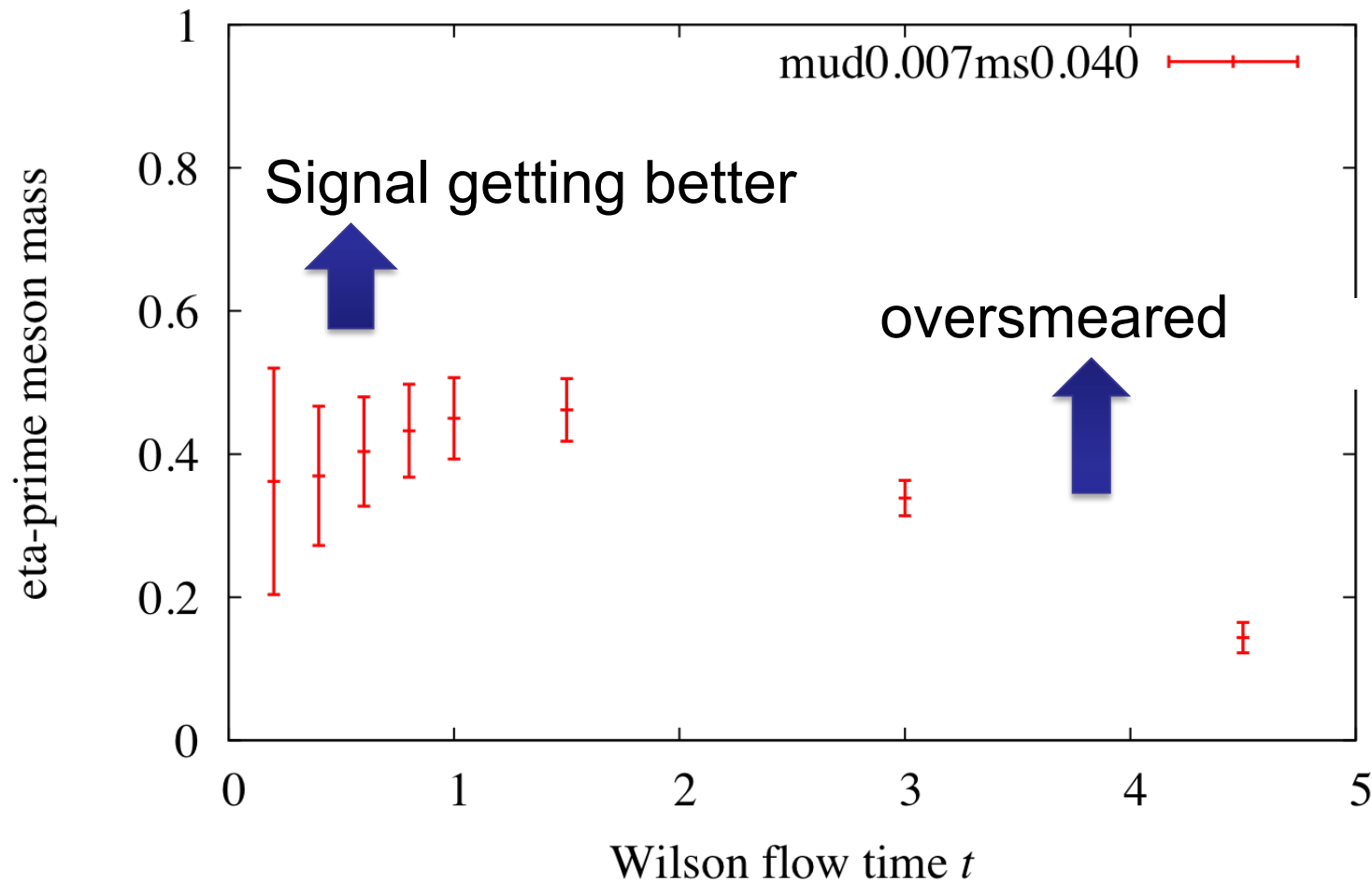


Fit range  
 $|x-y| > 8$  (  $\sim 0.7$  fm )

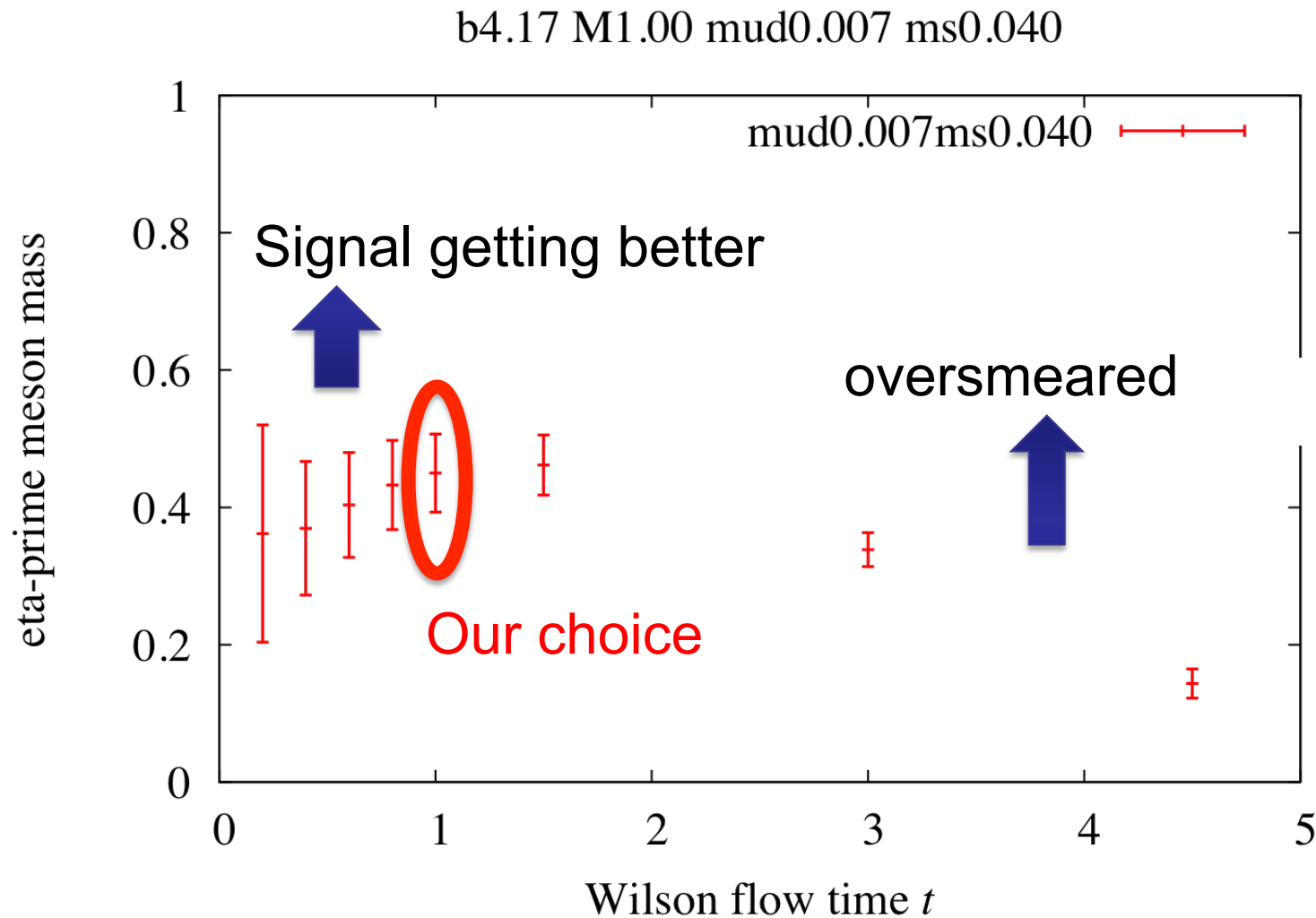


# Flow time dependence of fit

b4.17 M1.00 mud0.007 ms0.040

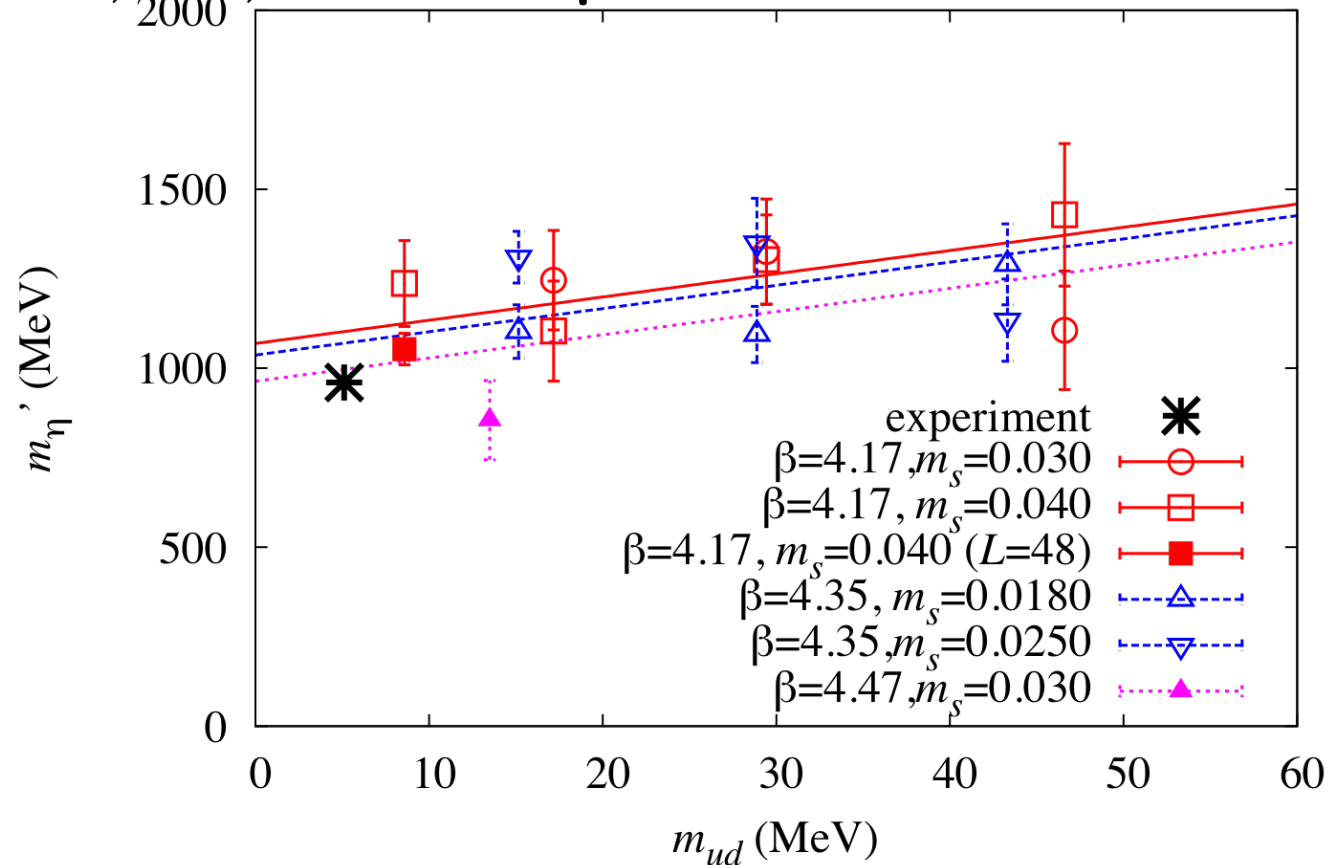


# Flow time dependence of fit



# Results

$m_{ud}, a, V$  dependences look mild.



linear fit in  $m_{ud}, a^2 \rightarrow m_{\eta'} = 1.00(11) \text{ GeV}$ .<sup>27</sup>

# Mixing with eta meson

SU(3) breaking  $\rightarrow$  mixing with eta:

$$\eta_0 = \frac{\bar{u}\gamma_5 u + \bar{d}\gamma_5 d + \bar{s}\gamma_5 s}{\sqrt{3}} = \eta' \cos \theta + \eta \sin \theta$$

So far we have neglected  $\theta$  .

**Cf. previous works**

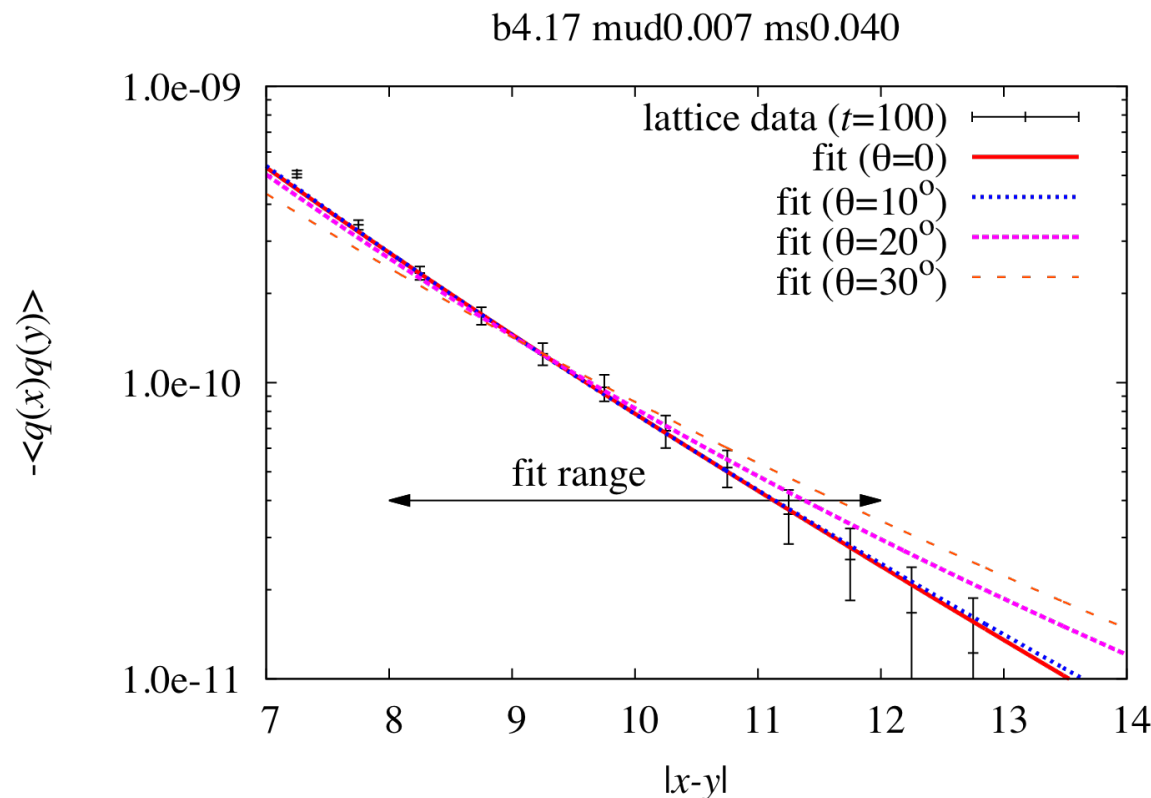
[RBC/UKQCD 2010, UKQCD 2012...]

**fit combined with**  $\eta_8 = \eta' \sin \theta + \eta \cos \theta$

$\rightarrow m_\eta$  and  $m_{\eta'}$

If  $|\theta| < 20^\circ$ , eta meson's contribution is negligible.

Modifying the fit curve w/ measured  $m_{\eta'}$  and  $m_{\eta}$  (using GMOR relation)



In our fit range,  
 contribution from eta  
 is small

if  $|\theta| < 20^\circ$

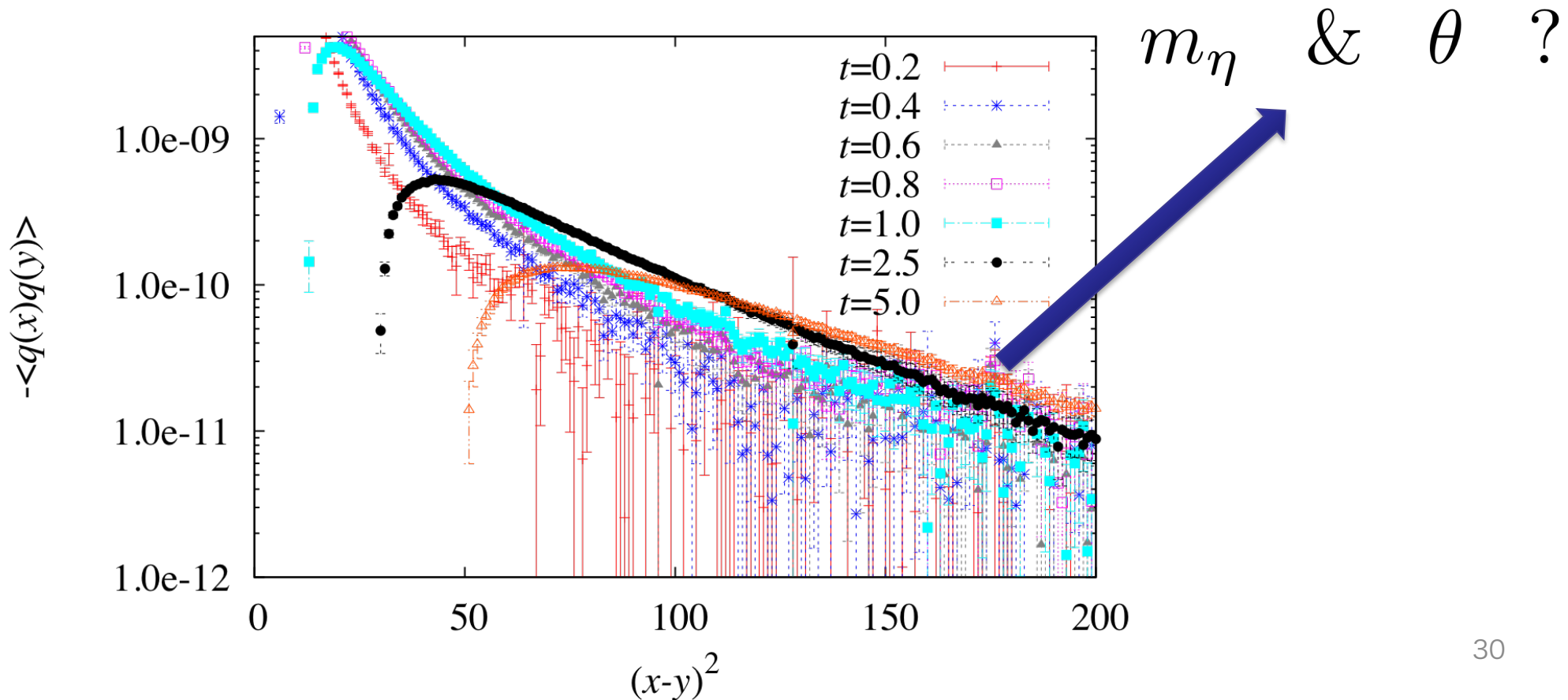
Cf. phenomenological  
 value (at phys.point)

$$-\theta = 10^\circ - 25^\circ$$

# Still interesting to look for eta

## Correlator at longer Wilson flow

b4.17 M1.00 mud0.0035 ms0.040





# 4. Summary

# Summary

Gluonic operator (topology density)

+

Short Wilson flow

numerical cost is negligible,

free from pion's contamination,

good noise reduction.

→ eta-prime meson mass  
( + eta meson mass & mixing angle )



# Numerical cost

$\langle q(x)q(y) \rangle + \text{Wilson flow}$   
= at most **1/10000**

of disconnected fermion integrals,  
( or **essentially zero** when you  
perform a Wilson flow for  
other measurements. )

# Backup slide 1

Correlator w/o momentum insertion

$$\langle q(x)q(0) \rangle = \frac{A}{|x|} K_1(m_{\eta'} |x|)$$

\* The data at same  $|x|$  are averaged.

# Backup slide 3

Non-trivial t dependence

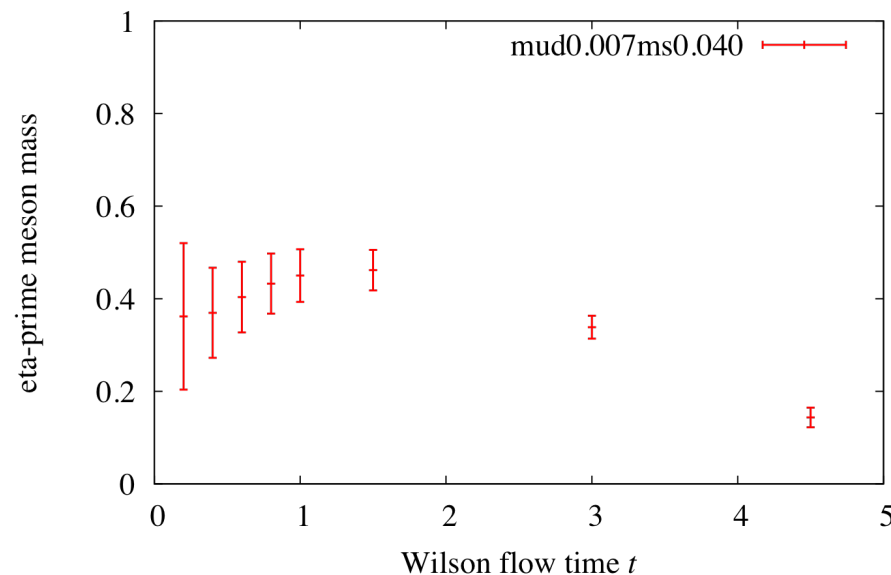
Gaussian model [Alpha collaboration 2014]

-> correction  $\sim e^{-\left(\frac{|x-y|}{\sqrt{8t}-m\sqrt{8t}}\right)^2} \frac{m(8t)^{3/2}}{|x-y|^2}$

b4.17 M1.00 mud0.007 ms0.040

< 1% in our case.

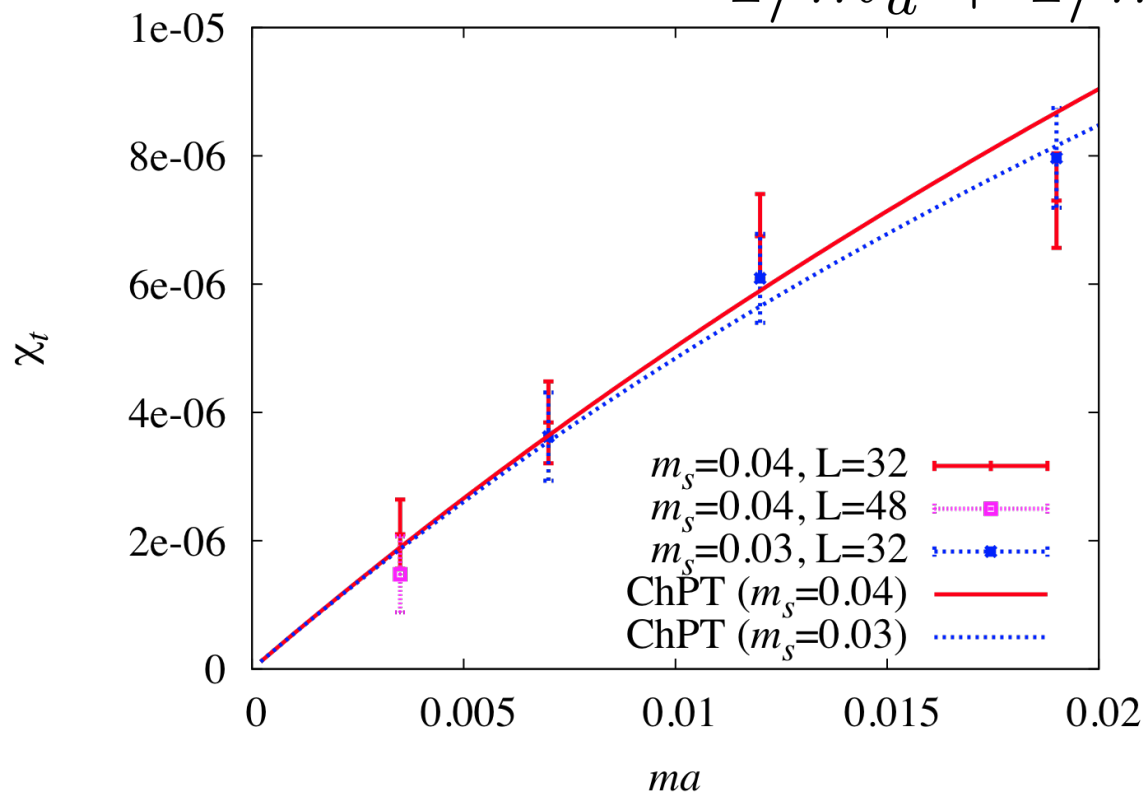
In fact, t dependence looks mild.



# Topological susceptibility

$\chi_t$   $Q$ -indep. agrees with ChPT prediction :

$$\chi_t^{\text{ChPT}} = \frac{\Sigma}{1/m_u + 1/m_d + 1/m_s}$$



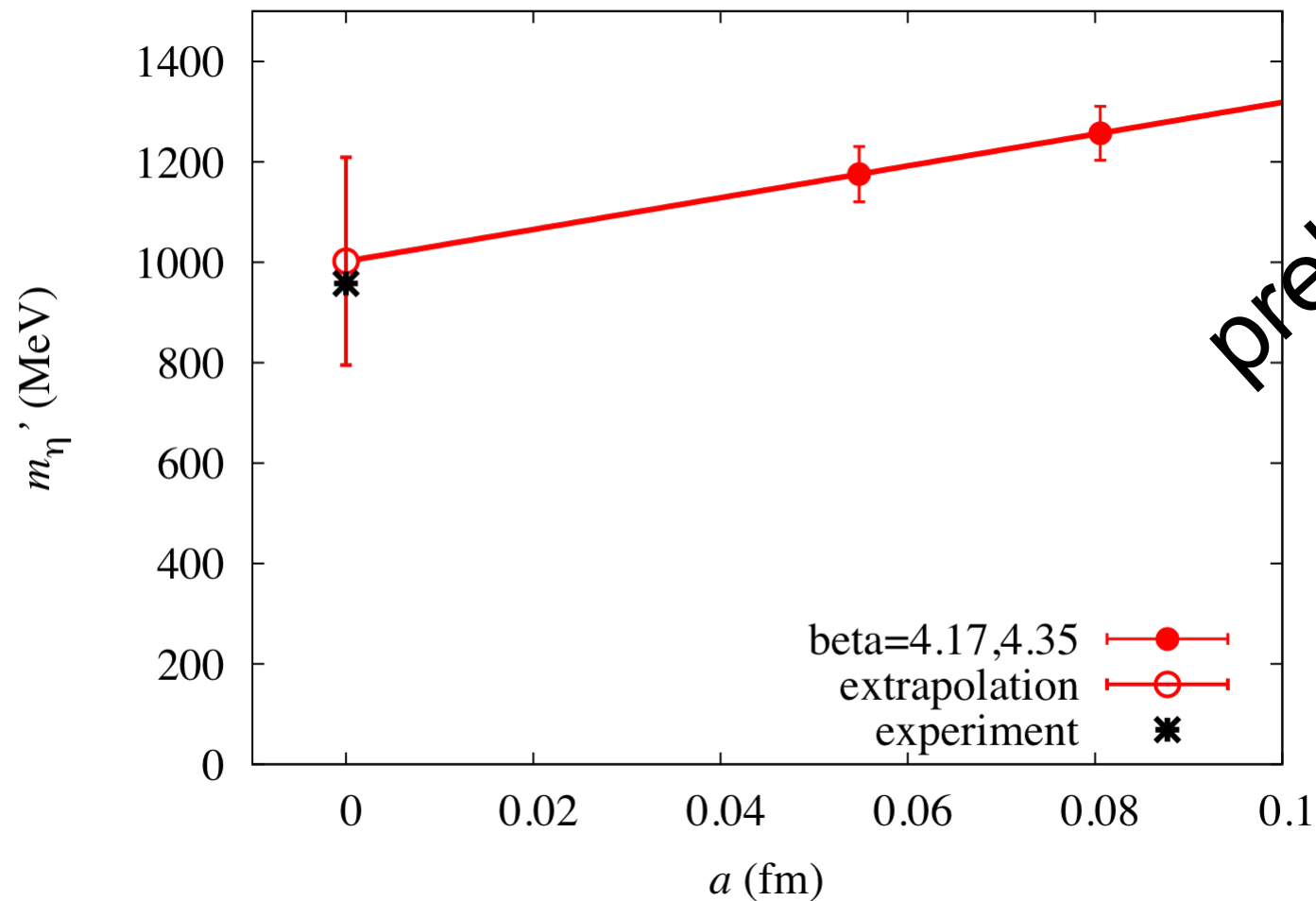
Fit with the chiral condensate

$$\Sigma^{1/3} = 250\text{MeV}.$$

$$\beta = 4.17 \quad (a^{-1} \sim 2.4\text{GeV})$$

# Other systematics

## Continuum limit

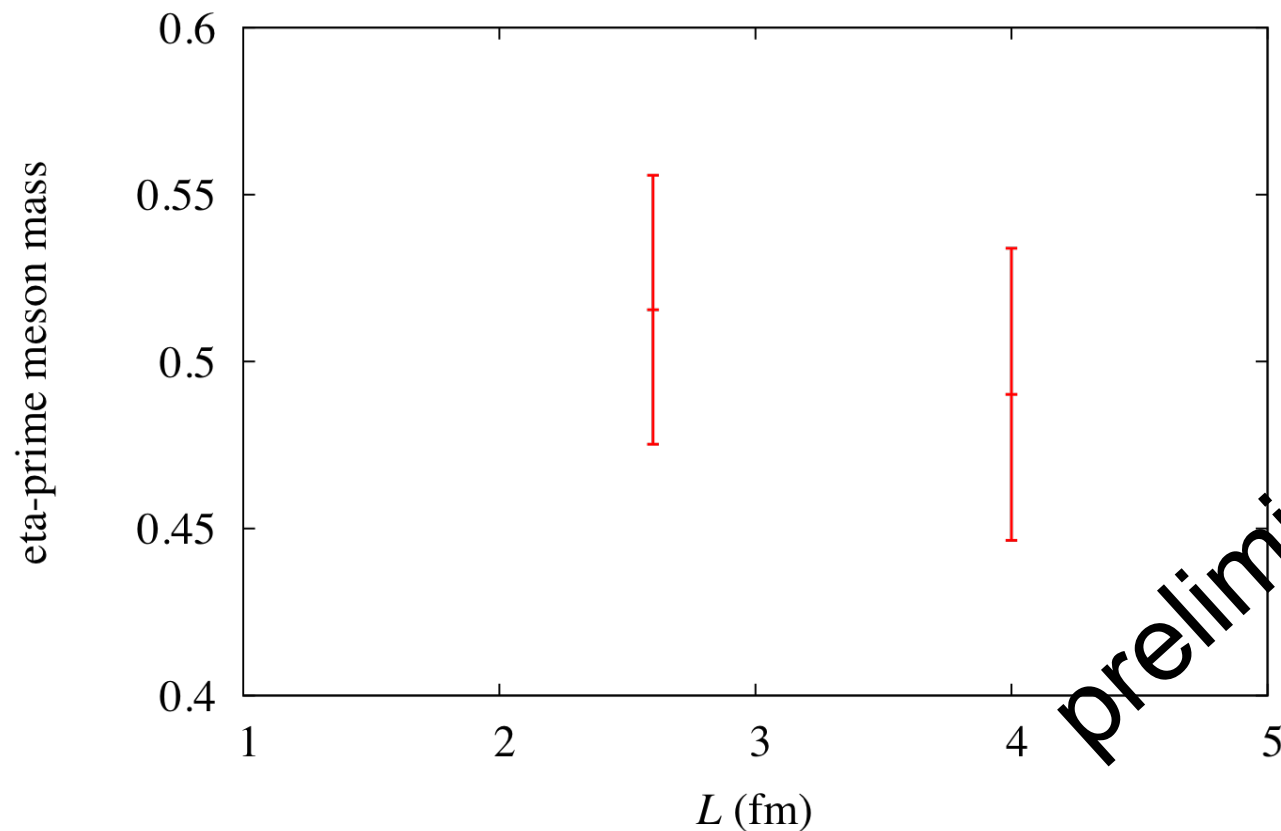


preliminary

# Other systematics

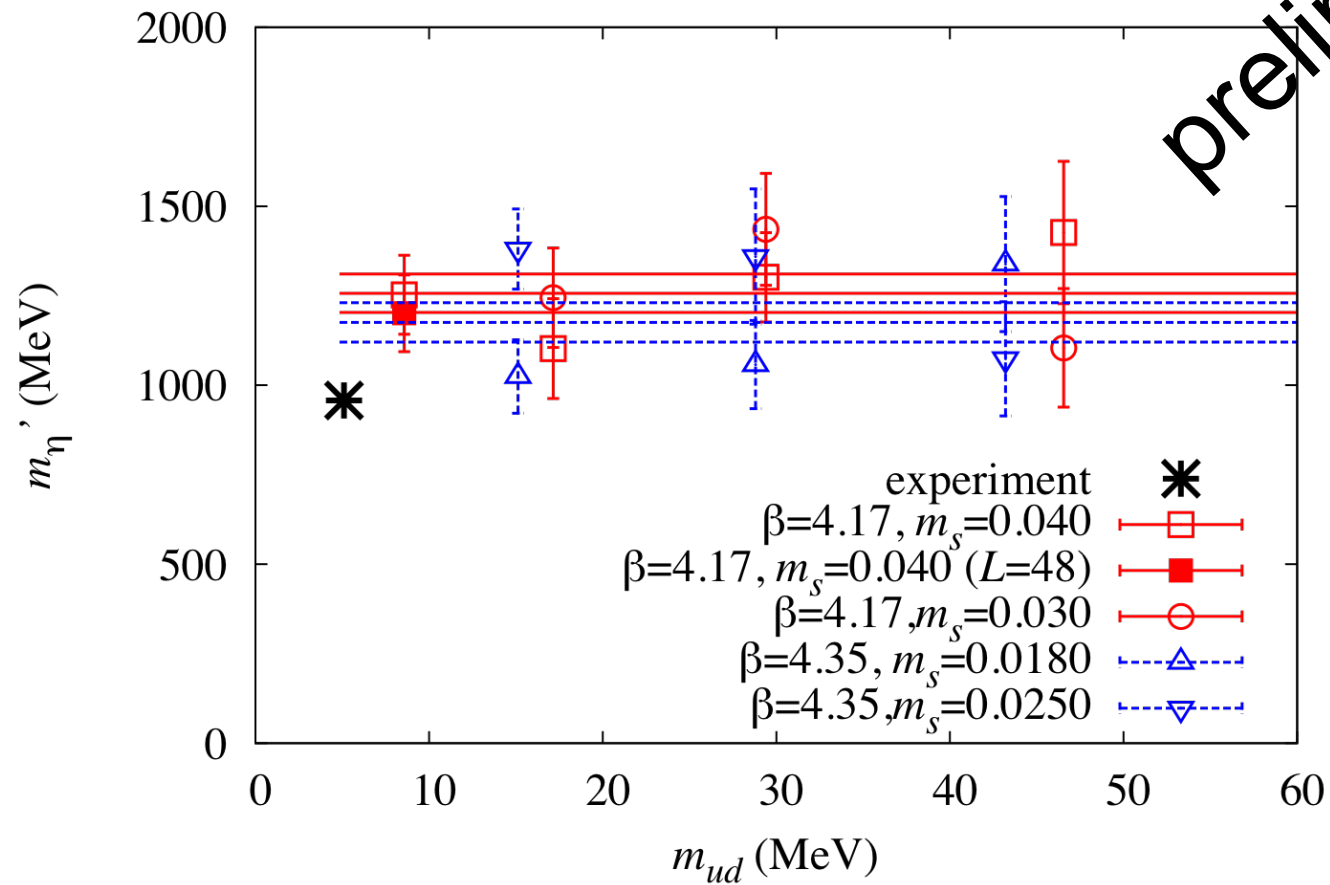
## Finite $V$ scaling at the lightest mass

b4.17 M1.00 mud0.0035 ms0.040



# Other systematics

## Chiral limit



preliminary