Combining ordinary and topological finite volume effects for fixed topology simulations

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Introduction	Physical results from fixed topology simulation	Including ordinary finite volume effects	Summary

Outline



2 Physical results from fixed topology simulation

- Correlators and partition function at fixed topology
- Extracting physical mass from fixed topology
- Numerical results: pure SU(2)
- Extracting topological susceptibility from fixed topology

Including ordinary finite volume effects

- Ordinary finite volume effects
- Combining ordinary and topological finite volume effects

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Motivation

Problem 1

Topology freezes for a too small lattice spacing $a < 0.05 fm^a$

^aLuscher, Martin JHEP 1008 (2010) 071

Problem 2

Topology fixed in purpose for some algorithms to avoid technical problems (Overlap)^a

^a H. Fukaya, et al. Phys. Rev. D 73 (2006) 014503.

Problem 3

Mixed action: different near-zero modes between sea Wilson quarks and valences overlap quarks: use only trivial topology $(Q=0)^{a}$

^a K. Cichy, G. Herdoiza and K. Jansen, Nucl. Phys. B 847, 179 (2011)

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Topology and Correlators

To take into account topology in your action

• Action with topological term:

$$S_E(\theta) = S_E - i\theta \frac{g^2}{32\pi^2} \int F_{\mu\nu} \tilde{F}_{\mu\nu} = S_E - i\theta Q[A]$$

- θ-term does not modify EOM
- Experimental measure of θ_{QCD} : $\theta_{QCD} \approx 0$

Path integral at fixed Q does not correspond to a physical theory

- No Hamiltonian!
- We can still defined a partition function, correlators and masses at fixed topology.

$$Z_Q = \int Z(\theta) e^{(-i\theta Q)}$$

$$Z_Q C_Q = \int Z(\theta) C(\theta) e^{(-i\theta Q)}$$

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Mass at fixed topological charge

• The mass at fixed topology of particles is given by the BCNW equation ¹:

$$M_Q = M(0) + \frac{M''(0)}{2\chi_T V} \left(1 - \frac{Q^2}{\chi_T V}\right) + \mathscr{O}\left(\frac{1}{(\chi_T V)^2}\right)$$

• Fixing the topology implies finite volume effects (TFV effect)

• Expansion on
$$\frac{1}{\chi_T V}$$
 and $\frac{Q^2}{\chi_T V}$
• $\chi_T V \in \{1, \dots, 10\}$ for our studies

¹Brower, R. et al. Phys.Lett. B560 (2003) 64-74, Aoki, Sinya et al□Phys.Rev. D76 (2007) 054508 → Q C

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Method

$$M_Q = M(0) + rac{M''(0)}{2\chi_T V} \left(1 - rac{Q^2}{\chi_T V}\right)$$

Method to extract physical mass:

- Compute M_Q using only configurations in a single topological sector for different volumes and topological charges.
- **2** Fit M(V, Q) with the BCNW-equation for $\chi_T V > max(|Q|, 1)$.

• get the parameters: $M(\theta = 0), \chi_T$, and M''(0).

To test this fixed topology method:

Compute M(0) using traditional method (unfixed and unfrozen topology)

- 2 Extract M(0) using fixed topology simulations
- Compare 1. and 2.

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SU(2) Yang-Mills Theory

$$\mathscr{L} = \frac{1}{4} F_{\mu\nu} F_{\mu\nu}$$

Set-up

- Observable: static potential $V_{qq}(R)$ for R = 1 to 6
- $\beta = 2.5$ which correspond to $a \approx 0.073 fm$
- Volumes: 14⁴,15⁴,16⁴ & 18⁴
- Number of configurations: 4000 per volume

0		Q	= 0		,	1	1	1	
1	-	н	Q =	1					-
2	ŀ		 0	Q = 2					-
3	ŀ			н С	Q = 3				-
4	-				—	Q = 4			-
5	ŀ					-		Q	= 5 -
all	╞		₩ ui	nfixed Q					-
0.	305	,	0.31	0.315	0.32	0.325	0.33	0.335	0.34
					$\mathcal{V}_{q\bar{q}}a($	r/a = 6)			

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SU(2) Yang-Mills Theory

$$M_Q = M(0) + rac{M''(0)}{2\chi_T V} \left(1 - rac{Q^2}{\chi_T V}
ight)$$



 $V_{qq}(R=6) = 0.3097(5)$ from fixed Q $V_{qq}(R=6) = 0.3101(3)$ unfixed topology simulation (ref)

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The AFHO² equation

• The $AFHO^2$ -equation

$$\langle q(x)q(0)\rangle_Q \approx_{x\to\infty} - \frac{\chi_T}{V}\left(1-\frac{Q^2}{\chi_T V}\right)$$

- The topological susceptibility can be extracted from the fit of a plateau for large x
- Need only one topological sector and one volume
- Results in agreement with unfixed and literature references



² S. Aoki, H. Fukaya, S. Hashimoto and T. Onogi, *Phys. Rev.* **D76** (2007) 054508.

Summary of these results

- Fixing Q results in topological finite volume effects (TFV).
- 2 The method is working well to extract the mass under the condition that $\chi_T V > max(|Q|, 1)$
- Interpretation of the state of the state
- Possibility to combine both:
 - to extract mass with a better precision
 - to extract mass from only one volume.

<u>Additional difficulties:</u> TFV effects are in competition with ordinary finite volume effects (OFV): short window to apply the method \Rightarrow Including OFV effects in the equation to increase the window

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Ordinary finite volume effects

Ordinary finite volume effects are due to:

- simulation at finite volume
- periodic boundary conditions

 $\Rightarrow The particle can interact with an image of itself !$

Ordinary finite volume effects (OFV) on a particle of mass M:

SU(N) equation:
$$M_L - M_{L=\infty} \propto \frac{1}{L} e^{-\frac{\sqrt{3}}{2}mL}$$

QCD equation:
$$m_{\pi,L} - m_{\pi,L=\infty} \propto \frac{1}{L} K_1(m_{\pi}L)$$

with m: mass of the lightest particle, L: length of the box

Ordinary finite volume effects in QCD:

- $\bullet~\mbox{Extremely costly to generate configurations} \Rightarrow \mbox{small volumes}$
- *m*_π is small in QCD.
 ⇒Difficulties to get rid of ordinary finite volume effects

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Combining ordinary and topological finite volume effects

- Working at fixed topology with ordinary finite volume effects: need to combine both kind of finite volume effects.
 - Need to calculate OFV in θ-vacuum using Lüscher method (for SU(N) and QCD)
 - Calculate the mass at fixed topology with OFV.
- Leading order (LO): BCNW-equation , OFV

$$M_{Q,L}^{SU(N)} = M(0) + \frac{M''(0)}{2\chi_T V} \left(1 - \frac{Q^2}{\chi_T V}\right) - \frac{A}{m^2(0)L} e^{-\frac{\sqrt{3}}{2}m(0)L} + \mathcal{O}\left(\frac{e^{-\frac{\sqrt{3}}{2}m(0)L}}{(\chi_T V)}\right)$$
$$m_{\pi,Q,L} = m_{\pi}(0) + \frac{m_{\pi}^{(2)}(0)}{2\chi_T V} \left(1 - \frac{Q^2}{\chi_T V}\right) + \frac{Bm_{\pi}(0)}{L} \mathcal{K}_1(m_{\pi}(0)L) + \mathcal{O}\left(\frac{\mathcal{K}_0(m_{\pi}(0)L)}{(\chi_T V)}\right)$$

- NLO: OFV effects will depend of Q.
- Our result is in agreement with result obtained for QCD with ChPT calculation in $\theta\text{-vacuum}^{3}$

³S. Aoki and H. Fukaya, Phys. Rev. D **81**, 034022 (2010) → < = → < = → へ ⊂

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Ordinary finite volume effects at fixed topology



- Ordinary finite volume effects for V < 14⁴ (discrepancy with the BCNW fit)
- Different OFV effects for different topological charges ⇒Need to go to next leading order

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Ordinary finite volume effects at fixed topology



• Fit of the next leading order (NLO) equation combining topological and ordinary finite volume effects

	$\hat{\mathscr{V}}_{q\bar{q},Q,V}(r=3a)$	ĥ	$\hat{\chi}_{T} imes 10^{5}$
fit results, NLO eq.	0.16437(15)	0.67(10)	9.5(2.0)
unfixed top. results	0.16455(7)	0.723(23)	7.0(0.9)

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Summary

- Study the possibility to work at fixed topology
 - Show the efficiency of the method to extract mass from frozen topology simulation
 - Precise results obtained on the mass
 - Good results to obtained topological susceptibility
- Combination of ordinary finite volume effects and topological finite volume effects
 - Equation combining both finite volume effects
 - Promising test on SU(2) Yang-Mills theory
 - Outlook
 - More tests on the combination of ordinary finite volume effects and topological finite volume effects.

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• Full QCD (in-going)