

Lattice calculation of the HVP contribution to the anomalous magnetic moment of muon

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HPQCD Collaboration*

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Current status of muon g-2

- BNL E821 0.54 ppm; planned 4-fold improvement in Fermilab E989.
- Overall theoretical uncertainty in a_μ : 0.42 ppm.

Contribution	Result ($\times 10^{-10}$)	Error
QED (leptons)	11658471.8	0.00 ppm
HVP(lo) [1]	692.3	0.36 ppm
HVP(ho)	-9.8	0.01 ppm
HLbL [2]	10.5	0.22 ppm
EW	15.4	0.02 ppm
Total SM	11659180.2	0.42 ppm

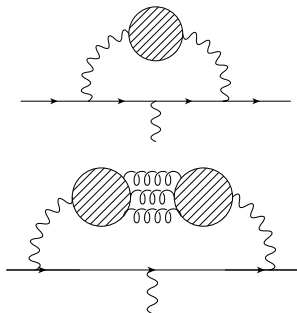
[1] Davier et. al., Eur.Phys.J. C71 (2011) 1515 : dispersion relation + cross section for e^+e^- (and τ) \rightarrow hadrons with 0.7% uncertainty; [2] Prades, et. al., 0901.0306.

- g-2 discrepancy around 3σ between the SM and the expt.:

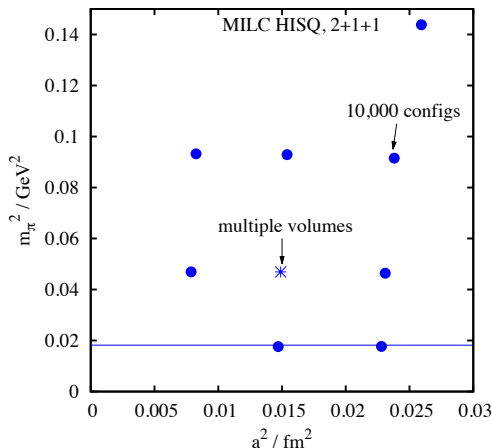
$$a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = 25(9) \times 10^{-10}.$$

- Calculate lowest order hadronic vacuum polarisation $a_\mu^{\text{HVP,lo}}$ from lattice QCD with $< 0.5\%$ uncertainty by 2017/18.

HVP (conn. + disc.)

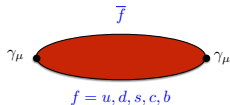


HPQCD analysis ingredients



- Light (up/down), strange, charm (2+1+1) quarks in sea:
 $m_l = m_s/5, m_s/10, m_s/27.5$ (physical).
- $a \approx 0.15\text{fm}$ (very coarse), 0.12fm (coarse), 0.09fm (fine).
- Large box size: 5-6 fm on the finest lattices; finite volume effects tested.
- High statistics: random color wall sources, 1000 (or 10,000) configurations and 16 time sources.

HPQCD analysis method



$$a_{\mu, \text{HVP}}^{(f)} = \frac{\alpha}{\pi} \int dq^2 f(q^2) (4\pi\alpha Q_f^2) \hat{\Pi}_f(q^2).$$

(Lautrup,'72, T. Blum,'02)

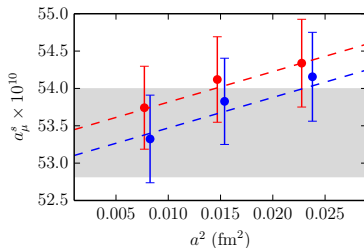
- Reconstruct $\hat{\Pi}(q^2)$ from its derivatives at $q^2 = 0$ + Padé approximants [2,2] to calculate vacuum polarization at arbitrary Euclidean q^2 :

$$\hat{\Pi}(q^2) = \sum_{j=1}^{\infty} q^{2j} \Pi_j$$

- Time moments ($n = 4, 6, 8, 10$) of (local-local) vector meson correlators give the derivatives:

$$\begin{aligned} G_{2n} &\equiv a^4 \sum_t \sum_{\vec{x}} t^{2n} Z_V^2 \langle j^i(\vec{x}, t) j^i(0) \rangle \\ &= (-1)^n \left. \frac{\partial^{2n}}{\partial q^{2n}} q^2 \hat{\Pi}(q^2) \right|_{q^2=0}. \end{aligned}$$

$$\text{And } \Pi_j = (-1)^{j+1} \frac{G_{2j+2}}{(2j+2)!}.$$



- 1% precision for strange (and charm) quark case; vector correlators reproduce accurately physical properties of ϕ and J/ψ mesons [1208.2855, 1410.8466].

[Ref: Chakraborty et. al., PRD89(2014)114501, arXiv:1403.1778]

New issues for 1% uncertainty in u/d connected case: Handling ρ meson on lattice

- Correlators much noisier: Use data-fit hybrid correlator to control noise at large t :

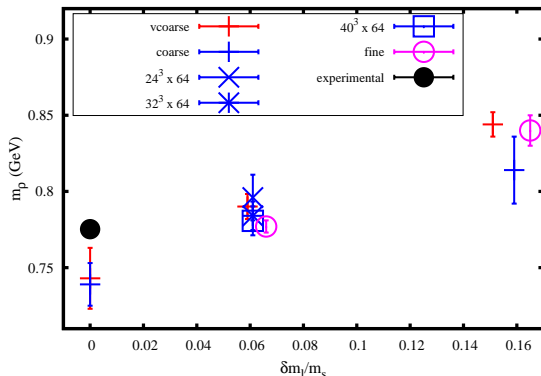
$$G(t) = \begin{cases} G_{data}(t) & \text{for } t \leq t^* \text{ from Monte Carlo,} \\ G_{fit}(t) & \text{for } t > t^* \text{ from multi-exponential fit.} \end{cases}$$

for $t^* = 1.5\text{fm} = 6/m_\rho$ (70% result from G_{data})
(same results to within $\pm\sigma/4$ with 0.75fm).

- Gaussian smearings to obtain precise ρ ground state (25-30% gain in uncertainty in $a_\mu^{u/d}$).
- 80% of light quark vacuum polarisation contribution from ρ meson pole, need to understand ρ better.

New issues for 1% uncertainty in u/d connected case: Handling ρ meson on lattice

- Finite volume, finite lattice spacing errors: m_ρ and f_ρ too low; Significant impact on $g-2$.



- Rescale Π_j^{latt} : Reduces or eliminates errors from scale setting, Z factors, a^2 , finite volume, m_l/m_s tuning (elaboration of ETMC idea in arXiv:1308.4327).

Rescaling to Taylor coefficients

- At large time we expect C_{ll} to behave as:

$$C_{ll,gs} = \frac{f_\rho^2 m_\rho}{2} e^{-m_\rho t}$$

- Π_j behaves as:

$$\Pi_{j,c} = (-1)^{j+1} \frac{f_\rho^2}{m_\rho^{2j+2}}$$

- Correction to the Taylor coefficients Π_j :

$$\Pi_j^{latt} \rightarrow \Pi_j^{latt} \left[\frac{m_\rho^{2+2j}}{f_\rho^2} \right]_{latt} \left[\frac{f_\rho^2}{m_\rho^{2+2j}} \right]_{expt}$$

- Experimental f_ρ is obtained from the decay width of $\rho \rightarrow e^+ e^-$ using:

$$\Gamma(\rho \rightarrow e^+ e^-) = \frac{4\pi}{3} \alpha_{QED}^2 \frac{f_\rho^2}{m_\rho} e_{u/d}^2$$

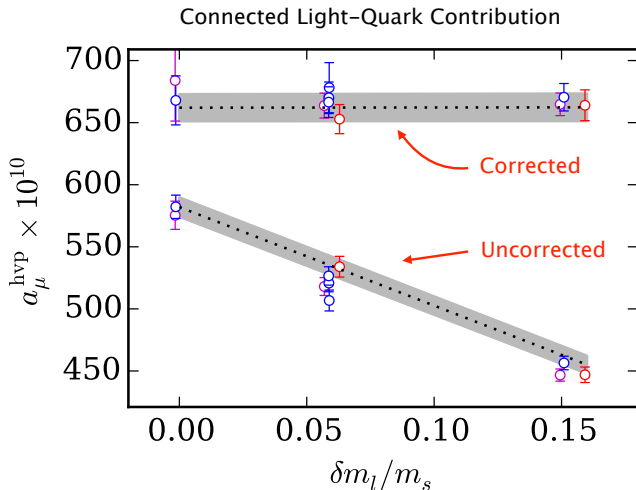
New issues for 1% uncertainty in u/d connected case: Handling $\pi\pi$ loop contribution on lattice

- $\pi\pi$ loop contribution $\sim 10\%$ of the total light quark contribution.
- Highly sensitive to m_π (contribution roughly proportional to $1/m_\pi^2$) and finite volume.
- Staggered quark introduces extra discretisation artefacts from different taste π meson.
- Remove lattice $\pi\pi$ contribution using one-loop, staggered quark, finite-volume ChPT.
- Restore $\pi\pi$ contribution using one-loop continuum ChPT, with physical π mass.
- Scaling gives:

$$\Pi_j^{\text{latt}} \rightarrow (\Pi_j^{\text{latt}} - \Pi_j^{\text{latt}}(\pi\pi)) \left[\frac{m_\rho^{2+2j}}{f_\rho^2} \right]_{\text{latt}} \left[\frac{f_\rho^2}{m_\rho^{2+2j}} \right]_{\text{expt}} + \Pi_j^{\text{cont}}(\pi\pi)$$

HPQCD preliminary Results: $a_\mu^{u/d}$ (Connected)

- Corrected results independent of m_l/m_s , a^2 , finite volume.
- [2, 2] Padé approximants converged to better than 0.2%: negligible error.
- Total uncertainty is around 1.7%; statistics dominate.



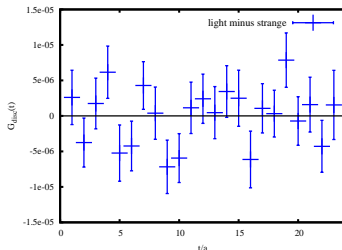
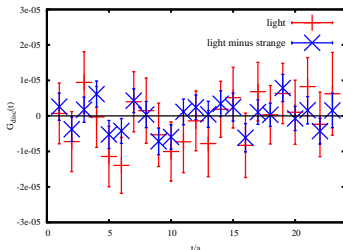
Disconnected correlators



- On lattice we calculate:

$$G_{disc}(x_0 - y_0) = -Z_V \langle \left(\sum_{\vec{x}} \text{Tr}[\gamma_k D^{-1}(x, x)] \right) \left(\sum_{\vec{y}} \text{Tr}[\gamma_k D^{-1}(y, y)] \right) \rangle$$

- All-to-all propagator method with 50 stochastic noise vectors on each configuration, one-link spatial vector currents for (l-s) on both sides.



- correlators from $l - s$ vector currents at each end are much less noisy (consistent with Mainz group, arXiv:1411.7592); about 60% reduction in uncertainty compared to light quark case on $m_l/m_s = 0.2$ very coarse lattice.

HPQCD Estimation of Disconnected contribution to

 $a_{\mu}^{HVP,lo}$

- Disconnected u/d piece D_{ll} dominates (Hadspec collaboration (arXiv:1309.2608)).
- Assuming D_{ll} provide difference between ω and ρ , at large t,

$$2D_{ll,gs} = -\frac{f_{\rho}^2 m_{\rho}}{2} e^{-m_{\rho} t} + \frac{f_{\omega}^2 m_{\omega}}{2} e^{-m_{\omega} t}$$

- Ratio of the moments from D_{ll} to that from C_{ll} :

$$R_j = \frac{\Pi_{j,d}}{\Pi_{j,c}} = \frac{1}{2} \left[\frac{m_{\rho}^{2+2j} f_{\omega}^2}{m_{\omega}^{2+2j} f_{\rho}^2} - 1 \right]$$

- From experiment, $m_{\rho} = 0.775$ GeV, $f_{\rho} = 0.217$, $m_{\omega} = 0.783$ GeV and $f_{\omega} = 0.195$ GeV:
 $R_1 = -0.025$; $R_2 = -0.026$; $R_3 = -0.028$ i.e. less than 3%.

$$a_{\mu, disc} / a_{\mu, conn} \approx -3\%.$$

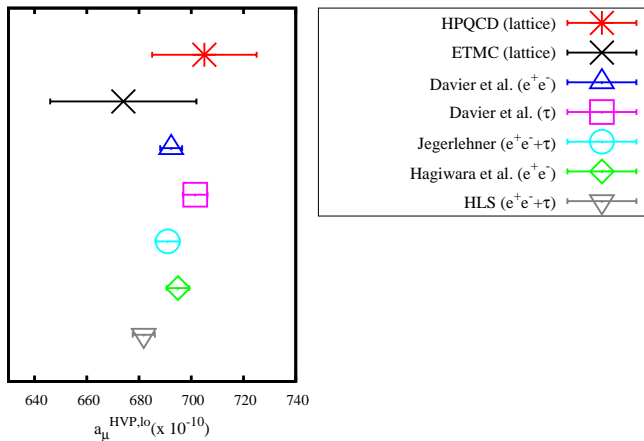
HPQCD Estimation for Total $a_\mu^{HVP,lo}$

Contribution	Result (x 10 ⁻¹⁰)	
light, conn	662(11)	(preliminary)
strange, conn	53.4 (6)	(1403.1778)
charm, conn	14.4(4)	(1403.1778, 1208.2855)
bottom, conn	0.27(4)	(1408.5768)
disconn. (estimate)	-25(15)	added -7 from $\pi\pi$ to simple estimate*
Total	705(20)	

[* $a_\mu^{(\pi\pi)} = 71 \times 10^{-10}$ and we considered -10% of this following partially quenched ChPT.]

The uncertainty in the total $a_\mu^{HVP,lo}$ is dominated by our estimation of disconnected piece.

Comparison of total $a_\mu^{HVP,lo}$ with phenomenology



Conclusion

To summarize:

- HPQCD preliminary for total $a_\mu^{HVP,lo}$ using HISQ : $705(20) \times 10^{-10}$ ($\sim 2.8\%$ uncertainty)
- First principle result including 2+1+1 dynamical sea quarks, physical valence quarks and pion masses, takes finite volume effects into account.

Future:

- Collaborating with MILC for using large ensemble size (10,000 on $\sim 0.12\text{fm}$ lattices).
- Need to understand ρ meson better on lattice.
- Get a better precision from disconnected diagrams : collaborating with Hadron Spectrum collaboration.