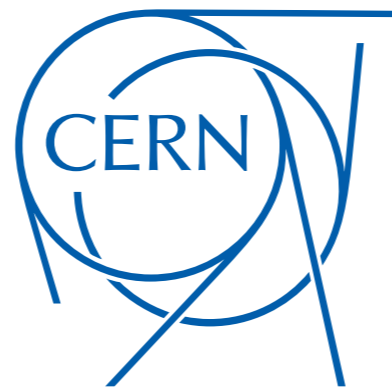


Leading isospin breaking correction to the HVP

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Motivation I: computing IB correction to the HVP

- Discrepancy between $a_\mu^{exp} - a_\mu^{th,SM}$ mainly coming from hadronic contributions
- Once the aimed precision ($<1\%$) for the connected HVP from the lattice is achieved (in the isosymmetric theory) \rightarrow the effects we neglected so far might become important:
 - disconnected contribution,
 - isospin breaking corrections,
 - charm in the sea, ...
- In the phenomenological determination of a_μ^{had} , model calculation of [Jegerlehner,Szafron '11]
 - \rightarrow correctly applied IB correction reduced the discrepancy between e^+e^- and τ data
- Not clear how this translates to the Euclidean
- It would be good to have a model independent estimate of IB effects: lattice QCD+QED
- Note: systematic analysis based on the τ data may also benefit knowing how big/small this effect is

Motivation II: the method to compute IBE

- All necessary ingredients are, in principle, there
- R123 method [[arXiv:1303.4896](#)] for computing leading isospin breaking corrections (LIBE)
 - ➔ Expanding an observable (in the isospin broken theory) with respect to the isosymmetric QCD result
- For a start: applying it to the connected part of the HVP
- Main advantage w. respect to simulating QED+QCD:
 - ➔ Diagrams obtained individually (before multiplying with $O(\alpha_{em})$, $O(m_u - m_d)$ coeff.)
 - ➔ No extrapolation in α_{em}

The method I: LIBE in practice (R123)

- Reusing the gauge configurations generated in the isosymmetric theory

- Reweighting:
$$\langle O \rangle_{\vec{g}} = \frac{\langle R[U, A; \vec{g}, \vec{g}^0] O[U, A; \vec{g}] \rangle^{A, \vec{g}^0}}{\langle R[U, A; \vec{g}, \vec{g}^0] \rangle^{A, \vec{g}^0}}$$

\vec{g}^0 - bare param. of isosymm. th
 \vec{g} - bare param. of the full th

- For simplicity, approximate sea quarks as electrically neutral: $R[U, A; \vec{g}, \vec{g}^0] = 1$

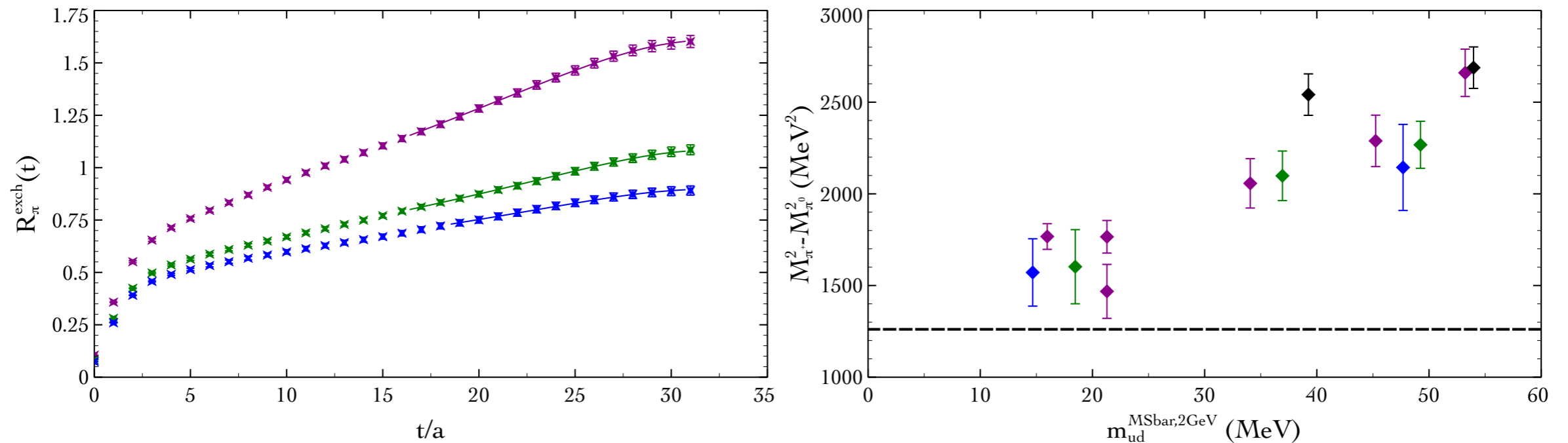
- ...once an appropriate renormalisation procedure is applied: $\Delta O = O(\vec{g}) - O(\vec{g}^0)$

- Example: $\Delta \longrightarrow \pm =$

$$(efe)^2 \xrightarrow{\text{wavy}} + (efe)^2 \xrightarrow{\text{star}} - [m_f - m_f^0] \text{---} \otimes \text{---} \mp [m_f^{cr} - m_0^{cr}] \text{---} \otimes \text{---}$$

$$\begin{aligned}
 & -e^2 e_f \sum_{f_1} e_{f_1} \xrightarrow{\text{wavy}} \text{---} \otimes \text{---} & -e^2 \sum_{f_1} e_{f_1}^2 \xrightarrow{\text{wavy}} \text{---} \otimes \text{---} & -e^2 \sum_{f_1} e_{f_1}^2 \xrightarrow{\text{wavy}} \text{---} \otimes \text{---} & + e^2 \sum_{f_1 f_2} e_{f_1} e_{f_2} \xrightarrow{\text{wavy}} \text{---} \otimes \text{---} \\
 & + \sum_{f_1} \pm [m_{f_1}^{cr} - m_0^{cr}] \text{---} \otimes \text{---} & + \sum_{f_1} [m_{f_1} - m_{f_1}^0] \text{---} \otimes \text{---} & + [g_s^2 - (g_s^0)^2] \xrightarrow{\text{---} \otimes \text{---}} \text{---} \otimes \text{---} .
 \end{aligned}$$

The method II: LIBE in practice (R123)



Review by N. Tantalo @Lat2013

- Previous results by Rome123 collaboration [[arXiv:1303.4896](https://arxiv.org/abs/1303.4896), [arXiv:1311.2797](https://arxiv.org/abs/1311.2797)]
- Leading correction to different hadronic observables: pion/kaon mass splitting, Dashen theorem breaking parameter, u-d quark mass difference ...
- Corrections function of the ratios of the correlators in the full and isosymmetric theory and give good numerical signal

Technicalities

- Leading correction: expanding in powers of the difference between bare param. in full and isosymm. th:

$$\Delta O = \left\{ e^2 \frac{\partial}{\partial e^2} + [g_s^2 - (g_s^0)^2] \frac{\partial}{\partial g_s^2} + [m_f - m_f^0] \frac{\partial}{\partial m_f} + [m_f^{cr} - m_0^{cr}] \frac{\partial}{\partial m_f^{cr}} \right\} O$$

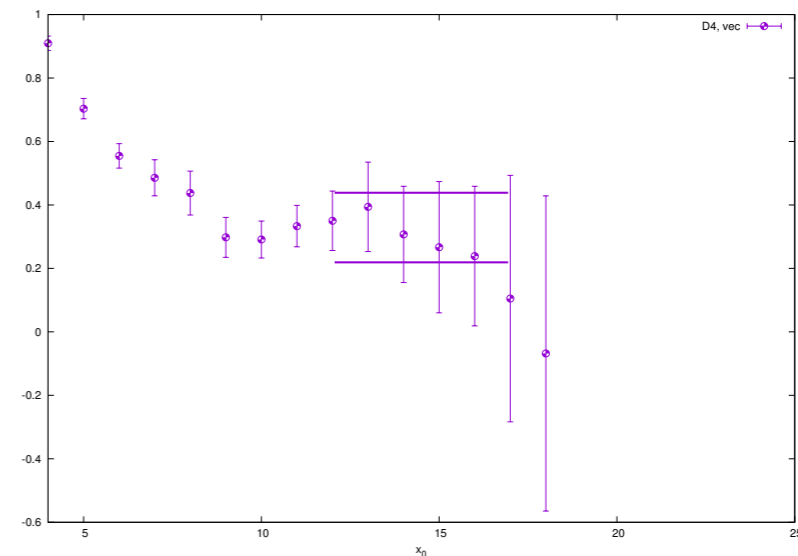
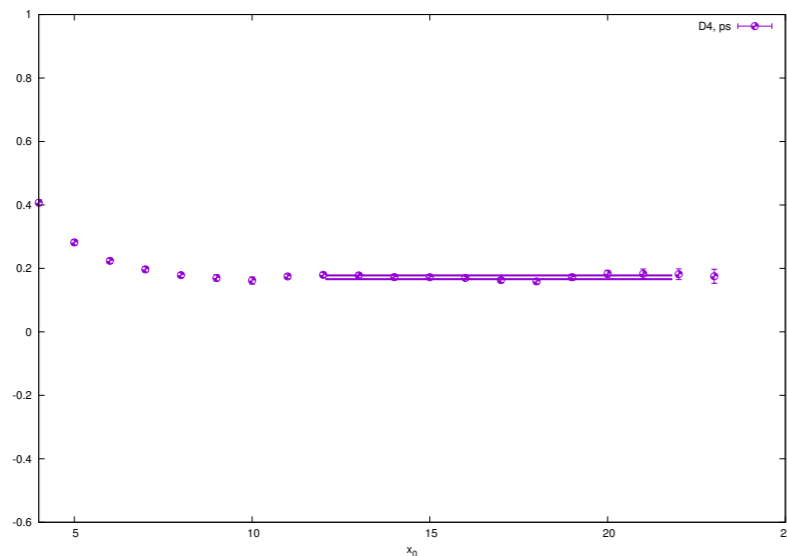
- Leading IB corrections are computed: also in QED+QCD simulations $O(\alpha(m_u - m_d))$ are neglected
- Main general obstacle in implementing this method
 - ➔ many diagrams need to be computed
 - ➔ including the 3-pt, 4-pt functions and the disconnected ones (beyond el-quenched approximaton)
- Implementation: requires careful organisation of the computation of the diagrams:

$$M_{K^+} - M_{K^0} = -2\Delta m_{ud} \partial_t \frac{\text{diagram 1}}{\text{diagram 2}} - (\Delta m_u^{cr} - \Delta m_d^{cr}) \partial_t \frac{\text{diagram 3}}{\text{diagram 4}}$$

$$+ (e_u^2 - e_d^2) e^2 \partial_t \frac{\text{diagram 5} - \text{diagram 6} - \text{diagram 7}}{\text{diagram 8}} + (e_u - e_d) e^2 \sum_f e_f \partial_t \frac{\text{diagram 9}}{\text{diagram 10}}$$

Pseudo-scalar vs. vector

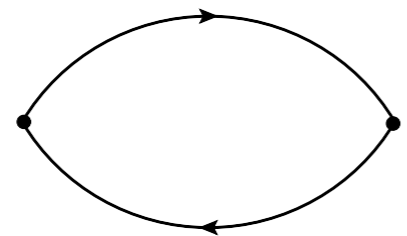
- We know that what works in pseudo-scalar channel
 - ➔ might not necessarily work that well in the vector one
- Example on two ensembles with $N_f=2$ $O(a)$ improved Wilson fermions (CLS based configurations)



- Lattice spacing $a \approx 0.07 fm$, pion masses:
 - ➔ D4: 48×24^3 , $m_\pi \approx 480 MeV$
 - ➔ E5: 64×32^3 , $m_\pi \approx 410 MeV$
 - ➔ D5: 48×24^3 , $m_\pi \approx 420 MeV$

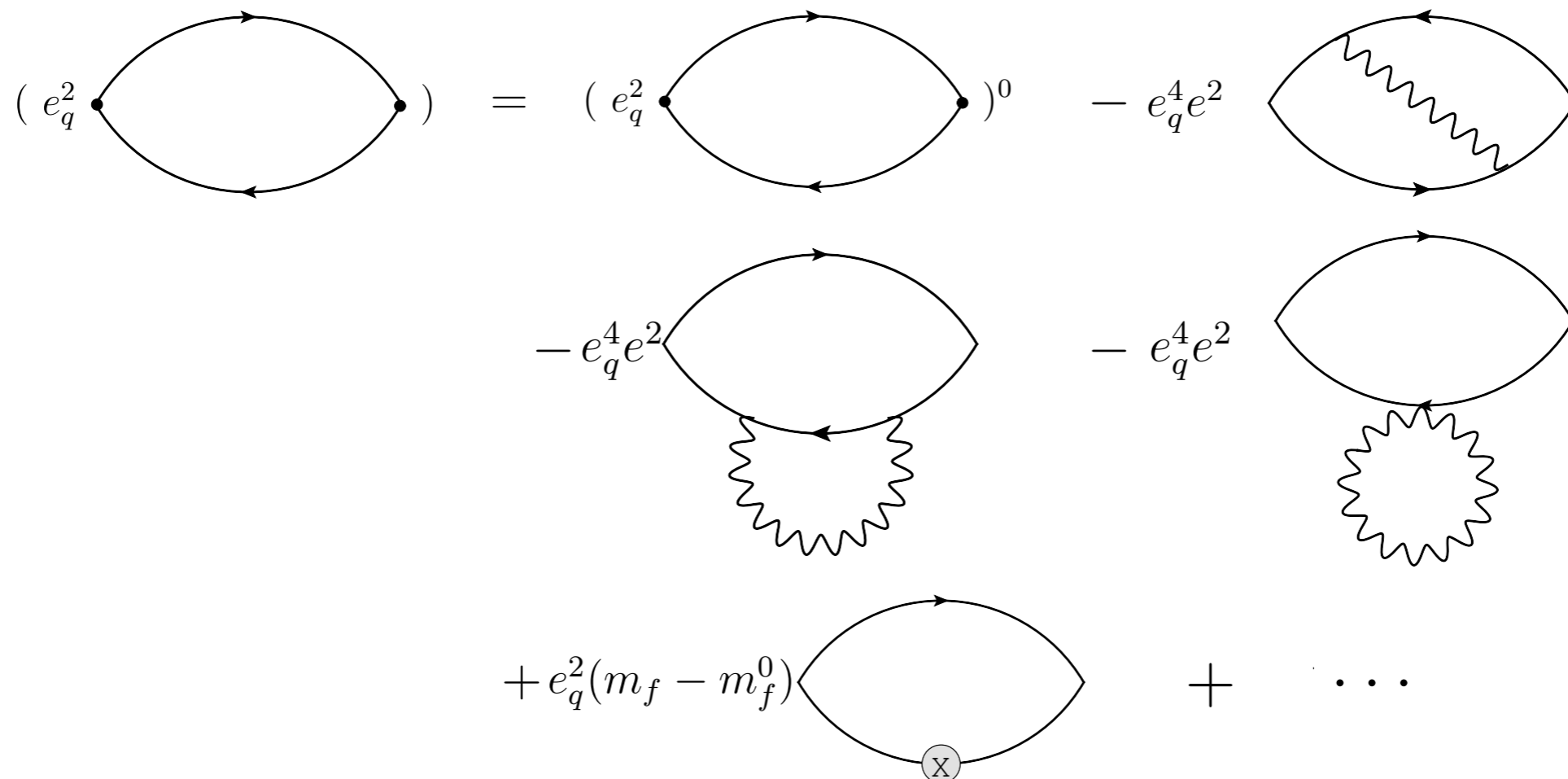
LIBE of the HVP in the electro-quenched approx.

- Expanding the connected part of the HVP



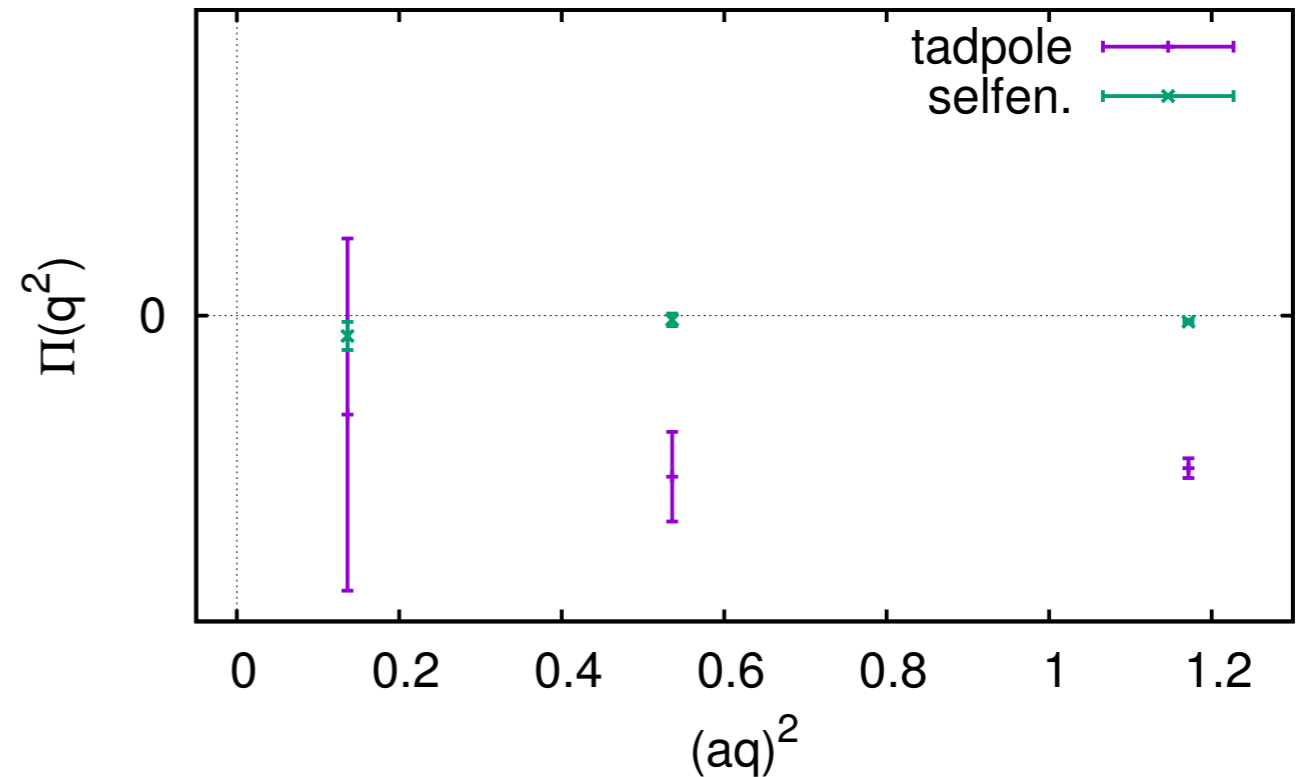
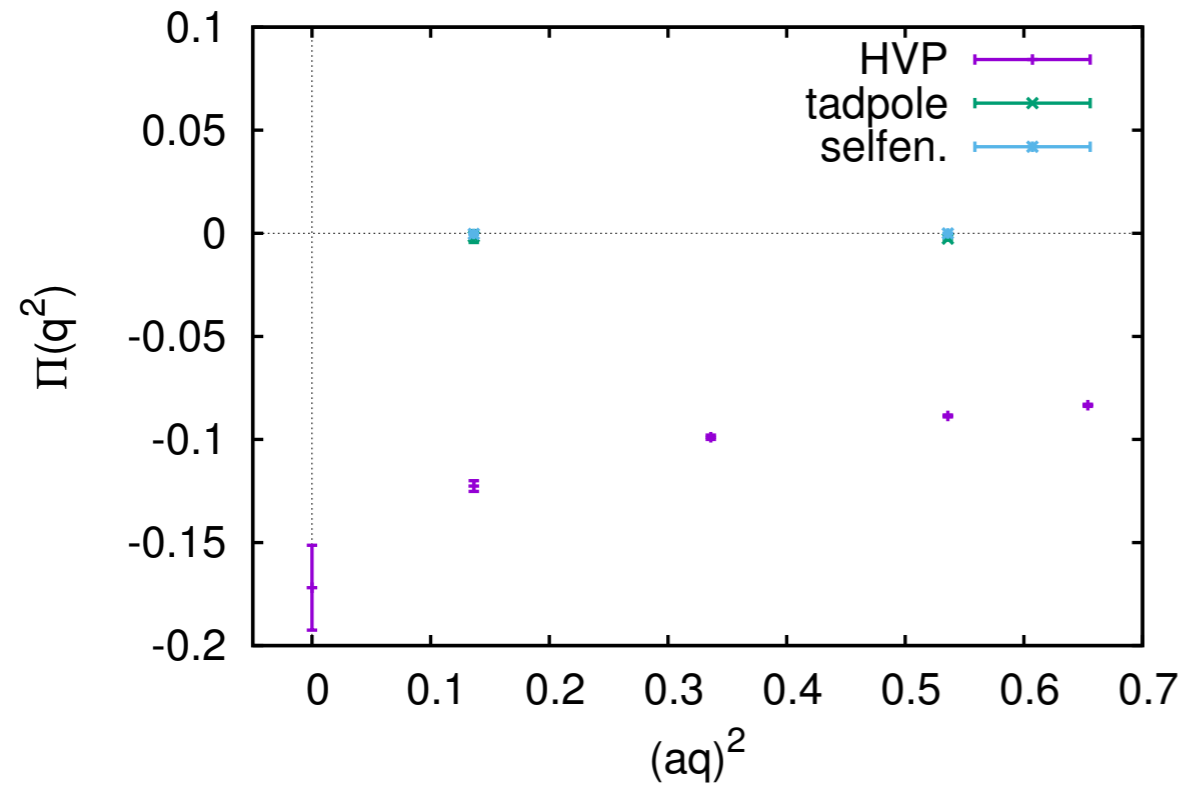
$$= \text{Tr}\{\gamma_\mu S_f \gamma_\nu S_f\}$$

- Electro-quenched approximation:



$$\begin{aligned}
 (e_q^2 \text{ loop}) &= (e_q^2 \text{ loop})^0 - e_q^4 e^2 \text{ loop with photon} \\
 &\quad - e_q^4 e^2 \text{ loop with gluon} - e_q^4 e^2 \text{ loop with ghost} \\
 &\quad + e_q^2 (m_f - m_f^0) \text{ loop with mass insertion} + \dots
 \end{aligned}$$

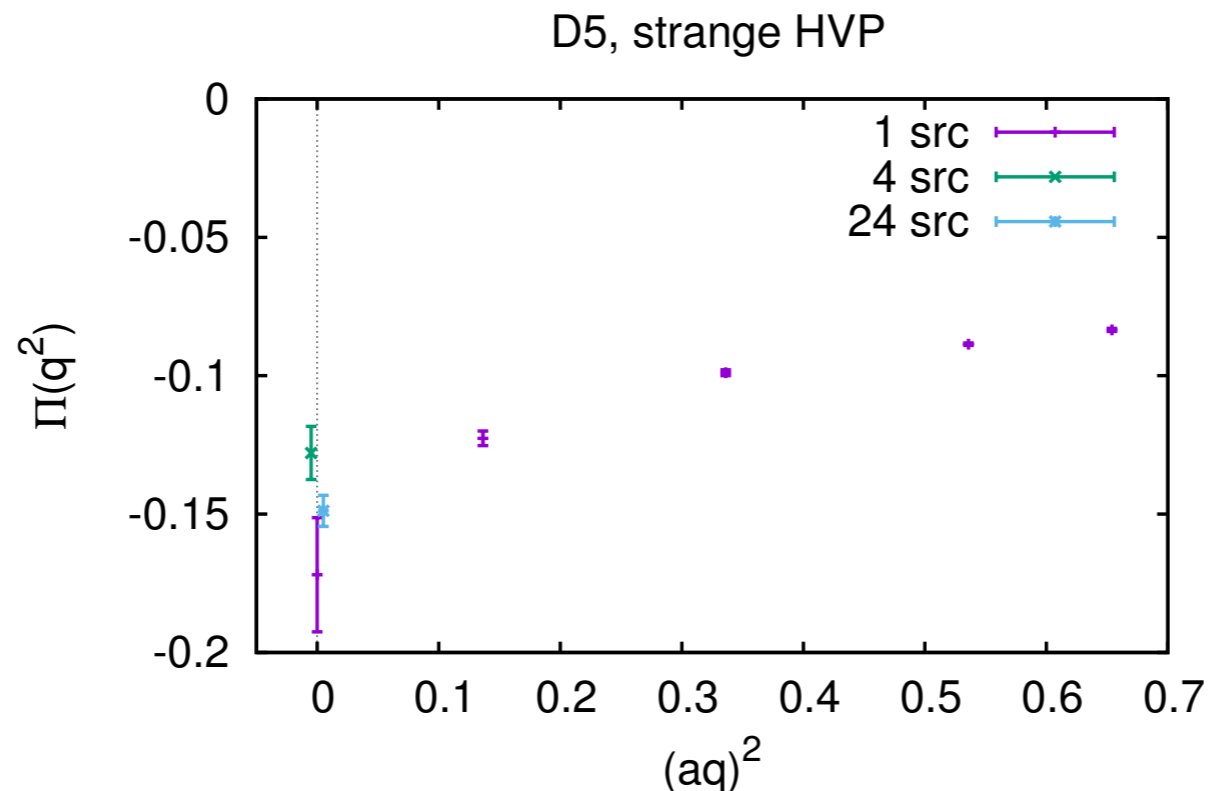
A first look at the signal/noise



- Strange HVP and EM corrections (exploratory study, same bare parameters)
- Their sum makes sense only after the renormalisation:
 - ➔ intermediate renormalisation prescription and matching procedure
 - ➔ using experimental determinations of the charged mesons to fix quark masses and the lattice spacing in the isosymmetric theory

Outline I: technical improvements

- Performing this first test on a moderate size cluster was possible partially due to the deflated solver from DD-HMC code [M. Lüscher '07]
- Nevertheless, upgrading the code to use the new openQCD solver: speedup needed for accumulating the statistics for the light contribution
- This study is performed with point source
 - ➔ indications that replacing the point with some extended source might be advantageous



- ➔ 1src, 4src, 24src on 83, 64, 41 configurations, $\Pi(0)$ as proposed in [1208.5914]

Outline II: conceptual improvements

- Reducing finite volume effects - they are expected to be strong
- Currently: global zero mode subtracted: $A_\mu(k=0) \equiv 0$
 - ➔ Violates reflection positivity and does not have a well defined $T \rightarrow \infty$ limit [1406.4088]
- Removing the zero mode of the field on each time slice separately [Hayakawa, Uno '08]
 - ➔ this explicitly violates the hypercubic symmetry of the lattice -> no trace of the violation in the inf. vol limit [1406.4088]
- Charged particles in QED/QED+QCD with C^* BCs \rightarrow FV effects even smaller
 - ➔ [talks by A. Patella and N. Tantalo: Thurs., 11.20, 11.40, s.402]
- This would be the way to go here as well, although the FV corrections for this quantity have not yet been explicitly computed in any of the above mentioned setups
- Getting the disconnected contributions (beyond el-quenched)

Conclusions

- Phenomenologically - IB plays an important role in the th.-exp. discrepancy
- Although the aimed precision of the HVP determinations from the lattice is not yet there to see this effect
 - ➔ useful to think ahead and work towards this estimate
 - ➔ all ingredients are there
- This is a first attempt to extract the IB correction to the HVP from first principles
- Difficult task, but R123 method should give better signal over simulating full theory (larger contributions are being computed)
 - ➔ QCD+QEDq certainly worth trying
- Even after (and if) getting the signal in the light sector —> many things to understand before a definite answer is known
 - ➔ 1. Finite volume effects
- Stay tuned ...

Thanks!

- Nazario for the (great!) code basis, discussions and motivation to look at the IB effects in this way
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BACKUP I: Different strategy than Pi(0)

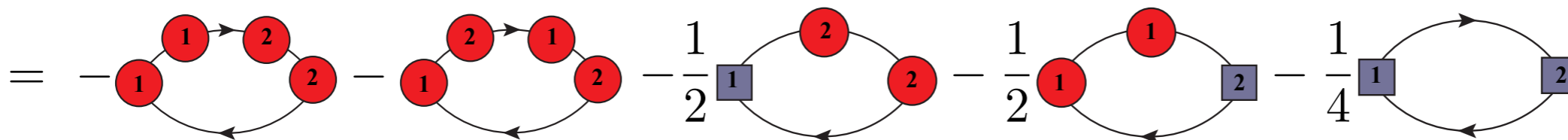
- But same machinery needs to be implemented [de Divitis, Petronzio, Tantalò1208.5914]:

- $\Pi_{12}(Q) = \sum_x \langle \text{Tr} \{ S[y, x; U] \Gamma_V^1(x, \vec{q}) S[x, y; U, \lambda^p] \Gamma_V^2(y, \vec{0}) \} \rangle.$

- $\Pi(0) = - \frac{\partial \Pi_{12}(Q)}{\partial Q_1 \partial Q_2} \Big|_{Q_s=0}$

$$= - \frac{1}{(TL^3)^2} \sum_{x,y} \langle \text{Tr} [S \Gamma_V^1 \frac{\partial^2 S}{\partial Q_1 \partial Q_2} \Gamma_V^2] - \frac{i}{2} \text{Tr} [S \Gamma_T^1 \frac{\partial S}{\partial Q_2} \Gamma_V^2]$$

$$- \frac{i}{2} \text{Tr} [S \Gamma_V^1 \frac{\partial S}{\partial Q_1} \Gamma_T^2] - \frac{1}{4} \text{Tr} [S \Gamma_T^1 S \Gamma_T^2] \rangle$$



- Consequently: tricks learned in one computation can be used in another and vice versa

BACKUP II: Tuning the critical mass

- Using WTI :
 - ➔ Dashen theorem: $\hat{m}_f = \{\hat{m}_u, \hat{m}_d, \hat{m}_s\} = 0$
 - ➔ Also with EM, in the massless theory:
 - $M_{\{\pi^0\}} = M_{\{K^0\}} = 0$
 - $M_{\{\pi^+\}} = M_{\{K^+\}} = 0$
 - $\kappa_s^{\text{crit}} = \kappa_d^{\text{crit}}$
- Need to be done with high accuracy, in order to cancel linear ultraviolet divergencies

$$\Delta m_f^{cr} = -\frac{e_f^2}{2} e^2 \lim_{\hat{m}_f \rightarrow 0} \frac{\partial_t \left(\text{Diagram 1} + 2 \text{Diagram 2} + \text{Diagram 3} \right)}{\partial_t \left(\text{Diagram 4} \right)}$$

The diagram shows the calculation of the critical mass shift Δm_f^{cr} . The numerator is the derivative with respect to the fermion mass \hat{m}_f of the sum of three diagrams: a fermion loop with a wavy photon line, two fermion loops with a starburst photon line, and a fermion loop with a starburst photon line. The denominator is the derivative with respect to \hat{m}_f of a fermion loop with a red cross on the photon line.

BACKUP II: Subtracting the zero mode

- Illustration of the reduced finite-volume effects, once zero mode of the photon field subtracted time-slice by time-slice (L) instead of the global zero mode subtraction (TL)
- Note that translation invariance is violated here, but it is recovered in the continuum limit [BMW,1406.4088]

