

Polyakov line actions from SU(3) lattice gauge theory with dynamical fermions: first results via relative weights

Roman Höllwieser^{ab}, Jeff Greensite^c

^aInstitute of Atomic and Subatomic Physics, Nuclear Physics Dept.,
Vienna University of Technology, Operngasse 9, 1040 Vienna, Austria

^bDepartment of Physics, New Mexico State University,
Las Cruces, NM 88003-8001, USA

^cPhysics and Astronomy Dept., San Francisco State University,
San Francisco, CA 94132, USA

Agenda

Motivation

Lattice QCD and the Sign problem

The Polyakov Line Action

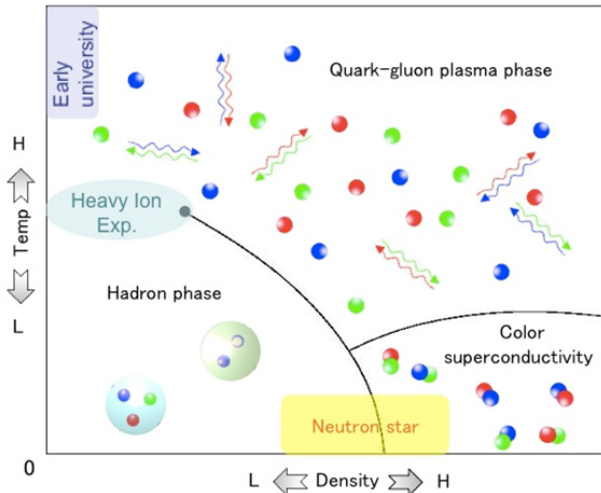
Preliminary Results

Conclusions & Outlook

Questions?

Motivation

The Phase Diagram of QCD



Lattice QCD and the Sign problem

- $Z = \int DUD\bar{\psi}D\psi e^{-S_{\text{YM}}(U) - S_{\text{F}}(U;\mu)}$

Lattice QCD and the Sign problem

- $Z = \int D U D \bar{\psi} D \psi e^{-S_{\text{YM}}(U) - S_{\text{F}}(U; \mu)}$
- $S_{\text{F}}(U; \mu) = - \int d^4 x \bar{\psi} M(U; \mu) \psi$

Lattice QCD and the Sign problem

- $Z = \int DU D\bar{\psi} D\psi e^{-S_{\text{YM}}(U) - S_{\text{F}}(U; \mu)}$
- $S_{\text{F}}(U; \mu) = - \int d^4x \bar{\psi} M(U; \mu) \psi$
- $Z = \int DU e^{-S_{\text{YM}}(U)} \det M(U; \mu)$

Lattice QCD and the Sign problem

- $Z = \int DU D\bar{\psi} D\psi e^{-S_{\text{YM}}(U) - S_{\text{F}}(U; \mu)}$
- $S_{\text{F}}(U; \mu) = - \int d^4x \bar{\psi} M(U; \mu) \psi$
- $Z = \int DU e^{-S_{\text{YM}}(U)} \det M(U; \mu)$
- numerical evaluation of bosonic integral with importance sampling

Lattice QCD and the Sign problem

- $Z = \int DUD\bar{\psi}D\psi e^{-S_{\text{YM}}(U)-S_{\text{F}}(U;\mu)}$
- $S_{\text{F}}(U; \mu) = - \int d^4x \bar{\psi} M(U; \mu) \psi$
- $Z = \int DU e^{-S_{\text{YM}}(U)} \det M(U; \mu)$
- numerical evaluation of bosonic integral with importance sampling
- observable $\langle O \rangle = \frac{\int DU e^{-S_{\text{YM}}} \det M O}{\int DU e^{-S_{\text{YM}}} \det M}$

Lattice QCD and the Sign problem

- $Z = \int DUD\bar{\psi}D\psi e^{-S_{\text{YM}}(U)-S_{\text{F}}(U;\mu)}$
- $S_{\text{F}}(U; \mu) = - \int d^4x \bar{\psi} M(U; \mu) \psi$
- $Z = \int DU e^{-S_{\text{YM}}(U)} \det M(U; \mu)$
- numerical evaluation of bosonic integral with importance sampling
- observable $\langle O \rangle = \frac{\int DU e^{-S_{\text{YM}}} \det M O}{\int DU e^{-S_{\text{YM}}} \det M}$
- lack of γ_5 -hermiticity, $\gamma_5 M(\mu) \gamma_5 = M^\dagger(-\mu^*) \neq M^\dagger(\mu)$

Lattice QCD and the Sign problem

- $Z = \int DU D\bar{\psi} D\psi e^{-S_{\text{YM}}(U) - S_{\text{F}}(U; \mu)}$
- $S_{\text{F}}(U; \mu) = - \int d^4x \bar{\psi} M(U; \mu) \psi$
- $Z = \int DU e^{-S_{\text{YM}}(U)} \det M(U; \mu)$
- numerical evaluation of bosonic integral with importance sampling
- observable $\langle O \rangle = \frac{\int DU e^{-S_{\text{YM}}} \det M O}{\int DU e^{-S_{\text{YM}}} \det M}$
- lack of γ_5 -hermiticity, $\gamma_5 M(\mu) \gamma_5 = M^\dagger(-\mu^*) \neq M^\dagger(\mu)$

determinant is complex and satisfies

$$[\det M(\mu)]^* = \det M(-\mu^*)$$

Importance of the Sign problem

- asymmetry between matter and anti-matter

Importance of the Sign problem

- asymmetry between matter and anti-matter
- free energy of particle q / anti-particle \bar{q}

Importance of the Sign problem

- asymmetry between matter and anti-matter
- free energy of particle q / anti-particle \bar{q}
- expectation value of Polyakov loop / adjoint:

$$\begin{aligned}\exp\left(-\frac{1}{T}F_q\right) &= \langle \text{Tr } P \rangle \\ &= \int \text{Re}(P) \times \text{Re}(d\varpi) - \text{Im}(P) \times \text{Im}(d\varpi)\end{aligned}$$

$$\begin{aligned}\exp\left(-\frac{1}{T}F_{\bar{q}}\right) &= \langle \text{Tr } P^* \rangle \\ &= \int \text{Re}(P) \times \text{Re}(d\varpi) + \text{Im}(P) \times \text{Im}(d\varpi)\end{aligned}$$

Importance of the Sign problem

- asymmetry between matter and anti-matter
- free energy of particle q / anti-particle \bar{q}
- expectation value of Polyakov loop / adjoint:

$$\begin{aligned}\exp\left(-\frac{1}{T}F_q\right) &= \langle \text{Tr } P \rangle \\ &= \int \text{Re}(P) \times \text{Re}(d\varpi) - \text{Im}(P) \times \text{Im}(d\varpi)\end{aligned}$$

$$\begin{aligned}\exp\left(-\frac{1}{T}F_{\bar{q}}\right) &= \langle \text{Tr } P^* \rangle \\ &= \int \text{Re}(P) \times \text{Re}(d\varpi) + \text{Im}(P) \times \text{Im}(d\varpi)\end{aligned}$$

- finite chemical potential μ favors propagation of quarks

Possible Solutions of the Sign problem

- **Reweighting:**

measurements of O are given a varying, oscillatory weight f/g in the ensemble average (“average sign”)

Possible Solutions of the Sign problem

- **Reweighting:**

measurements of O are given a varying, oscillatory weight f/g in the ensemble average (“average sign”)

- **Taylor expansion:**

of the observable in powers of μ/T at $\mu = 0$

Possible Solutions of the Sign problem

- **Reweighting:**
measurements of O are given a varying, oscillatory weight f/g in the ensemble average (“average sign”)
- **Taylor expansion:**
of the observable in powers of μ/T at $\mu = 0$
- **Imaginary μ :** analytic continuation of results to real μ

Possible Solutions of the Sign problem

- **Reweighting:**

measurements of O are given a varying, oscillatory weight f/g in the ensemble average (“average sign”)

- **Taylor expansion:**

of the observable in powers of μ/T at $\mu = 0$

- **Imaginary μ :** analytic continuation of results to real μ

- **|QCD|:**

$\det M = |\det M|e^{i\phi}$, simulations without $e^{i\phi}$ + reweighting

Possible Solutions of the Sign problem

- **Reweighting:**

measurements of O are given a varying, oscillatory weight f/g in the ensemble average (“average sign”)

- **Taylor expansion:**

of the observable in powers of μ/T at $\mu = 0$

- **Imaginary μ :** analytic continuation of results to real μ

- **|QCD|:**

$\det M = |\det M| e^{i\phi}$, simulations without $e^{i\phi}$ + reweighting

- **Complex Langevin:** stochastic quantization - evolution of fields in a fictitious time with Brownian noise and search for stationary solutions with correct measure

Possible Solutions of the Sign problem

- **Reweighting:**

measurements of O are given a varying, oscillatory weight f/g in the ensemble average (“average sign”)

- **Taylor expansion:**

of the observable in powers of μ/T at $\mu = 0$

- **Imaginary μ :** analytic continuation of results to real μ

- **|QCD|:**

$\det M = |\det M| e^{i\phi}$, simulations without $e^{i\phi}$ + reweighting

- **Complex Langevin:** stochastic quantization - evolution of fields in a fictitious time with Brownian noise and search for stationary solutions with correct measure

- **Worldline formalism and strong coupling limit:**

change order of integration, partial integration over loops and hopping parameter expansion

Effective Polyakov Line Action

- Indirect approach: Polyakov line action ($SU(3)$ spin) model

Effective Polyakov Line Action

- Indirect approach: Polyakov line action ($SU(3)$ spin) model
- fix Polyakov line holonomies $U_0(\vec{x}, 0) = U_x$ (temporal gauge) and integrate out all other d.o.f.

Effective Polyakov Line Action

- Indirect approach: Polyakov line action (SU(3) spin) model
- fix Polyakov line holonomies $U_0(\vec{x}, 0) = U_x$ (temporal gauge) and integrate out all other d.o.f.

$$e^{S_P(U_x)} = \int DU_0(\vec{x}, 0) DU_k D\psi \prod_x \delta[U_x - U_0(\vec{x}, 0)] e^{S_L}$$

Effective Polyakov Line Action

- Indirect approach: Polyakov line action (SU(3) spin) model
- fix Polyakov line holonomies $U_0(\vec{x}, 0) = U_x$ (temporal gauge) and integrate out all other d.o.f.

$$e^{S_P(U_x)} = \int DU_0(\vec{x}, 0) DU_k D\psi \prod_x \delta[U_x - U_0(\vec{x}, 0)] e^{S_L}$$

- derive S_P at $\mu = 0$, for $\mu > 0$ we have (true to all orders of strong coupling/hopping parameter expansion)

Effective Polyakov Line Action

- Indirect approach: Polyakov line action (SU(3) spin) model
- fix Polyakov line holonomies $U_0(\vec{x}, 0) = U_x$ (temporal gauge) and integrate out all other d.o.f.

$$e^{S_P(U_x)} = \int DU_0(\vec{x}, 0) DU_k D\psi \prod_x \delta[U_x - U_0(\vec{x}, 0)] e^{S_L}$$

- derive S_P at $\mu = 0$, for $\mu > 0$ we have (true to all orders of strong coupling/hopping parameter expansion)

$$S_P^\mu(U_x, U_x^\dagger) = S_P^{\mu=0}[e^{N_t\mu} U_x, e^{-N_t\mu} U_x^\dagger]$$

Effective Polyakov Line Action

- Indirect approach: Polyakov line action (SU(3) spin) model
- fix Polyakov line holonomies $U_0(\vec{x}, 0) = U_x$ (temporal gauge) and integrate out all other d.o.f.

$$e^{S_P(U_x)} = \int DU_0(\vec{x}, 0) DU_k D\psi \prod_x \delta[U_x - U_0(\vec{x}, 0)] e^{S_L}$$

- derive S_P at $\mu = 0$, for $\mu > 0$ we have (true to all orders of strong coupling/hopping parameter expansion)

$$S_P^\mu(U_x, U_x^\dagger) = S_P^{\mu=0}[e^{N_t\mu} U_x, e^{-N_t\mu} U_x^\dagger]$$

- hard to compute $\exp[S_P(U_x)]$, use relative weights...

Relative Weights Method

- S'_L ... lattice action in temporal gauge with $U_0(\vec{x}, 0) = U'_x$, compute the ratio

$$\begin{aligned} e^{\Delta S_P} &= \frac{\exp[S_P(U'_x)]}{\exp[S_P(U''_x)]} = \frac{\int DU_k D\psi e^{S'_L}}{\int DU_k D\psi e^{S''_L}} \\ &= \frac{\int DU_k D\psi \exp[S'_L - S''_L] e^{S''_L}}{\int DU_k D\psi e^{S''_L}} \equiv \langle \exp[S'_L - S''_L] \rangle'' \end{aligned}$$

Relative Weights Method

- S'_L ... lattice action in temporal gauge with $U_0(\vec{x}, 0) = U'_x$, compute the ratio

$$\begin{aligned} e^{\Delta S_P} &= \frac{\exp[S_P(U'_x)]}{\exp[S_P(U''_x)]} = \frac{\int DU_k D\psi e^{S'_L}}{\int DU_k D\psi e^{S''_L}} \\ &= \frac{\int DU_k D\psi \exp[S'_L - S''_L] e^{S''_L}}{\int DU_k D\psi e^{S''_L}} \equiv \langle \exp[S'_L - S''_L] \rangle'' \end{aligned}$$

- $U_x(\lambda)$ path through configuration space parametrized by λ

Relative Weights Method

- S'_L ... lattice action in temporal gauge with $U_0(\vec{x}, 0) = U'_x$, compute the ratio

$$\begin{aligned}
 e^{\Delta S_P} &= \frac{\exp[S_P(U'_x)]}{\exp[S_P(U''_x)]} = \frac{\int DU_k D\psi e^{S'_L}}{\int DU_k D\psi e^{S''_L}} \\
 &= \frac{\int DU_k D\psi \exp[S'_L - S''_L] e^{S''_L}}{\int DU_k D\psi e^{S''_L}} \equiv \langle \exp[S'_L - S''_L] \rangle''
 \end{aligned}$$

- $U_x(\lambda)$ path through configuration space parametrized by λ

$$U'_x = U_x(\lambda_0 + \Delta\lambda/2), U''_x = U_x(\lambda_0 - \Delta\lambda/2) \rightarrow \left(\frac{dS_P}{d\lambda}\right)_{\lambda_0} = \frac{\Delta S}{\Delta\lambda}$$

- derivatives of S_P w.r.t. Fourier components a_k of

- derivatives of S_P w.r.t. Fourier components a_k of

$$P_x \equiv \frac{1}{3} \text{Tr} U_x = \sum_k a_k e^{ikx}$$

- derivatives of S_P w.r.t. Fourier components a_k of

$$P_x \equiv \frac{1}{3} \text{Tr} U_x = \sum_k a_k e^{ikx}$$

- effective Polyakov line action motivated by heavy-dense action, where h is some inverse power of hopping parameter and satisfies the Pauli exclusion principle as $\mu \rightarrow \infty$ - no more than three (staggered) quarks per site

- derivatives of S_P w.r.t. Fourier components a_k of

$$P_x \equiv \frac{1}{3} \text{Tr} U_x = \sum_k a_k e^{ikx}$$

- effective Polyakov line action motivated by heavy-dense action, where h is some inverse power of hopping parameter and satisfies the Pauli exclusion principle as $\mu \rightarrow \infty$ - no more than three (staggered) quarks per site

$$S_{\text{eff}}[U_x] = \sum_{x,y} P_x K(x-y) P_y \\ + p \sum_x \log(1 + h e^{\mu/T} \text{Tr}[U_x] + h^2 e^{2\mu/T} \text{Tr}[U_x^\dagger] + h^3 e^{3\mu/T}) \\ \log(1 + h e^{-\mu/T} \text{Tr}[U_x] + h^2 e^{-2\mu/T} \text{Tr}[U_x^\dagger] + h^3 e^{-3\mu/T})$$

- derivatives of S_P w.r.t. Fourier components a_k of

$$P_x \equiv \frac{1}{3} \text{Tr} U_x = \sum_k a_k e^{ikx}$$

- effective Polyakov line action motivated by heavy-dense action, where h is some inverse power of hopping parameter and satisfies the Pauli exclusion principle as $\mu \rightarrow \infty$ - no more than three (staggered) quarks per site

$$S_{\text{eff}}[U_x] = \sum_{x,y} P_x K(x-y) P_y \\ + p \sum_x \log(1 + h e^{\mu/T} \text{Tr}[U_x] + h^2 e^{2\mu/T} \text{Tr}[U_x^\dagger] + h^3 e^{3\mu/T}) \\ \log(1 + h e^{-\mu/T} \text{Tr}[U_x] + h^2 e^{-2\mu/T} \text{Tr}[U_x^\dagger] + h^3 e^{-3\mu/T})$$

- determine $K(x-y)$ and h from fitting to lattice data

- derivatives of S_P w.r.t. Fourier components a_k of

$$P_x \equiv \frac{1}{3} \text{Tr} U_x = \sum_k a_k e^{ikx}$$

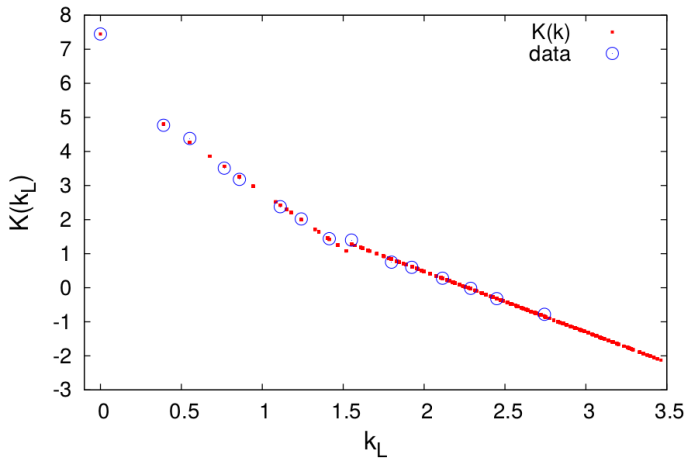
- effective Polyakov line action motivated by heavy-dense action, where h is some inverse power of hopping parameter and satisfies the Pauli exclusion principle as $\mu \rightarrow \infty$ - no more than three (staggered) quarks per site

$$S_{\text{eff}}[U_x] = \sum_{x,y} P_x K(x-y) P_y \\ + p \sum_x \log(1 + h e^{\mu/T} \text{Tr}[U_x] + h^2 e^{2\mu/T} \text{Tr}[U_x^\dagger] + h^3 e^{3\mu/T}) \\ \log(1 + h e^{-\mu/T} \text{Tr}[U_x] + h^2 e^{-2\mu/T} \text{Tr}[U_x^\dagger] + h^3 e^{-3\mu/T})$$

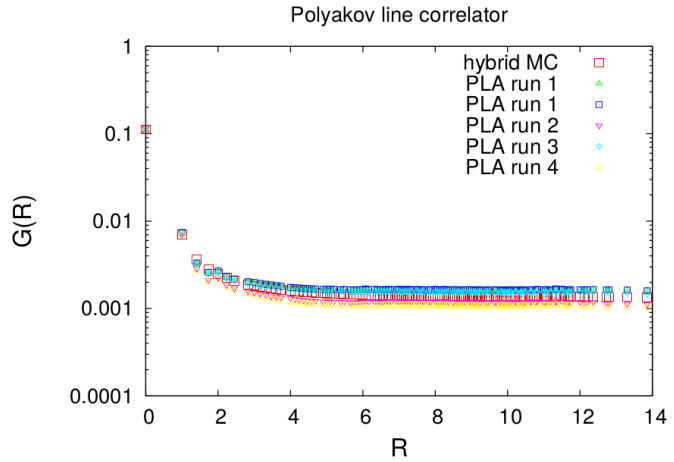
- determine $K(x-y)$ and h from fitting to lattice data

$$\frac{1}{L^3} \left(\frac{\partial S_P}{\partial a_k} \right)_{a_k=\alpha} = 2K(k)\alpha + \frac{p}{L^3} \sum_x (3h e^{ikx} + 3h^2 e^{-ikx} + \text{c.c.})$$

Preliminary Results



Preliminary Results



Solve sign problem for the effective action

- remaining sign problem can be solved by mean field theory

Solve sign problem for the effective action

- remaining sign problem can be solved by mean field theory
- treatment of $SU(3)$ spin models at finite μ is a minor variation of standard mean field theory at zero chemical potential

Solve sign problem for the effective action

- remaining sign problem can be solved by mean field theory
- treatment of $SU(3)$ spin models at finite μ is a minor variation of standard mean field theory at zero chemical potential
- two magnetizations introduced for $\text{Tr}U$ and $\text{Tr}U^\dagger$ determined by minimizing the free energy

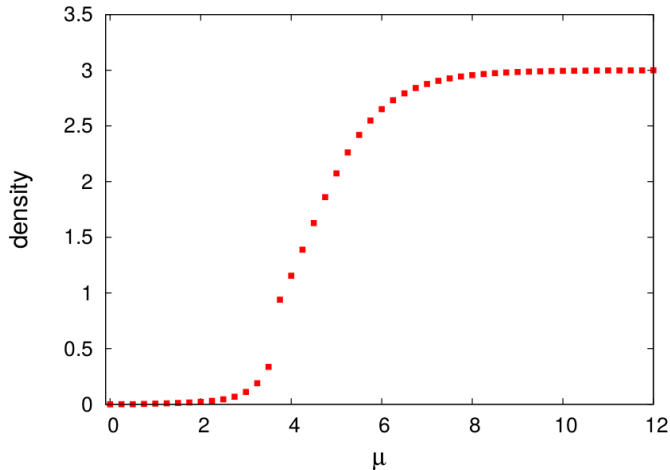
Solve sign problem for the effective action

- remaining sign problem can be solved by mean field theory
- treatment of $SU(3)$ spin models at finite μ is a minor variation of standard mean field theory at zero chemical potential
- two magnetizations introduced for $\text{Tr}U$ and $\text{Tr}U^\dagger$ determined by minimizing the free energy
- basic idea is that each spin is effectively coupled to the average spin on the lattice, not just nearest neighbors, through non-local kernel $K(x - y)$

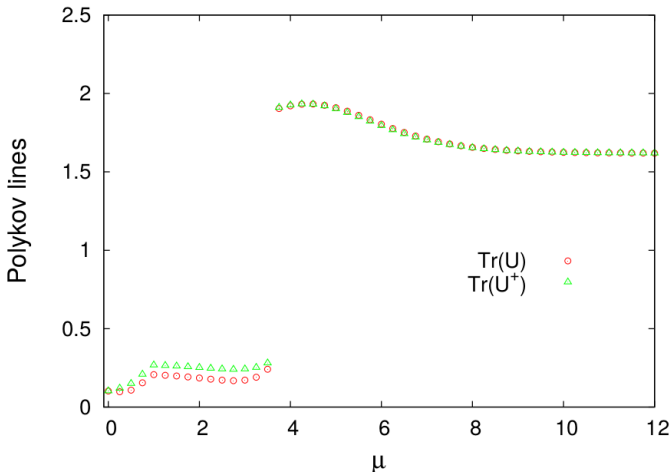
Solve sign problem for the effective action

- remaining sign problem can be solved by mean field theory
- treatment of $SU(3)$ spin models at finite μ is a minor variation of standard mean field theory at zero chemical potential
- two magnetizations introduced for $\text{Tr}U$ and $\text{Tr}U^\dagger$ determined by minimizing the free energy
- basic idea is that each spin is effectively coupled to the average spin on the lattice, not just nearest neighbors, through non-local kernel $K(x - y)$
- for details and comparison to complex Langevin see *Splitdorff and Greensite (2012)*

Preliminary Results



Preliminary Results



Conclusions & Outlook

- determined effective Polyakov line action for asqtad staggered fermions with $ma = 0.3$ and Symanzik one loop improved gauge action at $\beta = 7.0$ on $16^3 \times 6$ lattices

Conclusions & Outlook

- determined effective Polyakov line action for asqtad staggered fermions with $ma = 0.3$ and Symanzik one loop improved gauge action at $\beta = 7.0$ on $16^3 \times 6$ lattices
- good agreement for the Polyakov line correlators computed in the effective theory and underlying lattice gauge theory

Conclusions & Outlook

- determined effective Polyakov line action for asqtad staggered fermions with $ma = 0.3$ and Symanzik one loop improved gauge action at $\beta = 7.0$ on $16^3 \times 6$ lattices
- good agreement for the Polyakov line correlators computed in the effective theory and underlying lattice gauge theory
- solved sign problem for the effective theory by mean field and find a phase transition and correct density limit

Conclusions & Outlook

- determined effective Polyakov line action for asqtad staggered fermions with $ma = 0.3$ and Symanzik one loop improved gauge action at $\beta = 7.0$ on $16^3 \times 6$ lattices
- good agreement for the Polyakov line correlators computed in the effective theory and underlying lattice gauge theory
- solved sign problem for the effective theory by mean field and find a phase transition and correct density limit
- ...

Conclusions & Outlook

- determined effective Polyakov line action for asqtad staggered fermions with $ma = 0.3$ and Symanzik one loop improved gauge action at $\beta = 7.0$ on $16^3 \times 6$ lattices
- good agreement for the Polyakov line correlators computed in the effective theory and underlying lattice gauge theory
- solved sign problem for the effective theory by mean field and find a phase transition and correct density limit
- ...
- determine quadratic, quasi-local center symmetry breaking terms which may be important at finite chemical potential...

Conclusions & Outlook

- determined effective Polyakov line action for asqtad staggered fermions with $ma = 0.3$ and Symanzik one loop improved gauge action at $\beta = 7.0$ on $16^3 \times 6$ lattices
- good agreement for the Polyakov line correlators computed in the effective theory and underlying lattice gauge theory
- solved sign problem for the effective theory by mean field and find a phase transition and correct density limit
- ...
- determine quadratic, quasi-local center symmetry breaking terms which may be important at finite chemical potential...
- go on to smaller quark masses...

Questions?

Thank You &

Derar Altarawneh, Michael Engelhardt, Manfred Faber, Martin Gal,
 Jeff Greensite, Urs M. Heller, James Hettrick, Andrei Ivanov, Thomas
 Layer, Štefan Olejnik, Luis Oxman, Mario Pitschmann, Jesus Saenz,
 Thomas Schweigler, Wolfgang Söldner, David Vercauteren, Markus
 Wellenzohn





TECHNISCHE
UNIVERSITÄT
WIEN



Polyakov line actions from $SU(3)$ lattice gauge theory with dynamical fermions: first results via relative weights

Roman Höllwieser^{ab}, hroman@kph.tuwien.ac.at
Jeff Greensite^c, greensit@sfsu.edu

^aInstitute of Atomic and Subatomic Physics, Nuclear Physics Dept.,
Vienna University of Technology, Operngasse 9, 1040 Vienna, Austria

^bDepartment of Physics, New Mexico State University,
Las Cruces, NM 88003-8001, USA

^cPhysics and Astronomy Dept., San Francisco State University,
San Francisco, CA 94132, USA