ELECTROMAGNETIC STRUCTURE OF CHARMED BARYONS — EXTENDED TO SPIN-3/2

<u>K. Utku Can</u>¹, G. Erkol², M. Oka^{1,3}, T. T. Takahashi⁴

¹ Tokyo Institute of Technology, Tokyo, Japan
 ² Ozyegin University, Istanbul, Turkey
 ³ Advanced Science Research Center, Japan Atomic Energy Agency, Ibaraki, Japan
 ⁴ Gunma National Collage of Technology, Gunma, Japan

LATTICE 2015 — July 13-18, Kobe, JAPAN

OUTLINE

- Motivation
- Electromagnetic (EM) form factors
 - Parameterisation
 - Lattice Formulation
- Simulation Details
- Results
- Summary

MOTIVATION

- Probe the hadron structure
 - size, charge radius, magnetic moment
 - Effect of heavy quarks

kuc, G. Erkol, M. Oka, A. Ozpineci, T.T. Takahashi PLB 719 Previous works: heavy quark shrinks the mesons & baryons
 kuc, G. Erkol, B. Isildak, M. Oka, T.T. Takahashi JHEP05(2014) 125

• Effect of spin structure?





SIMULATION DETAILS

PACS-CS 2+1 flavor Clover Phys. Rev. D79 (034503)

$N_s^3 \times N_t$	κ_{ud}	$m_{\pi} [{ m MeV}]$	κ_s	a, a^{-1}	$L \ [\mathrm{fm}]$	$t_s = 12a$	# confs	# meas.
$32^3 \times 64$	0.13781	156	0.1364	$0.0907~{\rm fm}$	2.9	$1.09~\mathrm{fm}$	200	$200 \; ({\rm spin-1/2})$
				$2.176 \mathrm{GeV}$				$300 \; (\mathrm{spin-}3/2)$

Clover action for	all valance quarks
$\underline{\kappa_{\rm s}}=0.13640$	$m_{\Omega^{-}} = 1.790(17) \mathrm{GeV}$
PRD79 034503 (2009)	$m_{\Omega^-} = 1.772(7) \mathrm{GeV}$
PDG (2014)	$m_{\Omega^{-}}^{exp} = 1.673(29){ m GeV}$

 $\frac{\kappa_{\rm c} = 0.1246}{\rm tuned \ to \ 1S \ M_{\eta\text{-}J/\psi} \ , \ M_{D\text{-}D^*} \ , \ M_{D_{\rm s}\text{-}D_{\rm s}^*}}$

This work $m_{\Omega_c^*} - m_{\Omega_c} = 54 \pm 17 \,\text{MeV}$ Belle Collab. PLB 672 (2009) $m_{\Omega_c^*} - m_{\Omega_c} = 70.7 \pm 0.9^{+0.1}_{-0.9} \,\text{MeV}$

- 1. Lowest momenta, $|\mathbf{q}|=2\pi/L$
- 2. Point-split (conserved) vector current: renormalisation not necessary
- 3. Shell source Wall sink pairs
 - I. Wall sinks: no need for sequential inversions
- 4. Connected diagrams only
- 5. $\langle \mathcal{O} \rangle = N_s e_s \langle \mathcal{O}_s \rangle + N_c e_c \langle \mathcal{O}_c \rangle$
- 6. Single-elimination Jackknife analysis







- Strange quark charge radii insensitive to baryon composition
- Charm has a tendency to increase with the c-quark number





• Ω_c and Ω_c^* have similar charge radius



• Ω_{cc} and Ω^*_{cc} charge radii differ due to non-trivial cancellations between quark sectors



Assume E0 and M1 scales similarly for low Q^2

Estimate Extract $G_{M1}^{s,c}(0) = G_{M1}^{s,c}(Q_{min}^2) \frac{G_E^{s,c}(0)}{G_F^{s,c}(Q_{min}^2)}$

 m_N : physical mass of nucleon m_B : lattice mass of the baryon

 $=G_{M1}(0)\frac{m_N}{m}\mu_N$ μ_B



 Quark sector contributions are similar amongst spin-3/2 (expected from a quark model approach)





• Contributions within spin-3/2 systems are larger

Ω_c and Ω^{*}_c have similar magnetic moments in magnitude
 Ω^{*}_{cc} has a vanishing mag. moment unlike Ω_{cc}

• Ω^*_{cc} and Ω_{ccc} have oblate electric charge distributions

• Charge radii:

- $\langle r_E^2 \rangle_s$ is insensitive to the composition of the baryon
- $\langle r_E^2 \rangle_c$ slightly increases with the increasing number of c-quarks
- Magnetic moments:
 - $\label{eq:main_s} \quad \mu_s^{3/2} > \mu_s^{1/2}, \quad \mu_c^{3/2} > \mu_c^{1/2}$
 - \bullet u -> c , s-quark contribution is insensitive,
 - $\bullet \quad uu \rightarrow cc \ , effect \ of \ uu \ component \ is \ bigger$
 - Ω_c and Ω_c^* have similar mag. moments in magnitude
 - Ω^*_{cc} has a vanishing mag. moment unlike Ω_{cc}
- E2 moments:
 - Ω^{*+}_{cc} and Ω^{++}_{ccc} have oblate electric charge distributions

ACKNOWLEDGEMENTS

- Part of numerical calculations in this work were performed on National Center for High Performance Computing of Turkey (Istanbul Technical University) under project number 10462009.
- This work is supported in part by The Scientific and Technological Research Council of Turkey (TUBITAK) under project number 114F261 and in part by KAKENHI under Contract Nos. 25247036 and 24250294.
- This work is also supported by the Research Abroad and Invitational Program for the Promotion of International Joint Research, Category (C) and the International Physics Leadership Program at Tokyo Tech.

THANK YOU

BACKUP SLIDES

Lattice Formulation

$$\langle C^{\mathcal{B}}(t;\mathbf{p};\Gamma_{4})\rangle = \sum_{\mathbf{x}} e^{-i\mathbf{p}\cdot\mathbf{x}}\Gamma_{4}^{\alpha\alpha'} \langle \operatorname{vac}|T[\eta^{\alpha}_{\mathcal{B}}(x)\bar{\eta}^{\alpha'}_{\mathcal{B}}(0)]|\operatorname{vac}\rangle$$
$$\langle C^{\mathcal{B}\mathcal{V}_{\mu}\mathcal{B}'}(t_{2},t_{1};\mathbf{p}',\mathbf{p};\Gamma)\rangle = -i\sum_{\mathbf{x}_{2},\mathbf{x}_{1}} e^{-i\mathbf{p}\cdot\mathbf{x}_{2}} e^{i\mathbf{q}\cdot\mathbf{x}_{1}}\Gamma^{\alpha\alpha'} \langle \operatorname{vac}|T[\eta^{\alpha}_{\mathcal{B}}(x_{2})V_{\mu}(x_{1})\bar{\eta}^{\alpha'}_{\mathcal{B}'}(0)]|\operatorname{vac}\rangle$$

$$R(t_{2}, t_{1}; \mathbf{p}', \mathbf{p}; \mathbf{\Gamma}; \mu) = \frac{\langle C^{\mathcal{B}\mathcal{V}_{\mu}\mathcal{B}'}(t_{2}, t_{1}; \mathbf{p}', \mathbf{p}; \mathbf{\Gamma}) \rangle}{\langle C^{\mathcal{B}\mathcal{B}}(t_{2}; \mathbf{p}'; \mathbf{\Gamma}_{4}) \rangle} \\ \times \left[\frac{\langle C^{\mathcal{B}\mathcal{B}}(t_{2} - t_{1}; \mathbf{p}; \mathbf{\Gamma}_{4}) \rangle \langle C^{\mathcal{B}\mathcal{B}}(t_{1}; \mathbf{p}'; \mathbf{\Gamma}_{4}) \rangle \langle C^{\mathcal{B}\mathcal{B}}(t_{2}; \mathbf{p}'; \mathbf{\Gamma}_{4}) \rangle}{\langle C^{\mathcal{B}\mathcal{B}}(t_{2} - t_{1}; \mathbf{p}'; \mathbf{\Gamma}_{4}) \rangle \langle C^{\mathcal{B}\mathcal{B}}(t_{1}; \mathbf{p}; \mathbf{\Gamma}_{4}) \rangle \langle C^{\mathcal{B}\mathcal{B}}(t_{2}; \mathbf{p}; \mathbf{\Gamma}_{4}) \rangle} \right]^{1/2}$$

$$\eta_{\Xi_{cc}}(x) = \epsilon^{ijk} [c^{Ti}(x) C \gamma_5 \ell^j(x)] c^k(x)$$

$$\eta_{\Sigma_c}(x) = \epsilon^{ijk} [\ell^{Ti}(x) C \gamma_5 c^j(x)] \ell^k(x)$$

$$\eta_{\Omega_{cc}}(x) = \epsilon^{ijk} [c^{Ti}(x) C \gamma_5 s^j(x)] c^k(x)$$

$$\eta_{\Omega_c}(x) = \epsilon^{ijk} [s^{Ti}(x) C \gamma_5 c^j(x)] s^k(x)$$

 $\mathbf{\Sigma}$

$$R(t_{2}, t_{1}; \mathbf{p}', \mathbf{p}; \Gamma; \mu) \xrightarrow{t_{1} \gg a} \Pi(\mathbf{p}', \mathbf{p}; \Gamma; \mu)$$

$$\Pi(\mathbf{0}, -\mathbf{q}; \Gamma_{4}; \mu = 4) = \left[\frac{(E_{\mathcal{B}} + m_{\mathcal{B}})}{2E_{\mathcal{B}}}\right]^{1/2} G_{E,\mathcal{B}}(q^{2})$$

$$\Pi(\mathbf{0}, -\mathbf{q}; \Gamma_{j}; \mu = i) = \left[\frac{1}{2E_{\mathcal{B}}(E_{\mathcal{B}} + m_{\mathcal{B}})}\right]^{1/2} \epsilon_{ijk} q_{k} G_{M,\mathcal{B}}(q^{2})$$

$$G(0) \text{ should be extrapolated from higher momenta}}$$

Γ

spin-3/2

$$\langle \mathcal{B}_{\sigma}(p',s')|j_{\mu}|\mathcal{B}_{\tau}(p,s)\rangle = \sqrt{\frac{M_{\mathcal{B}}^{2}}{E E'}} \bar{u}_{\sigma}(p',s')\mathcal{O}^{\sigma\mu\tau}u_{\tau}(p,s)$$

$$\langle \mathcal{O}^{\sigma\mu\tau} = -g^{\sigma\tau} \left\{ a_{1}\gamma^{\mu} + \frac{a_{2}}{2M_{\mathcal{B}}}P^{\mu} \right\} - \frac{q^{\sigma}q^{\tau}}{(2M_{\mathcal{B}})^{2}} \left\{ c_{1}\gamma^{\mu} + \frac{c_{2}}{2M_{\mathcal{B}}}P^{\mu} \right\} \quad \stackrel{\mathsf{P} = \mathsf{p'} + \mathsf{p}}{\mathsf{q} = \mathsf{p'} - \mathsf{p}}$$

as drived in S. Nozawa and D. Leinweber, Phys.Rev. D42, 3567(1990)

$$\begin{array}{l} \mbox{Electric charge}\\ G_{E0}(q^2) = (1 + \frac{2}{3}\tau) \left\{ a_1 + (1 + \tau)a_2 \right\} \\ & -\frac{1}{3}\tau(1 + \tau) \left\{ c_1 + (1 + \tau)c_2 \right\}, \end{array} \begin{array}{l} \mbox{Magnetic dipole}\\ G_{M1}(q^2) = (1 + \frac{4}{3}\tau)a_1 - \frac{2}{3}\tau(1 + \tau)c_1, \\ G_{E2}(q^2) = \left\{ a_1 + (1 + \tau)a_2 \right\} \\ \mbox{Electric}\\ \mbox{quadrupole} \end{array} - \frac{1}{2}(1 + \tau) \left\{ c_1 + (1 + \tau)c_2 \right\}, \end{array} \right. \begin{array}{l} \mbox{Magnetic dipole}\\ G_{M1}(q^2) = (1 + \frac{4}{3}\tau)a_1 - \frac{2}{3}\tau(1 + \tau)c_1, \\ \mbox{G}_{E1}(q^2) = \left\{ a_1 + (1 + \tau)a_2 \right\} \\ \mbox{Electric}\\ \mbox{quadrupole} \end{array} - \frac{1}{2}(1 + \tau) \left\{ c_1 + (1 + \tau)c_2 \right\}, \end{array} \right.$$

spin-3/2

$$\begin{array}{l} (1) \\ \langle G^{\mathcal{BB}}_{\sigma\tau}(t;\mathbf{p};\Gamma_{4})\rangle = \sum_{\mathbf{x}} e^{-i\mathbf{p}\cdot\mathbf{x}}\Gamma_{4}^{\alpha\alpha'} \\ \times \langle \operatorname{vac}|T[\eta^{\alpha}_{\sigma}(x)\bar{\eta}^{\alpha'}_{\tau}(0)]|\operatorname{vac}\rangle \\ \langle G^{\mathcal{B}j^{\mu}\mathcal{B}}_{\sigma\tau}(t_{2},t_{1};\mathbf{p}',\mathbf{p};\Gamma)\rangle = -i\sum_{\mathbf{x}_{2},\mathbf{x}_{1}} e^{-i\mathbf{p}\cdot\mathbf{x}_{2}}e^{i\mathbf{q}\cdot\mathbf{x}_{1}} \\ \times \Gamma^{\alpha\alpha'}\langle \operatorname{vac}|T[\eta^{\alpha}_{\sigma}(x_{2})j_{\mu}(x_{1})\bar{\eta}^{\alpha'}_{\tau}(0)]|\operatorname{vac}\rangle \end{array} \right| \\ \begin{array}{l} (2) \\ R_{\sigma}^{\mu}{}_{\tau}(t_{2},t_{1};\mathbf{p}',\mathbf{p};\Gamma) = -i\sum_{\mathbf{x}_{2},\mathbf{x}_{1}} e^{-i\mathbf{p}\cdot\mathbf{x}_{2}}e^{i\mathbf{q}\cdot\mathbf{x}_{1}} \\ \frac{t_{1}\gg a}{t_{2}-t_{1}\gg a} \left(\frac{E_{p}+M_{\mathcal{B}}}{2E_{p}}\right)^{1/2} \left(\frac{E_{p'}+M_{\mathcal{B}}}{2E_{p'}}\right)^{1/2} \\ \times \Pi_{\sigma}^{\mu}{}_{\tau}(\mathbf{p}',\mathbf{p};\Gamma) \end{array}$$

$$\begin{split} \Im & G_{E0}(q^2) = \frac{1}{3} \left(\Pi_1^{\ 4}_1(\mathbf{q}_i, 0; \Gamma_4) + \Pi_2^{\ 4}_2(\mathbf{q}_i, 0; \Gamma_4) \right) \\ & + \Pi_3^{\ 4}_3(\mathbf{q}_i, 0; \Gamma_4) \right), \\ & G_{E2}(q^2) = 2 \frac{M(E+M)}{|\mathbf{q}_i|^2} \left(\Pi_1^{\ 4}_1(\mathbf{q}_i, 0; \Gamma_4) \right) \\ & + \Pi_2^{\ 4}_2(\mathbf{q}_i, 0; \Gamma_4) - 2\Pi_3^{\ 4}_3(\mathbf{q}_i, 0; \Gamma_4) \right), \end{split}$$

21	PRD87	094512	(2013)	
----	-------	--------	--------	--

[14] PRD79 034503 (2009) PRD90 074501 (2014)

PRD86 094504 (2012)

	J^P	This work	PACS-CS $[21]$	ETMC [13]	Briceno et al. [23]	Exp. [22]
		[GeV]	[GeV]	[GeV]	[GeV]	[GeV]
Ω_c	$\frac{1}{2}^+$	2.783(13)	2.673(17)	2.629(22)	2.681(48)	2.695(2)
Ω_{cc}	$\frac{1}{2}^+$	3.747(10)	3.704(21)	3.654(18)	3.679(62)	
Ω	$\frac{3}{2}^+$	1.790(17)	1.772(7) [14]	1.672(18)		1.673(29)
Ω_c^*	$\frac{3}{2}^+$	2.837(18)	2.738(17)	2.709(26)	2.764(49)	2.766(2)
Ω_{cc}^*	$\frac{3}{2}^+$	3.819(10)	3.779(23)	3.724(21)	3.765(65)	
Ω_{ccc}	$\frac{3}{2}^+$	4.769(6)	4.789(27)	4.733(18)	4.761(79)	

Quark charge radii increase in doubly charmed spin-3/2

RESULTS

D and D^* electric charge radius extrapolations

<u>-</u>	$\kappa^{u,d}_{val}$	$g_{DD ho}$	$\langle r_D^2 \rangle ~({\rm fm}^2)$	$\langle r_{D,d}^2 \rangle ~({\rm fm}^2)$	$\langle r_{D,c}^2 \rangle ~({\rm fm}^2)$
	0.13700	5.57(27)	0.098(8)	0.226(20)	0.040(9)
M-D	0.13727	5.03(28)	0.094(10)	0.232(29)	0.033(9)
▲ M=D*	0.13754	5.30(30)	0.124(13)	0.350(50)	0.033(12)
	0.13770	4.85(41)	0.133(19)	0.308(85)	0.059(13)
-	Lin. Fit	4.84(34)	0.138(13)	0.342(67)	0.051(11)
	Quad. Fit	4.90(56)	0.152(26)	0.320(118)	0.074(16)
	$\kappa^{u,d}_{val}$	$g_{D^*D^* ho}$	$\langle r_{D^*}^2 \rangle ~({\rm fm}^2)$	$\langle r_{D^*,d}^2 \rangle ~({\rm fm}^2)$	$\langle r_{D^*,c}^2 \rangle ~(\mathrm{fm}^2)$
····	0.13700	5.93(41)	0.106(12)	0.270(31)	0.035(13)
	0.13727	5.75(36)	0.113(17)	0.296(44)	0.036(14)
	0.13754	6.69(47)	0.167(25)	0.497(92)	0.044(21)
	0.13770	5.57(66)	0.169(30)	0.404(127)	0.075(26)
····	Lin. Fit	5.94(56)	0.185(24)	0.475(94)	0.071(16)
	Quad. Fit	5.42(94)	0.192(43)	0.406(156)	0.096(29)
	= -0.13	8 fm2 ~1	2 - 1 - 0.31	$2 \operatorname{fm}^2 < w^2 > 2$	- 0.051 fi
	$< r^2 >_D = 0.13$ $< r^2 >_D^* = 0.183$	5 fm² , <r< td=""><td>$2 >_{D,a} = 0.34$</td><td>$25 fm^2$, $< r^2 > D$</td><td>$f_{c} = 0.001 \text{Jm}$ $f_{c} = 0.071 \text{Jm}$</td></r<>	$2 >_{D,a} = 0.34$	$25 fm^2$, $< r^2 > D$	$f_{c} = 0.001 \text{Jm}$ $f_{c} = 0.071 \text{Jm}$
0.04 0.08 0.12 m_{π}^2		$< r^2$	$^{2}>_{\pi}=0.4$	$152 fm^2$	

30

 κ_{ud}

(137-)

stat.

8

7

6

5

4

3 0.3

0.2

0.1

0

0.6

0.4

0.2

0

0 0

 $m_{\pi} \sim 300 \ 410$

570

700

 $g_{MM\rho}$

 $\langle r^2_{\ M} \ \rangle \, (fm^2)$

 $\langle r^2_{M,d} \rangle (fm^2)$

70

70

54

50

27

50

00

45

samples

MeV

kuc, G. Erkol, M. Oka, A. Ozpineci, T. T. Takahashi, Phys.Lett. B719 (2013) 103-109

