

Density of states method for $SU(3)$ -spin system

Mario Giuliani

Christof Gatteringer, Pascal Törek

Karl Franzens Universität Graz



Der Wissenschaftsfonds.



Pascal Törek talk: Z_3 discrete variables \Rightarrow exact DoS parametrization
 $SU(3)$ -spin system: continuous variables \Rightarrow different DoS parametrization

- Monte-Carlo simulation of quantum field theories at finite density: Sign problem
- Different approaches to solve this problem:
 - Reweighting
 - Expansion methods
 - Stochastic differential equations
 - Mapping to dual variables
 - Et cetera...
- Density of states approach:
 - New development: Functional Fit Approach (FFA)
 - Tested in the $SU(3)$ -spin system
 - $SU(3)$ -spin system has a dual formulation [Gattringer Nuclear Physics B850 (2011) 242-252]

Density of states:

$$Z = \int \mathcal{D}[C] e^{S_R[C] + iS_I[C]} \text{ where } S_I[C] \text{ is bounded: } S_I[C] \in [S_{I_{min}}; S_{I_{max}}]$$



Use a Monte Carlo calculation to estimate the density:

$$\rho(S_I) = \int \mathcal{D}[C] e^{S_R[C]} \delta(S_I - S_I[C])$$

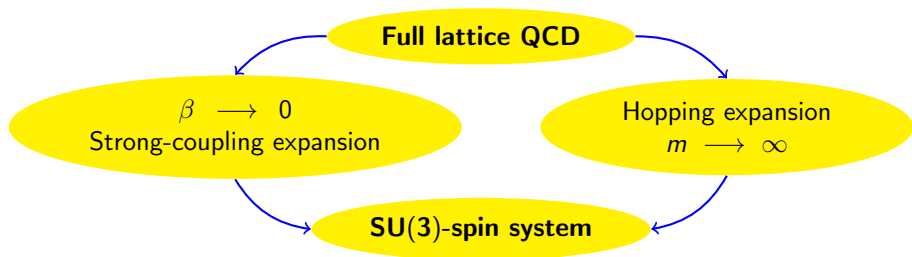


$$Z = \int_{S_{I_{min}}}^{S_{I_{max}}} d[S_I] \rho(S_I) \{\cos(S_I) + i \sin(S_I)\}$$

$$\langle \mathcal{O} \rangle = \int_{S_{I_{min}}}^{S_{I_{max}}} d[S_I] \rho(S_I) \{\cos(S_I) + i \sin(S_I)\} \mathcal{O}(S_I)$$



Sine and cosine are oscillating functions so we need to know $\rho(S_I)$ with high precision

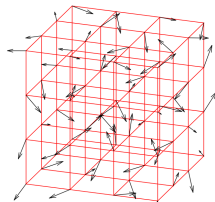


- We have a system with action:

$$S[P] = \tau \sum_x \sum_{\nu=1}^3 \left[\text{Tr}P(x) \text{Tr}P(x+\nu)^\dagger + \text{c.c.} \right] + k \sum_x \left[e^\mu \text{Tr}P(x) + e^{-\mu} \text{Tr}P(x)^\dagger \right]$$

The action depends only on the trace

$$P(x) = \begin{pmatrix} e^{i\theta_1} & 0 & 0 \\ 0 & e^{i\theta_2} & 0 \\ 0 & 0 & e^{-i(\theta_1+\theta_2)} \end{pmatrix}$$



- Term with chemical potential can be decomposed into real and imaginary parts:

$$k \sum_x [e^{\mu} \text{Tr}P(x) + e^{-\mu} \text{Tr}P(x)] = \\ 2k \cosh(\mu) \sum_x \text{Re}[\text{Tr}P(x)] + i2k \sinh(\mu) \sum_x \text{Im}[\text{Tr}P(x)]$$

- Definition:

$$- \sum_x \text{Im}[\text{Tr}P(x)] = F[P] = \sum_x [\sin(\theta_1(x)) + \sin(\theta_2(x)) - \sin(\theta_1(x) + \theta_2(x))]$$

$$- F[P] \in \left[-\frac{3\sqrt{3}}{2} V_3; \frac{3\sqrt{3}}{2} V_3\right] \equiv [-l_{max}, l_{max}]$$

- We define the weighted density of states:

$$\rho(l) = \int \mathcal{D}[P] e^{S_R[P]} \delta(l - F[P]) \quad l \in [-l_{max}, l_{max}]$$

- The density of states is an even function of l :

$$\rho(-l) = \rho(l)$$

because of symmetry properties under $P(x) \rightarrow P(x)^*$

- This simplifies the partition function:

$$Z = \int_{-l_{max}}^{l_{max}} dl \rho(l) \cos(2k \sinh(\mu)l) = 2 \int_0^{l_{max}} dl \rho(l) \cos(2k \sinh(\mu)l)$$

- We can also use $\rho(l) = \rho(-l)$ for the expectation value of the observables
- As observables we consider functions of $F[P]$, and define even and odd parts:

$$\begin{aligned}\mathcal{O}[F] &= \mathcal{O}_E[F] + \mathcal{O}_O[F] \\ \mathcal{O}_E[F] &= \frac{1}{2}(\mathcal{O}[F] + \mathcal{O}[-F]) \quad ; \quad \mathcal{O}_O[F] = \frac{1}{2}(\mathcal{O}[F] - \mathcal{O}[-F])\end{aligned}$$

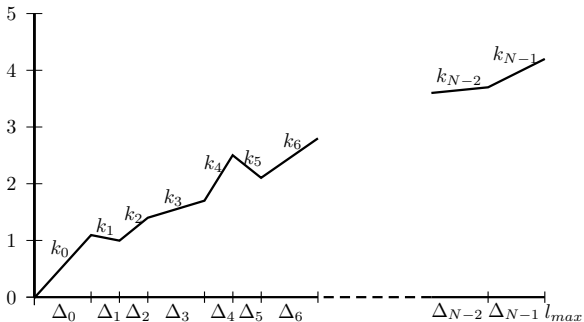
- We have for the expectation value of the observable:

$$\langle \mathcal{O}[F] \rangle = \frac{2}{Z} \int_0^{l_{\max}} dl \rho(l) \left[\mathcal{O}_E[l] \cos(2k \sinh(\mu)l) + i \mathcal{O}_O[l] \sin(2k \sinh(\mu)l) \right]$$

- We need $\rho(l)$ only for $l \geq 0$
- For large k and μ the sine and cosine are strongly oscillating functions

Parametrization of the density $\rho(l)$

- Ansatz for the density: $\rho(l) = e^{-L(l)}$, normalization $\rho(0) = 1 \Rightarrow L(0) = 0$
- We divide the interval $[0, l_{max}]$ into N intervals $n = 0, 1, \dots, N - 1$.
- $L(l)$ is continuous and linear on each of the intervals, with a slope k_n :



- How do we find the slopes k_n ? Restricted expectation values which depend on a parameter $\lambda \in \mathbb{R}$:

$$\langle\langle \mathcal{O} \rangle\rangle_n(\lambda) = \frac{1}{Z_n(\lambda)} \int \mathcal{D}[P] e^{S_R[P] + \lambda F[P]} \mathcal{O}[F[P]] \theta_n[F[P]]$$

$$Z_n(\lambda) = \int \mathcal{D}[P] e^{S_R[P] + \lambda F[P]} \theta_n[F[P]]$$

$$\theta_n[l] = \begin{cases} 1 & \text{for } l \in [l_n, l_n + 1] \\ 0 & \text{otherwise} \end{cases}$$

- Expressed in terms of the density of states ($l = F[P]$):

$$\langle\langle \mathcal{O} \rangle\rangle_n(\lambda) = \frac{1}{Z_n(\lambda)} \int_{-l_{\max}}^{l_{\max}} dl \rho(l) e^{\lambda l} \mathcal{O}[l] \theta_n[l] = \frac{1}{Z_n(\lambda)} \int_{l_n}^{l_{n+1}} dl \rho(l) e^{\lambda l} \mathcal{O}[l]$$

$$Z_n(\lambda) = \int_{-l_{\max}}^{l_{\max}} dl \rho(l) e^{\lambda l} \theta_n[l] = \int_{l_n}^{l_{n+1}} dl \rho(l) e^{\lambda l}$$

- For computing the slopes we use as observable $F[P]$:

$$\langle\langle F[P] \rangle\rangle_n(\lambda) = \frac{1}{Z_n(\lambda)} \int_{l_n}^{l_{n+1}} dl \rho(l) e^{\lambda l} l = \frac{\partial}{\partial \lambda} \ln[Z_n(\lambda)]$$

- Explicit evaluation:

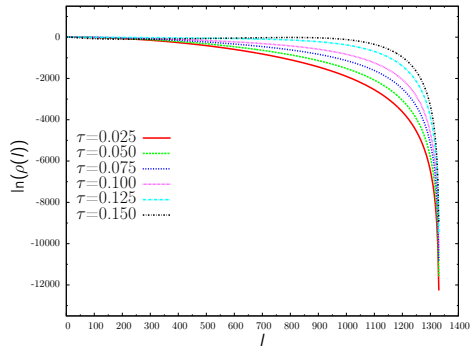
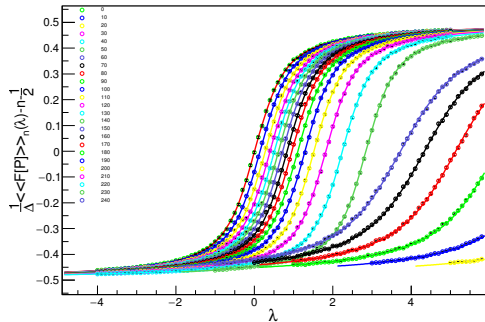
$$Z_n(\lambda) = e^{(\lambda - k_n)n\Delta} \frac{1 - e^{-\Delta(k_n - \lambda)}}{k_n - \lambda}$$
$$\langle\langle F[P] \rangle\rangle_n(\lambda) = n\Delta + \frac{\Delta}{1 - e^{\Delta(k_n - \lambda)}} + \frac{1}{k_n - \lambda}$$

- Evaluate $\langle\langle F[P] \rangle\rangle_n(\lambda)$ for different values of $\lambda \Rightarrow$ Fit these values to obtain k_n

Fit of slopes \Rightarrow density $\rho(l)$

Example \Rightarrow **Lattice:** 8^3 , $k = 0.005$, $\mu = 0.0$
 $\tau = 0.075$

$\tau = 0.025, 0.050, \dots, 0.150$



$$e^{-L(l)} = \rho(l)$$

- We use the density of states to compute observables:
- Particle number density n :

$$n = \frac{1}{iV_3} F[P] = \frac{1}{V_3} \left(-\frac{1}{2k} \right) \frac{\partial}{\partial \sinh(\mu)} \ln Z$$

- ... and the corresponding susceptibility χ_n :

$$\begin{aligned} \chi_n &= \frac{1}{2k} \frac{\partial}{\partial \sinh(\mu)} n = \frac{1}{V_3} \left(-\frac{1}{(2k)^2} \right) \frac{\partial^2}{(\partial \sinh(\mu))^2} \ln Z \\ &= \frac{1}{V_3} \left(-\frac{1}{(2k)^2} \right) \left[\frac{1}{Z} \frac{\partial^2 Z}{(\partial \sinh(\mu))^2} - \left(\frac{1}{Z} \frac{\partial Z}{\partial \sinh(\mu)} \right)^2 \right] \end{aligned}$$

- In terms of the density of states $\rho(l)$ we have:

$$Z = 2 \int_0^{l_{\max}} dl \rho(l) \cos(2k \sinh(\mu)l)$$

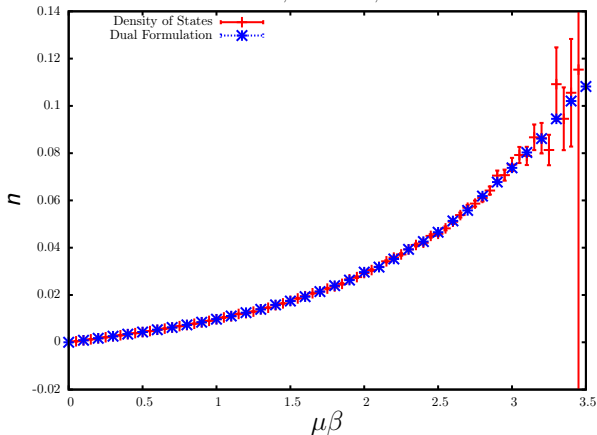
$$n = \frac{1}{V_3} \frac{2}{Z} \int_0^{l_{\max}} dl \rho(l) \sin(2k \sinh(\mu)l) l$$

$$\chi_n = \frac{1}{V_3} \left\{ \frac{2}{Z} \int_0^{l_{\max}} dl \rho(l) \cos(2k \sinh(\mu)l) l^2 + \left(\frac{2}{Z} \int_0^{l_{\max}} dl \rho(l) \sin(2k \sinh(\mu)l) \right)^2 \right\}$$

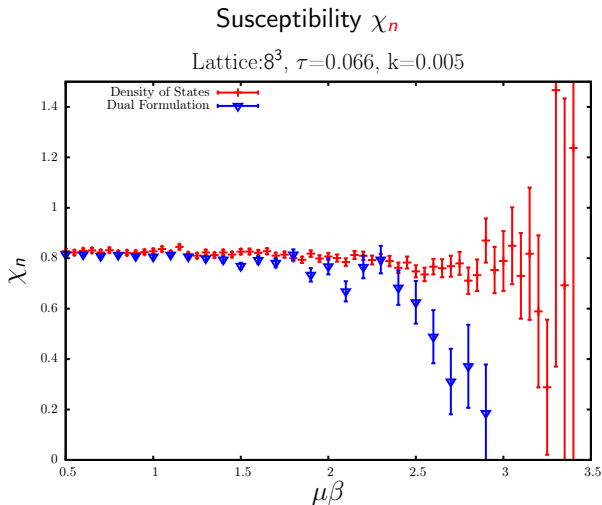
- For reliable results we need intervals of order $\Delta_n \approx \frac{1}{2k \sinh \mu}$

Particle number density n

Lattice: 8^3 , $\tau=0.066$, $k=0.005$



- For dual results see Y.D.Mercado, C. Gattringer Nuclear Physics B862 (2012) 737-750
- We find a good agreement for chemical potential up to $\mu\beta \approx 3$



- We find a good agreement for chemical potential up to $\mu\beta \approx 2$
- We need to reduce the size of the intervals Δ to explore larger $\mu\beta$

- We have developed further the density of states using the $SU(3)$ -spin system
- The density of states has been calculated with restricted Monte Carlo updates and the functional fit approach (FFA)
- Good agreement with results from the dual approach
- At very large μ the rapidly oscillating factor $\cos(2k \sinh(\mu)l)$ limits the accuracy of DoS