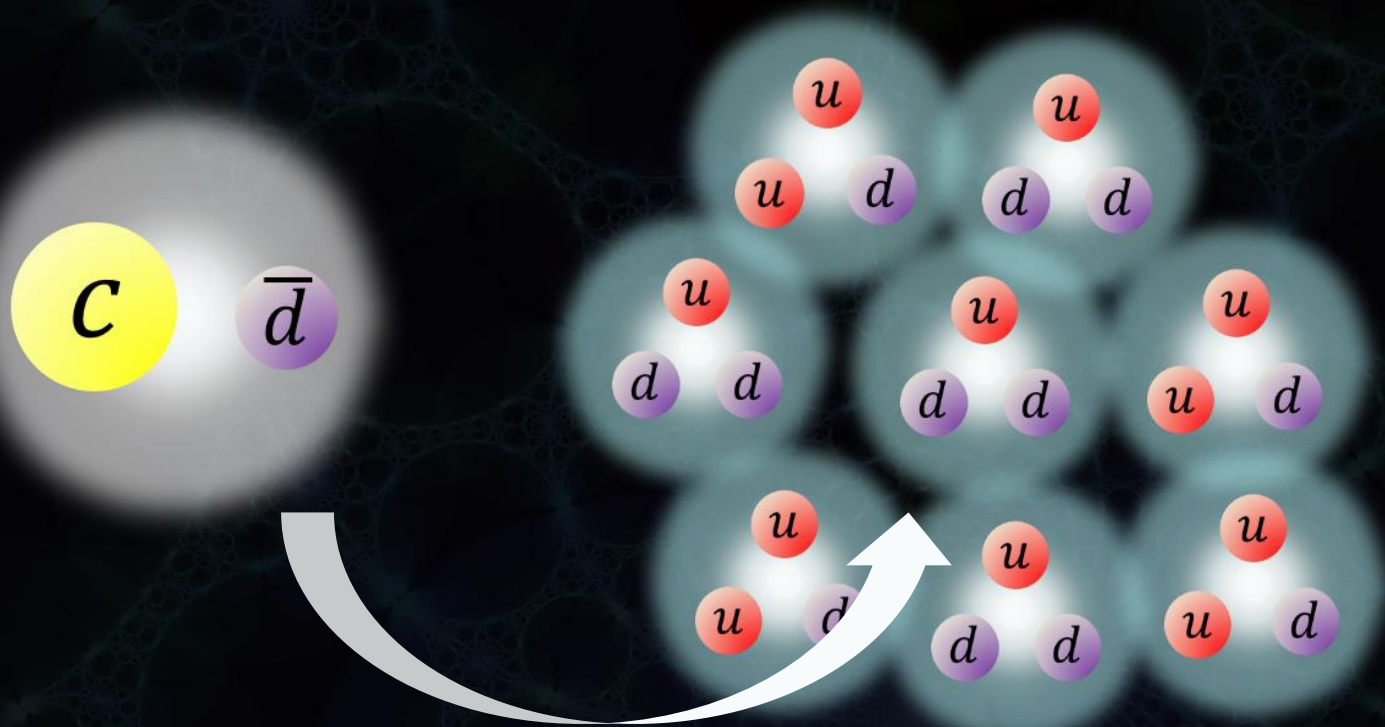


D meson properties in nuclear medium from QCD sum rules



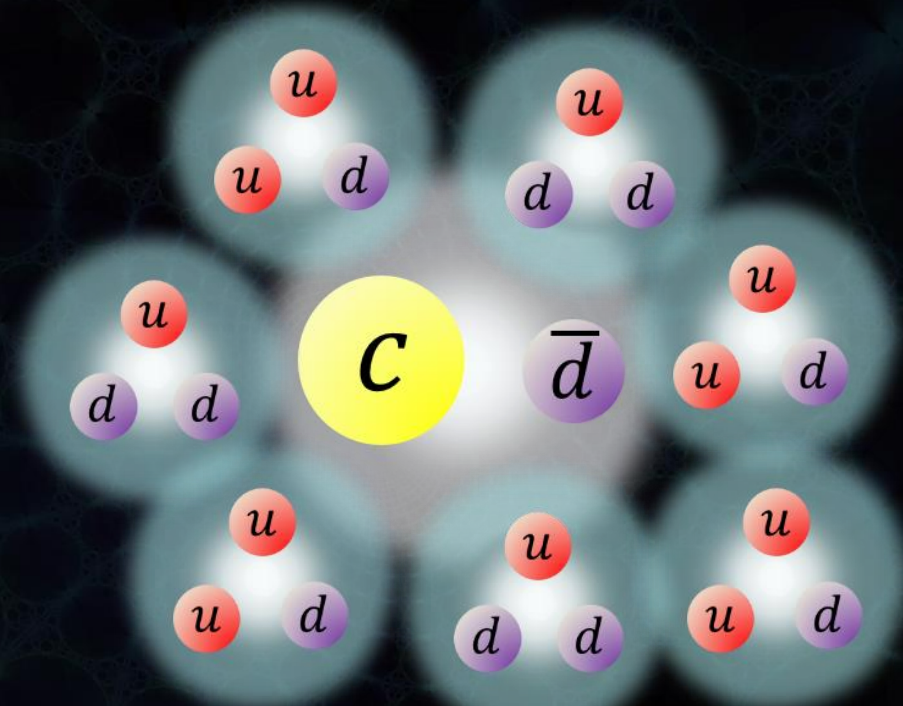
Kei Suzuki (RIKEN)

Collaboration with

Philipp Gubler (ECT*) and Makoto Oka (TITech)

Outline of talk

1. Introduction - Hadrons **in nuclear matter**
 - Chiral symmetry restoration
 - Charge symmetry breaking
2. QCD sum rules **in nuclear matter**
3. Results and summary



1. Introduction

Hadrons in nuclear matter

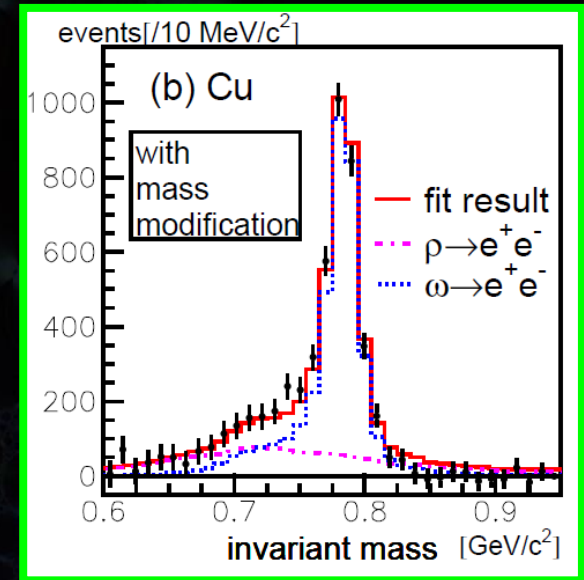
ρ , ω , ϕ mesons in nuclear matter

- Probe of chiral symmetry restoration
- Many theoretical predictions
- Some experimental indications



D meson in nuclear matter

- Probe of chiral symmetry restoration, D - N (\bar{D} - N) interaction, mesic nuclei
- Many theoretical predictions
- No experiment (future at J-PARC/FAIR)



M. Naruki et al., (KEK E-325),
PRL 96 (2006) 092301

Many theoretical works for **D meson** in matter

Coupled channel approach

- L. Tolos, J. Schaffner-Bielich, and A. Mishra, PRC70, 025203 (2004)
- M. Lutz and C. Korpa, PLB633, 43 (2006)
- T. Mizutani and A. Ramos, PRC74, 065201 (2006)
- L. Tolos, A. Ramos, and T. Mizutani, PRC77, 015207 (2008)
- L. Tolos, C. Garcia-Recio, and J. Nieves, PRC80, 065202 (2009)
- C. Jimenez-Tejero, A. Ramos, L. Tolos, and I. Vidana, PRC84, 015208 (2011)

Mean field approach

- A. Mishra, E. Bratkovskaya, J. Schaffner-Bielich, S. Schramm, and H. Stoecker, PRC69, 015202 (2004)
- A. Mishra and A. Mazumdar, PRC79, 024908 (2009)
- A. Kumar and A. Mishra, PRC81, 065204 (2010)
- A. Kumar and A. Mishra, EPJ. A47, 164 (2011)

Pion exchange model for $D\bar{N}$

- S. Yasui and K. Sudoh, PRC87, 015202 (2013)

QMC model

- K. Tsushima, D.-H. Lu, A. W. Thomas, K. Saito, and R. Landau, PRC59, 2824 (1999)
- A. Sibirtsev, K. Tsushima, and A. W. Thomas, EPJ. A6, 351 (1999)

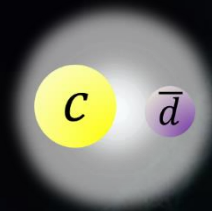
QCD sum rules

- P. Morath, W. Weise, and S.-H. Lee (1999)
- A. Hayashigaki, PLB487, 96 (2000)
- T. Hilger, R. Thomas, and B. Kampfer, Phys. Rev. C79,025202 (2009)
- K. Azizi, N. Er, and H. Sundu, EPJ. C74, 3021 (2014)
- W.Z. Gang (2015) arXiv:1501.05093 [hep-ph]

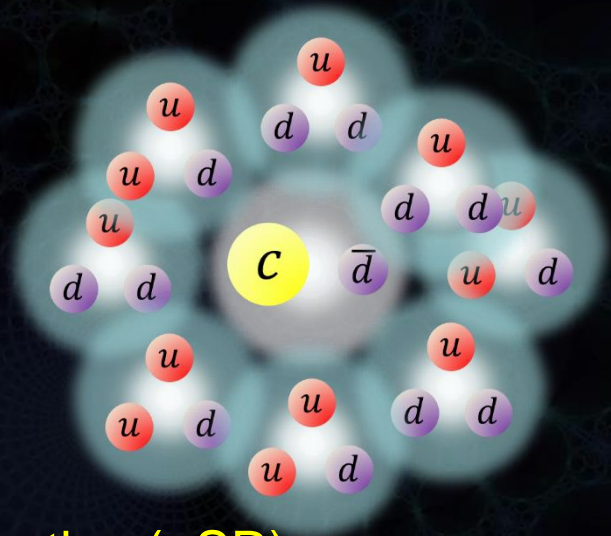
D meson in nuclear matter

If a D meson is put into nuclear matter, what will happen ?

In vacuum



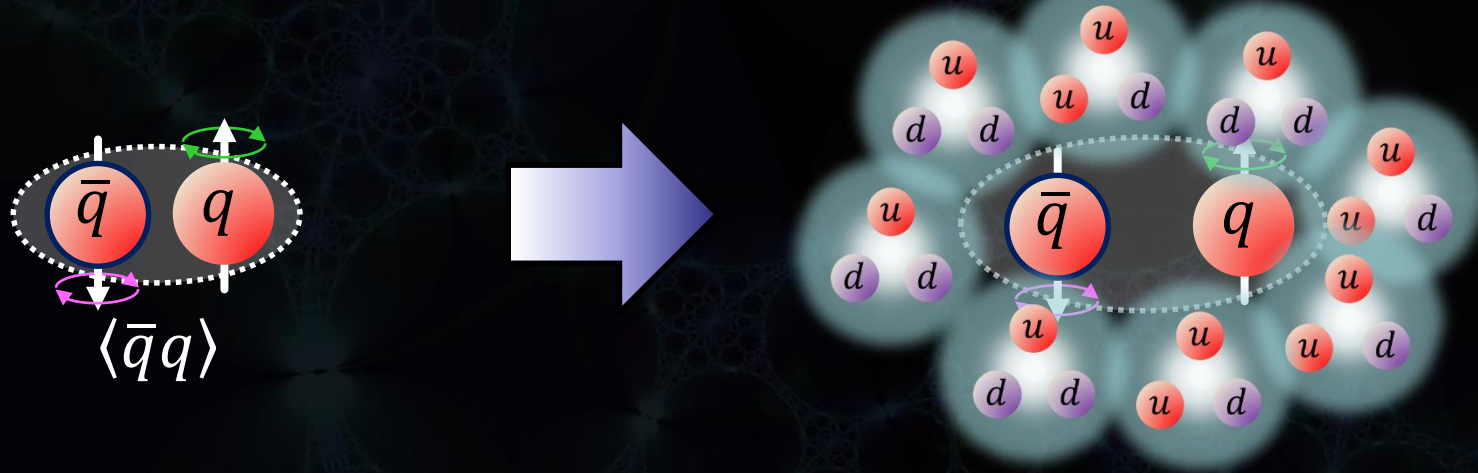
In nuclear matter



Key points

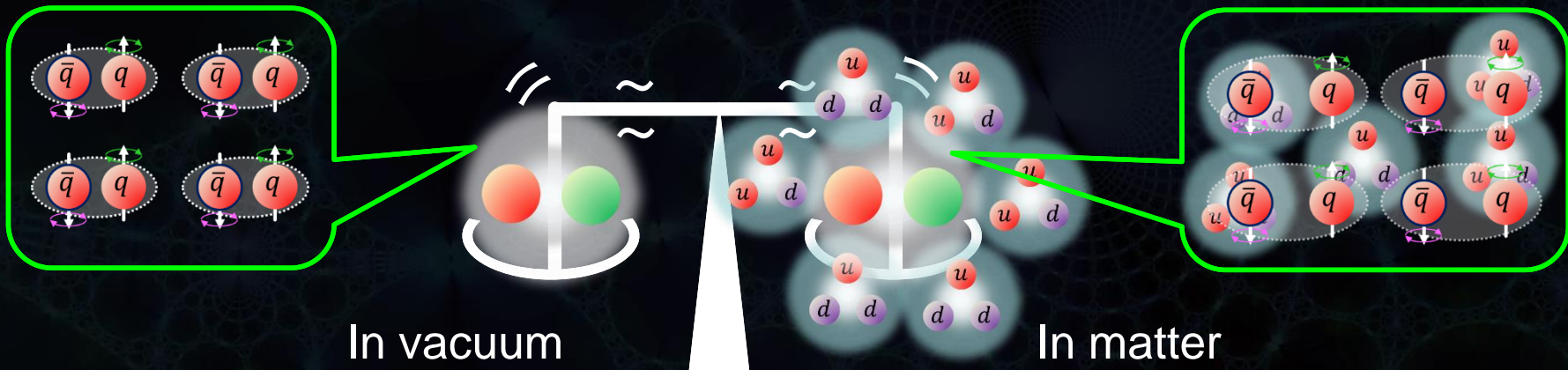
1. Chiral symmetry restoration (χ SR)
2. Charge symmetry breaking (CSB)
= Particle - anti-particle symmetry

Chiral symmetry restoration in nuclear matter



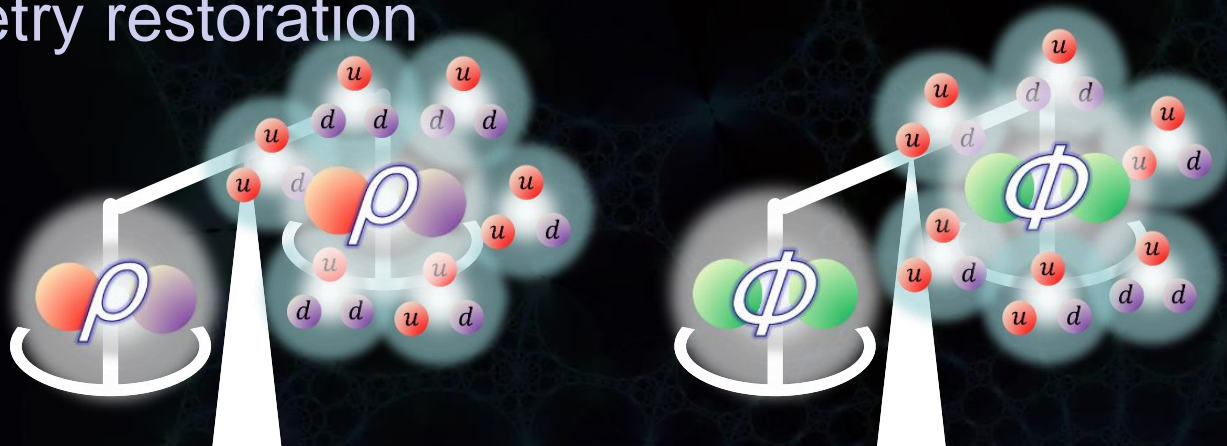
⇒ Chiral condensate is reduced in NM

- Chiral symmetry restoration changes hadron masses

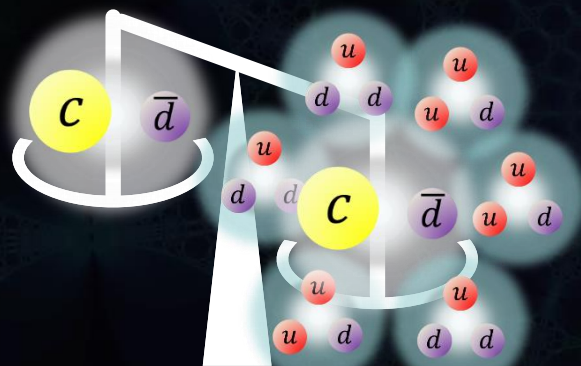


Different hadron mass shifts by Chiral symmetry restoration

- ρ and ϕ meson masses are decreased by chiral symmetry restoration



- D meson mass is increased by chiral symmetry restoration

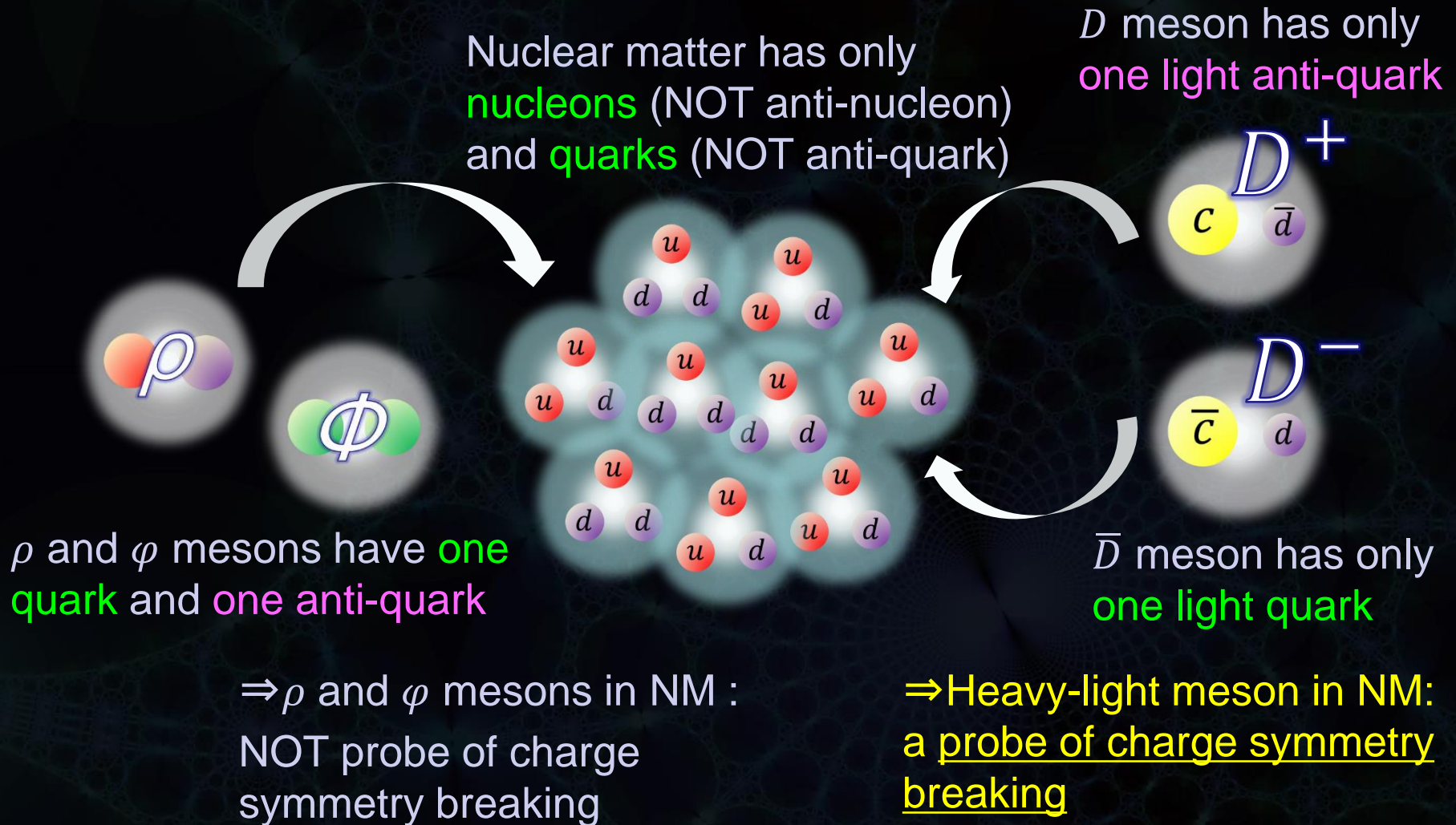


cf.) Chiral partner:
Pseudoscalar \leftrightarrow Scalar

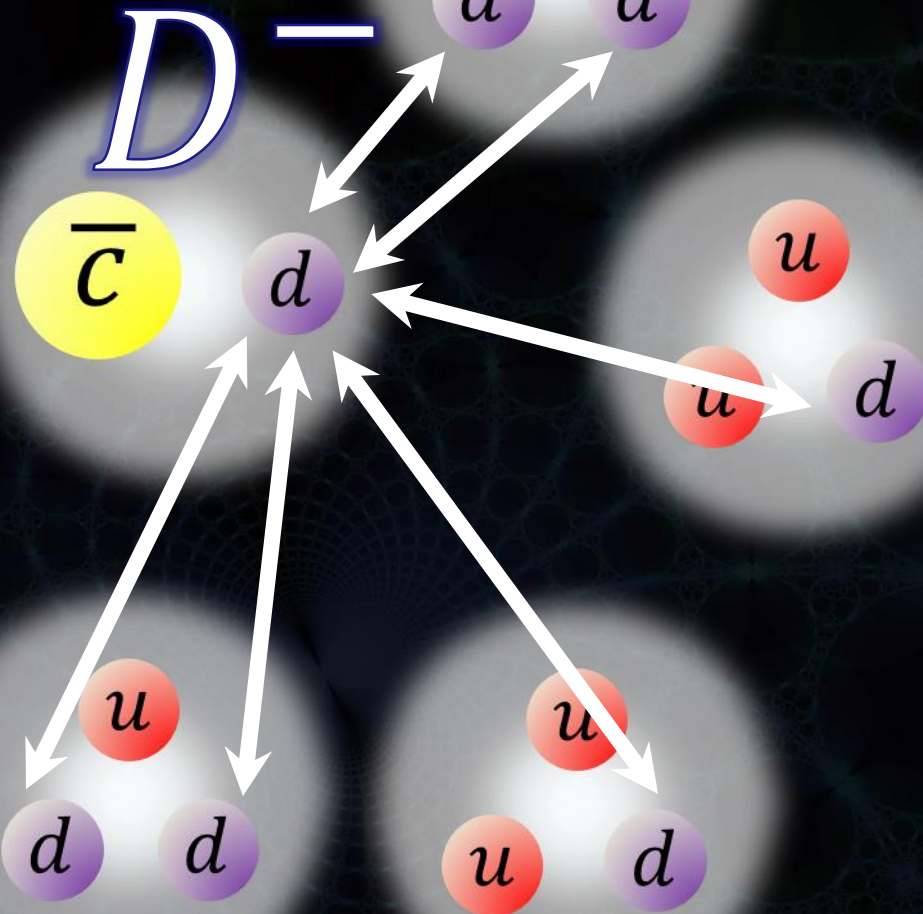


Chiral partner degenerates by chiral symmetry restoration

Charge Symmetry Breaking = imbalance b/w **particle** and **anti-particle**



Ex. Quark Pauli blocking



Only D^- feels repulsive forces from Pauli effect
 \Rightarrow positive mass shift

2. QCD sum rule in nuclear medium

QCD sum rule

Relation between operator product expansion (OPE) of QCD correlation function and hadron spectral function

$$\Pi_{\text{OPE}}(M^2) = \int_0^\infty K(s, M^2) \rho(s) ds$$

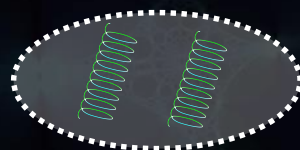
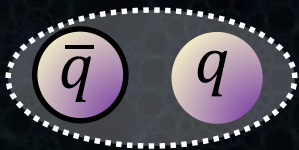
Quark and Gluon dynamics



QCD vacuum condensates

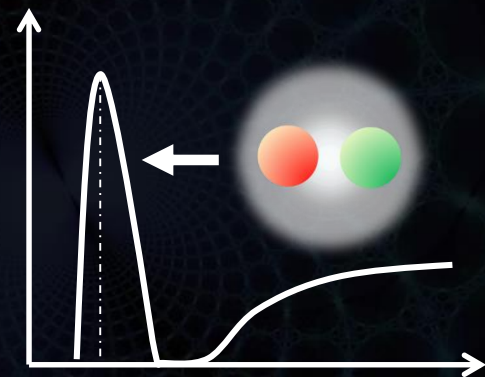
$$\langle \bar{q}q \rangle$$

$$\langle G_{\mu\nu} G^{\mu\nu} \rangle$$



etc...

Hadron properties
(mass, width...)

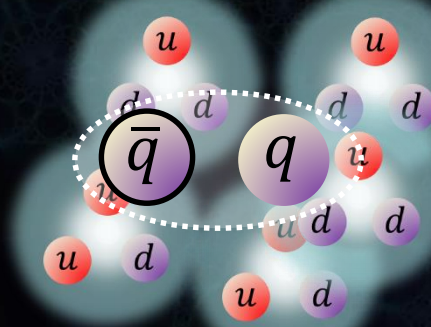
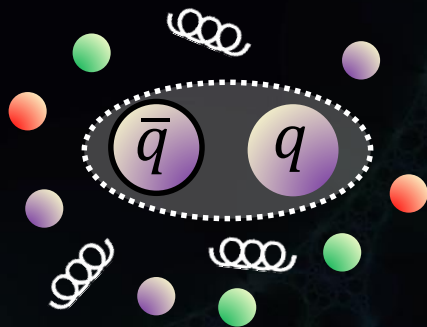


QCD sum rules in medium

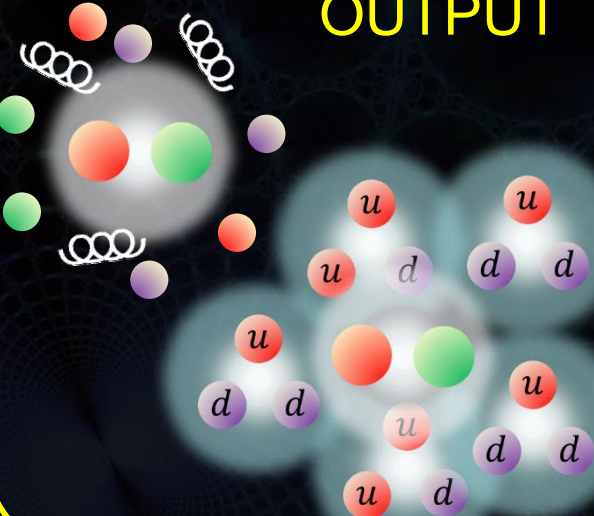
$$\Pi_{\text{OPE}}(M^2) = \int_0^\infty K(s, M^2) \rho(s) ds$$

Medium modification of OPE INPUT

T- depend. (ex. in hot π gas, QGP) density depend. (ex. in nuclear matter)



Hadron modification OUTPUT



⇒ QCD sum rule relates modification of OPE (or condensate) to modification of hadron state

QCD sum rules in nuclear matter

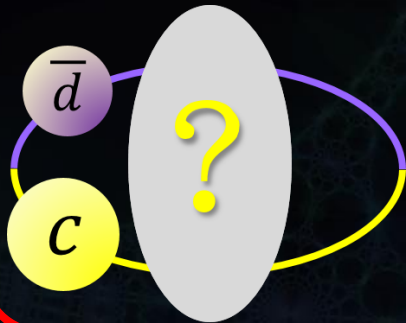
QCD sum rule

$$\Pi_{\text{OPE}}(M) = \int_0^\infty K(s, M) \rho(t) ds$$

② Kernel

Weight of spectral function
• Gaussian sum rule

① OPE



• dens.-dependence
of condensates up to
dim.5

③ Output spectral function

- Maximum entropy method (MEM)
P. Gubler and M. Oka, PTP124 (2010) 995
- Charge-conjugate projection
cf.) D. Jido, N. Kodama, and M. Oka,
PRD54, 4532 (1996)

cf.) A. Hayashigaki, PLB487 (2000) 96

T. Hilger, R. Thomas, B. Kampfer, PRC79 (2009) 025202

Condensates in nuclear matter

- Chiral-symmetry-breaking condensates

$\langle \bar{q}q \rangle_n = \langle \bar{q}q \rangle_{vac} + \frac{\sigma_N}{2m_q} n$	$\langle \bar{q}g\sigma Gq \rangle_n = \lambda^2 \langle \bar{q}q \rangle_n$
---	--

- Others (Gluon cond. , Twist cond. , ...)

$\left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle_n = \left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle_{vac} - \frac{8M_N^0}{9} n$	$\left\langle \frac{\alpha_s}{\pi} \left(\frac{(vG)^2}{v^2} - \frac{G^2}{4} \right) \right\rangle_n = (-0.05 \text{ GeV}) n$
$\langle q^\dagger iD_0 q \rangle_n = \frac{3}{8} M_N A_2^q(\mu^2) n$	$\left[\langle \bar{q} D_0^2 q \rangle_n - \frac{1}{8} \langle \bar{q} g \sigma G q \rangle_n \right] = -\frac{3}{4} M_N^2 e_2^q(\mu^2) n$

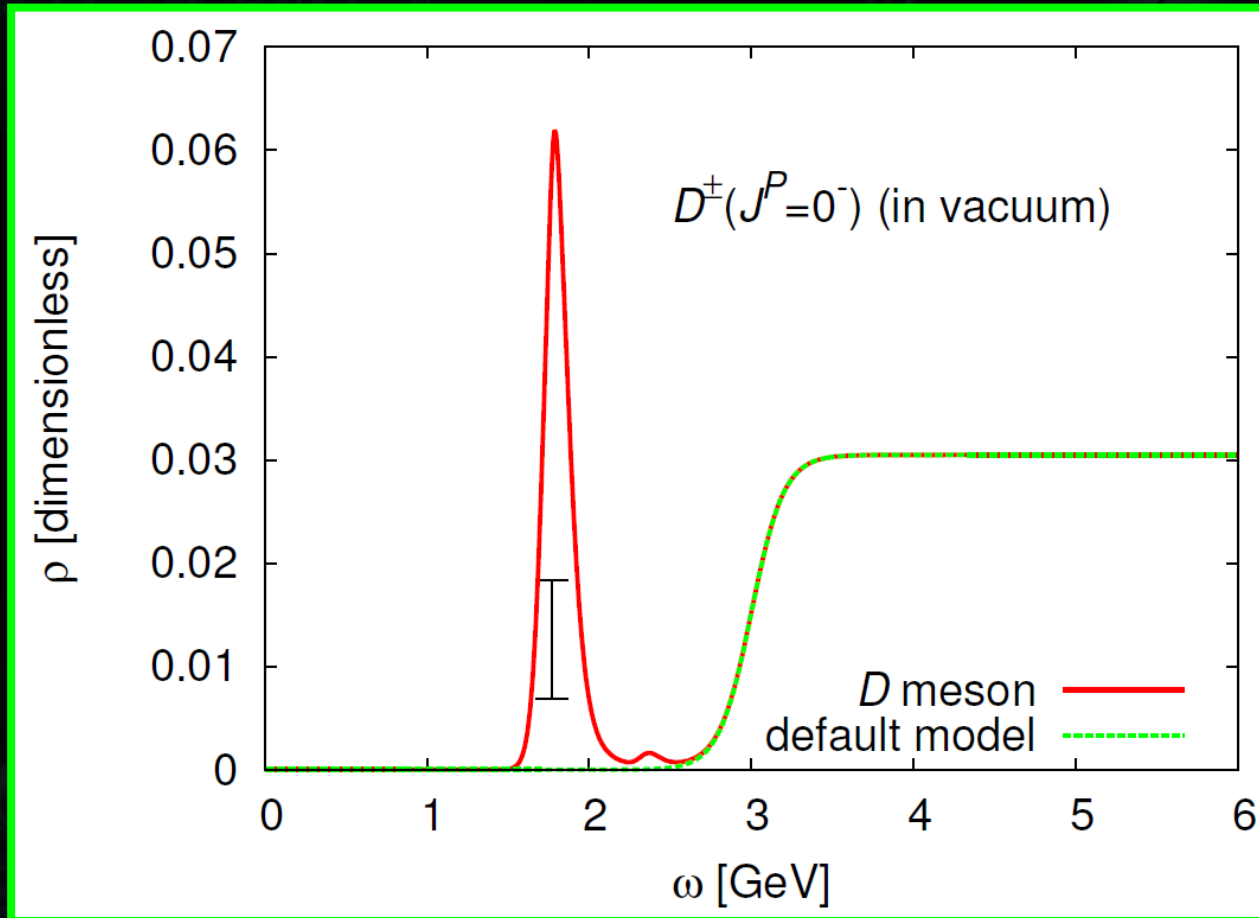
⇒ Same sign contribution to D and \bar{D} meson

- Charge-symmetry-breaking condensates

$\langle q^\dagger q \rangle_n = \frac{3}{2} n$	$\langle q^\dagger g \sigma G q \rangle_n = (0.33 \text{ GeV}^2) n$	$\langle q^\dagger D_0^2 q \rangle_n = -\frac{1}{4} M_N^2 A_3^q(\mu^2) n + \frac{1}{12} \langle q^\dagger g \sigma q \rangle_n$
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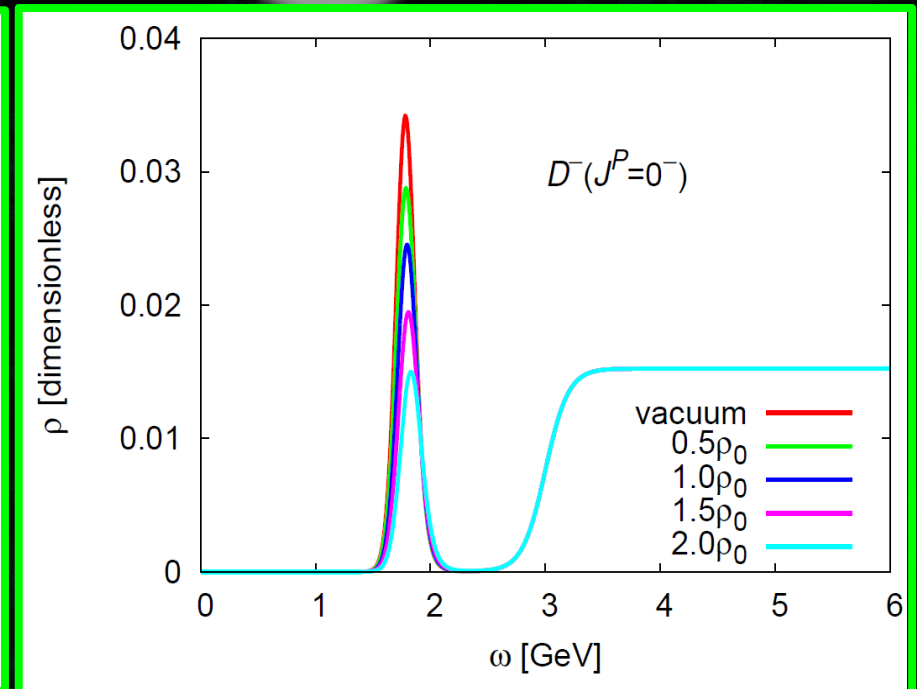
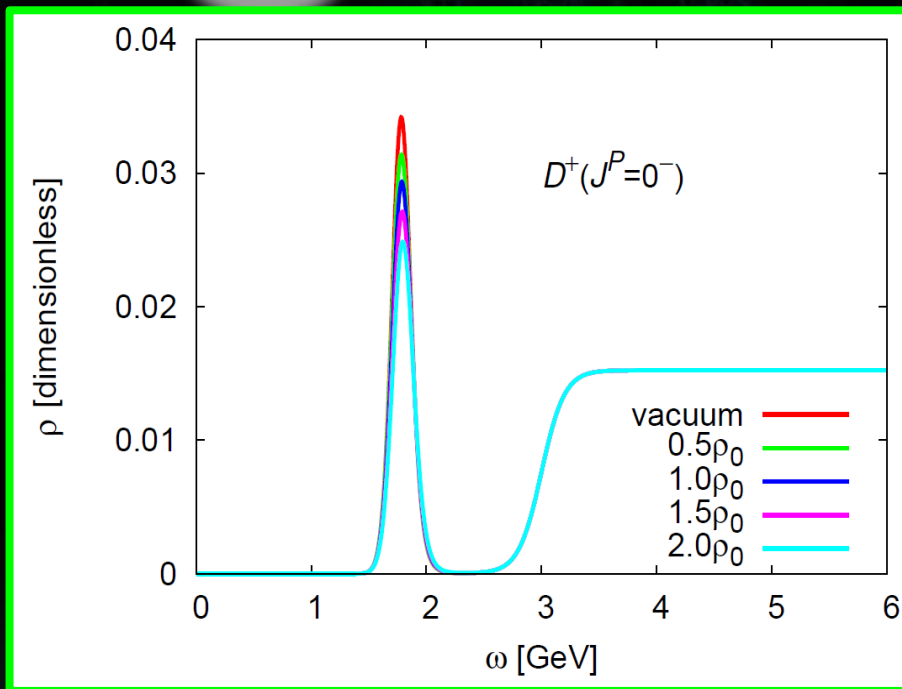
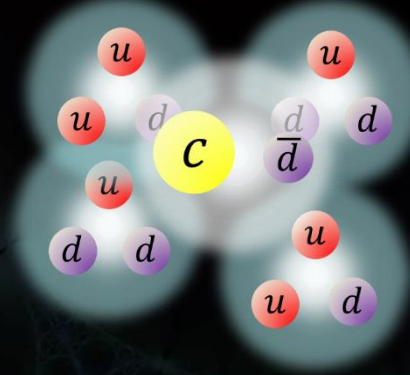
⇒ Opposite sign contribution to D and \bar{D} meson

D meson spectral function (in vacuum)



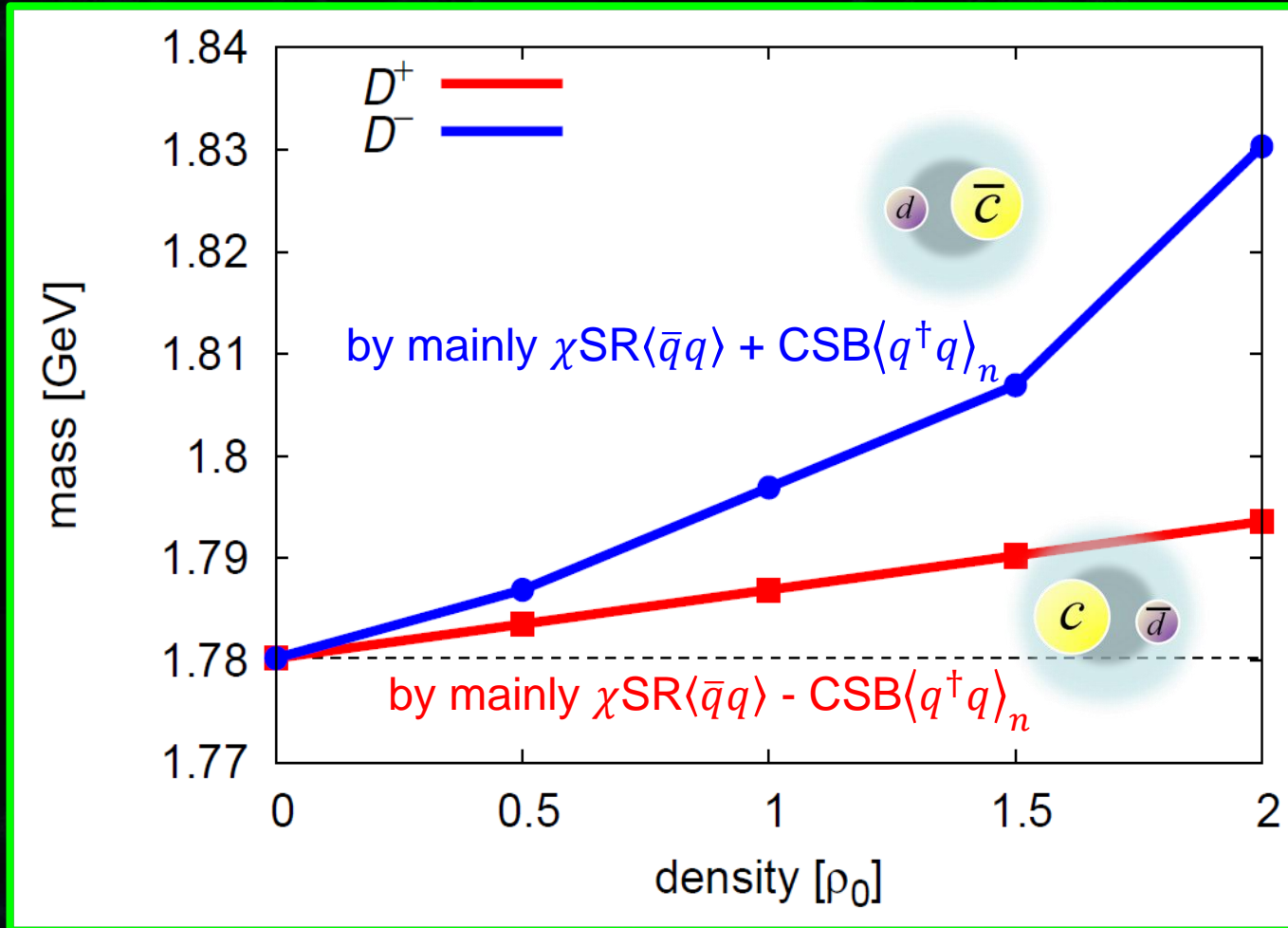
Mass : **1.78 GeV** Exp. : 1.87 GeV

D meson spectral function (in medium)





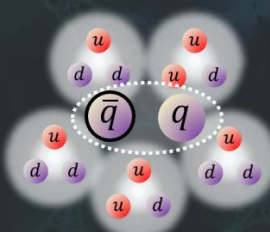
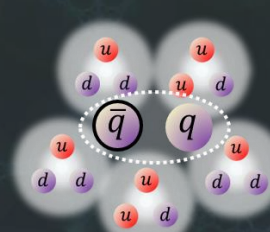
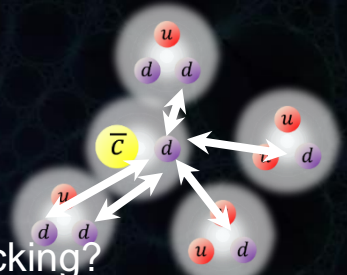
\Rightarrow Peak position in D^\pm shifts to higher energy side with increasing density (D^+ : $\sim 5\text{MeV}$ D^- : $\sim 15\text{MeV}$ at ρ_0)

Comparison of D^+ and D^-



$\Rightarrow D^+ - D^-$ mass splitting is about 10 MeV at ρ_0

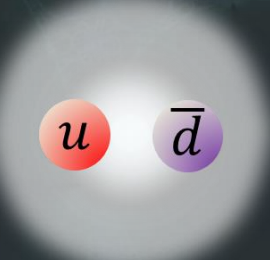

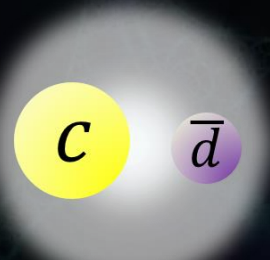

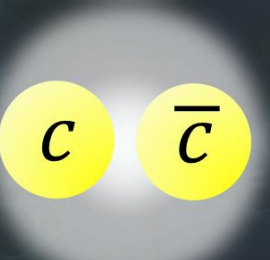
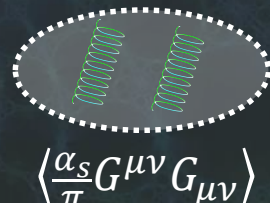
Summary of **D meson** in nuclear matter

	D+ 	D- 
$\chi_{SR} = \langle \bar{q}q \rangle$ reduction	Increase $\uparrow\uparrow$ 	Increase $\uparrow\uparrow$ 
CSB effect	Decrease \downarrow	Increase \uparrow 
Our results	Increase \uparrow (~5MeV)	More increase $\uparrow\uparrow$ (~15MeV)

\Rightarrow D meson is a good probe of χ_{SR} and CSB

Backup

Difference of meson systems

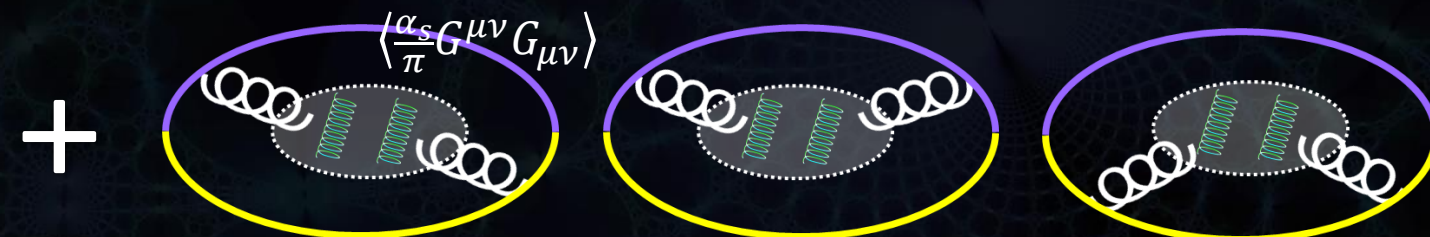
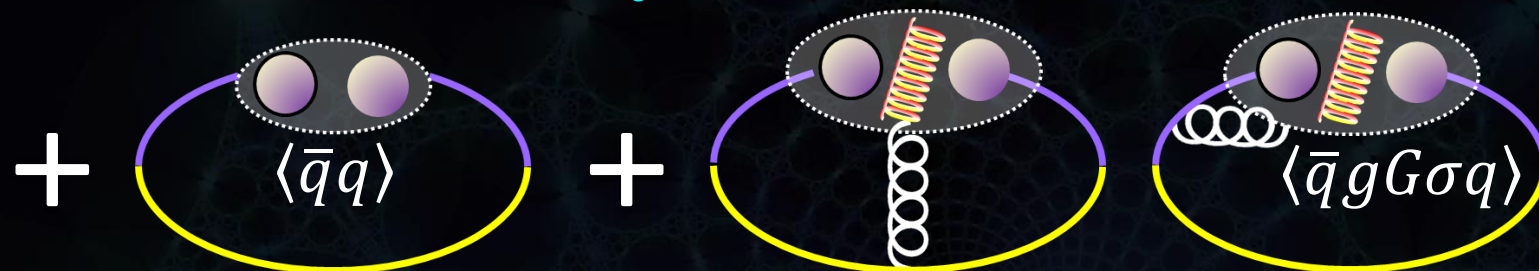
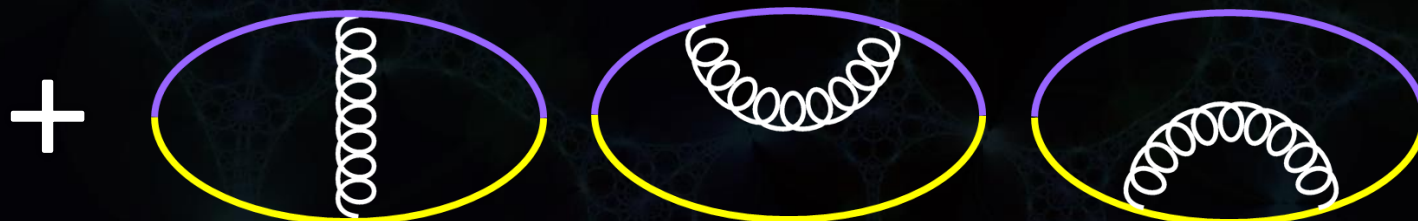
Meson	Dominant contributions in vacuum
Light-Light (ρ, ω meson) 	Probe of <u>4-quark</u> and <u>gluon condensates</u> (2-quark condensate is suppressed as $m_q \langle \bar{q}q \rangle$) 
Light-Heavy (D, B meson) 	Probe of <u>2-quark condensate</u> as $m_c \langle \bar{q}q \rangle$ 
Heavy-Heavy ($J/\psi, \Upsilon$) 	Almost <u>perturbative</u> object (Probe of gluon condensate) 

D meson OPE (in vacuum)

$\Pi_{\text{OPE}}(M^2) =$ perturbative term

$$+e^{-m_c^2/M^2} \left[-m_c \langle \bar{q}q \rangle + \frac{1}{2} \left(\frac{m_c^2}{2M^4} - \frac{1}{M^2} \right) m_c \langle \bar{q}g\sigma Gq \rangle \right. \\ \left. + \frac{1}{12} \left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle - \frac{16\pi}{27} \frac{1}{M^2} \left(1 + \frac{1}{2} \frac{m_c^2}{M^2} - \frac{1}{12} \frac{m_c^4}{M^4} \right) \alpha_s \langle \bar{q}q \rangle^2 \right]$$

- | | | |
|---|---|--|
| <ol style="list-style-type: none"> 1. Chiral condensate 2. Mixed condensate 3. Gluon condensate 4. 4-quark condensate | } | <p>Coefficients are proportional to <u>charm quark mass</u></p> <p>⇒ These terms are <u>enhanced</u></p> |
| | } | <p>Other condensates are relatively <u>suppressed</u></p> |



+ ...