

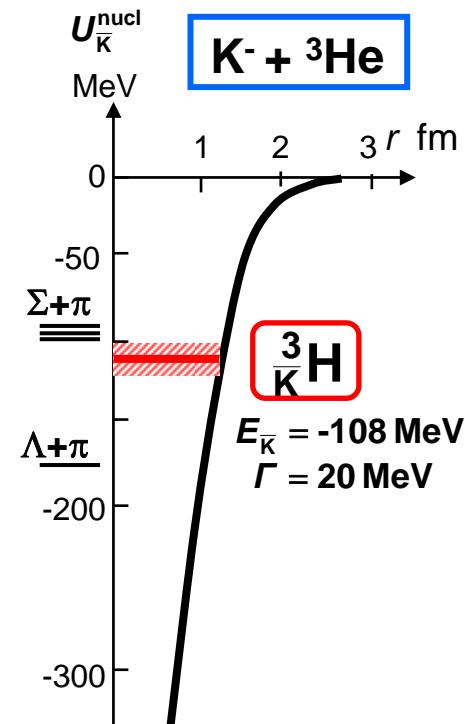
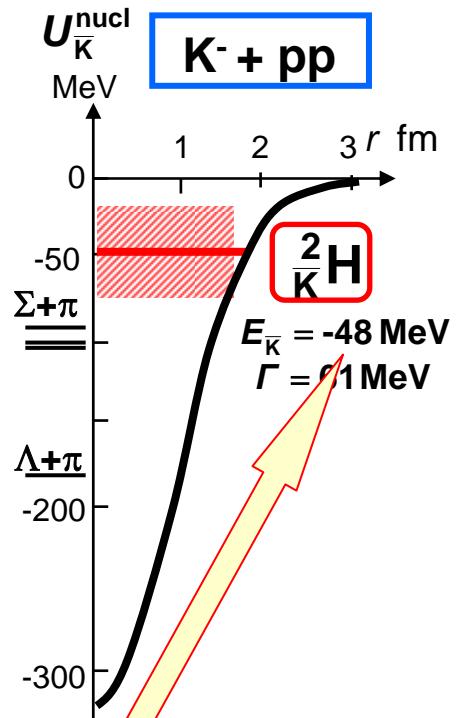
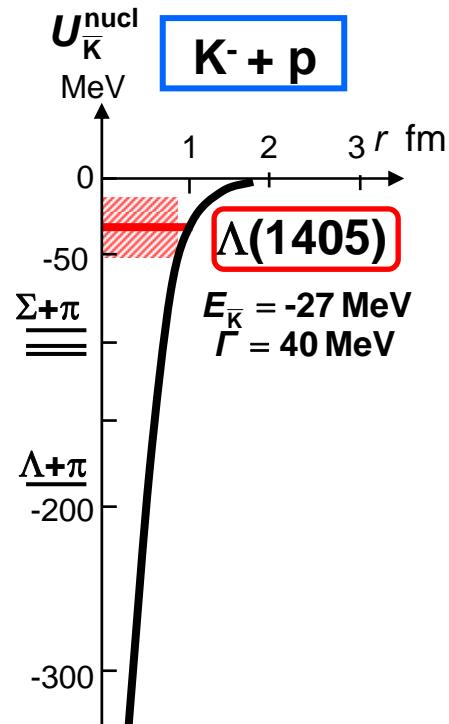
HYP2015, Sendai
September 7, 2015

Possible Existence of Λ^* Strangelets

— toward $\Lambda^* = (\Lambda\bar{p})^{l=0}$ condensed matter —

Yoshinori AKAISHI & Toshimitsu YAMAZAKI

" $\Lambda(1405)$ Ansatz"



N.V. Shevchenko, A. Gal & J. Mares, Phys. Rev. Lett. 98 (2007) 082301

$E = -55 \sim -70 \text{ MeV}$, $\Gamma = 90 \sim 110 \text{ MeV}$

Y. Ikeda & T. Sato, Phys. Rev. C 76 (2007) 035203

$E = -80 \text{ MeV}$, $\Gamma = 73 \text{ MeV}$

A. Dote, T. Hyodo & W. Weise, Phys. Rev. C 79 (2009) 014003

$E = -20 \pm 3 \text{ MeV}$, $\Gamma = 40 \sim 70 \text{ MeV}$

Y. Akaishi & T. Yamazaki, Phys. Rev. C 65 (2002) 044005

T. Yamazaki & Y. Akaishi, Phys. Lett. B 535 (2002) 70

K⁻p

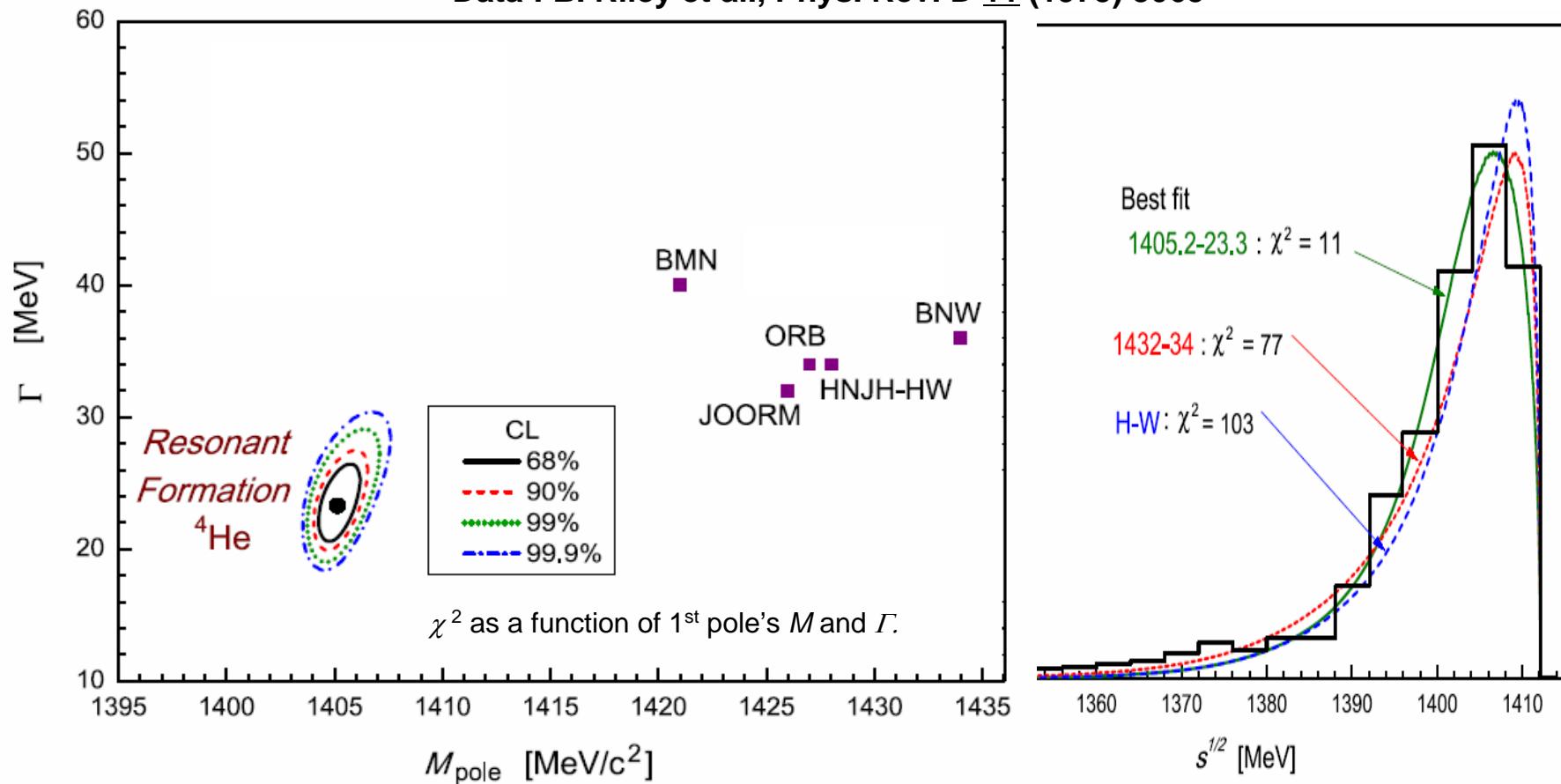
Which is
the $\Lambda(1405)$ mass,
1405 or 1420 MeV/c²?

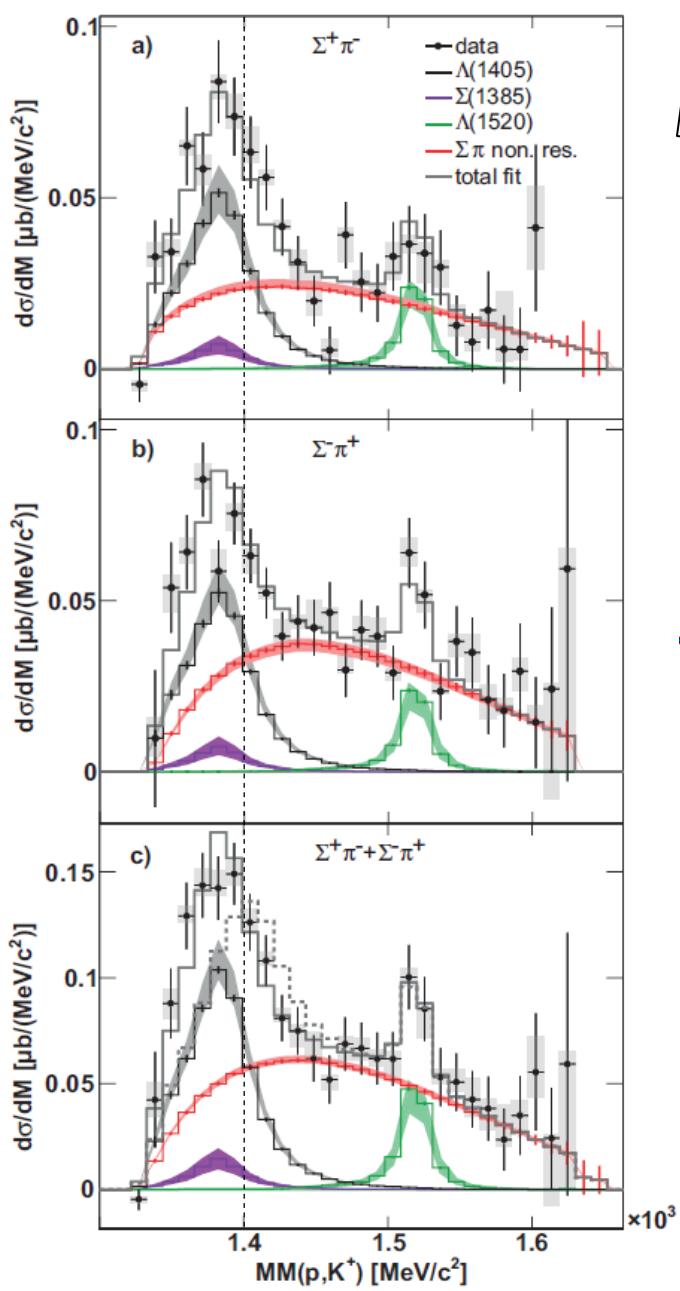
$\Sigma\pi$ invariant mass from stopped K⁻ on ${}^4\text{He}$

J. Esmaili, Y. Akaishi & T. Yamazaki, Phys. Lett. B686 (2010) 23

$$M = 1405.5^{+1.4}_{-1.0} \text{ MeV}/c^2 \quad \text{and} \quad \Gamma = 23.6^{+4}_{-3} \text{ MeV}$$

Data : B. Riley et al., Phys. Rev. D 11 (1975) 3065

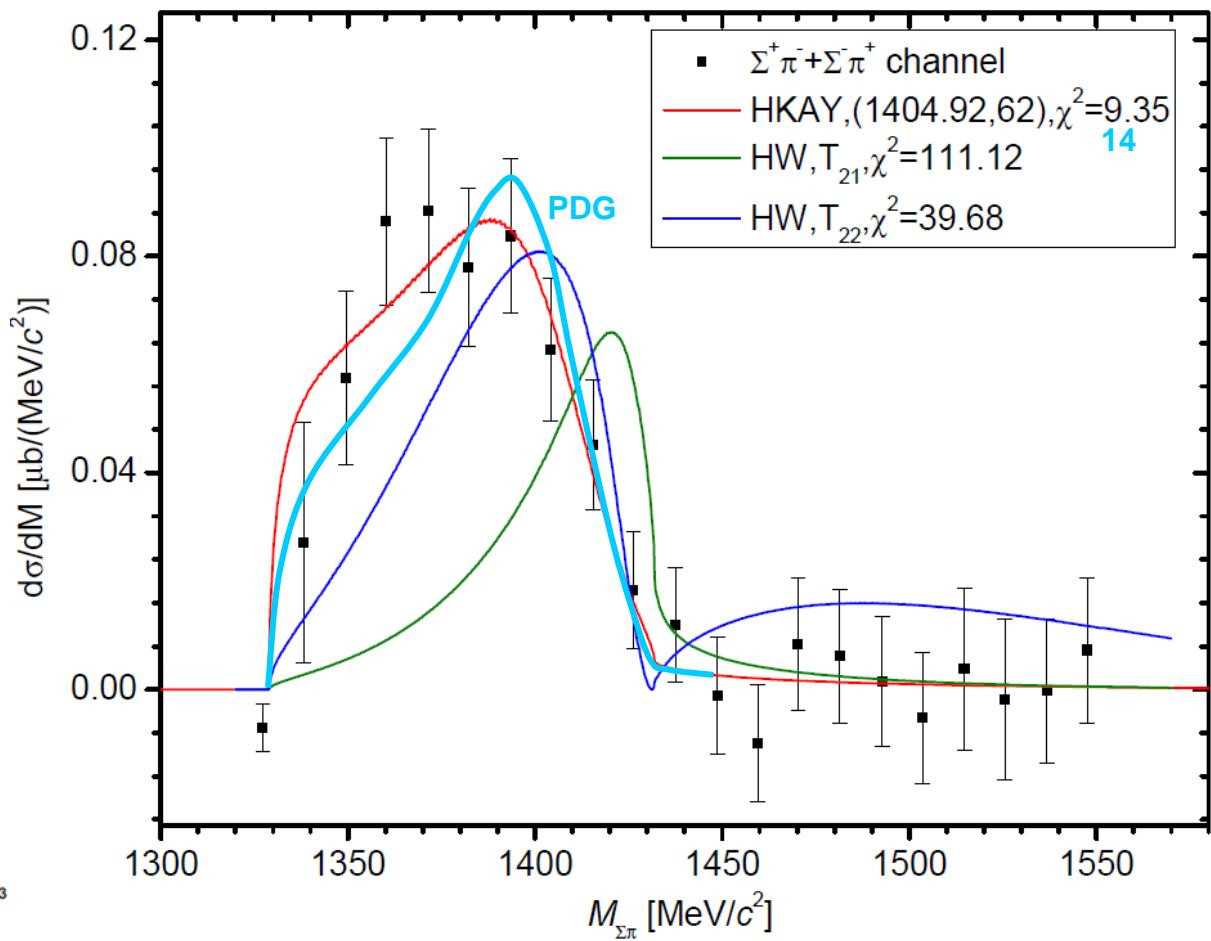




$\Lambda(1405)$ from HADES

G. Agakishiev et al., Phys. Rev. C **87** (2013) 025201

M. Hassanvand et al., Phys. Rev. C **87** (2013) 055202



$\Lambda(1405)$ $1/2^-$

$I(J^P) = 0(\frac{1}{2}^-)$ Status: ***

2015

$\Lambda(1405)$ MASS

PRODUCTION EXPERIMENTS

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	COMMENT
-------------	------	-------------	------	---------

1405.1 ± 1.3 OUR AVERAGE

1405	± 11	HASSANVAND 13	SPEC	$p p \rightarrow p\Lambda(1405)K^+$
1405	± 1.4	ESMAILI	RVUE	${}^4\text{He } K^- \rightarrow \Sigma^\pm \pi^\mp X$ at rest
1406.5	± 4.0	¹ DALITZ	91	M-matrix fit

• • • We do not use the following data for averages, fits, limits, etc. • • •

1391	± 1	700	¹ HEMINGWAY	85	HBC	$K^- p$ 4.2 GeV/c
~ 1405		400	² THOMAS	73	HBC	$\pi^- p$ 1.69 GeV/c
1405		120	BARBARO...	68B	DBC	$K^- d$ 2.1–2.7 GeV/c
1400	± 5	67	BIRMINGHAM	66	HBC	$K^- p$ 3.5 GeV/c
1382	± 8		ENGLER	65	HDBC	$\pi^- p, \pi^+ d$ 1.68 GeV/c
1400	± 24		MUSGRAVE	65	HBC	$\bar{p} p$ 3–4 GeV/c
1410			ALEXANDER	62	HBC	$\pi^- p$ 2.1 GeV/c
1405			ALSTON	62	HBC	$K^- p$ 1.2–0.5 GeV/c
1405			ALSTON	61B	HBC	$K^- p$ 1.15 GeV/c

$\Lambda(1405)$ WIDTH

PRODUCTION EXPERIMENTS

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	COMMENT
-------------	------	-------------	------	---------

50.5 ± 2.0 OUR AVERAGE

62	± 10	HASSANVAND 13	SPEC	$p p \rightarrow p\Lambda(1405)K^+$
50	± 2	¹ DALITZ	91	M-matrix fit

• • • We do not use the following data for averages, fits, limits, etc. • • •

K-pp

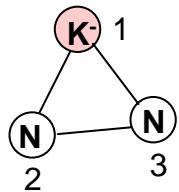
Variational wave function of K-pp

ATMS

Amalgamation of **T**wo-body correlations into **M**ultiple **S**cattering process

$$\Psi = \left[\left\{ f^{I=0}(r_{12}) \hat{P}_{12}^{I=0} + f^{I=1}(r_{12}) \hat{P}_{12}^{I=1} \right\} f_{NN}(r_{23}) f(r_{31}) + f(r_{12}) f_{NN}(r_{23}) \left\{ f^{I=0}(r_{31}) \hat{P}_{31}^{I=0} + f^{I=1}(r_{31}) \hat{P}_{31}^{I=1} \right\} \right] |T = 1/2\rangle$$

$$\hat{P}_{12}^{I=0} = \frac{1 - \bar{\tau}_K \bar{\tau}_N}{4}, \quad \hat{P}_{12}^{I=1} = \frac{3 + \bar{\tau}_K \bar{\tau}_N}{4}$$



$$|T = 1/2\rangle = \sqrt{\frac{3}{4}} \left[(\bar{K}_1 N_2)^{0,0} p_3 \right] + \sqrt{\frac{1}{4}} \left[-\sqrt{\frac{1}{3}} (\bar{K}_1 N_2)^{1,0} p_3 + \sqrt{\frac{2}{3}} (\bar{K}_1 N_2)^{1,1} n_3 \right]$$

A p*

Euler-Lagrange equation

$$\delta_f \left\{ \langle \Psi | H | \Psi \rangle - \lambda \langle \Psi | \Psi \rangle \right\} = 0$$

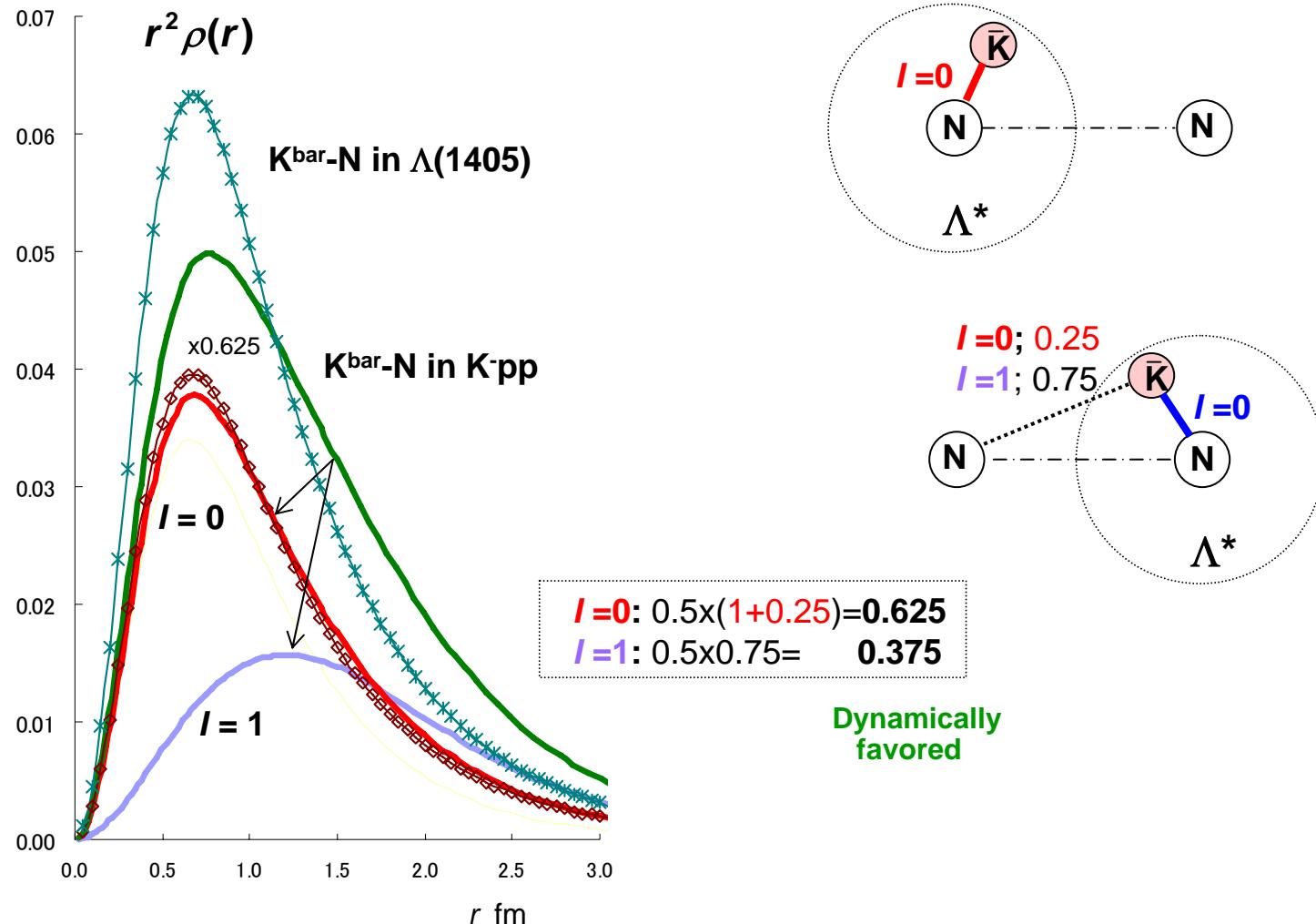
$$v_{\bar{K}N}^{T=0}(r) = \{-595 - i83\}_{\text{MeV}} \exp\left\{-\left(r/0.66_{\text{fm}}\right)^2\right\}$$

$$v_{\bar{K}N}^{T=1}(r) = \{-175 - i105\}_{\text{MeV}} \exp\left\{-\left(r/0.66_{\text{fm}}\right)^2\right\}$$

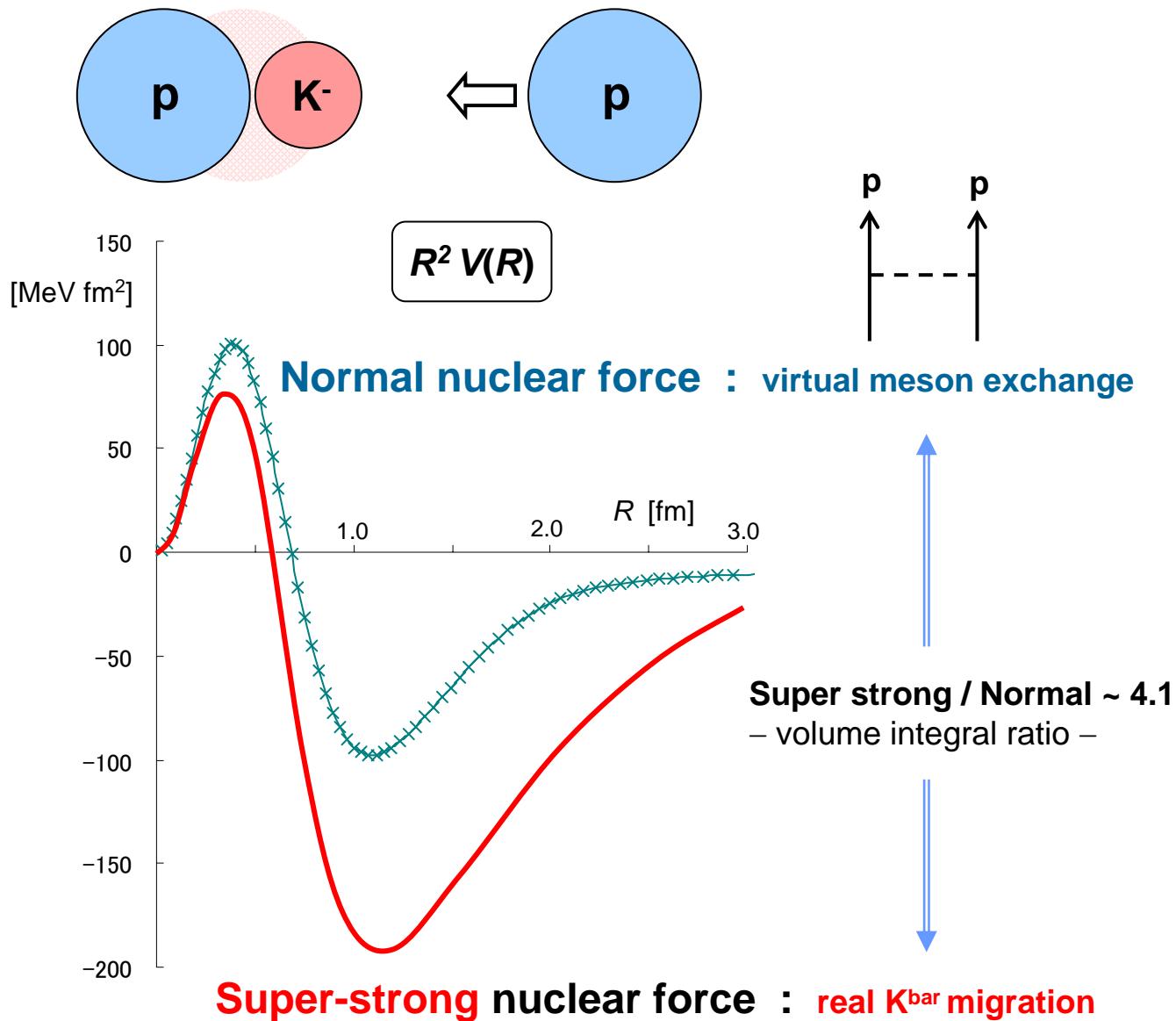
$$v_{NN}(r) = 2000_{\text{MeV}} \exp\left\{-\left(r/0.447_{\text{fm}}\right)^2\right\} - 270_{\text{MeV}} \exp\left\{-\left(r/0.942_{\text{fm}}\right)^2\right\} - 5_{\text{MeV}} \exp\left\{-\left(r/2.5_{\text{fm}}\right)^2\right\}$$

Density distributions of \bar{K} -N

T. Yamazaki and Y. Akaishi, Phys. Rev. C 76 (2007) 045201

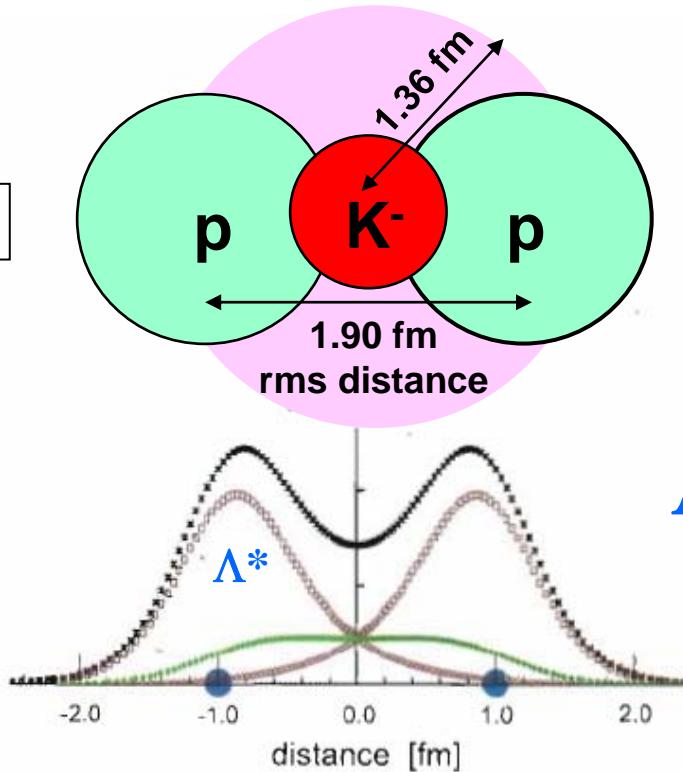


Adiabatic p-p potential in K⁻pp



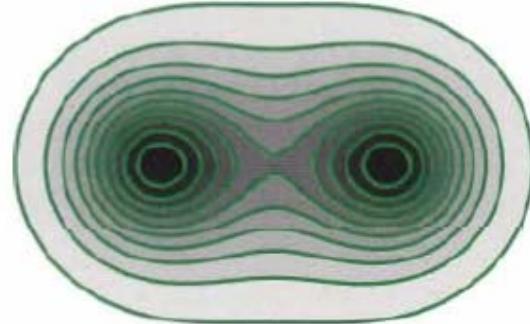
K⁻pp quasi-bound state

2002

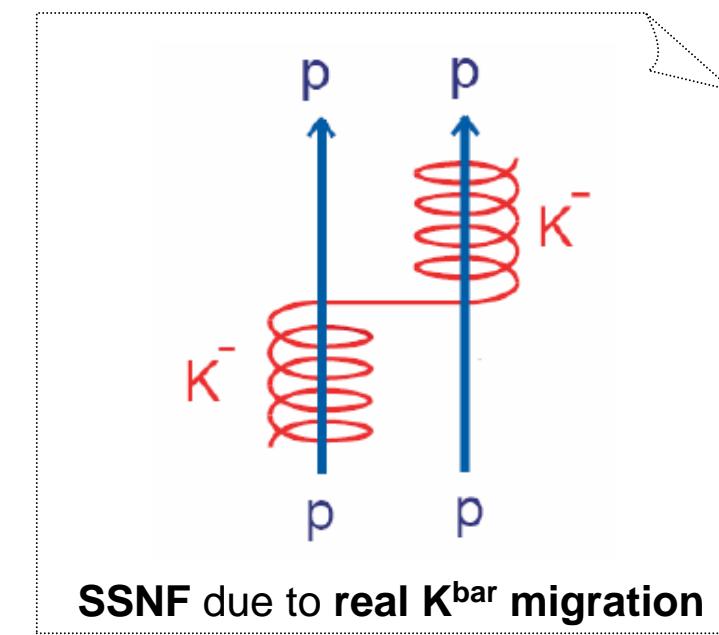


Λ^* -p structure

2007



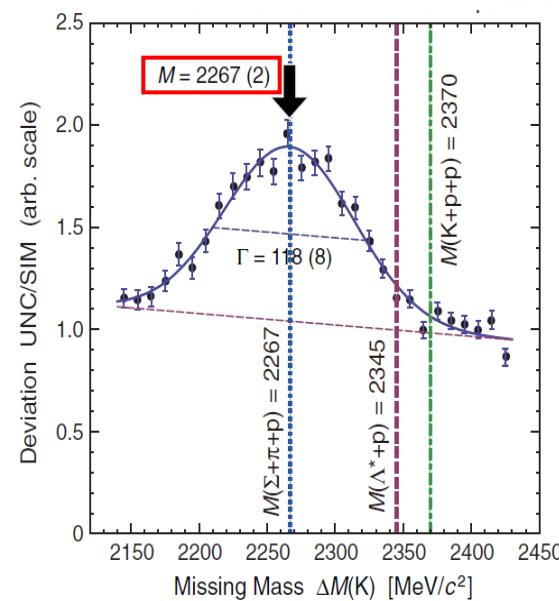
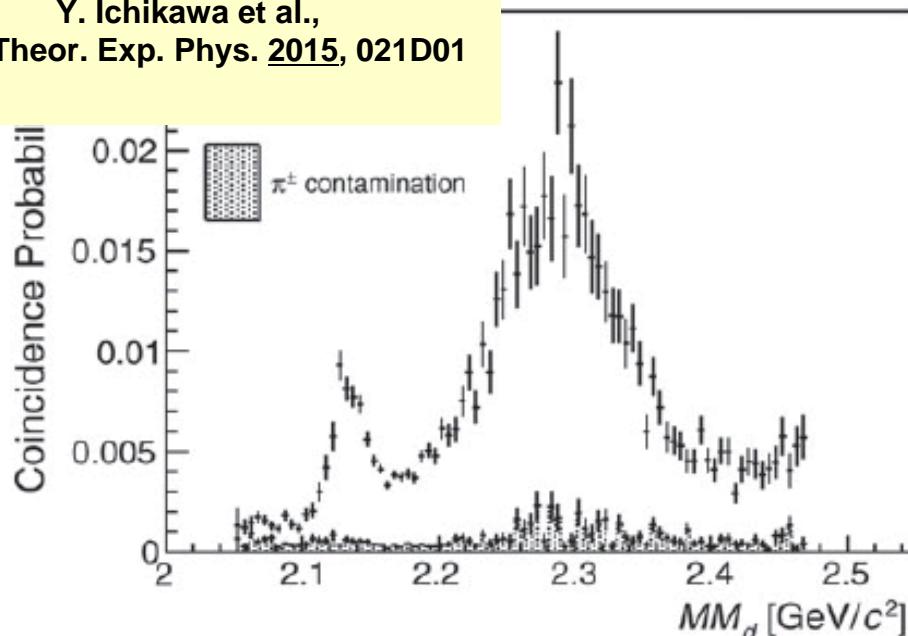
T. Yamazaki and Y. Akaishi, Phys. Rev. C 76 (2007) 045201



SSNF due to **real $K^{\bar{b}ar}$** migration

E27@J-PARC

Y. Ichikawa et al.,
Prog. Theor. Exp. Phys. 2015, 021D01



Λ^*_{1405}

Attractive interaction is enhanced;
 $s \rightarrow f s_0$, $s_0 = -1.37$

f	V	T	BE
1.00	-182.9	156.2	26.6
1.05	-219.2	182.7	36.4
1.10	-256.1	208.6	47.5
1.15	-293.6	233.9	59.7

Theor.
Yamazaki-Akaishi

$$E_K = -48 \text{ MeV}$$

$$\Gamma = 61 \text{ MeV}$$

17% enhanced
 $K^{\bar{N}}$ interaction

S. Maeda et al.,
Proc. Jpn. Acad. B 89
(2013) 418

$$E_K = -103 \text{ MeV}$$

$$\Gamma = 118 \text{ MeV}$$

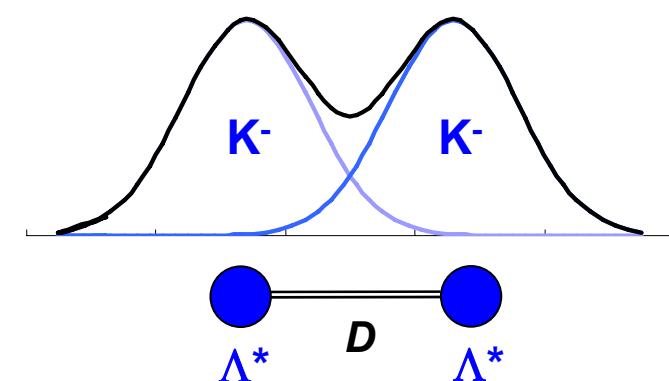
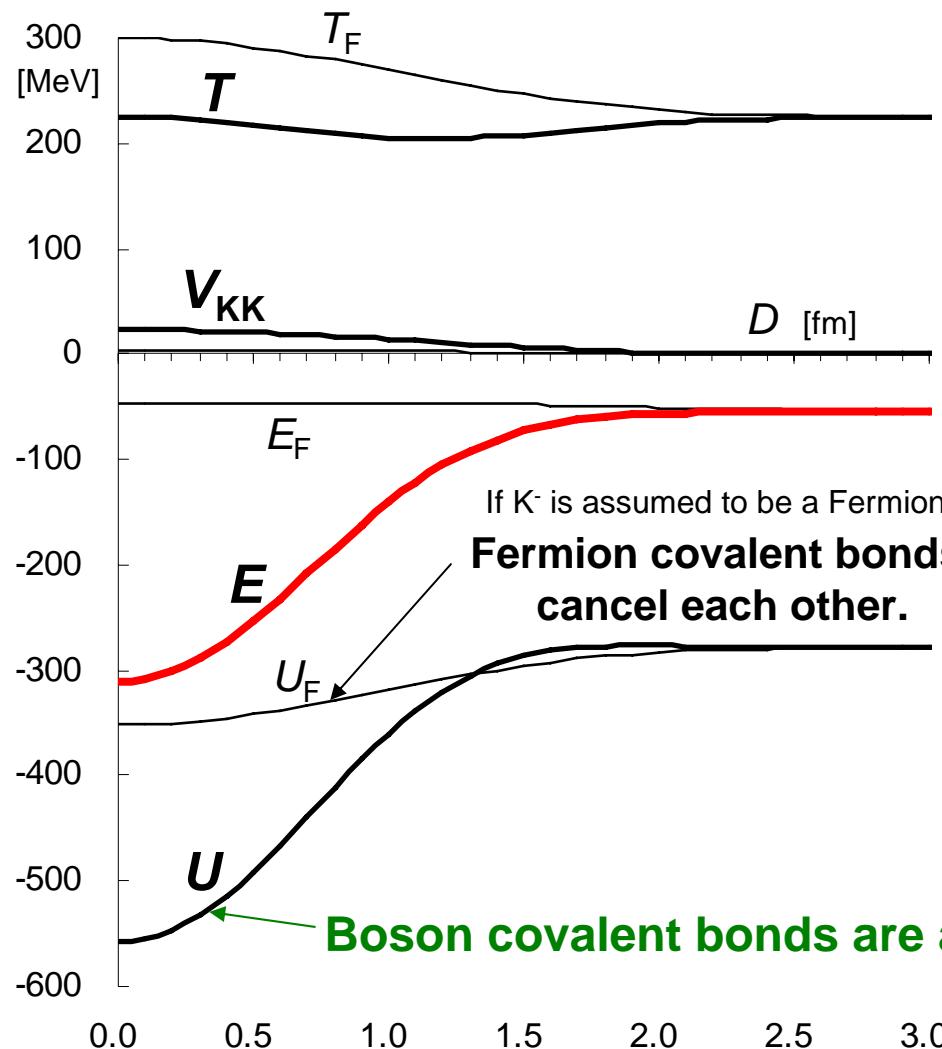
DISTO

T. Yamazaki et al.,
Phys. Rev. Lett.
104 (2010) 132502

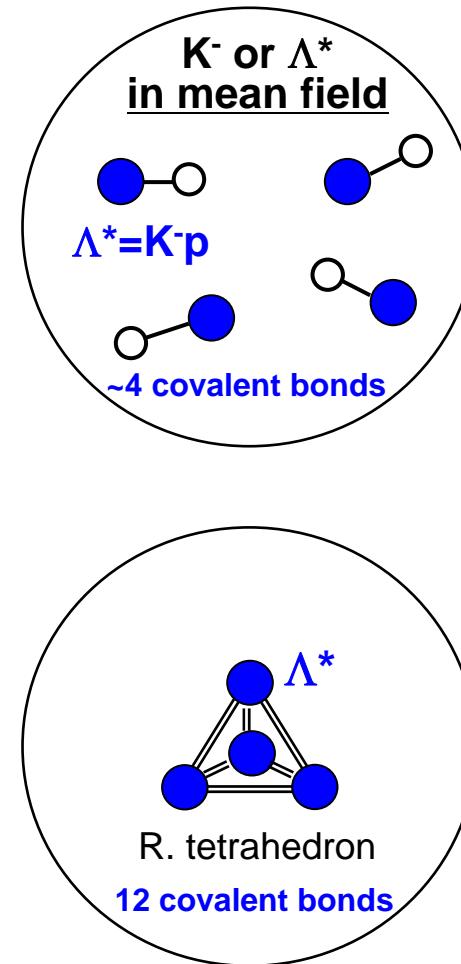
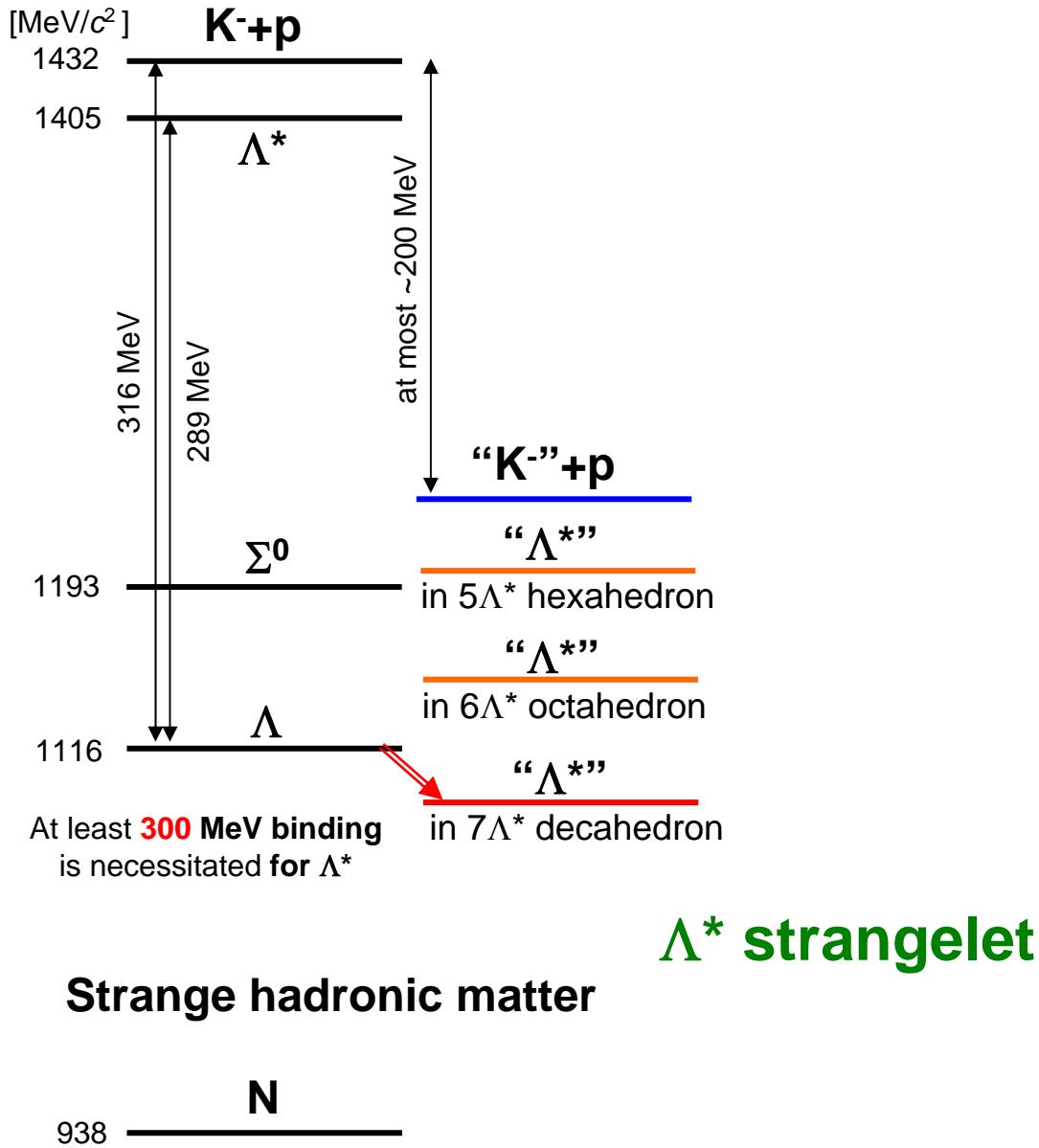
K-K⁻pp & strangelets

$\Lambda^*\Lambda^*$ model for K-K-pp

$$\Phi(\vec{r}_1, \vec{r}_2) = N\{\phi_a(\vec{r}_1)\phi_b(\vec{r}_2) \pm \phi_b(\vec{r}_1)\phi_a(\vec{r}_2)\}$$



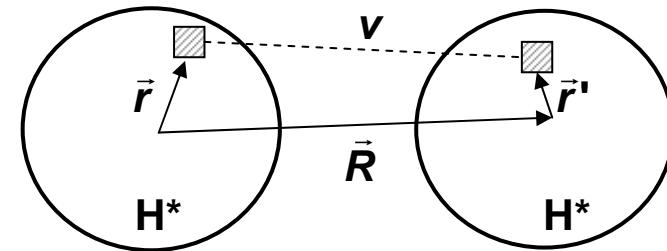
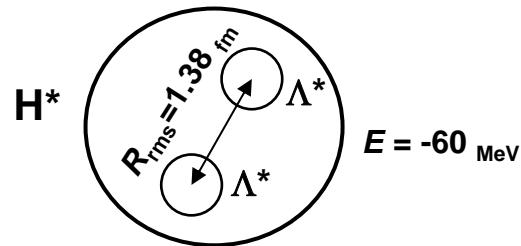
Possible existence of “few-body Λ^* strangelet”



H*-H* double folding potential

Λ^* - Λ^* potential

$$v(r_{\text{rel}}) = 751_{\text{MeV}} \exp\left\{-\left(\frac{r_{\text{rel}}}{0.6_{\text{fm}}}\right)^2\right\} - 416_{\text{MeV}} \exp\left\{-\left(\frac{r_{\text{rel}}}{1.2_{\text{fm}}}\right)^2\right\}$$



$$U(\vec{R}) = \int d\vec{r}' d\vec{r} \rho_B(\vec{r}') v(\vec{R} + \vec{r}' - \vec{r}) \rho_A(\vec{r})$$

H*-H* double folding potential

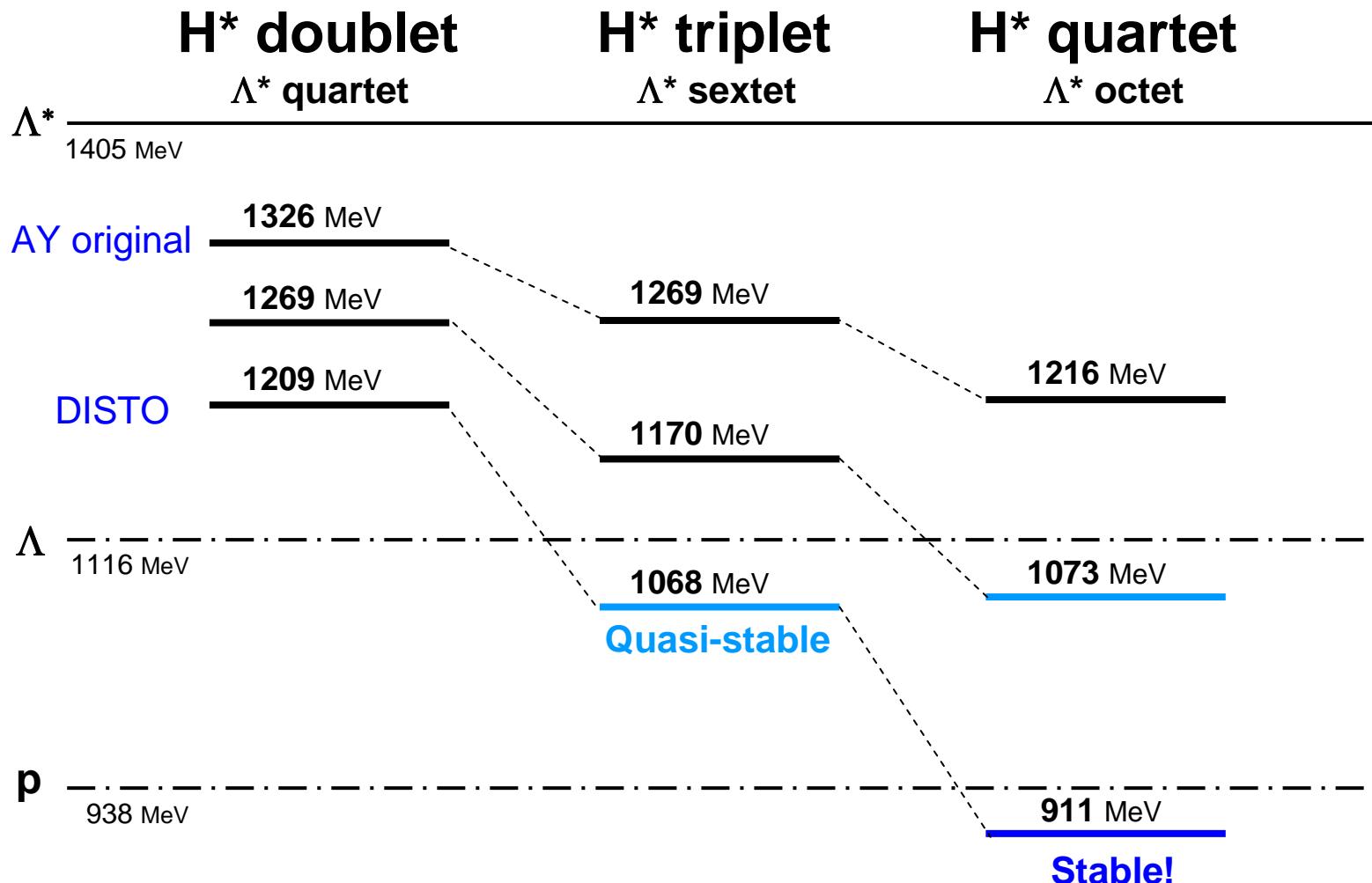
$$U(R) = 662_{\text{MeV}} \exp\left\{-\left(\frac{R}{0.99_{\text{fm}}}\right)^2\right\} - 967_{\text{MeV}} \exp\left\{-\left(\frac{R}{1.44_{\text{fm}}}\right)^2\right\}$$

Pauli forbidden state

$$\phi_{0s}(R) = \left(\frac{A}{\pi}\right)^{3/4} \exp\left\{-\frac{1}{2}AR^2\right\}, \quad A = \frac{3}{2R_{\text{rms}}} = 0.79 \text{ fm}^{-2}$$

Λ^* effective mass in H^* strangelet

$$H^* = \Lambda^* \Lambda^* = K^- K^- pp$$



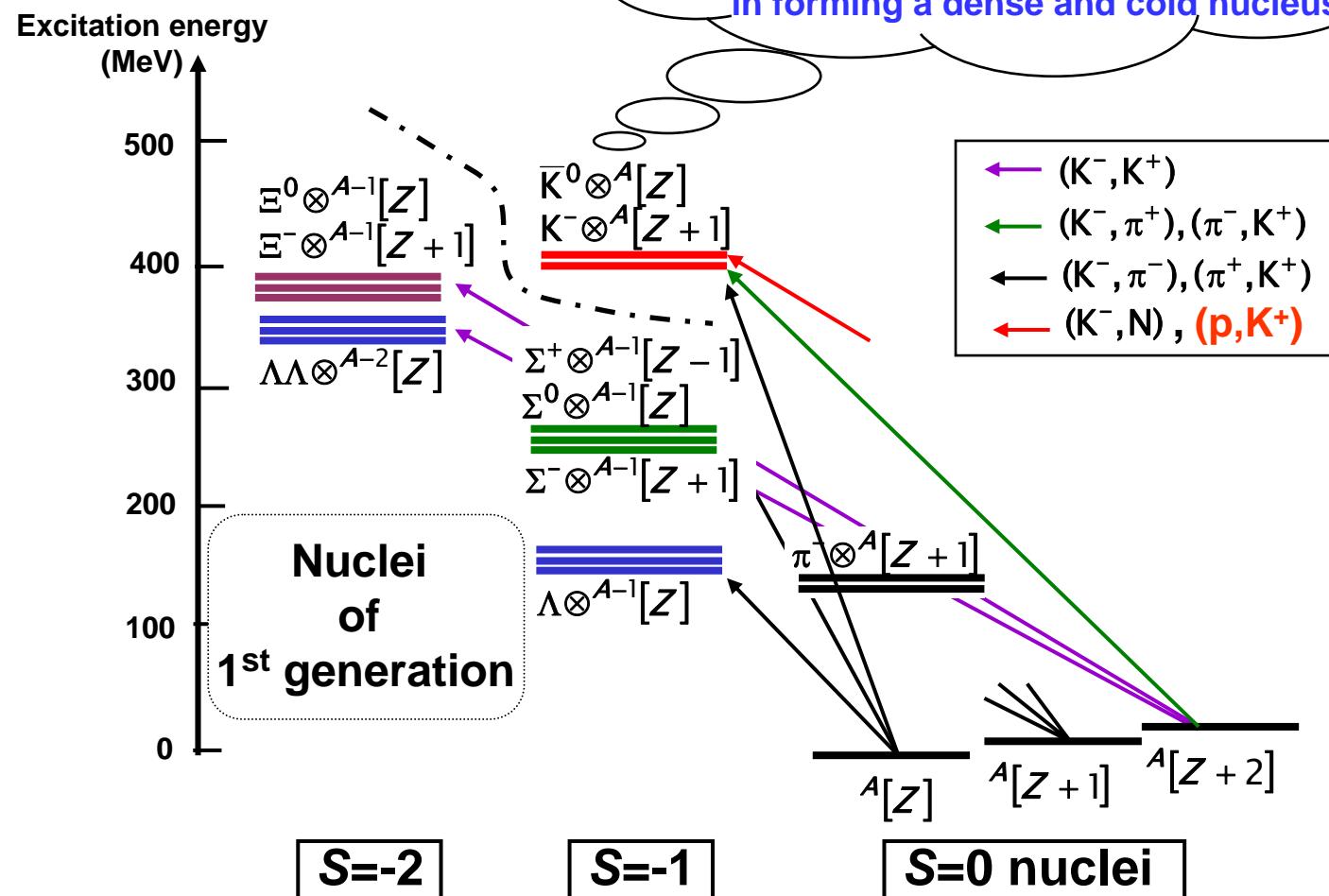
ATMS calculation

jade/HsHs.f jade/ATMS3Hs2.f jade/ATMS4Hs2.f

A new paradigm of nuclear physics

“Swan Nuclear Physics”

Yamazaki diagram

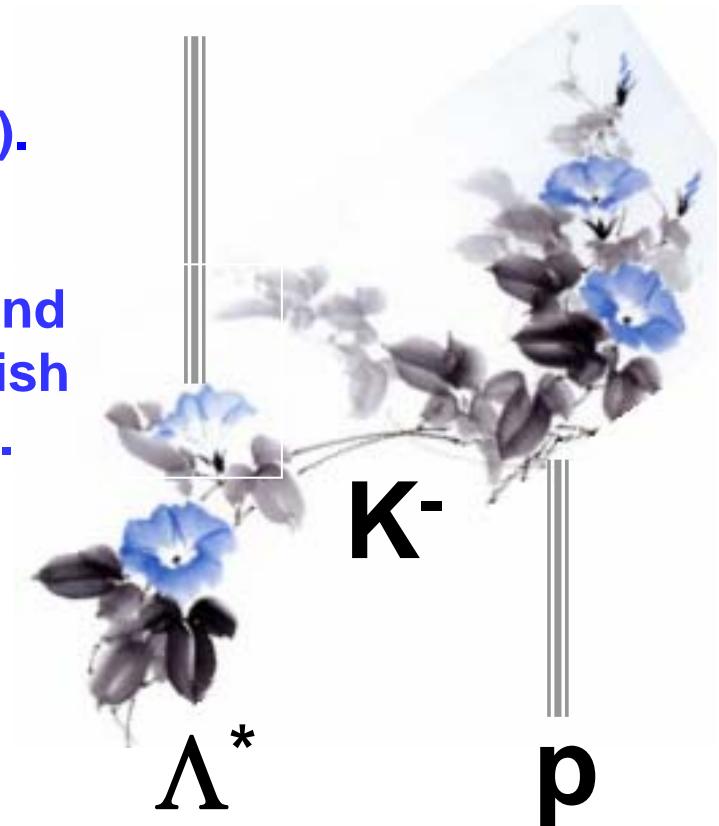


Remarks

The $\Lambda^*=\Lambda(1405)$ plays an essential role in forming " $K^{\bar{b}ar}$ Nuclear Clusters" (KNC).

Experimental information on $\Lambda^*p=K^-pp$ and $\Lambda^*\Lambda^*=K^-K^-pp$ is vitally important to establish the "super-strong nuclear force" (SSNF).

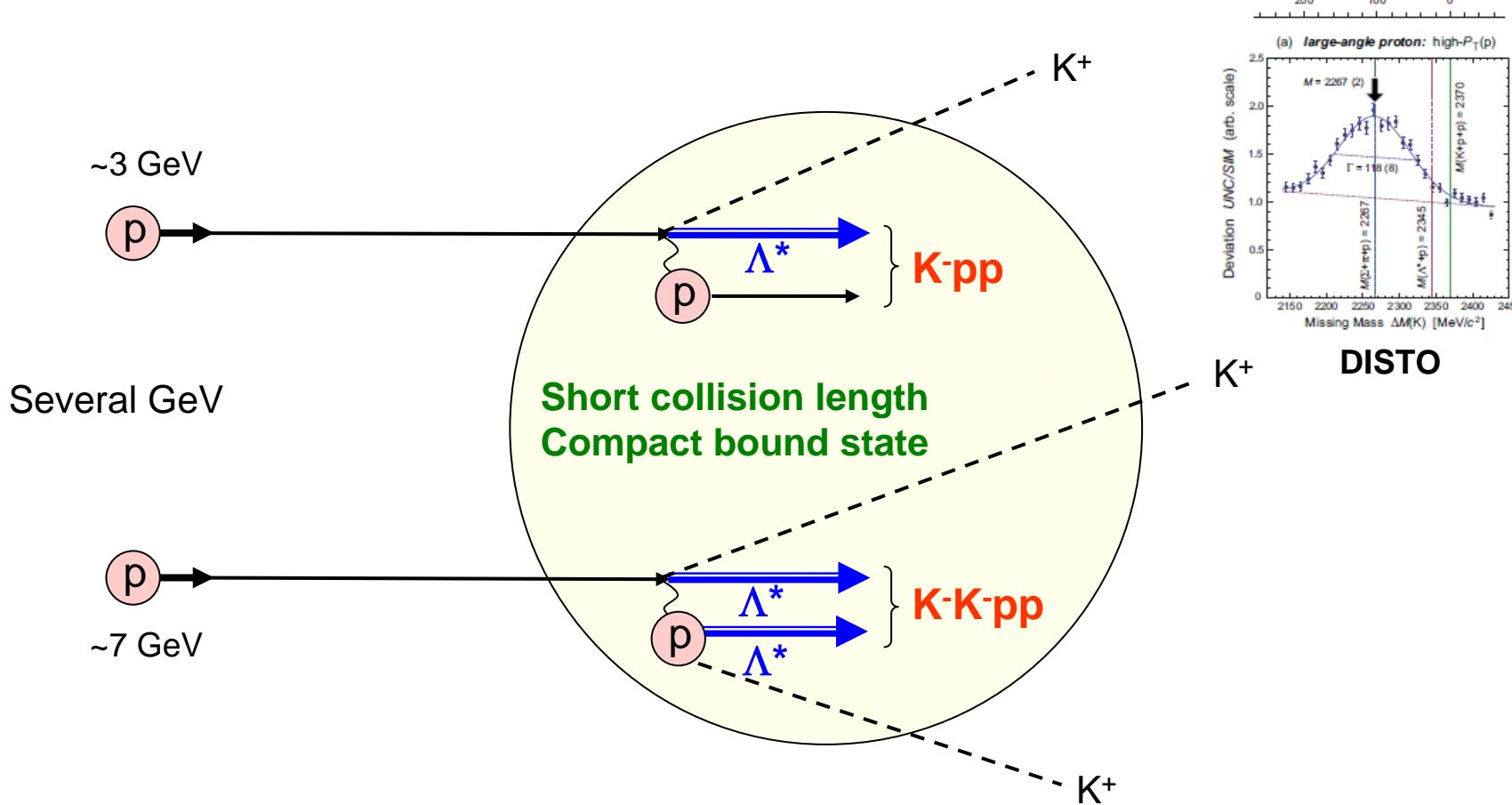
SSNF due to boson ($K^{\bar{b}ar}$) covalency predicts the possible existence of "finite Λ^* strangelets", which could be stable against some strong decays.



The late Prof. Nishijima

Gateway to "Swan Nuclear Physics"

T. Yamazaki, Y. Akaishi & M. Hassanvand, Proc. Jpn. Acad. B 87 (2011) 362



Thank you very much!