

Double-pole structure on a prototype of kaonic nuclei “ K^-pp ”



Akinobu Doté (KEK Theory Center, IPNS / J-PARC branch)

Takashi Inoue (Nihon university)

Takayuki Myo (Osaka Institute of Technology)

1. Introduction

2. Method

- *Feshbach projection of coupled-channel Complex Scaling Method*

3. Result

- *“ K^-pp ” as a $K^{\text{bar}}NN-\pi YN$ system with ccCSM+Feshbach method*

4. Double pole of “ K^-pp ”?

5. Summary and future plan

*12th International Conference on Hypernuclear and
Strange Particle Physics (HYP2015)*

*07. Sep. '15 @ Tohoku University Centennial Hall
(Kawauchi Hagi Hall), Sendai, Japan*

1. Introduction

1. Introduction

Kaonic nuclei = Exotic system !?

- Strong $K^{\text{bar}}N$ attraction $\leftarrow \Lambda(1405) \sim$ quasi-bound state of $l=0 K^{\text{bar}}N$
 - Deeply bound (Total B.E. $\sim 100\text{MeV}$)
 - Quasi-stable ($\pi\Sigma$ decay mode closed)
 - Highly dense state ... anti-kaon attracts nucleons

Y. Akaishi and T. Yamazaki, PRC65, 044005 (2002)

A. Dote, H. Horiuchi, Y. Akaishi and T. Yamazaki, PRC70, 044313 (2004)



“K⁻pp” = A prototype of kaonic nuclei

Experimental search for “K⁻pp”

- FINUDA : K stopped on Li, C, Al target
PRL 94, 212303 (2005)
- DISTO : $p + p \rightarrow p + \Lambda + K^+$ at $T_p=2.85\text{ GeV}$
PRL104, 132502 (2010)
- J-PARC E27 : $d(\pi^+, K^+) X$ at $P_{\pi}=1.69\text{ GeV}/c$
PTEP 021D01 (2015)

Signal at $\sim 100\text{ MeV}$ below K⁻pp threshold

- J-PARC E15 : ${}^3\text{He}(K^-, n) X$ at $P_K=1\text{ GeV}/c$
PTEP 2015, 061D01

Attraction below K⁻pp threshold

- LEPS/SPring8 : $d(\gamma, \pi^+K^+) X$ at $E_\gamma=1.5-2.4\text{ GeV}$
PLB 728, 616 (2014)

No evidence of deeply bound K⁻pp

Theoretical studies of “K⁻pp”

$B(K^{\text{-}}pp) < 100\text{ MeV}$

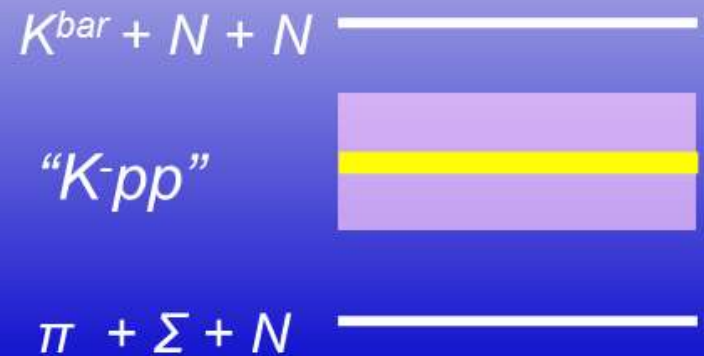
	Dote-Hyodo-Weise PRC79, 014003(2009)	Akaishi-Yamazaki PRC76, 045201(2007)	Barnea-Gal-Livertz PLB712, 132 (2012)	Ikeda-Kamano-Sato PTP124, 533 (2010)	Shevchenko-Gal- Mares PRC76, 044004(2007)
B(K ⁻ pp)	20 ± 3	47	16	9 ~ 16	50 ~ 70
Γ_{mesonic}	40 ~ 70	61	41	34 ~ 46	90 ~ 110

- $\Lambda(1405) = \text{Resonant state \& } K^{\text{bar}}N \text{ coupled with } \pi\Sigma$

- “K-pp” ... Resonant state of
 $K^{\text{bar}}NN\text{-}\pi YN$ coupled-channel system

*Doté, Hyodo, Weise, PRC79, 014003(2009). Akaishi, Yamazaki, PRC76, 045201(2007)
Ikeda, Sato, PRC76, 035203(2007). Shevchenko, Gal, Mares, PRC76, 044004(2007)
Barnea, Gal, Liverts, PLB712, 132(2012)*

- Resonant state
- Coupled-channel system



⇒ “coupled-channel
Complex Scaling Method”

Complex Scaling Method

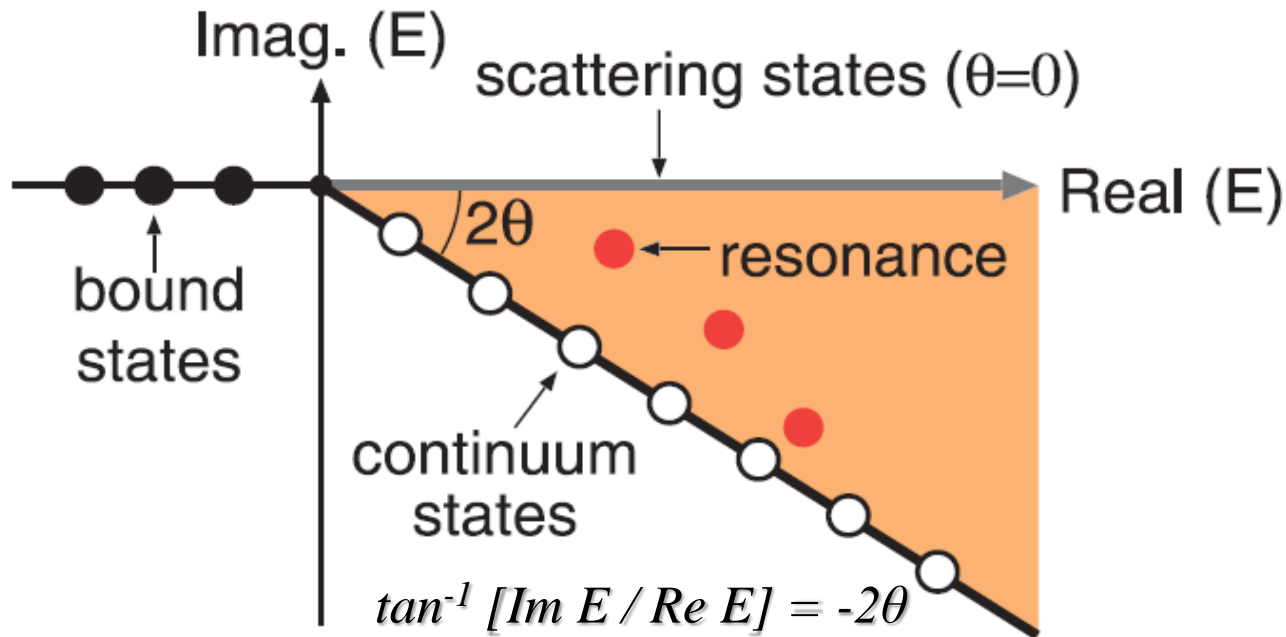
... Powerful tool for resonance study of many-body system

Complex rotation (Complex scaling) of coordinate

Resonance wave function $\rightarrow L^2$ integrable

$$U(\theta): \mathbf{r} \rightarrow \mathbf{r} e^{i\theta}, \quad \mathbf{k} \rightarrow \mathbf{k} e^{-i\theta}$$

Diagonalize $H_\theta = U(\theta) H U^{-1}(\theta)$ with Gaussian base,



- Continuum state appears on 2ϑ line.
- Resonance pole is off from 2ϑ line, and independent of ϑ . (ABC theorem)

Chiral SU(3) potential with a Gaussian form

A. D., T. Inoue, T. Myo, Nucl. Phys. A 912, 66 (2013)

• **Anti-kaon = Nambu-Goldstone boson**

⇒ Chiral SU(3)-based $K^{\text{bar}}N$ potential

- Weinberg-Tomozawa term of effective chiral Lagrangian
- Gaussian form in r -space
- Semi-rela. / Non-rela.
- Based on Chiral SU(3) theory
→ **Energy dependence**

A non-relativistic potential (NRv2c)

$$V_{ij}^{(I=0,1)}(r) = -\frac{C_{ij}^{(I=0,1)}}{8f_{\pi}^2} (\omega_i + \omega_j) \sqrt{\frac{1}{m_i m_j}} g_{ij}(r)$$

$$g_{ij}(r) = \frac{1}{\pi^{3/2} d_{ij}^3} \exp\left[-(r/d_{ij})^2\right] : \text{Gaussian form}$$

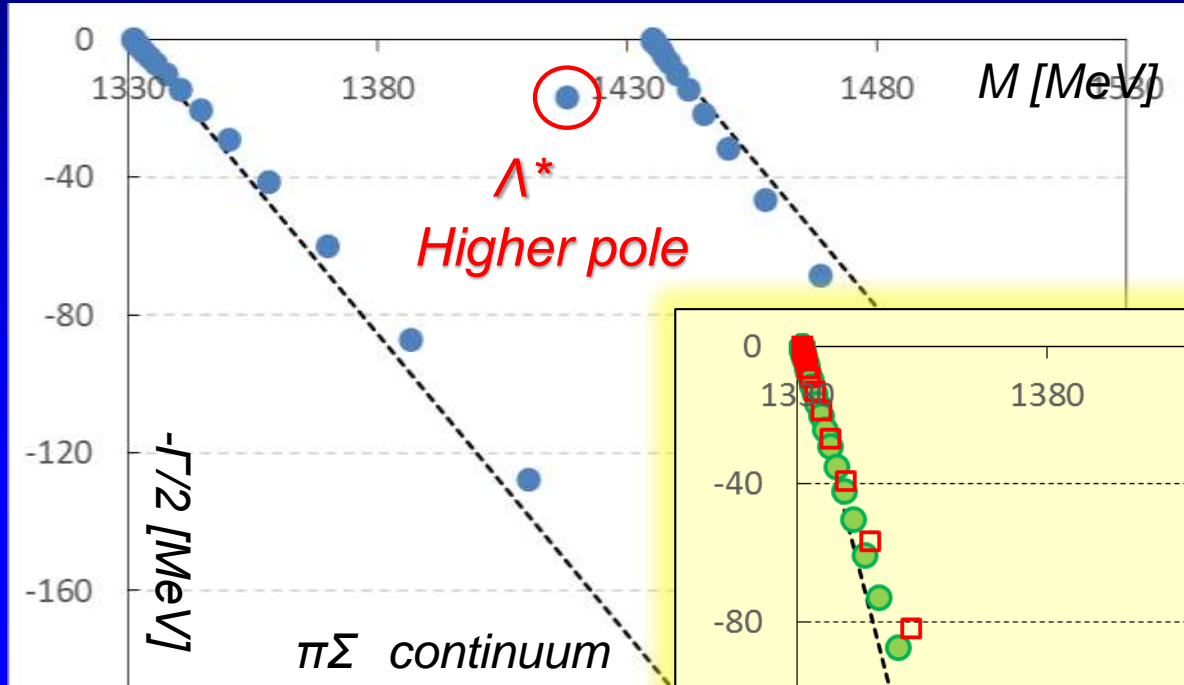
ω_i : meson energy

Constrained by $K^{\text{bar}}N$ scattering length

$$a_{KN(I=0)} = -1.70 + i0.67 \text{fm}, \quad a_{KN(I=1)} = 0.37 + i0.60 \text{fm}$$

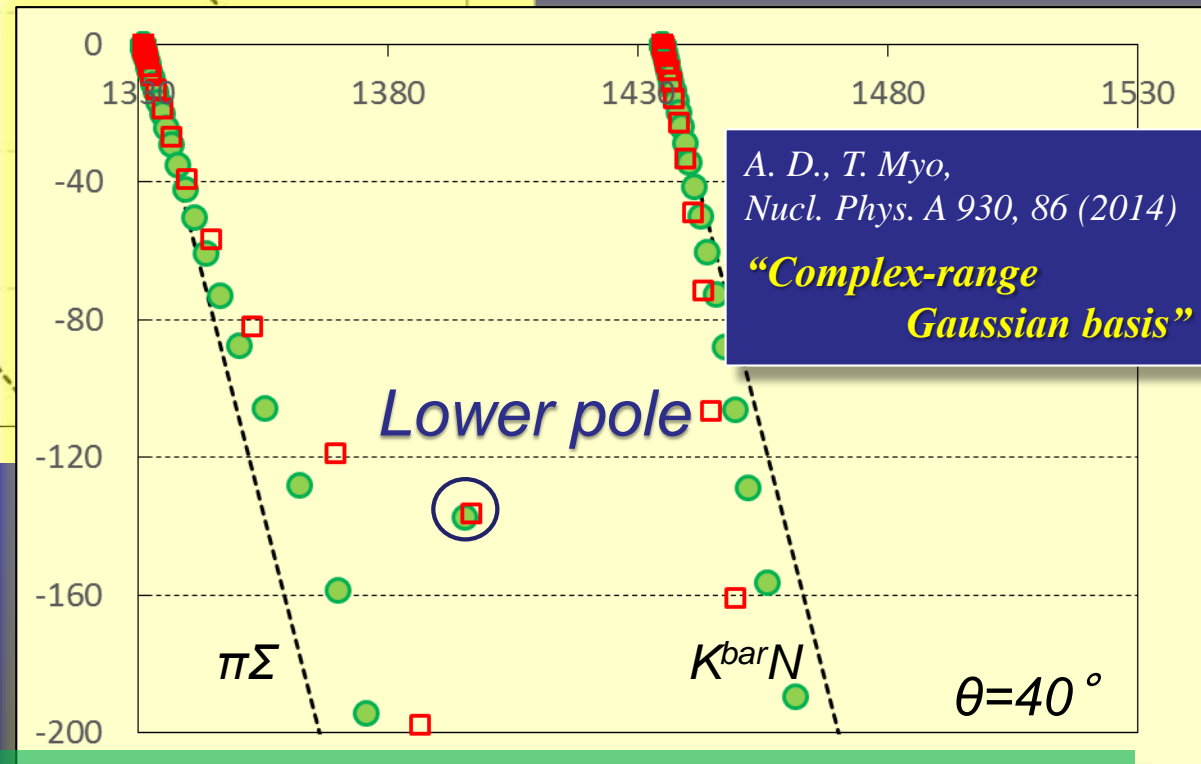
A. D. Martin, NPB179, 33(1979)

$\Lambda(1405)$ on coupled-channel Complex Scaling Method



$K^{bar}N$ potential:
a chiral SU(3) potential
(NRv2, $f_\pi=110$)

A. D., T. Inoue, T. Myo,
Nucl. Phys. A 912, 66 (2013)



A. D., T. Myo,
Nucl. Phys. A 930, 86 (2014)
"Complex-range
Gaussian basis"

Double-pole structure of $\Lambda(1405)$

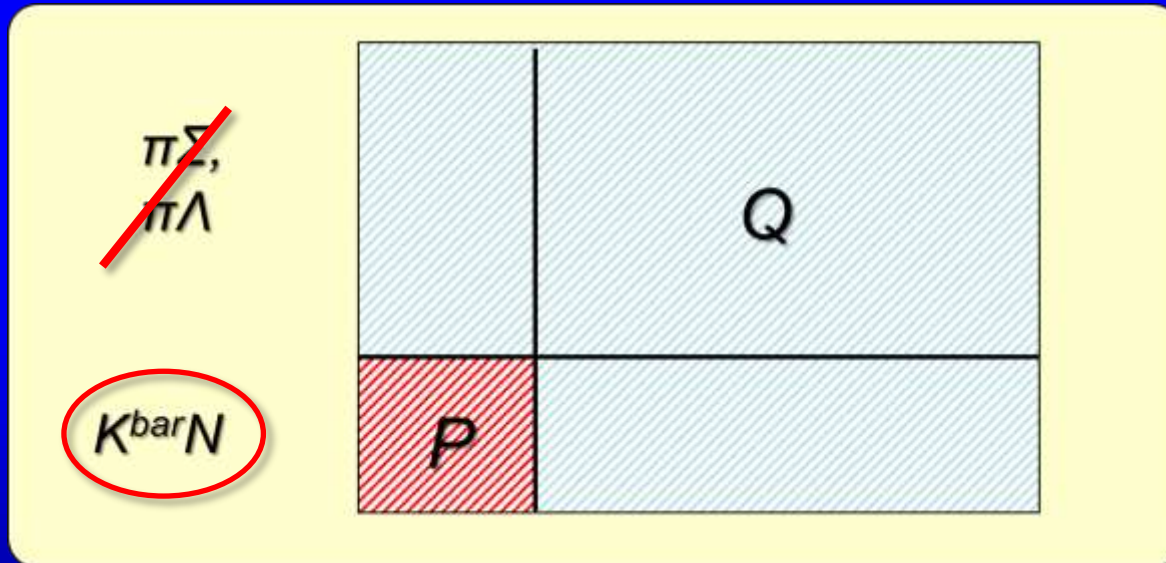
2. Method

- ***Feshbach projection on
coupled-channel Complex Scaling Method
“ccCSM+Feshbach method”***

***A. D., T. Inoue, T. Myo,
PTEP 2015, 043D02 (2015)***

ccCSM+Feshbach method

- $\Lambda(1405) =$ *two-body* system of $K^{\text{bar}}N-\pi\Sigma$
→ Explicitly treat coupled-channel problem
- “ K^-pp ” = *three-body* system of $K^{\text{bar}}NN-\pi YN$
... High computational cost



For economical treatment of “ K^-pp ”, we construct an effective $K^{\text{bar}}N$ single-channel potential by means of *Feshbach projection* on CSM.

Formalism of ccCSM + Feshbach method

Elimination of channels by Feshbach method

Schrödinger eq.
in model space "P" and out of model space "Q"

$$\begin{pmatrix} T_P + v_P & V_{PQ} \\ V_{QP} & T_Q + v_Q \end{pmatrix} \begin{pmatrix} \Phi_P \\ \Phi_Q \end{pmatrix} = E \begin{pmatrix} \Phi_P \\ \Phi_Q \end{pmatrix}$$

Schrödinger eq. in P-space : $(T_P + U_P^{Eff}(E))\Phi_P = E\Phi_P$

Effective potential for P-space

$$U_P^{Eff}(E) = v_P + V_{PQ} G_Q(E) V_{QP}$$

Q-space Green function:

$$G_Q(E) = \frac{1}{E - H_{QQ}}$$

Extended Closure Relation in Complex Scaling Method

$$H_{QQ}^\theta |\chi_n^\theta\rangle = \varepsilon_n^\theta |\chi_n^\theta\rangle$$

$$H_{QQ}^\theta = U(\theta) H_{QQ} U^{-1}(\theta)$$

$$\int_C \sum_{R+B} |\chi_n^\theta\rangle \langle \chi_n^\theta| = 1$$



Diagonalize H_{QQ}^θ with **Gaussian base**,

$$\sum_n |\chi_n^\theta\rangle \langle \chi_n^\theta| \approx 1 \quad \text{Well approximated}$$

T. Myo, A. Ohnishi and K. Kato, PTP 99, 801 (1998)
R. Suzuki, T. Myo and K. Kato, PTP 113, 1273 (2005)

Express the $G_Q(E)$ with **Gaussian base** using ECR

$$G_Q^\theta(E) = \frac{1}{E - H_{QQ}^\theta} \approx \sum_n |\chi_n^\theta\rangle \frac{1}{E - \varepsilon_n^\theta} \langle \chi_n^\theta|$$



$$U_P^{Eff}(E) = v_P + V_{PQ} \underbrace{U^{-1}(\theta) G_Q^\theta(E) U(\theta)}_{G_Q(E)} V_{QP}$$

$\{ |\chi_n^\theta\rangle \}$: expanded with Gaussian base.

Apply ccCSM + Feshbach method to K^-pp

“ K^-pp ” ... $K^{bar}NN - \pi\Sigma N - \pi\Lambda N$ ($J^\pi=0^-, T=1/2$)

For the two-body system, $P = K^{bar}N$, $Q = \pi Y$

$$\begin{matrix} V(K^{bar}N - \pi Y; I=0,1) \\ V(\pi Y - \pi Y'; I=0,1) \end{matrix} \xrightarrow{\text{Feshbach + ccCSM}} U_{K^{bar}N(I=0,1)}^{Eff}(E)$$

- Schrödinger eq. for $K^{bar}NN$ channel :

$$\left(T_{K^{bar}NN} + V_{NN} + \sum_{i=1,2} U_{K^{bar}N_i(I)}^{Eff}(E_{K^{bar}N}) \right) \Phi_{K^{bar}NN} = E \Phi_{K^{bar}NN}$$

- Trial wave function

$$\begin{aligned} |"K^-pp"> &= \sum_a C_a^{(KNN,1)} \left\{ G_a^{(KNN,1)}(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) + G_a^{(KNN,1)}(-\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) \right\} |S_{NN}=0> \left[[K[NN]_1]_{T=1/2} \right] \\ &+ \sum_a C_a^{(KNN,2)} \left\{ G_a^{(KNN,2)}(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) - G_a^{(KNN,2)}(-\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) \right\} |S_{NN}=0> \left[[K[NN]_0]_{T=1/2} \right] \end{aligned} \quad \begin{matrix} \text{Ch. 1: } K^{bar}NN, \quad NN:1E \\ \text{Ch. 2: } K^{bar}NN, \quad NN:1O \end{matrix}$$

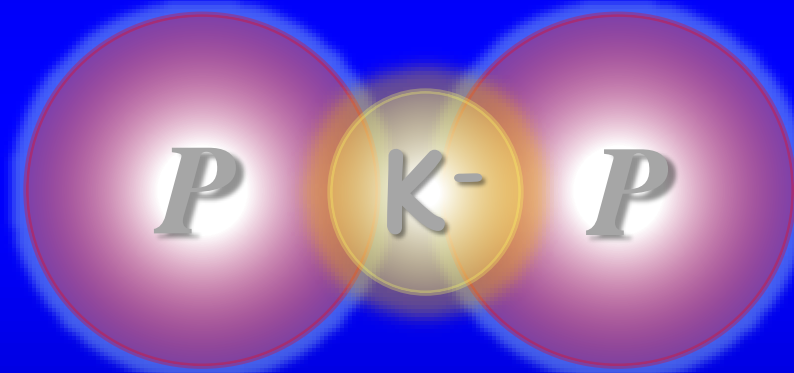
- Basis function = Correlated Gaussian
...including 3-types Jacobi-coordinates

$$G_a^{(KNN,i)}(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) = N_a^{(KNN,i)} \exp \left[-(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) A_a^{(KNN,i)} \begin{pmatrix} \mathbf{x}_1^{(3)} \\ \mathbf{x}_2^{(3)} \end{pmatrix} \right]$$

3. Result

Three-body “K⁻pp” resonance

on ccCSM+Feshbach projection

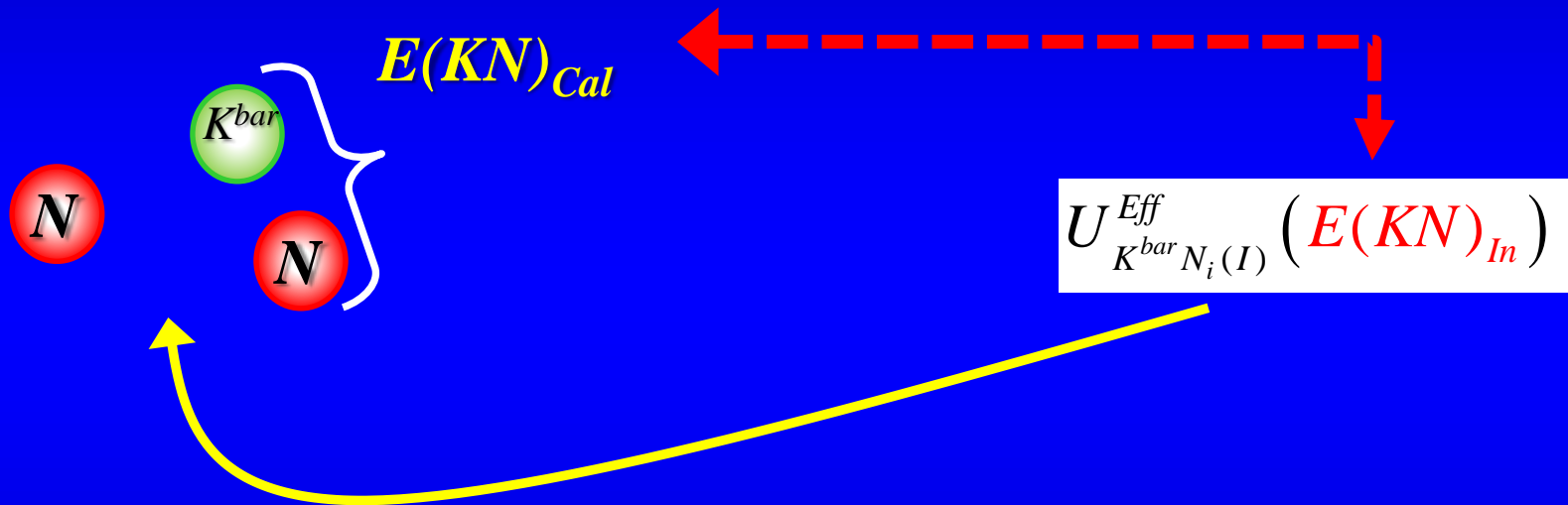


“K⁻pp” =

$K^{\text{bar}}NN - \pi\Sigma N - \pi\Lambda N$ ($J^\pi = 0^-, T=1/2$)

Self-consistency for **complex** $K^{\text{bar}}N$ energy

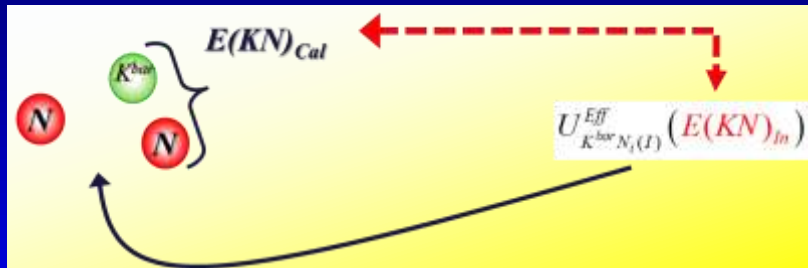
Effective $K^{\text{bar}}N$ potential has energy dependence...



- $E(KN)_{\text{In}}$: assumed in the $K^{\text{bar}}N$ potential
- $E(KN)_{\text{Cal}}$: calculated with the obtained K - pp

When $E(KN)_{\text{In}} = E(KN)_{\text{Cal}}$,
a self-consistent solution is obtained.

Self-consistency for **complex** K^{bar} N energy



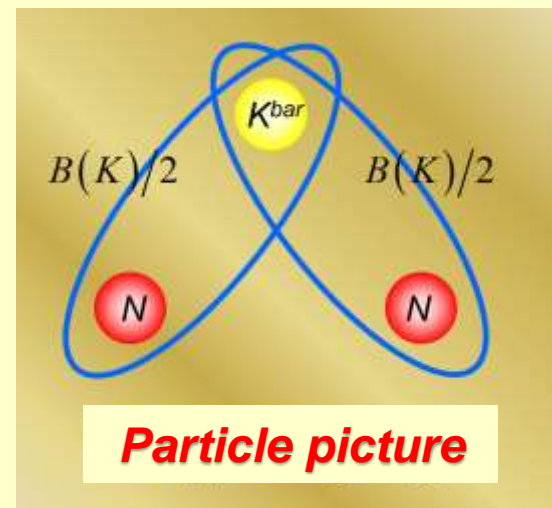
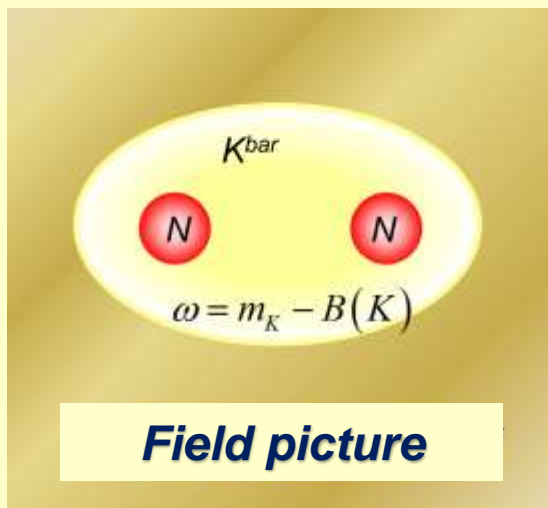
How to determine the two-body energy in the three-body system?

A. D., T. Hyodo, W. Weise,
PRC79, 014003 (2009)

1. Kaon's binding energy: $B(K) \equiv -\left\{ \langle H \rangle - \langle H_{NN} \rangle \right\}$ H_{NN} : Hamiltonian of two nucleons

2. Define a K^{bar} N-bond energy in two ways

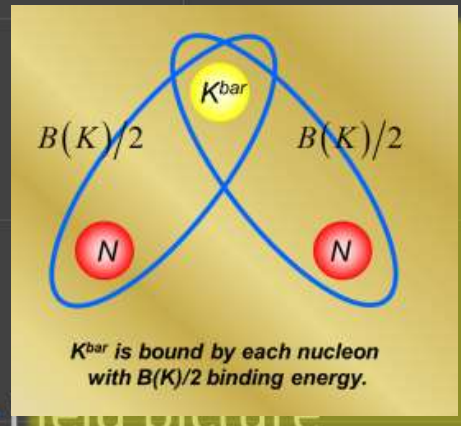
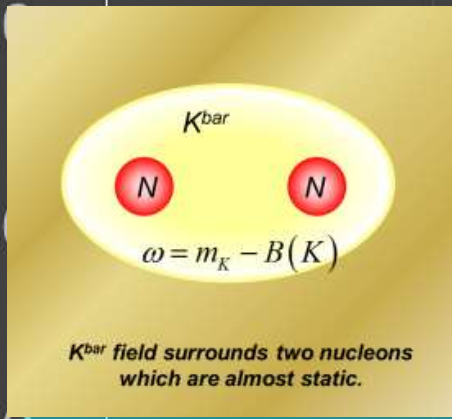
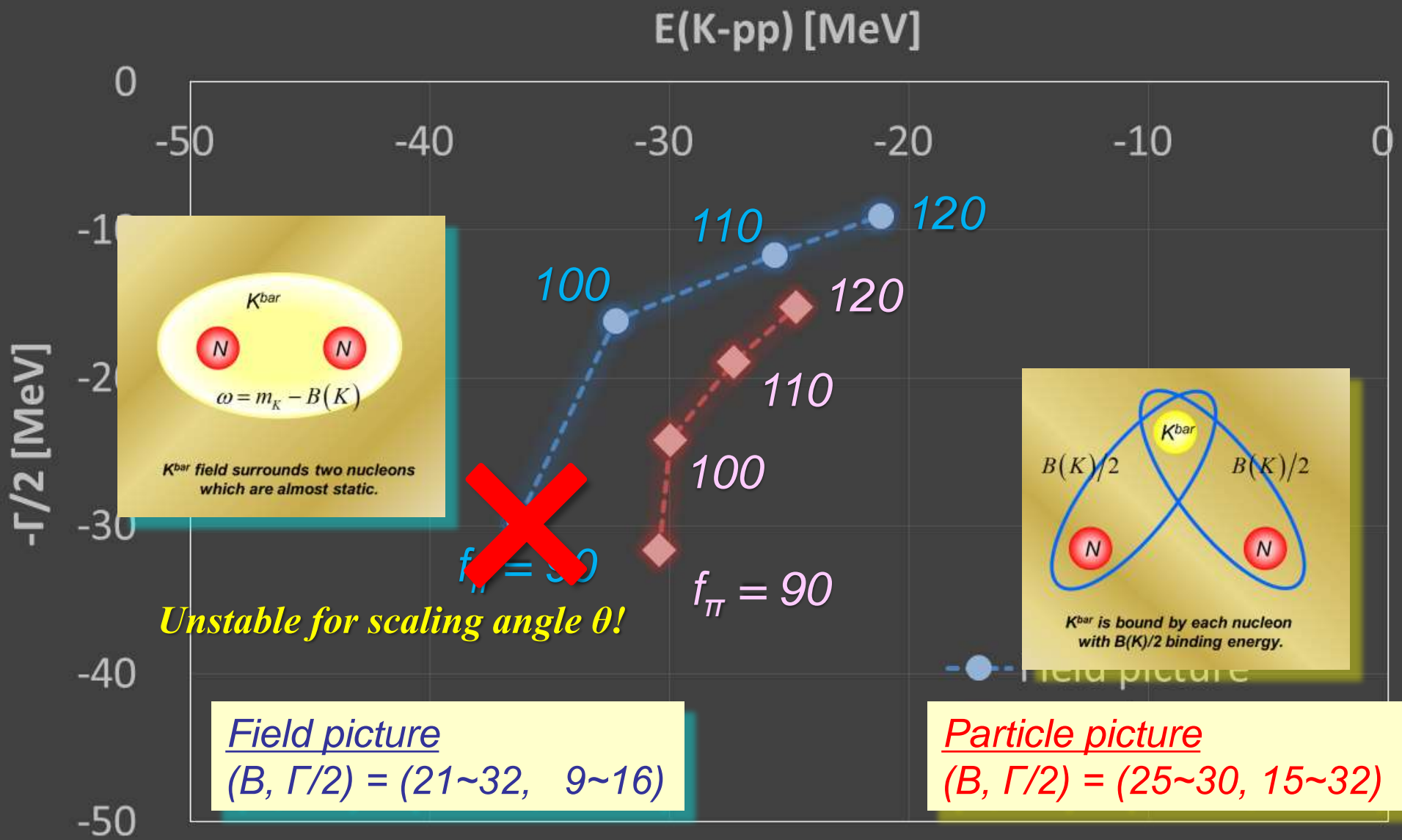
$$E_{KN} = M_N + \omega = \begin{cases} M_N + m_K - B(K) & : \text{Field picture} \\ M_N + m_K - B(K)/2 & : \text{Particle picture} \end{cases}$$



Self-consistent results

$f_\pi = 90 \sim 120 \text{ MeV}$

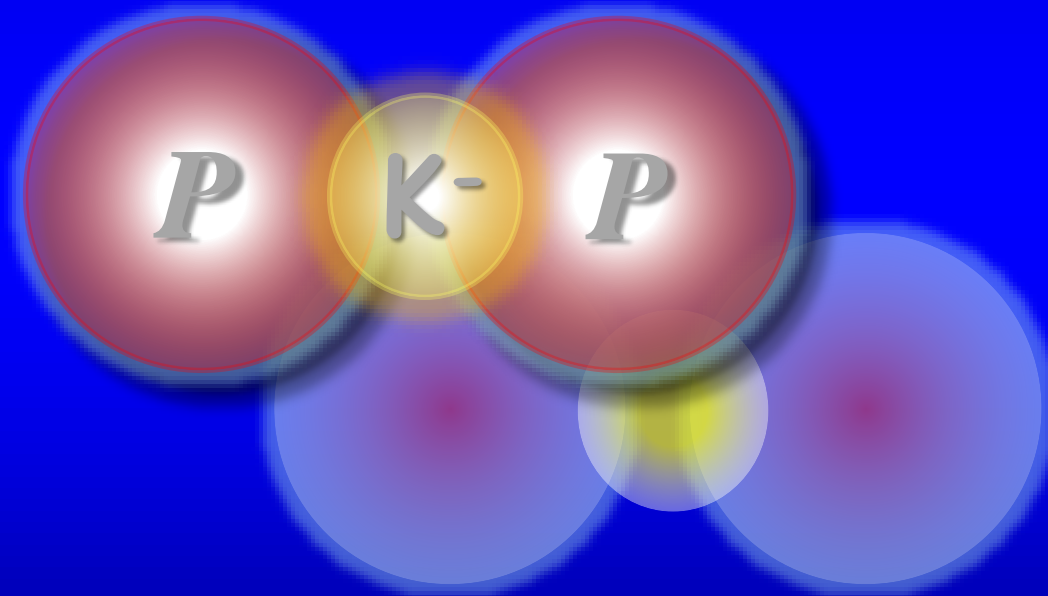
NN pot. : Av18 (Central)
 K^{bar} N pot. : NRv2c potential
 ($f_\pi = 90 - 120 \text{ MeV}$)



Field picture
 $(B, \Gamma/2) = (21 \sim 32, 9 \sim 16)$

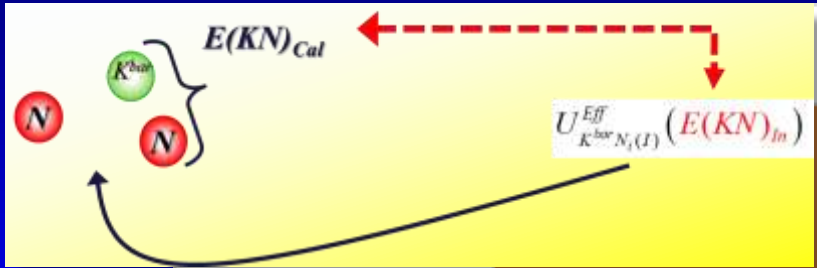
Particle picture
 $(B, \Gamma/2) = (25 \sim 30, 15 \sim 32)$

4. Double pole of “ K^-pp ”?

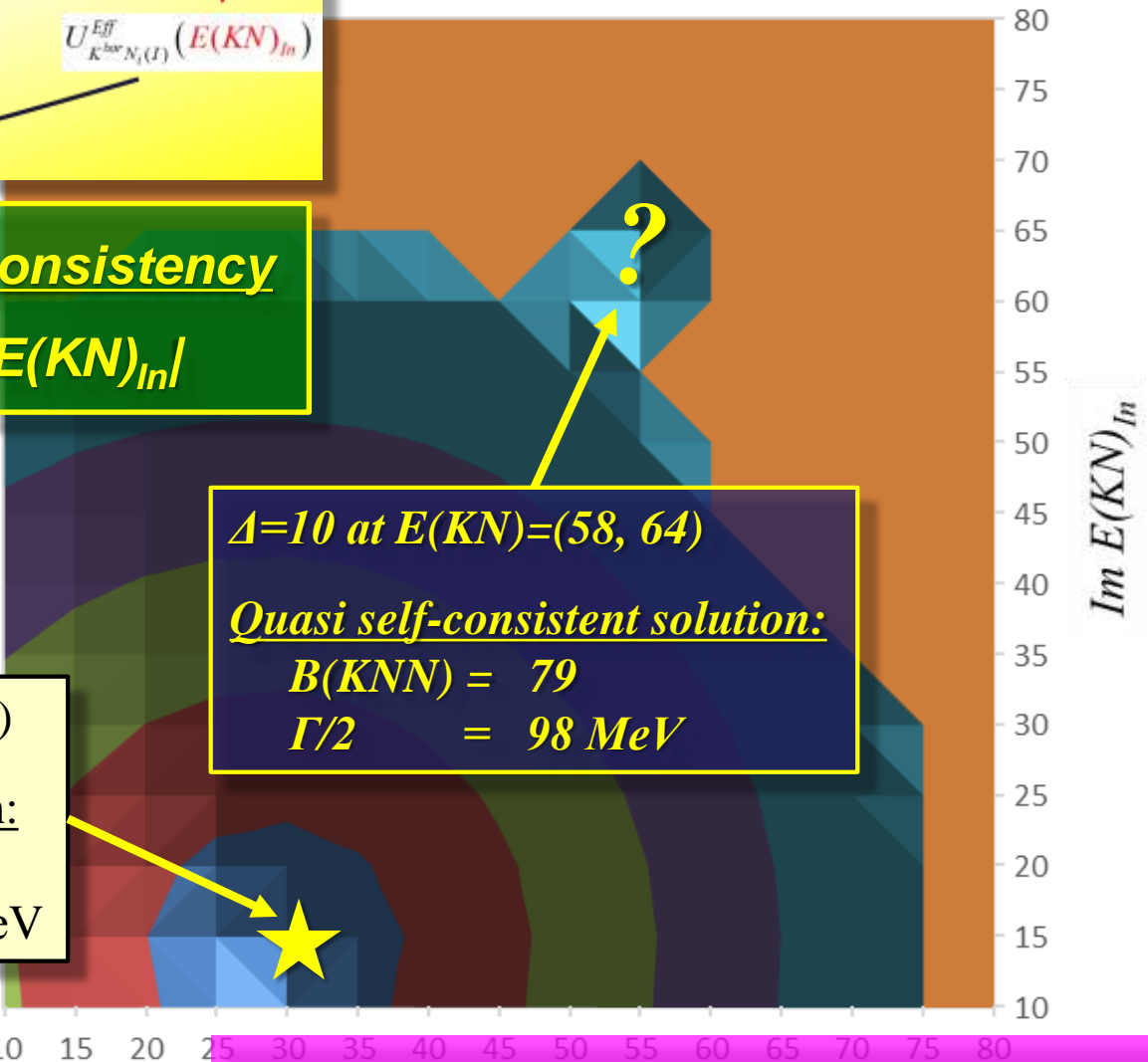


Quasi self-consistent solution

NRv2c ($f_\pi = 110$ MeV)
Particle picture



Indicator of self-consistency
 $\Delta = |E(KN)_{Cal} - E(KN)_{In}|$



$\Delta = 10$ at $E(KN) = (58, 64)$
Quasi self-consistent solution:
 $B(KNN) = 79$
 $\Gamma/2 = 98$ MeV

$\Delta = 0$ at $E(KN) = (29, 14)$
Self-consistent solution:
 $B(KNN) = 27.3$
 $\Gamma/2 = 18.9$ MeV

“Double pole of K^-pp ” ?

Double-pole structure in “K-pp”?

- ✓ Quasi self-consistent solution is obtained ...
 $(B(KNN), \Gamma/2) = (62 \sim 79, 74 \sim 104) \text{ MeV}$ for $f_\pi=90 \sim 120 \text{ MeV}$
with Particle picture
- ✓ Such solutions are not obtained with Field picture.

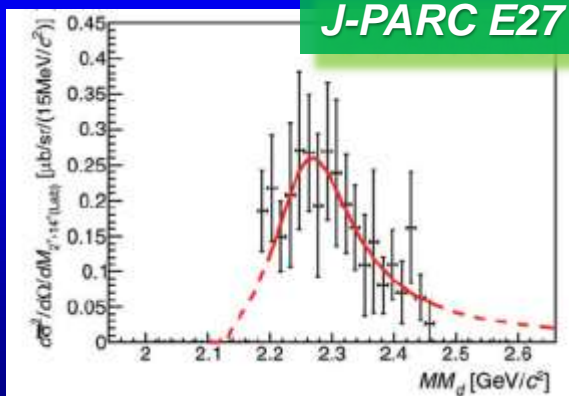
- A Faddeev-AGS calc. has predicted the double-pole structure of “K-pp”.

Lower pole : $(B(KNN), \Gamma/2) = (67 \sim 89, 122 \sim 160) \text{ MeV}$

Higher pole : $(B(KNN), \Gamma/2) = (9 \sim 16, 17 \sim 23) \text{ MeV}$

Y. Ikeda, H. Kamano, and T. Sato, PTP 124, 533 (2010)

- Relation to signals observed by J-PARC E27, DISTO?



Lower pole of “K-pp” ($J^\pi=0^-, l=1/2$)
... “K-pp” has two poles similarly to $\Lambda(1405)$.
The lower pole appears.

Partial restoration of chiral symmetry

... $K^{\text{bar}}N$ potential is enhanced by 17%.

S. Maeda, Y. Akaishi, T. Yamazaki, Proc. Jpn. Acad., Ser. B 89, 418 (2013)

Pion assisted dibaryon “ $\Upsilon = \pi\Sigma N - \pi\Lambda N$ ($J^\pi=2^+, l=3/2$)”

A. Gal, arXiv:1412.0198 (Proceeding of EXA2014)

Signal at $\sim 100 \text{ MeV}$ below $K^{\text{bar}}NN$ thr.

5. Summary

and future plans

5. Summary

A prototype of K^{bar} nuclei “ $K\text{-}pp$ ” = Resonance state of $K^{\text{bar}}NN\text{-}\pi YN$ coupled system

“coupled-channel Complex Scaling Method + Feshbach projection”

... Represent the **Q-space Green function** with the **Extended Complete Set** well approximated by **Gaussian base**

⇒ Eliminate πY channels to reduce the problem to a $K^{\text{bar}}NN$ single channel problem.

$K\text{-}pp$ studied with ccCSM+Feshbch method

- Used a Chiral $SU(3)$ -based potential (Gaussian form in r -space)
- Self-consistency for $K^{\text{bar}}N$ **complex** energy (Field and Particle pictures)

$$\underline{K\text{-}pp (J^\pi=0^-, T=1/2) \dots (B, \Gamma/2) \doteq (20\sim 30, 10\sim 30) \text{ MeV}}$$

- Quasi self-consistent solution in case of Particle picture
... Deeper binding and larger decay width

$$\underline{K\text{-}pp (J^\pi=0^-, T=1/2) \dots (B, \Gamma/2) \doteq (60\sim 80, 75\sim 105) \text{ MeV}}$$

“ $K\text{-}pp$ ” has a double-pole structure similarly to $\Lambda(1405)$?

- Relation to the $K\text{-}pp$ search experiments

*The signal observed in J-PARC E27 is considered to correspond to the lower pole of “ $K\text{-}pp$ ”??
J-PARC E15 may pick up the higher pole of “ $K\text{-}pp$ ”???*

5. Future plans

- Full-coupled channel calculation of $K\bar{p}p$
... Detailed study for the double pole structure of $K\bar{p}p$
- Application to resonances of other hadronic systems



Thank you for your attention!

References:

1. A. D., T. Inoue, T. Myo,
NPA 912, 66 (2013)
2. A. D., T. Myo, NPA 930, 86 (2014)
3. A. D., T. Inoue, T. Myo,
PTEP 2015, 043D02 (2015)

Cats in KEK