

Photoproduction of typical hypernuclei

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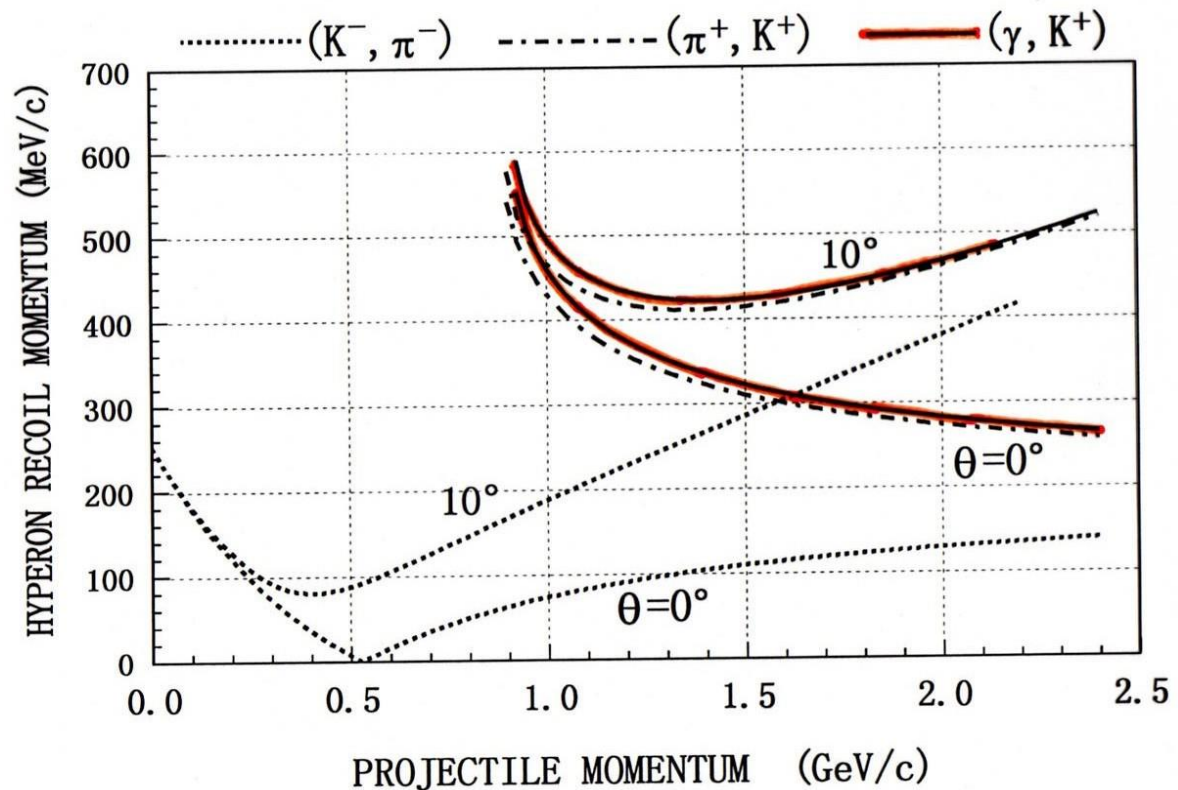
- (0) Properties of hypernuclear photoproduction
- (1) Characteristics utilized in ${}^4\text{He}(\gamma, \text{K}^+)$
- (2) p-shell high-resolution $(e, e' \text{K}^+)$ experiments and a new theoretical challenge
- (3) sd- and fp-shell hypernuclei as missing link of high resolution s.p.e.
- (4) Hyperon s.p.e $({}^{208}\text{Pb}(\gamma, \text{K}^+)_{\Lambda} {}^{208}\text{Al})$
- (5) Λ -rotation(deformation) coupling
- (6) Summary

(0) **BASICS:** Hyperon **recoil momentum** and the **transition operator itself** determine the reaction characteristics

(1)



Momentum transfers are both large and comparable.



$q_\Lambda = 350-420$ MeV/c at $E_\gamma = 1.3$ GeV

Microscopic treatment based on the elementary transition amplitudes (π, K) case

$$\frac{d\sigma(\theta_L)}{d\Omega_L} = \gamma \cdot \frac{(2\pi)^4 p_K^2 E_\pi E_K E_H}{p_\pi \{p_K(E_H + E_K) - p_\pi E_K \cos\theta_L\}} \overline{|T_{if}^L|^2},$$

$$|T_{if}^L|^2 = \sum_{M_f} R(if; M_f),$$

$$R(if; M_f) = \frac{1}{[J_i]} \sum_{M_i} \left| \langle J_f M_f | \int d^3r \chi^{(-)}(p_K; r)^* \cdot \chi^{(+)}(p_\pi; r) \right. \\ \left. \times \sum_{k=1}^A U_-(k) \delta(r - r_k) \cdot \lambda [f + ig(\sigma_k \cdot \hat{n})] |J_i M_i\rangle \right|^2,$$

(2) Elementary amplitude $N \rightarrow Y$ (π, K) case

$f =$ spin-nonflip, $g =$ spin-flip, $\sigma =$ baryon spin

Lab $d\sigma/d\Omega$ photoproduction case (2Lab)

$$\frac{d\sigma}{d\Omega}\Big|_{2\text{Lab}} = \frac{(2\pi)^4 p^2 E_K E_\gamma E_\Lambda}{k\{p(E_\Lambda + E_K) - kE_K \cos\theta_L\}} \left| \langle \mathbf{k} - \mathbf{p}, \mathbf{p} | t | \mathbf{k}, 0 \rangle_L \right|^2, \quad (2.4)$$

$$\langle \mathbf{k} - \mathbf{p}, \mathbf{p} | t | \mathbf{k}, 0 \rangle_L = a_1(\boldsymbol{\sigma} \cdot \boldsymbol{\epsilon}) + a_2(\boldsymbol{\sigma} \cdot \hat{\mathbf{k}})(\hat{\mathbf{p}} \cdot \boldsymbol{\epsilon}) + a_3(\boldsymbol{\sigma} \cdot \hat{\mathbf{p}})(\hat{\mathbf{p}} \cdot \boldsymbol{\epsilon}) + a_4\{(\hat{\mathbf{k}} \times \hat{\mathbf{p}}) \cdot \boldsymbol{\epsilon}\}. \quad (2.5)$$

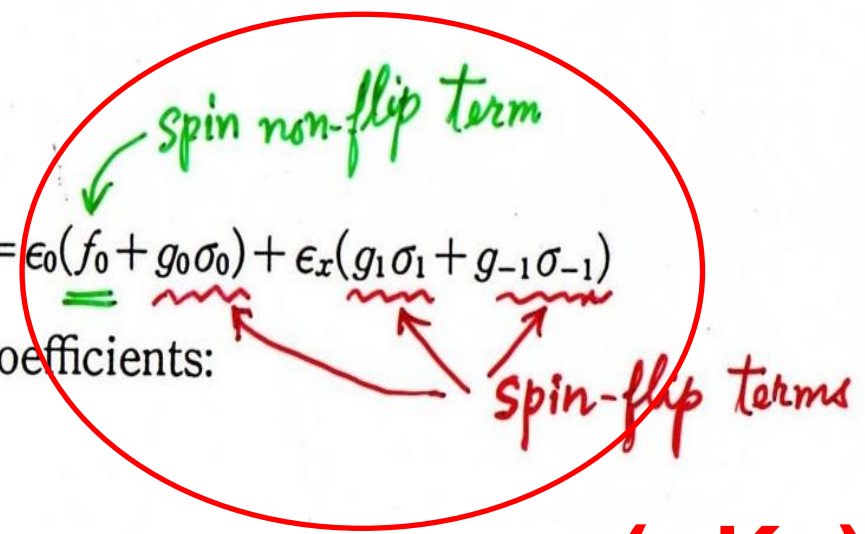
$$\langle \mathbf{k} - \mathbf{p}, \mathbf{p} | t | \mathbf{k}, 0 \rangle_L = \epsilon_0(f_0 + g_0\sigma_0) + \epsilon_x(g_1\sigma_1 + g_{-1}\sigma_{-1}) \quad (2.11)$$

with definitions of the coefficients:

$$f_0 = a_4 \sin\theta_L,$$

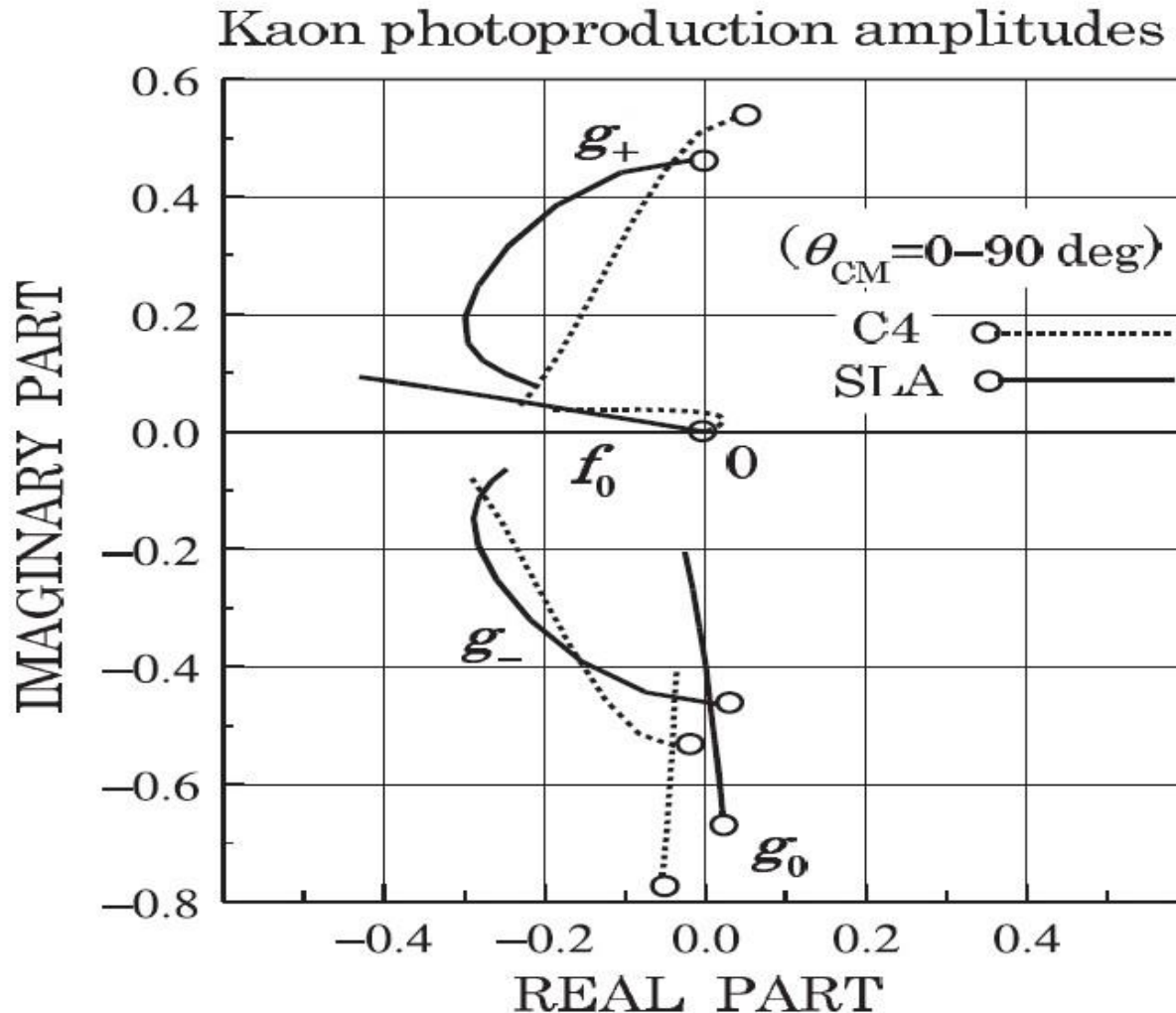
$$g_0 = a_1,$$

$$g_{\pm 1} = \frac{1}{\sqrt{2}} \{ \mp (a_1 + a_3 \sin^2\theta_L) - i \sin\theta_L (a_2 + a_3 \cos\theta_L) \}. \quad (2.12)$$



(γ, K^+) case

Elementary amplitudes (complex and p-dependent, θ -dependent)



Three spin-flip terms are all large in Kaon photoproduction

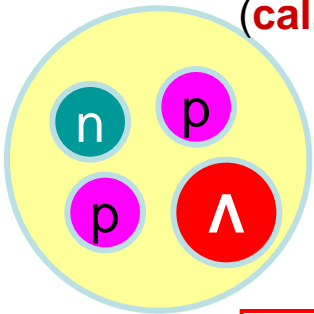
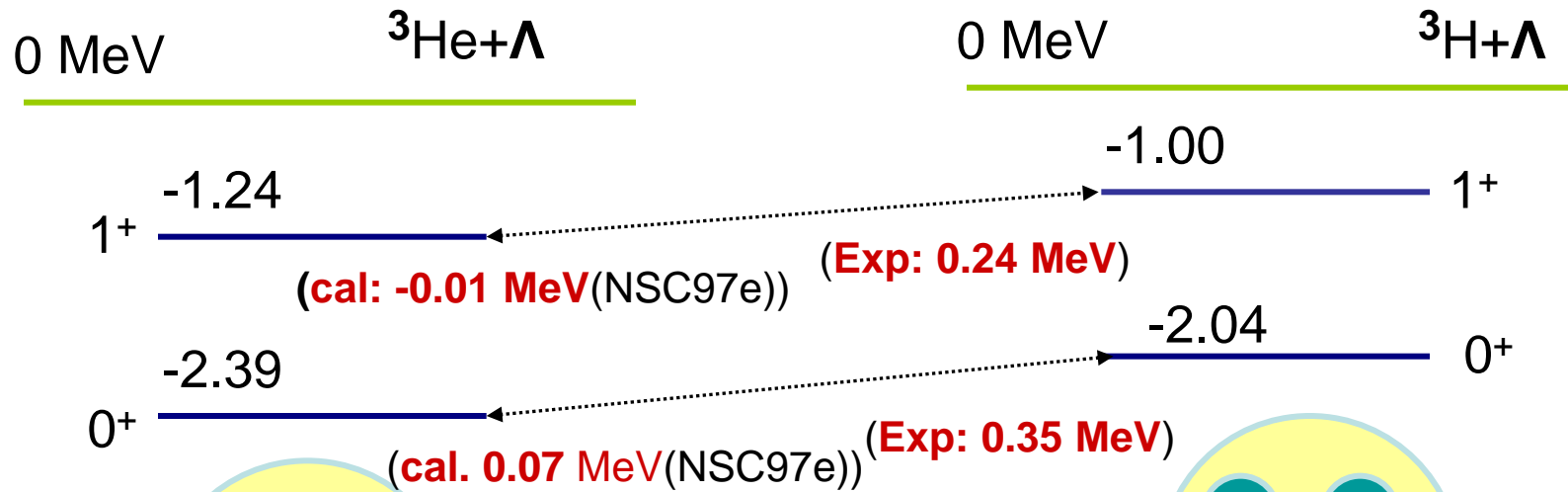
(1) Application to the lightest closed-shell target ${}^4\text{He}$
(A proposal from theory side)

Unique role of $(e,e'K^+)$ or (γ,K^+) reaction:
to excite ${}^4_{\Lambda}\text{H}(1+)$ state preferentially by making use of the spin-flip dominant nature.

(An important issue is to determine $1+$ energy position (update) for the study of CSB effect in Λ -N interaction.)

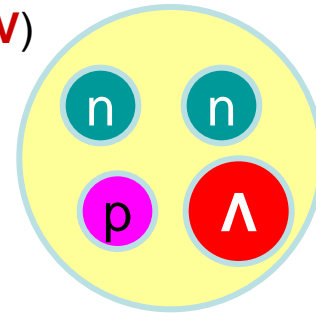
Many 4-body calculations have been done → Talks by Hiyama, Gal

Hiyama,
Gal



$^4_{\Lambda}\text{He}$

• A. Nogga, H. Kamada and W. Gloeckle,
Phys. Rev. Lett. 88, 172501 (2002)

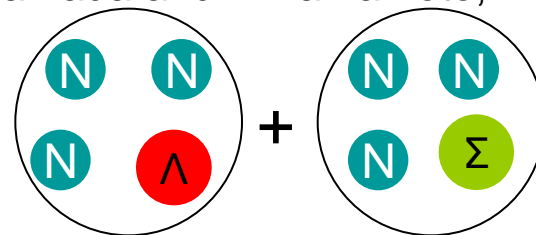


$^4_{\Lambda}\text{H}$

(slide from Hiyama)

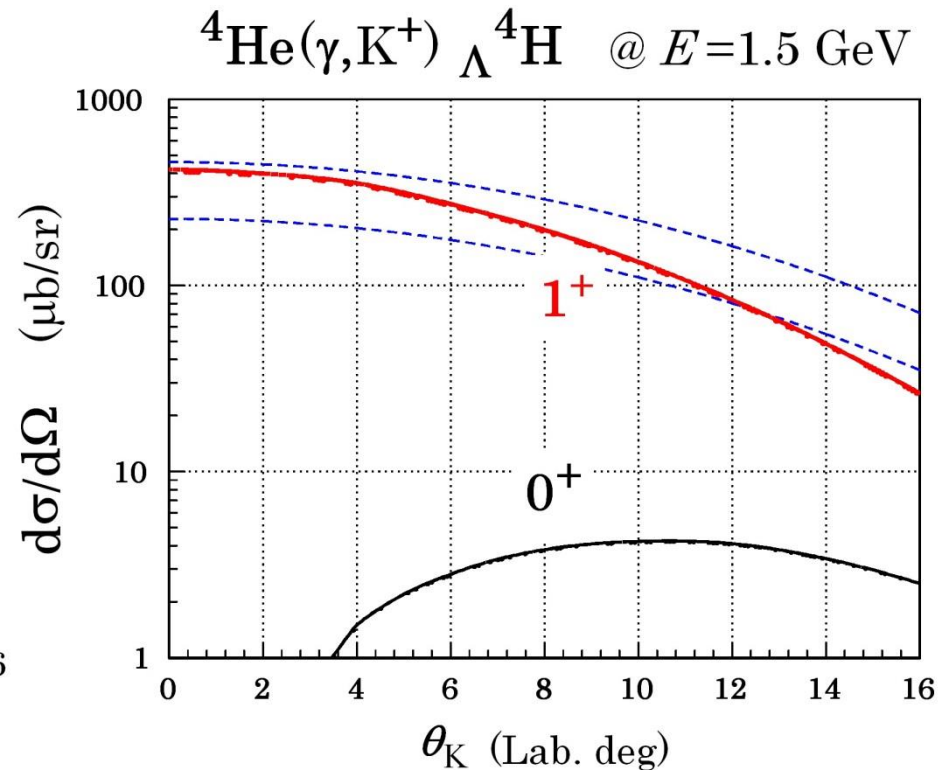
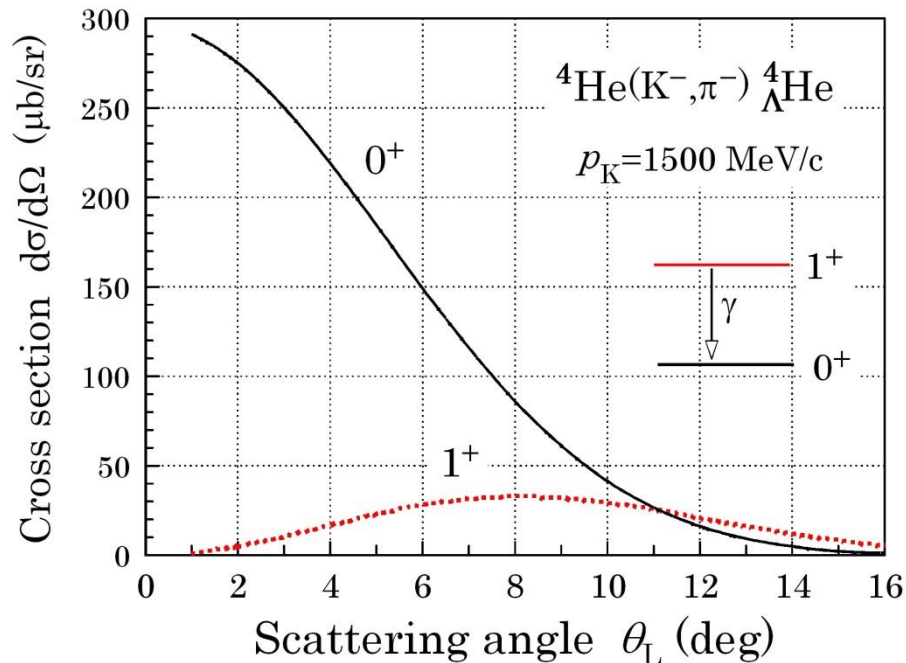
• E. Hiyama, M. Kamimura, T. Motoba, T. Yamada and Y. Yamamoto,
Phys. Rev. C65, 011301(R) (2001).

• H. Nemura, Y. Akaishi and Y. Suzuki,
Phys. Rev. Lett.89, 142504 (2002).



${}^4\text{He}(\text{K}^-, \pi^-)$ vs. ${}^4\text{He}(\gamma, \text{K}^+)$

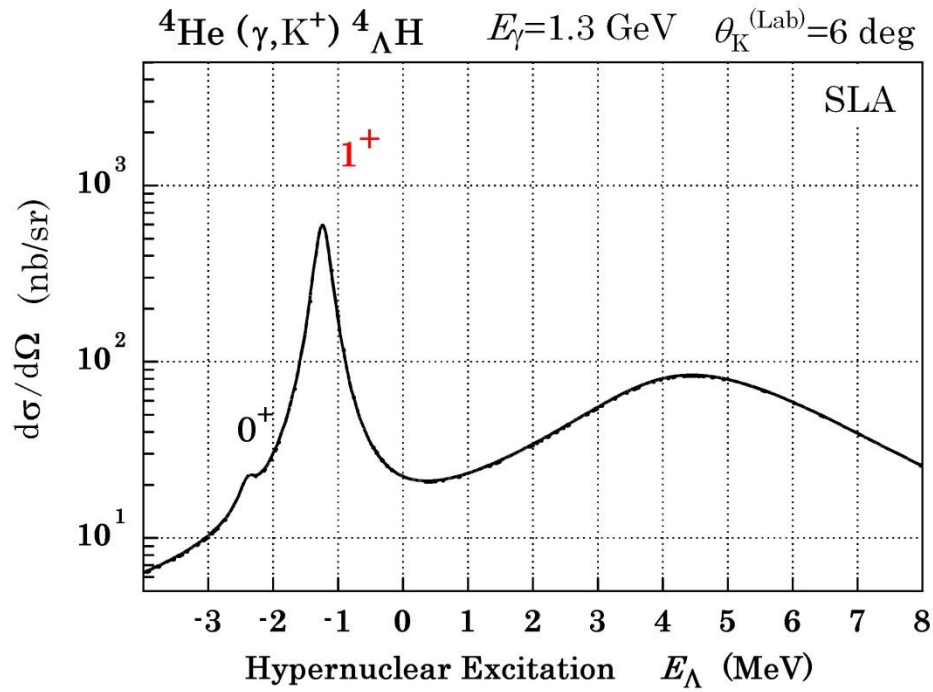
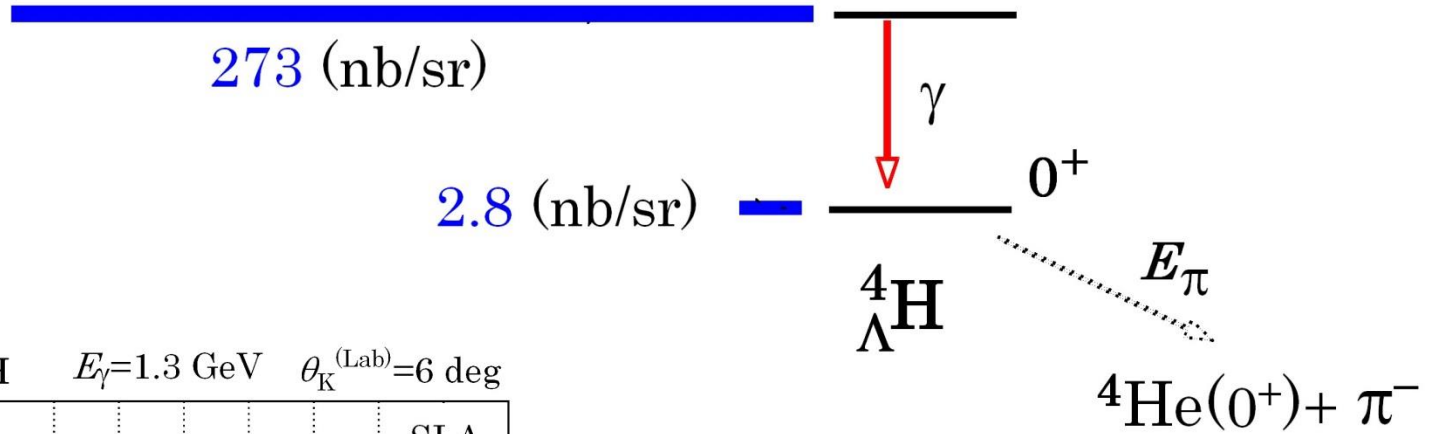
(A reaction theoretical view)



1^+ gets minor X-S but excited anyway, then $(\text{K}, \pi\gamma)$ coincidence method successful. → [Tamura's talk](#)

$X\text{-S}(1^+)$ is predominantly larger than $X\text{-S}(0^+)$

${}^4\text{He}(\gamma, K^+) \quad E_\gamma = 1.5 \text{ GeV}$
 cross section ($\theta_K = 6 \text{ deg}$)



Measure weak
 decay π energy:
($T=53.24$)
 $p=133.03 \text{ MeV}/c$

$E(1+)$ peak energy will be determined by ${}^4\text{He}(e, e' K^+)$
 10

(2) JLab (e,e'K+) experiments opened a new stage of hypernuclear reaction spectroscopy

- Success of JLab experiments (Hall A & C) on p-shell targets ---- high E resolution ~540keV
- Suggesting new theoretical aspects

The most typical one: $^{12}\text{C}(e, e'K^+)$

Tang et al. PRC 90(2014)

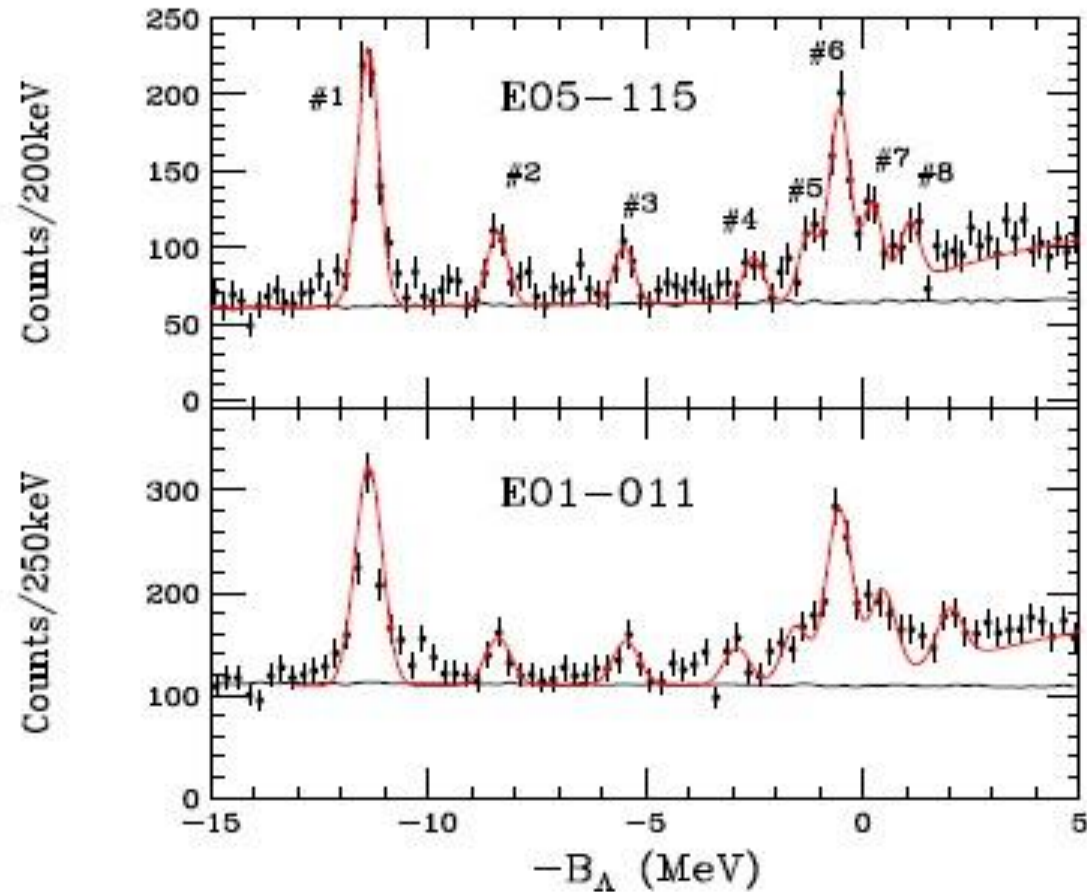
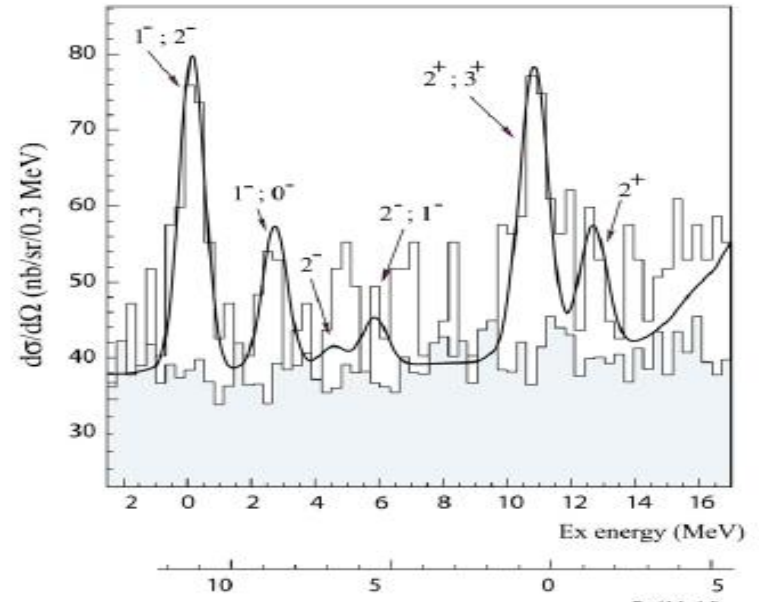
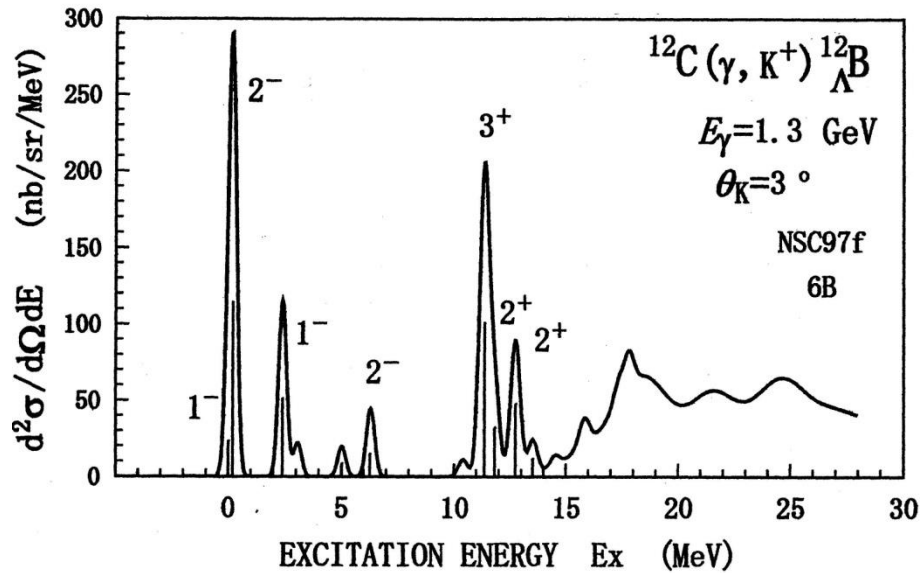


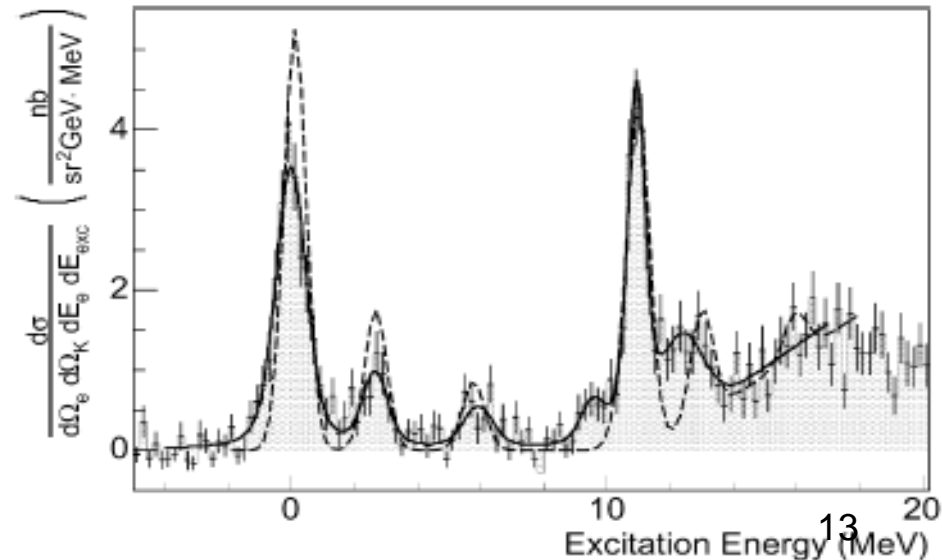
FIG. 10. (Color online) Spectroscopy of $^{12}_{\Lambda}\text{B}$ from the E05-115 and E01-011 experiments. The area below the black line is the accidental background.

Theor. prediction vs. (e,e'K⁺) experiments

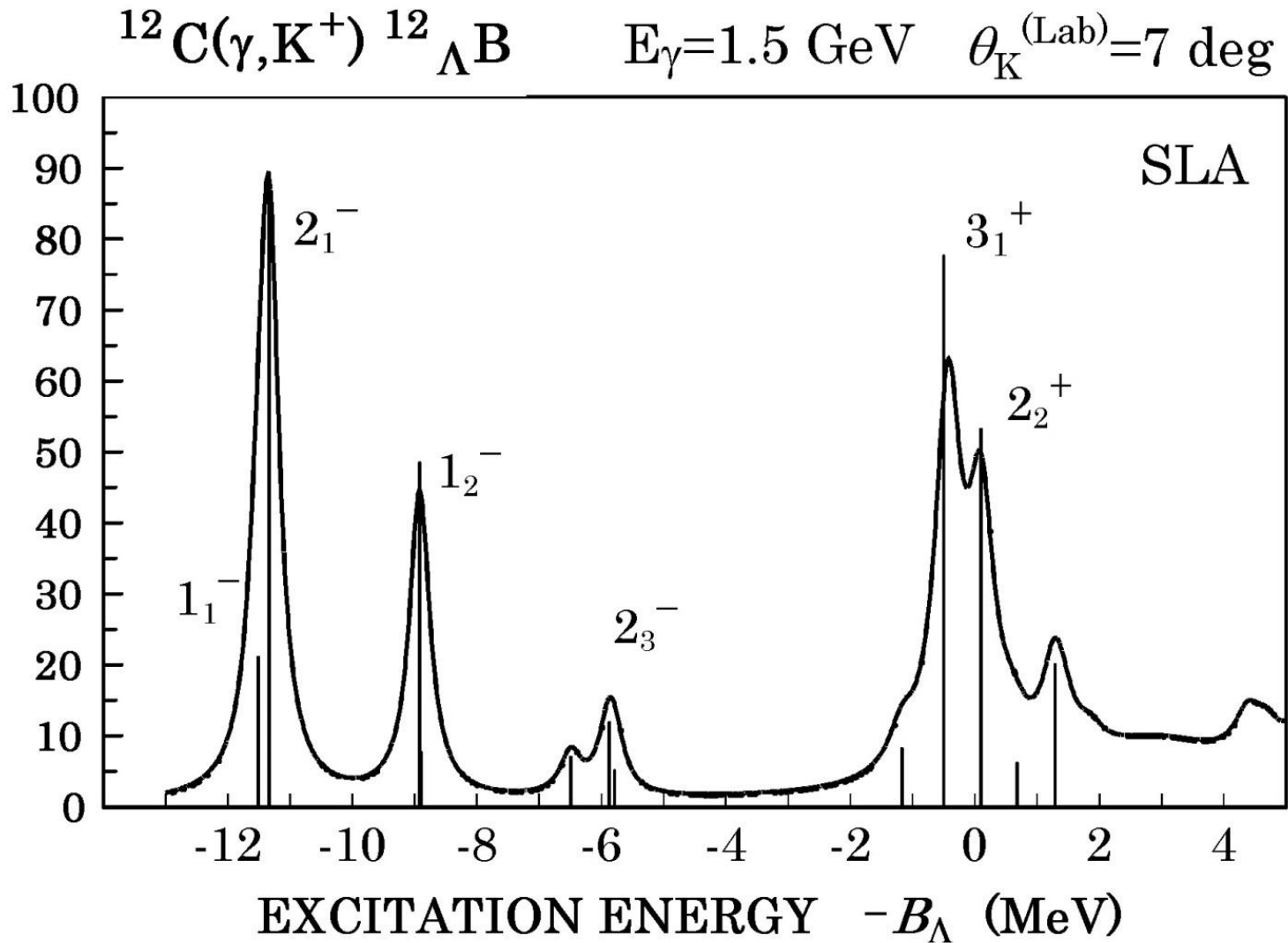


Motoba, Sotona, Itonaga,
*Prog.Theor.Phys.Sup.***117** (1994)
 T.M. *Mesons & Light Nuclei* (2000)
 updated w/NSC97f.

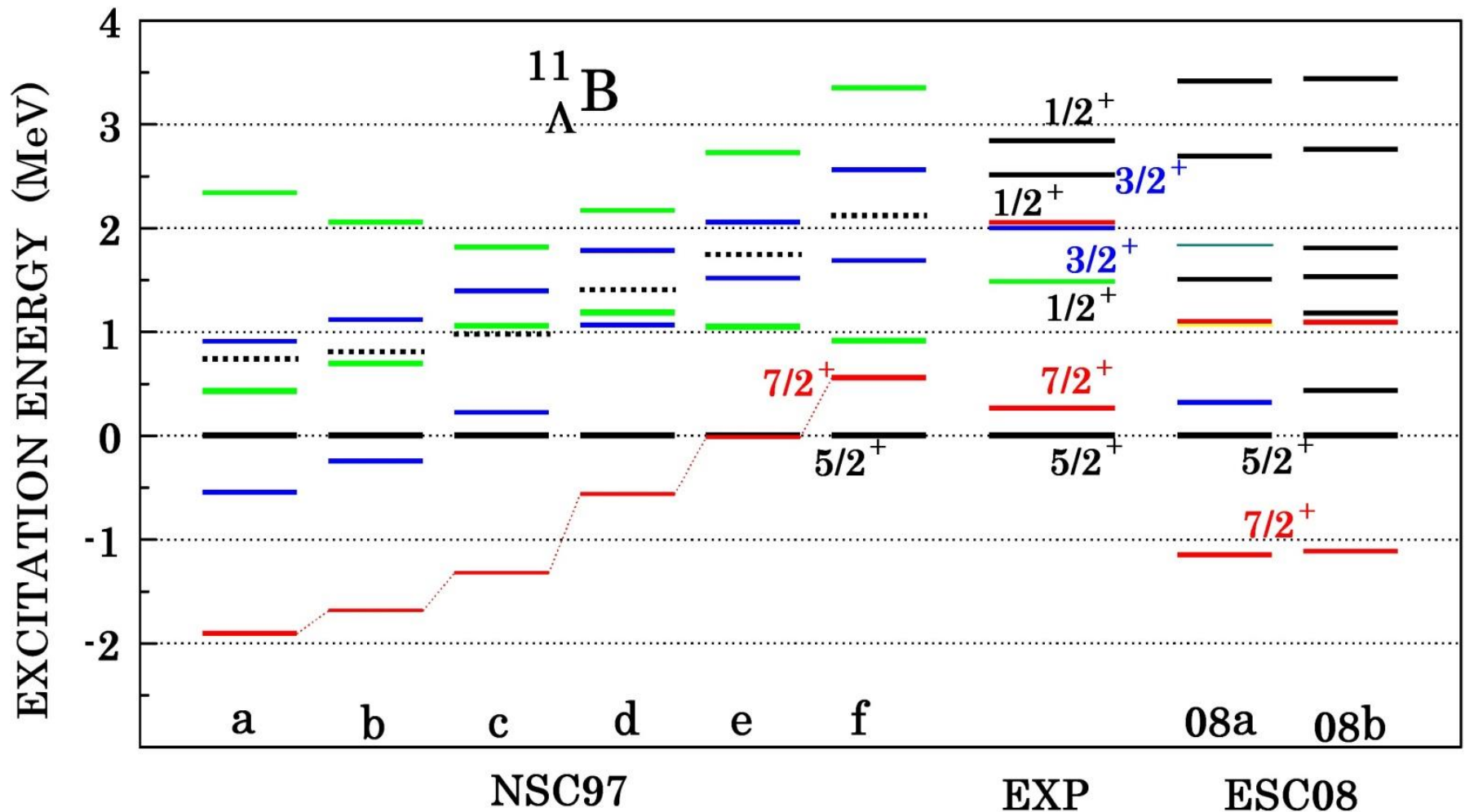
----->
Hall C (up) T. Miyoshi et al.
*P.R.L.***90** (2003) 232502. $\Gamma=0.75\text{keV}$
Hall A (bottom), J.J. LeRose et al.
N.P. **A804** (2008) 116. $\Gamma=0.67\text{keV}$



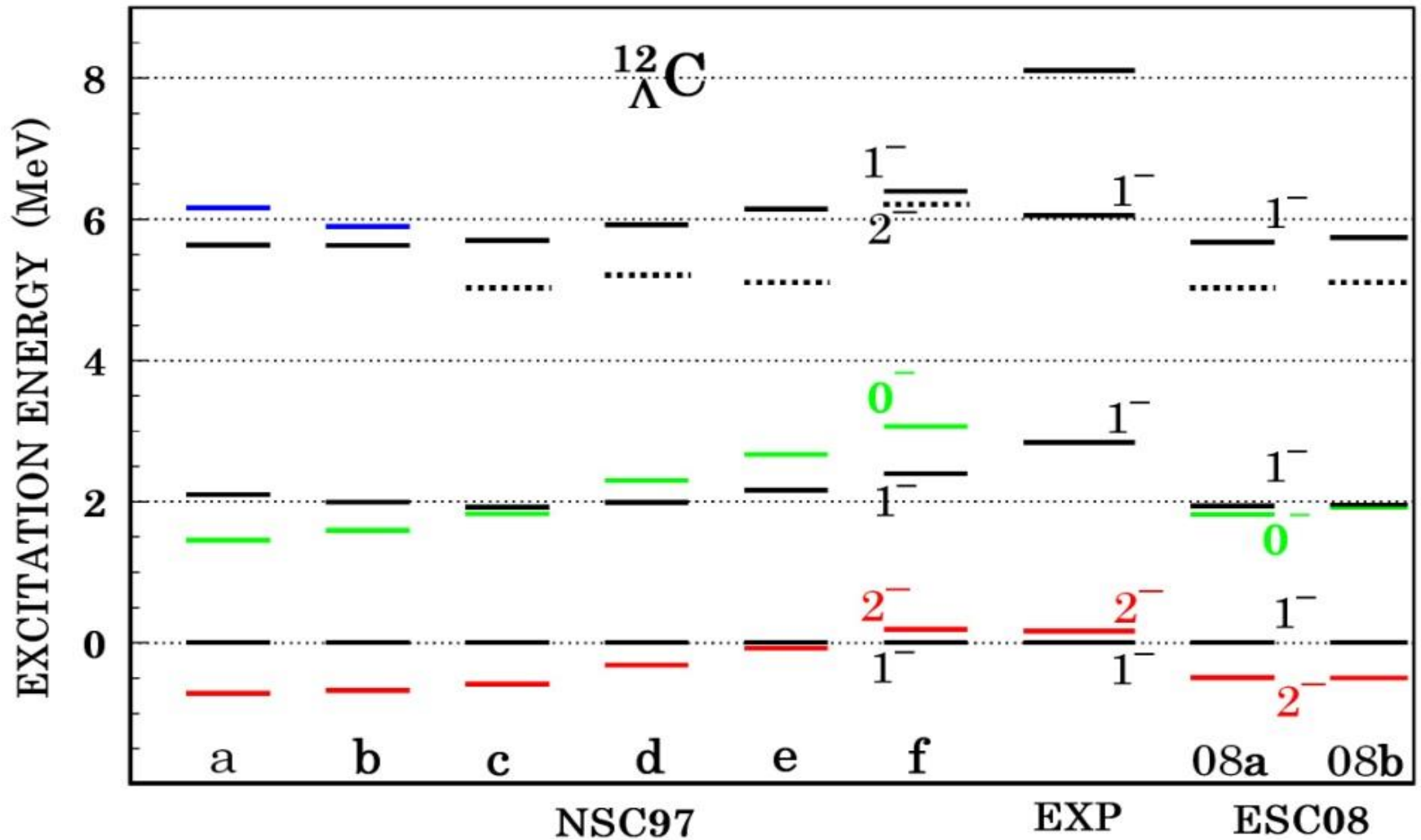
DWIA calculation



Nijmegen B-B interaction model improved by taking account of hypernuclear reaction data γ



Providing discrimination of several versions of meson theoretical Y-N interaction models



EXP. X-S and DWIA estimates are in good correspondence

12C(γ, K^+) Cross sec. calculated in DWIA at $E_\gamma = 1.5\text{GeV}$, $\theta_K(\text{Lab})=7\text{deg}$

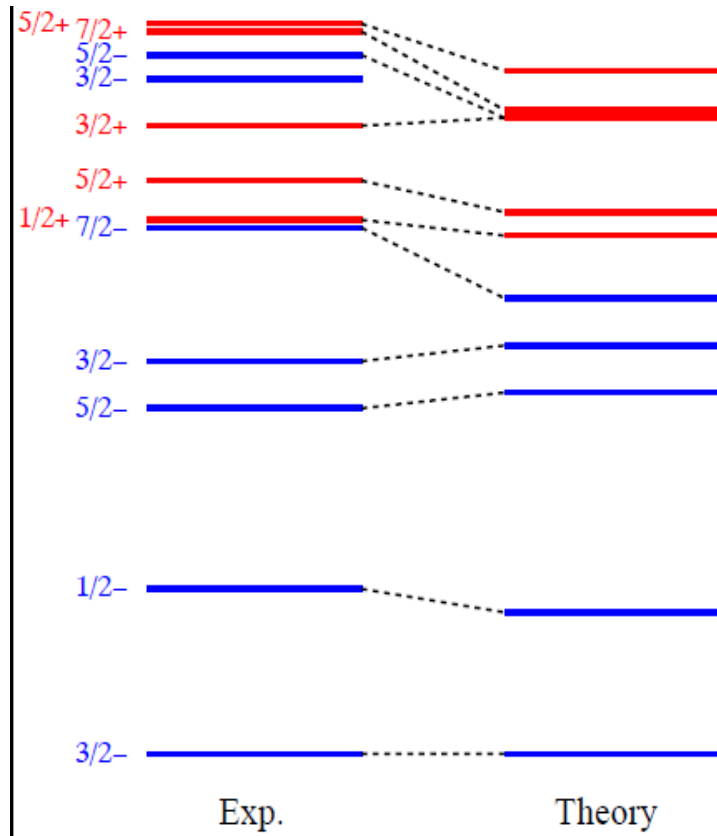
(Relative strengths with respect to the ground-state peak are also shown for reference)

Peak	J _i	Ex (MeV) NSC97f	Ex (MeV) DJM	EXP Ex	EXP E_Lam	Theory S6B [nb/sr] (sum)	Theory SLA [nb/sr] (sum)
#1	1 ₋ (1)	0.000	0.00	0.00	-11.508	10.54	21.04
	2 ₋ (1)	0.186	0.12	0.17	-11.340	63.05	89.33
#2	1 ₋ (2)	2.398	2.59	3.10	-8.390	18.96	48.44
	0 ₋ (1)	3.062	2.60			5.18	7.66
#3	2 ₋ (2)	5.022	5.02			4.90	6.96
	3 ₋ (1)	5.411				(-)	(-)
	4 ₋ (1)	6.228				(-)	(-)
	2 ₋ (3)	6.267	5.64	6.04	-5.488	8.34	11.84
	3 ₋ (2)	6.356				(-)	(-)
	1 ₋ (3)	6.389	5.72			2.29	5.02
#5	2 ₊ (1)	11.000	10.29			1.34	1.33
	1 ₊ (1)	11.120	10.34	10.22	-1.220	5.11	8.16
#6	3 ₊ (1)	11.081	11.01	10.98	-0.524	57.14	77.56
#8	2 ₊ (2)	11.610	10.93			23.95	53.17
	0 ₊ (1)	11.860	11.86			0.18	0.16
	1 ₊ (2)	12.129	12.13			7.03	6.08
	2 ₊ (3)	12.784	12.80	12.49	1.047	9.12	19.96
	1 ₊ (3)	13.176	12.91			4.22	3.74
	0 ₊ (2)	13.772	13.77			0.01	0.01

Relative XS w.r.t. GS=100(*)

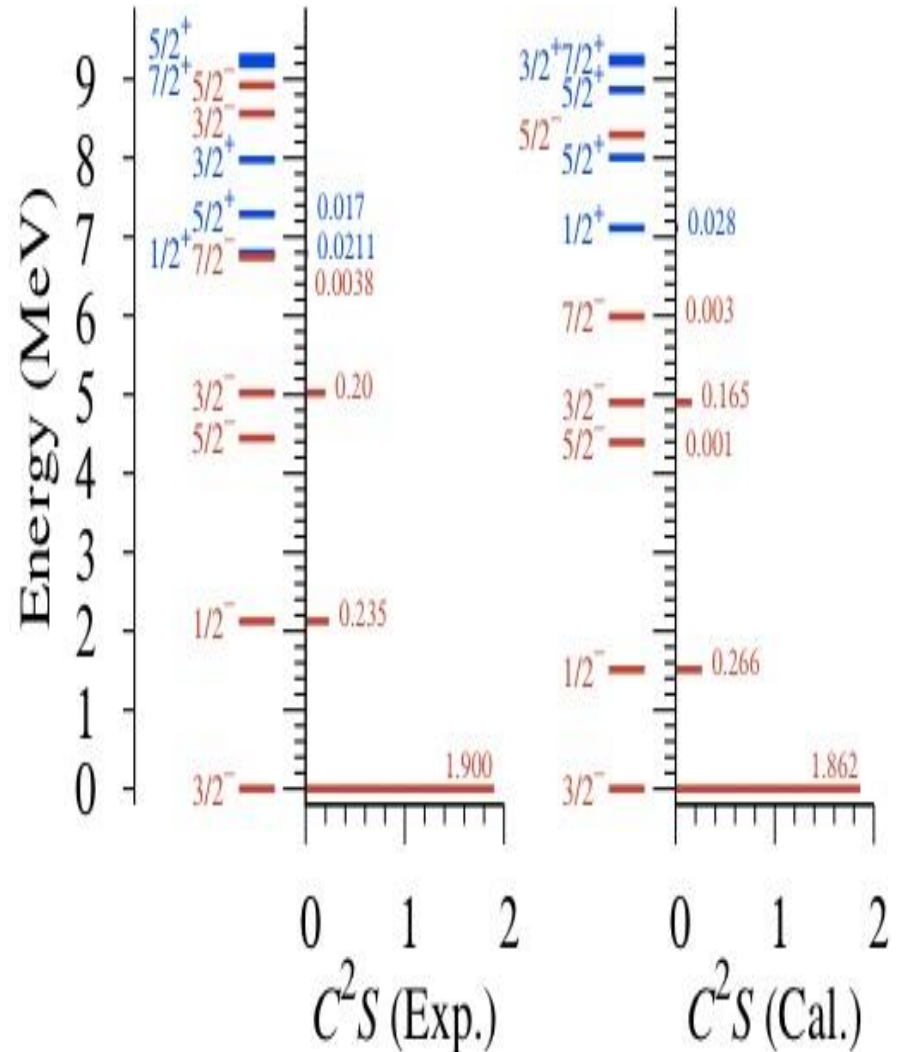
S6B		SLA		Kmaid	RpR
14.3	(Sum)	21.0	(Sum)		
85.7	100(*)	89.0	100(*)		
25.8	32.8	48.3	55.9		
7.0		7.6			
6.7		6.9			
		0.0			
		0.0			
11.3	21.1	11.8	23.7		
		0.0			
3.1		5.0			
1.8	8.8	1.3	9.5		
6.9		8.1			
77.7	110.3	77.3	130.3		
32.6		53.0			
0.2		0.2			
9.6		6.1			
12.4	28.0	19.9	29.8		
5.7		3.7			
0.0		0.0			

There are opposite parity excited states at low energy $E < 10 \text{ MeV}$. (Many theoretical attempts)



^{11}B

(2015.05.10 Umeya)



Proton pickup S-factor

Our new theoretical challenge
is to take both parity states into account.

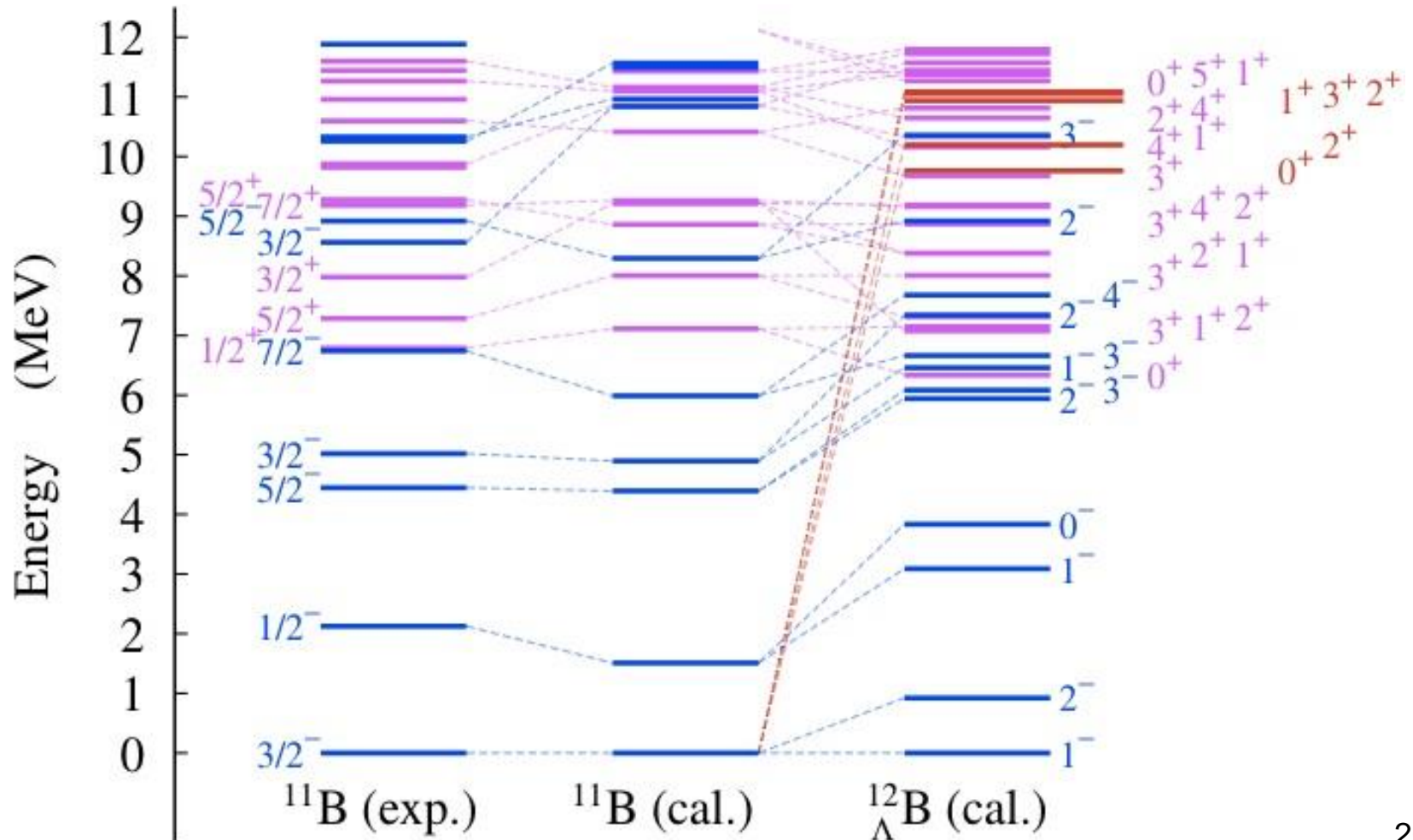
major shell mixing mediated by Λ

that is a new concept seen only in hypernucleus
(we have called it as parity mixing)

$${}^{12}_{\Lambda}\text{B}(\mathcal{J}_H^-) = \{ {}^{11}\text{B}(\mathcal{J}_C^-)_0 \times \Lambda_S \}^{(0)} + \{ {}^{11}\text{B}(\mathcal{J}_C^+)_1 \times \Lambda_P \}^{(2)}$$

$${}^{12}_{\Lambda}\text{B}(\mathcal{J}_H^+) = \{ {}^{11}\text{B}(\mathcal{J}_C^-)_0 \times \Lambda_P \}^{(1)} + \{ {}^{11}\text{B}(\mathcal{J}_C^+)_1 \times \Lambda_S \}^{(1)}$$

Results of parity mixing calculation (preliminary)



Components connected via (γ, K^+)

proton is converted $\rightarrow \Lambda$ in s or p orbits

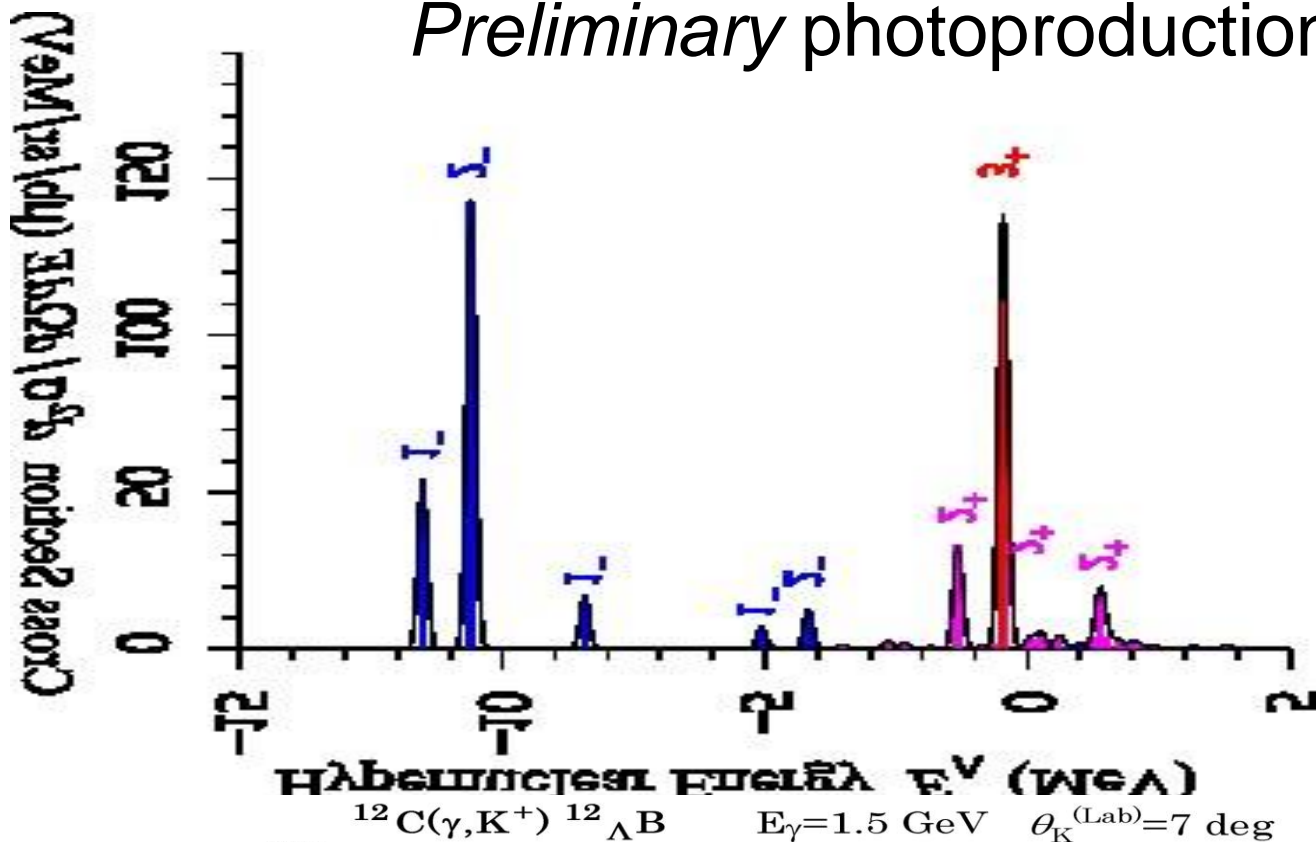
(So far **only green arrows** are taken into account.)

$$\begin{array}{l}
 {}^{12}\Lambda\text{B}(J_{\text{H}}^-) = \{ \underbrace{([s^4]p^7; J_c^-)_0 \times \Lambda_s}_{\text{green}} \}^{(0)} + \{ ([s^4]p^6(sd)^1; J_c^+)_1 \times \Lambda_p \}^{(2)} + \{ ([s^3]p^8; J_c^+)_1 \times \Lambda_p \}^{(2)} \\
 {}^{12}\text{C}(0^+)_{0+2h\omega} = \underbrace{|[s^4]p^8\rangle}_{\text{green}} + |[s^4]p^7(fp)^1\rangle + |[s^4]p^6(sd)^2\rangle + |[s^3]p^8(sd)^1\rangle + |[s^2]p^{10}\rangle \\
 {}^{12}\Lambda\text{B}(J_{\text{H}}^+) = \{ \underbrace{([s^4]p^7; J_c^-)_0 \times \Lambda_p}_{\text{green}} \}^{(1)} + \{ ([s^4]p^6(sd)^1; J_c^+)_1 \times \Lambda_s \}^{(1)} + \{ ([s^3]p^8; J_c^+)_1 \times \Lambda_s \}^{(1)}
 \end{array}$$

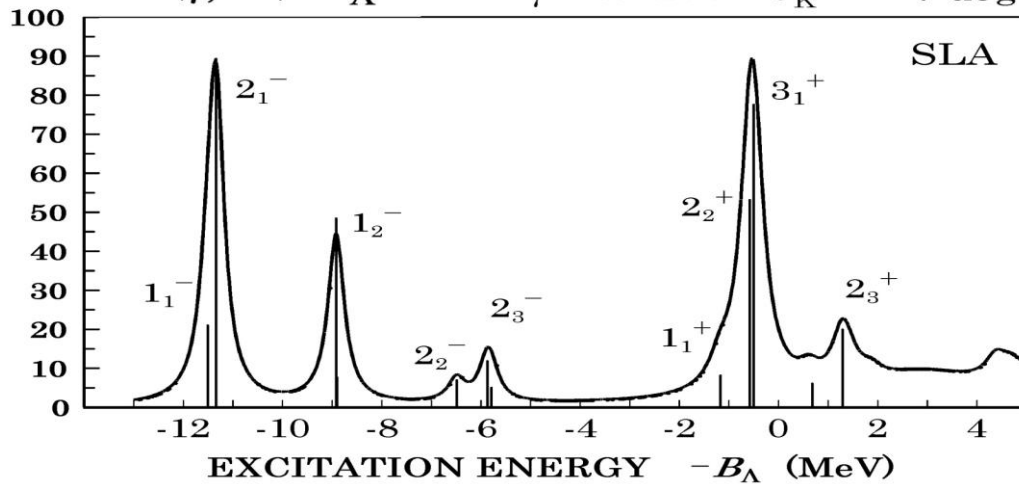
Problems:

What kind of effective interactions should be used
in describing those WF in the extended model space.

Preliminary photoproduction X-S



Down:
Previous
calculation



(3) Medium-heavy nuclear targets

A typical example of medium-heavy target : ^{28}Si : $(d_{5/2})^6$ and $(sd)^6\text{P}(sd)^6\text{N}$

to show characteristics of the (γ, K^+) reaction with DDHF w.f.

Spin-orbit splitting:
consistent with ${}_{\Lambda}^7\text{Li}$, ${}^9\text{Be}$, ${}^{13}\text{C}$, ${}^{89}\text{Y}$

These characteristic merits of the $\gamma p \rightarrow \Lambda K^+$ process (ability to excite high-spin unnatural-parity states) should be realized better in heavier systems involving large j_p and large j_Λ

$$(e, e' K^+) \quad d^3\sigma/dE_e d\Omega_e d\Omega_K = \Gamma \times d\sigma/d\Omega_K$$

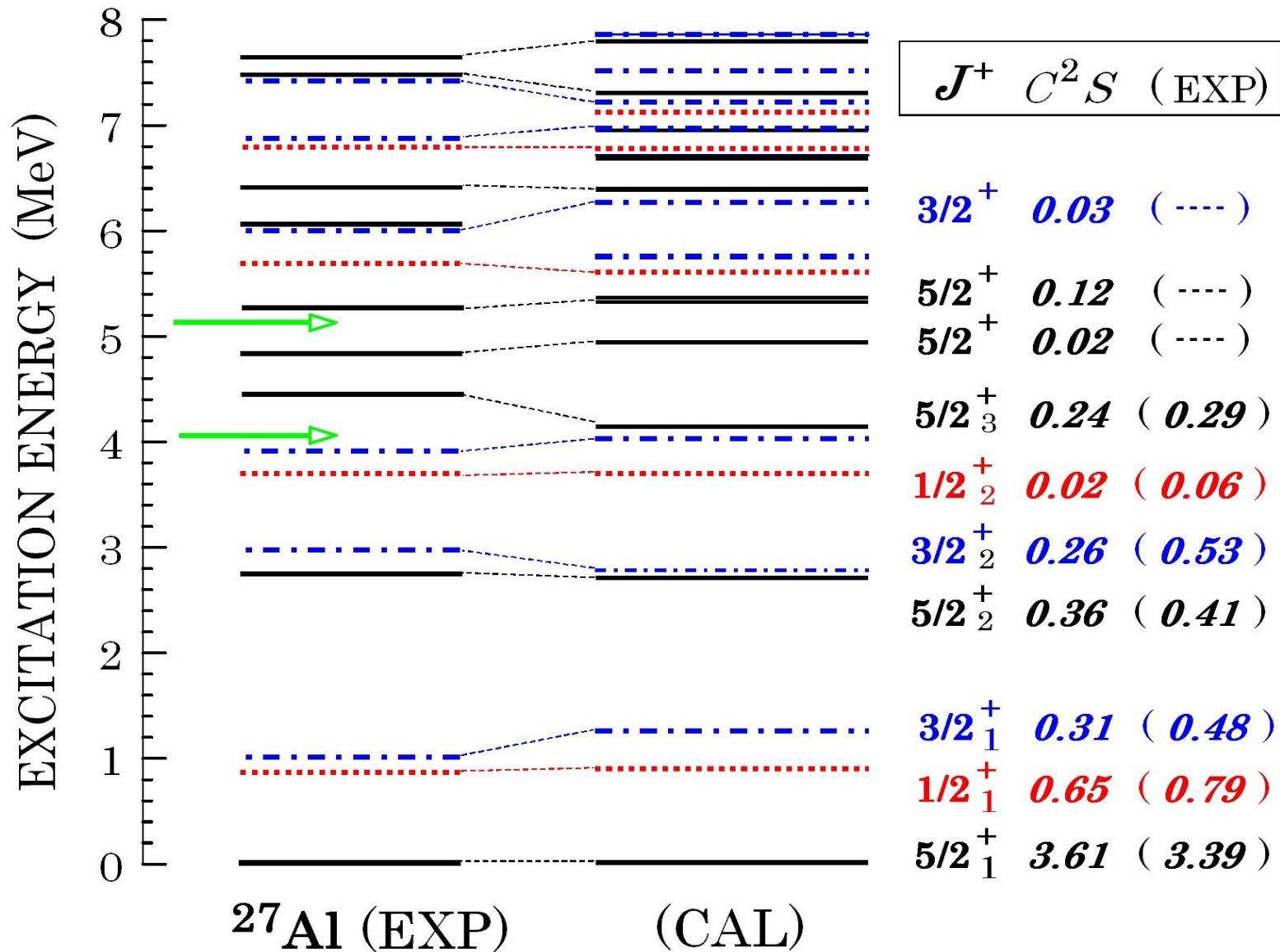
Γ : virtual photon flux (kinematics)

Hereafter we discuss $d\sigma/d\Omega_K$ for ${}^A Z (\gamma, K^+)_{\Lambda} {}^A Z'$

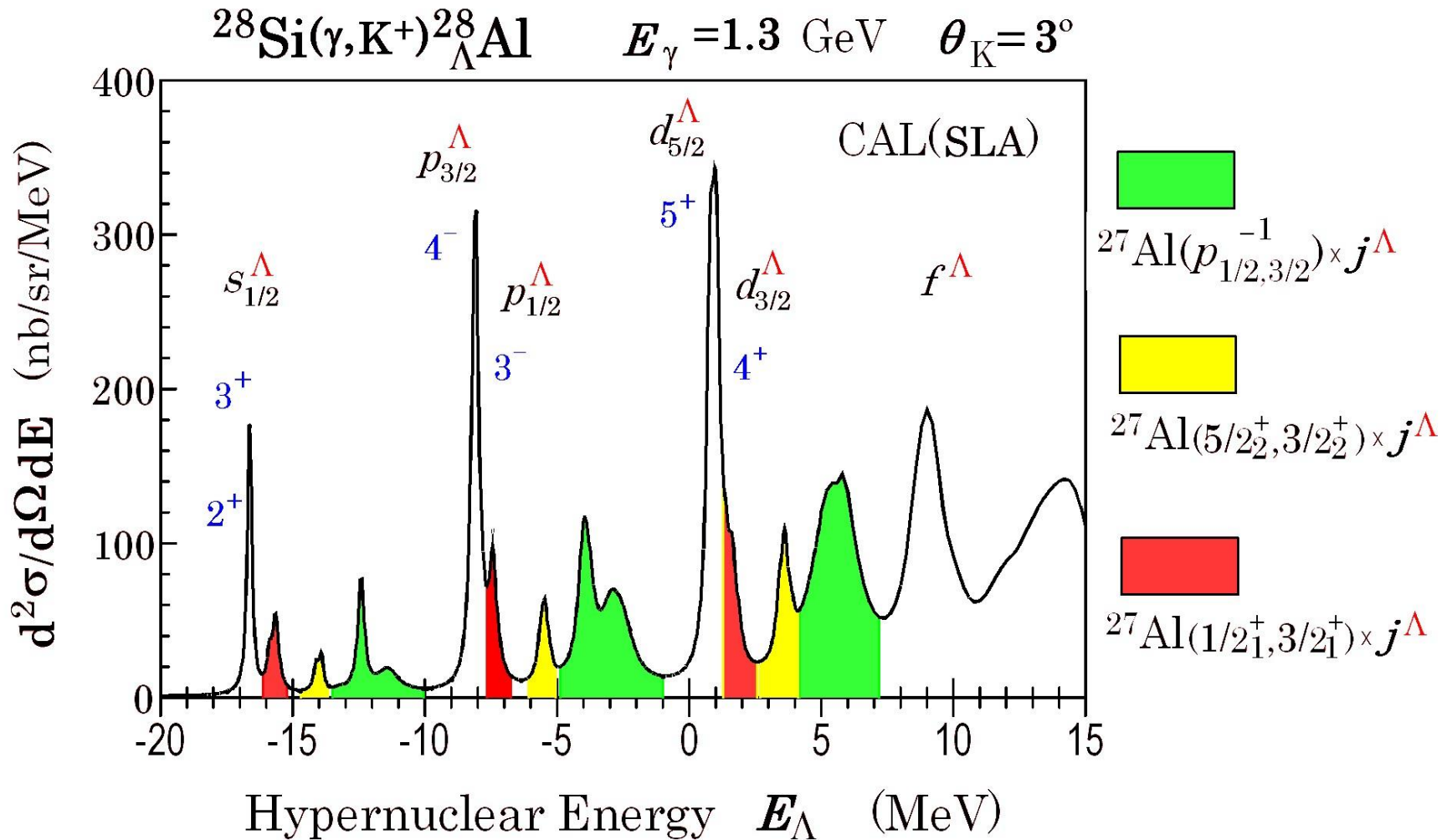
Theor. x-section for $(d_{5/2})^6 (\gamma, K^+) [j_h - j_\Lambda] J$

DWIA		[nb/sr]											
Lambda=		s1/2L		p3/2L		p1/2L		1s1/2L		d5/2L		d3/2L	
		(-16.92)		(-8.40)		(-8.40)		(0.32)		(0.69)		(0.69)	
Proton hole d5/2 (-16.17) (g.s.)				1-	5.4					0+	0.0		
				2-	7.1	2-	19.4	2+	2.2	1+	26.0	1+	8.9
	2+	29.2		3-	4.2	3-	76.2	3+	4.6	2+	0.3	2+	34.9
	3+	63.8		4-	141.8					3+	26.7	3+	30.4
										4+	0.5	4+	112.0
										5+	164.1		
p1/2 (-25.49)	0-	9.4				0+	0.0	0-	3.7				
	1-	30.5	1+	2.0	1+	28.3	1-	12.2			1-	1.4	
			2+	66.9					2-	10.7	2-	43.5	
									3-	76.9			
p3/2 (-29.84)			0+	0.0								0-	2.0
	1-	14.3	1+	8.9	1+	1.8	1-	5.9	1-	3.2	1+	5.7	
	2-	59.1	2+	0.4	2+	62.5	2-	24.8	2-	4.5	2+	17.5	
			3+	109.1					3-	4.5	3+	96.3	
									4-	148.6			
s1/2 (-44.55)	0+	0.1			0-	7.3	0+	0.3					
	1+	19.2	1-	12.1	1-	23.7	1+	51.4			1+	16.5	
			2-	50.0					2+	27.0	2+	40.1	
									3+	58.1			

Proton pickup from $^{28}\text{Si}(0^+):(sd)^6=(d_{5/2})^{4.1}(1s_{1/2})^{0.9}(d_{3/2})^{1.0}$

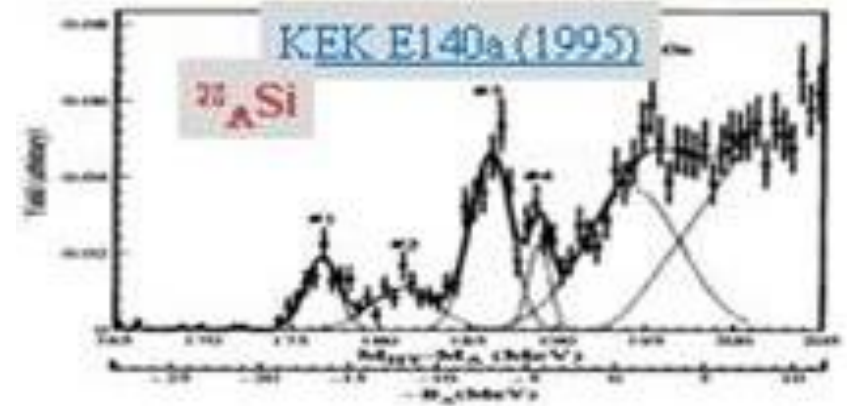
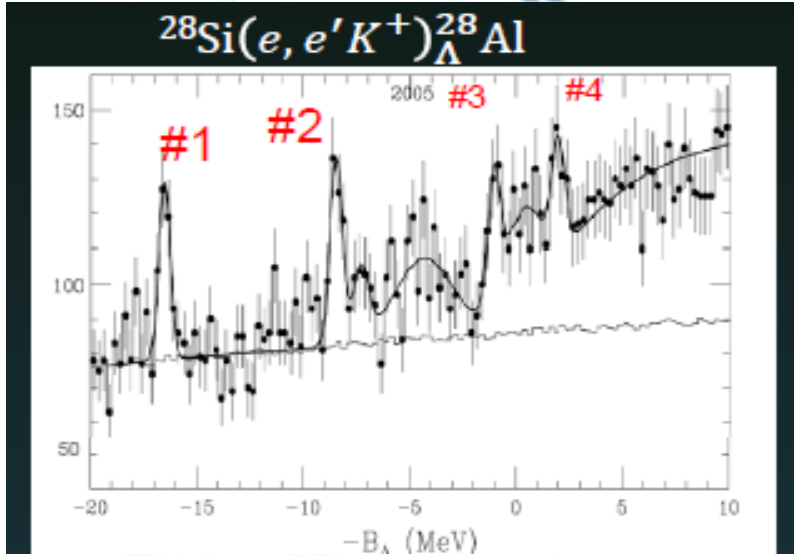


Peaks can be classified by the characters

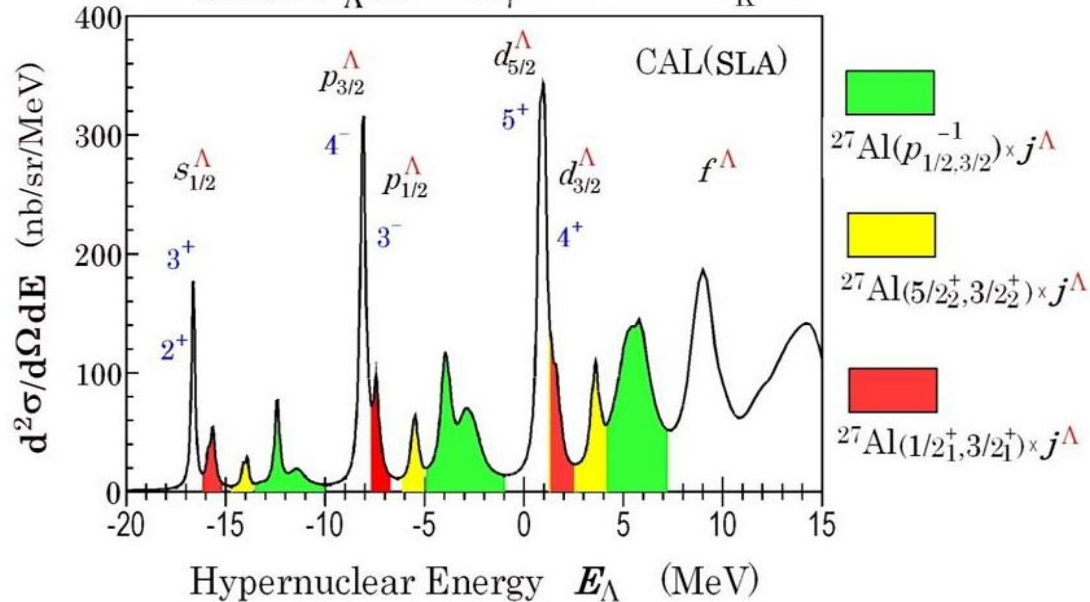


Major peak series : $[^{27}\text{Al}(5/2_1^+) \times j^{\Lambda}]_J$ with $j^{\Lambda} = s, p, d, \dots$

$^{28}\text{Si}(e, e'K^+)^{28}_{\Lambda}\text{Al}$ – First Spectroscopy of $^{28}_{\Lambda}\text{Al}$



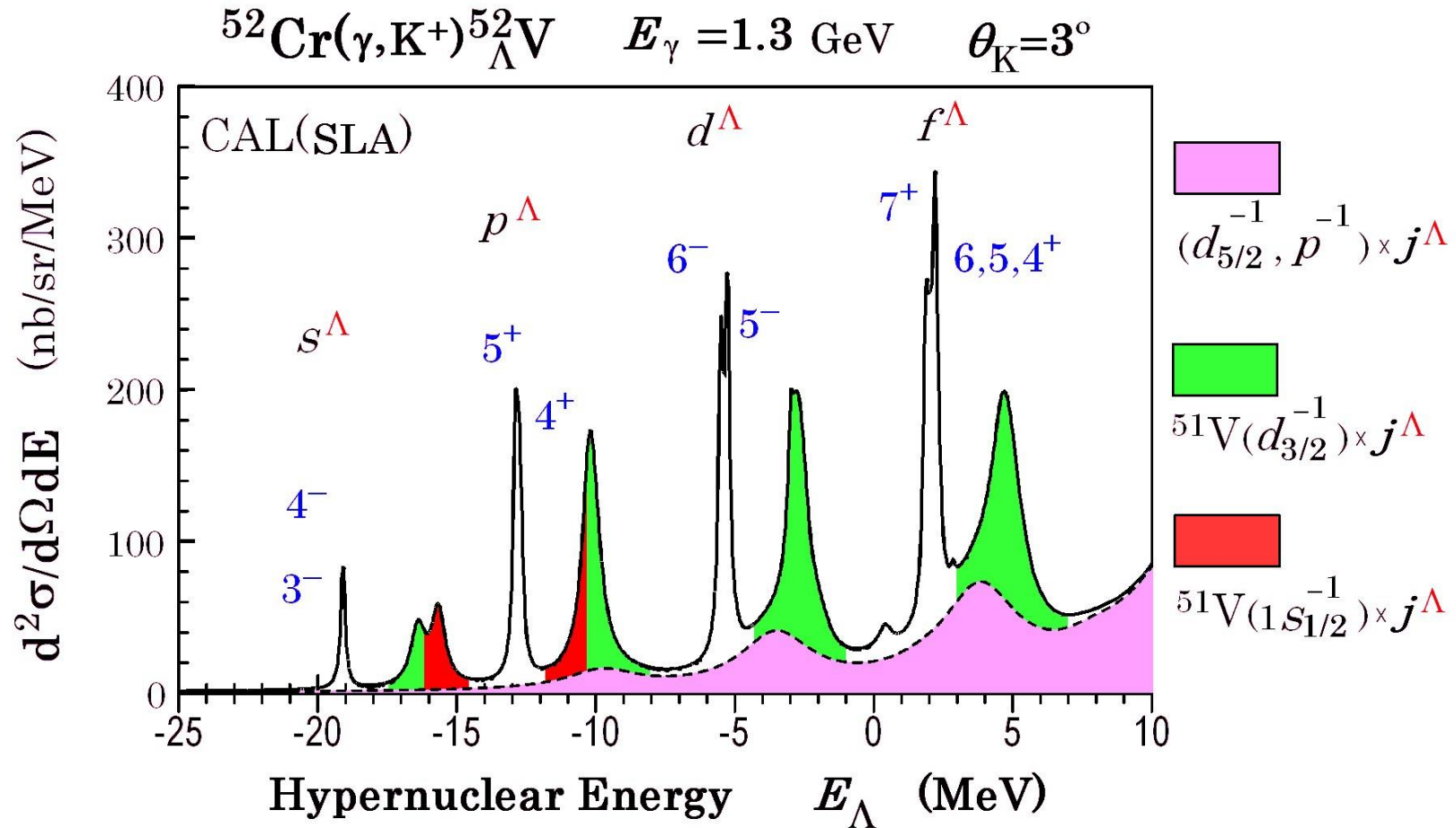
$^{28}\text{Si}(\gamma, K^+)^{28}_{\Lambda}\text{Al}$ $E_{\gamma} = 1.3 \text{ GeV}$ $\theta_K = 3^{\circ}$



Exp. data:
 Nakamura(HYP2015),
Seems promising,
 (waiting for finalization
 of exp. analysis)

^{52}Cr (j_{Λ} dominant target case)

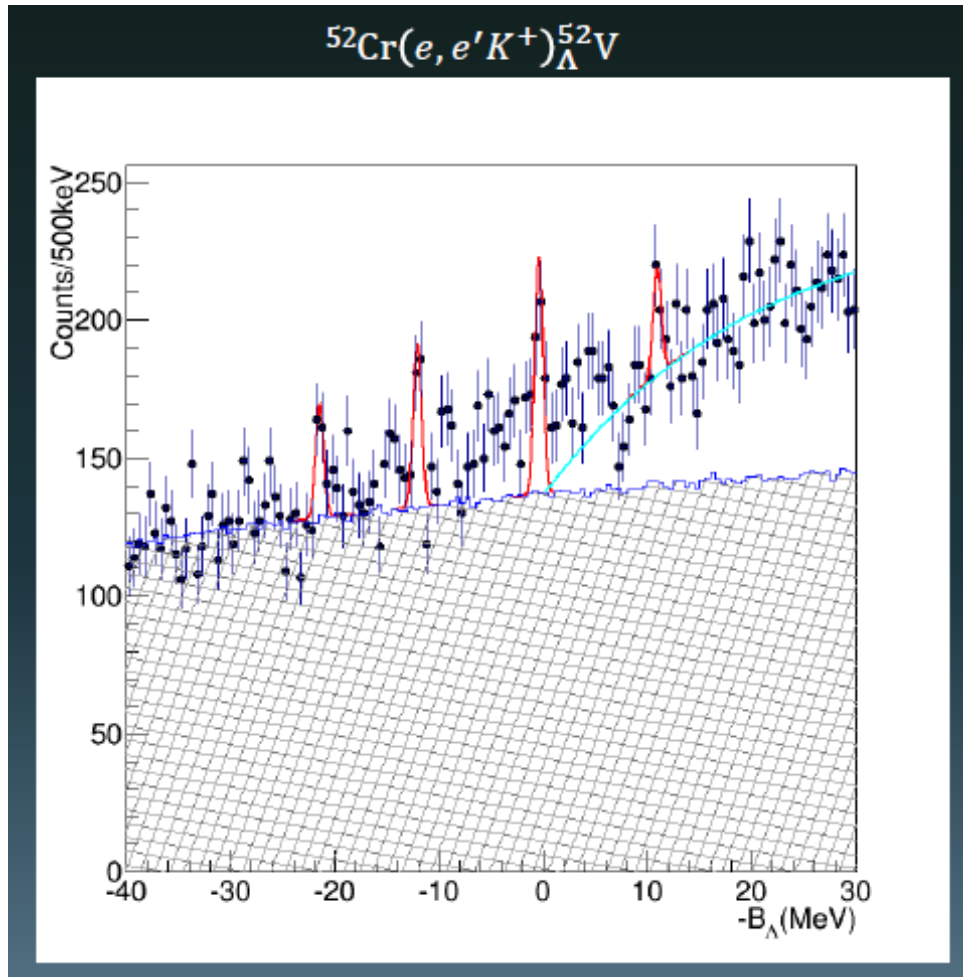
typical unnatural-parity high-spin states



Major peak series : $[^{51}\text{V}(7/2^{-}; \text{gs}) \times j^{\Lambda}]_J$ with $j^{\Lambda} = s, p, d, f, \dots$

$^{52}\text{Cr}(e, e'K^+) ^{52}_{\Lambda}\text{V}$ in analysis

Nakamura, priv. commun. (HYP2015)

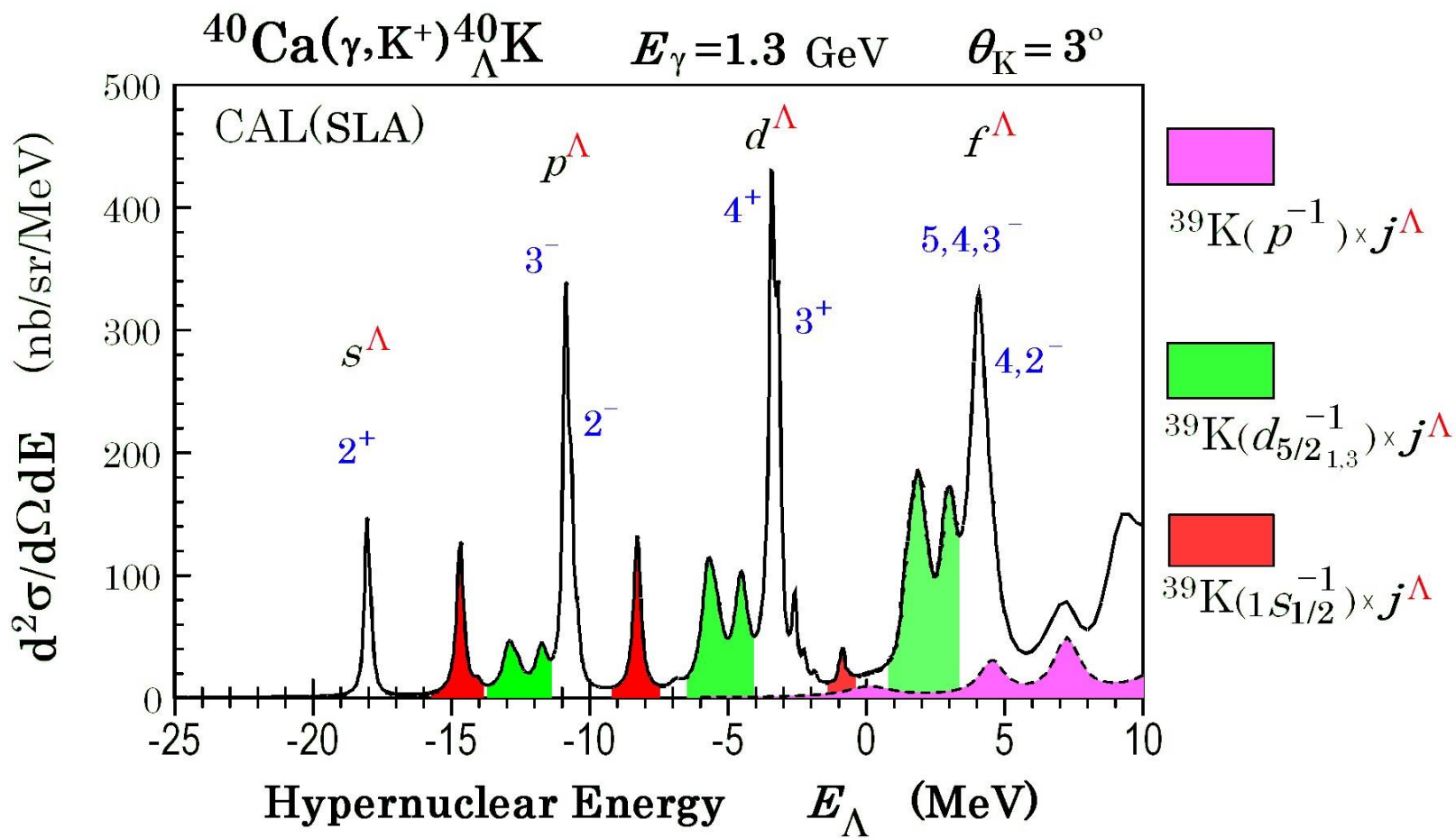


peak	B_{Λ} (MeV)
#1	-21.4
#2	-12.1
#3	-0.4
#4	+10.9

E01-115

^{40}Ca (LS-closed shell case):

high-spin states with natural-parity ($2^+, 3^-, 4^+$)



Major peak series : $[^{39}\text{K}(d_{3/2}^{-1}; \text{gs}) \times j^{\Lambda}]_J$ with $j^{\Lambda} = s, p, d, f, \dots$

Well-separated series of peaks due to large q and spin-flip dominance:

$$j_{>} = l + 1/2, \quad j_{<} = l - 1/2$$

$[(nlj)_p^{-1} (nlj)_{\Lambda}^{\Lambda}]_J$ a series of pronounced peaks

jj -closed target : (^{28}Si , ^{52}Cr)

$$[j_{>}^{-1} j_{>}^{\Lambda}]_J \quad J = j_{>} + j_{>}^{\Lambda} = l_p + l_{\Lambda} + 1 = L_{\max} + 1 \quad (\text{unnatural parity})$$

$$[j_{>}^{-1} j_{<}^{\Lambda}]_J \quad J = j_{>} + j_{<}^{\Lambda} = l_p + l_{\Lambda} = L_{\max} \quad (\text{natural parity})$$

LS -closed target : (^{40}Ca)

$$[j_{<}^{-1} j_{>}^{\Lambda}]_J \quad J = j_{<} + j_{>}^{\Lambda} = l_p + l_{\Lambda} = L_{\max} \quad (\text{natural parity})$$

(4) One of the major objects is to get systematics of Λ s.p.e.

Taken from: Millener-Dover-Gal, PRC18 (1988)

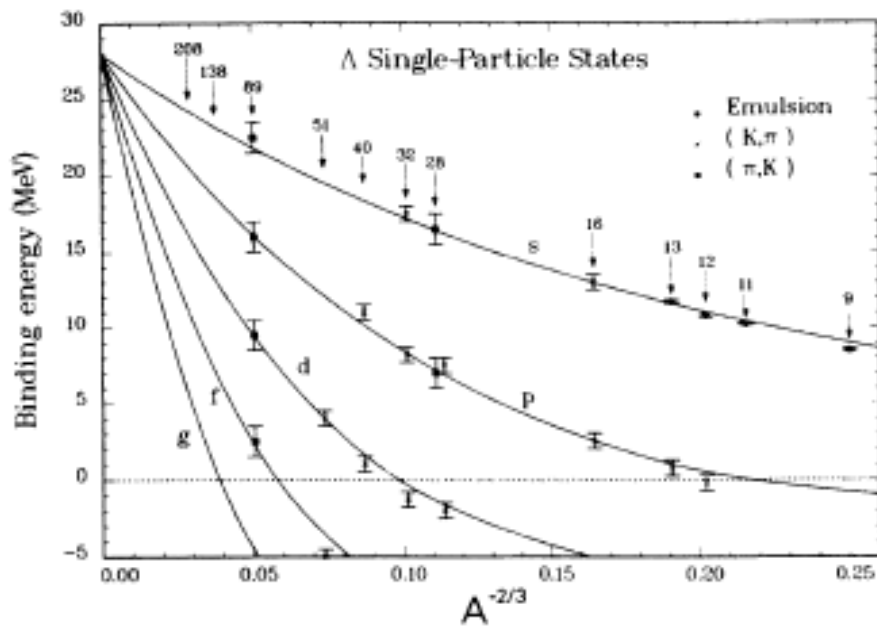


FIG. 1. The data on binding energies (B_Λ) of Λ sing

Woods-Saxon pot. $D=28\text{MeV}$
 $r_0=1.128+0.439A^{-2/3}$

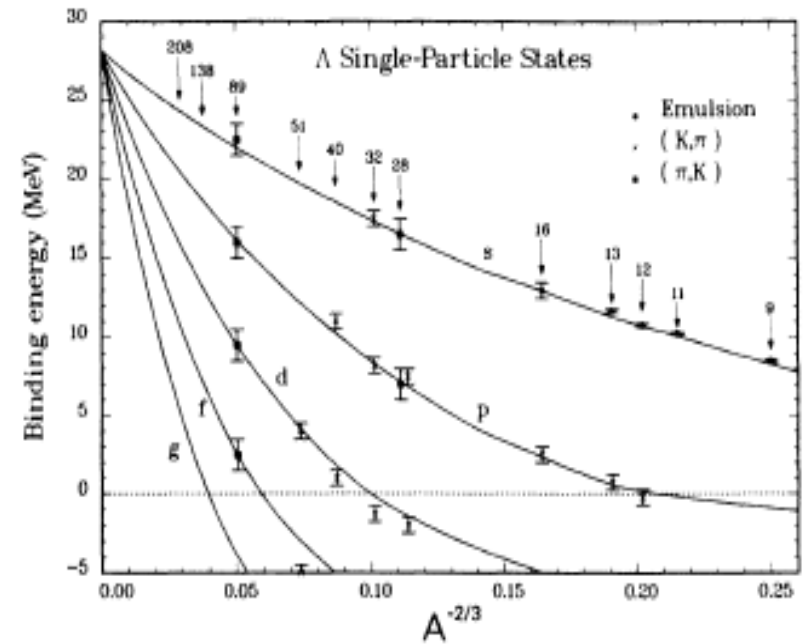


FIG. 5. Same as Fig. 4 but for the potential in Table III with $\rho^{4/3}$ density dependence.

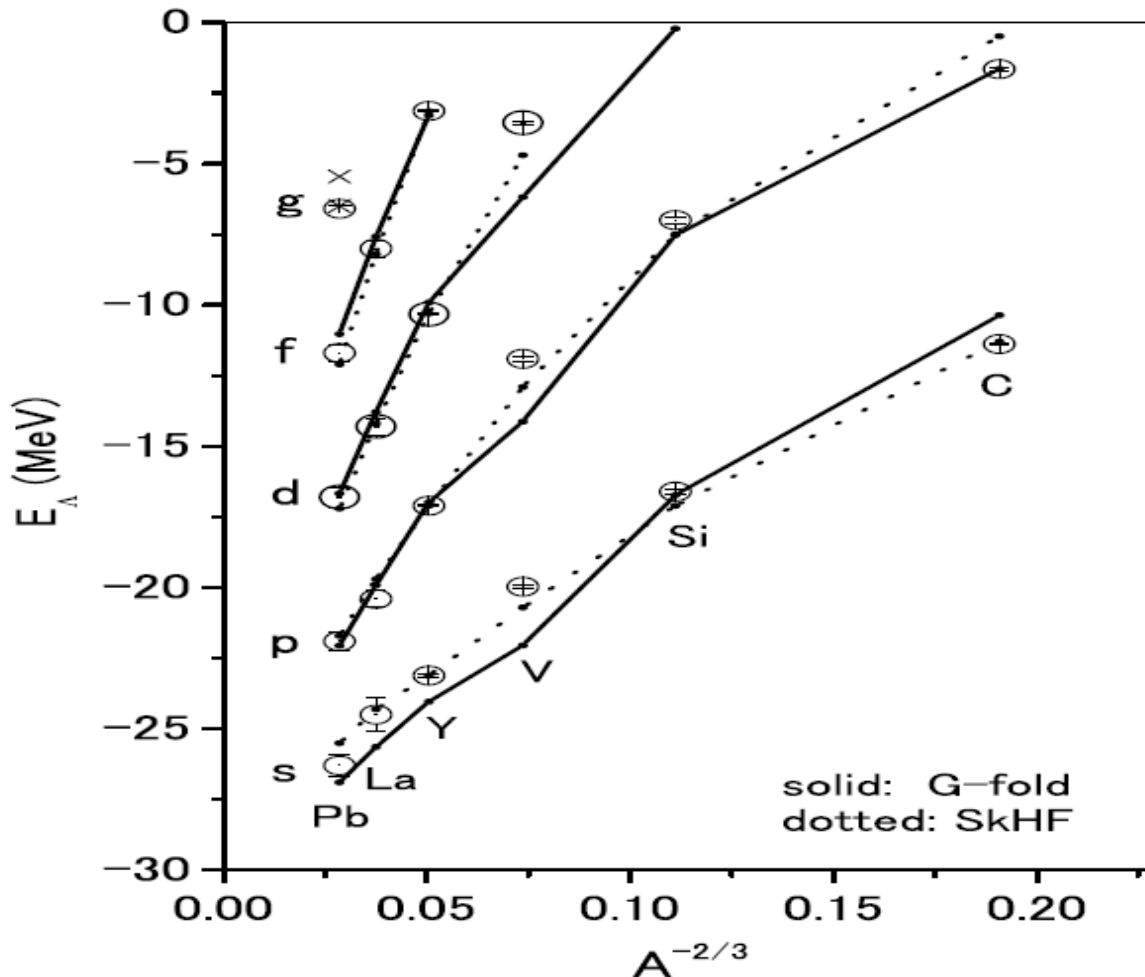
Skyrme HF with $\rho^{4/3}$
 Density dependent

Single-particle energies of Λ

G-matrix (ESC08c) results vs. experiments

(Y. Yamamoto et al.: PTP. S.185 (2010) 72 and priv. commun.)

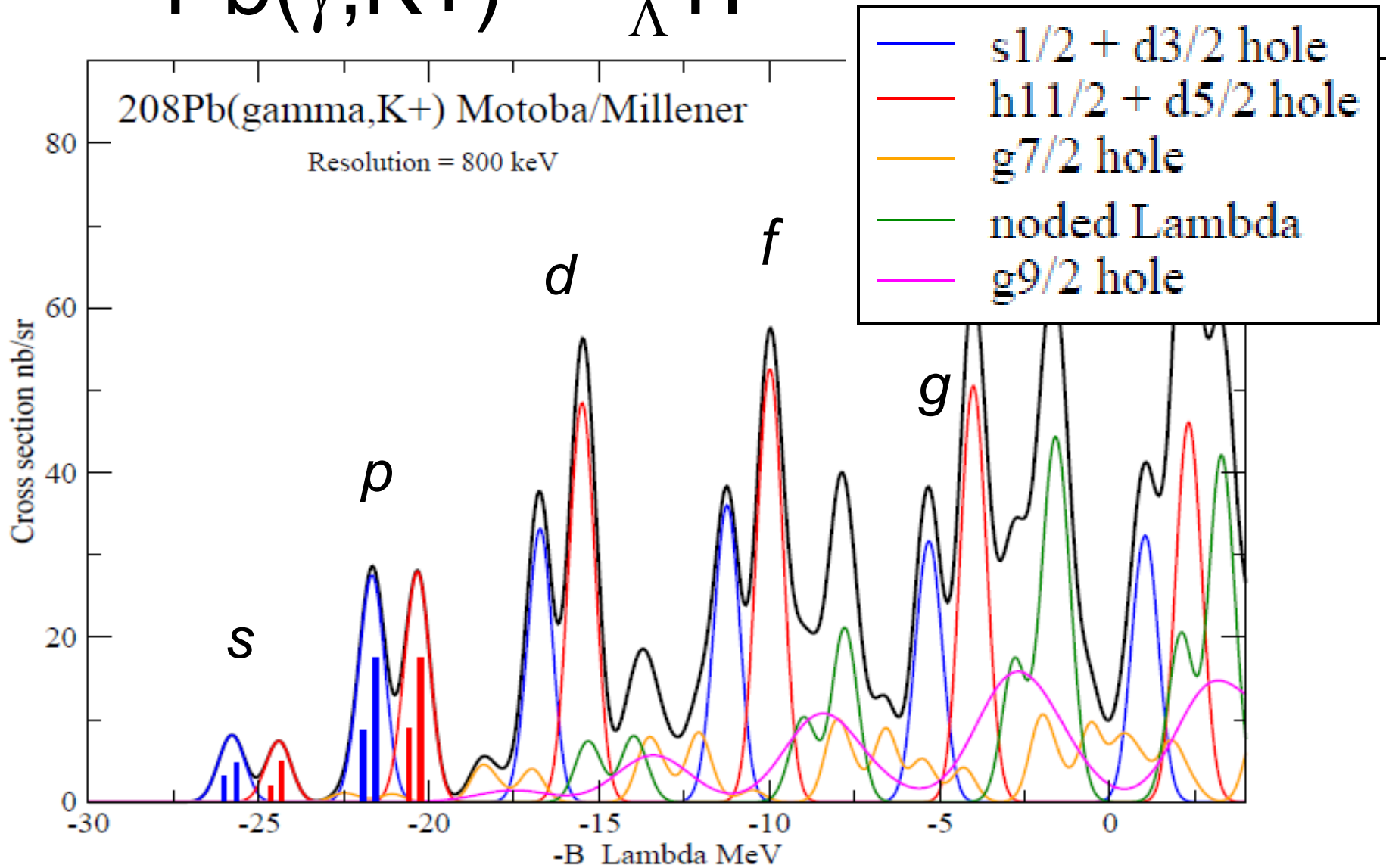
Y. Yamamoto, T. Motoba and Th. A. Rijken



*High resolution
exp. data over
wide A are
necessary.*

*sd, fp-shell and
heavier data are
quite Important to
extract the Λ
behavior in
nuclear matter.*

$^{208}\text{Pb}(\gamma, K^+) ^{208}_{\Lambda}\text{Tl}$



We have an opportunity to observe a series of Lambda orbits ?

(5) Another interesting topics related to medium-heavy hypernuclear structure includes

Λ -rotation(deformation) coupling

- Refer to Talks by Isaka and Hagino.

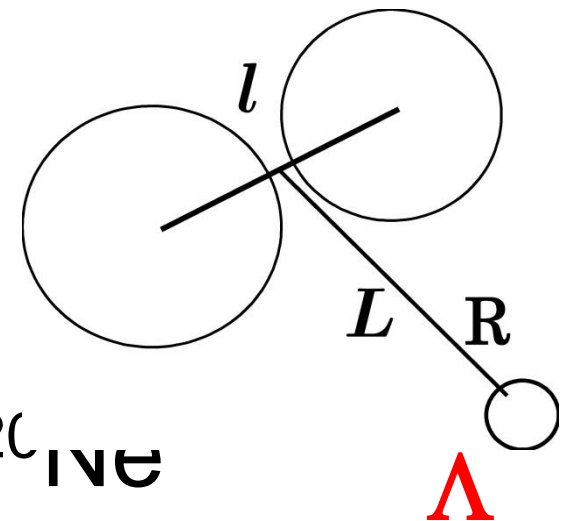
Why the strong coupling is realized between p-state Λ and $\alpha+\alpha$ core ?

Schematic consideration assuming the **SU(3) maximum configuration** for the nuclear g.s. rotational states:

$$(\lambda\mu)=(40) \quad \ell=0,2,4^+ \quad \text{for } {}^8\text{Be}$$

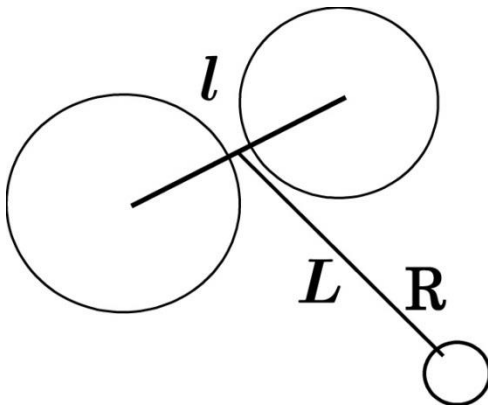
$$(\lambda\mu)=(04) \quad \ell=0,2,4^+ \quad \text{for } {}^{12}\text{C}$$

$$(\lambda\mu)=(80) \quad \ell=0,2,4,6,8^+ \quad \text{for } {}^{20}\text{Ne}$$



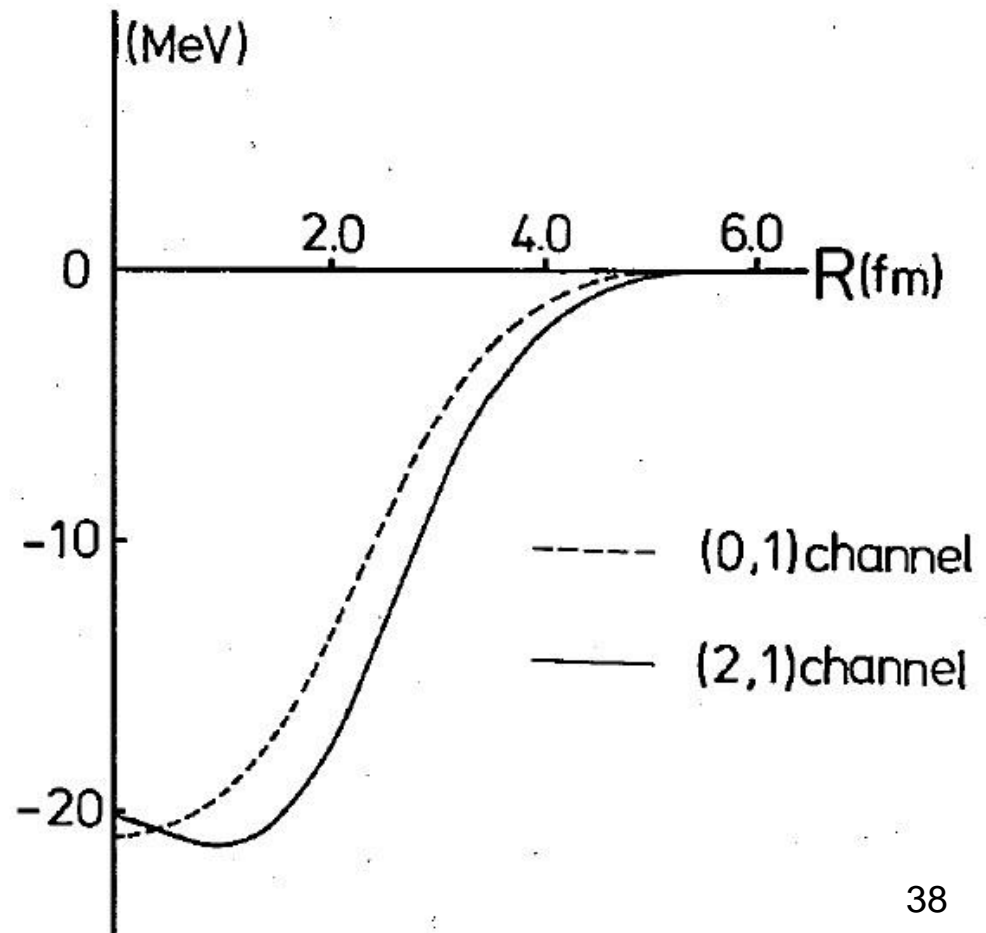
l -dependent folding potential

$$V_{lL,l'L'}(R) = \langle [\phi_l(\lambda\mu) Y_L(\hat{R})]_J | \sum_N v_{\Lambda N} | [\phi_{l'}(\lambda\mu) Y_{L'}(\hat{R})]_J \rangle$$



Diagonal potential
 $V_{lL}(R)$ for Λ p -state

$l=2$ is more
 attractive than $l=0$.



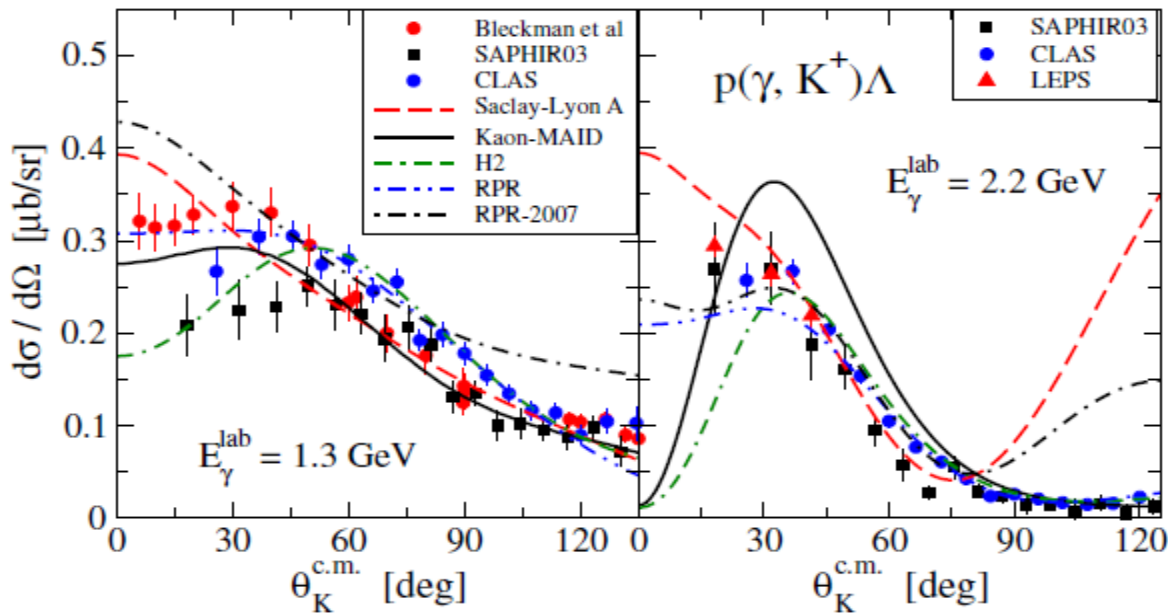
Possible test of $\gamma p \rightarrow \Lambda K$ ampl.

Comparison of isobar and Regge-plus-resonance models

H2: isobar model with hadronic f.f.; fit to CLAS data; nucleon resonances: $S_{11}(1650)$, $P_{11}(1710)$, $P_{13}(1720)$, $D_{13}(1895)$; hyperon resonances: $S_{01}(1670)$, $S_{01}(1800)$

RPR: fit to CLAS and LEPS data (cross sections) with resonances $S_{11}(1535)$, $S_{11}(1650)$, $P_{11}(1710)$, $P_{13}(1720)$, $D_{13}(1895)$; multidipole-Gauss hadronic f.f.;

motivated by RPR-2011B [Lesley De Cruz, PhD thesis, Ghent University, 2011]



RPR-2007:
Corthals et al, 2007,
version RPR-2+D13

P.B., M. Sotona, Nucl. Phys. A 835 (2010) 246

(From P. Bydzovsky)

Summary and outlook

- 1) Based on the $(\gamma, K+)$ reaction characteristics, typical physics contents are discussed by showing theoretical production X-sections.
- 2) Among others the DWIA predictions for p-shell, $^{28}_{\Lambda}\text{Al}$ and ^{52}Cr are well compared with the recent expt. ^{40}Ca and ^{208}Pb are also demonstrated.
- 3) In addition to the Λ s.p.e., dynamical coupling of Λ with collective nuclear motion is emphasized.
- 4) Remarks: For efforts to improve the theoretical elementary amplitudes, refer to Skoupil, Bydzovsky, Mart.