# Study of A-n Interaction via FSI in y+d Reaction

#### - Feasibility in the NKS2 Experiment -

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### Motivation

#### Study of nuclear force with strangeness

#### • Hypernuclear experiment

- Hadronic reaction
  - CERN-PS, BNL-AGS, KEK-PS, DA¢NE, J-PARC
- Electro-Magnetic reaction
  - JLab, MAMI
- Heavy ion reaction
  - GSI, BNL-RHIC, CERN-LHC

#### Hyperon-nucleon scattering

- Λ-p at Fermi Lab, Σ-p at J-PARC
- How about hyperon-n elastic scattering?



The first observation of a hypernucleus.



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→ Experimentally difficult



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  - $\rightarrow$  Experimentally difficult



The first observation of a hypernucleus.

#### We need the other data to understand NN, YN, YY force



• Using FSI in  $\gamma + d$  reaction for  $K^+ + \Lambda$  production



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 $\Lambda n$  interaction in Final State Interaction (FSI)



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 $\Lambda n$  interaction in Final State Interaction (FSI)

The other capability is  $K^-+d \rightarrow \gamma + \Lambda + n$ , but.....



### FSI Effect in the K<sup>+</sup> Cross-section

#### The shape of the curves

- Enhancement in forward K<sup>+</sup>
- variations: order of 10%
- Highly accurate measurements are required
  - in order to be able to distinguish among different potential models





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H. Yamamura et al., Phys. Rev C61 (1999) 014001



FIG. 2. The inclusive  $\gamma(d, K^+)$  cross section as a function of lab momenta  $p_K$  for  $\theta_K = 0^\circ$  and photon lab energy  $E_{\gamma} = 1.3$  GeV. The plane wave result is compared to two YN force predictions. The FSI effects are especially pronounced near the  $K^+\Lambda N$  and  $K^+\Sigma N$ thresholds, the locations of which are indicated by the arrows.

### Preceded Experiment



### Study about Ap Interaction

THE EUROPEAN

**PHYSICAL JOURNAL A** 

Eur. Phys. J. A **21**, 313–321 (2004) DOI 10.1140/epja/i2003-10203-3

#### Analysis of the $\Lambda p$ final-state interaction in the reaction $p+p \to K^+(\Lambda p)$

q

F. Hinterberger<sup>1,a</sup> and A. Sibirtsev<sup>2,3</sup>

#### 2.2 Final-state interaction

In the Watson-Migdal approximation [38–40] the FSI is taken into account by introducing a FSI enhancement factor  $|C_{\rm FSI}|^2$ ,

$$\frac{\mathrm{d}^2 \sigma}{\mathrm{d}\Omega_K \mathrm{d}M_{\Lambda p}} = |\mathcal{M}|^2 |C_{\mathrm{FSI}}|^2 \Phi_3, \qquad (3)$$

where now  $\mathcal{M}$  is a pure production matrix element and the FSI amplitude  $C_{\text{FSI}}$  depends on the internal momentum q of the  $\Lambda p$  subsystem. It converges to 1 for  $q \to \infty$  where the S-wave FSI enhancement vanishes.

$$C_{\rm FSI} = \frac{q - i\beta}{q + i\alpha}, \qquad |C_{\rm FSI}|^2 = \frac{q^2 + \beta^2}{q^2 + \alpha^2}. \tag{4}$$

The potential parameters  $\alpha$  and  $\beta$  can be used to establish phase-equivalent Bargmann potentials [44,45]. They are related to the scattering lengths a, and effective ranges rof the low-energy S-wave scattering

$$\alpha = \frac{1}{r} \left( 1 - \sqrt{1 - 2\frac{r}{a}} \right), \quad \beta = \frac{1}{r} \left( 1 + \sqrt{1 - 2\frac{r}{a}} \right). \quad (5)$$

The  $\Lambda p$  system can couple to singlet  ${}^{1}S_{0}$  and triplet  ${}^{3}S_{1}$  states. Near production threshold the singlet-triplet transitions due to the final-state interaction cannot occur. Therefore, the contributions of the spin-singlet and spin-triplet final states can be added incoherently. Taking the spin-statistical weights into account the unpolarized double differential cross-section may be written as

$$\frac{\mathrm{d}^{2}\sigma}{\mathrm{d}\Omega_{K}\mathrm{d}M_{Ap}} = \Phi_{3} \left[ 0.25 \, |\mathcal{M}_{s}|^{2} \, \frac{q^{2} + \beta_{s}^{2}}{q^{2} + \alpha_{s}^{2}} \right. \\ \left. + 0.75 \, |\mathcal{M}_{t}|^{2} \, \frac{q^{2} + \beta_{t}^{2}}{q^{2} + \alpha_{t}^{2}} \right].$$
(6)

This equation leaves six free parameters, the singlet and triplet potential parameters  $\alpha_s$ ,  $\beta_s$ ,  $\alpha_t$ ,  $\beta_t$  and the production matrix elements  $|\mathcal{M}_s|$  and  $|\mathcal{M}_t|$ . Instead of the parameters  $\alpha_s$ ,  $\beta_s$ ,  $\alpha_t$  and  $\beta_t$  one can equally well use the singlet and triplet scattering length and effective-range parameters  $a_s$ ,  $r_s$ ,  $a_t$  and  $r_t$ . The functional dependence on the invariant mass  $M_{\Lambda p}$  can be evaluated by inserting the corresponding expression for the internal momentum q of the  $\Lambda p$  system,

$$=\frac{\sqrt{M_{Ap}^{2}-(m_{A}+m_{p})^{2}}\sqrt{M_{Ap}^{2}-(m_{A}-m_{p})^{2}}}{2M_{Ap}}.$$
 (7)



Fig. 4. Same as in fig. 1. Solid lines: Fit curves with parameters given by eq. (15) from a combined five-parameter fit of the missing-mass spectrum and the total-cross-section data, dashed line: phase space distribution, dotted lines: singlet contributions, dash-dotted lines: triplet contributions.

#### Enhancement Factors of An FSI



	Refs.	a <sub>s</sub> [fm]	r <sub>s</sub> [fm]	a <sub>t</sub> [fm]	r <sub>t</sub> [fm]		
Nijmegen D	PRD15 (1977) 2547	-2.03	3.66	-1.84	3.32		
Nijmegen F	PRD20 (1979) 1633	-2.40	3.15	-1.84	3.37		
NSC89	PRC40 (1989) 2226	-2.86	2.91	-1.24	3.33		
NSC97a		-0.77	6.09	-2.15	2.71		
NSC97b		-0.97	5.09	-2.09	2.80		
NSC97c	PRC59 (1999) 3009	-1.28	4.22	-2.07	2.86		
NSC97d		-1.82	3.52	-1.94	3.01		
NSC97e		-2.24	3.24	-1.83	3.14		
NSC97f	1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	-2.68	3.07	-1.67	3.34		
Jülich A ( $\Lambda N$ )		-1.56	1.43	-1.59	3.16		
Jülich A~ ( $\Lambda N$ )	NIDA 570 (1004) 542	-2.04	0.64	-1.33	3.91		
Jülich B ( $\Lambda N$ )	NPA570 (1994) 545	-0.56	7.77	-1.91	2.43		
Jülich B~ ( $\Lambda N$ )		-0.40	12.28	-2.12	2.57		
Verma	PRC22 (1980) 229	-2.29	3.14	-1.77	3.25		
Bhaduri (Set I, $\Lambda N$ )	PR 155 (1967) 1671	-2.46	3.87	-2.07	4.50		



Note: It is assumed that the production matrix of single and triplet are the same.

NKS2 Experiment

(The recent results will be shown in Parallel 2b by H. Kanda)



### Photon Beam Line

in Research Center of Electron Photon Science (ELPH), Tohoku Univ.



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- Tagged photon beam
  - *E*γ = 0.80-1.08 GeV
- Liquid D<sub>2</sub> or H<sub>2</sub> target
- Magnetic spectrometer
  - Tracker
    - Two drift chambers
    - Charged particle momentum, trajectory, and decay vertex
  - Hodoscopes
    - Plastic scintillator + PMT
    - Time-Of-Flight (TOF)
    - Particle identification combined with momentum
  - Electron Veto

#### Acceptance





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• Covering large kinematic region including forward angle

0.42 T

Dipole

Magnet

Photon beam

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### Particle Identification



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### Particle Identification



# Feasibility in NKS2



### Feasibility Study by a Simulation

- FSI affects to cross section
  - Using effective range approximation
    - Large in low relative momentum region
    - How looks like in K and A distribution (p and  $\theta$ )?
- Acceptance
  - $K^+$  single, and  $K^++\Lambda$  coincidence
- Study by a Monte-Carlo simulation
  - GEANT4 based
  - $\gamma + d \rightarrow K^+ + \Lambda + n$ 
    - Kaon-MAID:  $\gamma + p \rightarrow K^+ \Lambda$
    - Applying the effect of Fermi motion inside of deuteron
  - New K<sup>+</sup> ID detector
    - same acceptance with current outer hodoscopes

#### K<sup>+</sup> angle vs. K<sup>+</sup> momentum



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the

#### K<sup>+</sup> Single Measurement (MC studied)



- 2 track trigger requested
- Geometrical acceptance for *K*<sup>+</sup>:
  - ~20-30% (varied by relative momentum)
  - note: ~1-2 % for  $K^++\Lambda$
- If requesting of < 3% error
  - in each 100 MeV step of Eγ bin
  - total number of  $K^+$  event: ~50000

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### Summary

#### • $\Lambda n$ interaction via FSI effect

•  $\gamma + d \rightarrow K^+ + \Lambda + n \text{ and } \gamma + p \rightarrow K^+ + \Lambda$ 

#### NKS2 experiment

- NKS2 spectrometer covers forward angle region
- Liquid Deuterium and Hydrogen target
- Strangeness photo-production near the threshold

#### • Feasibility of measurements of the $\Lambda n$ interaction

- Studied by the MC simulation based on GEANT4
  - K<sup>+</sup> single measurement
    - ~20-30% geometrical acceptance (for  $\Lambda n$  relative momentum <300 MeV/c)
    - ~150 days of beam time for 50000 event
  - $K^+ \Lambda$  coincidence measurement
    - ~1-3% geometrical acceptance (for  $\Lambda n$  relative momentum <300 MeV/c)
    - Capability of the complete measurement of kinematics
    - A recoil polarization may give us more information
      - need helps of theoretical study
      - •single/triplet separation?





## Scattering length and Effective range for Ap interaction

Model	$a_s$ (fm)	$r_s$ (fm)	$a_t$ (fm)	$r_t$ (fm)	$ \mathcal{M}_s ^2$ (b/sr)	$ \mathcal{M}_t ^2~~\mathrm{(b/sr)}$	$\chi^2$	$\chi^2/\text{n.d.f.}$
Nijm a	-0.71	5.86	-2.18	2.76	$0.\pm0.1$	$22.8\pm0.2$	47.6	1.22
Nijm b	-0.90	4.92	-2.13	2.84	$0.\pm 0.1$	$23.4\pm0.2$	53.9	1.38
Nijm c	-1.20	4.11	-2.08	2.92	$0.\pm 0.1$	$23.9\pm0.2$	62.2	1.60
Nijm d	-1.71	3.46	-1.95	3.08	$0.\pm 0.1$	$25.0\pm0.2$	83.3	2.14
Nijm e	-2.10	3.19	-1.86	3.19	$77.5\pm0.7$	$0.\pm 0.1$	76.6	1.97
Nijm f	-2.51	3.03	-1.75	3.32	$74.7\pm0.7$	$0.\pm 0.1$	44.8	1.15
Jül Ã	-2.04	0.64	-1.33	3.91	$7.3\pm0.5$	$7.8\pm1.6$	55.9	1.43
Jül $ ilde{\mathrm{B}}$	-0.40	12.28	-2.12	2.57	$0.\pm 0.1$	$21.3\pm0.2$	43.6	1.12
Jül A	-1.56	1.43	-1.59	3.16	$33.8\pm0.3$	$0.\pm 0.1$	80.2	2.06
Jül B	-0.56	7.77	-1.91	2.43	$0.\pm 0.1$	$20.1\pm0.2$	55.5	1.42
NSC	-2.78	2.88	-1.41	3.11	$71.6\pm0.7$	$0.\pm 0.1$	32.8	0.84
Nijm D	-1.90	3.72	-1.96	3.24	$0.\pm0.2$	$26.1\pm0.2$	91.6	2.35
Nijm F	-2.29	3.17	-1.88	3.36	$77.4\pm0.7$	$0.\pm 0.1$	62.6	1.61
Jül 03	-1.02	4.49	-1.89	2.57	$0.\pm0.1$	$21.3\pm0.2$	62.4	1.60

Table 2. Test of potential model results performed with two-parameter fit.

EPJ A21 (2004) 313

Production matrix elements of  $\Lambda p$  Singlet and Triplet state are varied by model

#### Effective Range and Scattering Length for An interaction

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