

エクスクルーシヴ Drell-Yan 過程  
 $\pi^- p \rightarrow \ell^+ \ell^- n$  のQCDメカニズム

田中和廣 (順天堂大/KEK)

# High momentum beam line at J-PARC

- Primary beam (proton)

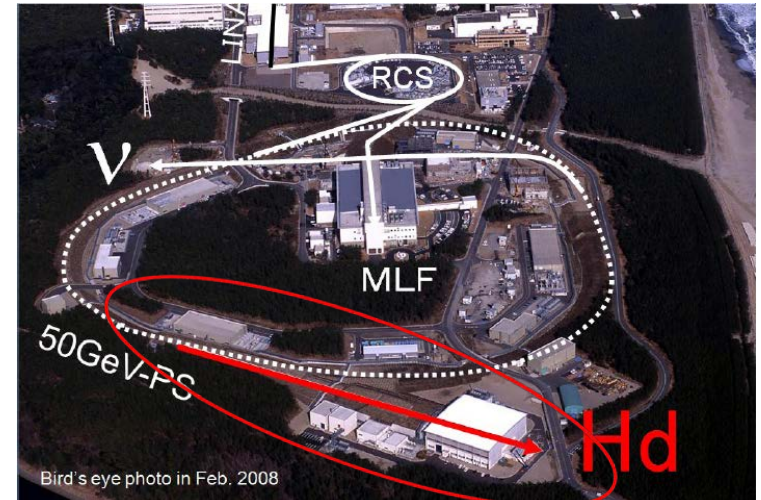
$$E = 30\text{GeV} \text{ (} \rightarrow 50\text{GeV?)}$$

$$L = 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$$

- ↔ PANDA (anti-proton)

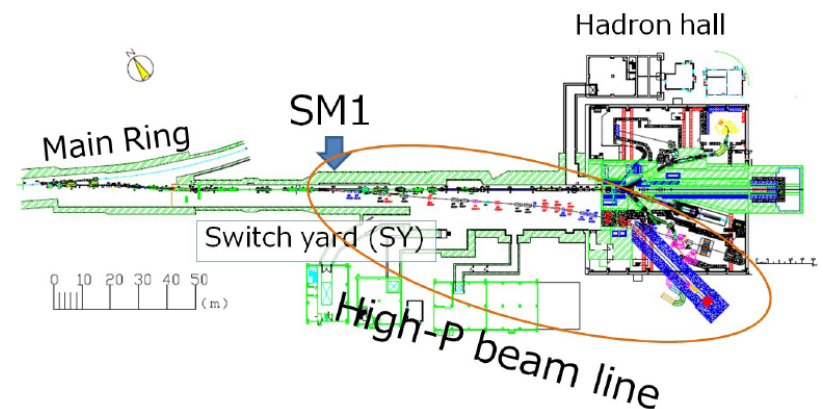
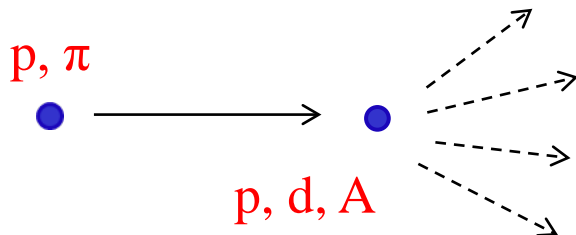
$$E \leq 15\text{GeV}, L = 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$$

## Hadron Facility at J-PARC



- Secondary beam (pion)

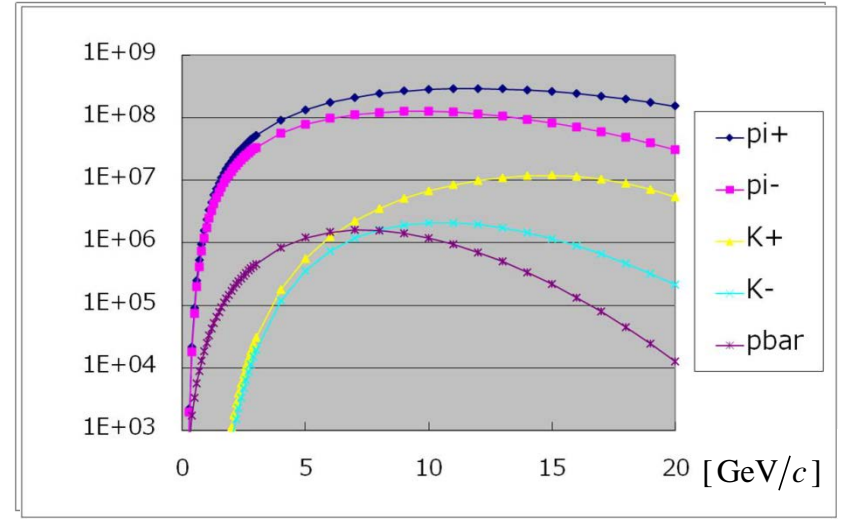
$$E = 15\text{-}20\text{GeV}$$





beam loss limit @ SM1:15kW

(limited by the thickness of the tunnel wall)



0° extraction angle

## High-momentum beamline

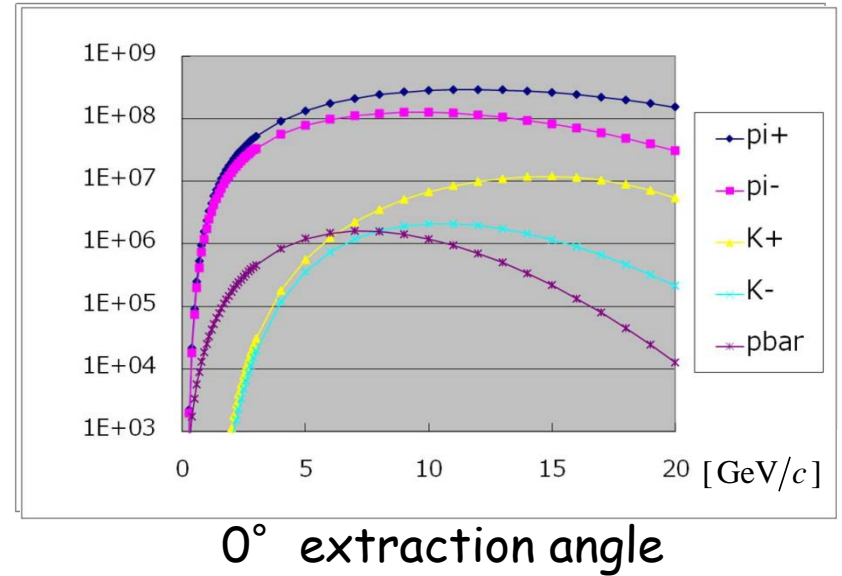
- 30 GeV proton
- ~15-20 GeV unseparated (mainly pions)

high intensity



beam loss limit @ SM1:15kW

(limited by the thickness of the tunnel wall)



# High-momentum beamline

- 30 GeV proton
- ~15-20 GeV unseparated (mainly pions)

high intensity

not too high energy

$$d\sigma \sim 1/s^a$$

best suited to study meson-induced  
hard exclusive processes

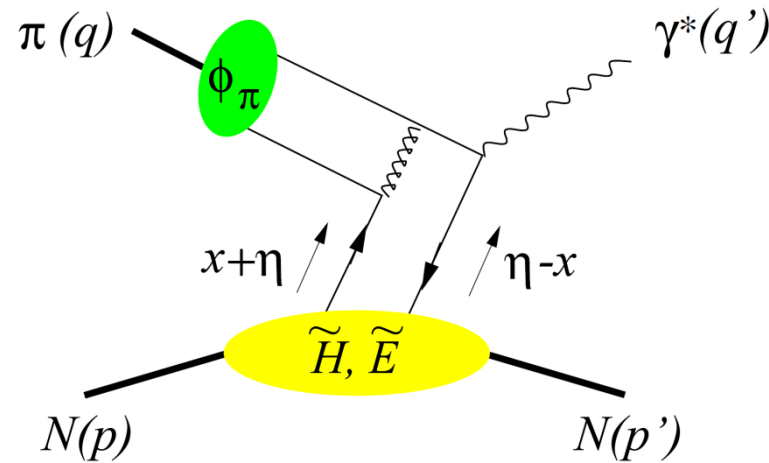
# Exclusive lepton pair production in $\pi N$ scattering

$$\pi^- p \rightarrow \gamma^* n \rightarrow \mu^+ \mu^- n$$

Berger, Diehl, Pire, PLB523(2001)265

“exclusive limit of DY”

small  $t = (q - q')^2$



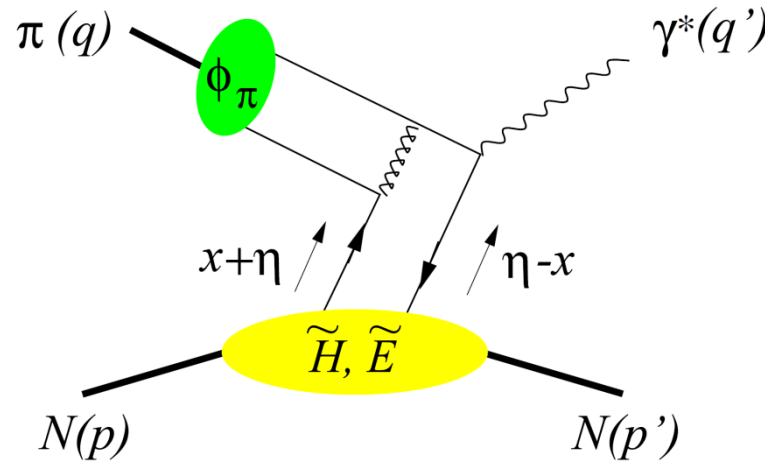
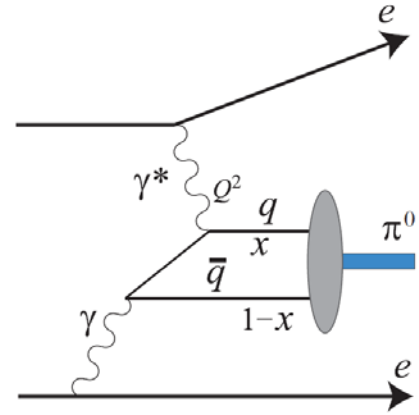
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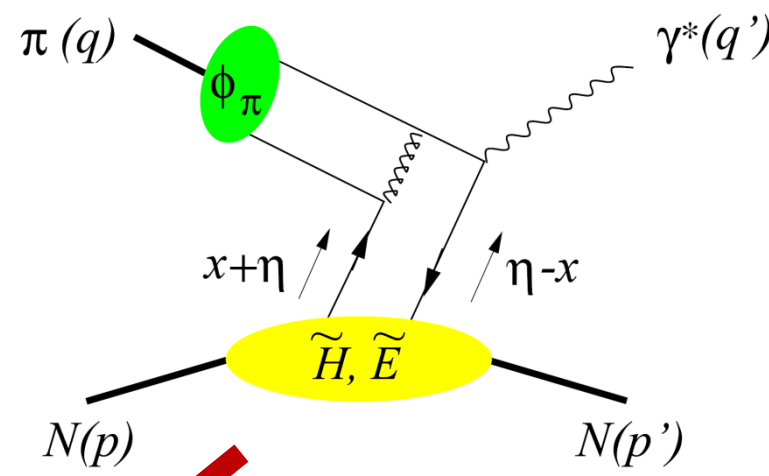
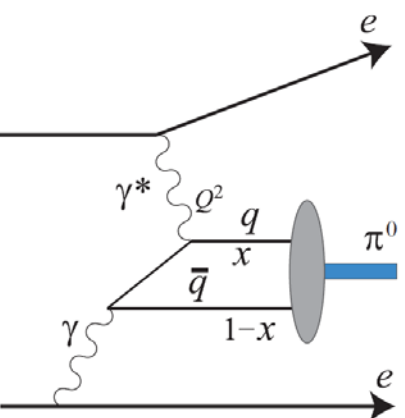
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small  $t = (q - q')^2$

$\Delta q(x)$   $\leftarrow$   $t \rightarrow 0$

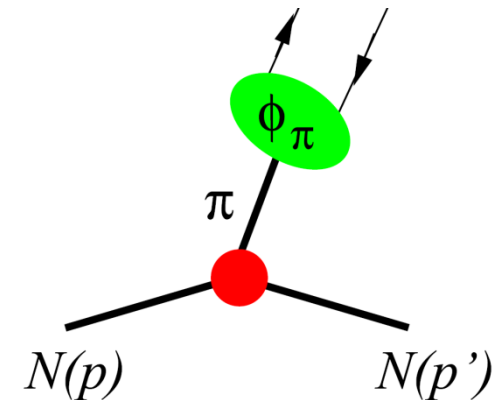
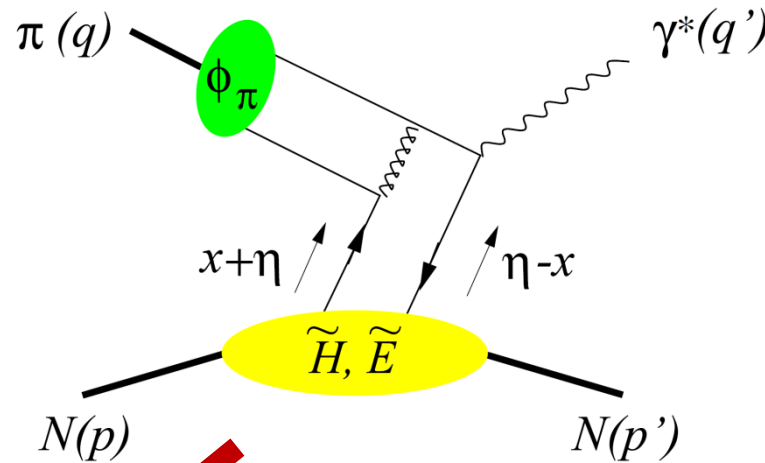
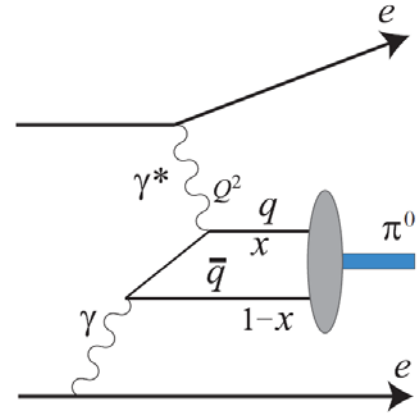
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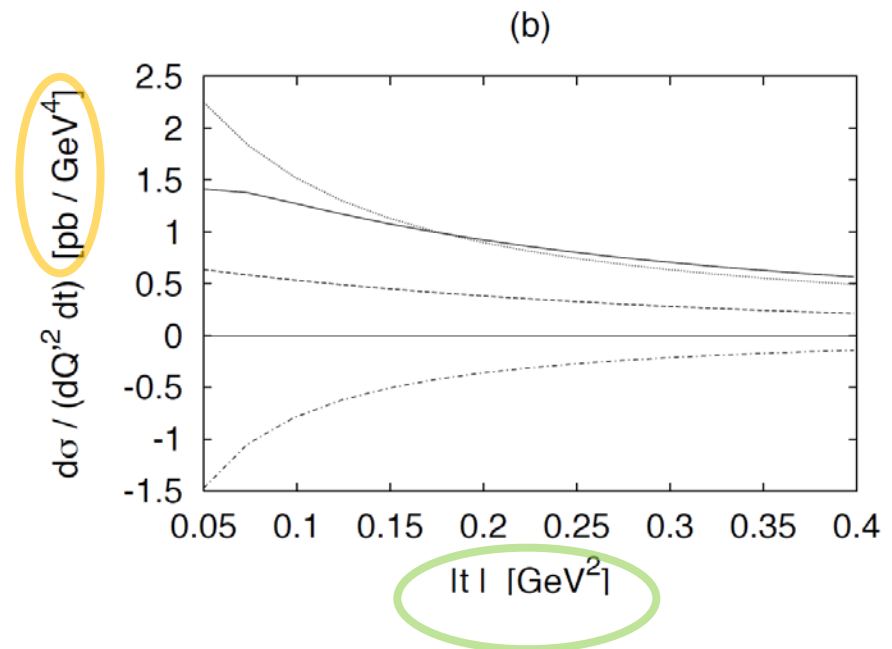
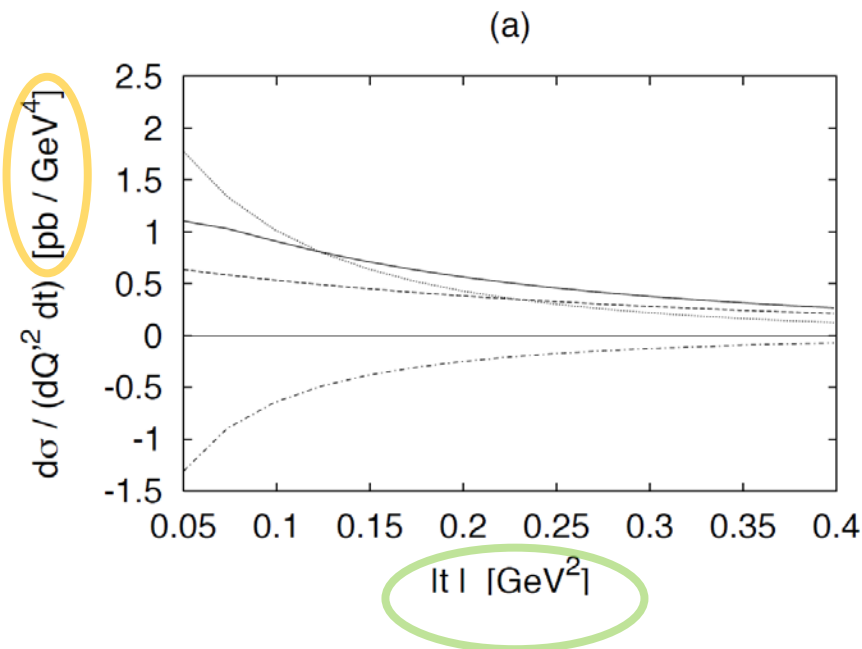


# LO Estimates

Bjorken variable  $\tau = \frac{Q'^2}{s-M^2}$

Berger, Diehl, Pire, PLB523(2001)265

$Q'^2 = 5\text{GeV}^2$        $\tau = 0.2$



(dashed) =  $|\tilde{\mathcal{H}}|^2$  ; (dash-dotted) =  $\text{Re}(\tilde{\mathcal{H}}^* \tilde{\mathcal{E}})$  ; (dotted) =  $|\tilde{\mathcal{E}}|^2$

$$\frac{d\sigma}{dQ'^2 dt}(\pi^- p \rightarrow \gamma^* n) = \frac{4\pi\alpha_{\text{em}}^2 \tau^2}{27 Q'^8} f_\pi^2 \left[ (1-\eta^2) |\tilde{\mathcal{H}}^{du}|^2 - 2\eta^2 \text{Re}(\tilde{\mathcal{H}}^{du*} \tilde{\mathcal{E}}^{du}) - \eta^2 \frac{t}{4M^2} |\tilde{\mathcal{E}}^{du}|^2 \right]$$

# LO Estimates

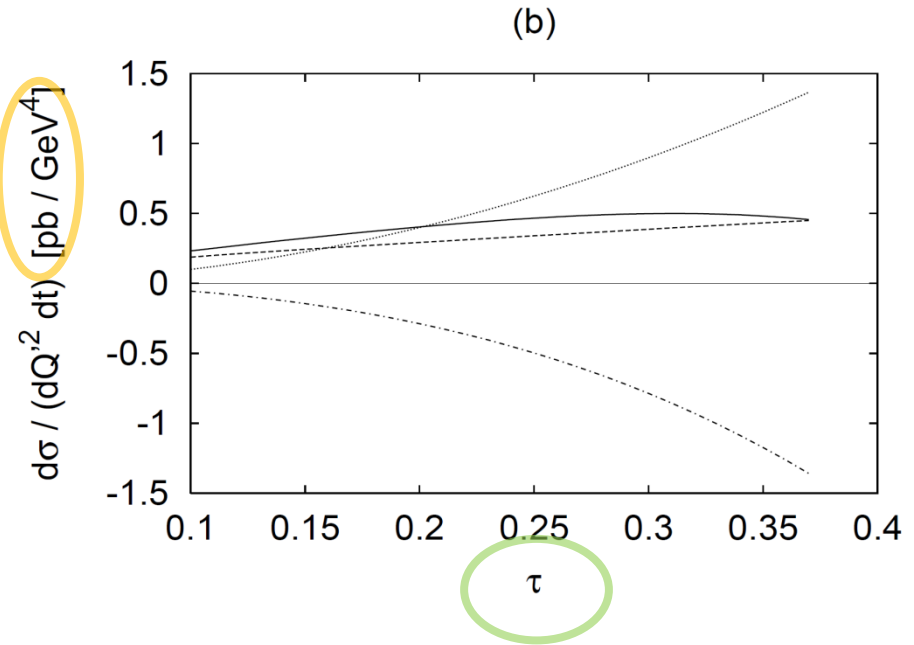
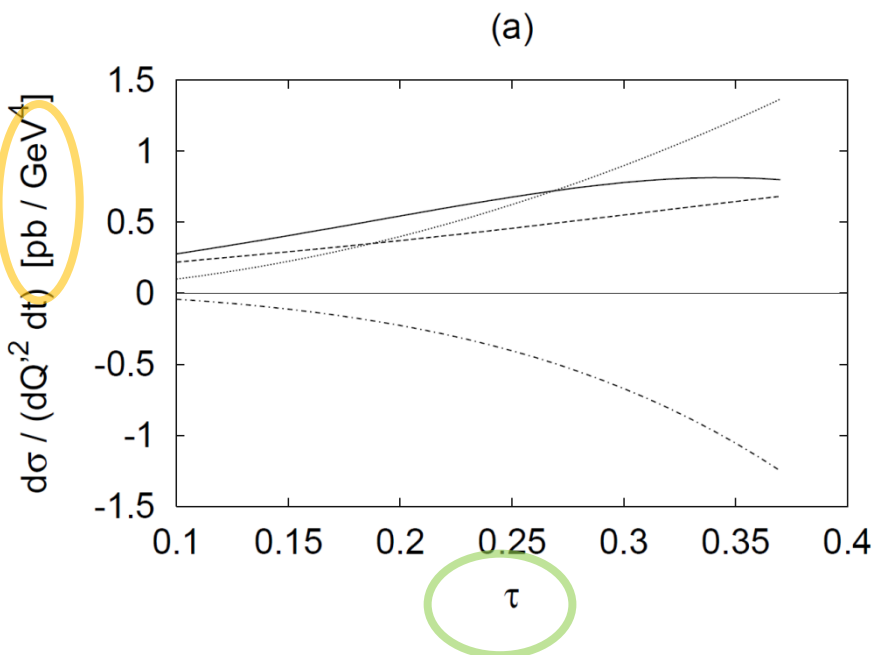
Bjorken variable

$$\tau = \frac{Q'^2}{s-M^2}$$

Berger, Diehl, Pire, PLB523(2001)265

$$Q'^2 = 5 \text{ GeV}^2$$

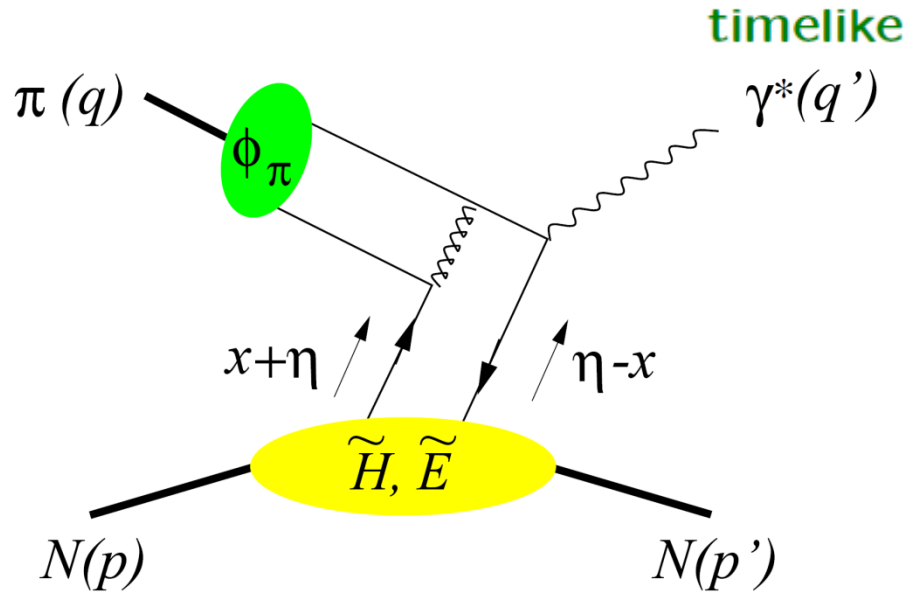
$$|t| = 0.2$$



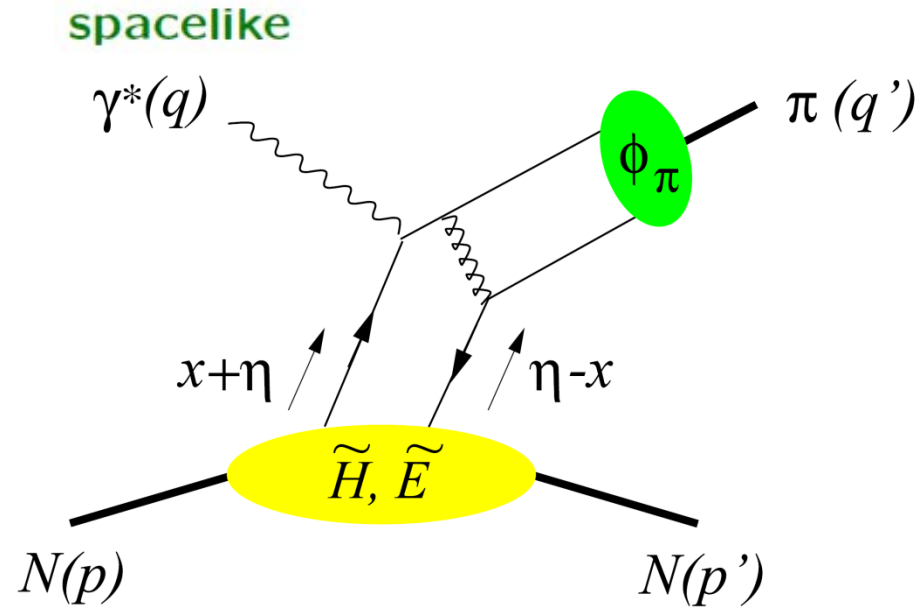
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# Pion beams reveal $\tilde{H}, \tilde{E}$ Generalized Parton distributions

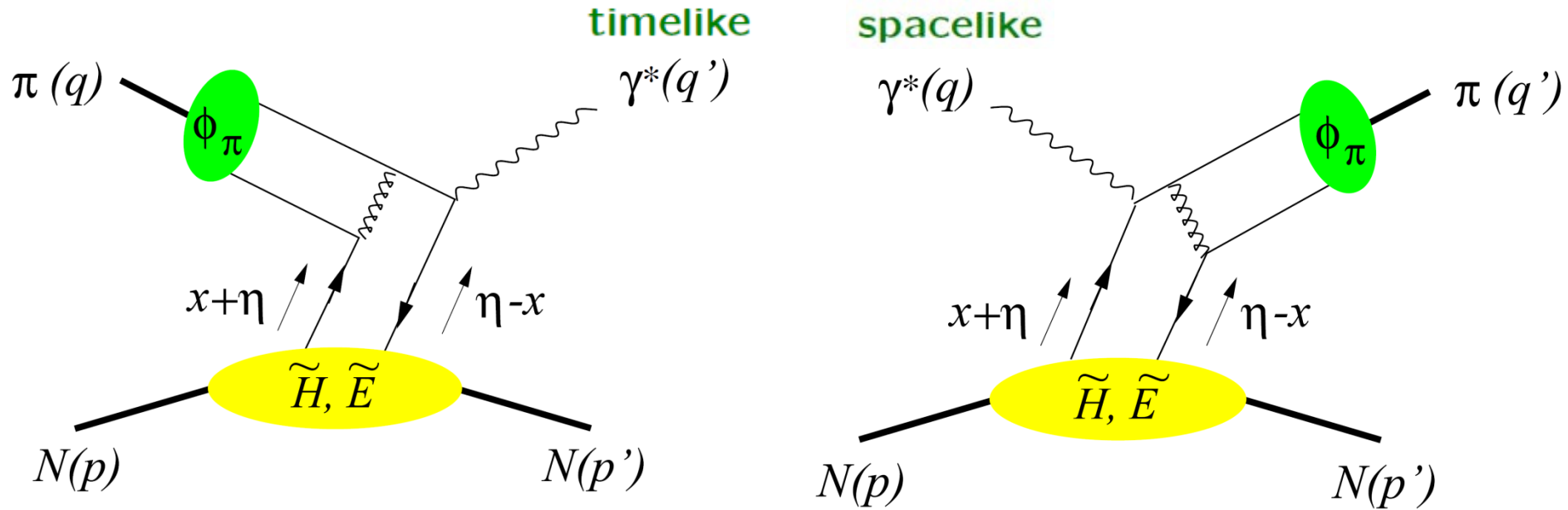


**exDY@J-PARC**



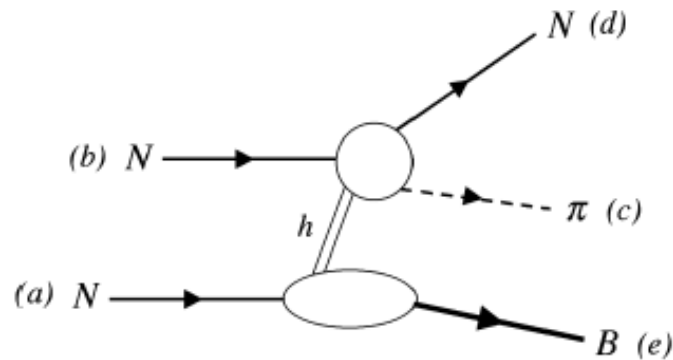
**DVMP@JLab**

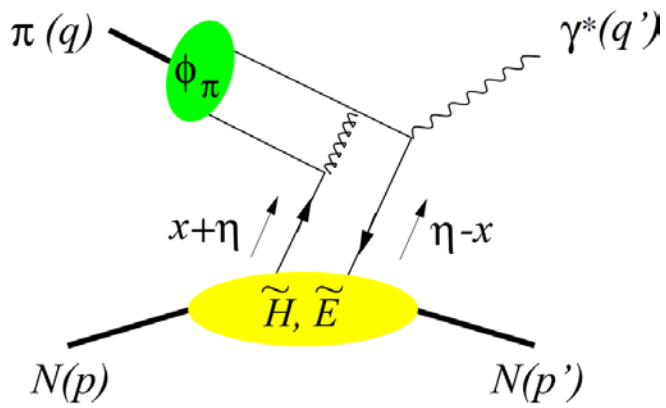
# Pion beams reveal $\tilde{H}, \tilde{E}$ Generalized Parton distributions



**exDY@J-PARC**

**DVMP@JLab**





**Bjorken variable:**  $\tau = \frac{Q'^2}{2p \cdot q}$

**Skewness:**  $\eta = \frac{p^+ - p'^+}{p^+ + p'^+}$

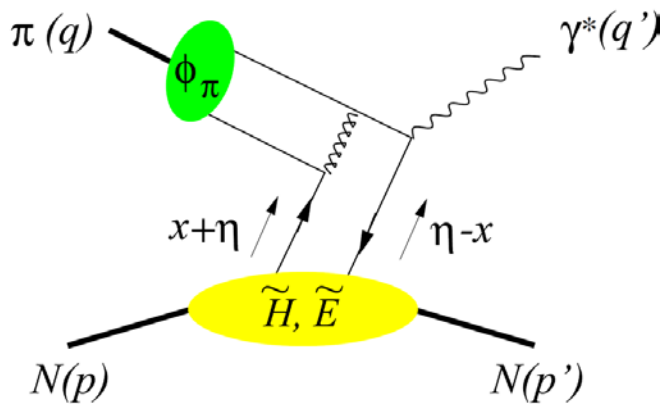
**Berger, Diehl, Pire, PLB523(2001)**

$$\frac{d\sigma}{dQ'^2 dt d(\cos\theta) d\varphi} = \frac{\alpha_{em}}{256 \pi^3} \frac{\tau^2}{Q'^6} \sum_{\lambda', \lambda} |M^{0\lambda', \lambda}|^2 \sin^2 \theta$$

$$M^{0\lambda', \lambda}(\pi^- p \rightarrow \gamma^* n) = -ie \frac{4\pi}{3} \frac{f_\pi}{Q'} \frac{1}{(p+p')^+} \bar{u}(p', \lambda') \left[ \gamma^+ \gamma_5 \tilde{\mathcal{H}}^{du}(\eta, t) + \gamma_5 \frac{(p'-p)^+}{2M} \tilde{\mathcal{E}}^{du}(\eta, t) \right] u(p, \lambda)$$

$$\tilde{\mathcal{H}}^{du}(\eta, t) = \frac{8\alpha_S}{3} \int_0^1 du \frac{\phi_\pi(u)}{4u(1-u)} \int_{-1}^1 dx \left[ \frac{e_d}{-\eta-x-i\epsilon} - \frac{e_u}{-\eta+x-i\epsilon} \right] [\tilde{H}^d(x, \eta, t) - \tilde{H}^u(x, \eta, t)]$$

$$\int \frac{dz^-}{2\pi} e^{ix\bar{P}^+ z^-} \langle p' | \bar{\psi}(-\frac{z^-}{2}) \gamma^+ \gamma_5 \psi(\frac{z^-}{2}) | p \rangle = \frac{1}{\bar{P}^+} \left[ \tilde{H}^q(x, \eta, t) \bar{u}(p') \gamma^+ \gamma_5 u(p) + \tilde{E}^q(x, \eta, t) \bar{u}(p') \frac{\gamma_5 (p'-p)^+}{2M} u(p) \right]$$



**Bjorken variable:**  $\tau = \frac{Q'^2}{2p \cdot q}$

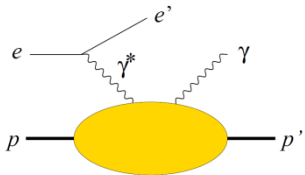
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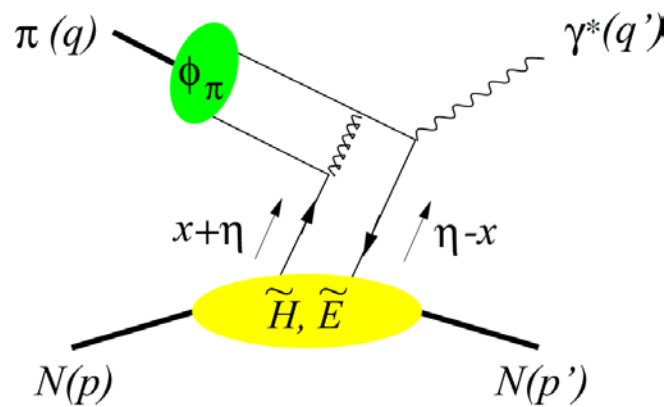
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$$\int \frac{dz^-}{2\pi} e^{ix\bar{P}^+ z^-} \langle p' | \bar{\psi}(-\frac{z^-}{2}) \gamma^+ \gamma_5 \psi(\frac{z^-}{2}) | p \rangle = \frac{1}{\bar{P}^+} \left[ \tilde{H}^q(x, \eta, t) \bar{u}(p') \gamma^+ \gamma_5 u(p) + \tilde{E}^q(x, \eta, t) \bar{u}(p') \frac{\gamma_5 (p'-p)^+}{2M} u(p) \right]$$



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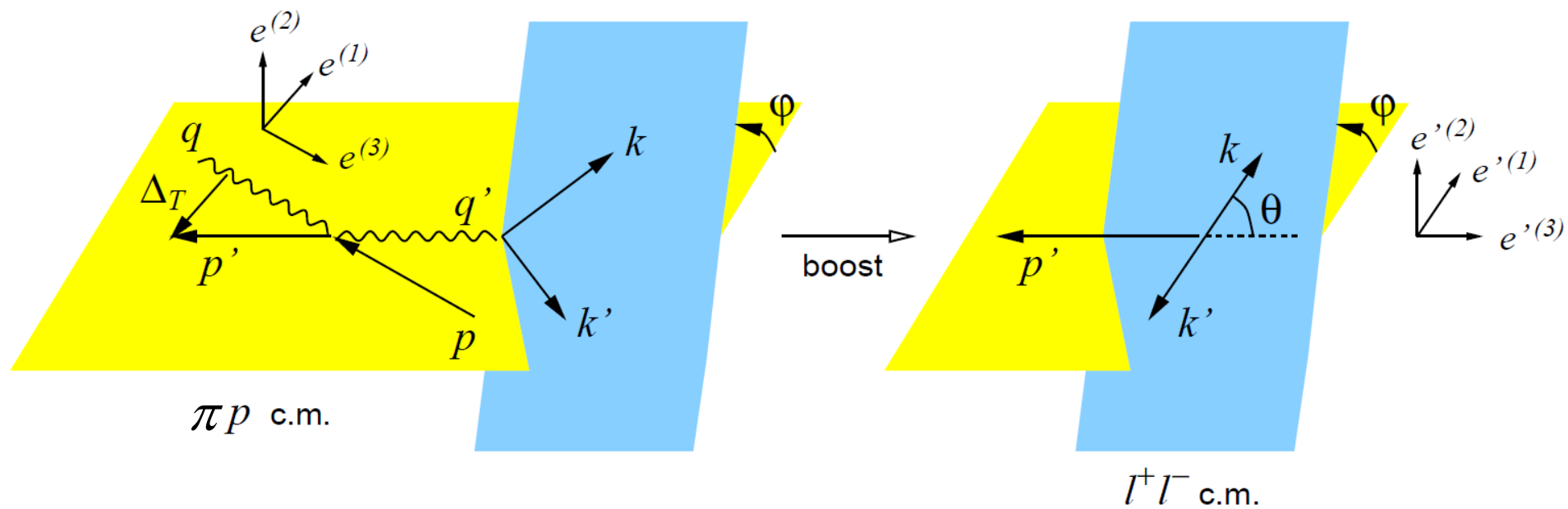
$$J_q = \frac{1}{2} \int_{-1}^1 dx x (H^q(x, \eta, 0) + E^q(x, \eta, 0))$$

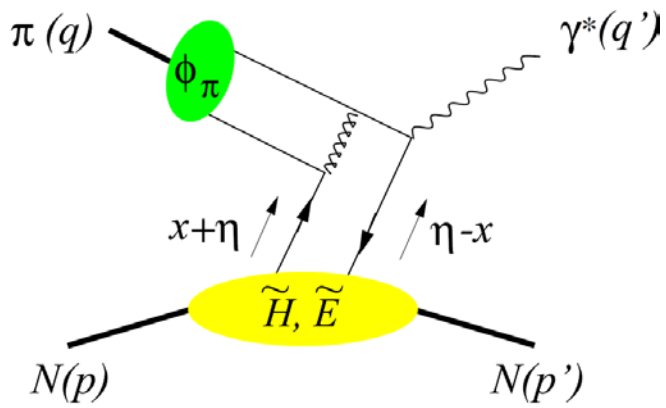


**Bjorken variable:**  $\tau = \frac{Q'^2}{2p \cdot q}$

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$$\frac{d\sigma}{dQ'^2 dt d(\cos\theta) d\varphi} = \frac{\alpha_{em}}{256 \pi^3} \frac{\tau^2}{Q'^6} \sum_{\lambda', \lambda} |M^{0\lambda', \lambda}|^2 \sin^2 \theta$$





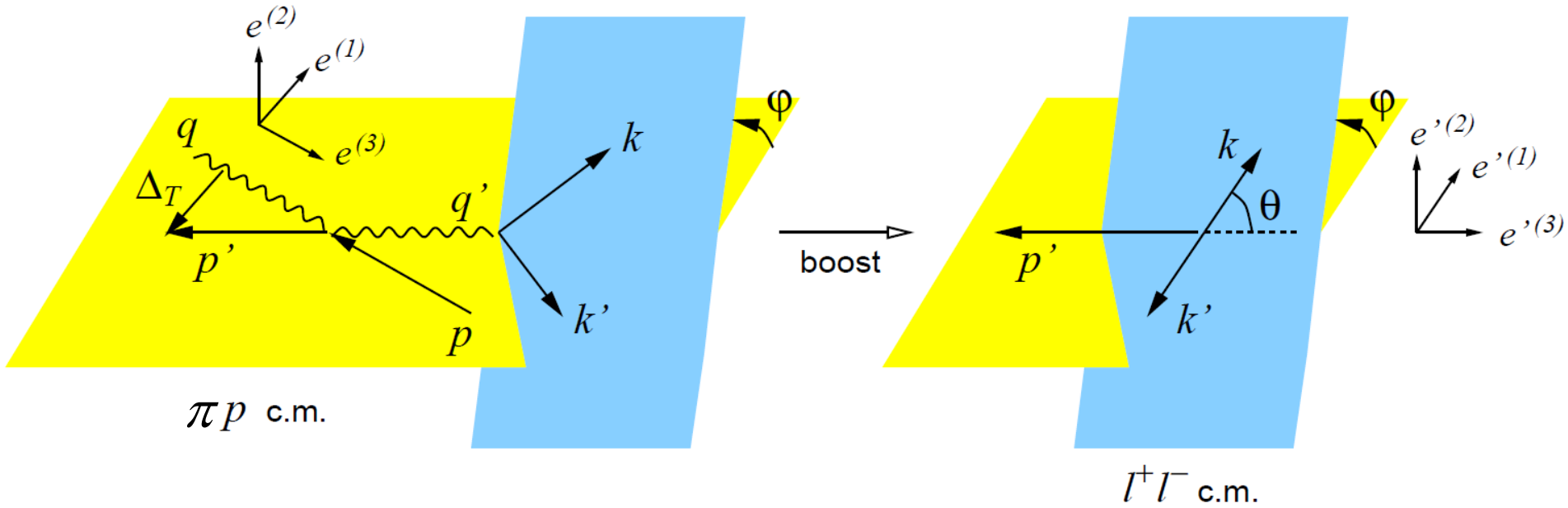
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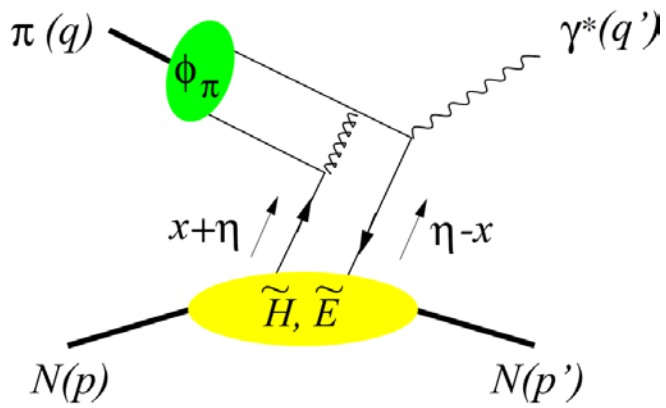
**long. photon**

$$|d_{-10}^1(\theta)|^2 + |d_{10}^1(\theta)|^2$$

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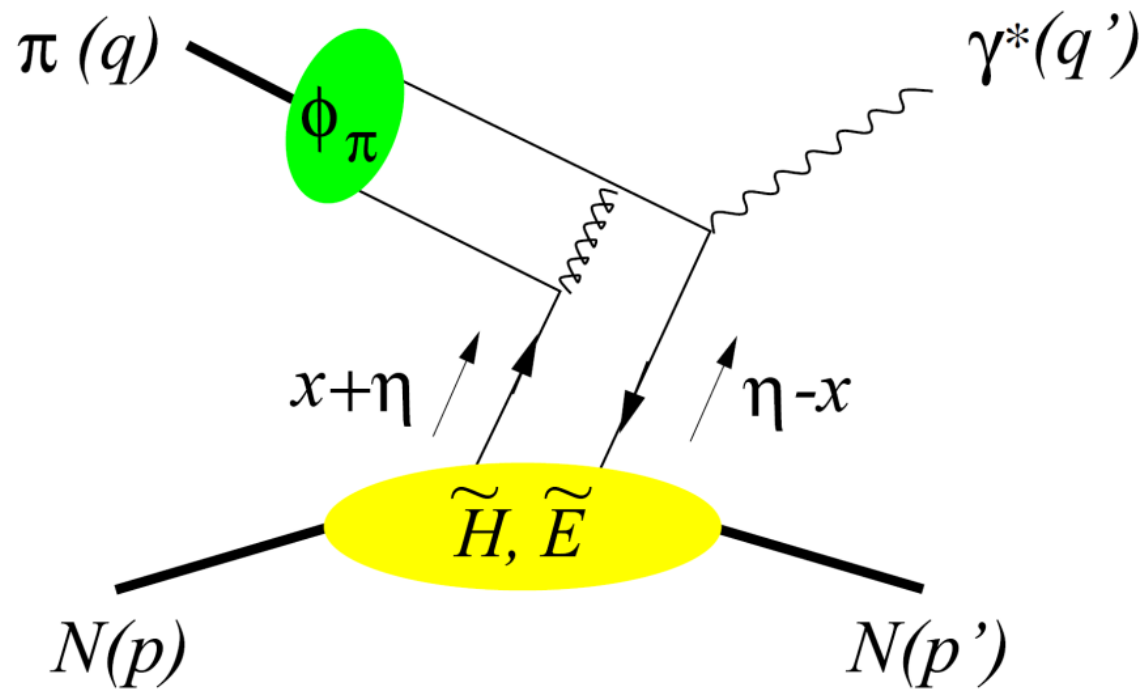
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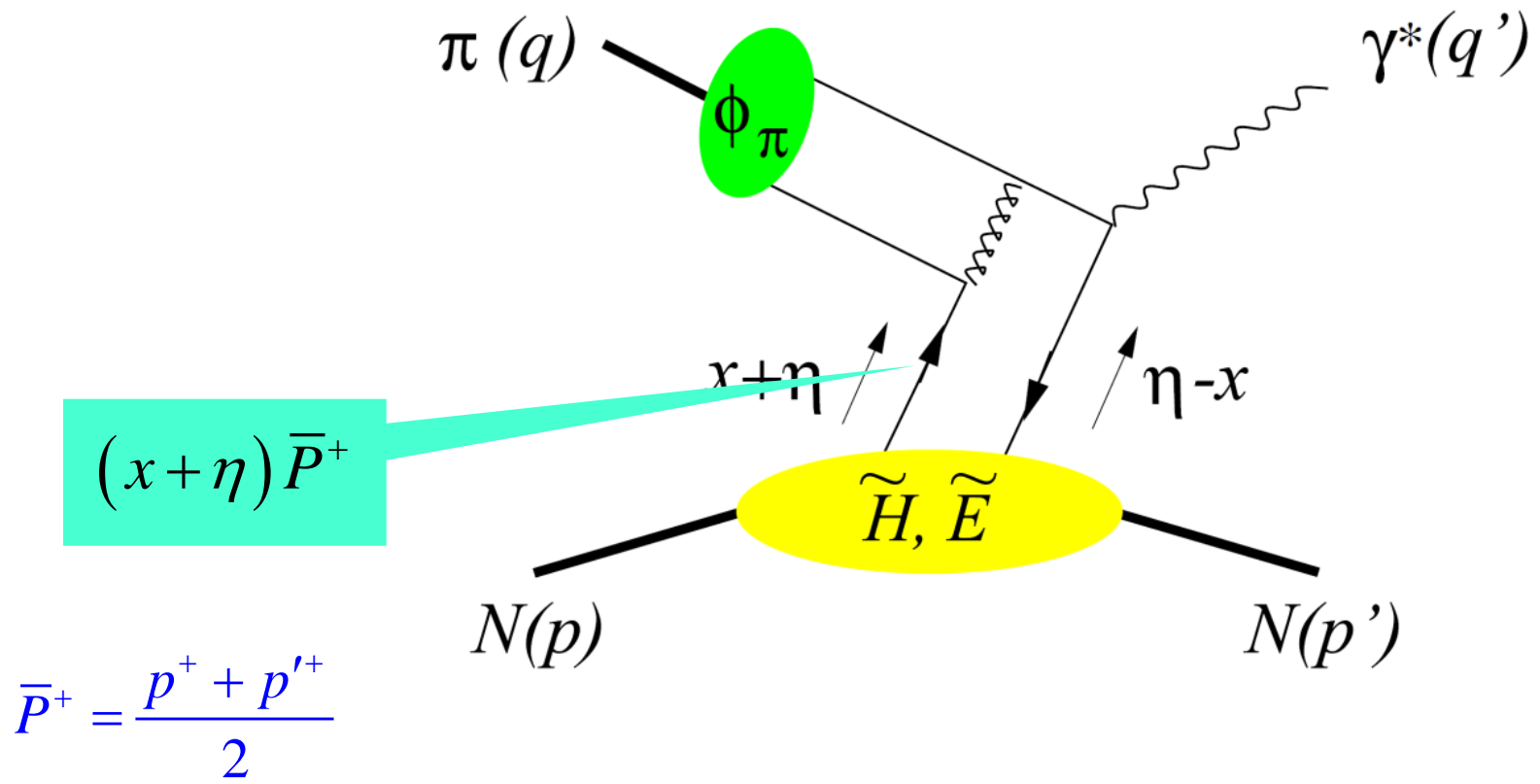
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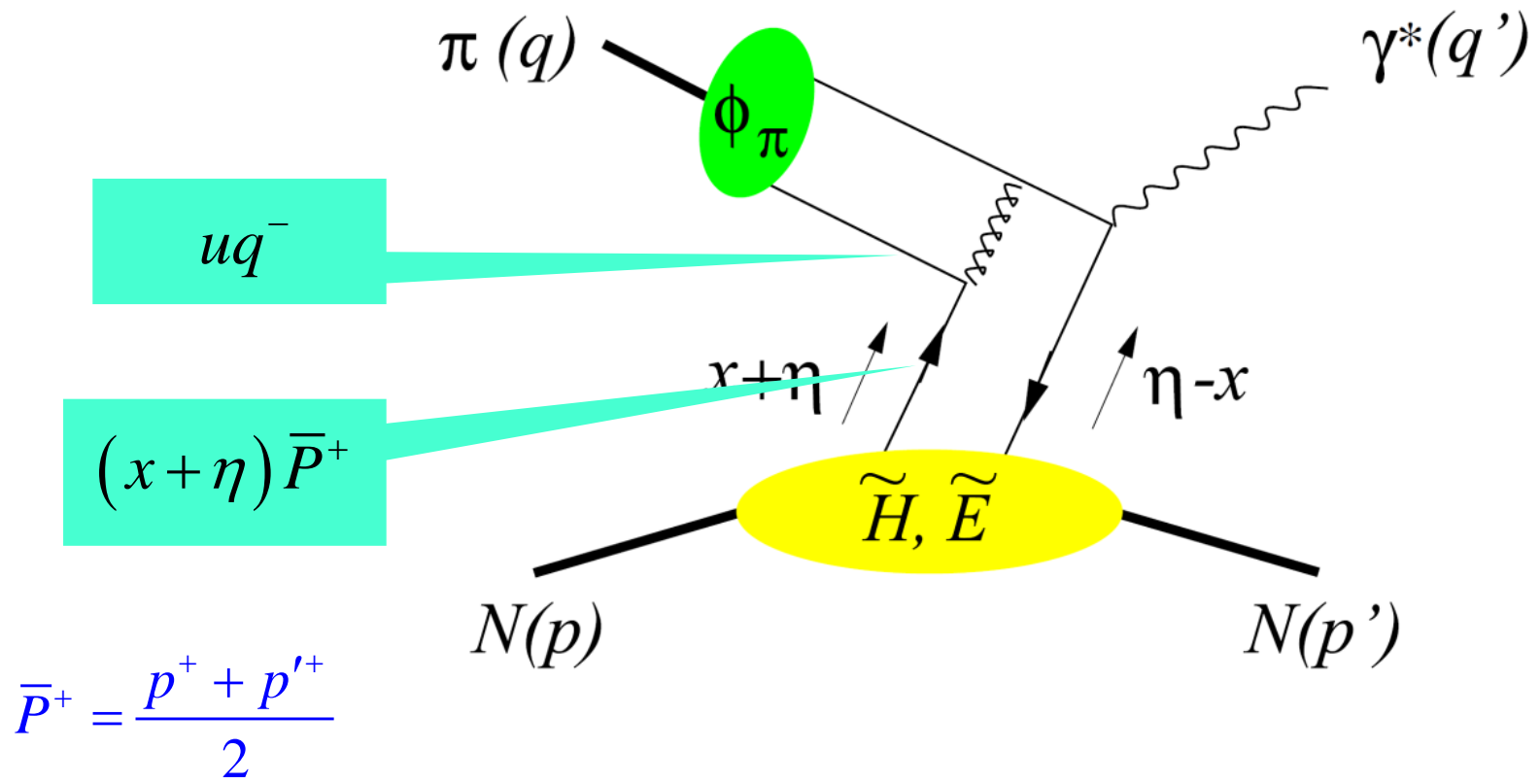
$$M^{0\lambda', \lambda}(\pi^- p \rightarrow \gamma^* n) = -ie \frac{4\pi}{3} \frac{f_\pi}{Q'} \frac{1}{(p+p')^+} \bar{u}(p', \lambda') \left[ \gamma^+ \gamma_5 \tilde{\mathcal{H}}^{du}(\eta, t) + \gamma_5 \frac{(p'-p)^+}{2M} \tilde{\mathcal{E}}^{du}(\eta, t) \right] u(p, \lambda)$$

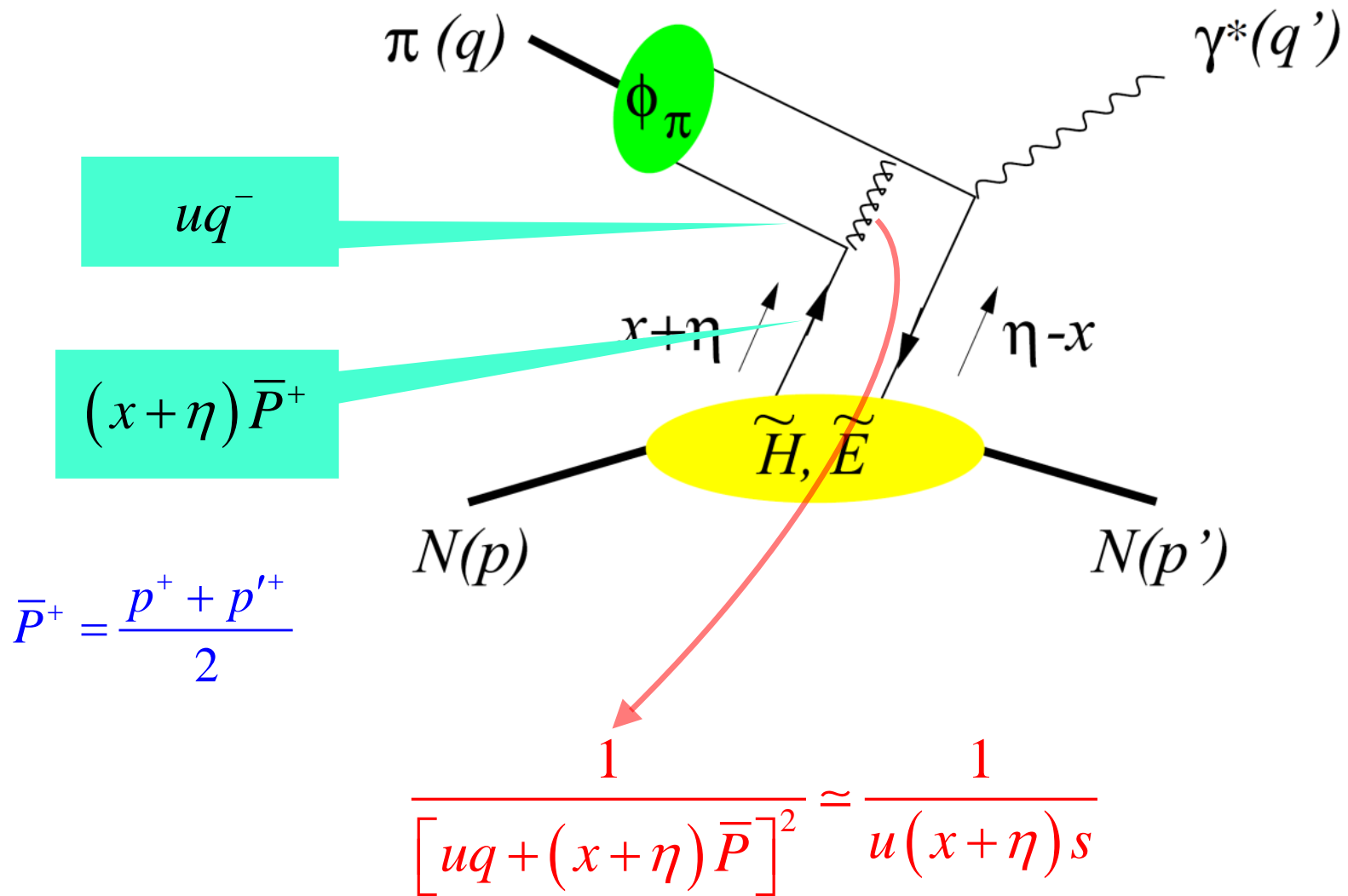
$$\tilde{\mathcal{H}}^{du}(\eta, t) = \frac{8\alpha_S}{3} \int_0^1 du \frac{\phi_\pi(u)}{4u(1-u)} \int_{-1}^1 dx \left[ \frac{e_d}{-\eta-x-i\epsilon} - \frac{e_u}{-\eta+x-i\epsilon} \right] [\tilde{H}^d(x, \eta, t) - \tilde{H}^u(x, \eta, t)]$$

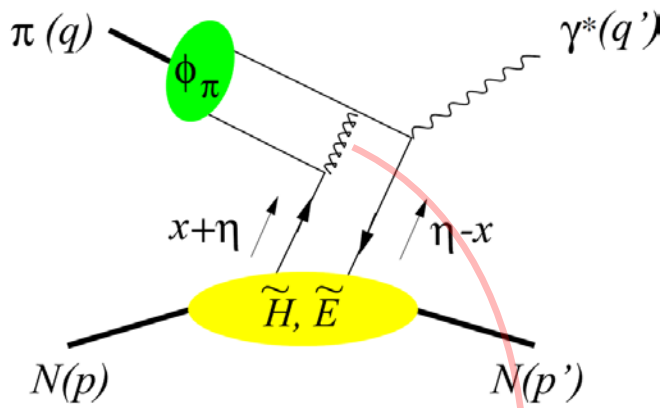
$$\int \frac{dz^-}{2\pi} e^{ix\bar{P}^+z^-} \langle p' | \bar{\psi}(-\frac{z^-}{2}) \gamma^+ \gamma_5 \psi(\frac{z^-}{2}) | p \rangle = \frac{1}{\bar{P}^+} \left[ \tilde{H}^q(x, \eta, t) \bar{u}(p') \gamma^+ \gamma_5 u(p) + \tilde{E}^q(x, \eta, t) \bar{u}(p') \frac{\gamma_5 (p'-p)^+}{2M} u(p) \right]$$











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**long. photon**

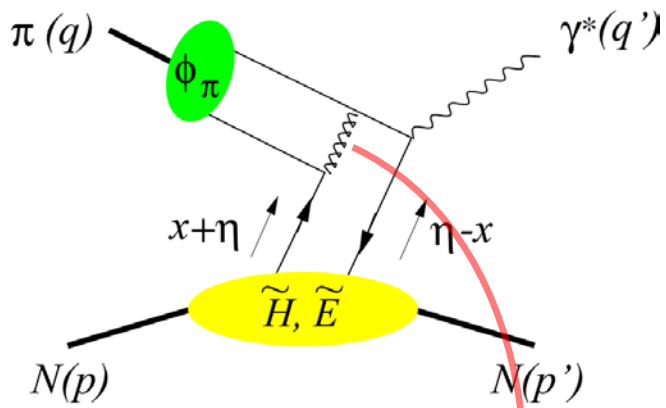
$$|d_{-1 0}^1(\theta)|^2 + |d_{1 0}^1(\theta)|^2$$

$$\frac{d\sigma}{dQ'^2 dt d(\cos\theta) d\varphi} = \frac{\alpha_{em}}{256 \pi^3} \frac{\tau^2}{Q'^6} \sum_{\lambda', \lambda} |M^{0\lambda', \lambda}|^2 \sin^2 \theta$$

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**long. photon**

$|d_{-10}^1(\theta)|^2 + |d_{10}^1(\theta)|^2$

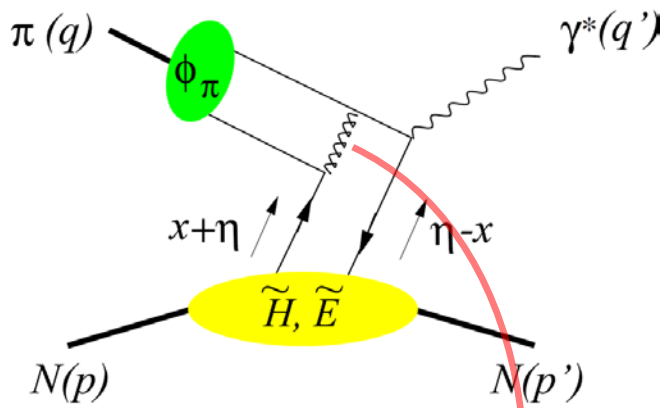
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$$\phi_\pi(u) \sim u(1-u)$$



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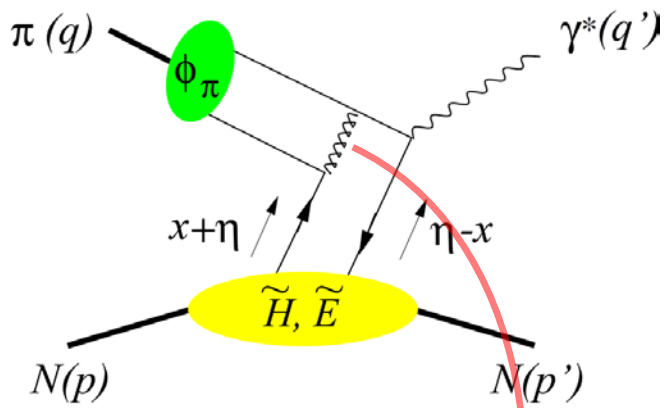
$$\tilde{H}^{du}(\eta, t) = \frac{8\alpha_S}{3} \int_0^1 du \frac{\phi_\pi(u)}{4u(1-u)} \int_{-1}^1 dx \left[ \frac{e_d}{-\eta-x-i\epsilon} - \frac{e_u}{-\eta+x-i\epsilon} \right] [\tilde{H}^d(x, \eta, t) - \tilde{H}^u(x, \eta, t)]$$

$$\int \frac{dz^-}{2\pi} e^{ix\bar{P}^+ z^-} \langle p' | \bar{\psi}(-\frac{z^-}{2}) \gamma^+ \gamma_5 \psi(\frac{z^-}{2}) | p \rangle = \frac{1}{\bar{P}^+} \left[ \tilde{H}^q(x, \eta, t) \bar{u}(p') \gamma^+ \gamma_5 u(p) + \tilde{E}^q(x, \eta, t) \bar{u}(p') \frac{\gamma_5 (p'-p)^+}{2M} u(p) \right]$$

$$\phi_\pi(u) \sim u(1-u)$$

$$M^{\pm 1, \lambda; \lambda'}(\pi^- p \rightarrow \gamma^* n) \sim \frac{\alpha_s}{Q'^2} \int_0^1 du \frac{\phi_p(u)}{u(1-u)} \otimes \frac{1}{(\eta \pm x + i\epsilon)^2} \otimes \{H_T^q(x, \eta, t), \tilde{H}_T^q(x, \eta, t), E_T^q(x, \eta, t), \tilde{E}_T^q(x, \eta, t)\}$$





**Bjorken variable:**  $\tau = \frac{Q'^2}{2p \cdot q}$

**Skewness:**  $\eta = \frac{p^+ - p'^+}{p^+ + p'^+}$

**long. photon**

$$|d_{-10}^1(\theta)|^2 + |d_{10}^1(\theta)|^2$$

$$\frac{d\sigma}{dQ'^2 dt d(\cos\theta) d\varphi} = \frac{\alpha_{em}}{256\pi^3} \frac{\tau^2}{Q'^6} \sum_{\lambda', \lambda} |M^{0\lambda', \lambda}|^2 \sin^2\theta$$

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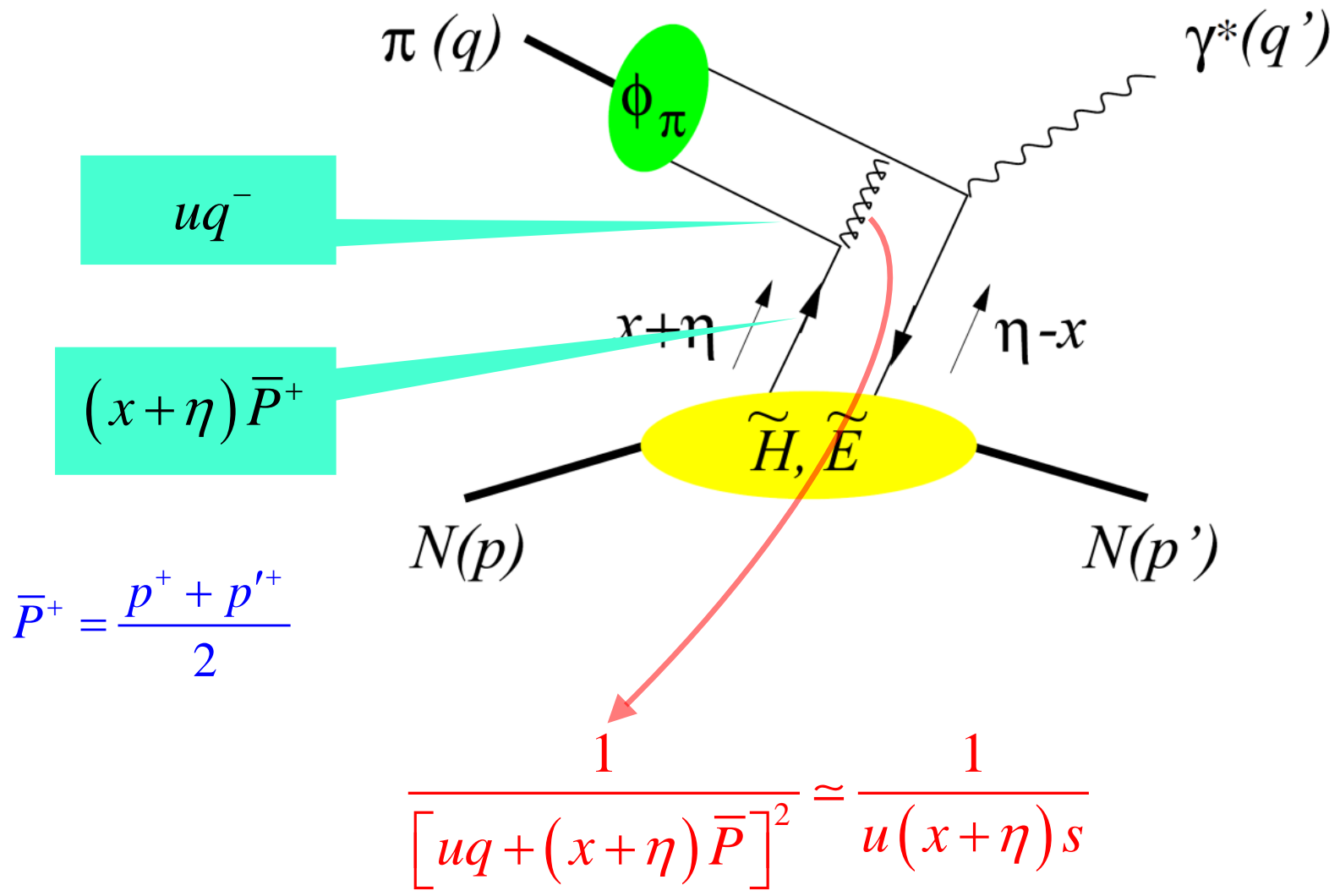
$$\tilde{H}^{du}(\eta, t) = \frac{8\alpha_S}{3} \int_0^1 du \frac{\phi_\pi(u)}{4u(1-u)} \int_{-1}^1 dx \left[ \frac{e_d}{-\eta-x-i\epsilon} - \frac{e_u}{-\eta+x-i\epsilon} \right] [\tilde{H}^d(x, \eta, t) - \tilde{H}^u(x, \eta, t)]$$

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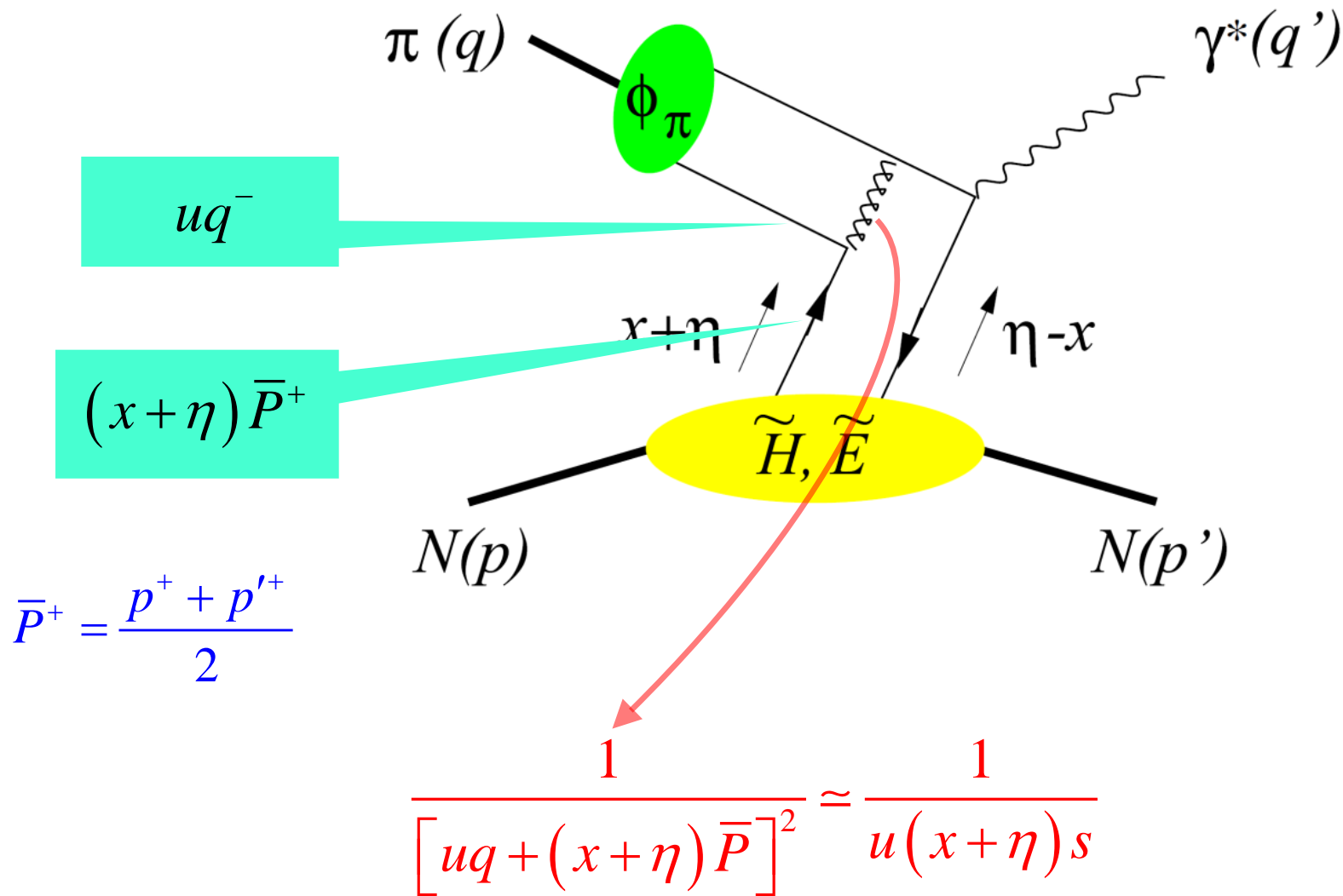
$$\phi_p(u) \sim 1$$

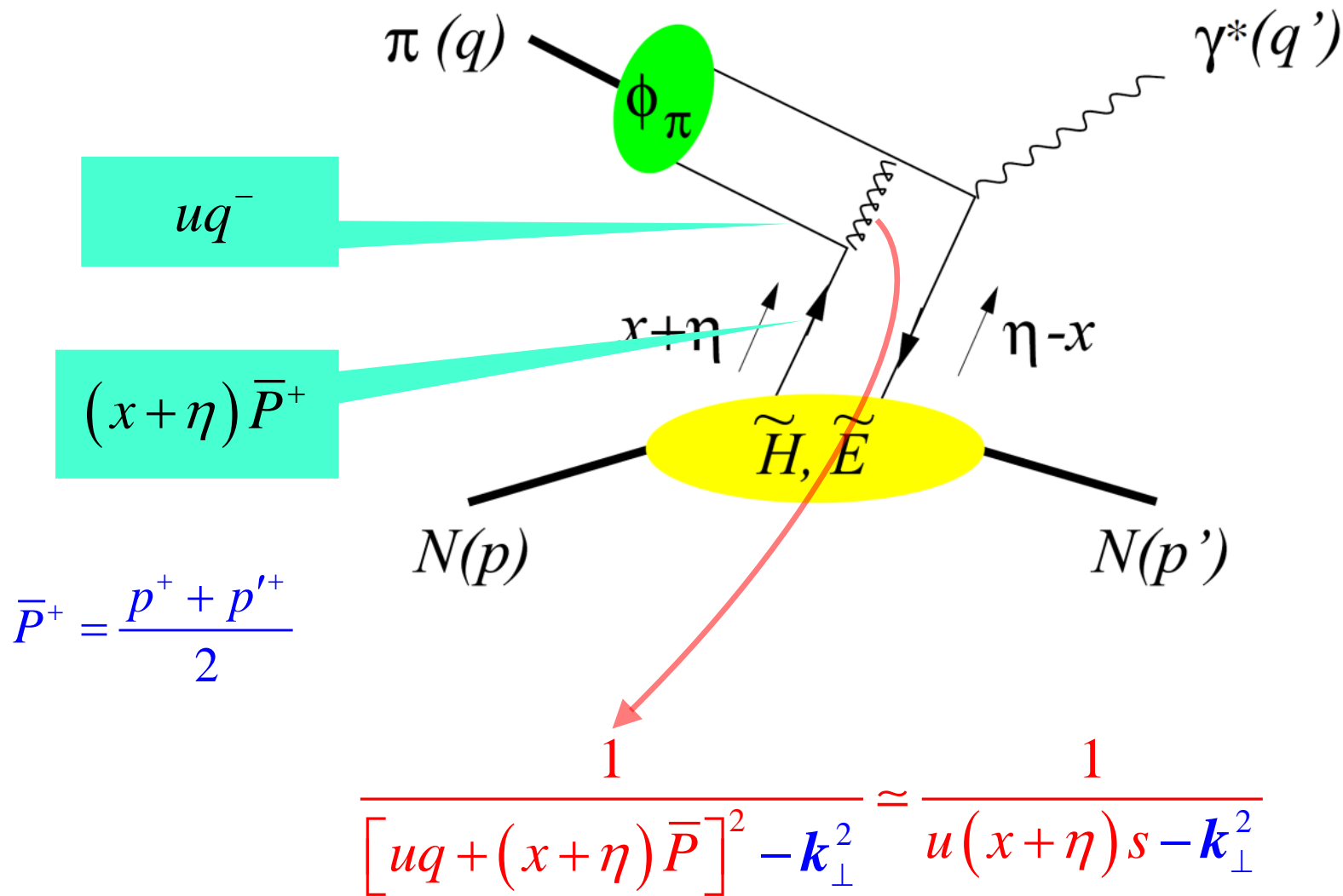
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# Collinear factorization does not work at twist-3:

- quark  $k_{\perp}$  (“ $k_T$ -factorization”) Goloskokov, Kroll  
with Sudakov resummation  
Li, Sterman



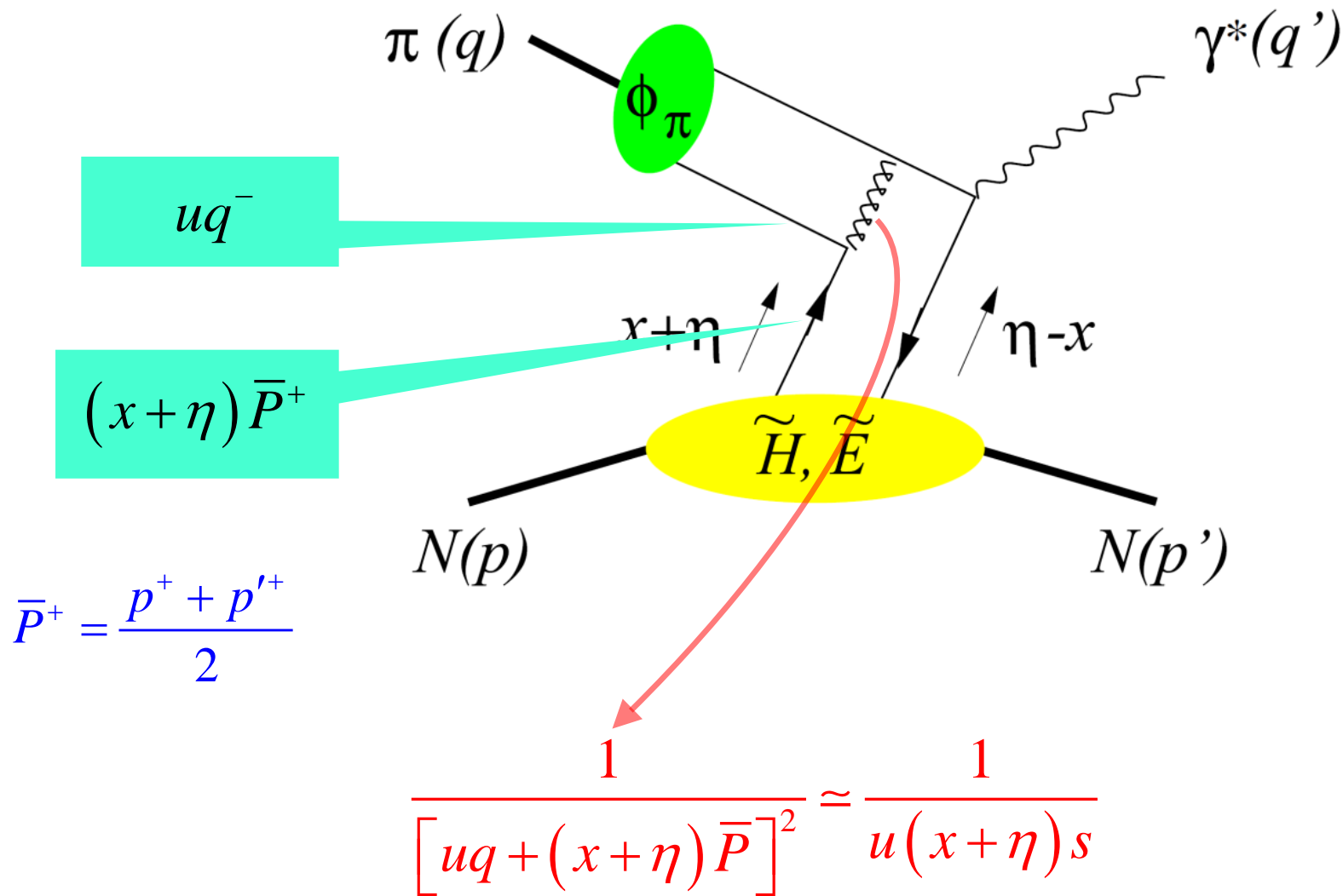


# Collinear factorization does not work at twist-3:

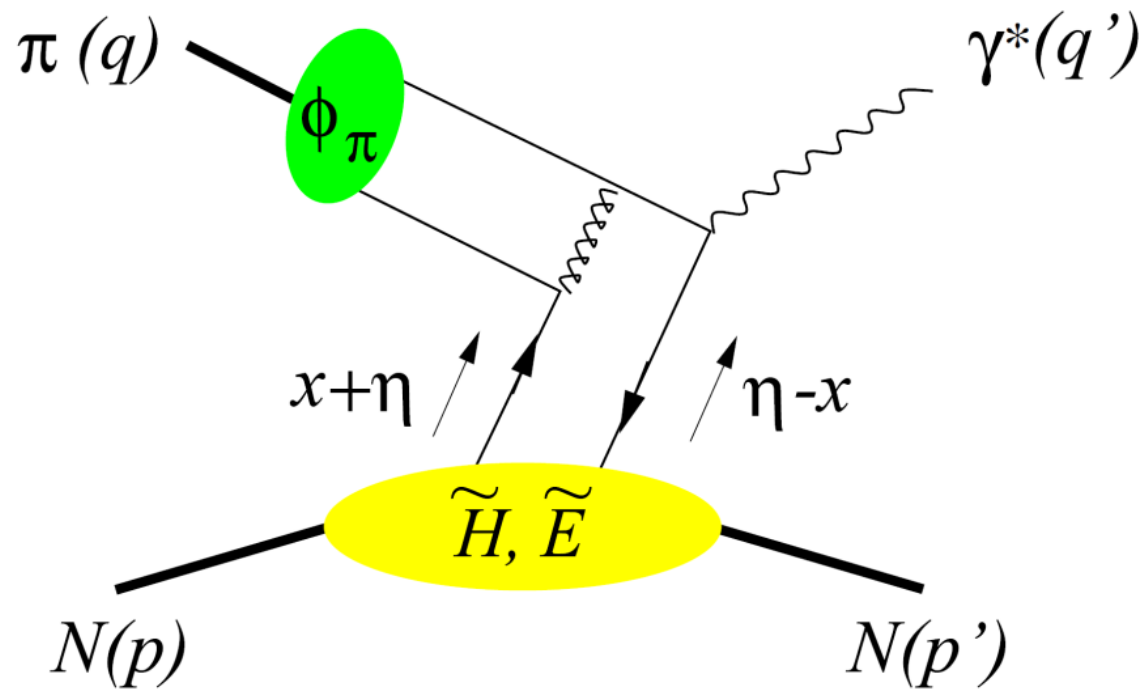
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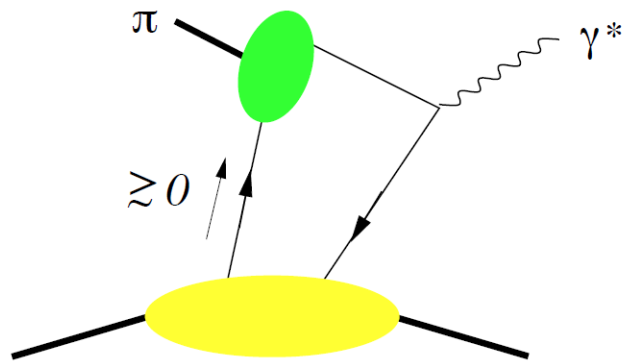
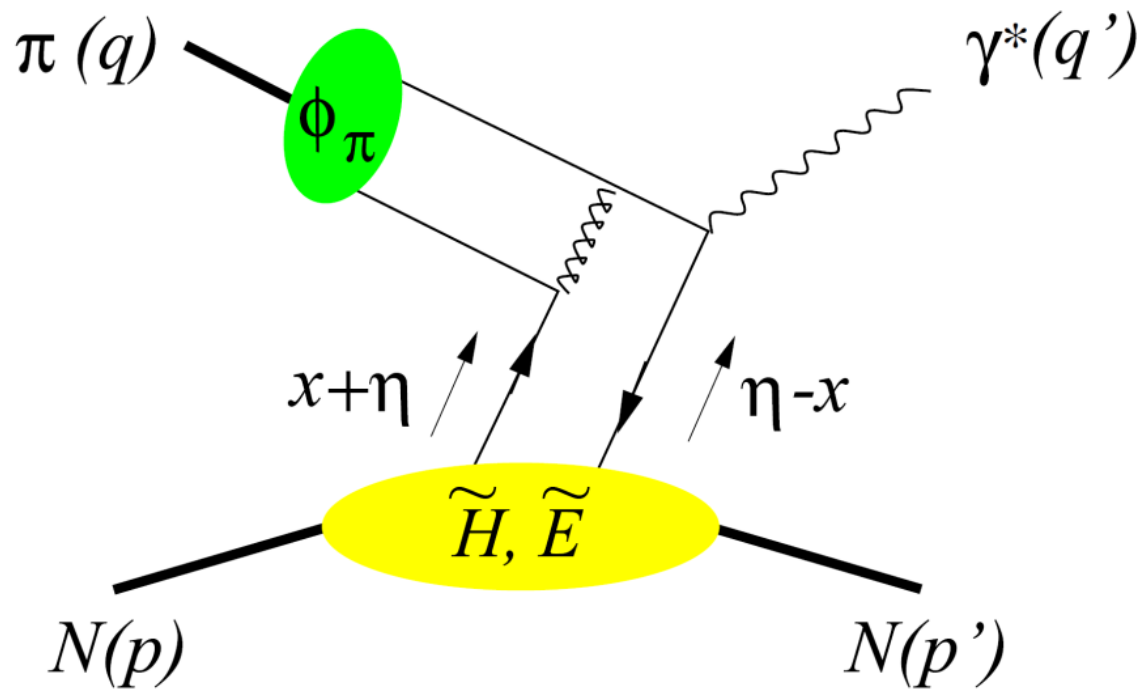
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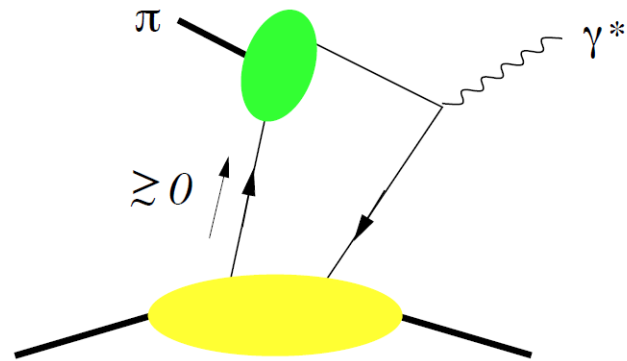
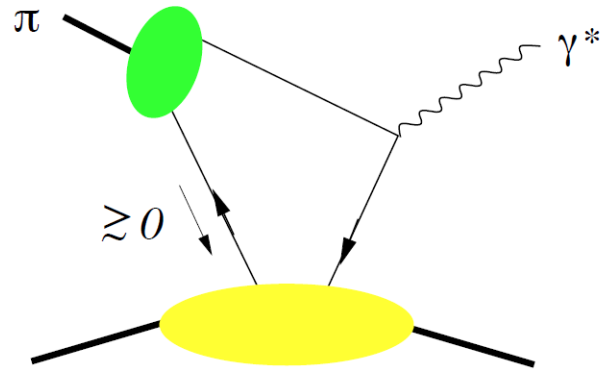




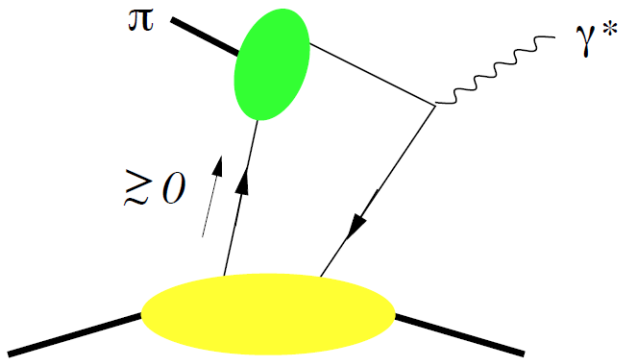
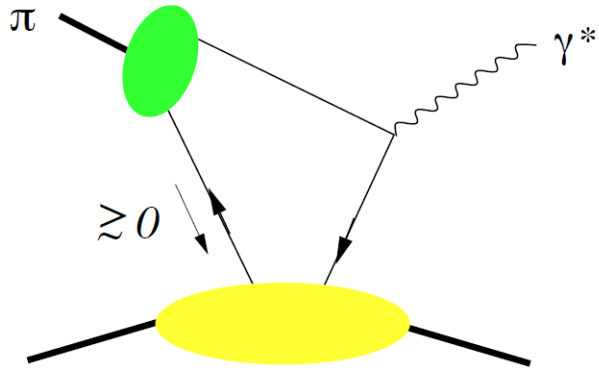




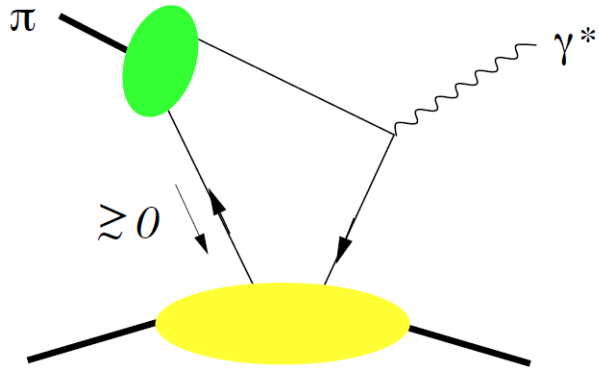
# "nonfactorizable" mechanism



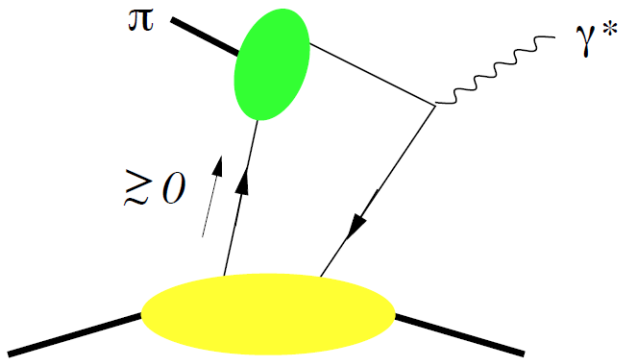
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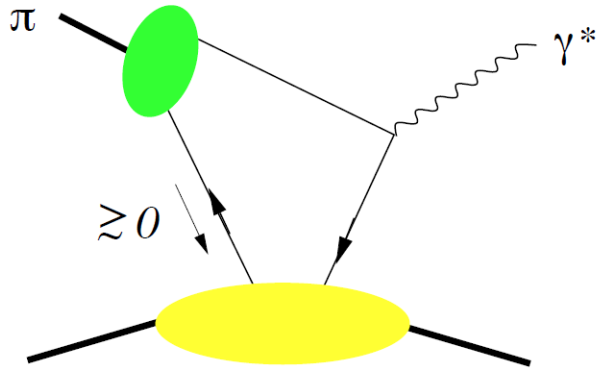
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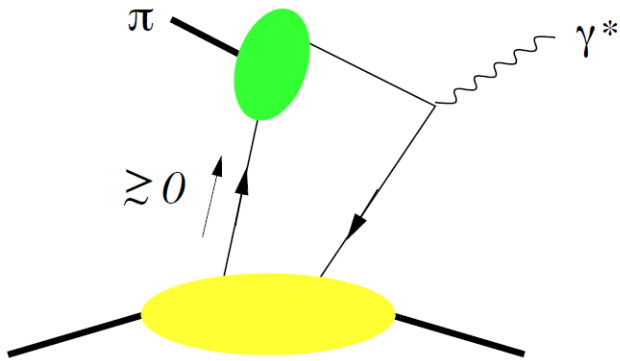
lower order in  $\alpha_s$



# "nonfactorizable" mechanism

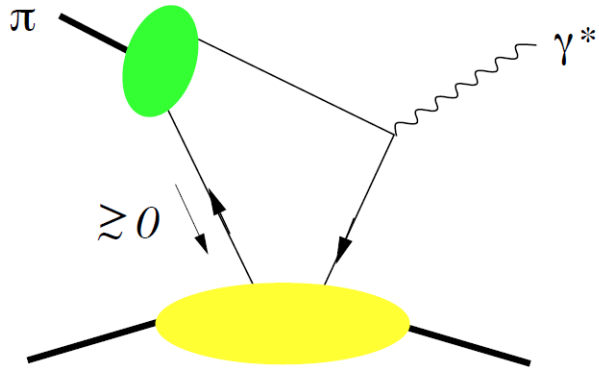


lower order in  $\alpha_s$

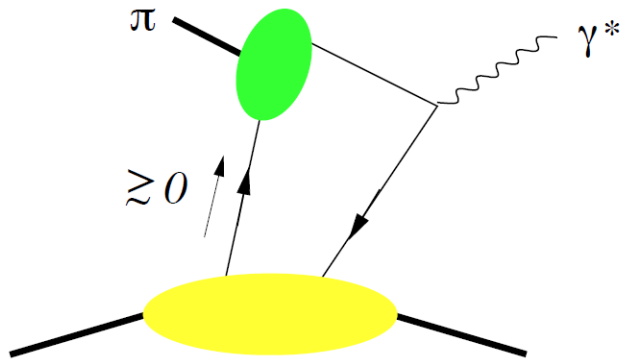


"Feynman mechanism"

# "nonfactorizable" mechanism

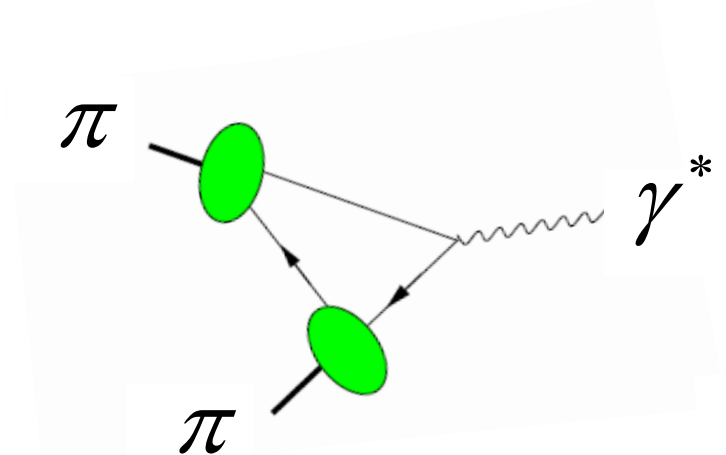


lower order in  $\alpha_s$



(d)

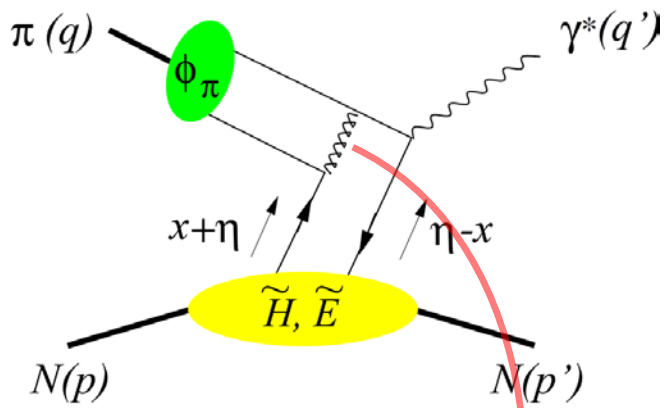
# "Feynman mechanism"



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**Li, Sterman**
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nonfactorizable “Feynman mechanism”  
at lower order in  $\alpha_s$





**Bjorken variable:**  $\tau = \frac{Q'^2}{2p \cdot q}$

**Skewness:**  $\eta = \frac{p^+ - p'^+}{p^+ + p'^+}$

**long. photon**

$$|d_{-10}^1(\theta)|^2 + |d_{10}^1(\theta)|^2$$

$$\frac{d\sigma}{dQ'^2 dt d(\cos\theta) d\varphi} = \frac{\alpha_{em}}{256\pi^3} \frac{\tau^2}{Q'^6} \sum_{\lambda', \lambda} |M^{0\lambda', \lambda}|^2 \sin^2\theta$$

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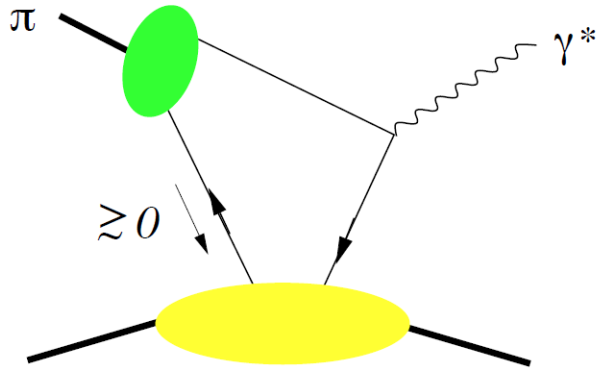
$$\int \frac{dz^-}{2\pi} e^{ix\bar{P}^+ z^-} \langle p' | \bar{\psi}(-\frac{z^-}{2}) \gamma^+ \gamma_5 \psi(\frac{z^-}{2}) | p \rangle = \frac{1}{\bar{P}^+} \left[ \tilde{H}^q(x, \eta, t) \bar{u}(p') \gamma^+ \gamma_5 u(p) + \tilde{E}^q(x, \eta, t) \bar{u}(p') \frac{\gamma_5 (p'-p)^+}{2M} u(p) \right]$$

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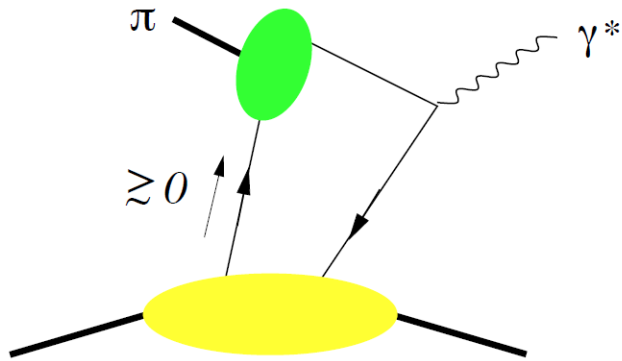
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# "nonfactorizable" mechanism

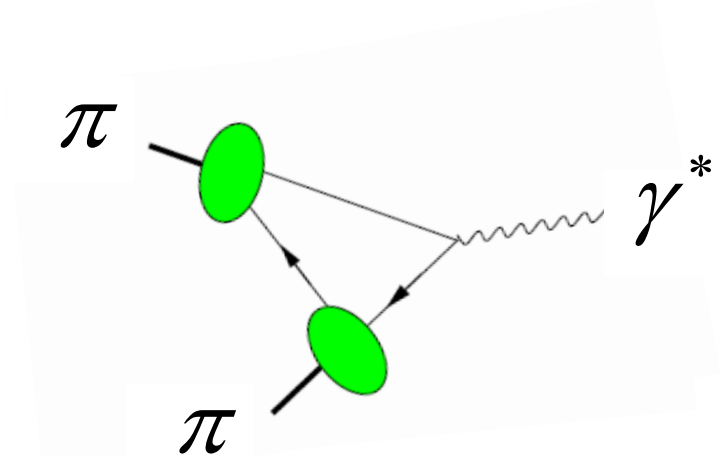


lower order in  $\alpha_s$



(d)

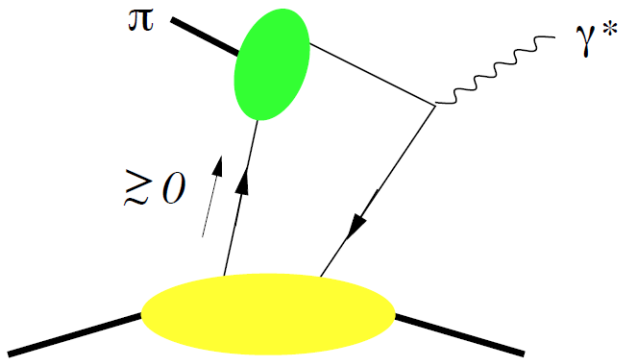
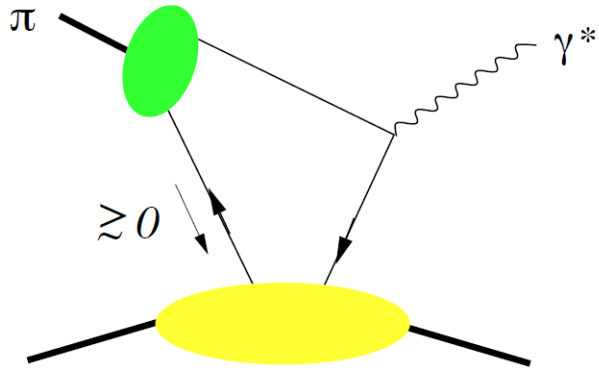
# "Feynman mechanism"



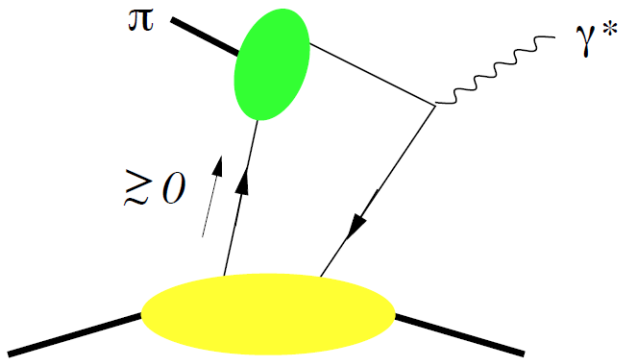
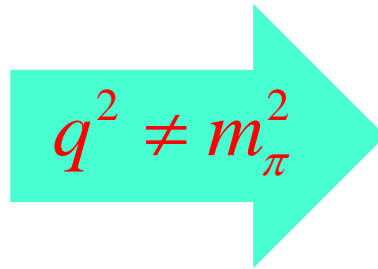
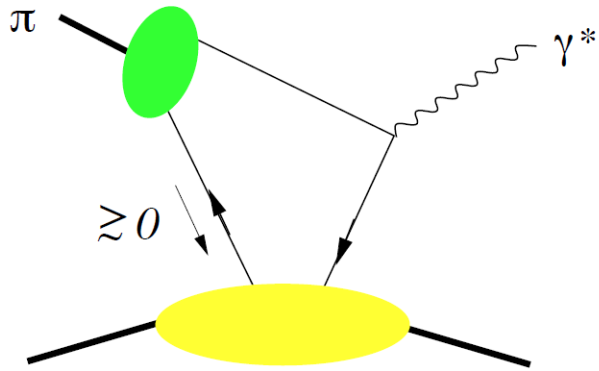
# Collinear factorization does not work at twist-3:

- quark  $k_{\perp}$  (“ $k_T$ -factorization”) **Goloskokov, Kroll**  
with Sudakov resummation  
**Li, Sterman**
- include “soft” propagator in long-distance part  
nonfactorizable “Feynman mechanism”  
at lower order in  $\alpha_s$   
**relevant also for leading twist!**

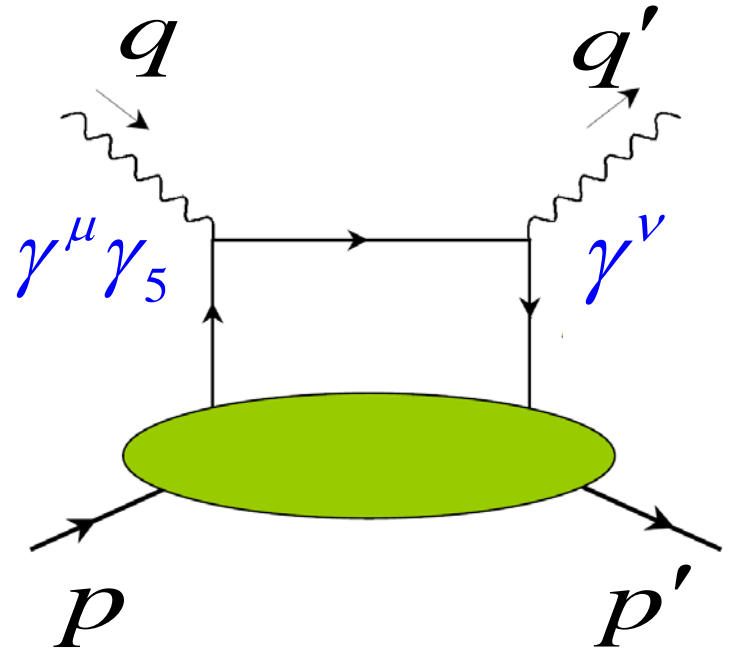
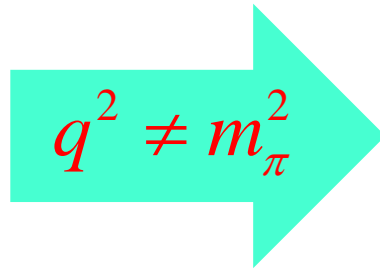
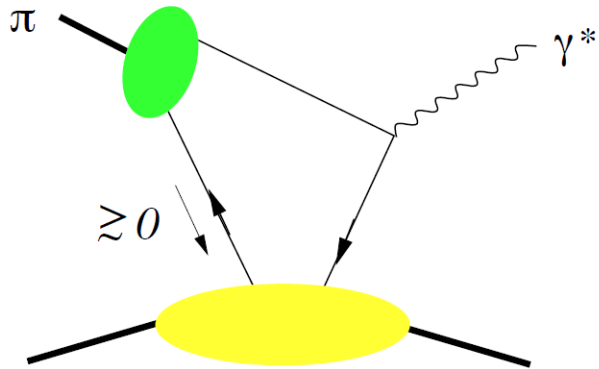
# "nonfactorizable" mechanism



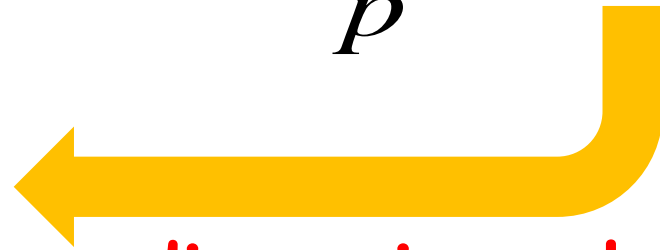
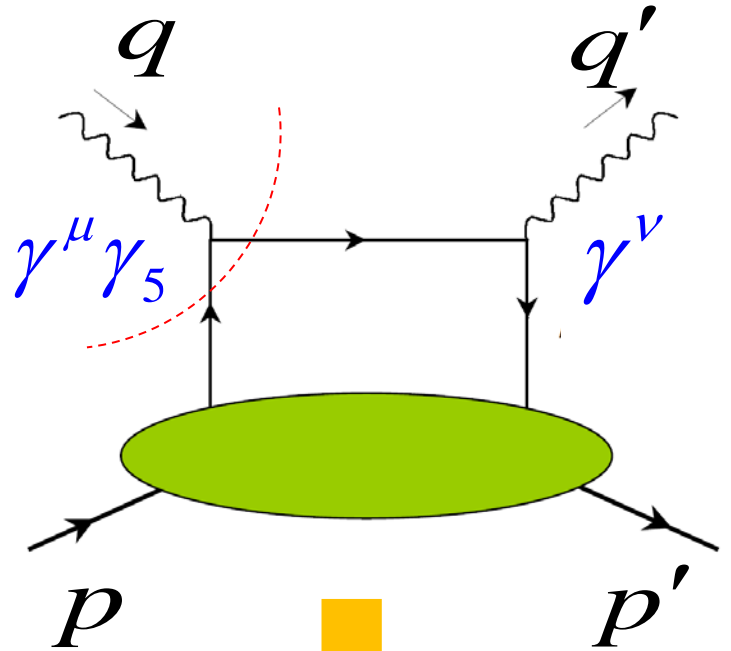
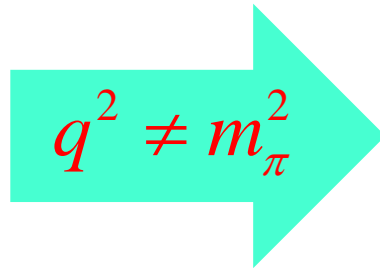
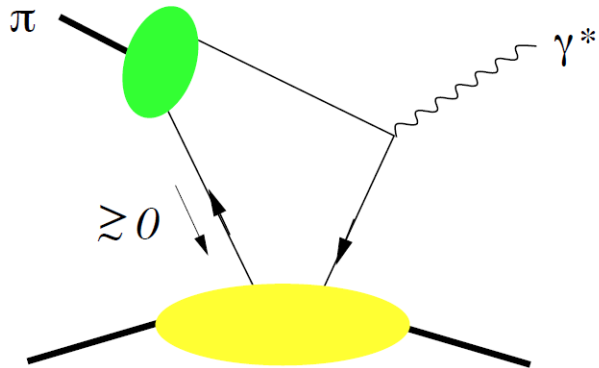
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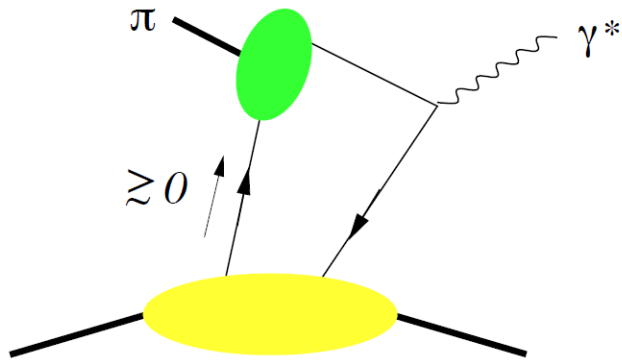
# "nonfactorizable" mechanism



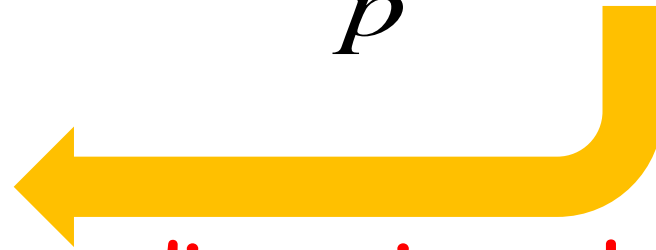
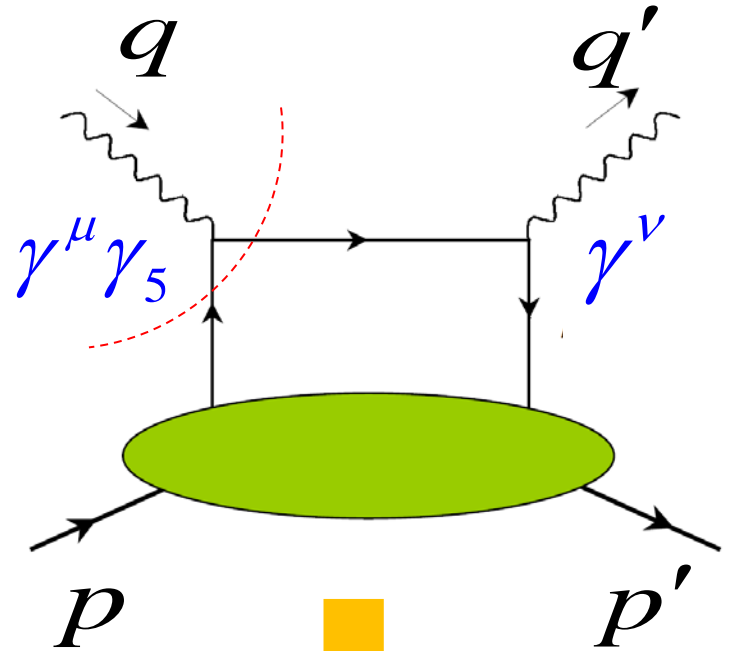
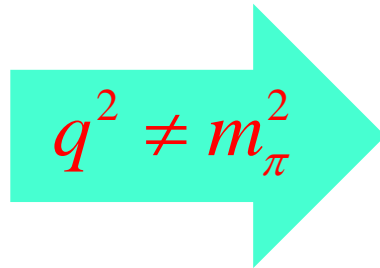
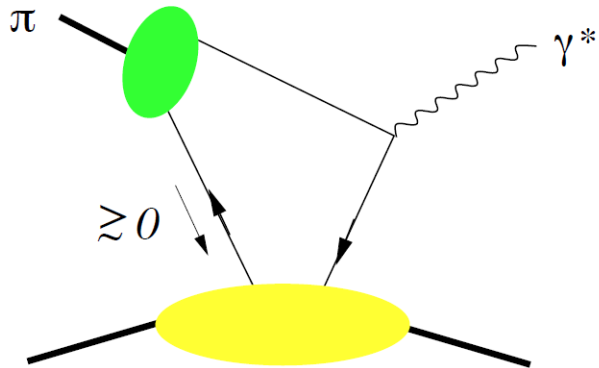
# "nonfactorizable" mechanism



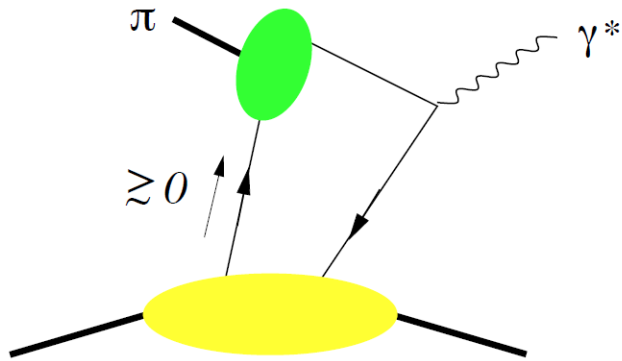
dispersion relation



# "nonfactorizable" mechanism



dispersion relation  
quark-hadron duality



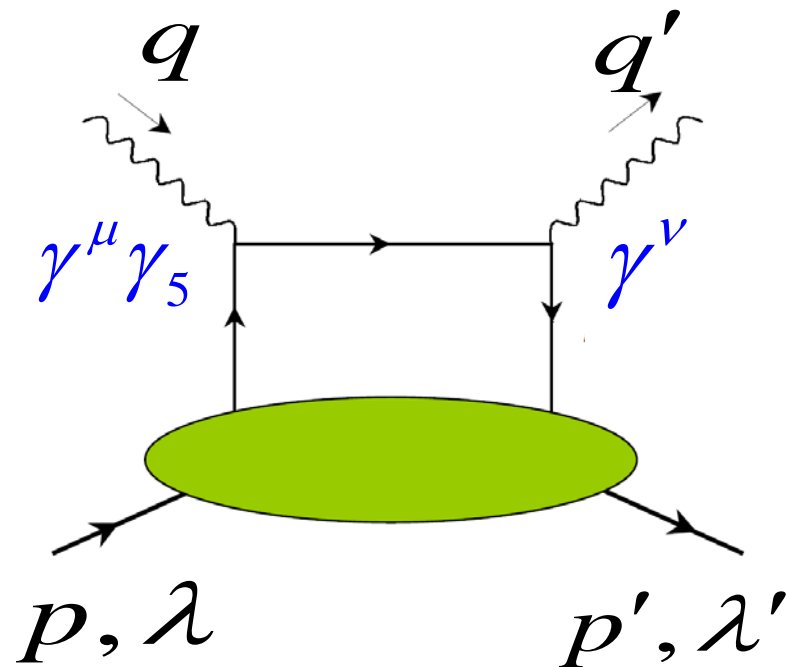


$$\int d^4x e^{iq' \cdot x} \langle p' \lambda' | \mathbf{T} j_\mu^5(0) j_\nu^{\text{em}}(x) | p \lambda \rangle$$

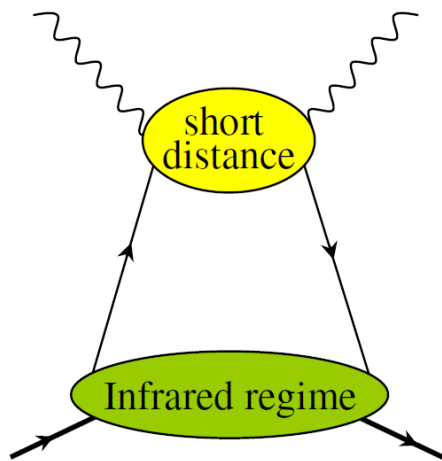
$$\equiv -iT_{\mu\nu}$$

$$j_\mu^5 = \bar{d} \gamma_\mu \gamma_5 u$$

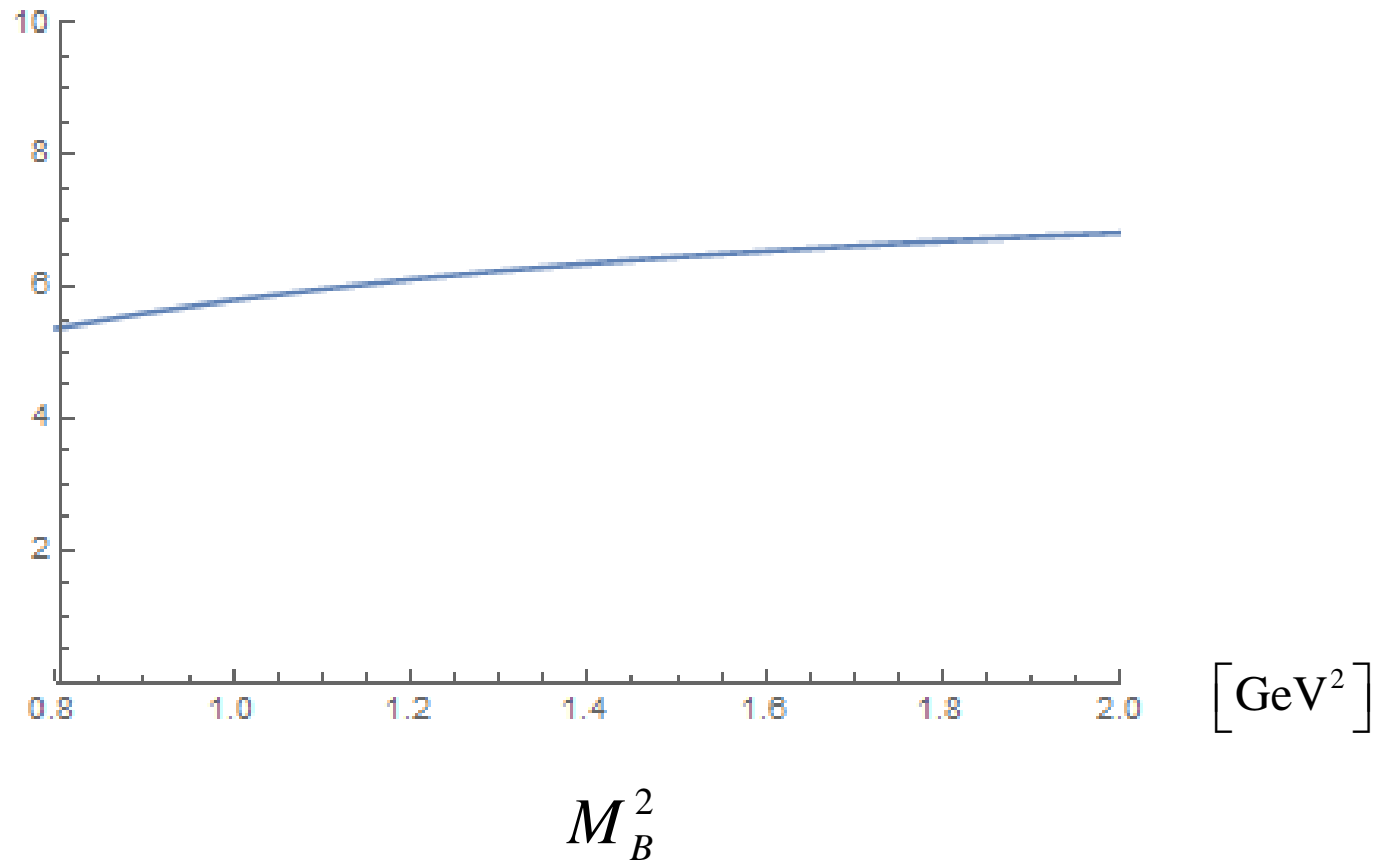
$$j_\nu^{\text{em}} = e_u \bar{u} \gamma_\nu u + e_d \bar{d} \gamma_\nu d$$



$$|q^2|, |q'^2| \gg \Lambda_{\text{QCD}}^2$$



# "Light-cone QCD SR (LCSR)"



**Borel transf.:** 
$$\hat{L}_{M_B} \left( \frac{1}{m^2 - q^2} \right) = \frac{1}{M_B^2} e^{-\frac{m^2}{M_B^2}}$$

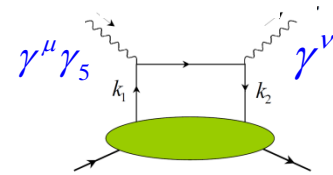
# Summary

$\pi^- p \rightarrow \gamma^* n \rightarrow \mu^+ \mu^- n$  at J-PARC GPDs

LO ( $O(\alpha_s^2)$ ) factorization formula is known, but it misses soft nonfactorizable mechanism (SNM)

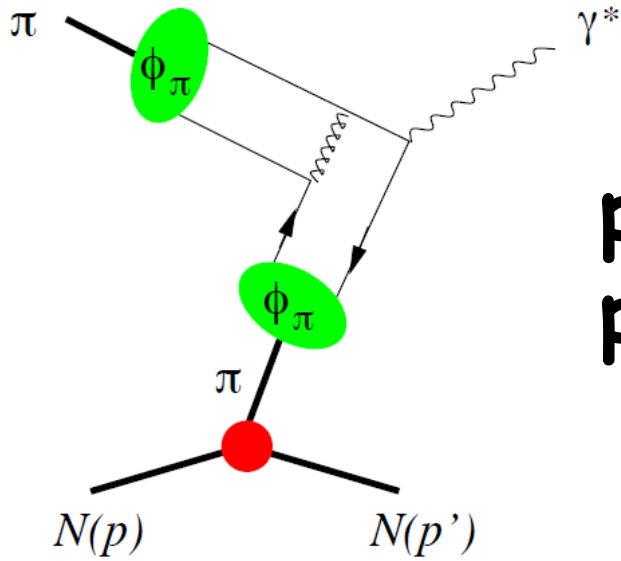
LCSR at LO ( $O(\alpha_s^0)$ ) is derived for largely model-independent estimate for SNM

$$\tilde{H}, \tilde{E}, q_{\text{th}}^2 (\sim 0.7 \text{ GeV}^2)$$



- numerical calculation at LO!
- NLO LCSR  $\longleftrightarrow$  quark  $k_\perp$ , pion pole contri.
- twist-3 LCSR  $\longrightarrow$   $M_{LCSR}^{\pm 1 \lambda', \lambda} (\pi^- p \rightarrow \gamma^* n)$

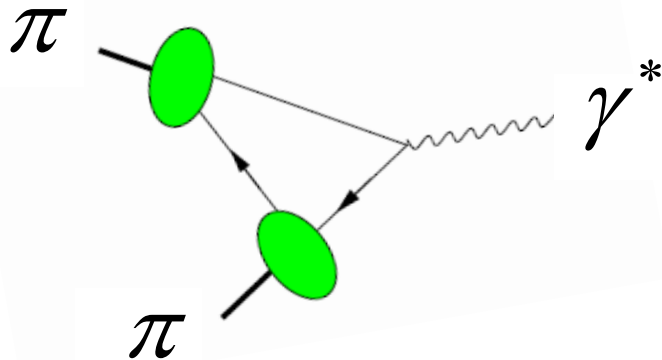
*interplay of soft/hard QCD mechanism*



pion-pole contribution using  
pion form factor  $F_\pi(Q'^2)$

Goloskokov, Kroll

$F_\pi(Q'^2)$  : important soft nonfactorizable  
contr. was shown with LCSR



Braun, Khodjamirian, Maul