Exotic baryons as a hadronic molecule in the heavy quark region

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Yasuhiro Yamaguchi and Elena Santopinto , arXiv:1606.08330 [hep-ph].

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Outline

Hidden-charm pentaquarks as a molecules

Introduction

- Exotic hadron in the heavy quark region
- Observed Pentaquarks
- Heavy Quark Spin Symmetry and Coupled channels
- 2 Meson-Baryon molecules: $\bar{D}^{(*)}\Lambda_c^{(*)} - \bar{D}^{(*)}\Sigma_c^{(*)}$



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③ Summary

Hadrons in the heavy quark region

- Hadron: Composite particle of Quarks and Gluons
- Constituent quark model (Baryon(qqq) and Meson $q\bar{q}$) has been successfully applied to the hadron spectra!



Exotic hadrons in the heavy quark region Introduction

Observation of the Exotic Hadron in the heavy quark (c, b) sectors!

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Exotic hadrons in the heavy quark region Introduction

Observation of the Exotic Hadron in the heavy quark (c, b) sectors!

e.g. Spectra of Charmonia



Charmonium cc

N. Brambilla, et al. Eur. Phys. J.C 71(2011)1534

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S. Godfrey and N. Isgur, PRD32(1985)189

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Exotic hadrons in the heavy quark region Introduction

Observation of the Exotic Hadron in the heavy quark (c, b) sectors!

e.g. Spectra of Charmonia



▷ Why are many exotic hadrons found in the heavy quark region?

Exotic structure: Hadronic molecules



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Exotic structure: Hadronic molecules



• Loosely bound states (resonances) of hadrons

 \rightarrow Appearing near the thresholds (M-M, M-B,...)

 \Rightarrow Analogous to Atomic Nuclei

• Molecules are formed by the Hadron-Hadron interaction dynamically.

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Exotic structure: Hadronic molecules Introduction



- Theoretical researches
 - X(3872) as DD̄*.

M. T. AlFiky, et al., PLB640(2006)238, M. B. Voloshin, Prog. Part. Nucl. Phys. 61(2008)455

- Z_h as BB^{*}, J. R. Zhang, et al., PLB704(2011)312,S.Ohkoda, et al., PRD86(2012)014004
- $\Lambda(1405)$ as $\bar{K}N$, T. Hyodo and D. Jido, Prog. Part. Nucl. Phys. 67 (2012) 55
- Λ^{*}_c as DN, T.Mizutani, A.Ramos, PRC74(2006)065201, C. Garcia-Recio, et al., PRD79(2009)054004

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Observation of two hidden-charm pentaguarks !! Introduction

week enung PHYSICAL REVIEW LETTERS PRL 115, 072001 (2015) 14 AUGUST 2015 Ś Observation of $J/\psi p$ Resonances Consistent with Pentaguark States in $\Lambda_b^0 \to J/\psi K^- p$ Decays R. Aaij et al. (LHCb Collaboration) (Received 13 July 2015; published 12 August 2015) Observations of exotic structures in the $J/\psi p$ channel, which we refer to as charmonium-pentaquark states, in $\Lambda_h^0 \to J/\psi K^- p$ decays are presented. The data sample corresponds to an integrated luminosity of 3 fb⁻¹ acquired with the LHCb detector from 7 and 8 TeV pp collisions. An amplitude analysis of the three-body final state reproduces the two-body mass and angular distributions. To obtain a satisfactory fit of the structures seen in the $J/\psi p$ mass spectrum, it is necessary to include two Breit-Wigner amplitudes that each describe a resonance the significance of each of these resonances is more than 9 standard deviations. One has a matrix 4380 \pm 8 \pm 29 MeV and a width of 205 \pm 18 \pm 86 MeV, while the second is narrower, with a map $4449.8 \pm 1.7 \pm 2.5$ MeV and a width of $39 \pm 5 \pm 19$ MeV. The preferred J^P assignments are of op e parity, with one state having spin 3/2 and the other 5/2. DOI: 10.1103/PhysRevLett.115.072001 PACS numbers: 14.40.Pg, 13.25.Gv



27 July 2016

Observation of two hidden-charm pentaquarks !! Introduction

PRL 115, 072001 (2015)

PHYSICAL REVIEW LETTERS

14 AUGUST 2015

\mathfrak{F} Observation of $J/\psi p$ Resonances Consistent with Pentaquark States

in $\Lambda_b^0 \to J/\psi K^- p$ Decays

R. Aaij et al.^{*} (LHCb Collaboration) (Received 13 July 2015; published 12 August 2015)

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1. $M = 4380 \pm 8 \pm 29$ MeV, $\Gamma = 205 \pm 18 \pm 86$ MeV (Broad)

2. $M = 4449.8 \pm 1.7 \pm 2.5$ MeV, $\Gamma = 39 \pm 5 \pm 19$ MeV (Narrow)

• J^P Assignment: 3/2⁻, 5/2⁺; 3/2⁺, 5/2⁻; 5/2⁺, 3/2⁻

What is the structure of the pentaquarks? Introduction

Theoretical discussions

- Compact pentaquark?
 - ▶ Quark model

W.L.Wang et al., PRC84(2011)015203

G. Yang and J. Ping, (2015), arXiv:1511.09053Pentaquark ... (Compact)

- Hadronic molecule?
 - SU(4) flavor symmetry
 J.-J.Wu *et al.*, PRL105(2010)232001
 C.W.Xiao *et al.*, PRD88(2013)056012
 T. Uchino,*et al.*, Eur.Phys.J.A52(2016)43



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Important issue of the heavy pentaguarks Introduction

1. Pentaguarks are close to the meson-baryon thresholds \Rightarrow Hadronic molecules appears near the thresholds!



2. Heavy Quark Spin Symmetry

 \Rightarrow SU(4) symmetry is broken in the charm guark sector.

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③ Summary

Heavy Quark Spin Symmetry and Mass degeneracy Introduction

Heavy Quark Spin Symmetry (HQS) N.Isgur, M.B.Wise, PLB232(1989)113

- Suppression of Spin-spin force in $m_Q \rightarrow \infty$.
- ⇒ Decomposition of Heavy quark spin and Light components $\vec{J} = \vec{L} + \vec{S} = \vec{S}_0 + \vec{i}$



 \Rightarrow Mass degeneracy of hadrons with the different J

• Mass degeneracy of $\{D, D^*\}(Q\bar{q})$, $\{\eta_c, J/\psi\}(Q\bar{Q})$, $\{\Sigma_c, \Sigma_c^*\}(Qqq)$ (baryons)...

Mass degeneracy of heavy hadrons

• Mass difference between vector and pseudoscalar mesons. $(Q\bar{q}, q = u, d)$



- $\triangleright \Delta m$ decreases when the quark mass increases.
- Mass degeneracy of heavy hadrons appears!

Mass degeneracy of heavy hadrons

• Mass difference between $1/2^+$ and $3/2^+$ baryons. (*Qqq*, q = u, d) $J^{\prime\prime} = 3/2^+$



- $\triangleright \Delta m$ decreases when the quark mass increases.
- Mass degeneracy of heavy hadrons appears!

 \Rightarrow Small mass splitting leads to **Channel couplings**!

Small mass splitting and Interactions (KN and DN)



Small mass splitting and Interactions (KN and DN)

- In the heavy (c, b) sector, the Heavy Quark Spin Symmetry induces the $\overline{D} \overline{D}^*$ mixing.

 $m_{K^*}-m_K\sim 400~{
m MeV} \Leftrightarrow m_{D^*}-m_D\sim 140~{
m MeV}$

• Appearance of the the one π exchange potential .

T.D.Cohen, P.M.Hohler, R.F.Lebed PRD72 (2005)074010, S.Yasui, K.Sudoh, PRD80 (2009)034008,

Y.Yamaguchi,S.Ohkoda,S.Yasui,A.Hosaka PRD84(2011)014032,...

π exchange potential (OPEP) and Coupled channel

> OPEP is important to bind atomic nuclei.

Tensor force of the OPEP generates a strong attraction.



π exchange potential (OPEP) and Coupled channel

> OPEP is important to bind atomic nuclei.

Tensor force of the OPEP generates a strong attraction.



▷ HQS: mixing of $D - D^*$

▷ Tensor force: mixing of S - D(P - F)



Unit: MeV

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Y.Y., S.Ohkoda, S.Yasui and A.Hosaka, PRD84 014032 (2011) and PRD85 054003 (2012)

Energy spectra of D meson-Nucleon (DN) states

Introduction

One bound state



Y.Y., S.Ohkoda, S.Yasui and A.Hosaka, PRD84 014032 (2011) and PRD85 054003 (2012)

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• One bound state, and resonances in charm



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• One bound state, and resonances in charm and bottom sectors!



● Many states near the thresholds. ⇔ No KN bound state

• One bound state, and resonances in charm and bottom sectors!



Y.Y., S.Ohkoda, S.Yasui and A.Hosaka, PRD84 014032 (2011) and PRD85 054003 (2012)

- Many states near the thresholds. ⇔ No KN bound state
- The tensor force from the D

 D
 * mixing is important to produce a strong attraction!

Channel coupling of Hidden-charm meson-baryon state



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Coupled channels of the hidden-charm pentaguark Introduction



Mass degeneracy of $\bar{D}\Sigma_c$, $\bar{D}\Sigma_c^*$, $\bar{D}^*\Sigma_c$ and $\bar{D}^*\Sigma_c^*$ 0

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$$ar{D} - ar{D}^*$$
 and $\Sigma_{
m c} - \Sigma_{
m c}^*$ mixings

• Mass degeneracy of $\bar{D}\Sigma_c$, $\bar{D}\Sigma_c^*$, $\bar{D}^*\Sigma_c$ and $\bar{D}^*\Sigma_c^*$

•
$$\Lambda_{c}$$
 (cqq): $ar{D}^{(*)}\Lambda_{c}$ channel?

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$$\bar{D}^{(*)}\Lambda_{c} - \bar{D}^{(*)}\Sigma_{c}^{(*)}$$
 mixing (analogous to $\Lambda N - \Sigma N$)
 $m_{\Sigma_{c}} - m_{\Lambda_{c}} \sim 170 \text{ MeV}$
 $\boxed{\bar{D}^{(*)}\Sigma_{c}^{(*)} \text{ coupling}}$ $\boxed{\bar{D}^{(*)}\Lambda_{c}^{(*)} - \bar{D}^{(*)}\Sigma_{c}^{(*)} \text{ coupling}}$
 $\bar{D}^{(*)}$
 π, ρ, ω, \dots $\Sigma_{c}^{(*)}$ $\bar{D}^{(*)}$
 π, ρ Λ_{c}

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• $\bar{D}^{(*)}\Lambda_c - \bar{D}^{(*)}\Sigma_c^{(*)}$ mixing (analogous to $\Lambda N - \Sigma N$) $m_{\Sigma_c}-m_{\Lambda_{
m c}}\sim 170~{
m MeV}$ $|\bar{\mathsf{D}}^{(*)} \mathbf{\Sigma}_{\mathrm{c}}^{(*)}|$ coupling $\left| ar{\mathsf{D}}^{(*)} \mathsf{\Lambda}^{(*)}_{\mathrm{c}} - ar{\mathsf{D}}^{(*)} \mathbf{\Sigma}^{(*)}_{\mathrm{c}} ext{ coupling}
ight|$ $\Rightarrow \bar{D}\Lambda_{c}, \bar{D}^{*}\Lambda_{c}, \bar{D}\Sigma_{c}, \bar{D}\Sigma_{c}^{*}, \bar{D}^{*}\Sigma_{c}, \bar{D}^{*}\Sigma_{c}^{*}$ (6 thresholds!)

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• $\bar{D}^{(*)}\Lambda_c - \bar{D}^{(*)}\Sigma_c^{(*)}$ mixing (analogous to $\Lambda N - \Sigma N$) $m_{\Sigma_c}-m_{\Lambda_c}\sim 170~{
m MeV}$ $\left| \bar{\mathsf{D}}^{(*)} \boldsymbol{\Sigma}_{c}^{(*)} \right|$ coupling $\bar{\mathsf{D}}^{(*)} \mathsf{\Lambda}_{c}^{(*)} - \bar{\mathsf{D}}^{(*)} \boldsymbol{\Sigma}_{c}^{(*)}$ coupling

 $\Rightarrow \bar{D}\Lambda_{\rm c}, \bar{D}^*\Lambda_{\rm c}, \bar{D}\Sigma_{\rm c}, \bar{D}\Sigma_{\rm c}^*, \bar{D}^*\Sigma_{\rm c}, \bar{D}^*\Sigma_{\rm c}^* \text{ (6 thresholds!)}$

Coupling to a state with ℓ ≠ 0 (D-wave,...)
 ⇒ Tensor force producing a strong attraction!

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• Allowed channels
$$({}^{2S+1}L)$$
 for $J^P=3/2^\pm,5/2^\pm$

J^P	Channels
3/2-	
3/2+	
$5/2^{-}$	
$5/2^{+}$	

•
$$J^P = 3/2^{\pm}$$
:
• $J^P = 5/2^{\pm}$:

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• Allowed channels
$$({}^{2S+1}L)$$
 for $J^P=3/2^\pm,5/2^\pm$

J^P	Channels
3/2-	$ \begin{array}{l} \bar{D}\Lambda_{\rm c}(^2D), \ \bar{D}^*\Lambda_{\rm c}(^4S, ^2D, ^4D), \ \bar{D}\Sigma_{\rm c}(^2D), \ \bar{D}\Sigma_{\rm c}^*(^4S, ^4D), \\ \bar{D}^*\Sigma_{\rm c}(^4S, ^2D, ^4D), \ \bar{D}^*\Sigma_{\rm c}^*(^4S, ^2D, ^4D, ^6D, ^6G) \end{array} $
3/2+	
$5/2^{-}$	
5/2+	

•
$$J^P = 3/2^{\pm}$$
: **15 channels!**
• $J^P = 5/2^{\pm}$:

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• Allowed channels
$$(^{2S+1}L)$$
 for $J^P=3/2^\pm,5/2^\pm$

J^P	Channels
3/2-	$\overline{D}\Lambda_{\rm c}(^2D), \ \overline{D}^*\Lambda_{\rm c}(^4S, ^2D, ^4D), \ \overline{D}\Sigma_{\rm c}(^2D), \ \overline{D}\Sigma_{\rm c}(^4S, ^4D),$
	$D^*\Sigma_{\rm c}({}^4S, {}^2D, {}^4D), \ D^*\Sigma_{\rm c}^*({}^4S, {}^2D, {}^4D, {}^6D, {}^6G)$
$3/2^{+}$	$\bar{D}\Lambda_{c}(^{2}P), \ \bar{D}^{*}\Lambda_{c}(^{2}P, ^{4}P, ^{4}F), \ \bar{D}\Sigma_{c}(^{2}P), \ \bar{D}\Sigma_{c}^{*}(^{4}P, ^{4}F),$
·	$\bar{D}^*\Sigma_{\rm c}({}^2P, {}^4P, {}^4F), \ \bar{D}^*\Sigma_{\rm c}^*({}^2P, {}^4P, {}^6P, {}^4F, {}^6F)$
$5/2^{-}$	
- /	
5/2+	
5/2	

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• Allowed channels
$$({}^{2S+1}L)$$
 for $J^P=3/2^\pm,5/2^\pm$

J^P	Channels
3/2-	$\bar{D}\Lambda_{\rm c}(^2D), \ \bar{D}^*\Lambda_{\rm c}(^4S,^2D,^4D), \ \bar{D}\Sigma_{\rm c}(^2D), \ \bar{D}\Sigma_{\rm c}^*(^4S,^4D),$
	$D^{*}\Sigma_{c}(3, 2D, D), D^{*}\Sigma_{c}(3, 2D, D, 0), G$
3/2+	$ \begin{array}{l} \bar{D}\Lambda_{\rm c}(^2P), \ \bar{D}^*\Lambda_{\rm c}(^2P, ^4P, ^4F), \ \bar{D}\Sigma_{\rm c}(^2P), \ \bar{D}\Sigma_{\rm c}^*(^4P, ^4F), \\ \bar{D}^*\Sigma_{\rm c}(^2P, ^4P, ^4F), \ \bar{D}^*\Sigma_{\rm c}^*(^2P, ^4P, ^6P, ^4F, ^6F) \end{array} $
$5/2^{-}$	$ \begin{array}{l} \bar{D}\Lambda_{\rm c}(^2D), \ \bar{D}^*\Lambda_{\rm c}(^2D, ^4D, ^4G), \ \bar{D}\Sigma_{\rm c}(^2D), \ \bar{D}\Sigma_{\rm c}^*(^4D, ^4G), \\ \bar{D}^*\Sigma_{\rm c}(^2D, ^4D, ^4G), \ \bar{D}^*\Sigma_{\rm c}^*(^6S, ^2D, ^4D, ^6D, ^4G, ^6G) \end{array} $
$5/2^{+}$	

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• Allowed channels
$$({}^{2S+1}L)$$
 for $J^P=3/2^\pm,5/2^\pm$

J^P	Channels
3/2-	$\bar{D}\Lambda_{\rm c}(^2D), \ \bar{D}^*\Lambda_{\rm c}(^4S, ^2D, ^4D), \ \bar{D}\Sigma_{\rm c}(^2D), \ \bar{D}\Sigma_{\rm c}^*(^4S, ^4D),$
	$D^*\Sigma_{\rm c}({}^4S, {}^2D, {}^4D), \ D^*\Sigma_{\rm c}^*({}^4S, {}^2D, {}^4D, {}^6D, {}^6G)$
$3/2^{+}$	$\bar{D}\Lambda_{c}(^{2}P), \ \bar{D}^{*}\Lambda_{c}(^{2}P, ^{4}P, ^{4}F), \ \bar{D}\Sigma_{c}(^{2}P), \ \bar{D}\Sigma_{c}^{*}(^{4}P, ^{4}F),$
	$D^*\Sigma_{\rm c}({}^2P, {}^4P, {}^4F), \ D^*\Sigma_{\rm c}^*({}^2P, {}^4P, {}^6P, {}^4F, {}^6F)$
$5/2^{-}$	$\bar{D}\Lambda_{\rm c}(^2D), \ \bar{D}^*\Lambda_{\rm c}(^2D,^4D,^4G), \ \bar{D}\Sigma_{\rm c}(^2D), \ \bar{D}\Sigma_{\rm c}^*(^4D,^4G),$
	$ar{D}^*\Sigma_{ m c}(^2D,^4D,^4G),\ ar{D}^*\Sigma_{ m c}^*(^6S,^2D,^4D,^6D,^4G,^6G)$
$5/2^{+}$	$\bar{D}\Lambda_{c}(^{2}F), \ \bar{D}^{*}\Lambda_{c}(^{4}P, ^{2}F, ^{4}F), \ \bar{D}\Sigma_{c}(^{2}F), \ \bar{D}\Sigma_{c}^{*}(^{4}P, ^{4}F),$
	$ar{D}^* \Sigma_{ m c}({}^4P, {}^2F, {}^4F), \ ar{D}^* \Sigma_{ m c}^*({}^4P, {}^6P, {}^2F, {}^4F, {}^6F, {}^6H)$

J^P = 3/2[±]: 15 channels!
 J^P = 5/2[±]: 16 channels!

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Main Subject: Pentaquarks

• Hadronic molecules formed by hidden-charm meson-baryon.



- Bound and resonant states of $ar{D}^{(*)} \Lambda_{
 m c} ar{D}^{(*)} \Sigma_{
 m c}^{(*)}$
- ▷ Coupling to $\overline{D}^{(*)}\Lambda_c$ and $\overline{D}^{(*)}\Sigma_c^{(*)}$
- \triangleright Coupling to the state with $\ell \neq 0$
- ▷ Negative and Positive parity states $(P = \pm)$

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- \triangleright Coupling to $\overline{D}^{(*)}\Lambda_c$ and $\overline{D}^{(*)}\Sigma_c^{(*)}$
- ▷ Coupling to the state with $\ell \neq 0$
- ▷ Negative and Positive parity states $(P = \pm)$

The full-coupled channel analysis of $\bar{D}^{(*)}\Lambda_{\rm c} - \bar{D}^{(*)}\Sigma_{\rm c}^{(*)}$ has never performed so far !

Results of $\overline{D}^{(*)}B$ states (2-body)



Bound state and Resonance

- We solve the coupled-channel Schrödinger equations with $J^P = 3/2^{\pm}, 5/2^{\pm}$ and isospin I = 1/2.
- Interaction: $\pi\rho\omega\sigma$ exchange potentials

$\bar{D}^{(*)}B$ Interaction: Meson exchange potential

• Effective Lagrangian with heavy quark symmetry

R.Casalbuoni *et al.*, Phys.Rept.**281** (1997)145, T.M.Yan *et al.*, PRD**46**(1992)1148, Y.-R.Liu and M.Oka, PRD**85**(2012)014015

Meson:
$$\mathcal{L}_{\pi HH} = g_{\pi} \text{Tr} \left[H_b \gamma_{\mu} \gamma_5 A^{\mu}_{ba} \bar{H}_a \right]$$

Baryon:

$$\mathcal{L}_{\pi BB} = \frac{3}{2} g_1 i v_{\kappa} \varepsilon^{\mu \nu \lambda \kappa} \text{tr} \left[\bar{S}_{\mu} A_{\nu} S_{\lambda} \right] + g_4 \text{tr} \left[\bar{S}^{\mu} A_{\mu} \Lambda_c \right]$$

Heavy meson and baryon fields

$$\begin{aligned} \mathcal{H}_{a} &= \frac{1+\not{\!\!\!/}}{2} \left[\mathbf{P}_{a\,\mu}^{*} \gamma^{\mu} - \mathbf{P}_{a} \gamma^{5} \right] \quad (1^{-} \text{ and } 0^{-}) \\ \mathcal{S}_{\mu} &= \boldsymbol{\Sigma}_{\mu}^{*} + \frac{\delta}{\sqrt{3}} \left(\gamma_{\mu} + \boldsymbol{v}_{\mu} \right) \gamma_{5} \boldsymbol{\Sigma} \quad (3/2^{+} \text{ and } 1/2^{+}) \end{aligned}$$

•
$$ar{D} - ar{D}^*$$
 and $\Sigma_{
m c} - \Sigma_{
m c}^*$ mixings

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$\bar{D}^{(*)}B$ Interaction: Meson exchange potential

• Effective Lagrangian with heavy quark symmetry

R.Casalbuoni *et al.*, Phys.Rept.**281** (1997)145, T.M.Yan *et al.*, PRD**46**(1992)1148, Y.-R.Liu and M.Oka, PRD**85**(2012)014015

Meson: $\mathcal{L}_{\pi HH} = g_{\pi} \operatorname{Tr} \left[\mathbf{H}_{\mathbf{b}} \gamma_{\mu} \gamma_{5} A_{ba}^{\mu} \bar{\mathbf{H}}_{a} \right]$ Baryon:

$$\mathcal{L}_{\pi BB} = \frac{3}{2} g_1 i v_{\kappa} \varepsilon^{\mu \nu \lambda \kappa} \text{tr} \left[\overline{\mathbf{S}}_{\mu} A_{\nu} \mathbf{S}_{\lambda} \right] + g_4 \text{tr} \left[\overline{\mathbf{S}}^{\mu} A_{\mu} \Lambda_c \right]$$

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$$S_{\mu} = \Sigma_{\mu}^{*} + \frac{\delta}{\sqrt{3}} \left(\gamma_{\mu} + v_{\mu} \right) \gamma_{5} \Sigma \quad (3/2^{+} \text{ and } 1/2^{+})$$

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D^(*)B Interaction: Meson exchange potential

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$$\begin{array}{c|c} \bar{D}^{(*)} & B \\ \mathcal{L}_{\mathbf{m}\bar{D}^{(*)}\bar{D}^{(*)}} & \pi, \rho, \omega, \sigma \\ \bar{D}^{(*)}\bar{D}^{(*)} & \overline{D}^{(*)}\mathcal{L}_{\mathbf{m}\mathbf{B}\mathbf{B}} \\ \bar{D}^{(*)} & B \\ \bar{D}^{(*)} & B \\ \bar{D}^{(*)} & B \\ \bar{D}^{(*)}\mathcal{L}_{\mathbf{C}} & \sigma \\ Fig: \text{ Meson exchange diagram} \\ V^{\pi}_{\bar{D}^{(*)}B-\bar{D}^{(*)}B} = G \left[\vec{\mathcal{O}}_{1} \cdot \vec{\mathcal{O}}_{2}C(r) + S_{\mathcal{O}_{1}\mathcal{O}_{2}}T(r) \right] \\ C(r): \text{ Central force, } T(r): \text{ Tensor force} \\ \bullet \text{ Form factor with common cutoff } \Lambda \leftarrow \text{Free parameter} \\ \end{array}$$

 $F(\Lambda, \vec{q}\,) = \frac{\Lambda^2 - m_{\alpha}^2}{\Lambda^2 + |\vec{q}\,|^2} \quad \text{(fixed by the observed mass of } P_{\rm c}\text{)}$

Determination of cutoff Λ by observed $P_{\rm c}^+$

▷ Narrow resonance $P_{\rm c}^+(4450)$ (12 σ)

 $\exists \rightarrow$

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Determination of cutoff Λ by observed P_c^+

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 \rightarrow In our results, only the $J^P=5/2^-$ state appears above the $\bar{D}\Sigma_{\rm c}^*$ threshold!



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 $J^P = 3/2^+$ state: M = 4339.7 MeV in $\Lambda = 1400$ MeV

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Other predicted states



• In $\Lambda = 1400$ MeV,

 $J^P = 3/2^-$: 4136.0 MeV, 4307.9 MeV and 4348.7 MeV $J^P = 3/2^+$: 4206.7 MeV

New states are predicted!

(can be decayed to $J/\psi p$, $\bar{D}^{(*)}\Lambda_{c},...$)

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• Obtained mass with Full channel coupling, without $\bar{D}^{(*)}\Lambda_c$ and without $\ell>0$ $(\ell>1)$



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• Obtained mass with Full channel coupling, without $\bar{D}^{(*)}\Lambda_c$ and without $\ell>0$ $(\ell>1)$



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• Obtained mass with Full channel coupling, without $\bar{D}^{(*)}\Lambda_c$ and without $\ell > 0$ ($\ell > 1$)



• $\overline{D}^{(*)}\Lambda_c$ and $\ell > 0$ ($\ell > 1$) components are not negligible.

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• Obtained mass with Full channel coupling, without $\bar{D}^{(*)}\Lambda_{\rm c}$ and without $\ell>0$



• $\overline{D}^{(*)}\Lambda_c$ and $\ell > 0$ components are important!

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Subject: Hidden-charm pentaquarks as a meson-baryon molecule

- Coupled-channel analysis is peformed, taking into account...
 - ▷ Meson-baryon components $\bar{D}^{(*)}\Lambda_c \bar{D}^{(*)}\Sigma_c^{(*)}$
 - ▷ Couplings to states with $\ell > 0$ (S D G, P F H) ⇒ Tensor force
 - $\triangleright~$ Negative and positive parity states with $J^{P}=3/2^{\pm},5/2^{\pm}$
- The meson exchange potential respecting to **the heavy quark spin symmetry** is employed.
- The J^P assignment of $P_c^+(4380)$ and $P_c^+(4450)$ is $3/2^+$ and $5/2^-$, respectively.
- New states are predicted in $J^P = 3/2^{\pm}$.

Outlook

• Coupling to $J/\psi p$, cutoff Λ , $1/m_Q$ correction,...

Back up

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Table : Obtained masses in $\Lambda = 1400$ MeV.

JP	$E = E_{\rm re} - i\Gamma/2$ [MeV]
3/2-	4136.0, 4307.9 - <i>i</i> 18.8, 4348.7 - <i>i</i> 21.1
$3/2^{+}$	4206.7 — <i>i</i> 41.2, 4339.7 — <i>i</i> 26.8
$5/2^{-}$	4428.6 - <i>i</i> 89.1

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Table : Obtained masses with full channel coupling (Full), without $\bar{D}^{(*)}\Lambda_c$ (w/o $\bar{D}^{(*)}\Lambda_c$) and without large orbital angular momentum ℓ (w/o $\ell > 0$ or w/o $\ell > 1$) in $\Lambda = 1400$ MeV.

J^P	Channels	Mass [MeV]
3/2-	Full	4136.0, 4307.9, 4348.7
	w/o $ar{D}^{(*)} \Lambda_{ m c}$	4278.4, 4400.4
	w/o $\ell > 0$	4220.4, 4376.6
3/2+	Full	4206.7, 4339.7
	w/o $ar{D}^{(*)} \Lambda_{ m c}$	—
	w/o $\ell > 1$	4275.3
5/2-	Full	4428.6
	w/o $ar{D}^{(*)} \Lambda_{ m c}$	—
	w/o $\ell > 0$	—

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Table : Comparison of the lowest mass of hidden-charm meson-baryon molecules with $I(J^P) = 1/2(3/2^-)$ by this work with the early works. The obtained masses are shown in the second column in the unit of MeV. The value of this work is in $\Lambda = 1400$ MeV. The third column gives the channels which are considered in those works.

Ref.	Mass [MeV]	Channels
This work	4136.0	$ar{D} \Lambda_{ m c}, ar{D}^* \Lambda_{ m c}, ar{D} \Sigma_{ m c}, ar{D} \Sigma_{ m c}^*, ar{D}^* \Sigma_{ m c}, ar{D}^* \Sigma_{ m c}^*$
PRL 105 (2010)232001	4415	$ar{D}^*\Sigma_{ ext{c}},ar{D}^*\Sigma_{ ext{c}}^*$ with only S -wave
PRC 84 (2010)015202	4454	$ar{D}^*\Sigma_{ m c},ar{D}^*\Sigma_{ m c}^*$ with only S -wave
PRD88(2013)056012	4334.5	$J/\psi N, ar{D}^* \Lambda_{\mathrm{c}}, ar{D}^* \Sigma_{\mathrm{c}}, ar{D} \Sigma_{\mathrm{c}}^*, ar{D}^* \Sigma_{\mathrm{c}}^*$
		with only <i>S</i> -wave

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