# Review on recent results from RHIC polarized collider; unexpected forward neutron asymmetry 

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## Outline

- RHIC
- Polarized Proton
- STAR \& PHENIX
- Transverse
- Soft Process (very forward n)

- Longitudinal
- Gluon Spin
- Sea Quark Spin


Versus


## The Relativistic Heavy Ion Collider accelerator complex at Brookhaven National Laboratory





STAR Detector

High resolution and Large acceptance and rates

$$
|\eta|<0.35, \Delta \phi \sim \pi
$$

$\gamma, \pi^{0}$ and $\eta$
Jet and correlation

$$
1.2<|\eta|<2.4
$$

$$
2<|\eta|<4
$$

Target observables


Forward $\pi^{0}$ and very forward neutron

## TRANSVERSE SINGLE SPIN ASYMMETRY

## Transverse Asymmetry Observables



## Production Mechanism of Forward Neutron

## Cross Section



Momentum Transfer ~ 100 MeV


Well Explained by One-Pion Exchange

# $p^{\uparrow}+p$ Forward Neutron $A_{N}$ 


$\delta:$ phase shift
Unpolarized Cross Section

# $p^{\uparrow}+p$ Forward Neutron $A_{N}$ 

Spin non-flip

Neutron $\uparrow$
Spin flip
Neutron


$$
A_{N} \approx \frac{\left(\phi_{\text {non-flip }}^{*} \phi_{\text {flip }} \sin \delta\right)}{\left|\phi_{\text {non-fip }}\right|^{2}+\left|\phi_{f i p}\right|^{2}}
$$

$\delta:$ phase shift

## $p^{\uparrow}+p$ Forward Neutron $A_{N}$



Data are well reproduced by the interference between $\pi$ and $a_{1}$ Reggeon

## The First Time Ever

 High Energy $p^{\uparrow}+A$ CollisionsRun15 (2015)
$100 \mathrm{GeV} /$ nucleon




## What is going on?



|  | \# of protion | \# of neutrion |
| :---: | :---: | :---: |
| P | 1 | 0 |
| Al | 13 | 14 |
| Au | 79 | 118 |

- Isospin Symmetry
- Surface Structure of Nucleus
- QED Process
- Gluon Saturation
- Else...


## Primakoff Effect Electro-Magnetic neutron



A


## Primakoff Effect (Electro-Magnetic) <br> neutron

A



## Full Description of $A_{N}$

$$
\begin{aligned}
& A_{N} \propto 2 \operatorname{Im}\left\{\phi_{n o n-f i p}^{*} \phi_{f l i p}\right\} \\
& \phi_{f l i p}=\phi_{\text {flip }}^{h a d}+\phi_{f l i p}^{E M} \\
& \phi_{\text {non-flip }}=\phi_{\text {non-flip }}^{\text {ad }}+\phi_{\text {non-flip }}^{E M}
\end{aligned}
$$

$$
A_{N} \propto 2 \operatorname{Im}\left(\phi_{\text {non-flip }}^{\text {had }}+\phi_{\text {non-fipp }}^{E M}\right)\left(\phi_{f l i p}^{\text {had* }}+\phi_{f l i p}^{E M^{*}}\right)
$$

$$
=2\left(\phi_{\text {non-fip }}^{\text {had }} \phi_{f l i p}^{\text {had }} \sin \delta_{1}+\phi_{\text {non-flip }}^{E M} \phi_{f i i p}^{\text {had }} \sin \delta_{2}+\phi_{\text {non-fip }}^{\text {had }} \phi_{f i p}^{E M} \sin \delta_{3}+\phi_{\text {non-fip }}^{E M} \phi_{f l i p}^{E M} \sin \delta_{4}\right)
$$

$$
\rightarrow 0(\text { For } p p)
$$

large $\quad-\boldsymbol{- t}$
small

## Beam-Beam Counter



## Can we identify Primakoff events?

## Semi-Inclusive



Primakoff MC : SOPHIA
G. Mitsuka, Eur. Phys. J.C. (2015) 75:614


## BBC Tagging and Vetoing



$$
A_{N} \sim h a d * h a d+h a d * E M+E M * h a d+E M * E M
$$

## BBC Tagging and Vetoing



## BBC Tagging and Vetoing



## Coulomb-Nuclear Interference in

 Forward Neutron Production

$$
A_{N} \sim h a d * h a d+h a d * E M+E M * h a d+E M * E M
$$

## Non-Diffractive Events



$$
A_{N} \sim h a d * h a d+\frac{h a d * E M+E M * h a d}{\text { Suppressed? }}+E M * E M
$$

## Transverse Asymmetry Observables



## Single Transverse Spin Asymmetries

Left Right


Theory Expectation:
Small asymmetries at high energies (Kane, Pumplin, Repko, PRL 41, 1689-1692 (1978))
$A_{N} \propto \frac{m_{q}}{\sqrt{S}} \quad \begin{aligned} & A_{N} O\left(10^{-4}\right) \\ & \text { Theory }\end{aligned}$

## Experiment:

(E704, Fermi National Laboratory, 1991)

$$
\begin{aligned}
& p p^{\uparrow} \rightarrow \pi+X \\
& \sqrt{s}=20 \mathrm{GeV}
\end{aligned} \quad \mathrm{~A}_{N} \mathrm{O}\left(10^{-1}\right) \text { Measured }
$$

E704: Left-right asymmetries


## Energy Dependence of $A_{N}$

RHIC


Non-perturbative
Perturbative

## What Could be the Cause?



- Sivers Effect +higher twist
- transverse-momentum dependence of partons inside the transversely-polarized nucleon
- Collins Effect +higher twist
- correlation between transversely-polarized nucleon and transversely polarized partons inside
- Collins fragmentation function


## Activities around $\pi^{0}$




More Isolated


Less Isolated

The more isolated $\pi^{0}$, the larger the $A_{N}$. Smaller $A_{N}$ for jet-like events?

## Event Topology: Forward TSSAs



## Event Topology: Forward TSSAs



## Transverse Summary

- Very drastic and unexpected A-dependence was observed in very forward neutron $A_{N}$.
- According to MC study, EM process likely to be in charge of the large positive asymmetry in large $A$.
- A-dependence in strong interaction is likely to be moderate as existing model predicts.
- Forward $\pi$ has been studied in PQCD framework, but recent data indicate possibility of soft process may be (partially) playing a role.

So far, "diffractiveness" can be the common key word of both asymmetries.

$\Delta G, \Delta \bar{q} \quad$ Polarization
LONGITUDINAL PROGRAM

## Proton Spin Decomposition

Proton spin $=\frac{1}{2} \quad$ Longitudinal Spin Sum Rule

$$
\begin{gathered}
S_{z}=\frac{1}{2} \Delta \Sigma+\Delta G+L_{z} \\
\downarrow \\
\sim 30 \% \quad ? \quad ? \\
\hline
\end{gathered}
$$

Quark Spin $\Delta \Sigma=\int(\Delta u+\Delta d+\Delta s+\underset{\text { Sea quark }}{\Delta \bar{u}+\Delta \bar{d}}+\Delta \bar{s}) d x$
Gluon Spin $\Delta G(x)=\int \Delta g(x) d x$

## Longitudinal Spin Structure

 Longitudinally Polarized proton-protonHadron
Asymmetry

$p p \rightarrow h X$


Fragmentation Function
Double Helicity Asymmetry $A_{L L}$

$$
A_{L L}=\frac{d \sigma^{++}-d \sigma^{+-}}{d \sigma^{++}+d \sigma^{+-}}=\frac{\sum_{a, b} \Delta f_{a} \otimes \Delta f_{b} \otimes d \hat{\sigma}^{f_{a} f_{b} \rightarrow f X} \cdot \hat{a}_{L L}^{f_{a} f_{b} \rightarrow f X} \otimes D_{f}^{h}}{\sum_{a, b} f_{a} \otimes f_{b} \otimes d \hat{\sigma}^{f_{a} f_{b} \rightarrow f X} \otimes D_{f}^{h}}
$$

## Subprocess of pp



## Latest $A_{\text {LL }}$ Results



## $\Delta g(x)$ extraction : Global Fit

arXiv:1602.03922

$\int_{0.05}^{0.2} \Delta g(x) d x=0.2 \pm_{0.07}^{0.06}$

## Sea Quark Polarization



## Single Lepton Measurement




## Single Lepton Asymmetry $A_{L}$




## Summary

- RHIC spin program has been providing new insights to spin structure of proton
- Absolutely unexpected strong A-dependence in very forward neutron TSSA of $p+A$.
- EM process likely to be taken into account.
- More sophisticated measurements on forward $\pi$ TSSA to disentangle competing effects (ISI or FSI).
- Diffractiveness may connects between forward $\pi$ and very forward $n$.
- Latest data shows non-zero gluon spin fraction
- Ongoing W analyses will improve precision of seaquark polarized parton distribution

BACKUP SLIDES

## Event Generation of UPCs



## Event Generation of UPCs II

## Flux of quasi photons

Cross section of $\mathrm{p}-\mathrm{y}$

## Event Generation of p-y

## Flux of quasi photons

## Weizsacker-Williams method

 $\frac{d N_{\gamma^{*}}}{d E_{\gamma^{*}}}=\frac{2 Z^{2} \alpha}{\pi E_{\gamma^{*}}}\left[x K_{0}(x) K_{1}(x)+\frac{x^{2}}{2}\left(K_{0}^{2}(x)-K_{1}^{2}(x)\right)\right]$


## The LHGf experiment



# UPC Monte Carlo 

## Forward hadron production in ultra-peripheral proton-heavy-ion collisions at the LHC and RHIC

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Received: 26 April 2015 / Accepted: 15 December 2015 / Published online: 26 December 2015 © The Author(s) 2015. This article is published with open access at Springerlink.com

Abstract We present a hadron production study in the forward rapidity region in ultra-peripheral proton-lead ( $p+\mathrm{Pb}$ ) collisions at the LHC and proton-gold ( $p+\mathrm{Au}$ ) collisions at RHIC. The present paper is based on the Monte Carlo simulations of the interactions of a virtual photon emitted by a fast moving nucleus with a proton beam. The simulation consists of two stages: the STARLIGHT event generator simulates the virtual photon flux, which is then coupled to the SOPHIA, DPMJET, and PYTHIA event generators for the simulation of particle production. According to these Monte Carlo simulations, we find large cross sections for ultra-peripheral collisions particle production, especially in the very forward region. We show the rapidity distributions for charged and neutral particles, and the momentum distributions for neutral pions and neutrons at high rapidities. These processes lead to substantial background contributions to the investigations of
virtual photons eı may anyway ints usually referred Ref. [1,2] for a r UPCs, so far, the gluon distrib photoproduction collisions can pı density in proto momentum fract measurements al tion at the CERl $p+\mathrm{Pb} \rightarrow p+\mathrm{Pb}+J / \psi$ [3]. has been paid, in UPCs, to par photon-proton interactions, i.e., Pseudorapidity

## Coulomb-Nuclear Interference

 (CNI)


$$
A_{N} \sim h a d * h a d+h a d * E M+E M * h a d+E M * E M
$$

## Elastic $A_{N}$ at Coulomb Nuclear Interference

$$
A_{N} \sim 2 \operatorname{Im}\left\{\phi_{\text {non-flip }}^{E M *} \phi_{\text {flip }}^{\text {had }} \sin \delta_{2}+\phi_{\text {non-fip }}^{\text {had }} \phi_{\text {flip }}^{E M} \sin \delta_{3}\right\}
$$




had

$16 / 06 / 27 \mathrm{E}_{\text {beam }}=21.7 \mathrm{GeV}$

## Run15 Au,AI beam + $\mathrm{p}^{\uparrow}$ target



- Strong A-Dependence
- Flips sign of $A_{N}$ in $A u+\hat{p}$
- $0.002<-t<0.014(\mathrm{Gev} / \mathrm{c})^{2}$


## Forward Neutron $0.02<-\mathrm{t}<0.5(\mathrm{Gev})$ c) ${ }^{2}$

## Underlying Mechanism Comparison

## Polarimeter

## Forward n



> | Elastic | Inelastic |
| :---: | :---: |
| $\sqrt{ } \mathrm{s}=14 \mathrm{GeV}$ | $\sqrt{ } \mathrm{s}=200 \mathrm{GeV}$ |
| $0.002<\cdot \mathrm{t}<0.014$ | $0.02<\cdot \mathrm{t}<0.5$ |
| $\Delta I=0$ | $\Delta I=1$ |

## Interim Summary (CNI)

- CNI is not necessarily new.
- Elastic in much smaller - $t$ than present neutron results observed similar behavior, but the asymmetry is in the order of $5 \%$ or less.
- There are possible diagram to hadron and EM interfere for neutron production as well.


## Measurement of the Analyzing Power in the Primakoff Process with a High-Energy Polarized Proton Beam

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# The origin of Asymmetry 

cording to Eq. (1), the asymmetry seen in photoproduction due to the interference between $\Delta$ and $N^{*}$ is expected in coherent Coulomb $\pi^{0}$ production by polarized protons, using the same region of the $\pi^{2}-p$ invariant mass. Therefore this process may be used to measure the polarization of the proton at high energies. ${ }^{6}$ Until now, there has been no measurement of the asymmetry in the nuclear coherent process.

The cross section for the Coulomb coherent process (1) has a sharp peak at $t^{\prime} \sim 10^{-5}(\mathrm{GeV} / c)^{2}$ and decreases rapidly as $t^{\prime} / t^{2}$. The "width" of the Coulomb peak is determined by the detector resolution. Diffractive dissociation due to the strong interaction is also present, but it has a much slower $t^{\prime}$ dependence.

We have measured the analyzing power (azimuthal asymmetry) of nuclear Coulomb coherent production from a Pb target by using the newly constructed 185$\mathrm{GeV} / c$ Fermilab polarized proton beam. ${ }^{7}$ The beam polarization is $45 \%$ and this is further described in Ref. 7. To reduce certain systematic errors, the spin direction of the incident proton was flipped every 10 min using a spin-rotator system. ${ }^{7}$
monitored to a $1 \%$ accuracy with a xenon-flash-tube system. A $30-\mathrm{GeV}$ positron beam was used to calibrate the calorimeter. The measured energy resolution is $3 \%$ (rms) at 30 GeV and the position resolution is 2 mm (rms). The measured $\pi^{0}$ energies in this experiment ranged from 25 to 75 GeV .

A set of thin plastic scintillation counters (TP1) is placed downstream of the magnet and provides the trigger for the scattered protons. The set consists of four counters arranged to distinguish protons scattered to the left, right, up, and down. The calorimeter also has left, right, up, and down sections, and signals from each section are summed for the trigger. In the coherent process where $t^{\prime}$ is almost zero, the $\pi^{0}$ and scattered protons are coplanar. Thus the trigger logic is such that the energy deposit is larger than 25 GeV in the left half of the calorimeter, less than 5 GeV in the right half, and a proton hits the right segment of TP1. There are four such combinations to cover the whole range of azimuthal angles. To reject the events which have any extra particle besides a proton and $\pi^{0}$, veto counters are included in the trigger logic.

## Origin of the Asymmetry


$\Delta^{*}$ (1232) P33
$\leftrightarrows \mathrm{N} \pi \quad(>99 \%)$

$N^{*}(1440) P 11$
$\leftrightarrows N \pi \quad(60 \sim 70 \%)$
Nлл (30~40\%)
$A_{N} \propto \phi_{\text {non-fip }}^{\text {had }} \phi_{f t i p}^{\text {had }} \sin \delta_{1}+\phi_{\text {non-flip }}^{E M} \phi_{f l i p}^{\text {had }} \sin \delta_{2}+\phi_{\text {non-flip }}^{\text {had }} \phi_{f i p}^{E M} \sin \delta_{3}+\phi_{\text {non-fip }}^{E M} \phi_{f t i p}^{E M} \sin \delta_{4}$

## Comparison between two

|  | Fermi | PHENIX | STAR17 |
| :---: | :---: | :---: | :---: |
| Beam Energy [GeV] | 185 | 100 | 255 |
| $\sqrt{ }$ [ [GeV] | 19.5 | 200 | 22 |
| Target | Pb | Au | $\mathrm{Al} / \mathrm{Sn} / \mathrm{Au}$ |
| Observables | $\mathrm{p}+\pi^{0}$ | n ( + charged) | $\begin{gathered} \mathrm{n}(+\underset{\text { charged }}{ }) \\ \pi^{0} ? \end{gathered}$ |
| t' | < 0.001 | $0.02<-\mathrm{t}<0.5$ |  |
| M | $1.36<M\left(\pi^{0} \mathrm{p}\right)<1.52$ | ? | ? |
| $\mathrm{A}_{\mathrm{N}}$ | $\begin{gathered} 0.57 \pm(0.12)_{\mathrm{sta}}+ \\ 0.21 \cdot 0.18 \end{gathered}$ | $+\underset{(B B C \text { veto })}{0.27 \pm 0.003}$ |  |

## Origin of the Asymmetry


$A_{N} \propto \phi_{\text {non-flip }}^{\text {had }} \phi_{f i p}^{\text {had }} \sin \delta_{1}+\phi_{\text {non-flip }}^{E M} \phi_{f i p}^{\text {had }} \sin \delta_{2}+\phi_{\text {non-fip }}^{\text {had }} \phi_{f i p}^{E M} \sin \delta_{3}+\phi_{\text {non-flip }}^{E M} \phi_{f i p}^{E M} \sin \delta_{4}$



Primakoff


Fermi ( $p+\pi^{0}$ channel)



Primakoff


RHIC ( $n+\pi^{+}$channel)

PH F ENIX

## Theory: $\mathrm{K}_{\mathrm{T}}$ vs Collinear Factorization

- Tran. Mom. Dep. Funs
- Sivers Fun
- Collins Fun


$$
\frac{d^{3} \sigma^{\uparrow}\left(p p^{\uparrow} \rightarrow h+X\right)}{d x_{1} d x_{2} d z} \propto q_{i}^{\uparrow}\left(x_{1}, k_{q, T}\right) \cdot q_{j}\left(x_{2}\right) \times \frac{d^{3} \hat{\sigma}^{\uparrow}\left(q_{i} q_{j} \rightarrow q_{k} q_{l}\right)}{d x_{1} d x_{2}} \times F F_{q_{k, l}}\left(z, p_{h, T}\right)
$$

- Twist-3 collinear
- Quark-gluon correl.
- Gluon-gluon correl.



## Future



## Present and Future

arXiv: 1501.01220



arXiv: 1212.1701


## Transverse Summary



## Large + positive <br> Diffractive?

Negative + Moderate
A-dependence


## $x_{F}$ and $p_{T}$ Dependence


$\mathrm{s}^{1 / 2}=200 \mathrm{GeV}$ and 500 GeV show same rise of $A_{N}$ vs. $x_{F}$ as lower $s^{1 / 2}$ measurements


Steve Heppelmann - CIPANP 2012

- Collins, Sivers, Twist-3 suggest $A_{N} \sim 1 / p_{T}$
$\square$ Flat $P_{T}$-dependence observed and raises the question as to what causes it


## Forward $\pi^{0} A_{N}$ in $p+A u$



Did not observe A-dependence unlike very forward neutron $A_{N}$ in PHENIX

