New Measurements of Hyperon Production in the Charmonium Region

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Introduction

- The study of the production of hadrons gives unique insight into their structure. The strangeness-containing **hyperons** are especially interesting since they provide a family of states that allow us to study how their properties evolve with:
 - Strange quark content (1–3 strange quarks)
 - Isospin (I=0,1)
 - Momentum transfer (|Q²|)
- Hadronization in e⁺e⁻ annihilation is an exceptionally clean laboratory for hyperon production.
 - Long history of studying these processes at Cornell (CLEO), SLAC (HRS,Mark II,TPC), DESY (CELLO,TASSO), LEP (DELPHI,L3,OPAL)
 - Can study the difference between quark-initiated (e+e⁻ → qq̄) versus gluon-initiated production (e.g, e+e⁻ → Y/ψ → ggg)

Hyperon Production in e+e- annihilation

 CLEO III measurements of hadron production in
 e⁺e⁻ → qq̄ @ √s ~ 10 GeV and Y(nS) → ggg show enhancement of ∧ production in Y(nS) decay.



Hyperon Production in e⁺e⁻ annihilation



- Measurements from LEP at the Z⁰ peak show:
 - Suppression of Ξ^- (2 s-quark) compared to Λ^0 (1 s-quark)
 - Suppression of $I=1 \Sigma^0$ compared to $I=0 \Lambda^0$, possible **diquark** effect

Data Sets & Hyperon Reconstruction

- We measure e⁺e⁻ → hyperon + X using e⁺e⁻ annihilation data taken with the CLEO-c detector at
 - $\psi(2S) [\sqrt{s} = 3686 \text{ MeV}]:$ 48 pb⁻¹, 25 M $\psi(2S)$
 - ψ(3770) [√s = 3772 MeV]: 805 pb⁻¹
 - $\sqrt{s} = 4170 \text{ MeV}$: 586 pb⁻¹
- ψ(2³S₁,3686) decays primarily to ggg, lighter charmonia
 - Expect hyperons to be produced primarily via
 e⁺e⁻ → (cc̄) → gluons → hyperon + X
- ψ(1³D₁,3770), ψ(2³D₁,4160) decay primarily to D mesons
 - Expect hyperons to be produced primarily from
 e⁺e⁻ → χ^{*} → qq̄ → hyperon + X
- We identify the hyperons by their dominant decays:

 $\begin{array}{ll} \Lambda^{0} \rightarrow p\pi^{-} & (64 \%) & \Sigma^{+} \rightarrow p\pi^{0} & (52 \%) & \Sigma^{0} \rightarrow \Lambda^{0} \mathbf{y} & (100 \%) \\ \hline \Xi^{-} \rightarrow \Lambda^{0} \pi^{-} & (100 \%) & \Xi^{0} \rightarrow \Lambda^{0} \pi^{0} & (100 \%) & \Omega^{-} \rightarrow \Lambda^{0} \text{ K}^{-} & (68 \%) \end{array}$

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Inclusive Hyperon Mass Spectra From ψ(2S)

 $\psi(2S) \rightarrow B \text{ or } \overline{B} + X$, (Preliminary, Statistical uncertainties only)



- Well reconstructed particles with a vertex displaced from the e+e-IP are selected. Strong peaks are seen in each channel.
- We also measure yields as a function of hyperon momentum.

Hyperon Yields From ψ(2S)

 $\psi(2S) \rightarrow B \text{ or } \overline{B} + X$, (Preliminary, Statistical uncertainties only)



- Pair-production peaks are clearly seen.
- Yields from Σ^{0}/Σ^{+} and Ξ^{0}/Ξ^{-} have similar shapes.

Hyperon Yields From $\sqrt{s} = 3772, 4170 \text{ MeV}$

 $e^+e^- \rightarrow B \text{ or } \overline{B} + X$, (Preliminary, Statistical uncertainties only)



- 4170 data have yields
 1.5—2x less than 3772 data (1.4x smaller luminosity).

Hyperon Yields From $\sqrt{s} = 3772, 4170 \text{ MeV}$

 $e^+e^- \rightarrow B \text{ or } \overline{B} + X$, (Preliminary, Statistical uncertainties only)

√s = 3772 MeV



 3772 data and 4170 data have similar shapes, differ from ψ(2S) data.

Hyperon Yields From $\sqrt{s} = 3772, 4170 \text{ MeV}$

 $e^+e^- \rightarrow B \text{ or } \overline{B} + X$, (Preliminary, Statistical uncertainties only) Open hist: $\sqrt{s}=3772 \text{ MeV}$ Filled hist: $\sqrt{s}=4170 \text{ MeV}$ data (1/s² norm.)



- Yields from both data are consistent after normalizing √s=4170 MeV data by additional factor of 1/s².
- Note that Spacelike proton FF has 1/√s dependence, opposed to expected 1/s dependence

Theoretical Models

- Hyperon production is generally modeled by the LUND string model as implemented in JETSET/PYTHIA.
 - To more accurately model string fragmentation at charm energies, we use the **LUNDCHARM** model, which includes C- and G-parity conservation.
- Baryon production in the LUND model arises from the combination of a diquark with a quark. There are three main parameters that affect this:
 - The probability of qq diquark to quark production
 - The probability to produce a strange su/sd diquark to a ud diquark
 - The probability to produce a spin-1 to a spin-0 diquark
- Properly tuning this model requires many careful measurements, but to start we can compare our measurements with predictions obtained from the standard CLEO-c parameters.

Baryon Production in LUND string model



Hyperon Yields From $\psi(2S)$: Data and MC

 $\psi(2S) \rightarrow ggg \rightarrow B \text{ or } \overline{B} + X$, (Preliminary, Statistical uncertainties only) Open hist: $\psi(2S) DATA$ Filled hist: $\psi(2S) MC$



- $\psi(2S)$ MC is mix of LUND and known b.f.'s fixed to PDG.
- MC reproduces some features of data, gives generally poor agreement.

Hyperon Yields From $\sqrt{s} = 4170$ MeV: Data and MC

 $e^+e^- \rightarrow q\bar{q} \rightarrow B \text{ or }\bar{B} + X$, (Preliminary, Statistical uncertainties only) Open hist: 4170 MeV DATA Filled hist: 4170 MeV MC



Reasonable agreement between data and MC.
 Data spectra are generally slightly softer than MC data.

Unequal Baryons — Production of $\Lambda^0 \Sigma^0$

 Production of Λ⁰ Σ⁰ probes two interesting phenomena: isospin violation in charmonium decays and hyperon transition form factors [see presentation of K. Seth for more details on hyperon pair production]



- The Λ⁰ (I=0) and Σ⁰ (I=1) states have the same uds quark content but opposite isospin.
 - Λ⁰ Σ⁰ (total l=1) pair production in e⁺e⁻ annihilation (total l=0) can only proceed through isospin-violating processes.
- In $\psi(2S)$ decay, the strong force conserves isospin, so the decay must be through a virtual photon: $\psi(2S) \rightarrow \chi^* \rightarrow \Lambda^0 \Sigma^0$
- In EM production, this process probes the $\Sigma^0 \rightarrow \Lambda^0$ transition form factor.

Unequal Baryons — Production of $\Lambda^0 \Sigma^0$

 $\psi(2S) \rightarrow ggg \rightarrow \Lambda^0 \Sigma^0$ $e^+e^- \rightarrow q\bar{q} \rightarrow \Lambda^0 \Sigma^0$ 12 12 ψ**(2S)** √s=3770 MeV 0.96 0.96 1.02 0.98 0.98 1.04 1.02 1.04 1 1 $X(\Lambda^0\Sigma^0)$ $X(\Lambda^0\Sigma^0)$

- $\psi(2S)$: N = 30 ± 5, Br = (0.12 ± 0.02(stat)) x 10⁻⁴ This is 20 – 30 times smaller than Br($\psi(2S) \rightarrow \Lambda^0 \Lambda^0, \Sigma^0 \Sigma^0$)
- Form factor @ $|Q^2| = 14.2 \text{ GeV}^2$:

$$\begin{split} N &= 30 \pm 5, \ G_{M}(\Lambda^{0} \Sigma^{0}) = (0.79 \pm 0.07(\text{stat})) \times 10^{-2} \\ \text{This is consistent with} & G_{M}(\Sigma^{0}) = (0.79 \pm 0.07) \times 10^{-2} \\ \text{and is factor} \sim 0.6 \text{ smaller than } G_{M}(\Lambda^{0}) = (1.31 \pm 0.05) \times 10^{-2} \end{split}$$

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Summary

- Detailed study of hyperon production can yield useful insights into their structure.
- We have measured yields of the hyperons Λ^0 , Σ^+ , Σ^0 , Ξ^- , Ξ^0 , Ω^- produced in the decay of $\psi(2S)$ and in e⁺e⁻ annihilation at two energy points above $D\overline{D}$ threshold, $\sqrt{s} = 3772$, 4170 MeV.
 - Detailed studies of efficiencies and MC simulations are being performed.
 - These measurements open the door for measurements of other properties of inclusive hyperon production.
- We have measured for the first time $\Lambda^0 \Sigma^0$ production from $\psi(2S)$ and at $\sqrt{s} = 3772$ MeV, probing isospin violating effects in these processes.

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CLEO-c Detector



Hyperon Yields From $\psi(2S)$ and $\sqrt{s}=3772$ MeV

 $\psi(2S) \rightarrow B \text{ or } \overline{B} + X$, (Preliminary, Statistical uncertainties only) Open hist: $\psi(2S)$ data Filled hist: $\sqrt{s}=3772$ MeV data (arb. norm.)



 Results from ψ(2S) and √s=3772 MeV data have different shapes, as expected from different production process.

CLEO-c MC JETSET Default Parameters

- Diquark suppression factor P(qq)/P(q): PARJ(1) = 0.065 [default = 0.1]
- Strange quark suppression factor P(s)/P(u): PARJ(2) = 0.26 [default = 0.30]
- Strange diquark suppression factor [P(us)/P(ud)]/[P(s)/P(d)]: PARJ(3) = 0.4
- Suppression of spin-1 diquarks relative to spin-0: (1/3)P[ud₁]/P[ud₀]: PARJ(4) = 0.05
- Popcorn production probability: PARJ(5) = 0.5