

Modern dynamical coupled-channels calculations for extracting and understanding the nucleon spectrum

**Hiroyuki Kamano
(KEK)**

Outline

PART I:

Background & motivation for spectroscopic study of N^* & Δ^* resonances

PART II:

Recent results from ANL-Osaka Dynamical Coupled-Channels (DCC) analysis

PART I

Background & motivation for spectroscopic study of N^* & Δ^* resonances

Introduction: N^* & Δ^* spectroscopy

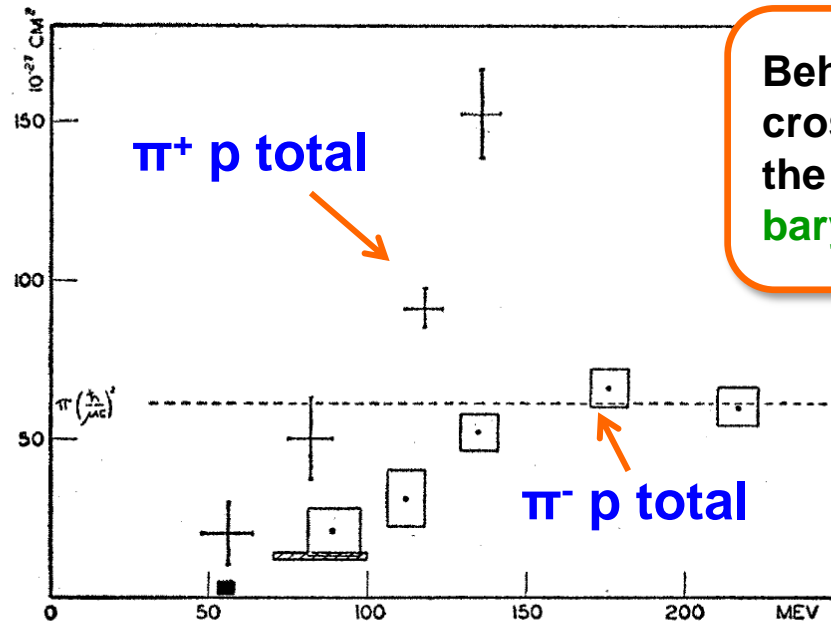
Discovery of the Δ baryon (1952)

Total Cross Sections of Positive Pions in Hydrogen*

H. L. ANDERSON, E. FERMI, E. A. LONG,† AND D. E. NAGLE

*Institute for Nuclear Studies, University of Chicago,
Chicago, Illinois*

(Received January 21, 1952)



Behavior of the $\pi^+ p$ & $\pi^- p$ cross sections implies the existence of a new baryon with isospin 3/2 !!

FIG. 1. Total cross sections of negative pions in hydrogen (sides of the rectangle represent the error) and positive pions in hydrogen (arms of the cross represent the error). The cross-hatched rectangle is the Columbia result. The black square is the Brookhaven result and does not include the charge exchange contribution.

Introduction: N^* & Δ^* spectroscopy

N^*			Δ^*		
Particle	J^P	overall	Particle	J^P	overall
N	$1/2^+$	****	$\Delta(1232)$	$3/2^+$	****
$N(1440)$	$1/2^+$	****	$\Delta(1600)$	$3/2^+$	***
$N(1520)$	$3/2^-$	****	$\Delta(1620)$	$1/2^-$	****
$N(1535)$	$1/2^-$	****	$\Delta(1700)$	$3/2^-$	****
$N(1650)$	$1/2^-$	****	$\Delta(1750)$	$1/2^+$	*
$N(1675)$	$5/2^-$	****	$\Delta(1900)$	$1/2^-$	**
$N(1680)$	$5/2^+$	****	$\Delta(1905)$	$5/2^+$	****
$N(1700)$	$3/2^-$	***	$\Delta(1910)$	$1/2^+$	****
$N(1710)$	$1/2^+$	****	$\Delta(1920)$	$3/2^+$	***
$N(1720)$	$3/2^+$	****	$\Delta(1930)$	$5/2^-$	***
$N(1860)$	$5/2^+$	**	$\Delta(1940)$	$3/2^-$	**
$N(1875)$	$3/2^-$	***	$\Delta(1950)$	$7/2^+$	****
$N(1880)$	$1/2^+$	**	$\Delta(2000)$	$5/2^+$	**
$N(1895)$	$1/2^-$	**	$\Delta(2150)$	$1/2^-$	*
$N(1900)$	$3/2^+$	***	$\Delta(2200)$	$7/2^-$	*
$N(1990)$	$7/2^+$	**	$\Delta(2300)$	$9/2^+$	**
$N(2000)$	$5/2^+$	**	$\Delta(2350)$	$5/2^-$	*
$N(2040)$	$3/2^+$	*	$\Delta(2390)$	$7/2^+$	*
$N(2060)$	$5/2^-$	**	$\Delta(2400)$	$9/2^-$	**
$N(2100)$	$1/2^+$	*	$\Delta(2420)$	$11/2^+$	****
$N(2120)$	$3/2^-$	**	$\Delta(2750)$	$13/2^-$	**
$N(2190)$	$7/2^-$	****	$\Delta(2950)$	$15/2^+$	**
$N(2220)$	$9/2^+$	****			
$N(2250)$	$9/2^-$	****			
$N(2300)$	$1/2^+$	**			
$N(2570)$	$5/2^-$	**			
$N(2600)$	$11/2^-$	****			
$N(2700)$	$13/2^+$	**			

PDG(2015):
<http://pdg.lbl.gov>

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All of these studies essentially agree on the existence and (most) properties of the 4-star states. For the 3-star and lower states, however, even a statement of existence is problematic.

— Arndt, Briscoe, Strakovsky, Workman PRC74(2006)045205

Introduction: N^* & Δ^* spectroscopy

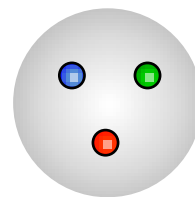
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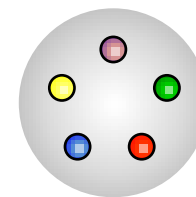
N^* & Δ^* spectroscopy remains as central issue in the hadron physics !!

- Mass, width, spin, parity ...?
- quark-gluon structure (form factors)?
- How produced in reaction processes?

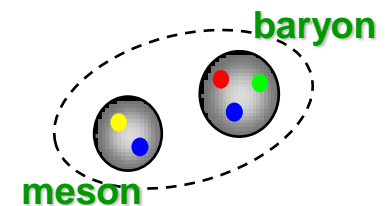
...



3-quark state



multiquark state



"molecule-like" state

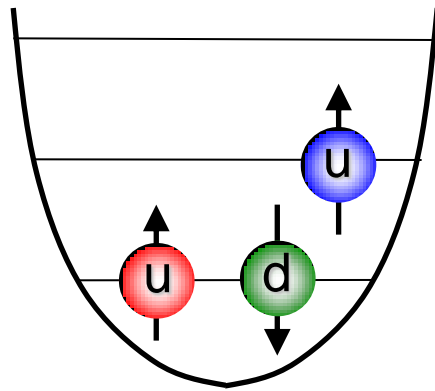
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$\Delta(2950)$	$15/2^+$	**

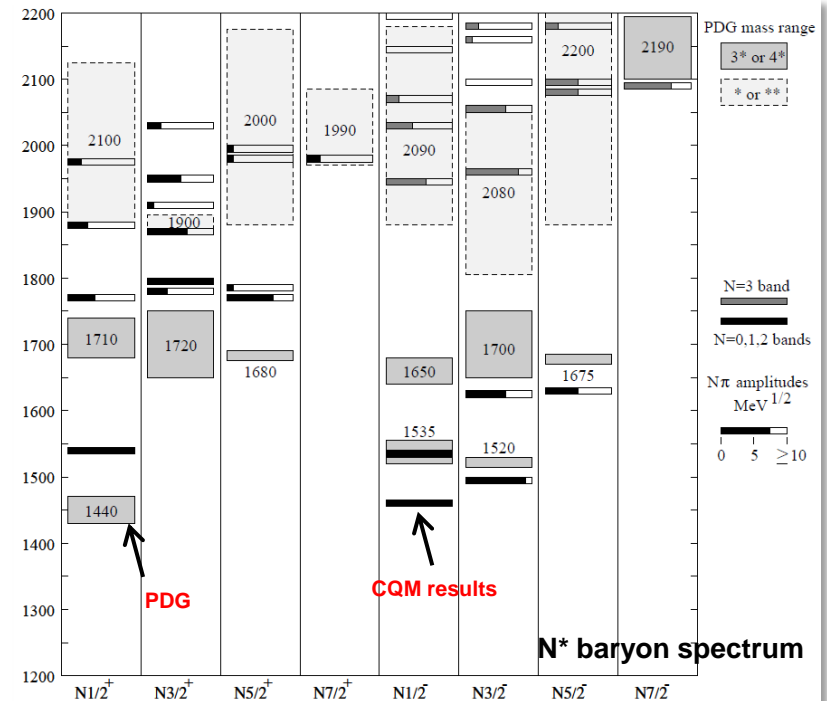
PDG(2015):
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Static hadron models and reaction dynamics

- ✓ Various **static hadron models** have been proposed to calculate hadron spectrum and form factors.
 - **Constituent quark models, soliton models, Dyson-Schwinger Eq. approaches,...**
 - **Excited hadrons** are treated as **stable particles**.



Constituent quark model

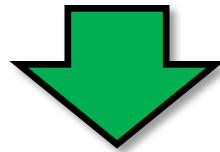


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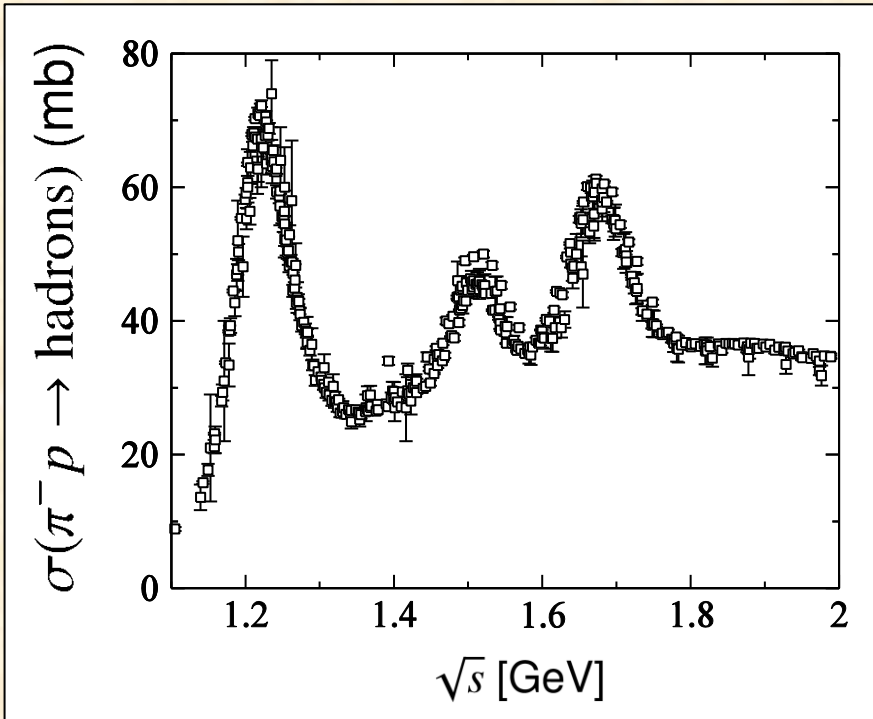
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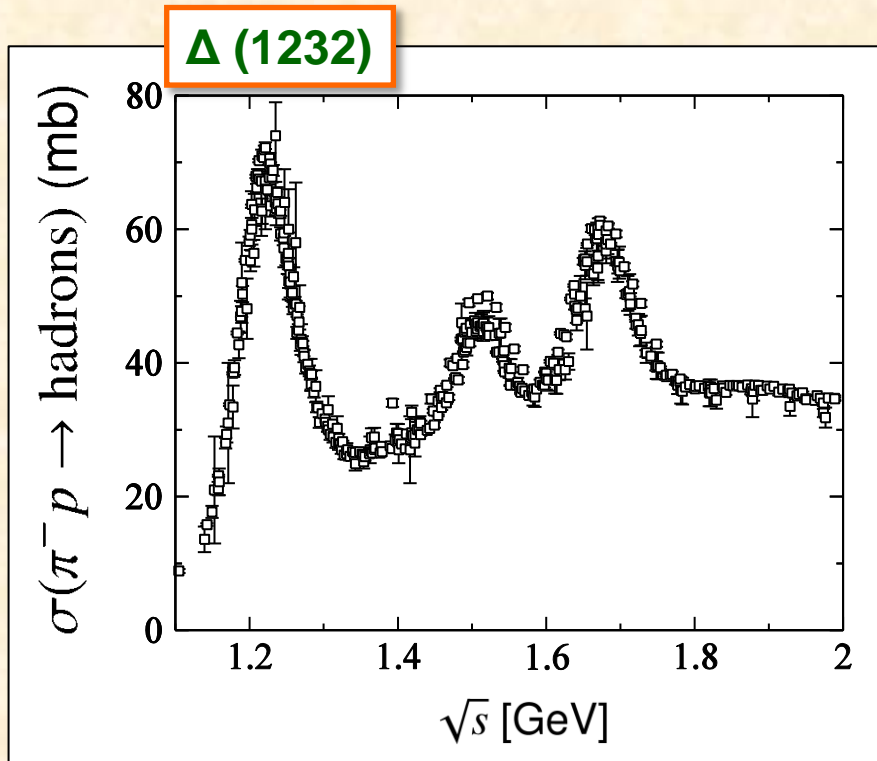
What is the role of **reaction dynamics** in interpreting the hadron spectrum, structures, and dynamical origins ??

Light-quark hadron spectroscopy : Physics of broad & overlapping resonances



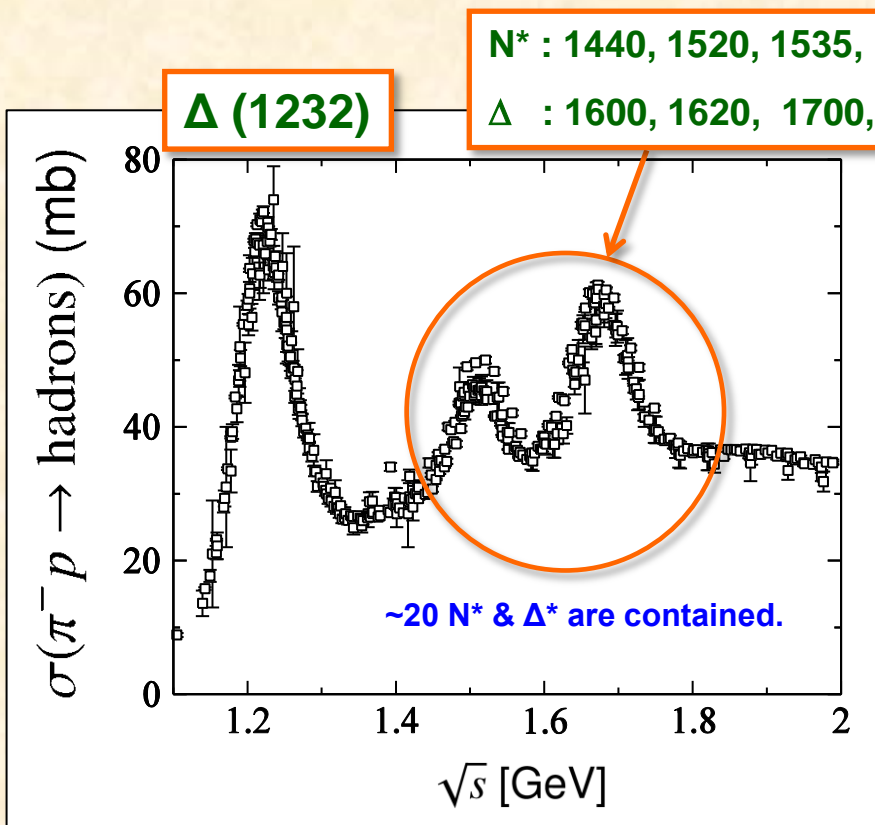
- ✓ Width: a few hundred MeV.
- ✓ Resonances are highly overlapping in energy except $\Delta(1232)$.

Light-quark hadron spectroscopy : Physics of broad & overlapping resonances



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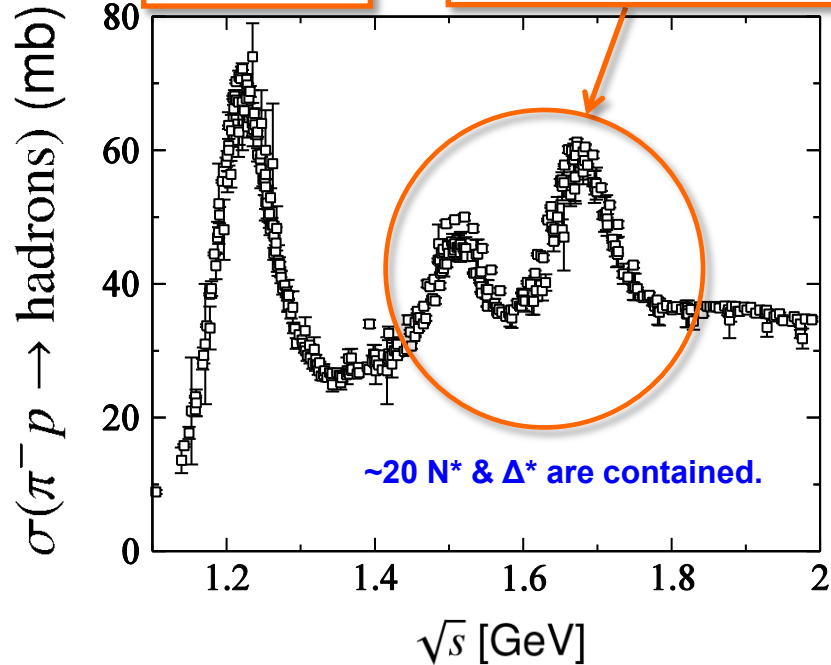
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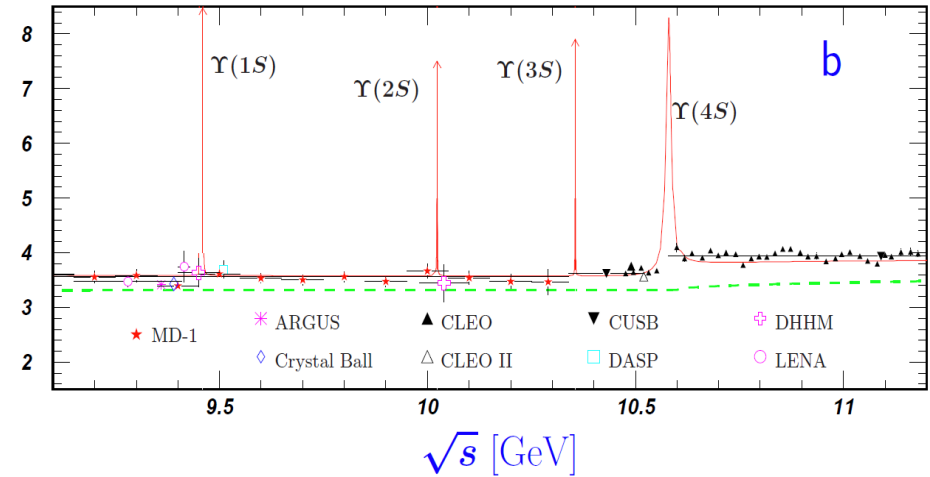
N^* : 1440, 1520, 1535, 1650, 1675, 1680, ...

Δ : 1600, 1620, 1700, 1750, 1900, ...

Δ (1232)



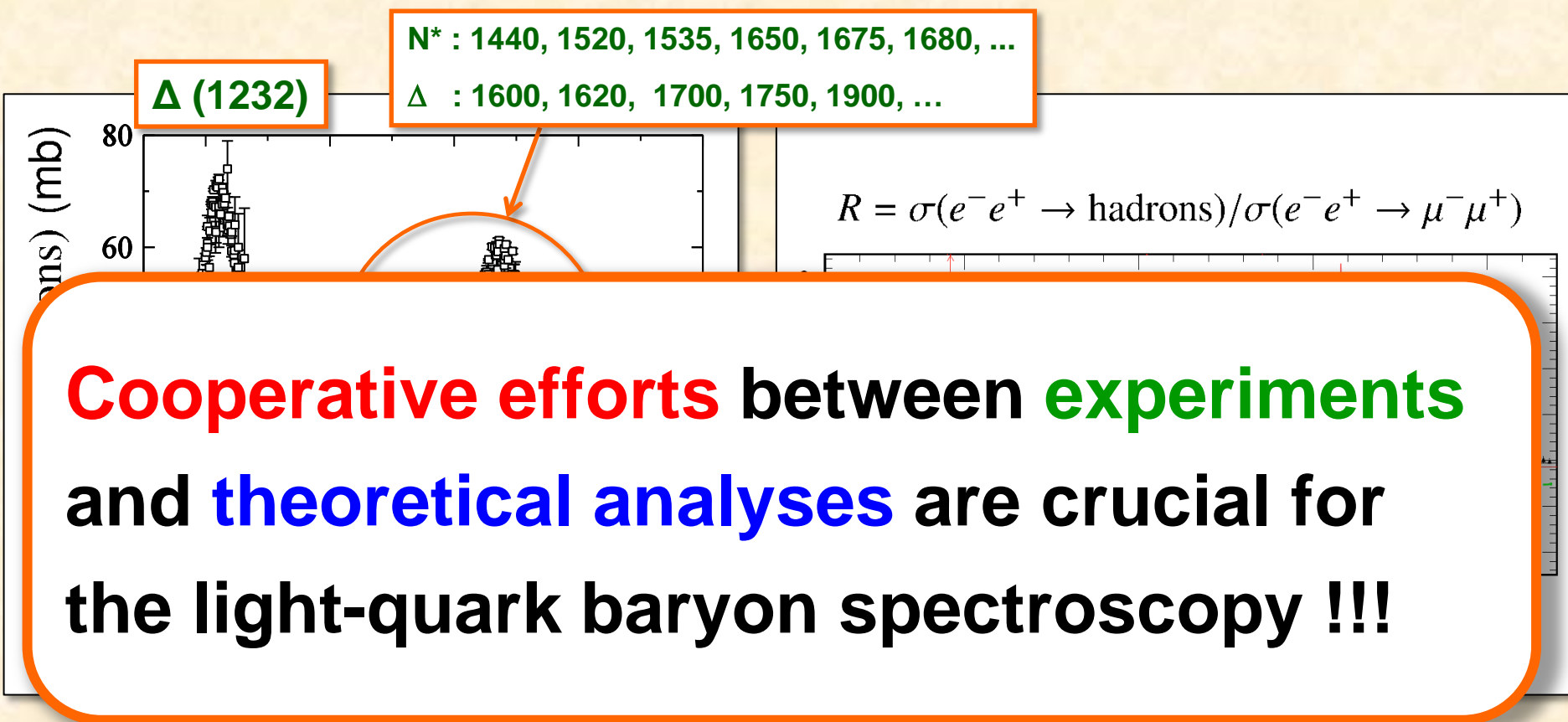
$$R = \sigma(e^-e^+ \rightarrow \text{hadrons})/\sigma(e^-e^+ \rightarrow \mu^-\mu^+)$$



- ✓ Width: a few hundred MeV.
- ✓ Resonances are highly overlapping in energy except $\Delta(1232)$.

- ✓ Width: ~10 keV to tens of MeV
- ✓ Each resonance peak is clearly separated.

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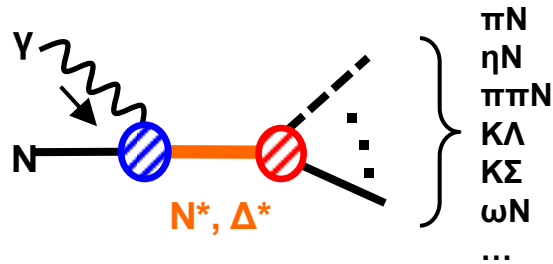
Cooperative efforts between experiments and theoretical analyses

Experiments



Since late 90s

➤ e.g.) Meson photoproductions



ELSA, ELPH, GRAAL, JLab,
MAMI, SPring-8, ...

Theoretical analyses with coupled-channels framework

ANL-Osaka
 Argonne-Pittsburgh
 Bonn-Gatchina
 Carnegie-Mellon-Berkeley
 Dubna-Mainz-Taipei
 EBAC
 Giessen
 GWU/VPI
 Juelich-Bonn
 Karlsruhe-Helsinki
 KSU
 Zagreb
 ...

✓ Multichannel unitary condition:

$$T_{ab}(E) - T_{ab}^\dagger(E) = -2\pi i \sum_c T_{ac}^\dagger \delta(E - E_c) T_{cb}(E)$$

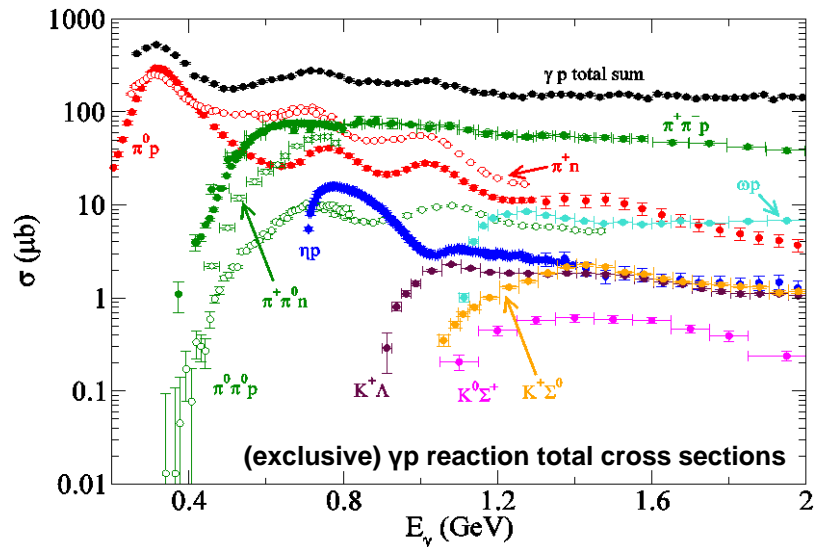
Why multichannel unitarity is so important ??

$$T_{ab}(E) - T_{ab}^\dagger(E) = -2\pi i \sum_c T_{ac}^\dagger \delta(E - E_c) T_{cb}(E)$$

1) Ensures **conservation of probabilities** in multichannel reaction processes

Key essential to

- **simultaneous analysis** of various inelastic reactions.
- increasing **predictivity** of constructed reaction models



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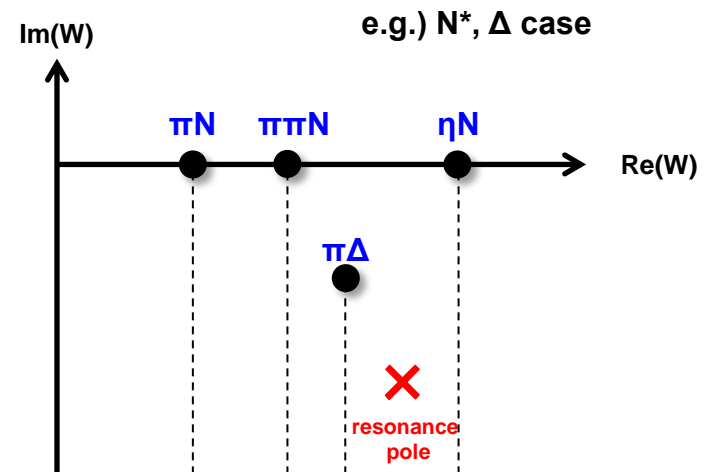
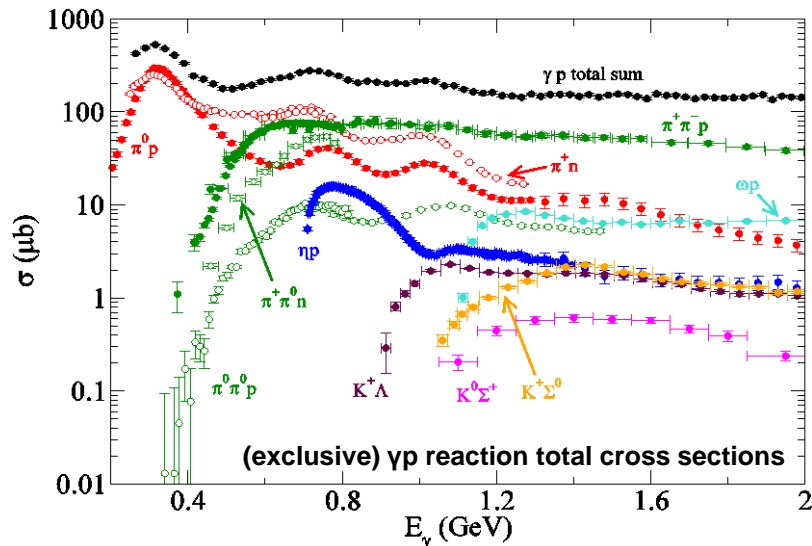
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Key essential to

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- increasing **predictivity** of constructed reaction models

2) Defines **proper analytic structure** (branch points, cuts,...) of scattering amplitudes in the **complex energy plane**, as required by scattering theory

- Crucial for extracting resonances “correctly”, and avoiding **WRONG** resonance signals !!
[e.g., Ceci et al, PRC84(2011)015205]



Approaches for coupled-channels analysis

- ✓ Multichannel unitary condition:

$$T_{ab}(E) - T_{ab}^\dagger(E) = -2\pi i \sum_c T_{ac}^\dagger \delta(E - E_c) T_{cb}(E)$$

- ✓ Heitler equation:

K-matrix: must be hermitian for real E.

$$T_{ab}(E) = K_{ab}(E) + \sum_c K_{ac}(E) [-i\pi\delta(E - E_c)] T_{cb}(E)$$

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➤ **(on-shell) K-matrix approach:**

$K_{ab}(E) \equiv$ (Polynomials of E) + (Pole terms)

✓ Heitler equation can be solved **algebraically**.

Argonne-Pittsburgh, Bonn-Gatchina, Carnegie Mellon-Berkely, GWU/VPI, Karlsruhe-Helsinki, KSU, ...

➤ **Dynamical-model approach:**

✓ Heitler equation becomes **identical to Lippmann-Schwinger integral equation**.

✓ Potential V is derived from a model Hamiltonian.

$$K_{ab}(E) \equiv K_{ab}(\vec{p}_a, \vec{p}_b; E) = V_{ab}(\vec{p}_a, \vec{p}_b; E) + \sum_c \mathcal{P} \int d\vec{q} V_{ac}(\vec{p}_a, \vec{q}; E) \frac{1}{E - H_c^0 + i\epsilon} K_{cb}(\vec{q}, \vec{p}_b; E)$$

off-shell rescattering effect

ANL-Osaka, Dubna-Mainz-Taipei, Juelich-Bonn, ...

Why dynamical coupled-channels approach??

	(on-shell) K-matrix	Dynamical model
Numerical cost	Cheap 😊 - solve on-shell algebraic eq. (Analysis can be done quickly on PC.)	(Very) Expensive 😞 - solve off-shell integral eq. (Supercomputers are needed.)
Data fitting	Efficient 😊 - $K(E)$ can be parametrized as one likes	Not so efficient 😞 - Form of V is severely constrained by theoretical input (model Hamiltonian)

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To understand the **physics of reaction dynamics** behind formation, structure, etc. of hadron resonances, one needs:

- **Modeling reaction processes appropriately with a model Hamiltonian.**
(→ not a simple “pole + polynomial” parametrization, etc.)
- **Solving proper quantum scattering equation (LS eq.)** in which off-shell rescattering effects are also appropriately contained.

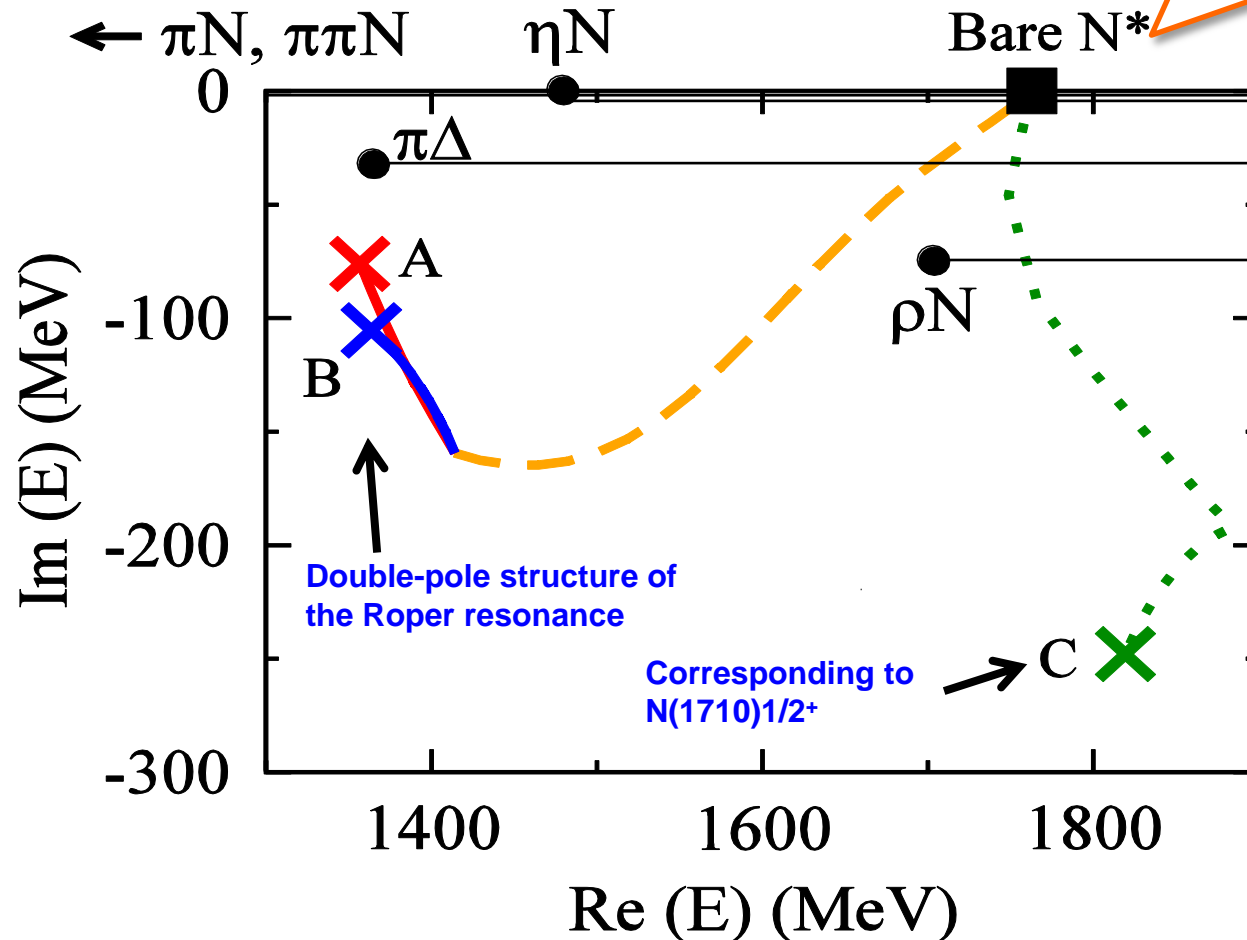


This can be achieved **only by using the dynamical-model approach !!**

Dynamical origin of $P_{11} N^*$ resonances

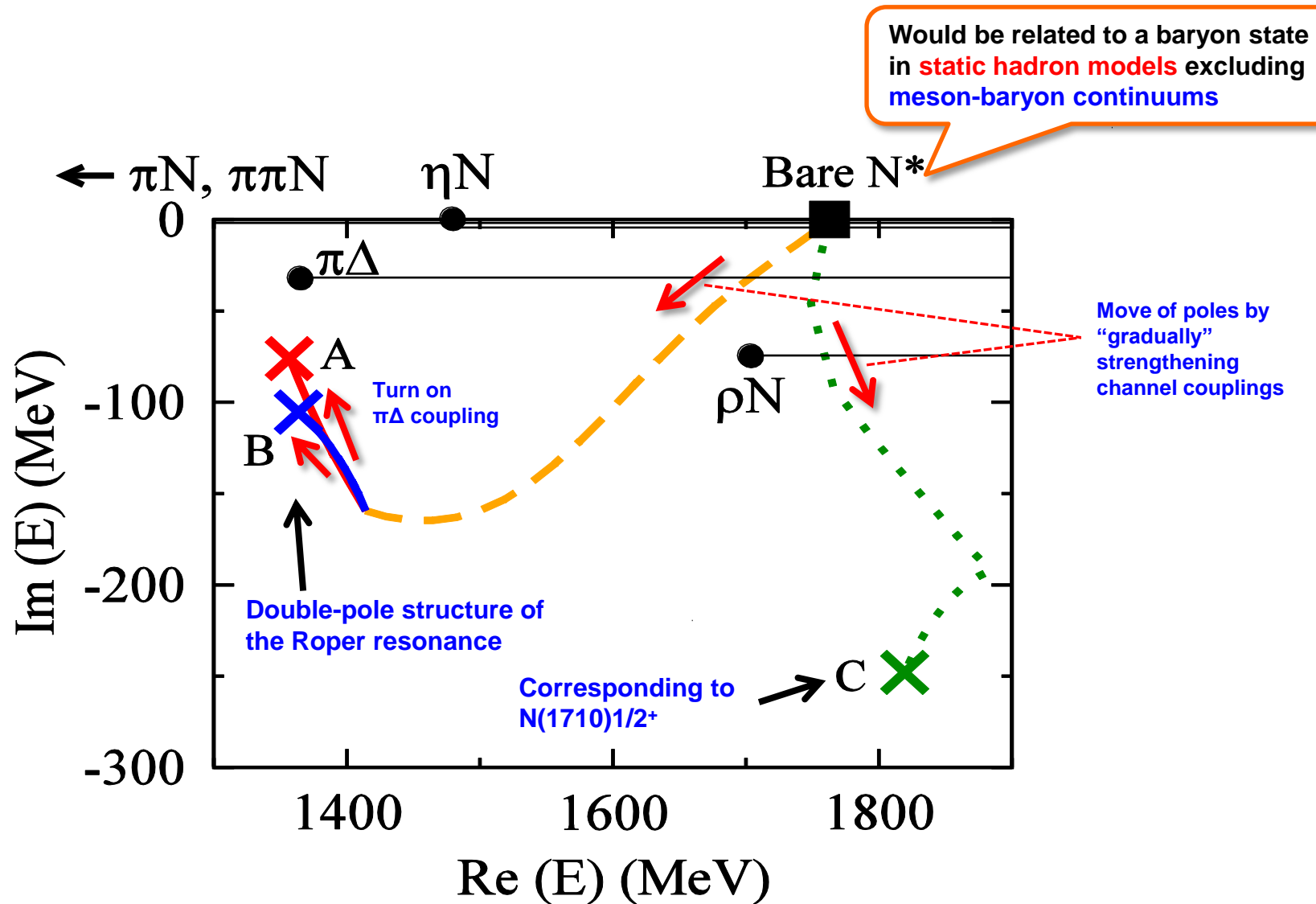
(nontrivial feature of multichannel reaction dynamics)

Would be related to a baryon state in **static hadron models** excluding meson-baryon continuums



Dynamical origin of $P_{11} N^*$ resonances

(nontrivial feature of multichannel reaction dynamics)



PART II

Recent results from ANL-Osaka DCC analysis

Dynamical coupled-channels (DCC) model for meson production reactions

For details see Matsuyama, Sato, Lee, Phys. Rep. 439(2007)193
 HK, Nakamura, Lee, Sato, PRC(2013)035209

- ✓ Partial-wave (LSJ) amplitudes of $a \rightarrow b$ reaction:

$$T_{a,b}^{(LSJ)}(p_a, p_b; E) = \underbrace{V_{a,b}^{(LSJ)}(p_a, p_b; E)}_{\text{coupled-channels effect}} + \sum_c \underbrace{\int_0^\infty q^2 dq V_{a,c}^{(LSJ)}(p_a, q; E) G_c(q; E) T_{c,b}^{(LSJ)}(q, p_b; E)}_{\text{off-shell effect}}$$

- ✓ Reaction channels:

$$a, b, c = (\gamma^{(*)}N, \pi N, \eta N, \underbrace{\pi\Delta, \sigma N, \rho N}_{\pi\pi N}, K\Lambda, K\Sigma, \omega N \dots)$$

- ✓ Transition Potentials:

$$V_{a,b} = \underbrace{v_{a,b}}_{\text{Exchange potentials}} + \underbrace{Z_{a,b}}_{\text{Z-diagrams}} + \sum_{N^*} \frac{\Gamma_{N^*,a}^\dagger \Gamma_{N^*,b}}{E - M_{N^*}} \underbrace{\text{bare } N^* \text{ states}}$$

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✓ Meson-Baryon Green functions G_{MB}

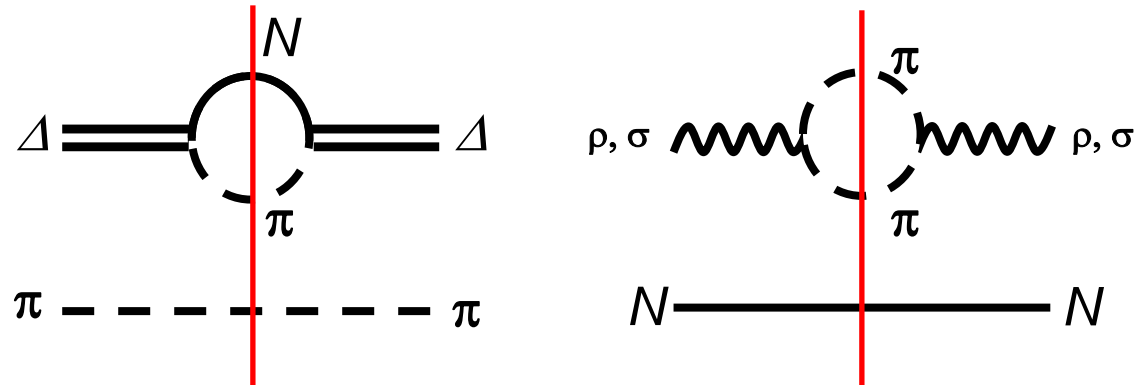
$MB = \pi N, \eta N, K\Lambda, K\Sigma, \omega N$

Stable channels



$MB = \pi\Delta, \rho N, \sigma N$

Quasi 2-body channels



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$\pi\pi N$

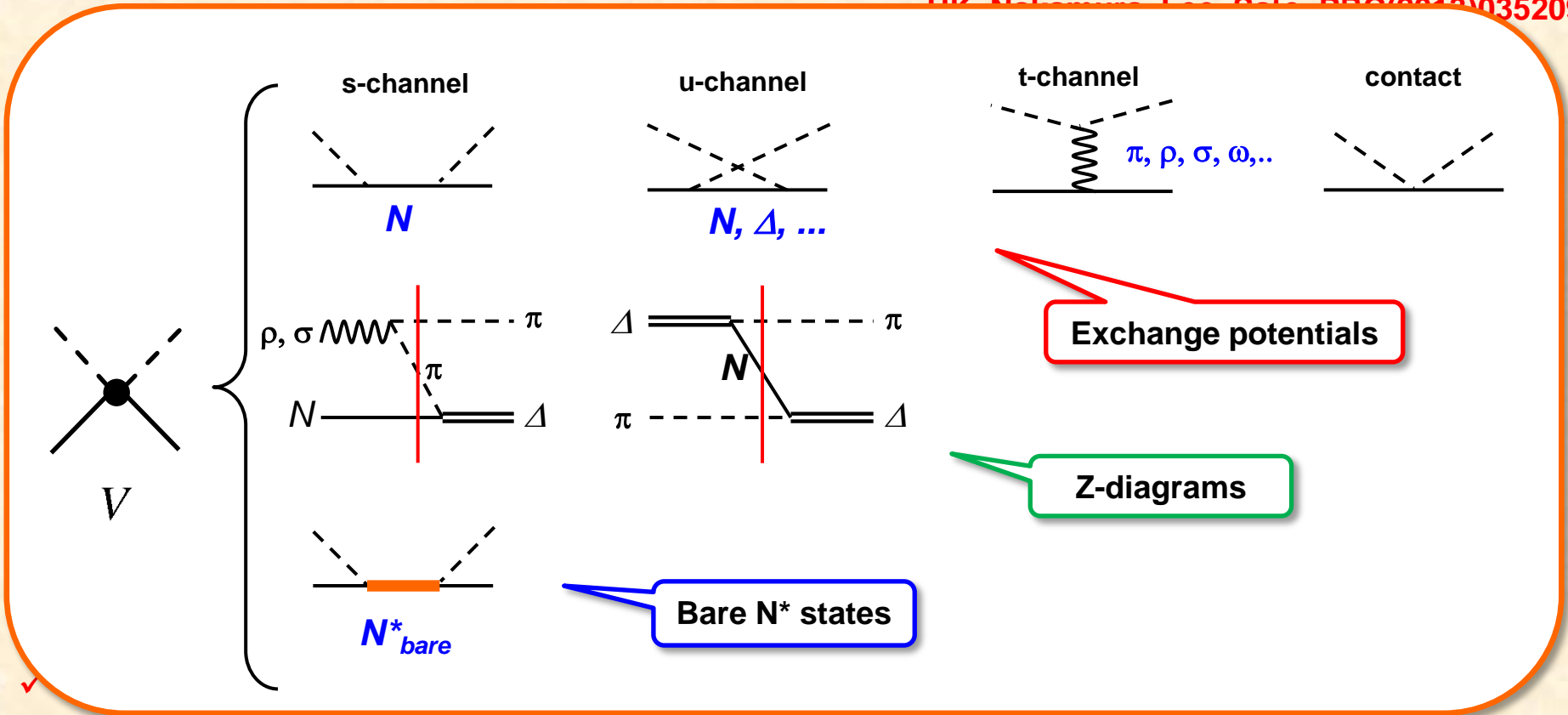
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$$V_{a,b} = v_{a,b} + Z_{a,b} + \sum_{N^*} \frac{\Gamma_{N^*,a}^\dagger \Gamma_{N^*,b}}{E - M_{N^*}}$$

Exchange potentials
 Z-diagrams
 bare N^* states

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HK, Nakamura, Lee, Sato, PRC(2013)035209

- ✓ Partial-wave (LSJ) amplitudes of $a \rightarrow b$ reaction:

$$T_{a,b}^{(LSJ)}(p_a, p_b; E) = \underbrace{V_{a,b}^{(LSJ)}(p_a, p_b; E)}_{\text{coupled-channels effect}} + \sum_c \underbrace{\int_0^\infty q^2 dq V_{a,c}^{(LSJ)}(p_a, q; E) G_c(q; E) T_{c,b}^{(LSJ)}(q, p_b; E)}_{\text{off-shell effect}}$$

- ✓ Reaction channels:

$$a, b, c = (\gamma^{(*)} N)$$

Would be related with hadron states of the **static hadron models** (quark models, DSE, etc.) **excluding meson-baryon continuums.**

- ✓ Transition Potentials:

$$V_{a,b} = \underbrace{v_{a,b}}_{\text{Exchange potentials}} + \underbrace{Z_{a,b}}_{\text{Z-diagrams}} + \sum_{N^*} \frac{\Gamma_{N^*,a}^\dagger \Gamma_{N^*,b}}{E - M_{N^*}} \underbrace{\text{bare } N^* \text{ states}}_{\text{bare } N^* \text{ states}}$$

Dynamical coupled-channels (DCC) model for meson production reactions

For details see Matsuyama, Sato, Lee, Phys. Rep. 439(2007)193
 HK, Nakamura, Lee, Sato, PRC(2013)035209

- ✓ Partial-wave (LSJ) amplitudes of $a \rightarrow b$ reaction:

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- ✓ Reaction channels:

$$a, b, c = (\gamma^{(*)}N, \pi N, \eta N, \underbrace{\pi\Delta, \sigma N, \rho N}_{\pi\pi N}, K\Lambda, K\Sigma, \omega N \dots)$$

- ✓ Transition Potentials:

$$V_{a,b} = \underbrace{v_{a,b}}_{\text{Exchange potentials}} + \underbrace{Z_{a,b}}_{\text{Z-diagrams}} + \sum_{N^*} \frac{\Gamma_{N^*,a}^\dagger \Gamma_{N^*,b}}{E - M_{N^*}} \underbrace{\text{bare } N^* \text{ states}}$$

Strategy for N^* and Δ^* spectroscopy

1) Construct a model by making χ^2 -fit of the world data of meson production reactions.

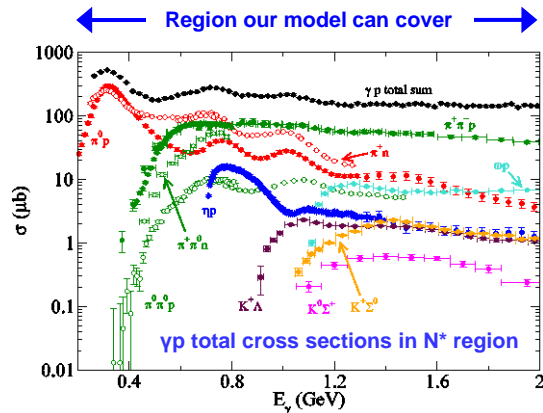
✓ Latest published model (8-channel):

HK, Nakamura, Lee, Sato, PRC88(2013)035209
[updated in PRC94(2016)015201]

Made simultaneous analysis of

- $\pi N \rightarrow \pi N$ (SAID amp) ($W < 2.3$ GeV)
- $\pi p \rightarrow \eta N, K\Lambda, K\Sigma$ ($W < 2.1$ GeV)
- $\gamma p \rightarrow \pi N, \eta N, K\Lambda, K\Sigma$ ($W < 2.1$ GeV)
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→ ~27,000 data points of both $d\sigma/d\Omega$ & spin-pol. obs.



➤ Use supercomputers to accomplish coupled-channels analyses:



Strategy for N^* and Δ^* spectroscopy

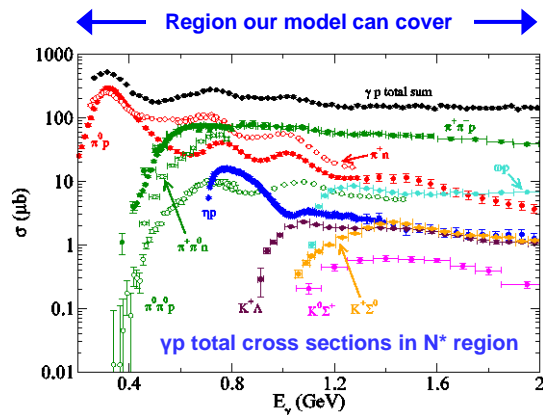
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 HK, Nakamura, Lee, Sato, PRC88(2013)035209
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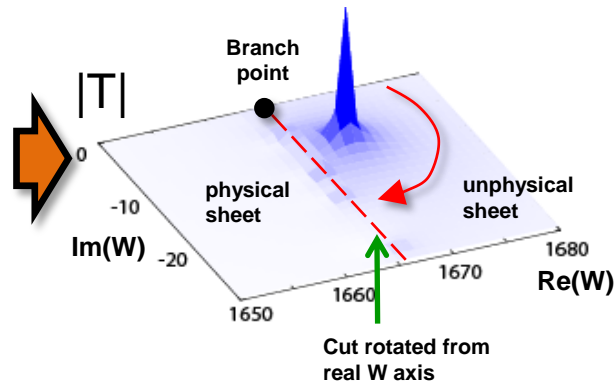
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2) Search **poles** of scattering amplitudes by **analytic continuation** to a **complex energy plane**.



Pole position → (complex) resonance mass

Residues → coupling strengths between resonance and meson-baryon channel

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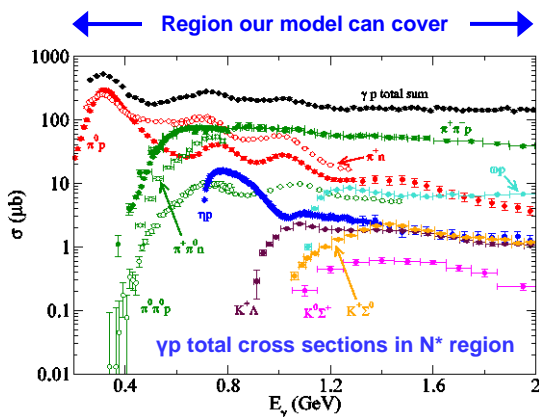
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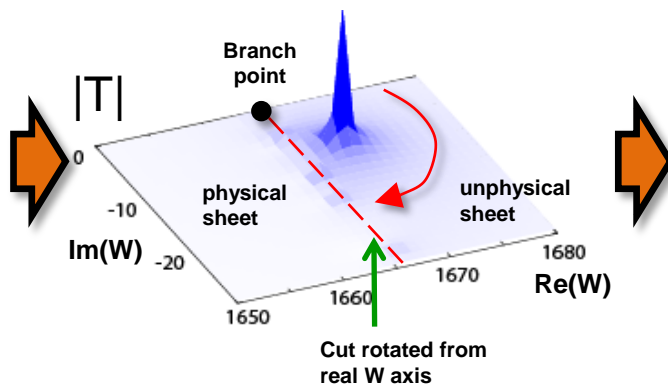
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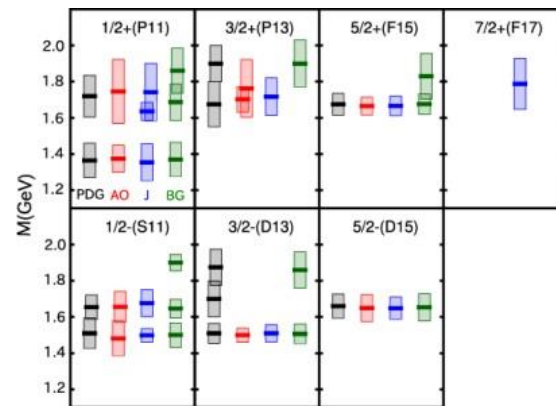


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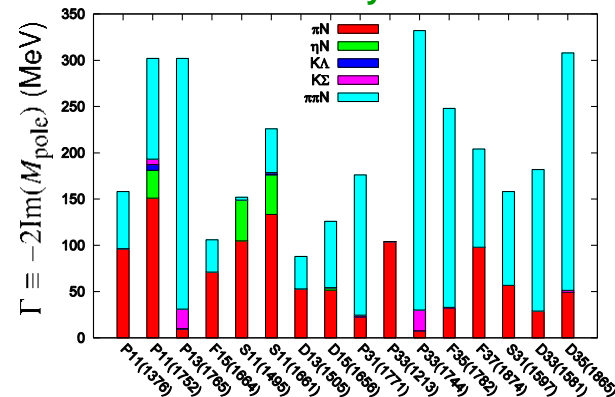
Residues → coupling strengths between resonance and meson-baryon channel

3) Extract resonance parameters **defined by poles**.

Mass spectrum



Partial decay widths

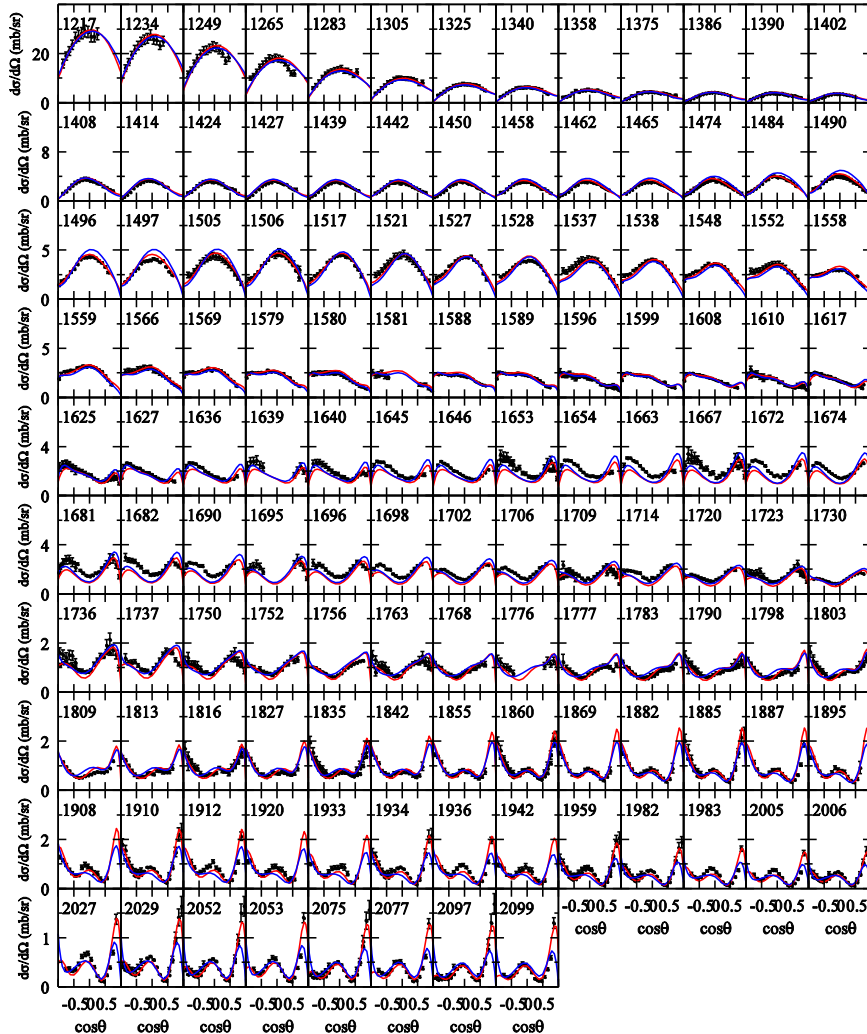


ANL-Osaka DCC analysis

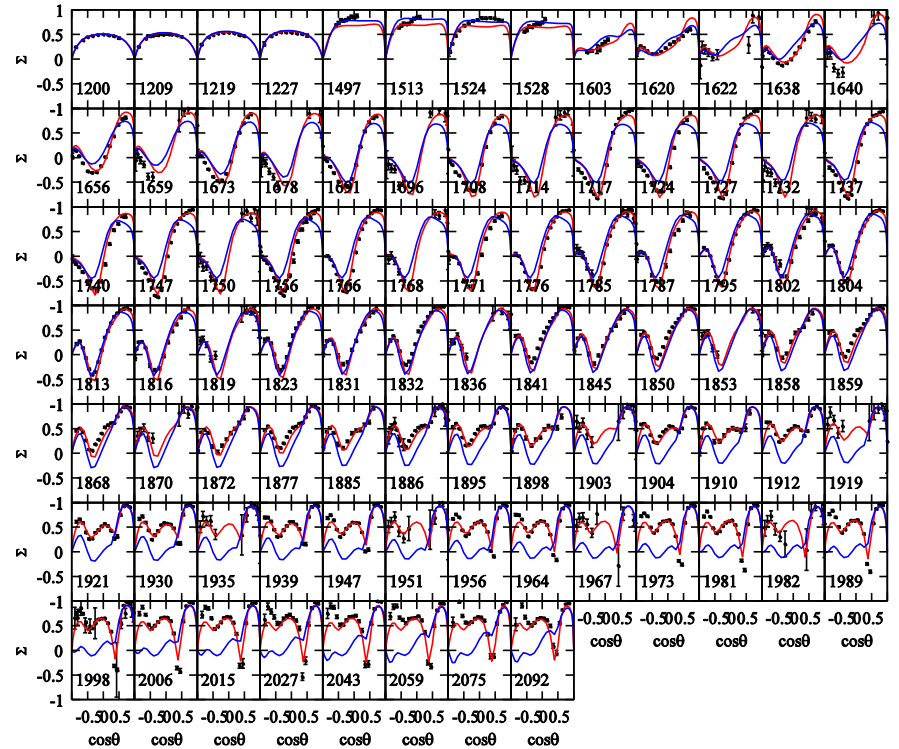
HK, Nakamura, Lee, Sato, PRC88(2013)035209; 94(2016)015201

$\gamma p \rightarrow \pi^0 p$

$d\sigma/d\Omega$ for $W < 2.1$ GeV



Σ for $W < 2.1$ GeV



Red: Updated model [PRC94(2016)015201]

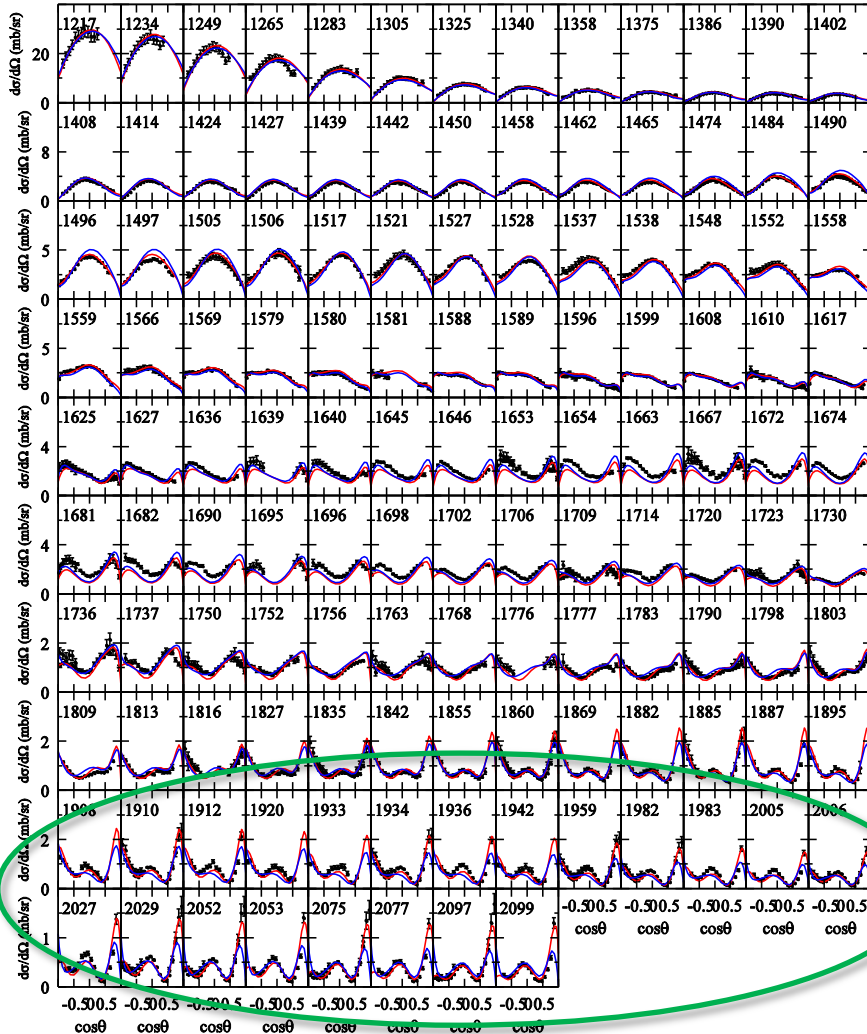
Blue: Original model [PRC88(2013)035209]

ANL-Osaka DCC analysis

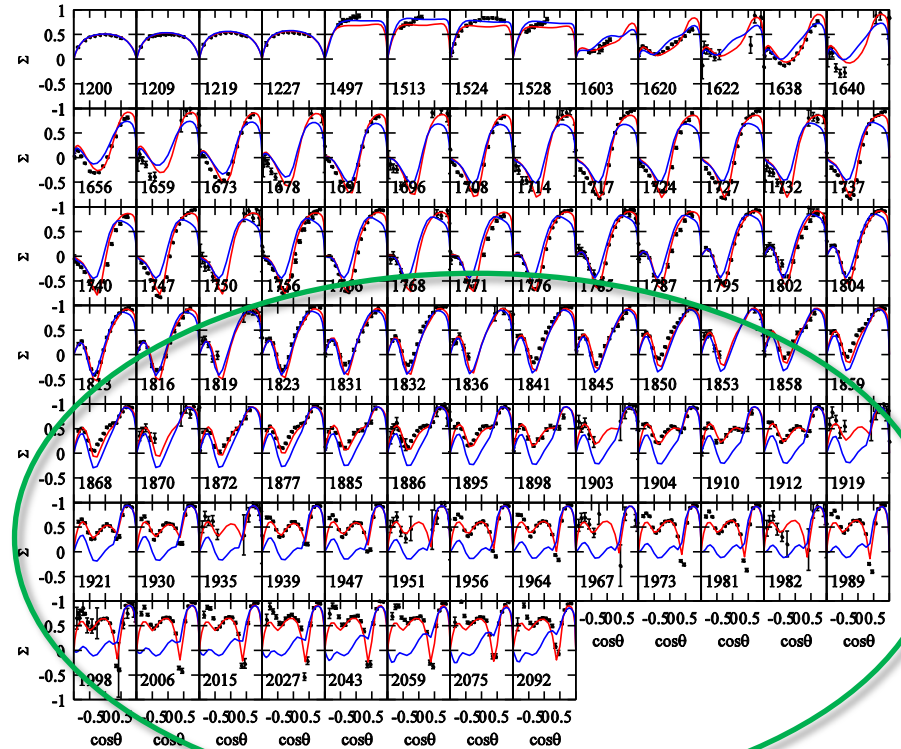
HK, Nakamura, Lee, Sato, PRC88(2013)035209; 94(2016)015201

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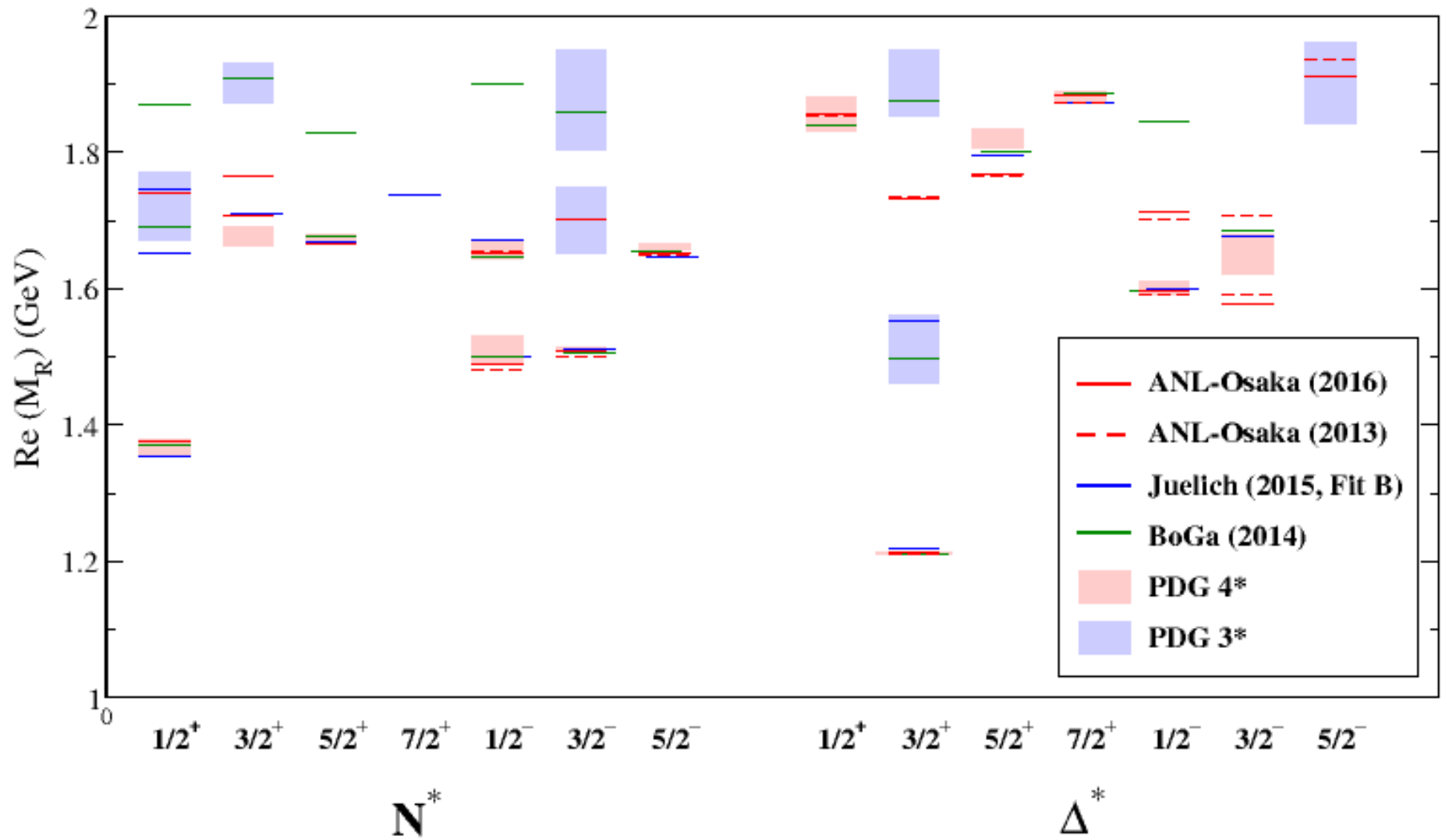
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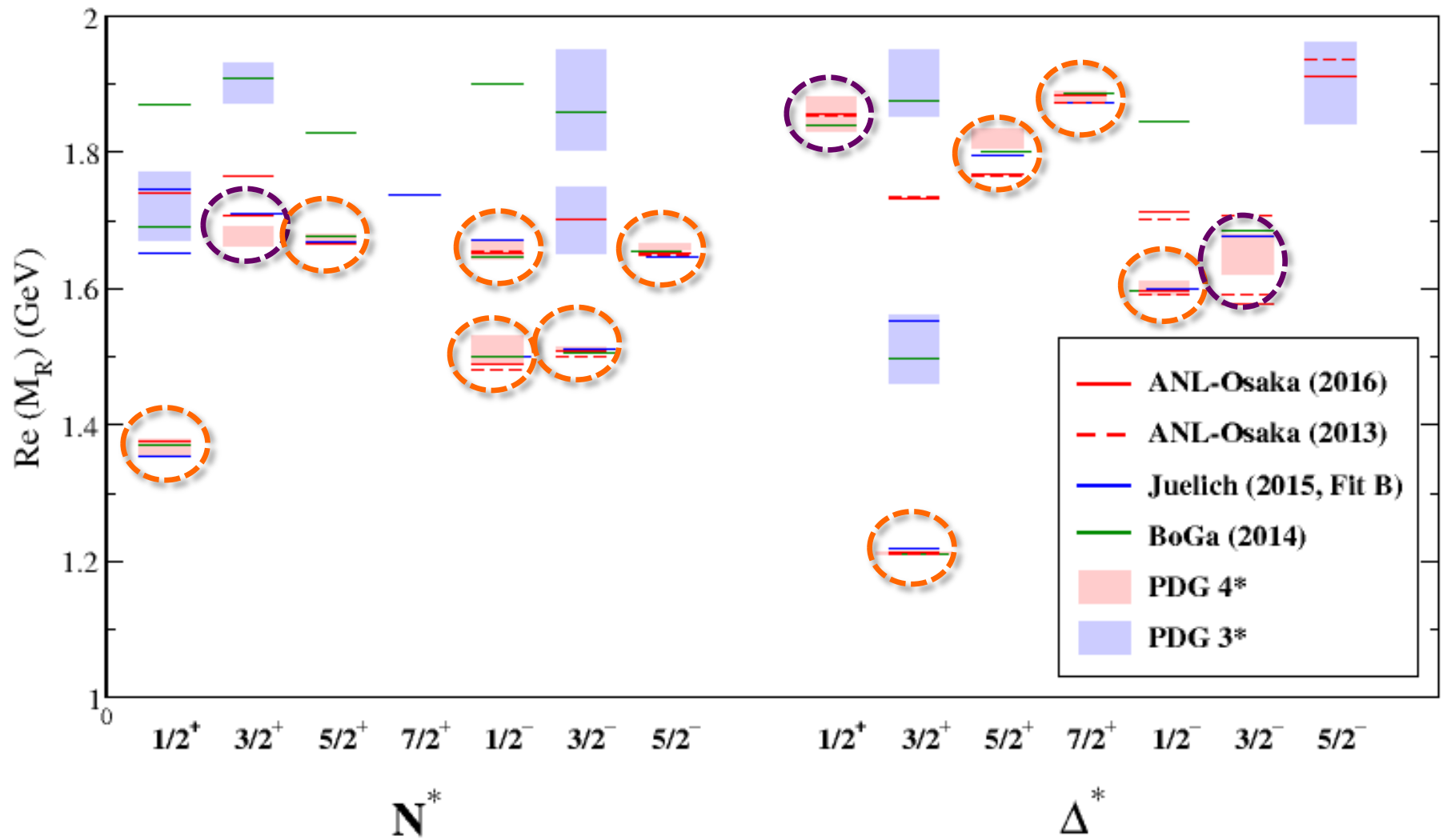
Red: Updated model [PRC94(2016)015201]

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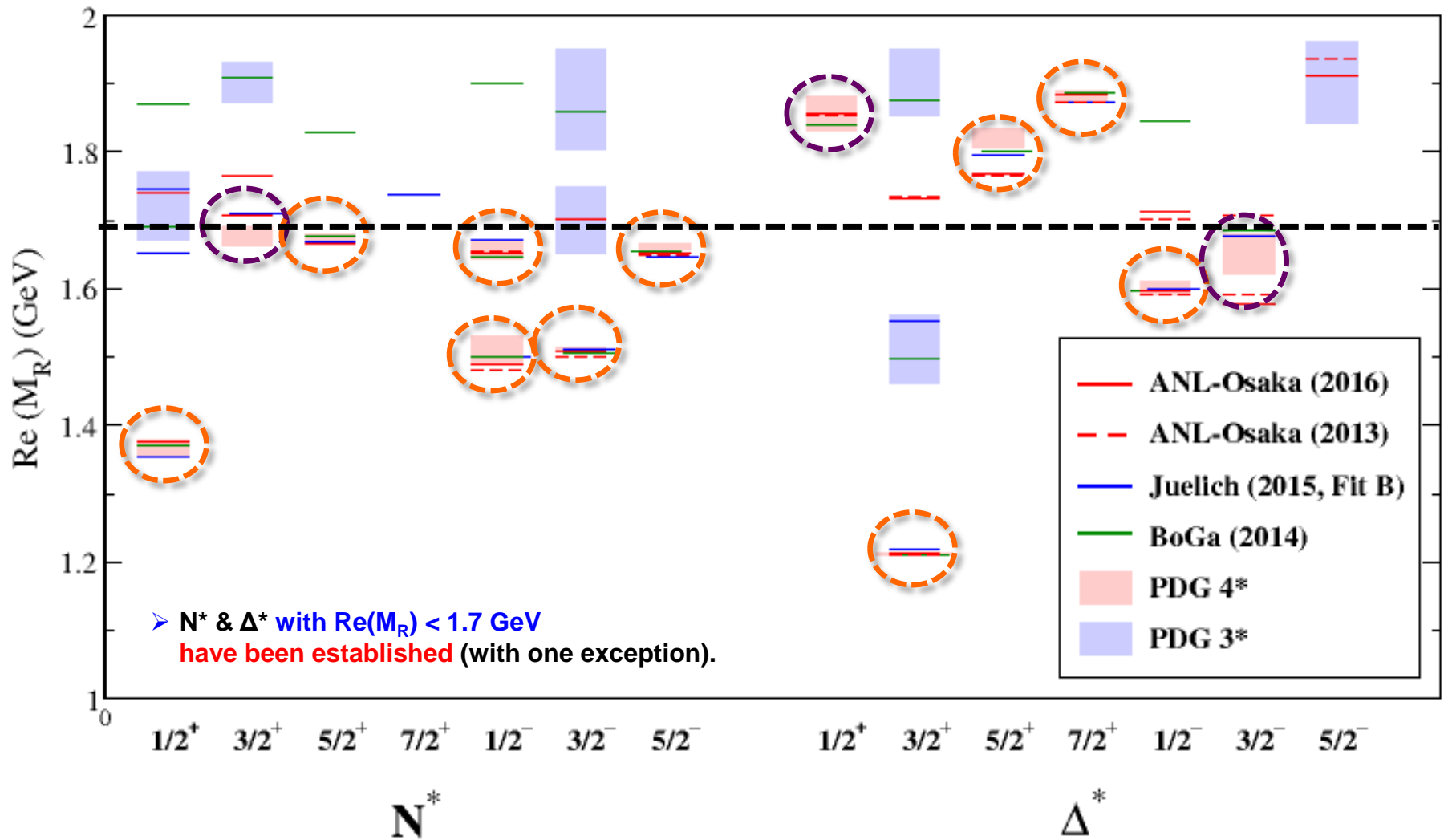
Extracted N^* & Δ^* mass spectrum (pole mass M_R)



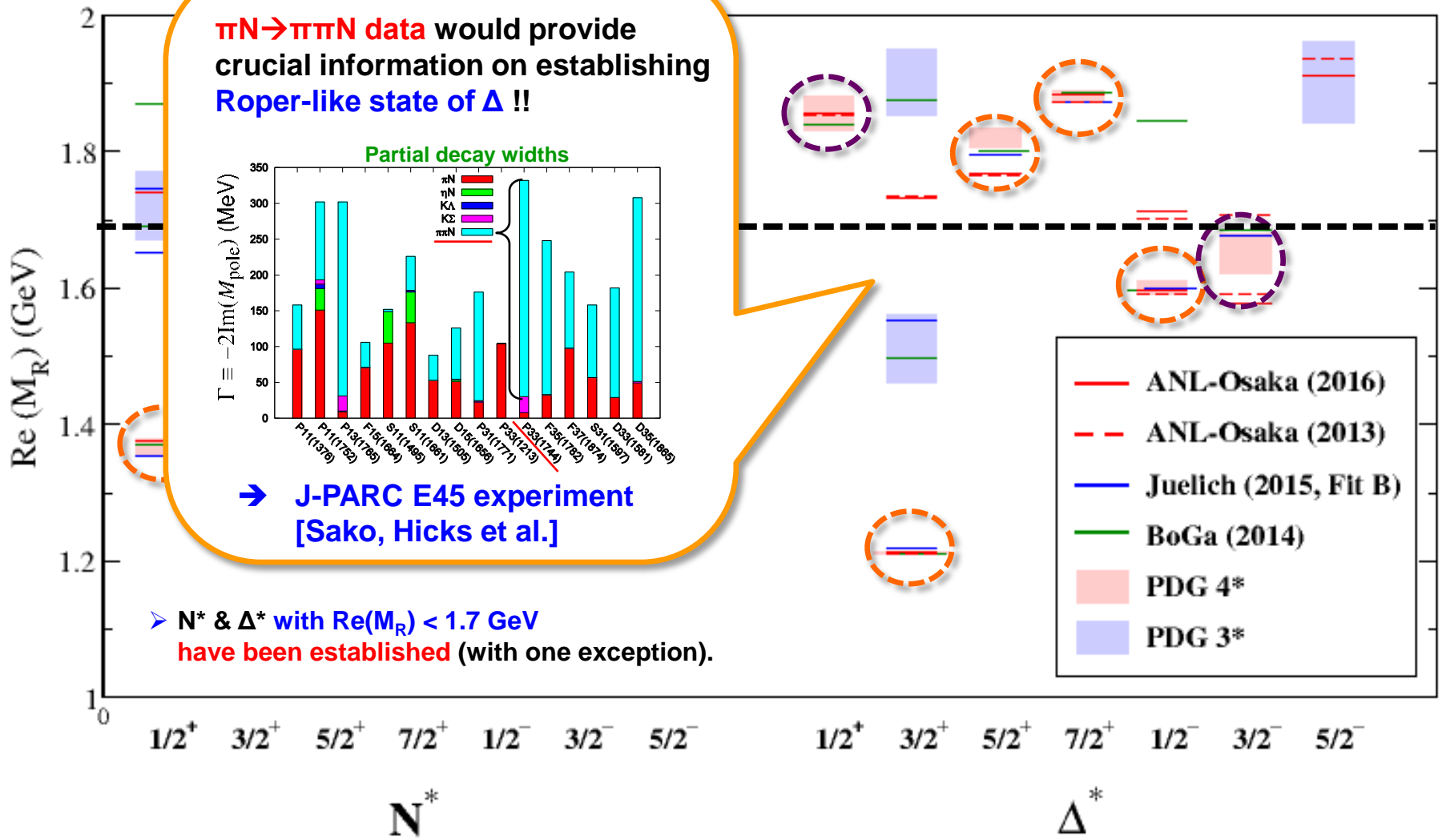
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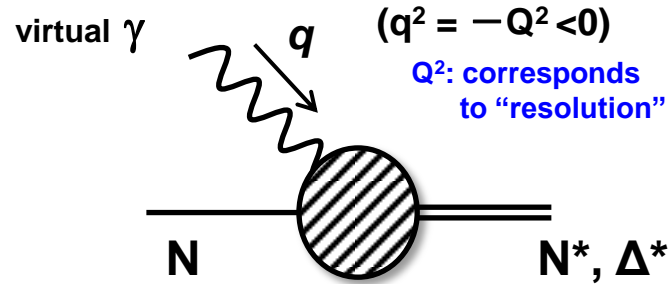
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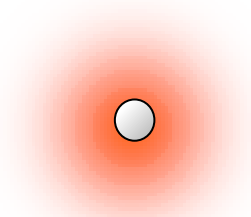


Electromagnetic transition form factors: quantitative understanding of N^* & Δ^* structure



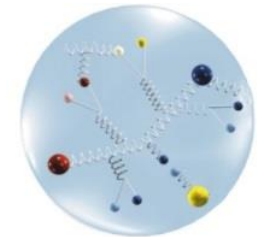
N-N* e.m. transition form factor

“dressed”-quark core
obscured by dense meson clouds



**How effective d.o.f.s of baryon
constituents change with Q^2 ??**

“Partons”

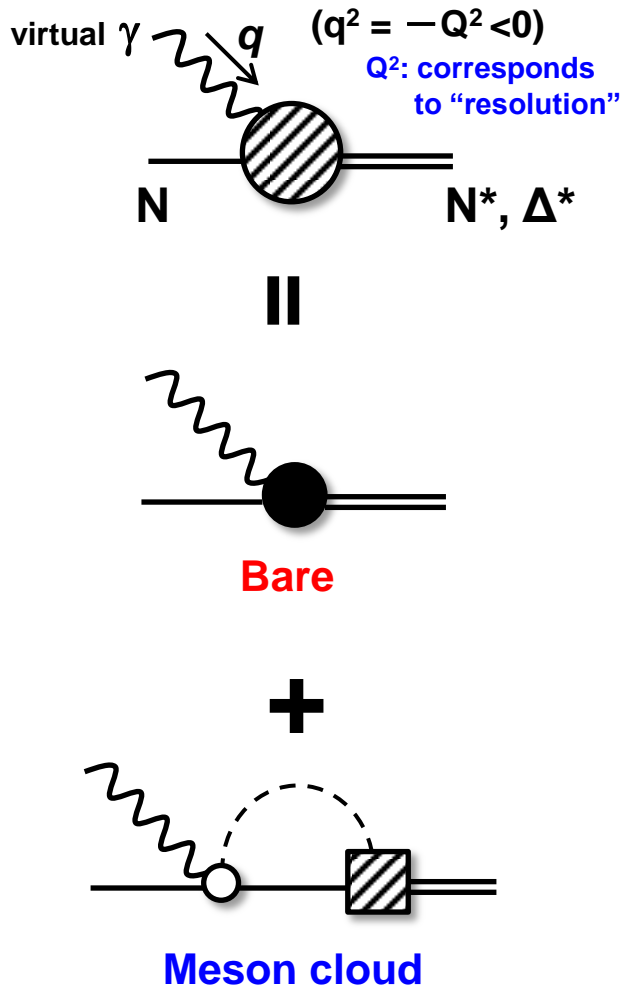


Q^2 : small
(low “resolution”)

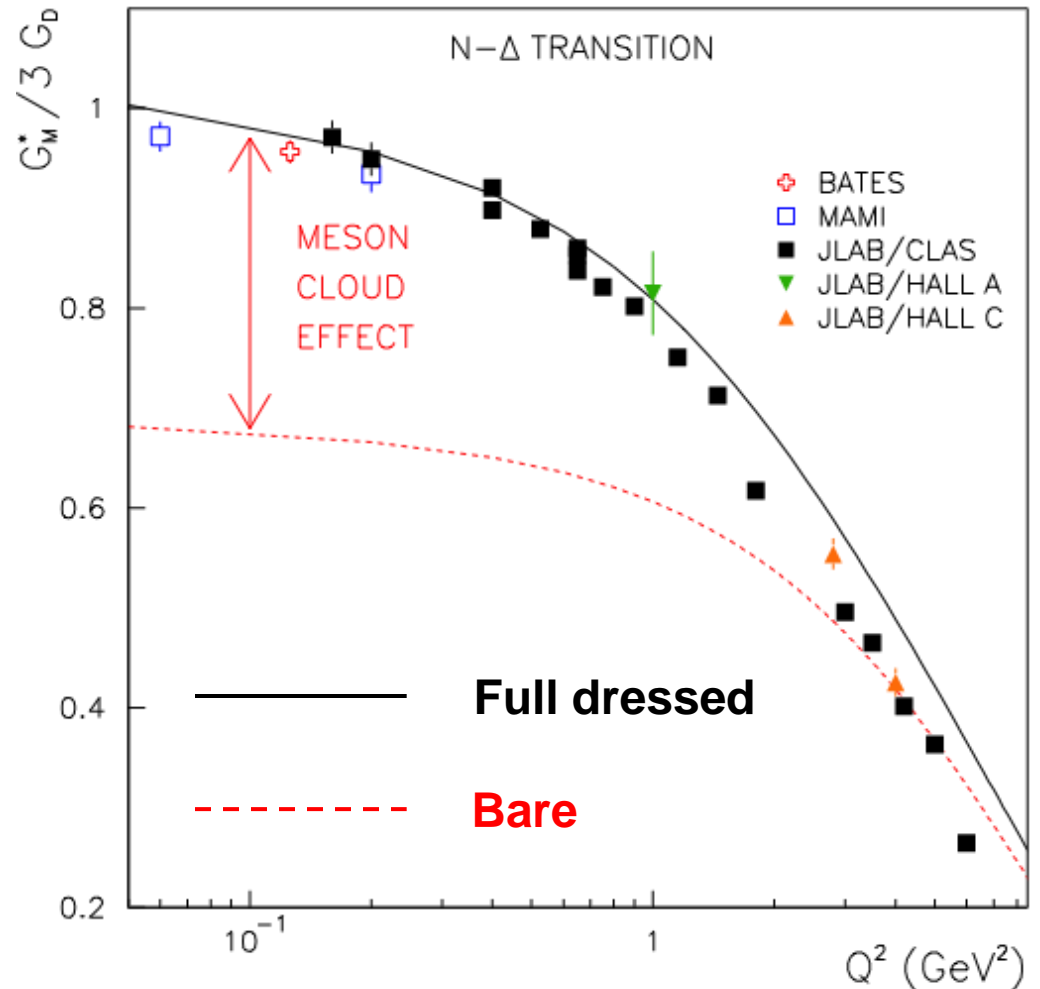
Q^2 : large
(high “resolution”)

Role of reaction dynamics in form factors: Meson-cloud effect

N-N* e.m. transition form factor



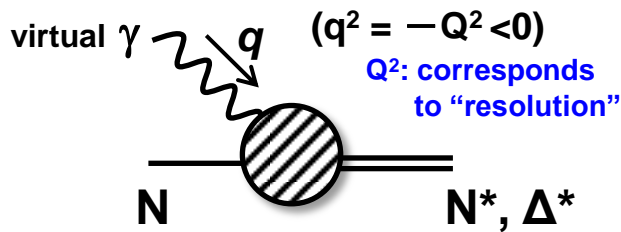
Re[$G_M(Q^2)$] for $\gamma N \rightarrow \Delta(1232)$ M1 transition



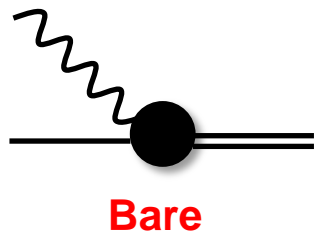
Julia-Diaz, et al., PRC75 015205 (2007)

Role of reaction dynamics in form factors: Meson-cloud effect

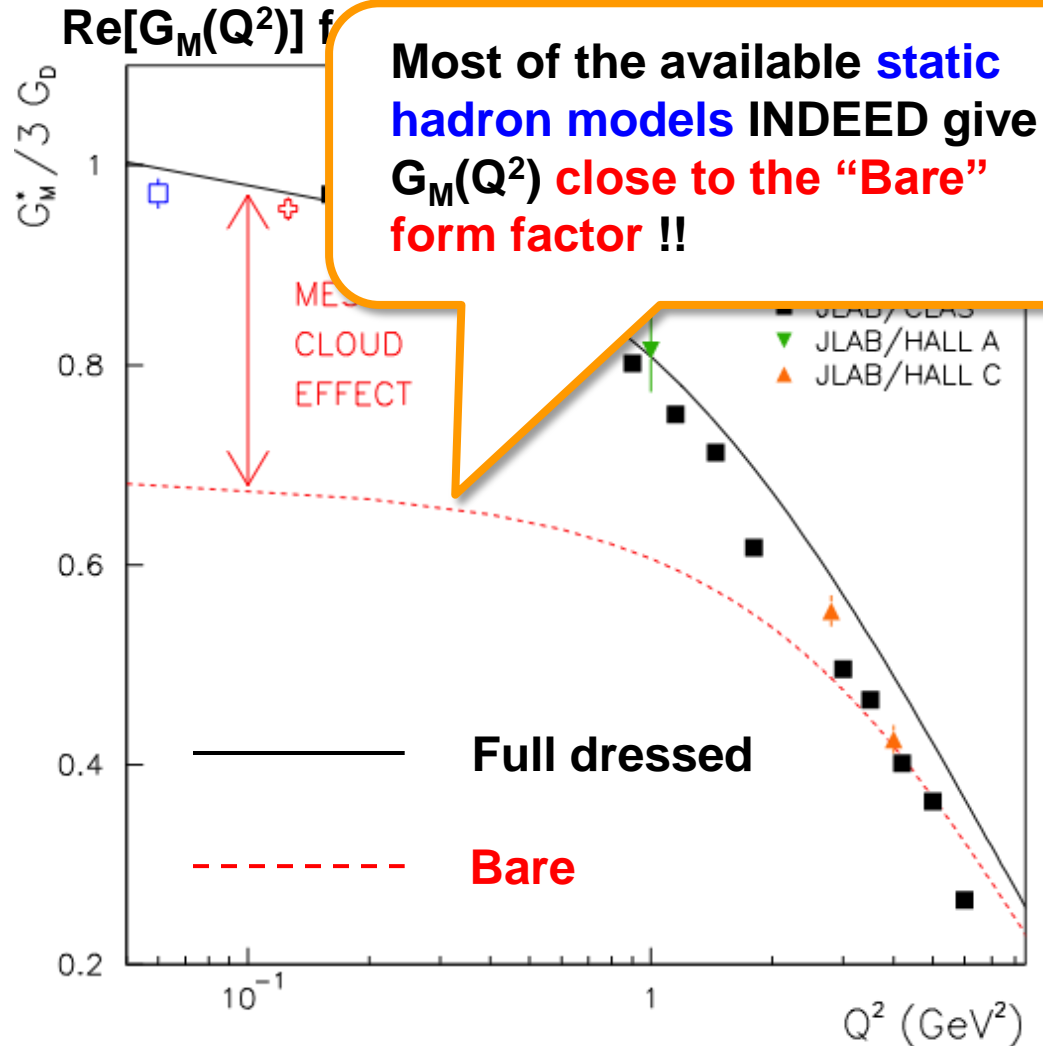
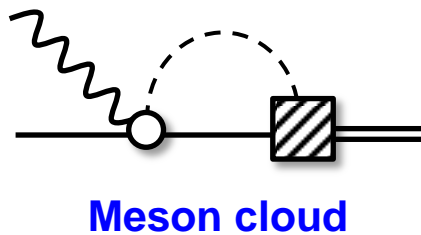
N-N* e.m. transition form factor



||



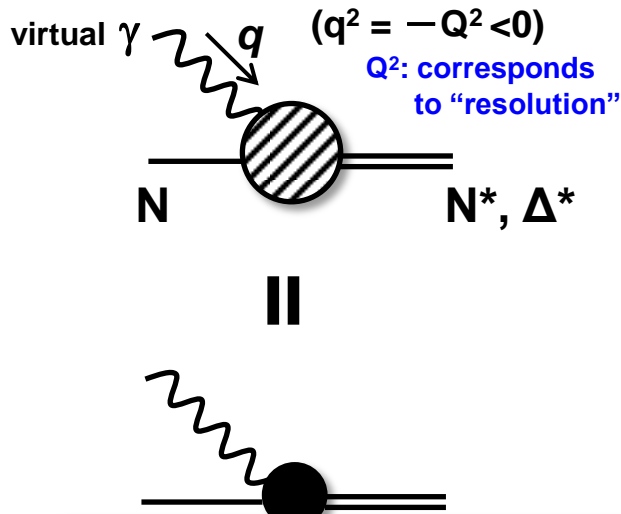
+



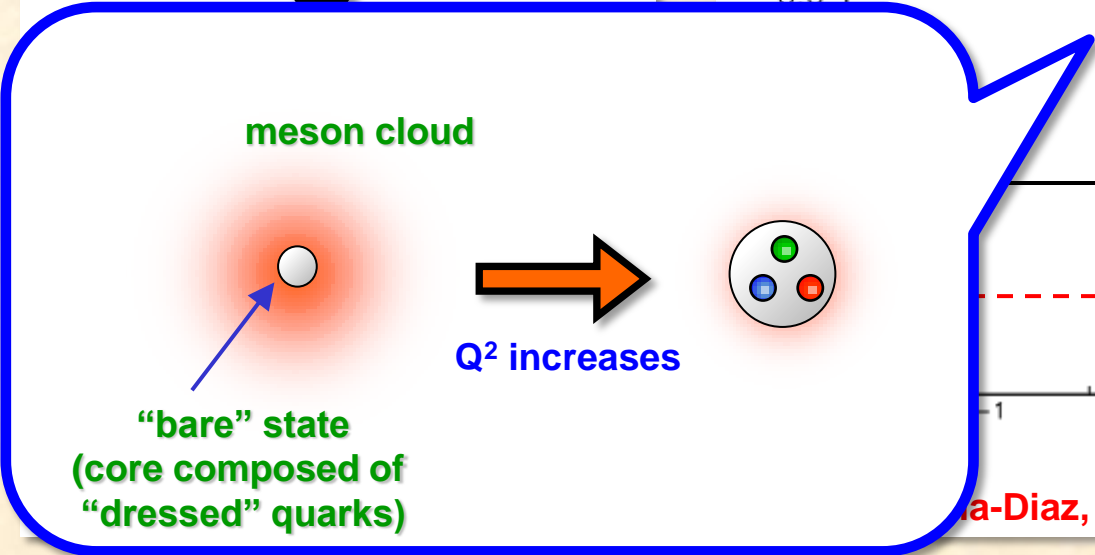
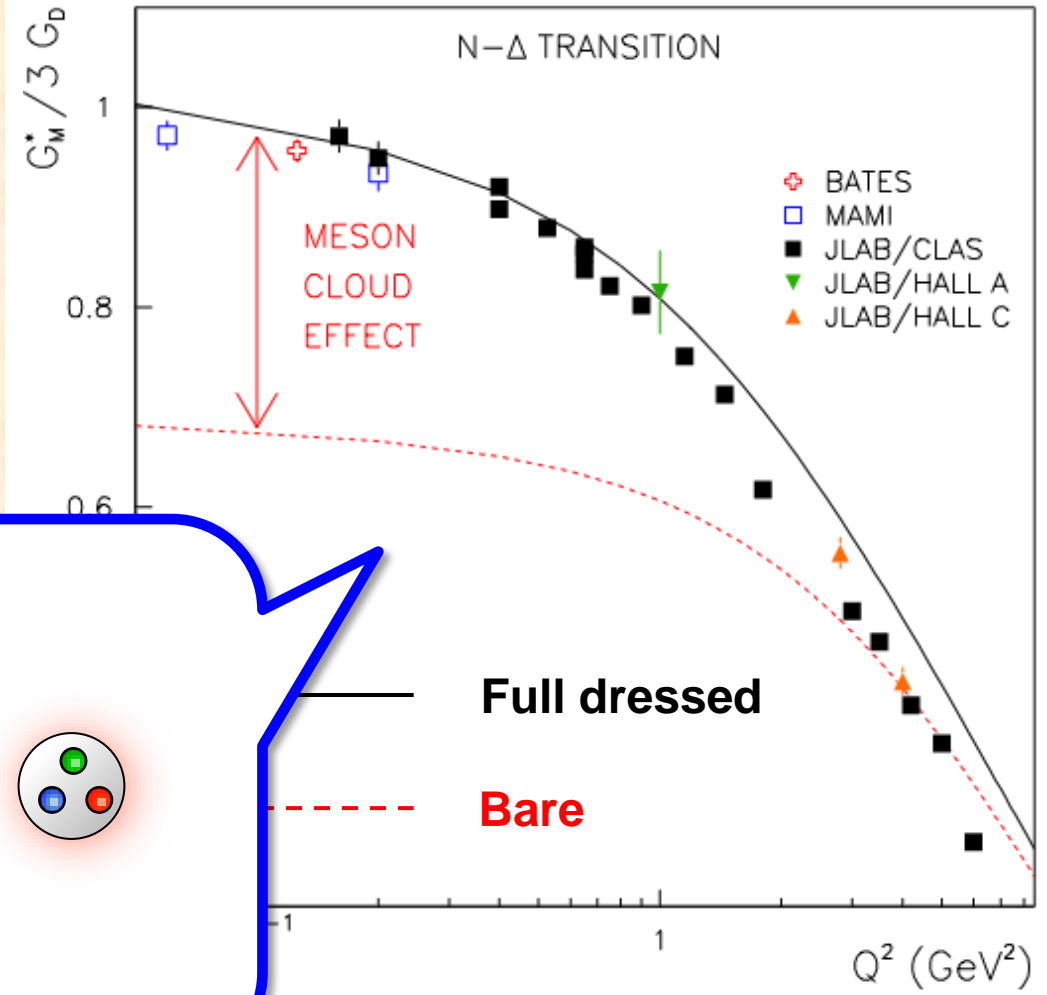
Julia-Diaz, et al., PRC75 015205 (2007)

Role of reaction dynamics in form factors: Meson-cloud effect

N-N* e.m. transition form factor



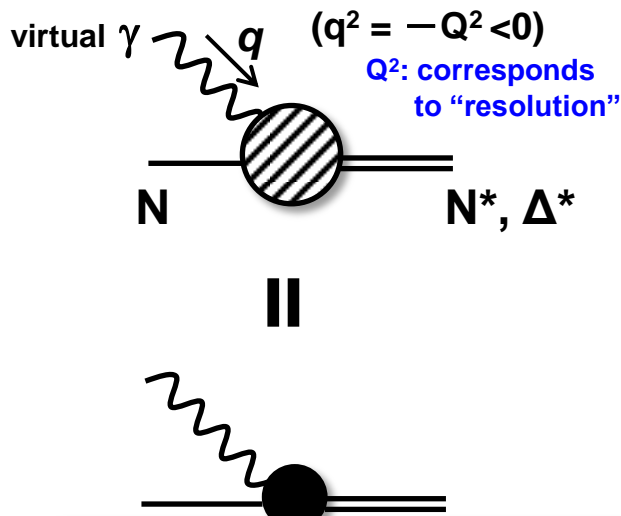
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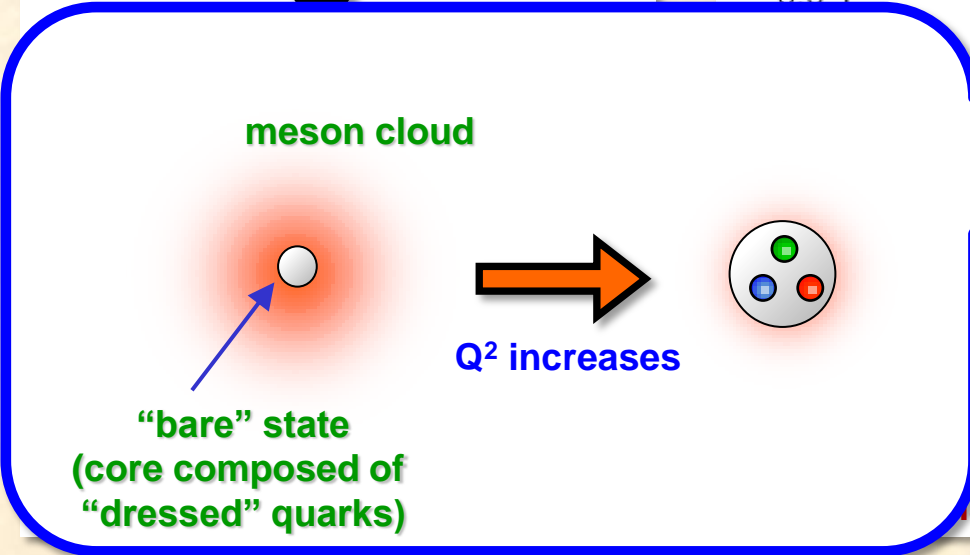
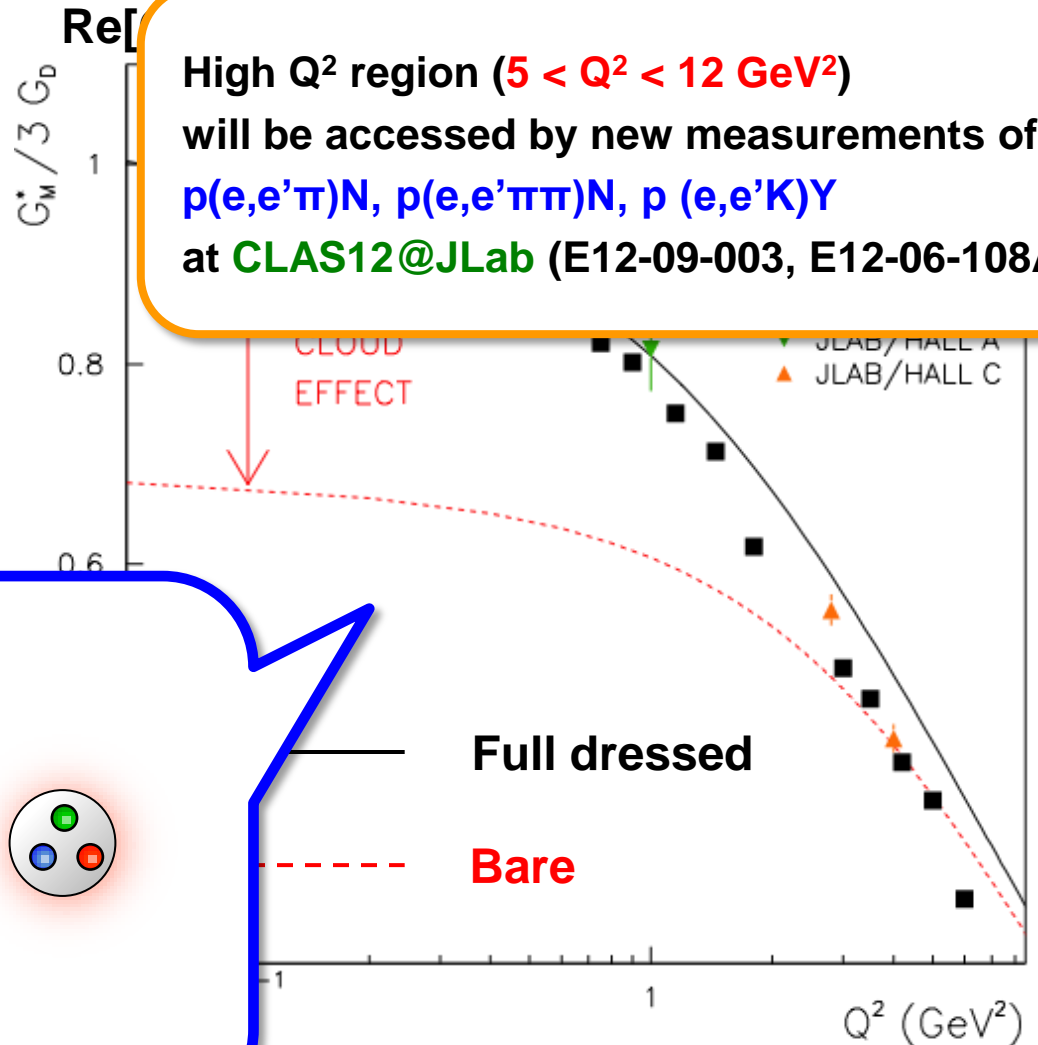
a-Diaz, et al., PRC75 015205 (2007)

Role of reaction dynamics in form factors: Meson-cloud effect

N-N* e.m. transition form factor



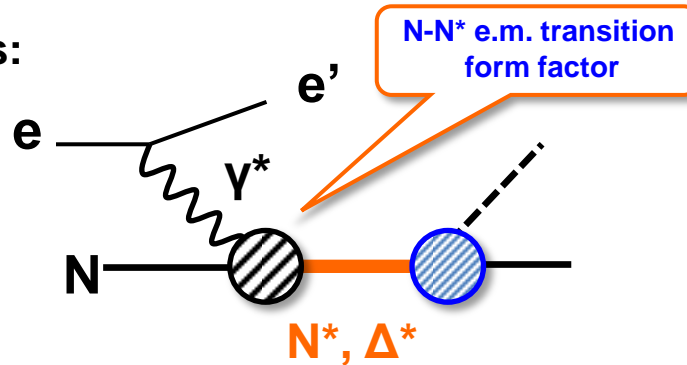
High Q^2 region ($5 < Q^2 < 12 \text{ GeV}^2$)
 will be accessed by new measurements of
 $p(e,e'\pi)N$, $p(e,e'\pi\pi)N$, $p(e,e'K)Y$
 at **CLAS12@JLab** (E12-09-003, E12-06-108A)



da-Diaz, et al., PRC75 015205 (2007)

Analysis of electroproduction reactions: Determining N-N* e.m. transition form factors

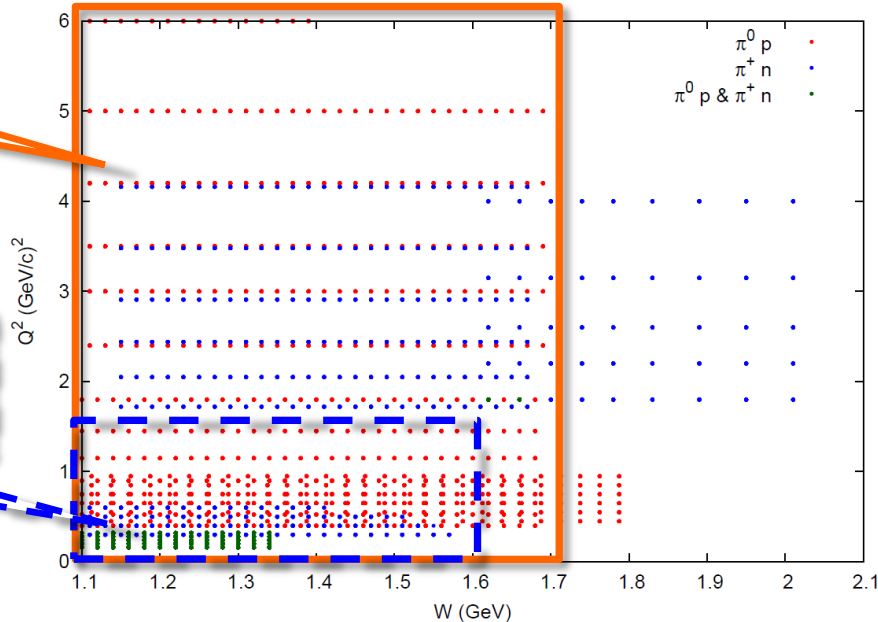
➤ Meson **electro**productions:



Database for 1π electroproduction@CLAS6
($Q^2 < 6 \text{ GeV}^2$)

(W, Q^2) region in
the current analysis
($Q^2 < 6 \text{ GeV}^2, W < 1.7 \text{ GeV}$)

(W, Q^2) region in the early
analysis:
Julia-Diaz, HK, Lee, Matsuyama,
Sato, Suzuki, PRC80(2009)025207



+ $K^+\Lambda, K^+\Sigma^0,$
 $\pi\pi N$
electro-
production
data

Analysis of electroproduction reactions: Determining N-N* e.m. transition form factors

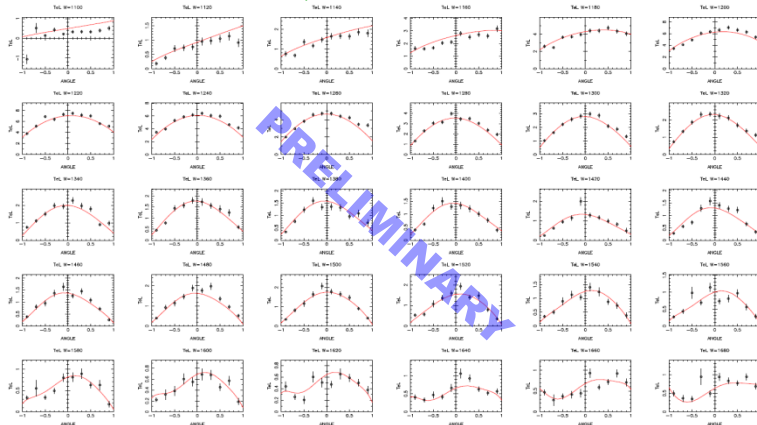
$$\frac{d\sigma^5}{dE_{e'}d\Omega_{e'}d\Omega_{\pi}^*} = \Gamma_{\gamma} \left[\sigma_T + \epsilon\sigma_L + \sqrt{2\epsilon(1+\epsilon)}\sigma_{LT} \cos\phi_{\pi}^* + \epsilon\sigma_{TT} \cos 2\phi_{\pi}^* + h_e \sqrt{2\epsilon(1-\epsilon)}\sigma_{LT'} \sin\phi_{\pi}^* \right].$$

$\sigma_T + \epsilon\sigma_L$ for $ep \rightarrow e\pi^0 p$

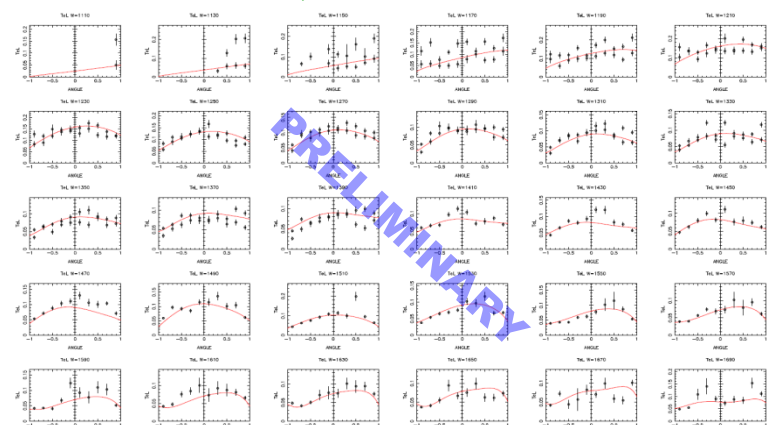
Data for structure functions are obtained with the help of K. Joo and L. C. Smith.

$$\sigma_{\alpha} = \sigma_{\alpha}(W, Q^2, \cos\theta_{\pi}^*)$$

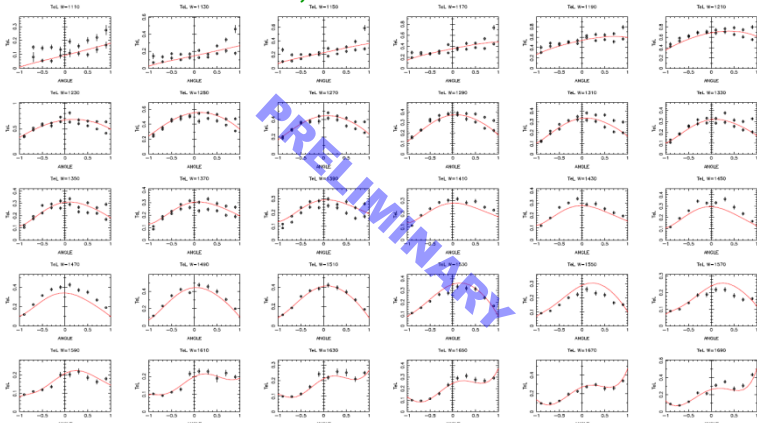
$Q^2 = 1.15 \text{ GeV}^2, 1.10 < W < 1.69 \text{ GeV}$



$Q^2 = 5.0 \text{ GeV}^2, 1.11 < W < 1.69 \text{ GeV}$

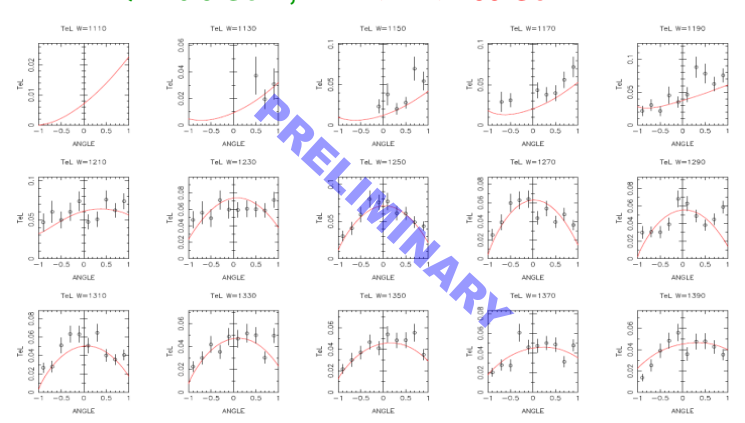


$Q^2 = 3.0 \text{ GeV}^2, 1.11 < W < 1.69 \text{ GeV}$



cos θ

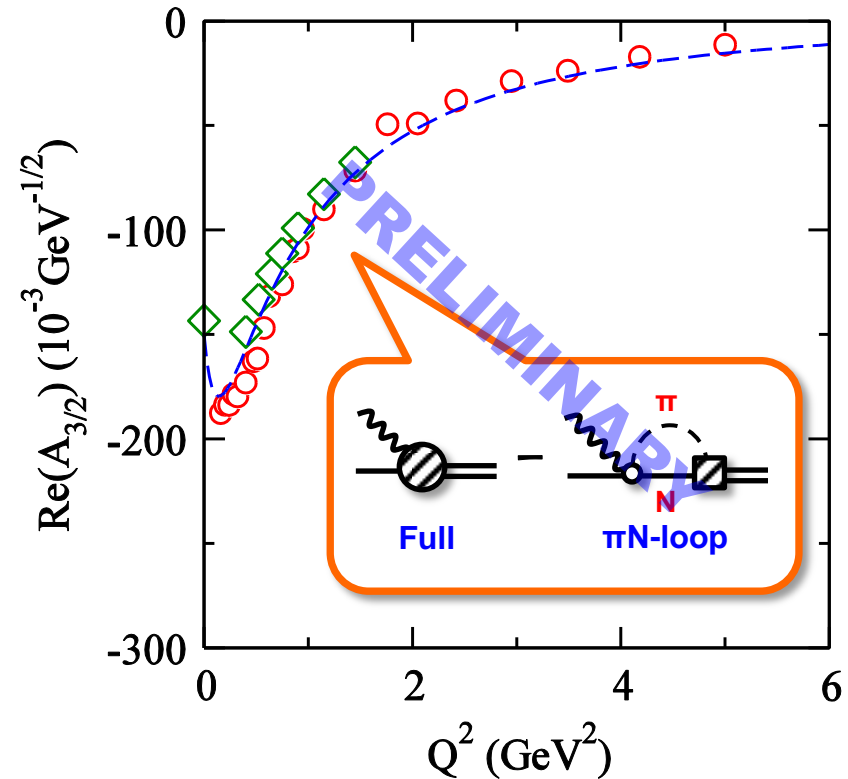
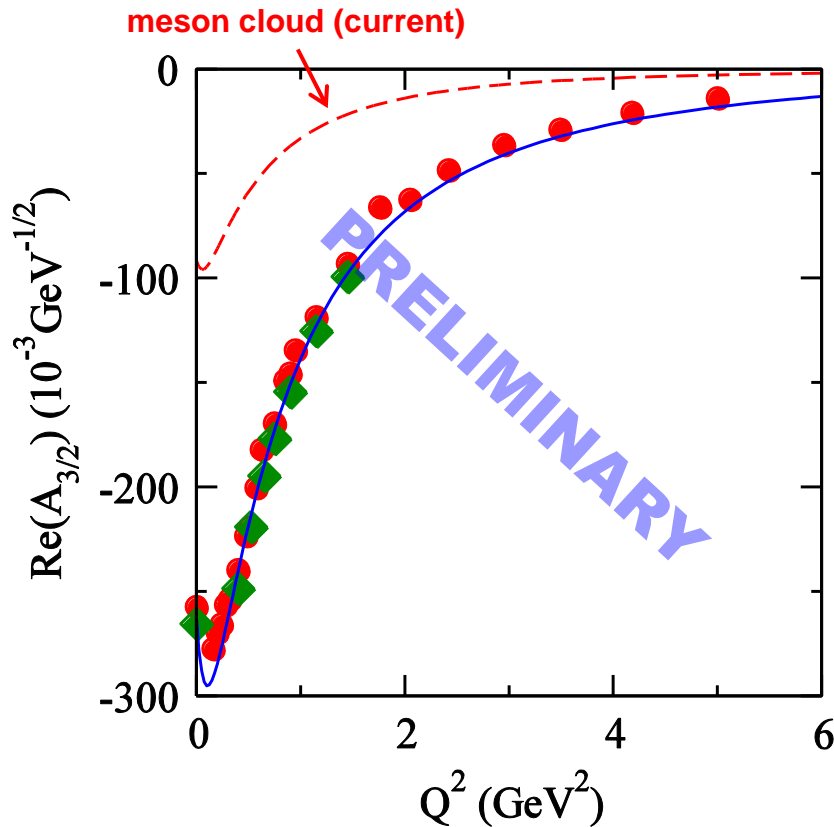
$Q^2 = 6.0 \text{ GeV}^2, 1.11 < W < 1.39 \text{ GeV}$



cos θ

Extracted e.m. transition form factors

- ✓ $N \rightarrow$ "1st P33($J^P=3/2^+$) Δ " transition form factor $A_{3/2}$
[evaluated at Δ pole mass: $M_R = 1210 - i 50$ MeV]



- Current** = $\pi N, \pi\pi N, \eta N, K\Lambda, K\Sigma$; 2 bare states in P33
- JLMS** = $\pi N, \pi\pi N, \eta N$; 2 bare states in P33
[PRC80(2009)025207; 82(2010)045206]
- Sato-Lee** = πN ; 1 bare state in P33
[PRC63(2001)055201; 75(2007)015205]

Summary

- ✓ **N^* & Δ^* spectroscopy as physics of broad & overlapping resonances**
 - **Cooperative efforts** between **experiments** and **theoretical analyses with coupled-channels framework** are indispensable to establishing the spectrum.
 - **Reaction dynamics** is a crucial part of understanding the spectrum, dynamical origin, and structure, ... of N^* & Δ^* .
 - **Dynamical coupled-channels approach** is a suitable one to study the role of reaction dynamics.
 - ➔ Multichannel reaction dynamics in **the origin of P_{11} N^* resonances**.
 - ➔ **Meson-cloud effect** on the transition form factors.

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✓ Major topics in N^* & Δ^* spectroscopy

- Establishing **high-mass N^* & Δ^* resonances** [$\text{Re}(M_R) > 1.7 \text{ GeV}$]
 - ➔ “(over-)complete” **experiments** for photoproduction reactions (**CLAS6, ELSA, MAMI,...**)
- Determining Q^2 dependence of **electromagnetic transition form factors** for well-established low-lying N^* & Δ^* resonances.
 - ➔ Measurements of electroproduction reactions **over wide Q^2 range** (**CLAS6, CLAS12**)

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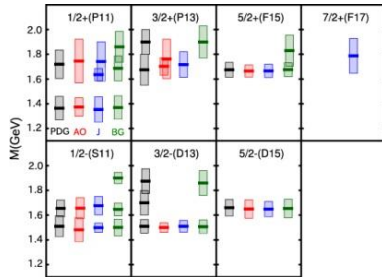
Electroproduction analysis & extension of our DCC model are underway !!

Back up

Applications of ANL-DCC approach

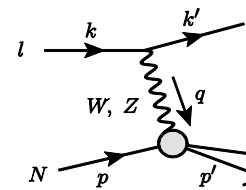
N* & Δ* spectroscopy

- Early analyses of πN & γN reactions:
 PRC76(2007)065201; 77(2008)045205; 78(2008)025204
 PRC79(2009)025206; 80(2009)065203; 81(2010)065207
 PRL104(2010)042302
- Latest analysis of πN & γN reactions:
 PRC88(2013)035209; 88(2013)045203; 94(2016)015201
- Electroproduction analysis & Form factor extraction:
 PRC80(2009)025207; 82(2010)045206

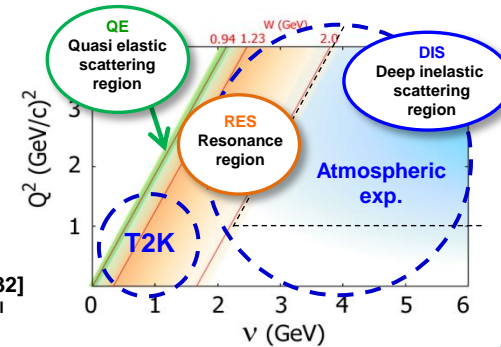


Neutrino reactions

- Calculation in $Q^2 = 0$ limit:
 PRD86(2012)097503
- Full DCC-model calculation up to $W = 2$ GeV, $Q^2 = 3$ GeV²:
 PRD92(2015)074024

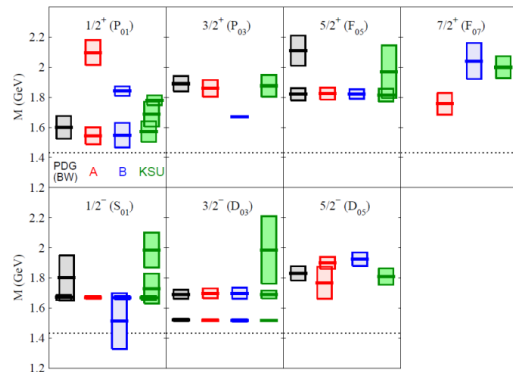


Collaboration@J-PARC Branch,
 KEK Theory Center [arXiv:1303.6032]
<http://j-parc-th.kek.jp/html/English/e-index.html>



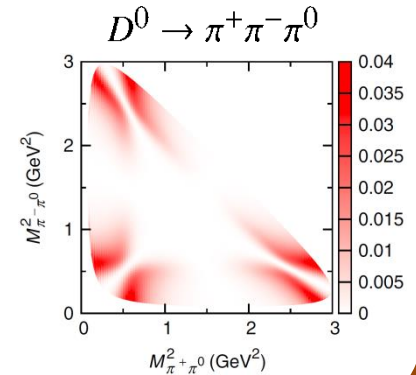
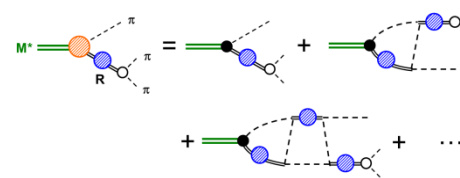
ANL-Osaka DCC approach

- Λ^* , Σ^* resonance extractions via analysis of K - p reactions:
 PRC90(2014)065204; 92(2015)025205



Λ* & Σ* spectroscopy

- Formulation of 3-body unitary model for decays of mesons:
 PRD84(2011)114019
- Application to $\gamma p \rightarrow M^* N \rightarrow (3\pi)N$:
 PRD86(2012)114012



Meson spectroscopy

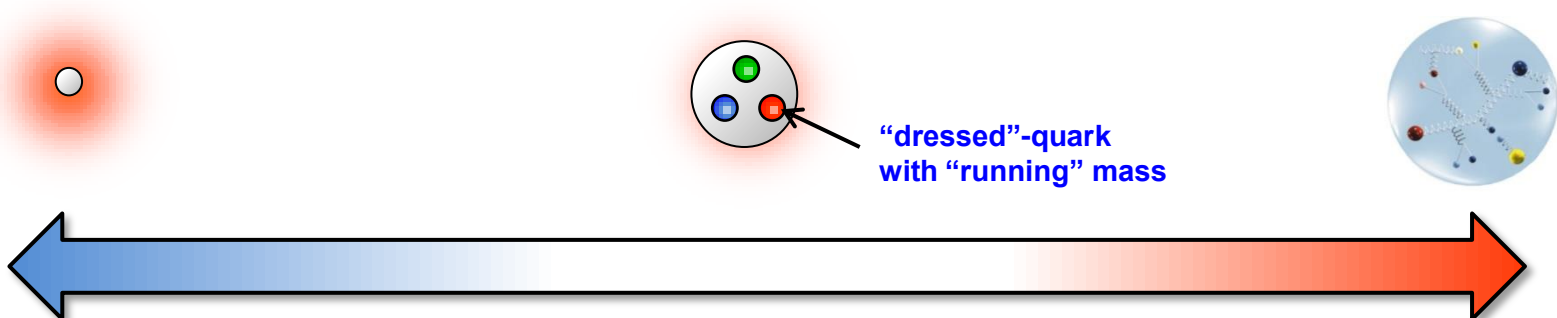
Substructure of N^* & Δ^*

See, e.g., ECT* workshop “Nucleon Resonances: From Photoproduction to High Photon Virtualities”, Oct. 2015
<http://www.ectstar.eu/node/1227>

“dressed”-quark core
 obscured by dense meson clouds

Meson clouds become small;
 “dressed”-quark core dominates

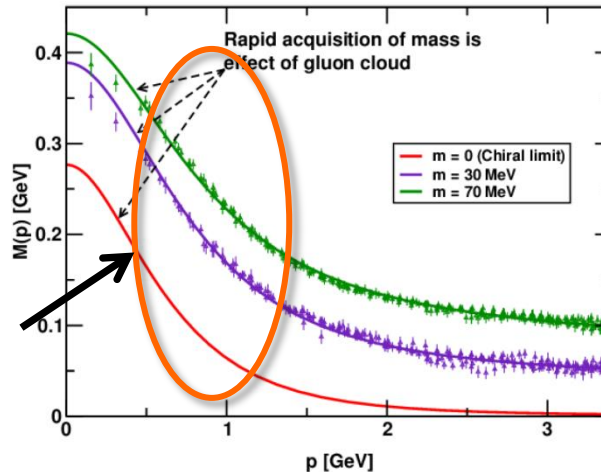
“Partons”



Q^2 : small
 (low “resolution”)

Q^2 : large
 (high “resolution”)

Running dressed quark mass
 “constituent” quark ↔ “current” quark



Curves: a model based on
 Dyson-Schwinger
 equations (Landau gauge)

Points: Lattice QCD

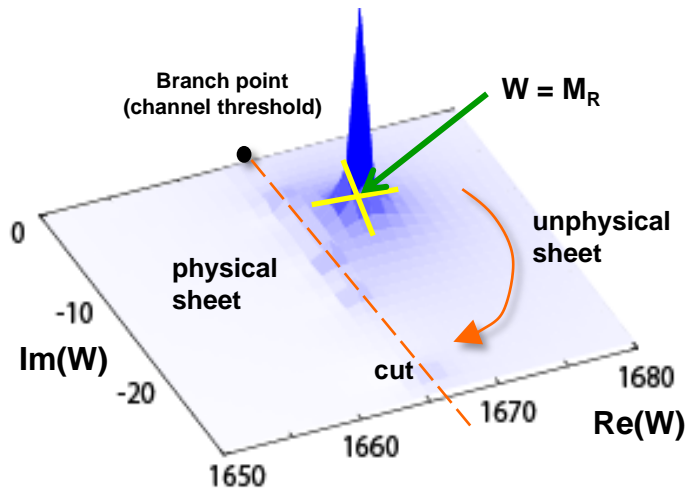
e.g.) Cloet, Roberts,
 Prog.Part.Nucl.Phys.77(2014)1

Will be accessed by
 CLAS12@JLab
 (E12-09-003, E12-06-108A)

Baryon resonances as poles of scattering amplitudes

PROPER definition of

- ✓ **Hadron resonance masses** (complex) → **Pole positions** of **scattering amplitudes** in the lower-half of complex-W plane
- ✓ **Transition amplitudes** between resonance and scattering states → ~ **Residues**^{1/2} at **the pole**



$$\langle p_b | \hat{T}(E) | p_a \rangle \Big|_{E \rightarrow M_R} \rightarrow \frac{R_{b,a}}{E - M_R} + (\text{regular terms})$$

Residue at the pole
 $\sim \langle p_b | V | R \rangle \times \langle R | V | p_a \rangle$

Resonance pole position
 ($\text{Im}(M_R) < 0$)

Extracting poles of amplitudes from analyzing reaction data is nothing but obtaining **“exact” energy eigenvalues of QCD !!!**

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Resonance theory based on Gamow vectors:

[G. Gamow (1928), R. E. Peierls (1959), ...]

“Quantum resonance state is an **(complex-)energy eigenstate** of the **FULL** Hamiltonian of the **underlying theory** under the Purely Outgoing Boundary Condition (POBC).”

Energy eigenvalue = **pole energy**

Transition matrix elements between resonance and scattering states ~ **Residues^{1/2}** at **the pole**

Extracting poles of amplitudes from analyzing reaction data is nothing but obtaining **“exact” energy eigenvalues of QCD !!!**

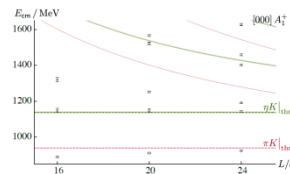
Baryon resonances as poles of scattering amplitudes

PROPER definition of

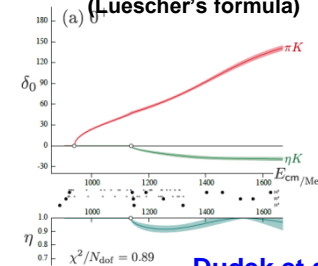
- ✓ **Hadron resonance masses** (complex)
- ✓ **Transition amplitudes** between resonance and scattering states

There are attempts to link **real energy spectrum of QCD in the finite volume** to **resonance pole masses**.

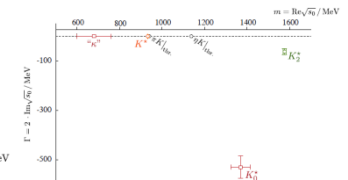
QCD spectrum in finite volume



Extracted phase shifts & inelasticity (Luescher's formula)



Extracted resonance pole positions (K-matrix analysis)



Dudek et al, PRL113(2014)182001

Resonance theory based on
[G. Gamow (1928), R. E. Peierls

“Quantum resonance state is an eigenstate of the **FULL** Hamiltonian of the **underlying** system with **Purely Outgoing Boundary Condition**”

Energy eigenvalue

Transition matrix elements between resonance and scattering states

pole energy

Residues^{1/2} at the pole

Extracting poles of amplitudes from analyzing reaction data is nothing but obtaining **“exact” energy eigenvalues of QCD !!!**