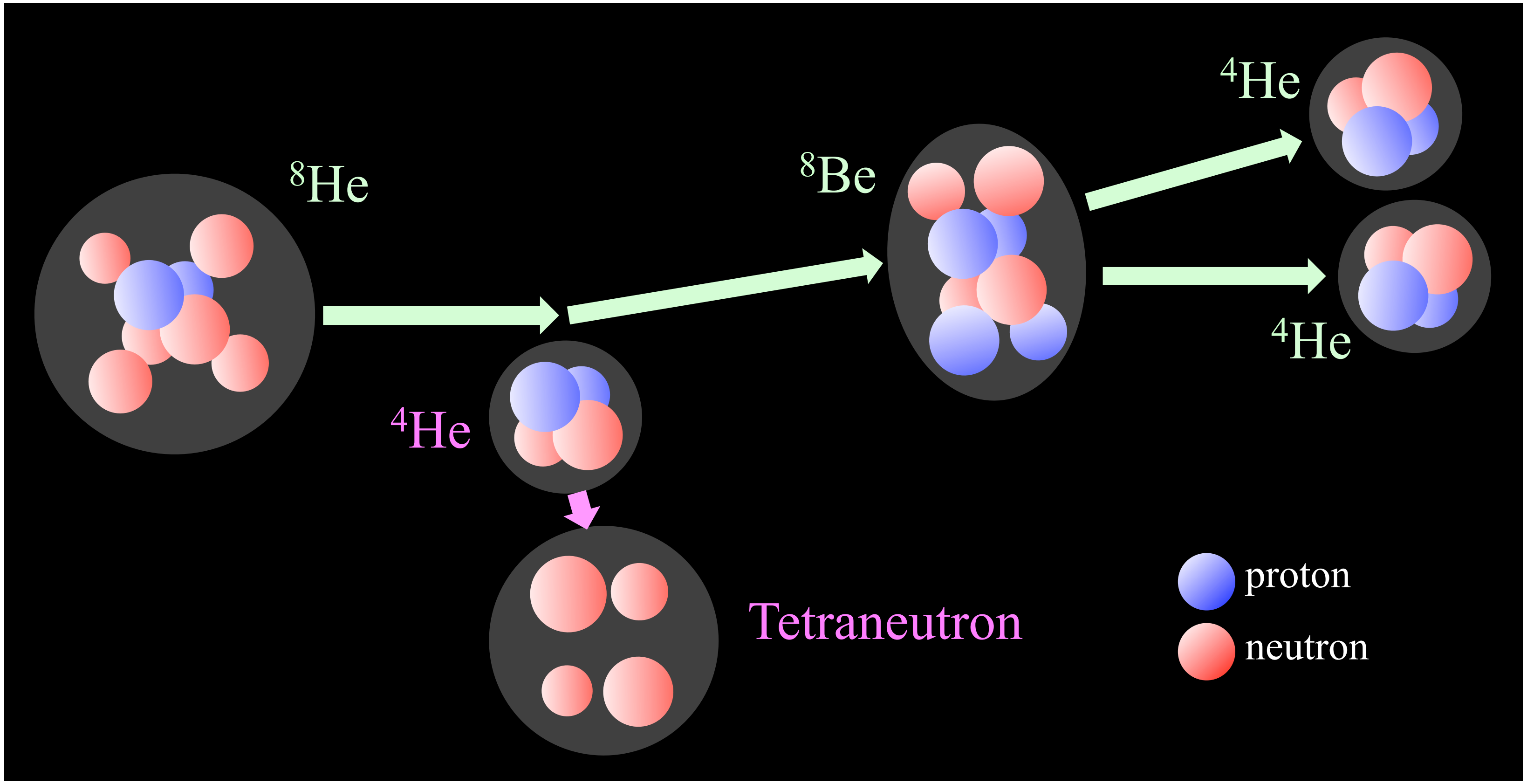
Tetra-neutron system studied by double-charge exchange reaction $^4He(^8He,^8Be)$



S. Shimoura

CNS, University of Tokyo

Motivation

Tetra-neutron

- Multi-neutron System
  - Neutron cluster (?) in fragmentation of ^{14}Be
    PRC65, 044006 (2002)
  - NN, NNN, NNNN interactions
    - T=3/2 NNN force
      -> 3-body force in neutron matter
    - Ab initio type calculations
  - Multi-body resonances
  - Correlations in multi-fermion scattering states

Historical Review

~ search for a bound state of 4n~

1960s

fission of Uranium

- No evidence for particle stable state of tetra-neutron

J. P. Shiffer Phys. Lett. 5, 4, 292 (1963)

1980s

^{4}He(π^-, π^+) reaction

- Only upper limit of cross section was decided.

J. E. Unger, et al., Phys. Lett. B 144, 333 (1984)

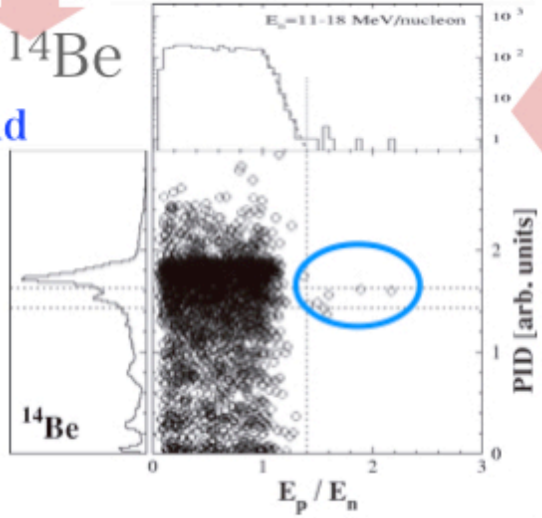
Bound state: No clear evidence.

2000s

Breakup of ^{14}Be

- Candidates of **bound tetra-neutron** were observed.

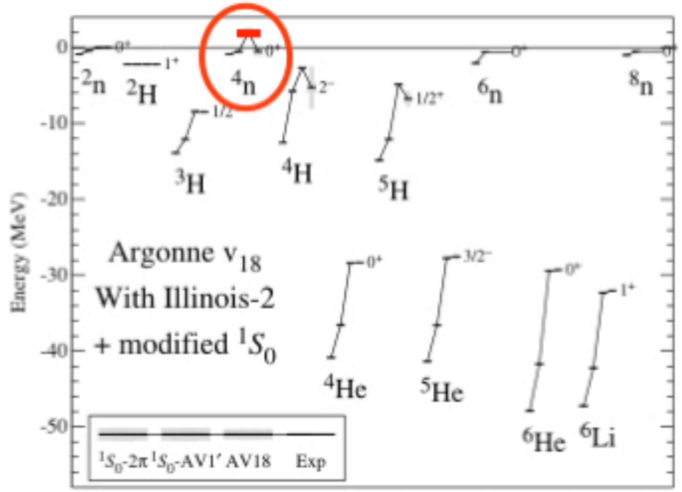
F. M. Marques, et al, Phys. Rev. C 65, 044006 (2002)



2000s

Theoretical work

- ab-initio calculation
  NN, NNN interaction

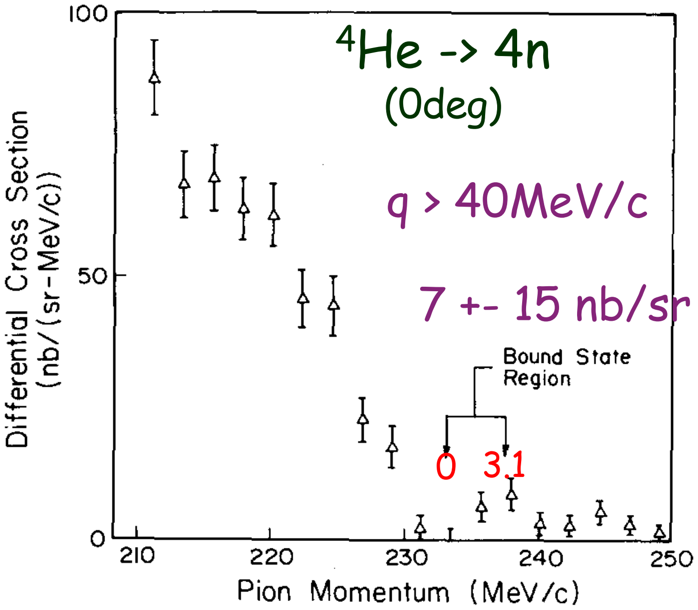


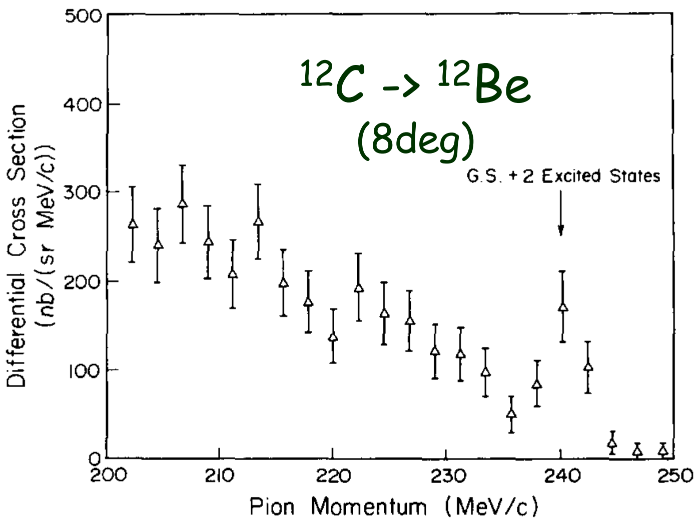
S. C. Piper, Phys. Rev. Lett. 90, 252501 (2003)

- **Bound ^{4}n cannot exist**
- **Possible resonance stete ~2 MeV**

Resonance state : Possibility of the state is still an open and fascinating question.

(π^-,π^+) reaction @ 165 MeV; $\theta_{\pi+}$ = 0 degree





The peak is due primarily to the transition to the ^{12}Be ground state, with some contribution from the first two excited states as well.

We have measured the momentum spectrum of π^+ produced at 0° by 165 MeV π^- on ^{4}He. A $\Delta P/P$ = 1% beam of 10^6 π^- per second was provided by the P^3 line of the Los Alamos Meson Physics Facility, and a cell of 910 mg/cm^2 liquid ^{4}He with windows of 18 mg/cm^2 Kapton served as the target [15]. An

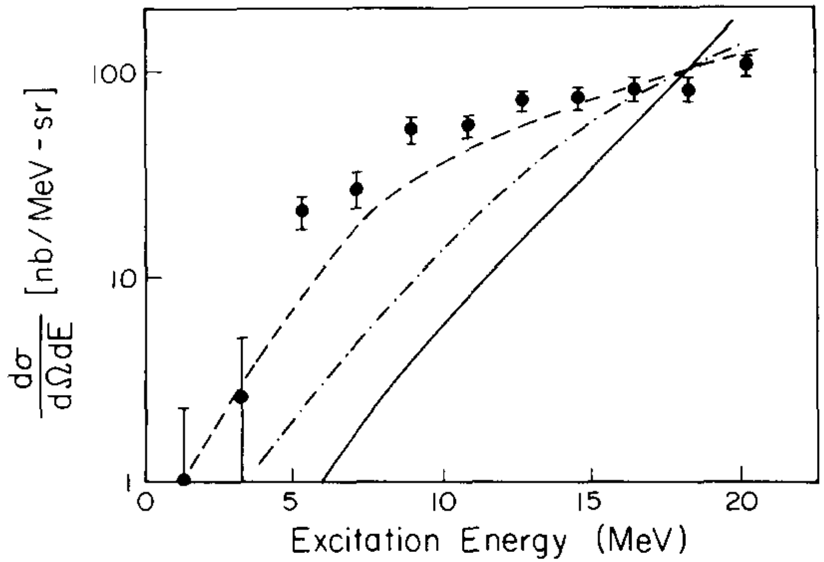
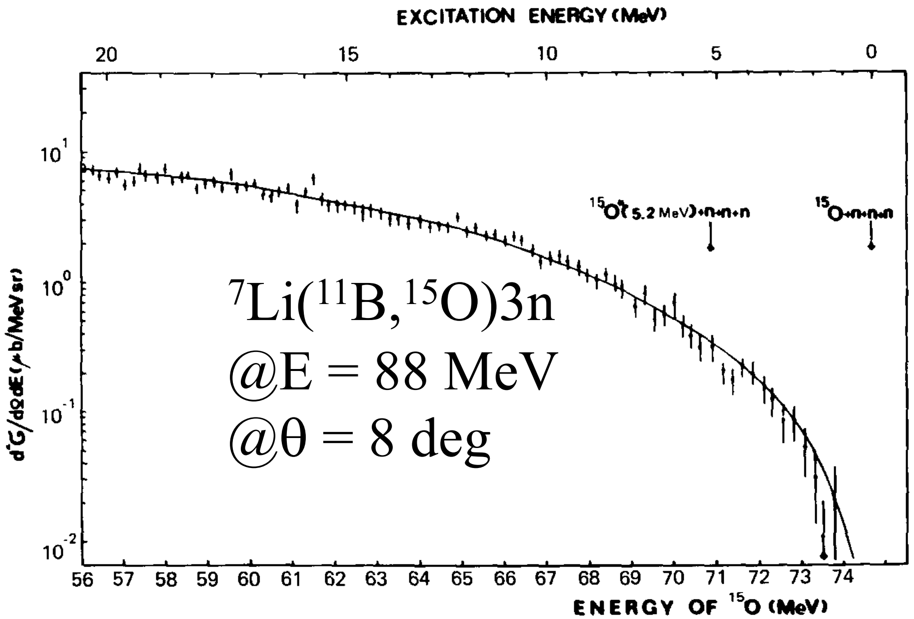


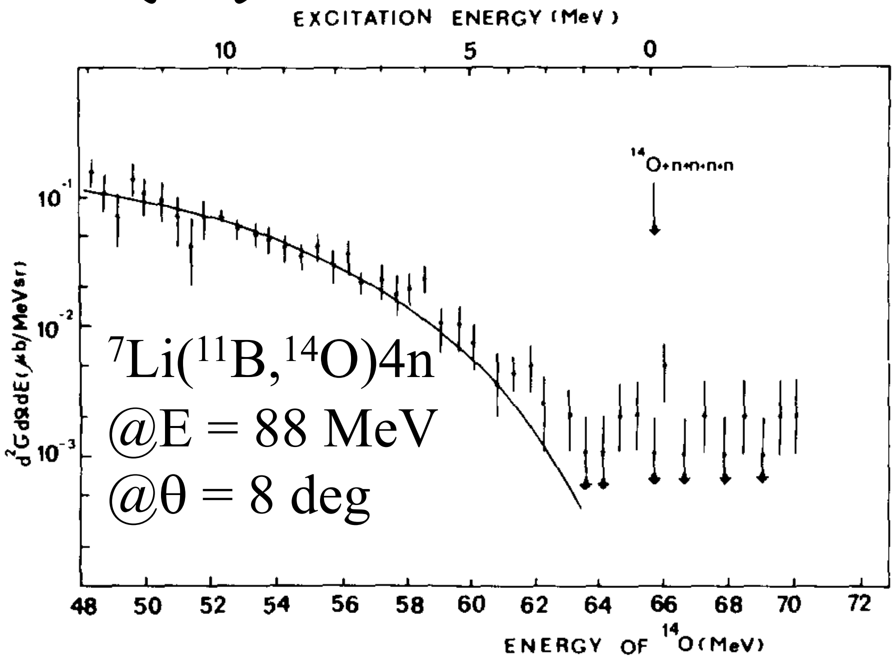
Fig. 3. The experimental results are plotted against the excitation of the final four-neutron state. The solid curve corresponds to the pure four-neutron phase space, while the dot-dashed and dashed curves are the four-neutron phase space curves with singlet state interactions in, respectively, one and both of the final state neutron pairs.

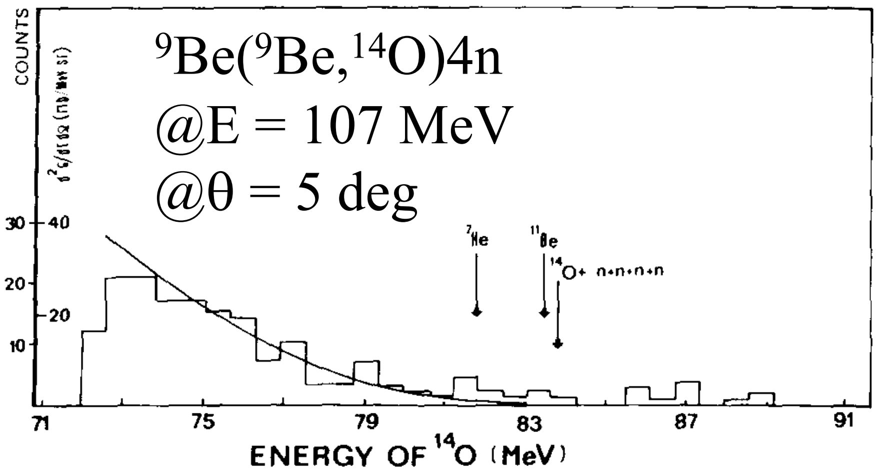
J.E. Ungar et al., PLB 144 (1987) 333

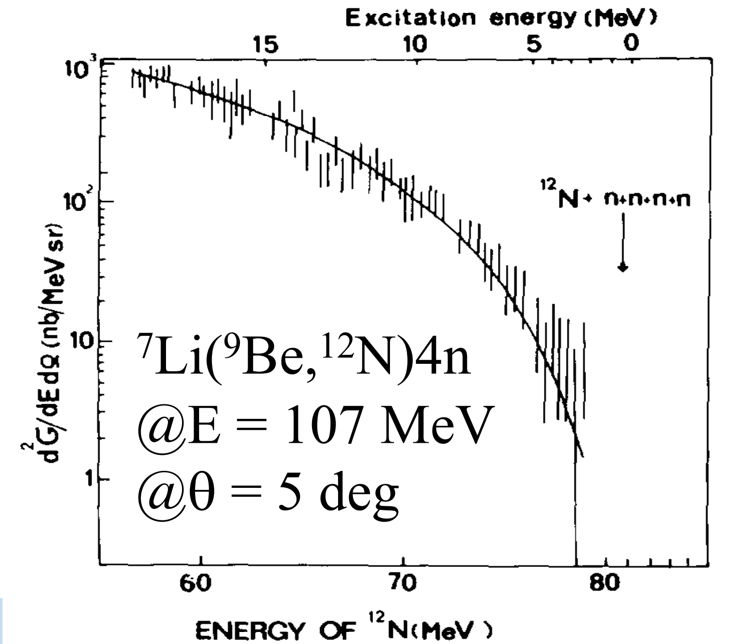
Historical review (2)

Nucl. Phys. A477 (1988) 131



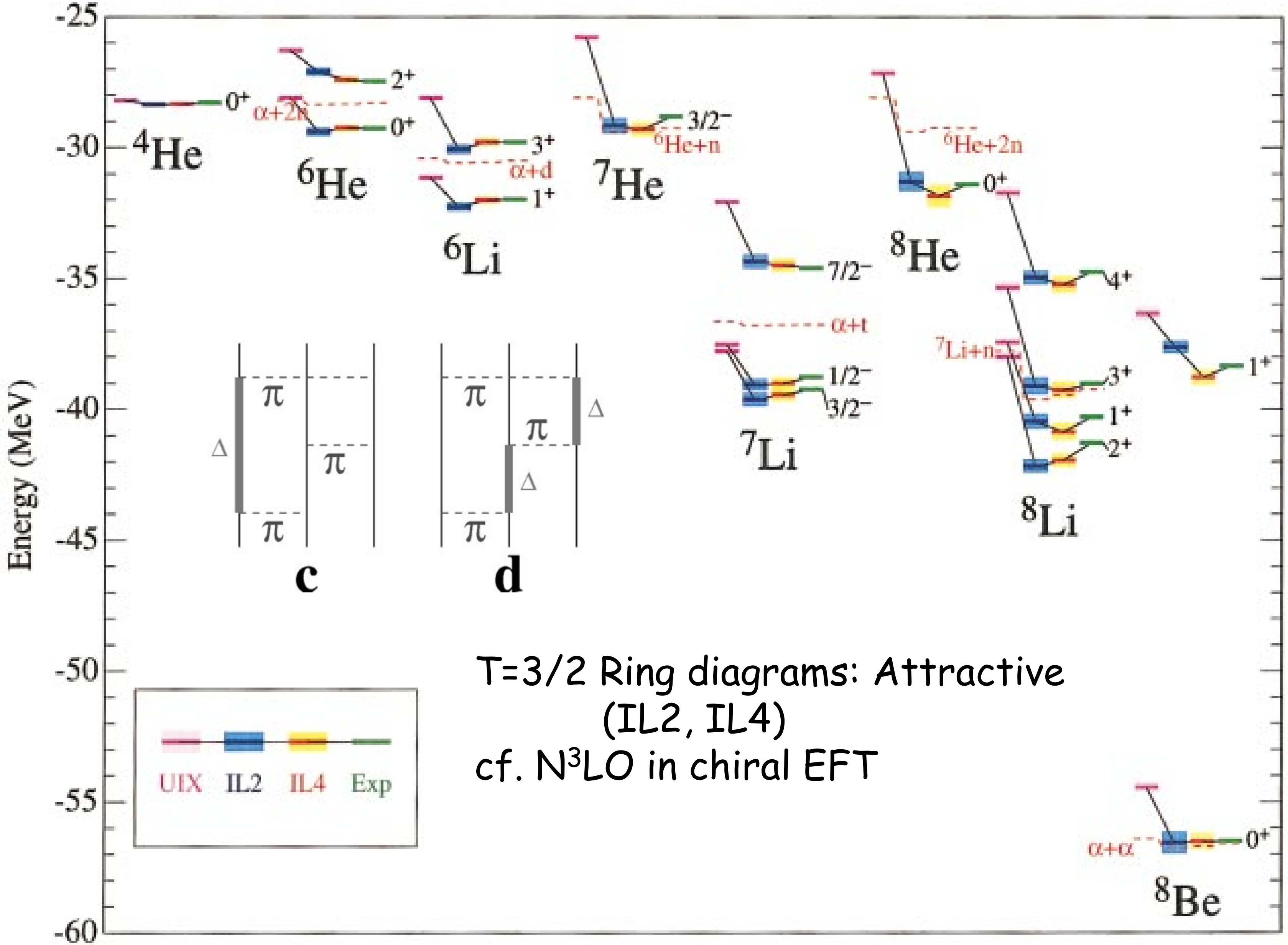






3-body force

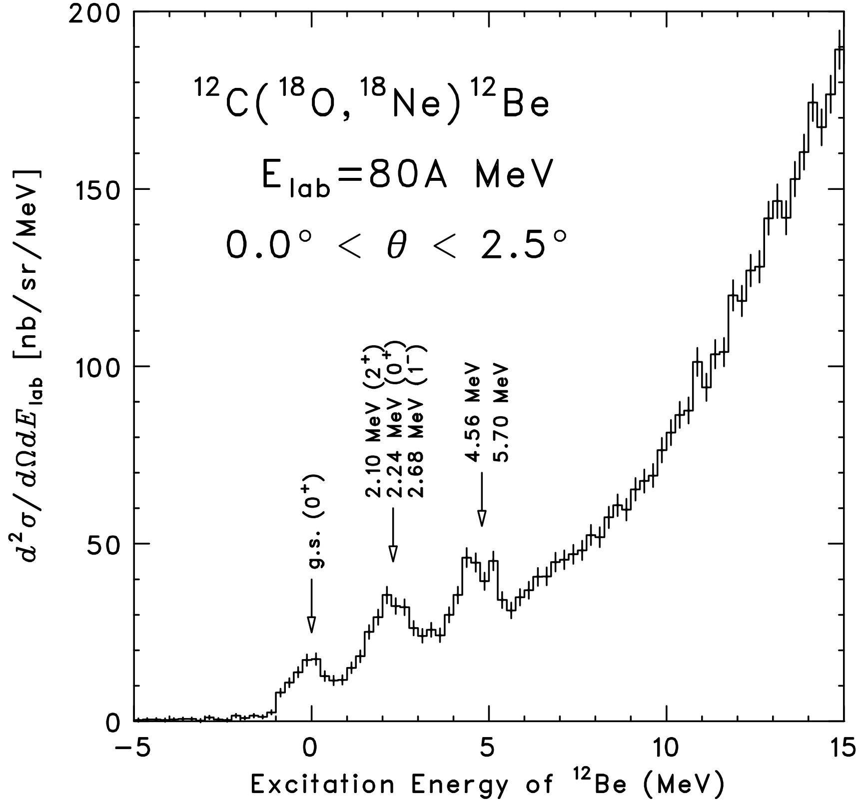
S. C. Pieper, et al., PRC 64, 014001



T=3/2 Ring diagrams: Attractive (IL2, IL4)

cf. N^{3}LO in chiral EFT

Double charge exchange (DCX) reaction of HI



^{12}C -> ^{12}Be

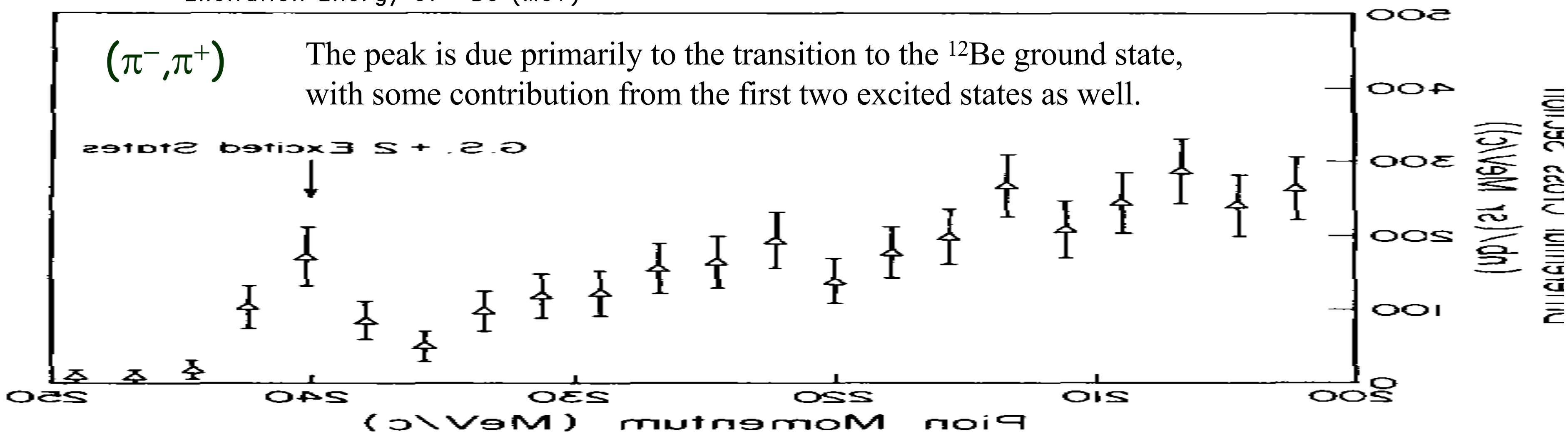
Stable ^{18}O beam (80A MeV) (Takaki et al.)

~70nb/sr (Gnd)
~200nb/sr (~2MeV)

HI DCX reaction can be used for spectroscopy for exotic nuclei
(q is not so small >80 MeV/c)

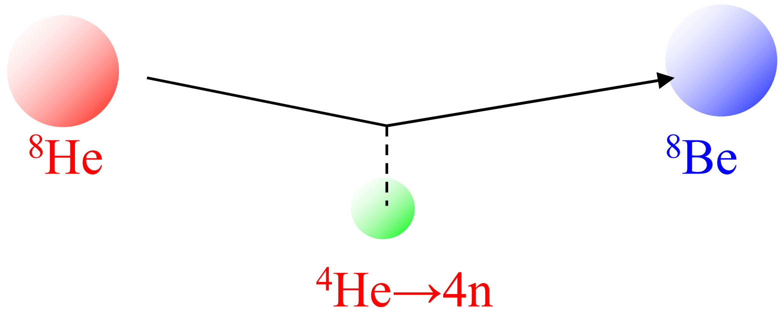
(π^-,π^+)

The peak is due primarily to the transition to the ^{12}Be ground state, with some contribution from the first two excited states as well.



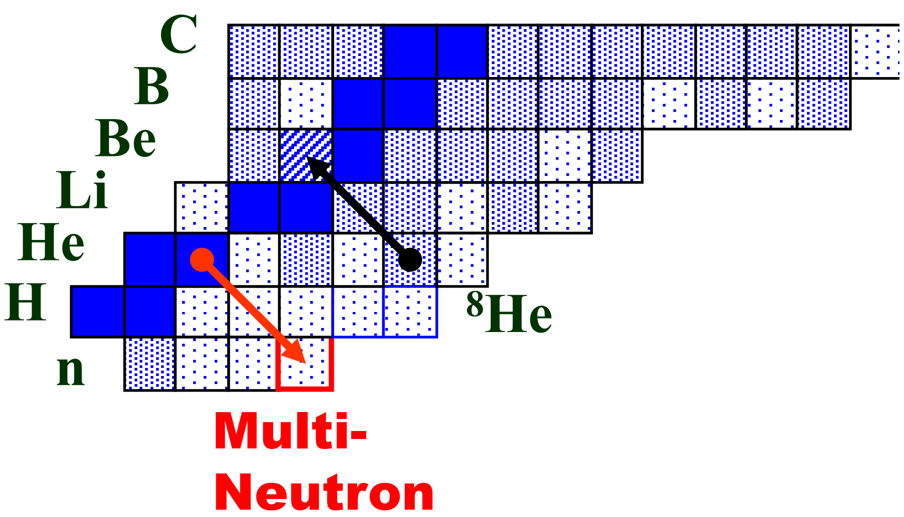
Exothermic DCX Reaction

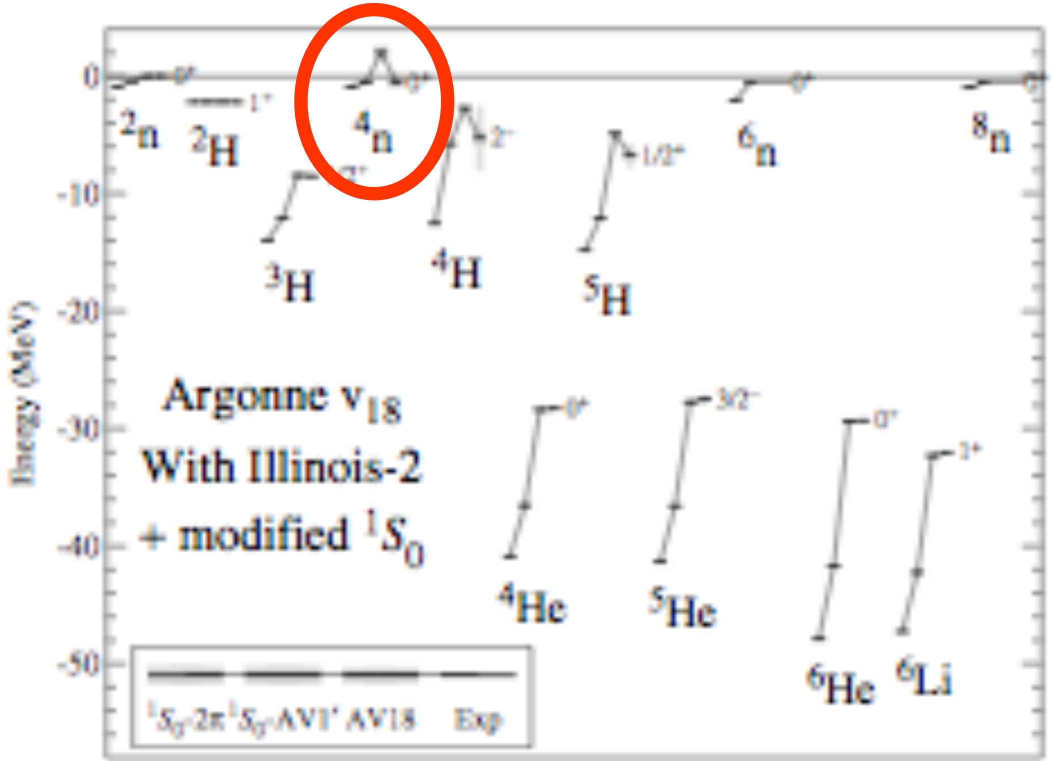
Tetra-neutron system produced by exothermic double-charge exchange reaction



Almost recoil-less condition with ^{4}He(^{8}He,^{8}Be)4n reaction at 200 A MeV

Recoil-less 4n system via DCX using internal energy of ^{8}He

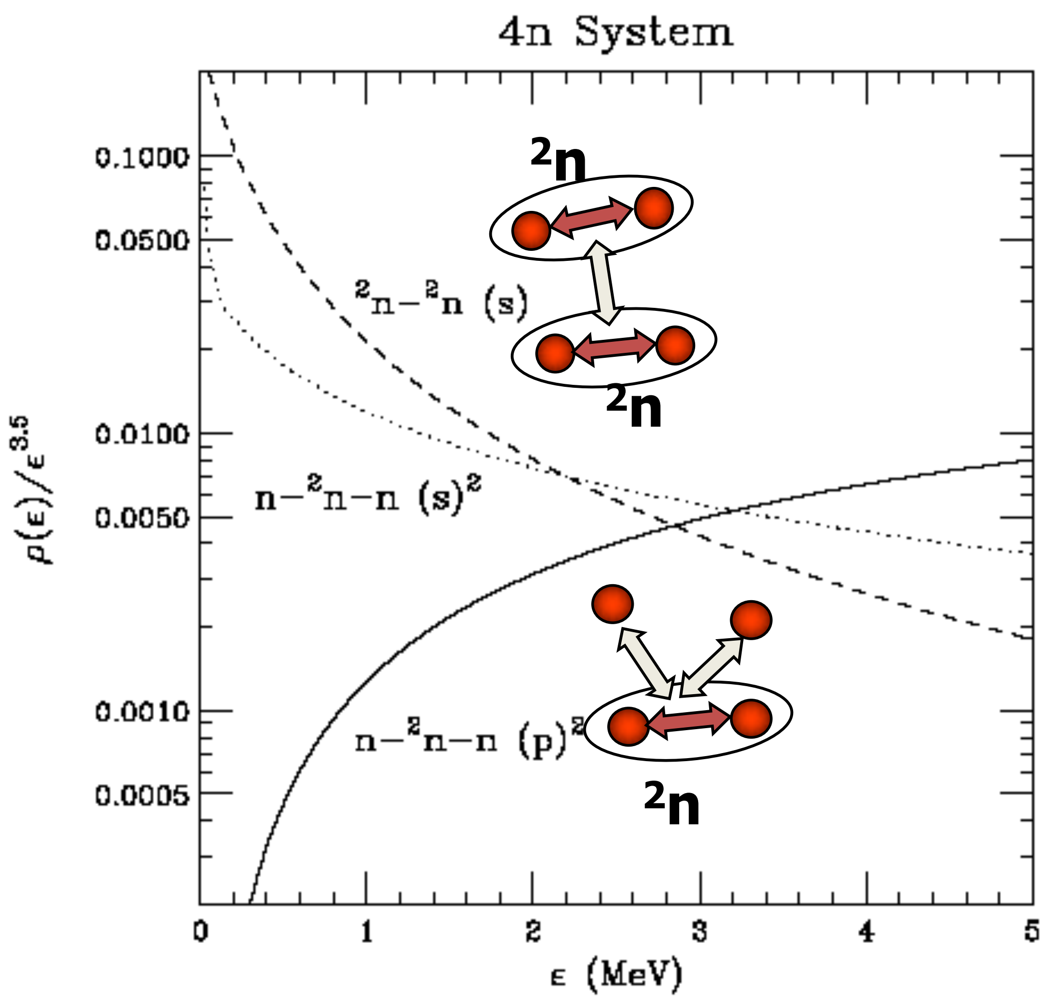




S.C. Pieper et al., PRL 90, 252501 (2003)

4n in breakup of ^{14}Be : Marques et al. PRC 65 (2002) 044006

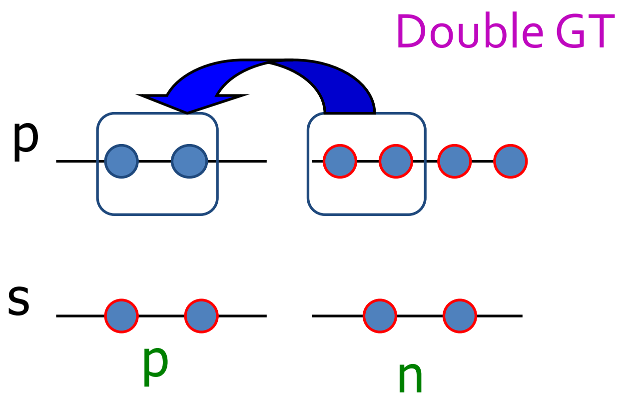
Correlation in multi-body continuum



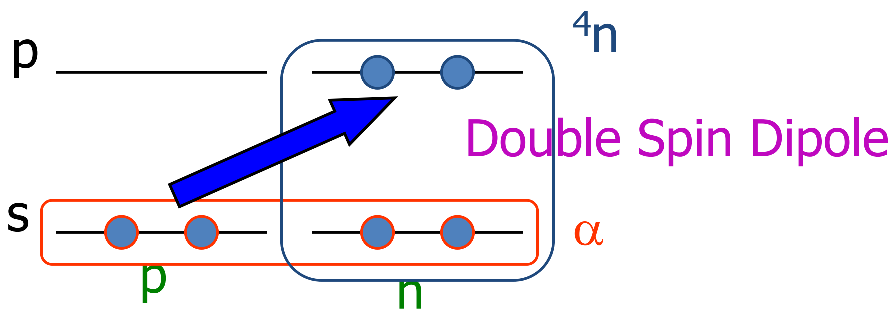
- Deviation from four-body phase space informs us the final state interaction(s) of sub-system

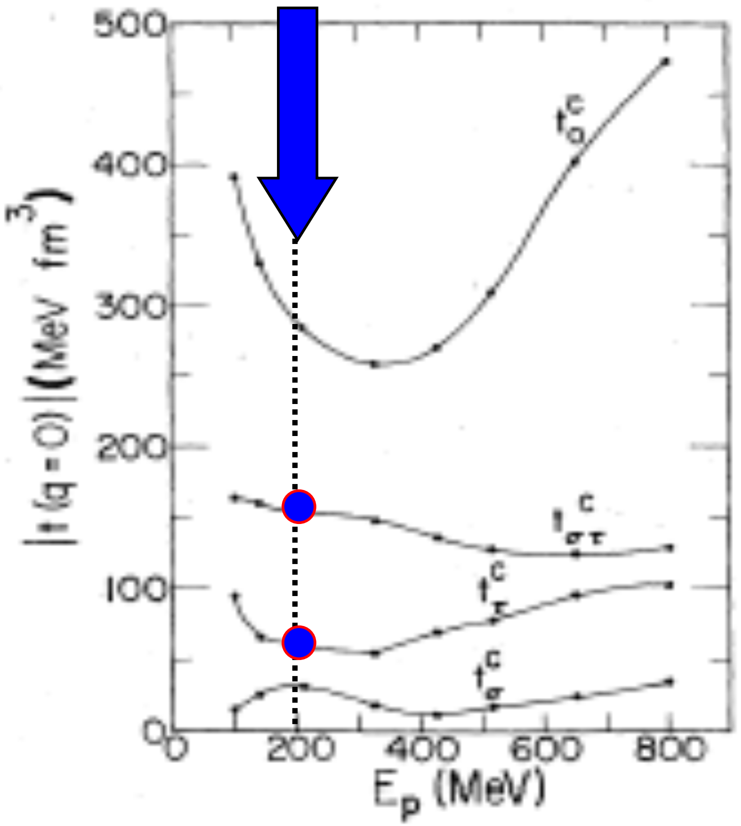
Reaction Mechanism

$^{8}He \rightarrow {}^{8}Be$



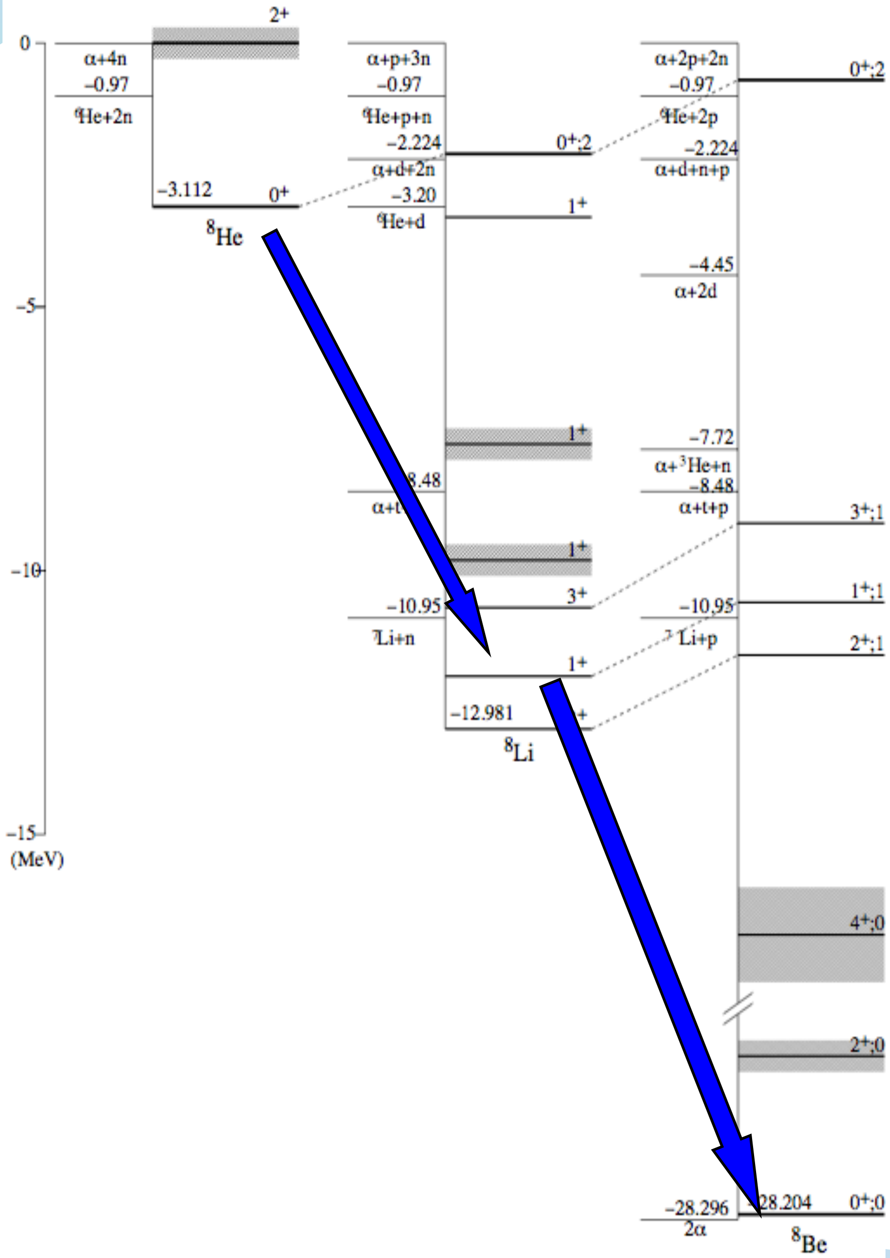
$^{4}He \rightarrow 4n$

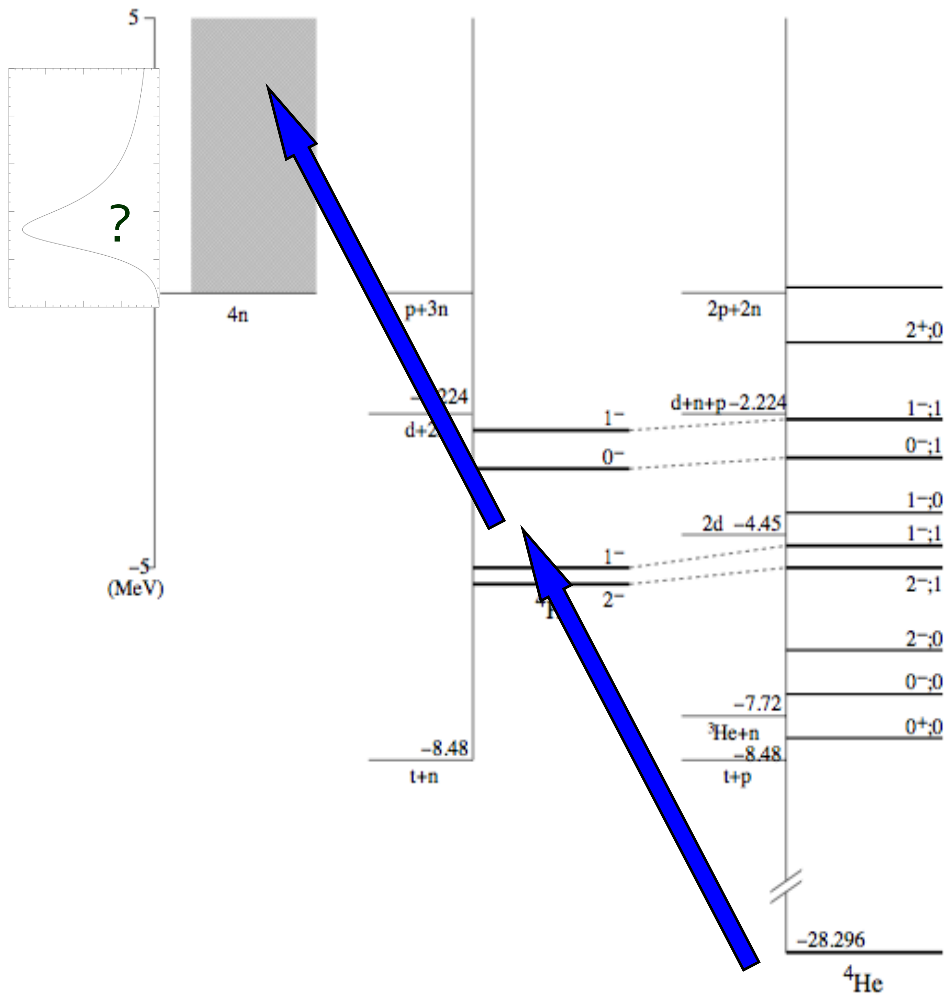




$$\left[\left(\vec{\tau}_{\mathrm{p}}\cdot\vec{\tau}_{\mathrm{t}}\right)\left(\vec{\sigma}_{\mathrm{p}}\cdot\vec{\sigma}_{\mathrm{t}}\right)r_{\mathrm{t}}Y_{1}\left(\hat{r}_{\mathrm{t}}\right)\right]^{2}$$

Level diagrams



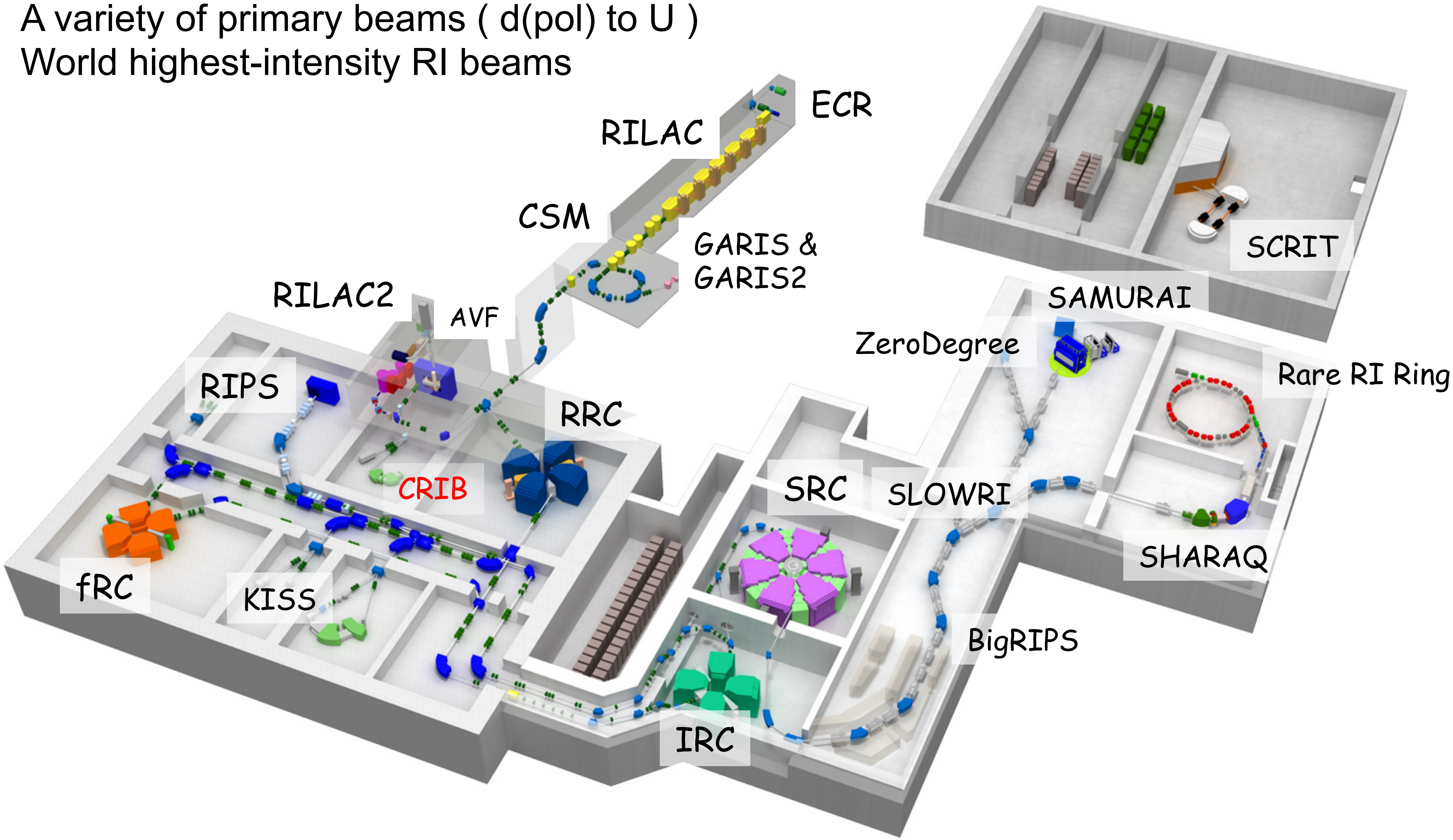


q_{min} ~ 10MeV/c

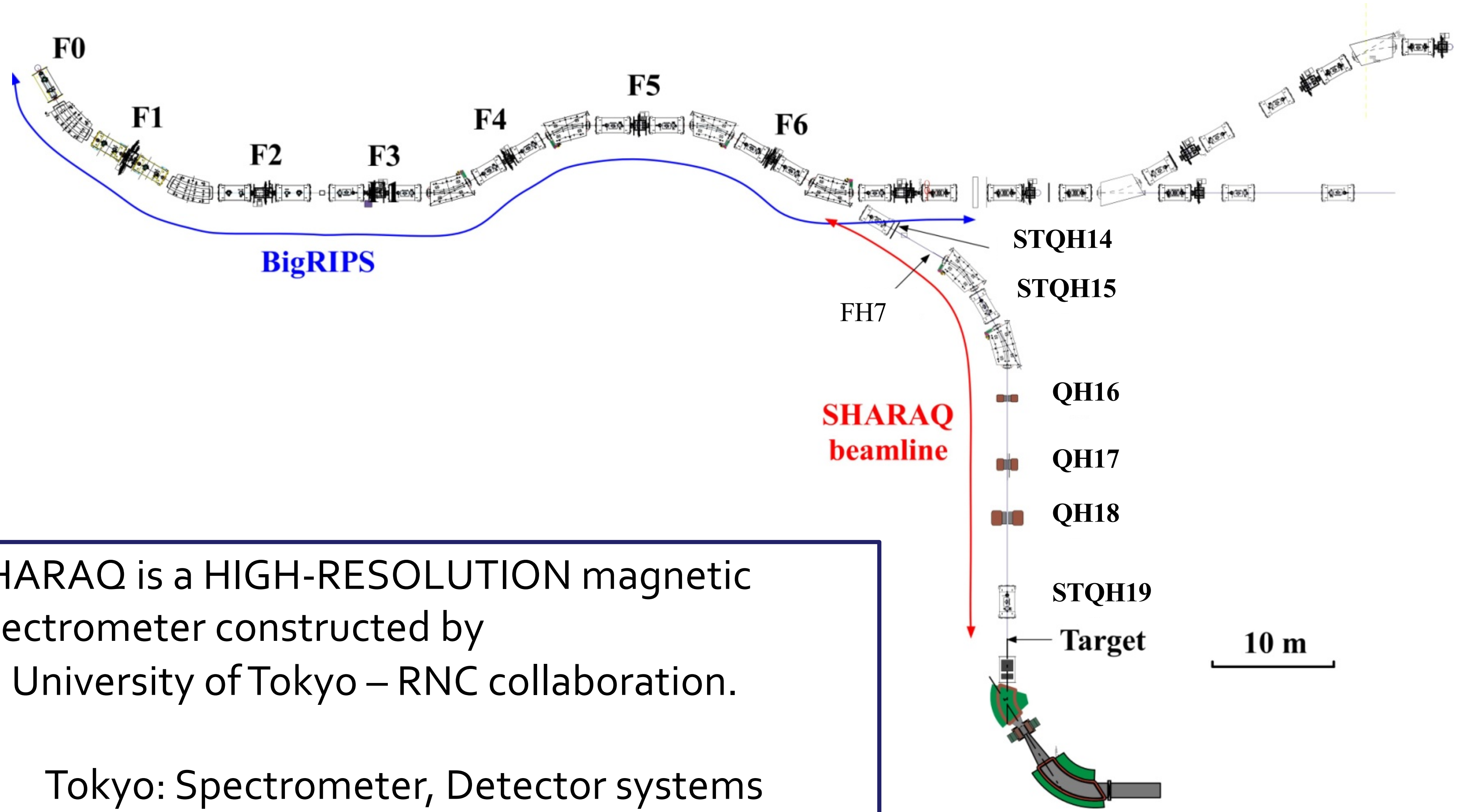
Experiment

RI Beam Factory at RIKEN

3 injectors + cascade of 4 cyclotrons
⇒ several to 345 MeV/nucleon
A variety of primary beams (d(pol) to U)
World highest-intensity RI beams



SHARAQ @ RI beam factory

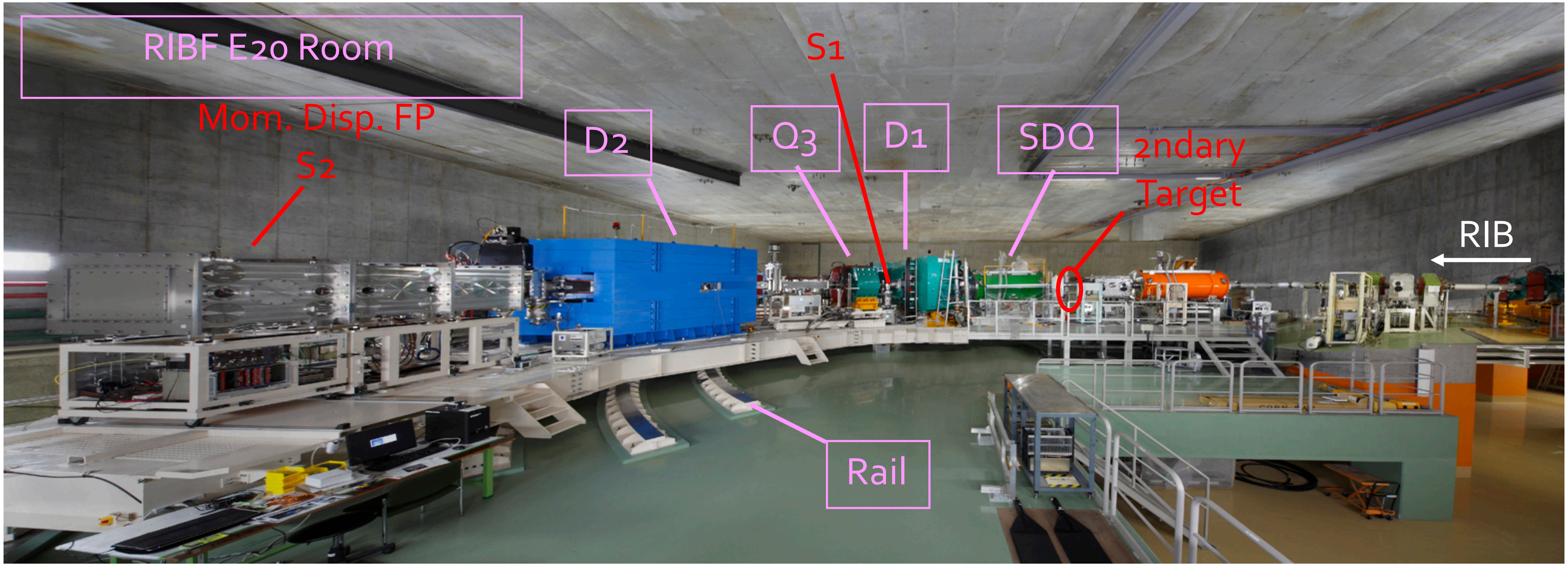


SHARAQ is a HIGH-RESOLUTION magnetic spectrometer constructed by
University of Tokyo – RNC collaboration.

Tokyo: Spectrometer, Detector systems
RNC: Beam-line, Infrastructure

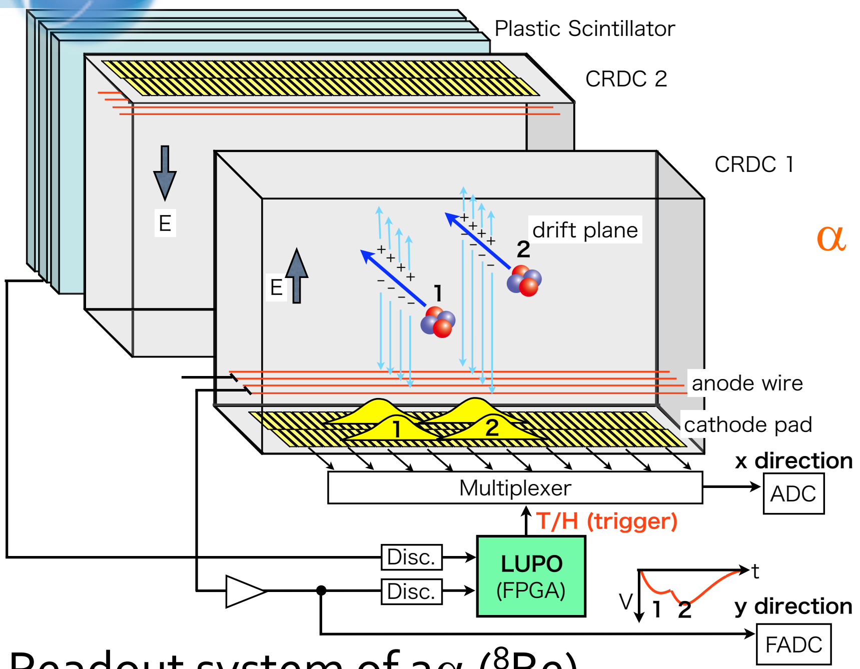
SHARAQ spectrometer

T. Uesaka et al.,
NIMB B **266** (2008) 4218.
PTEP 2012, 03C007 (2012)



Maximum rigidity	**6.8 Tm**
Momentum resolution	**d*p*/*p* = 1/14700**
Angular resolution	**~ 1 mrad**
Momentum acceptance	**± 1%**
Angular acceptance	**~ 5 msr**

銀杏 ICHOR

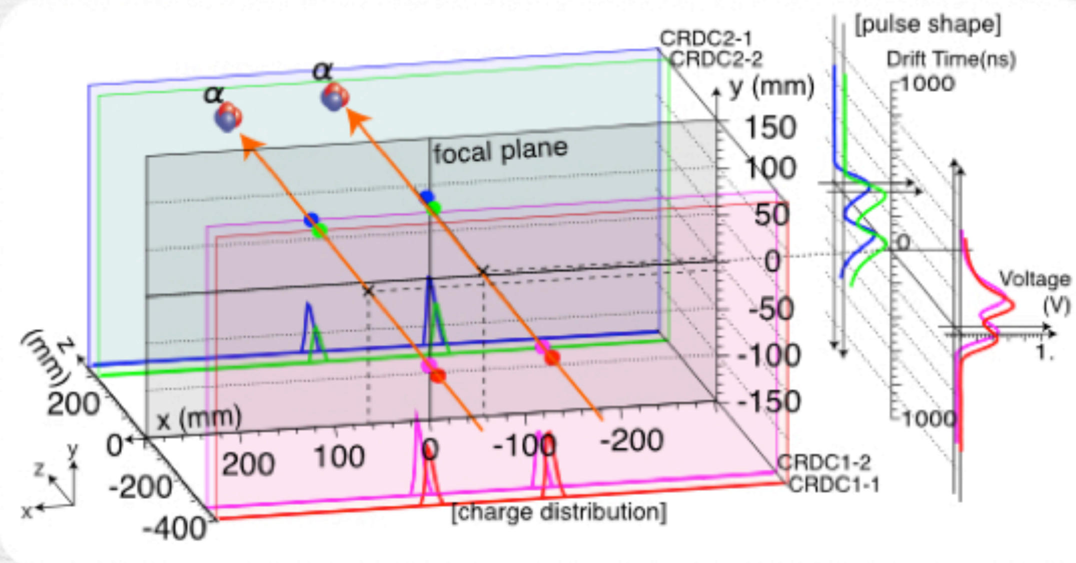


Readout system of 2α (^{8}Be)

Analysis

Selection of 4n Events

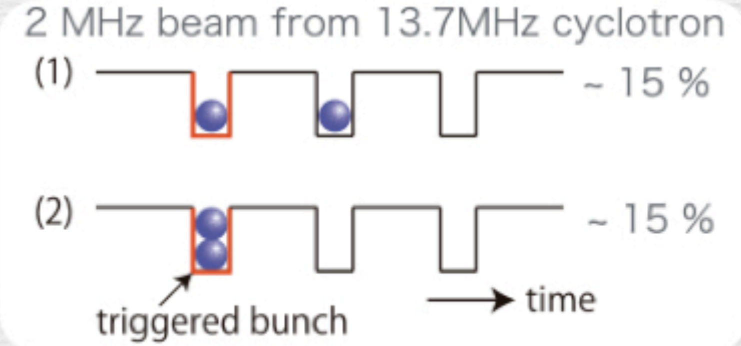
- Extracting 2α events @SHARAQ
- Multi-particle in high-intensity beam



Background process:
Breakup of two ^{8}He in the same beam bunch to two alpha particle
Identified by multi-hit in F6-MWDC

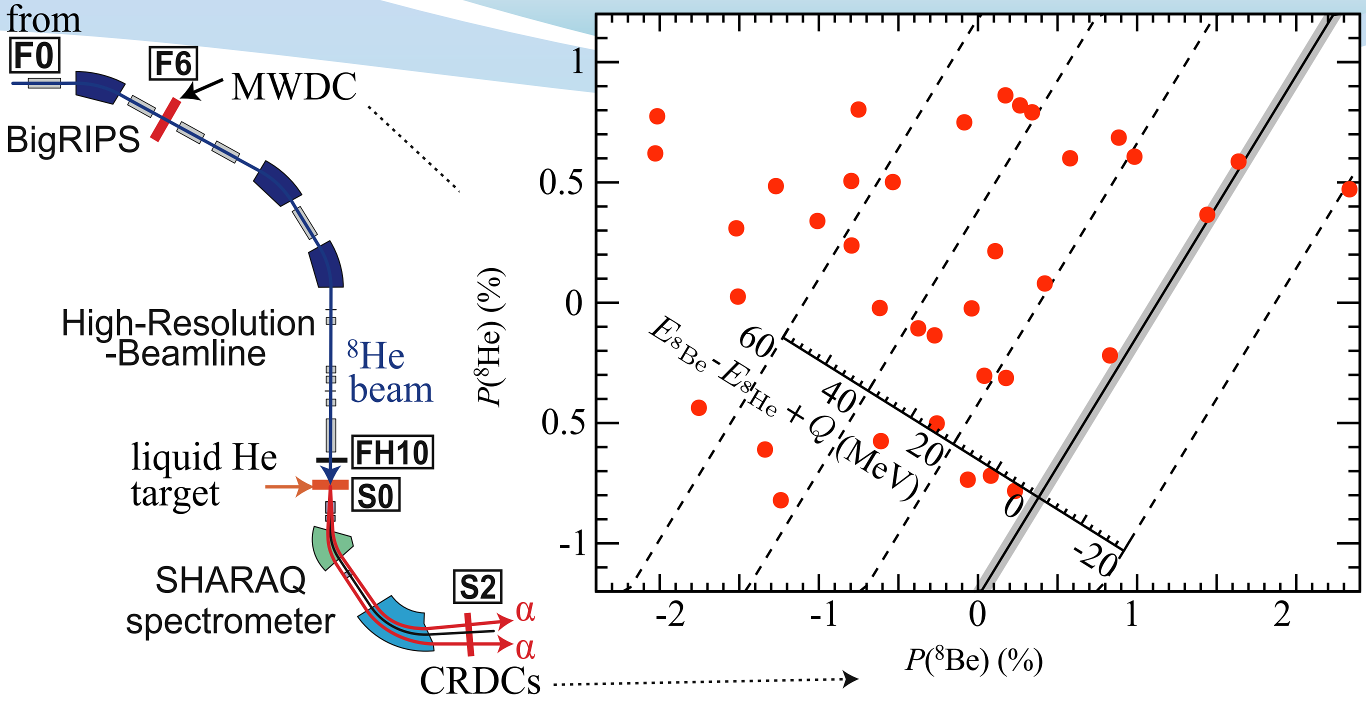
Background Estimation

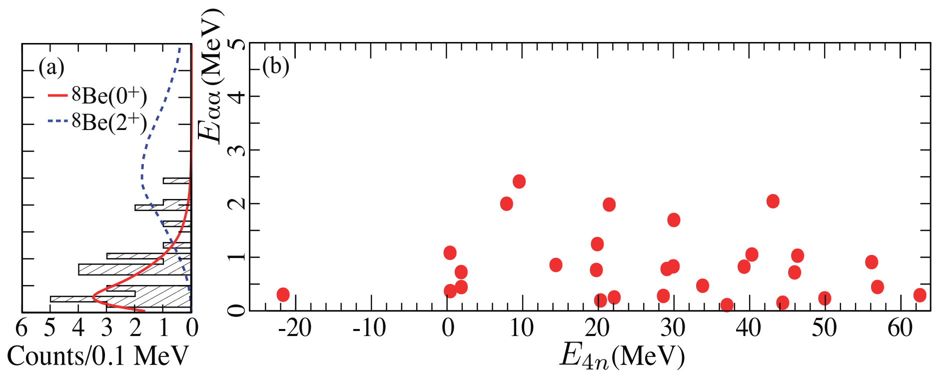
- Shape in spectrum: random 2α
- Number of events:
  - failure of the multi-particle rejection at MWDC
  - multi-particle in one cell of MWDC



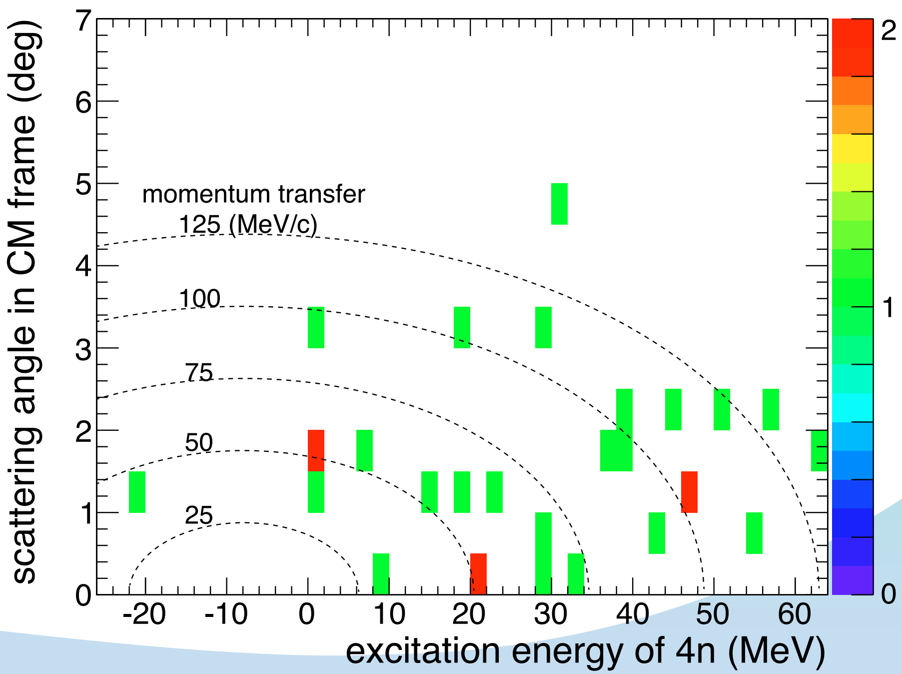
Backgrounds after analysis:
Finite efficiency of multi-hit events at F6-MWDC

Experimental Results

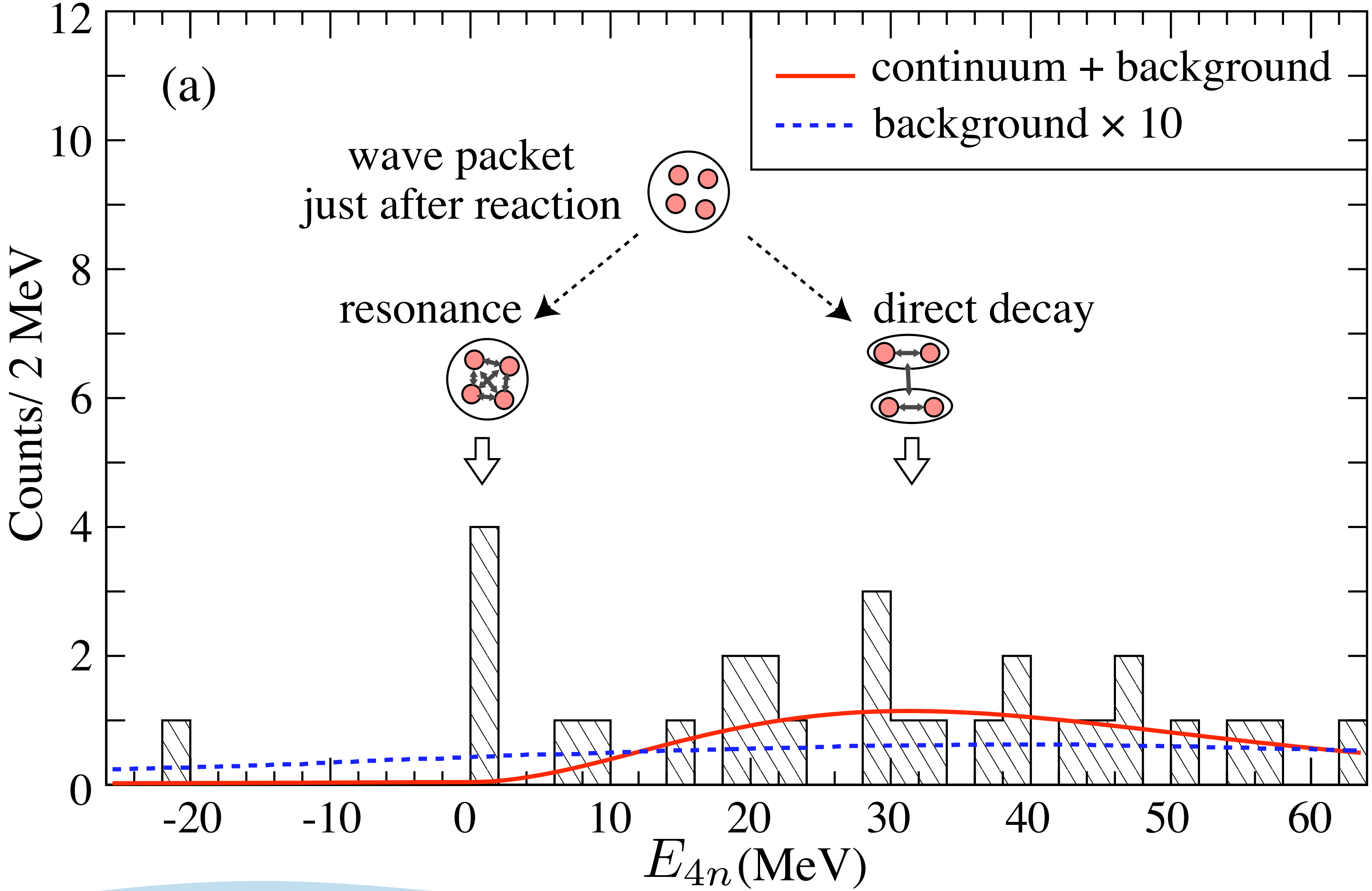




Acceptance for ^{8}Be(2+) was 13 % of that for ^{8}Be(0+)
A few events could be from ^{8}Be(2+).



Experimental Results

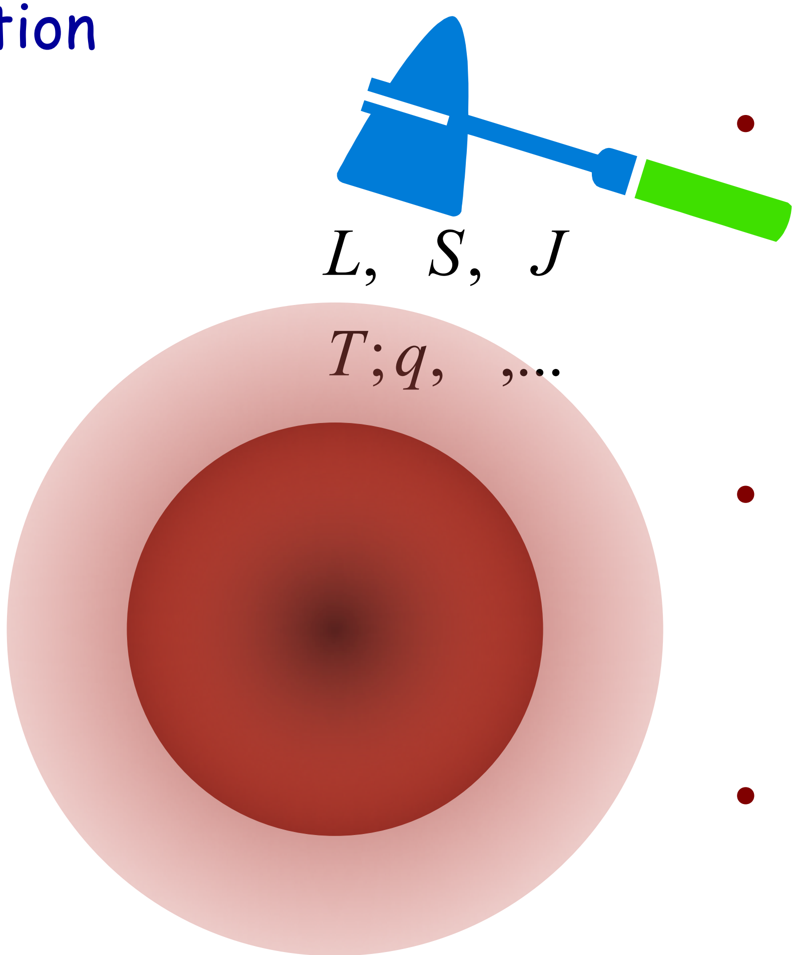


Analysis & Discussion

Studies of Nuclei via Direct reactions

- Size/ρ-distribution
  - Skin/Halo
- Shell Structure
  - New magic #
  - Isospin / Deformation
- New modes
  - IVE1
  - ISE0, ISE1
- Correlation
  - Pairing
  - Clustering
  - etc
- etc.

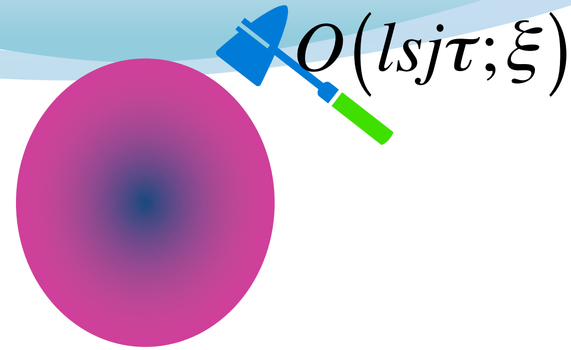
Direct Reactions



- Size/ρ-distribution
  - σ_R, elastic scat.
- Shell Structure
  - Mass / S_n, S_{2n}
  - Inelastic scatt.
    - Low lying states
  - Knockout / Transfer
- New modes
  - Coulex
  - Inelastic scatt.
  - CEX
- Correlation
  - Knockout/Transfer
  - Breakup
  - CEX
- etc.

"Hit and analyze the sound"

Transition Probabilities



$$M_{if} = \langle E_f J_f \pi_f T_f ; \xi_f \| O(lsj\tau;\xi) \| E_i J_i \pi_i T_i ; \xi_i \rangle$$

if distortion is insensitive to ω

$$\text{Cross Section} \propto |M_{if}|^2 \quad ; \quad \text{Lifetime} \propto 1/|M_{if}|^2$$

$O(lsj\tau;\xi)$: Propety of Reaction / Aciton / Decay Processes

sum of
one-body operator

e.g.

$$O(lsj\tau;\vec{r}) = \sum_i f(r_i) T(\tau_i) \left[S(\sigma_i) \otimes Y_l(\hat{r}_i) \right]_j$$

$|E_i J_i \pi_i T_i ; \xi_i\rangle$ and/or $|E_f J_f \pi_f T_f ; \xi_f\rangle$ energy eigen functions

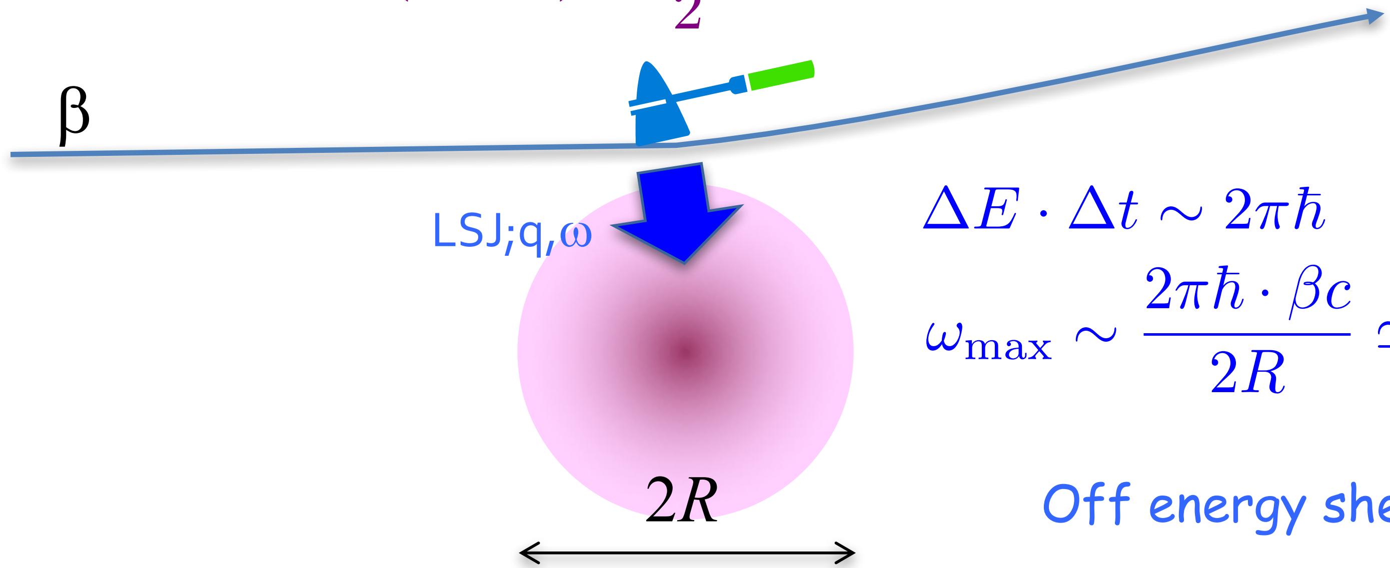
$$O(lsj\tau;\xi) |E_i J_i \pi_i T_i ; \xi_i\rangle = \sum_f\!\!\!\!\!\!\int M_{if}(E_f) |E_f J_f \pi_f T_f ; \xi_f\rangle \quad \text{Response}$$

$|M_{if}(E_f)|^2$: Energy Spectrum

coherent sum of wave packets made by one-body action
"Collective wave packet" (not always energy eigen state),
e.g. coherent sum of 1p-1h for inelastic-type excitation

Reaction time & excitation energy
for intermediate-energy "inelastic-type scattering"

$$\omega \ll \mu c^2 (\gamma - 1) \simeq \frac{1}{2}\mu c^2 \beta^2$$

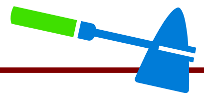


$$\Delta E \cdot \Delta t \sim 2\pi\hbar$$

$$\omega_{\max} \sim \frac{2\pi\hbar \cdot \beta c}{2R} \simeq 100\beta \ \ \mathrm{MeV}$$

Off energy shell

E/A ~ 200 MeV : β ~ 0.6 : ω_{max} ~ 60 MeV

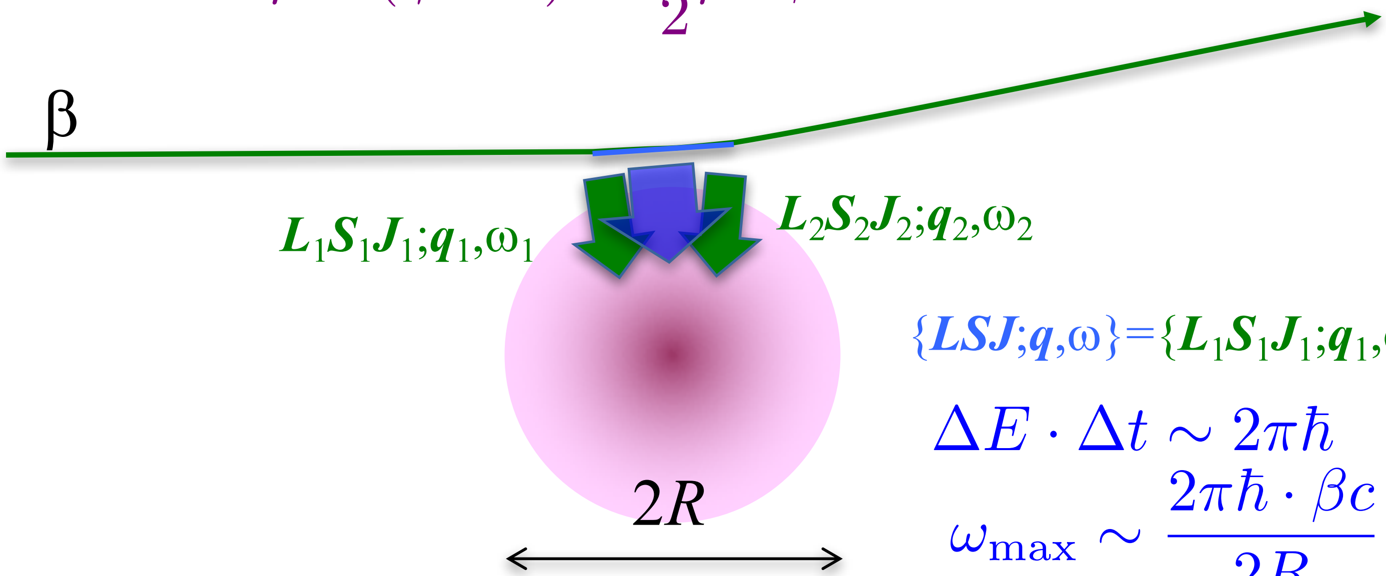


$$O(lsj\tau;\xi)\left|E_i J_i \pi_i T_i;\xi_i\right\rangle = \sum_f\!\!\!\!\!\!\int M_{if}(E_f)\left|E_f J_f \pi_f T_f;\xi_f\right\rangle$$ Response

$\left|M_{if}(E_f)\right|^2$: Energy Spectrum

Two step?

$$\omega \ll \mu c^2 (\gamma - 1) \simeq \frac{1}{2}\mu c^2 \beta^2$$



$$\{LSJ;q,\omega\}=\{L_1S_1J_1;q_1,\omega_1\} \oplus \{L_2S_2J_2;q_2,\omega_2\}$$

$$\Delta E \cdot \Delta t \sim 2\pi\hbar$$

$$\omega_{\rm max} \sim \frac{2\pi\hbar \cdot \beta c}{2R} \simeq 100\beta \ \ \text{MeV}$$

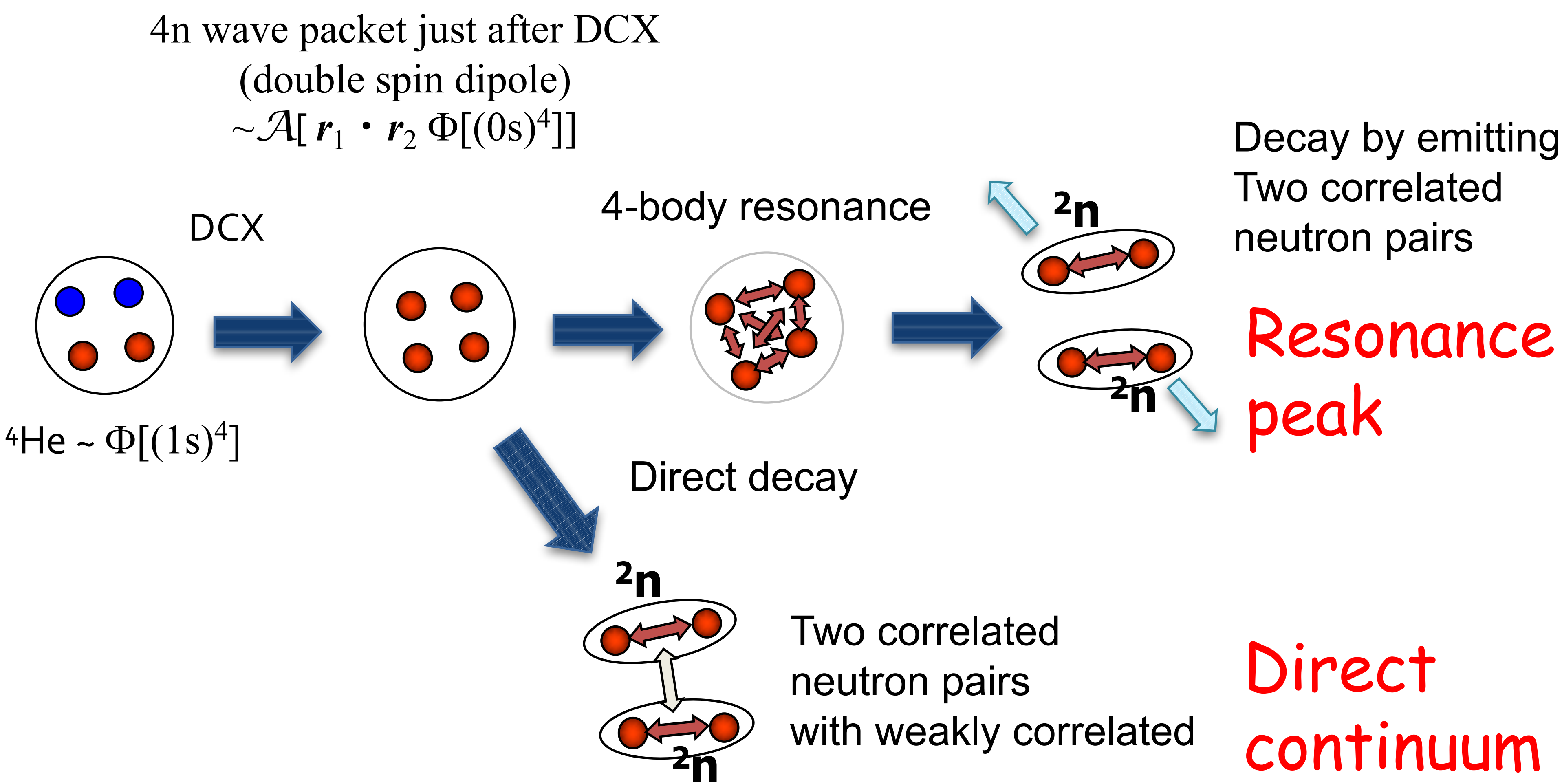
$$\Delta t = \Delta t_1 + \Delta t_2$$

"Intermediate state": Not energy eigen state

~ wave packet consists of "eigen states" over 200β MeV

~ closure approximation ~ almost one-step

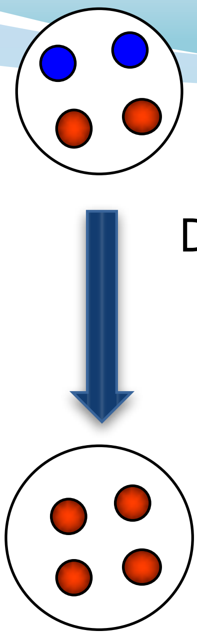
Picture of ^{4}He DCX reaction @ 200 A MeV



Direct Part

Continuum spectrum with n-n FSI

c.f.
L.V. Grigorenko, N.K. Timofeyuk, M.V. Zhukov, Eur. Phys. J. A 19, 187 (2004)



^{4}He ~ $\Phi[(0s)^4]$

$q << 200$ MeV/c

4n wave packet just after DCX
$\Phi_0 \sim \boldsymbol{r}_1 \cdot \boldsymbol{r}_2\, \Phi[(0s)^4]$

$$\mathcal{A}\Phi_0(\boldsymbol{r}_{12}, \boldsymbol{r}_{34}, \boldsymbol{r}_\alpha) \sim$$

$$\left[\left(\frac{r_{12}^2}{a^2} - \frac{3}{2}\right) - \left(\frac{r^2}{a^2} - \frac{3}{4}\right)\right] \exp\left[-\frac{r^2}{a^2} - \frac{r_{12}^2}{2a^2} - \frac{r_{34}^2}{2a^2}\right] \chi(1,2)\chi(3,4)$$

$$\left[\left(\frac{r_\alpha^2}{(a/\sqrt{2})^2} - \frac{3}{2}\right) - \frac{2\vec{r}_{12}\cdot\vec{r}_{34}}{a^2}\right] \exp\left[-\frac{r_\alpha^2}{a^2} - \frac{r_{12}^2}{2a^2} - \frac{r_{34}^2}{2a^2}\right] \chi(1,3)\chi(4,2)$$

$$\left[\left(\frac{r_\alpha^2}{(a/\sqrt{2})^2} - \frac{3}{2}\right) + \frac{2\vec{r}_{12}\cdot\vec{r}_{34}}{a^2}\right] \exp\left[-\frac{r_\alpha^2}{a^2} - \frac{r_{12}^2}{2a^2} - \frac{r_{34}^2}{2a^2}\right] \chi(1,4)\chi(2,3)$$

$$\vec{r}_\alpha = \frac{\vec{r}_1 + \vec{r}_2}{2} - \frac{\vec{r}_3 + \vec{r}_4}{2} \qquad \chi(i,j) = \frac{1}{\sqrt{2}}(\uparrow(i)\downarrow(j) - \downarrow(i)\uparrow(j))$$

Fourier Transform: $(\boldsymbol{r}_{12}, \boldsymbol{r}_{34}, \boldsymbol{r}_\alpha) \rightarrow (\boldsymbol{k}_{12}, \boldsymbol{k}_{34}, \boldsymbol{k})$

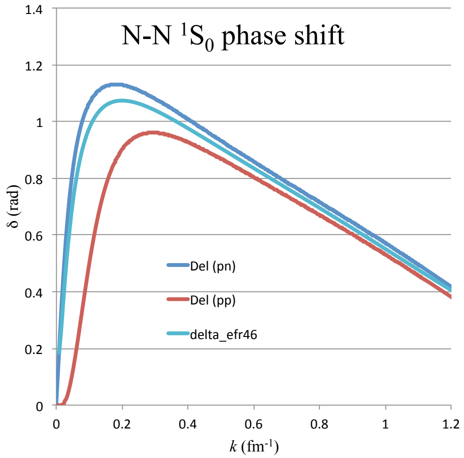
$$\int \left|\mathcal{A}\tilde{\Phi}_0\right|^2 d^3k\, d^3k_{12}\, d^3k_{34}\, \delta(E - \epsilon - \epsilon_{12} - \epsilon_{34}) \propto X^{11/2} \exp(-X)$$

Peak at $X = 11/2$; $E \sim 60$ MeV

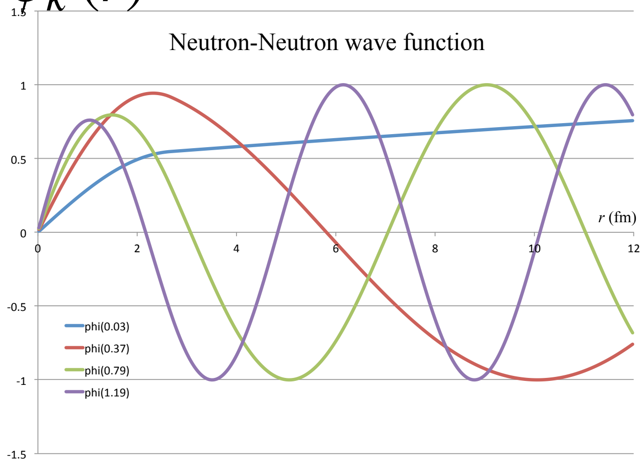
$$X = E/\epsilon_a \qquad \epsilon_a = \frac{\hbar^2}{m_N a^2} = 11\,\text{MeV}$$

NN FSI

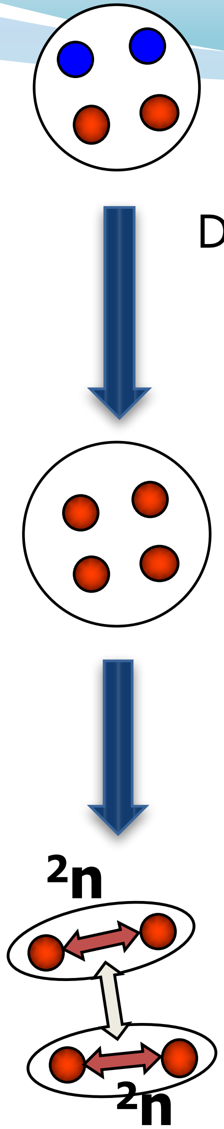
Continuum spectrum with n-n FSI



$\phi_k(r)$



^{4}He ~ $\Phi[(0s)^4]$



DCX

$q > 15$ MeV/c

4n wave packet just after DCX
$\Phi_0 \sim \boldsymbol{r}_1 \cdot \boldsymbol{r}_2\, \Phi[(0s)^4]$

Two correlated neutron pairs with weakly correlated

Density of State

$$D_{ns}(\epsilon_{\rm nn}) = \frac{\left|\hat{A}_{ns}(k)\right|^2}{k} \quad (\text{for } n = 1,2) \quad ; \quad \epsilon_{\rm nn} = \frac{\hbar^2 k^2}{m_{\rm N}}$$

$$\hat{A}_{1s}(k) = \int_0^\infty dr\, r\, \psi_{1s}(r)\, \phi_k(r) = 2\left(\frac{1}{\sqrt{\pi}a^3}\right)^{1/2} k\, A_{1s}(k)$$

$$\hat{A}_{2s}(k) = \int_0^\infty dr\, r\, \psi_{2s}(r)\, \phi_k(r) = 2\sqrt{\frac{2}{3}}\left(\frac{1}{\sqrt{\pi}a^3}\right)^{1/2} k\, A_{2s}(k)$$

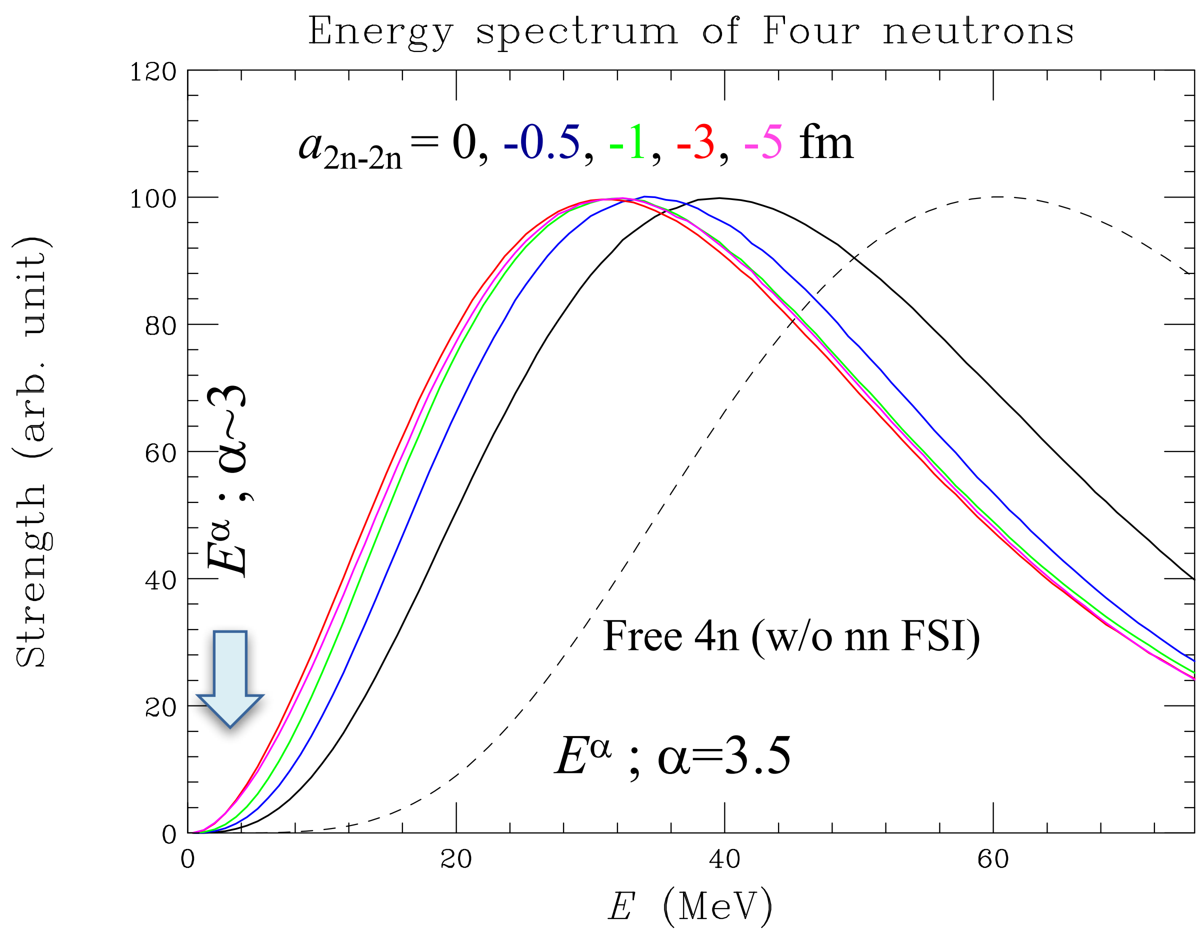
Expand $\mathcal{A}\Phi_0$ with correlated n-n scattering wave $\phi_k(r)$
$A(k)$'s are used instead of Fourier component

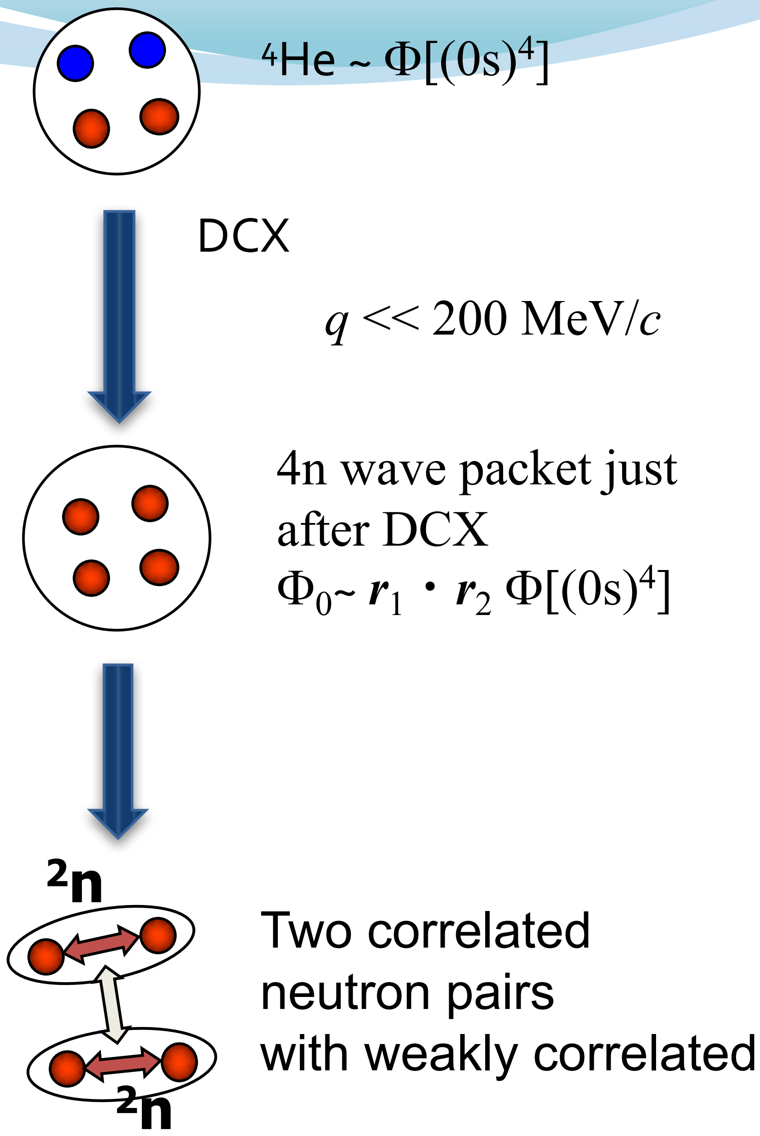
Direct Part

Continuum spectrum with n-n FSI

c.f.

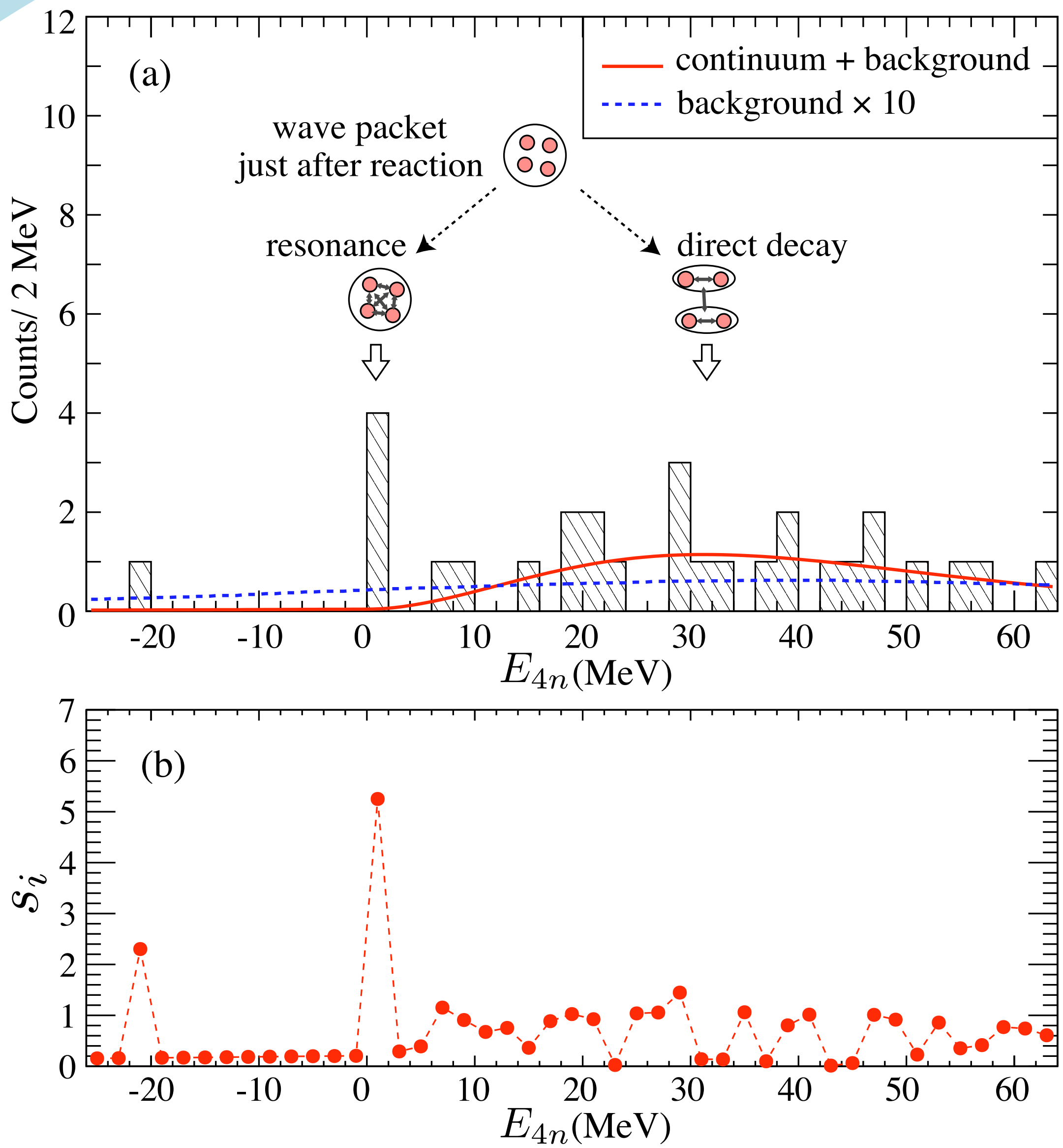
L.V. Grigorenko, N.K. Timofeyuk, M.V. Zhukov, Eur. Phys. J. A 19, 187 (2004)





Correlation is taking into account for 2n-2n relative motion by using scattering length

Fit with direct component & BG



Energy spectrum is expressed by the continuum from the direct decay and (small) experimental background except for four events at $0<E_{4n}<2$ MeV

The Four events suggest a possible resonance at

0.83 ± 0.65(stat.) $\pm$ 1.25(sys.) MeV

with width narrower than 2.6 MeV (FWHM). [4.9σ significance]

Integ. cross section $\theta_{cm}< 5.4$deg:

$3.8\ ^{+2.9}_{-1.8}$ nb

- likelihood ratio test

$$\chi^2_\lambda = -2\ln\left[L(\boldsymbol{y};\boldsymbol{n})/L(\boldsymbol{n};\boldsymbol{n})\right]$$

- Significance:

$$s_i = \sqrt{2[y_i - n_i + n_i \ln(n_i/y_i)]}$$

n_i : num. of events in the i-th bin

y_i : trial function in the i-th bin

- Look Elsewhere Effect

$$\mu^n e^{-\mu}/n! \simeq 10^{-6} \quad \text{for} \quad \mu = 0.07,\ n = 4$$