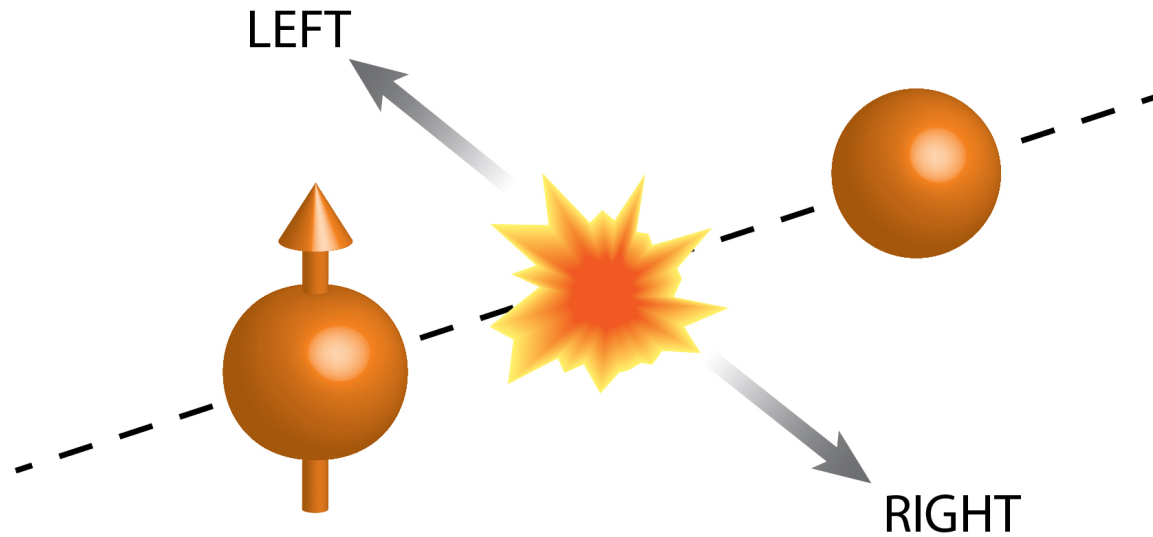


# Very Forward Neutron Asymmetry in Transversely Polarized p+A collision

Itaru Nakagawa

on behalf of PHENIX Collaboration

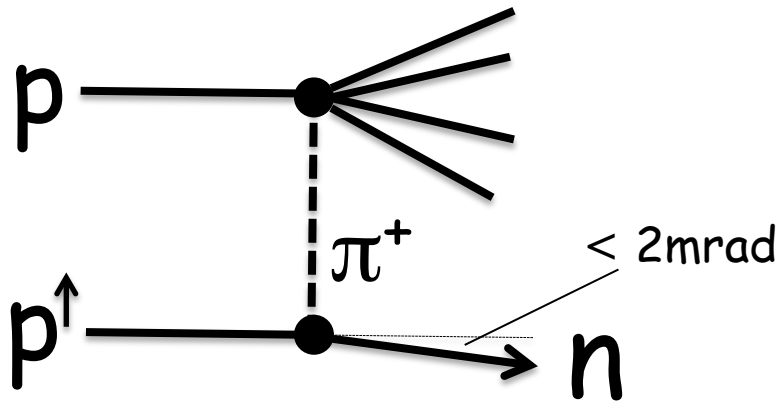


Forward  $\pi^0$  and very forward neutron

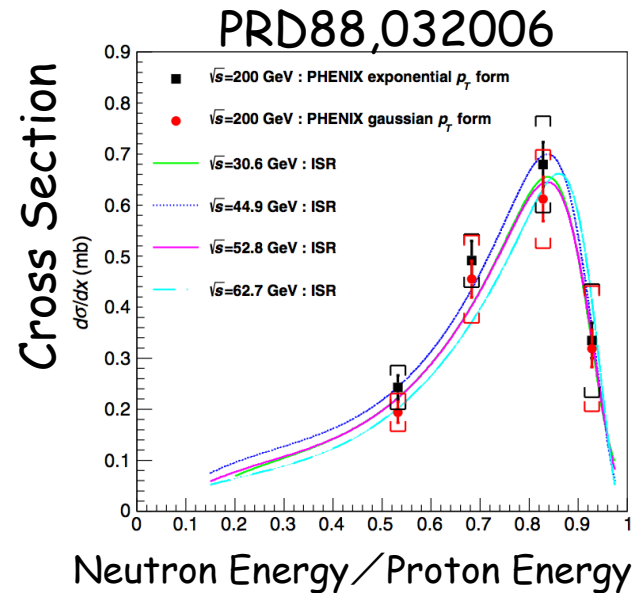
# TRANSVERSE SINGLE SPIN ASYMMETRY

# Production Mechanism of Forward Neutron

## Cross Section

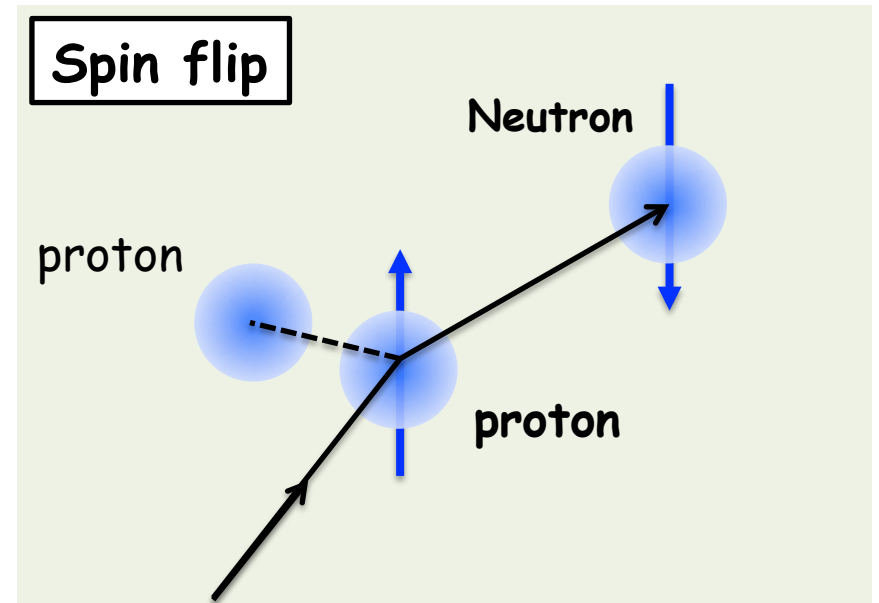
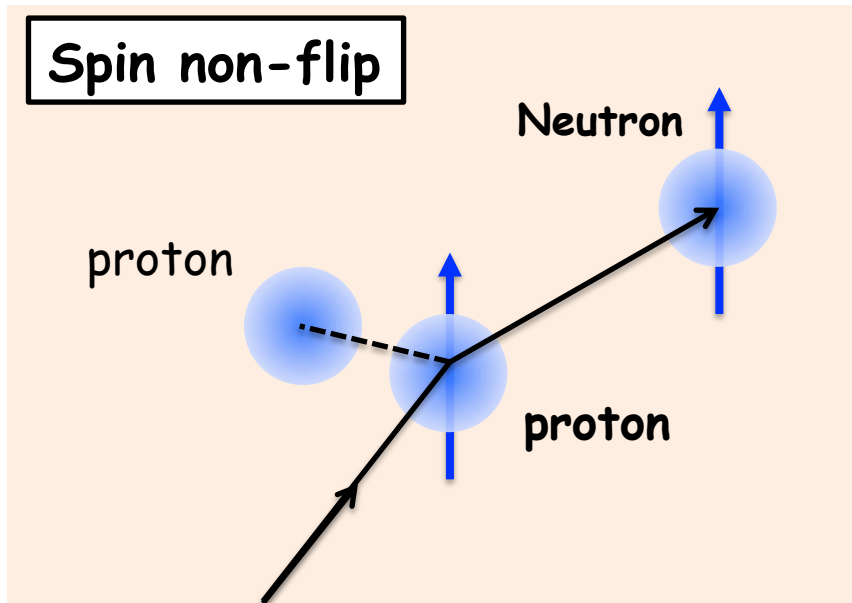


Momentum Transfer  $\sim 100\text{MeV}$



Well Explained by  
**One-Pion Exchange**

# $p^\uparrow + p$ Forward Neutron $A_N$



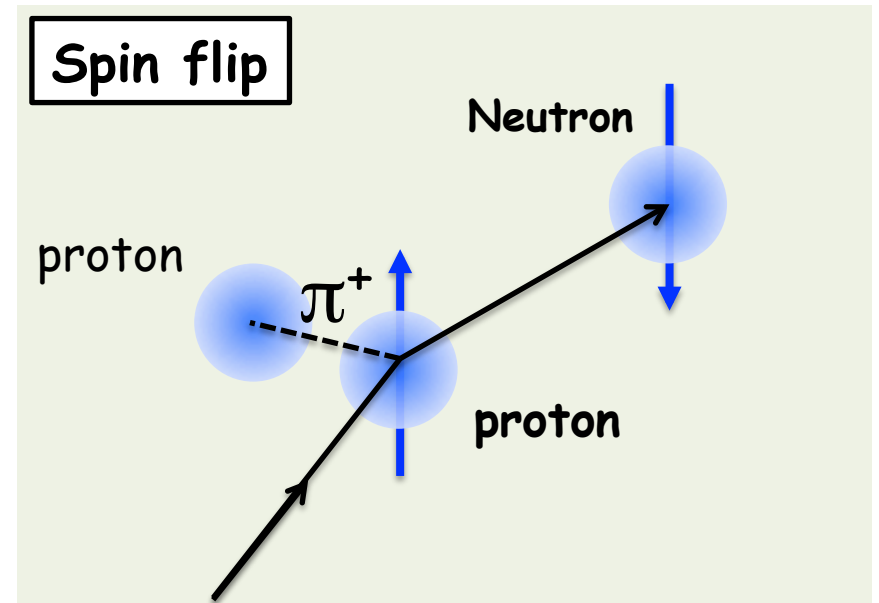
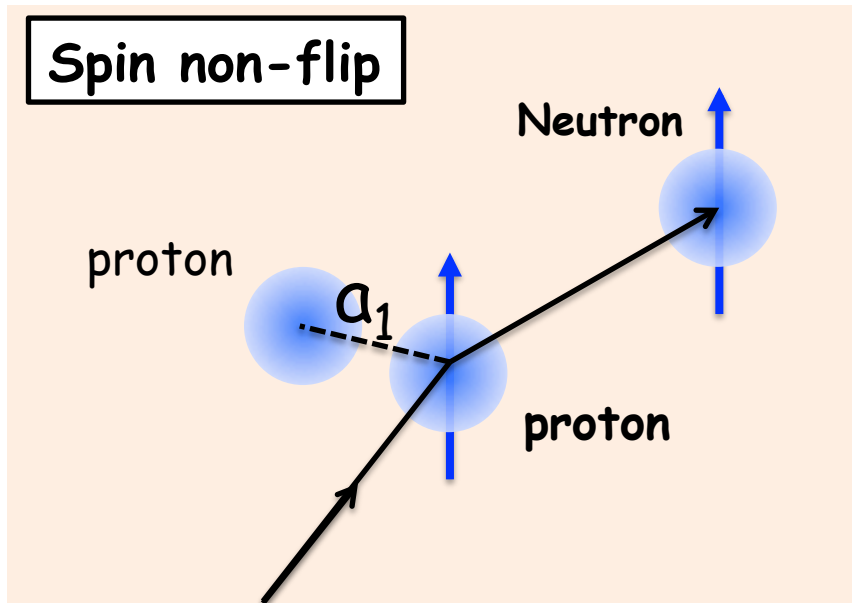
$$A_N \approx \frac{\left( \phi_{non-flip}^* \phi_{flip} \sin \delta \right)}{\left| \phi_{non-flip} \right|^2 + \left| \phi_{flip} \right|^2}$$

$\delta$  : phase shift

Unpolarized Cross Section



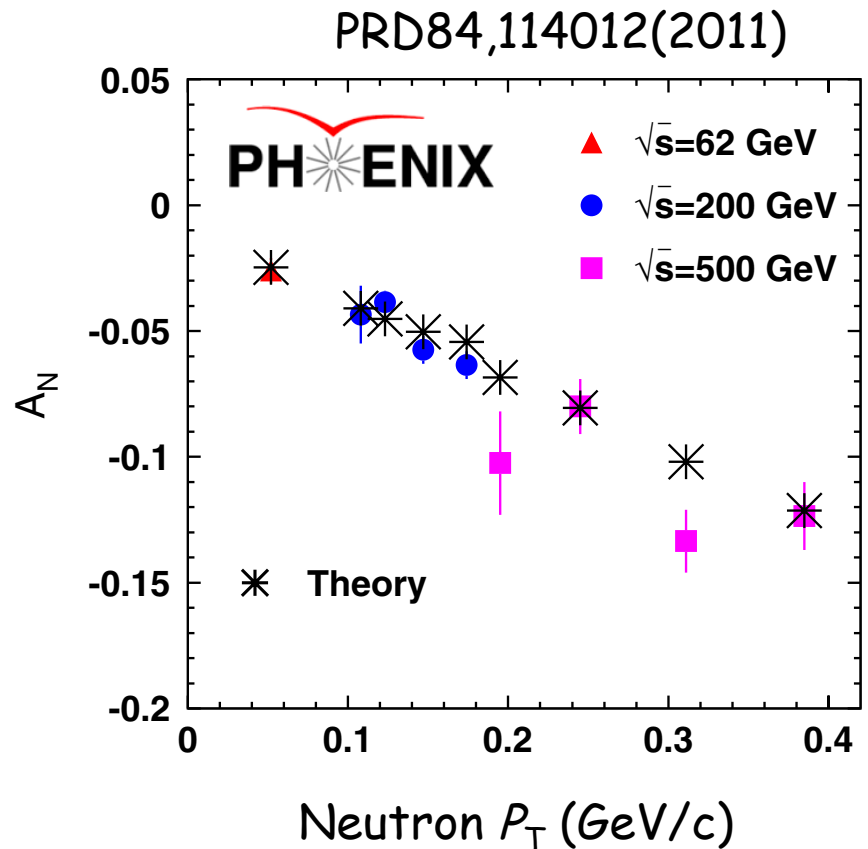
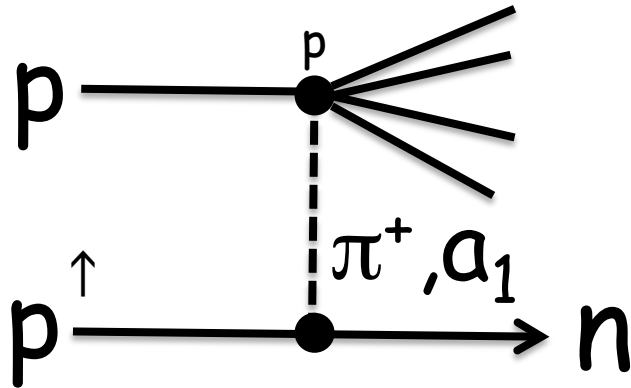
# $p^\uparrow + p$ Forward Neutron $A_N$



$$A_N \approx \frac{\left( \phi_{non-flip}^* \phi_{flip} \sin \delta \right)}{\left| \phi_{non-flip} \right|^2 + \left| \phi_{flip} \right|^2}$$

$\delta$  : phase shift

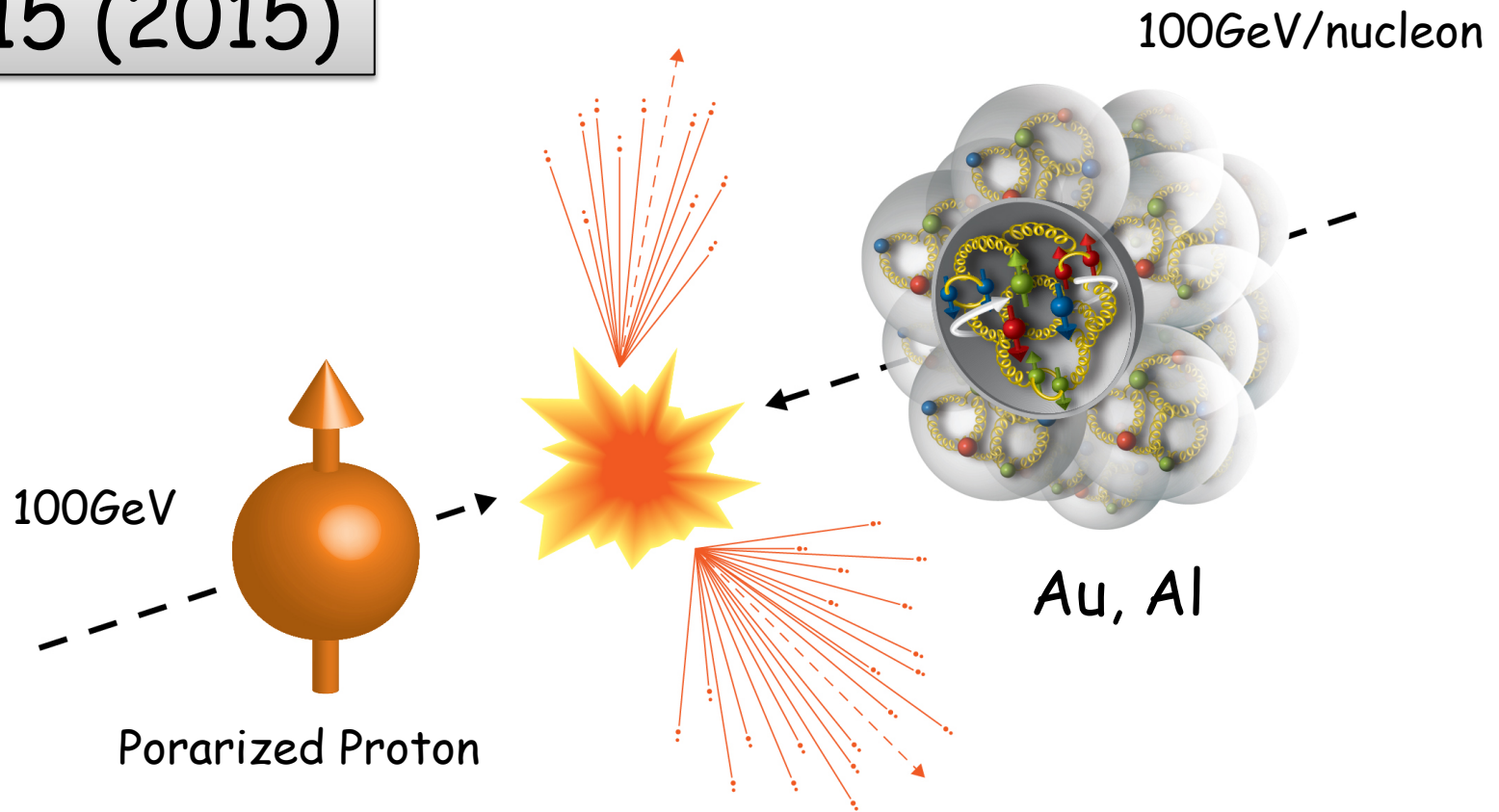
# $p^\uparrow + p$ Forward Neutron $A_N$



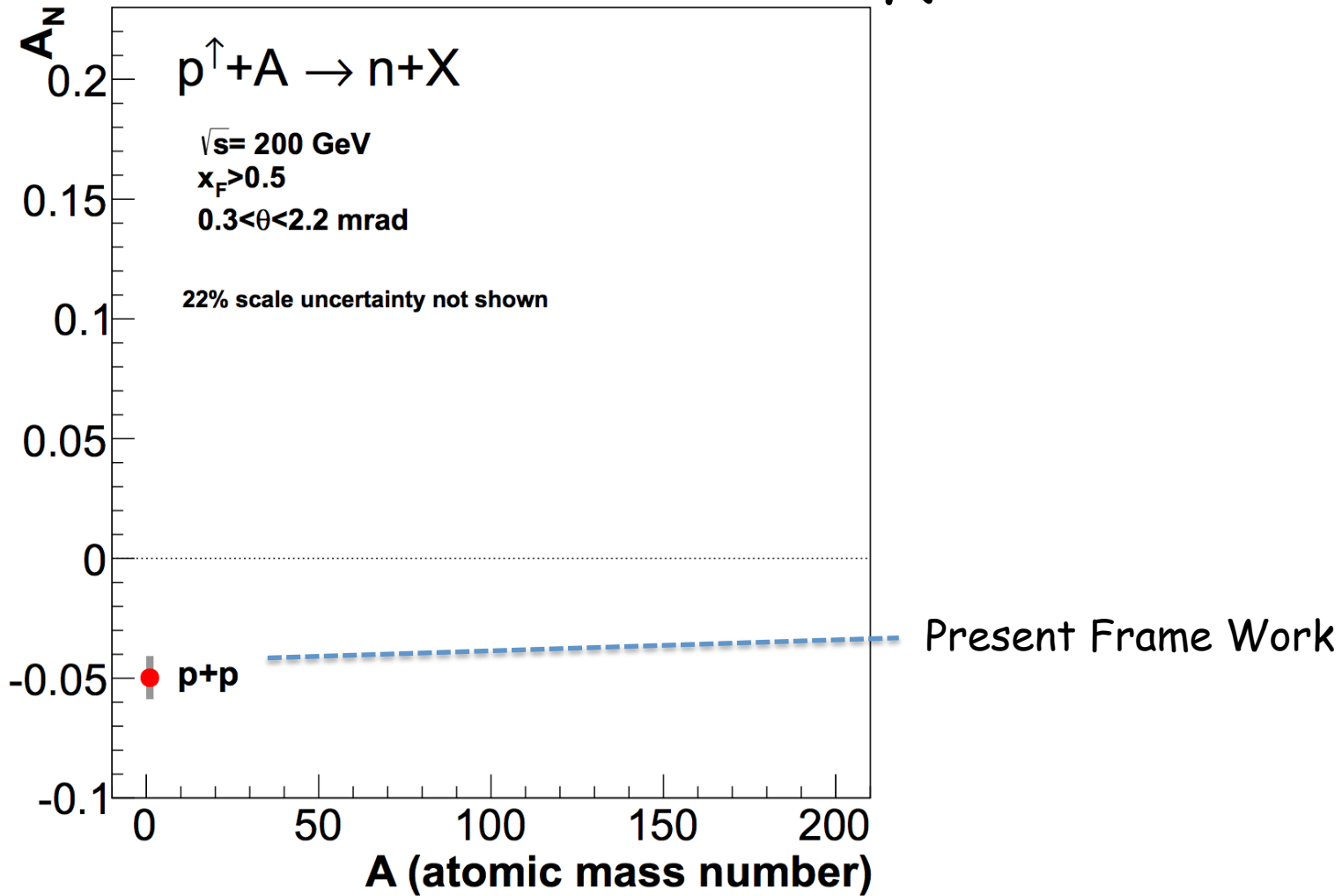
Data are well reproduced by the interference between  $\pi$  and  $a_1$  Reggeon

# The First Time Ever High Energy $p^\uparrow + A$ Collisions

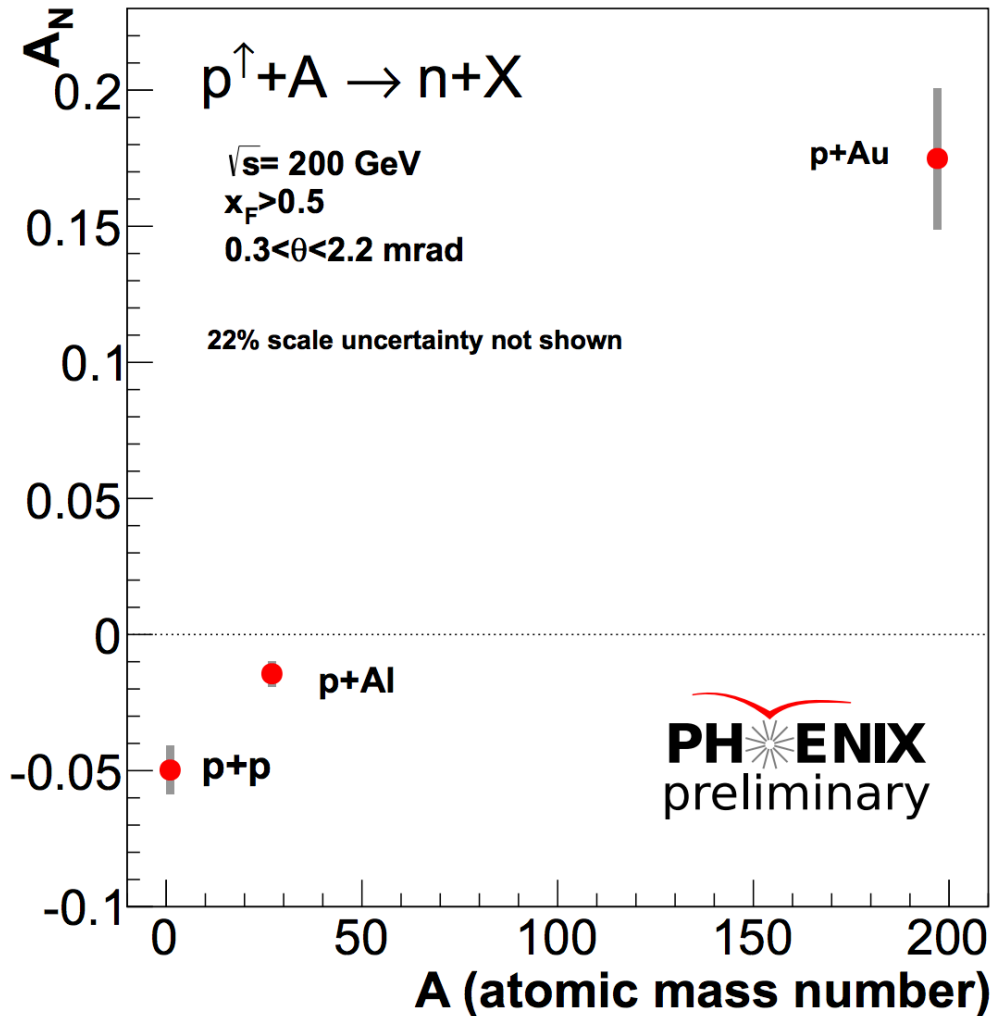
Run15 (2015)



# A-Dependent $A_N$ (inclusive)

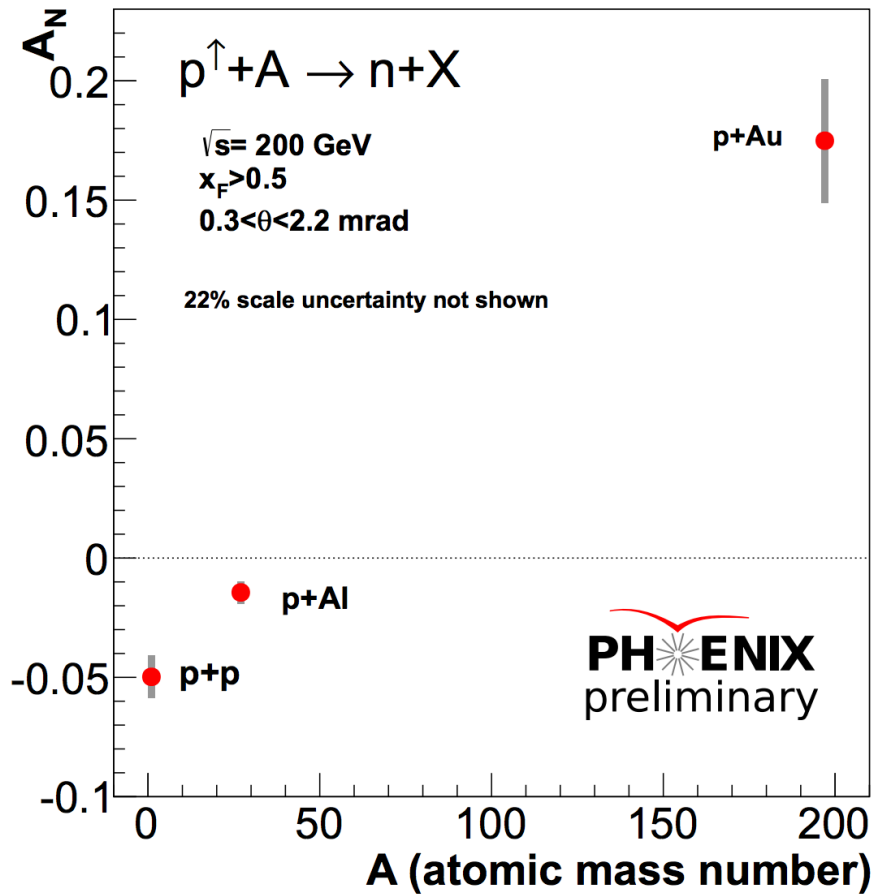


# A-Dependent $A_N$ (inclusive)



Analysis by  
Minjung Kim (SNU/RIKEN)

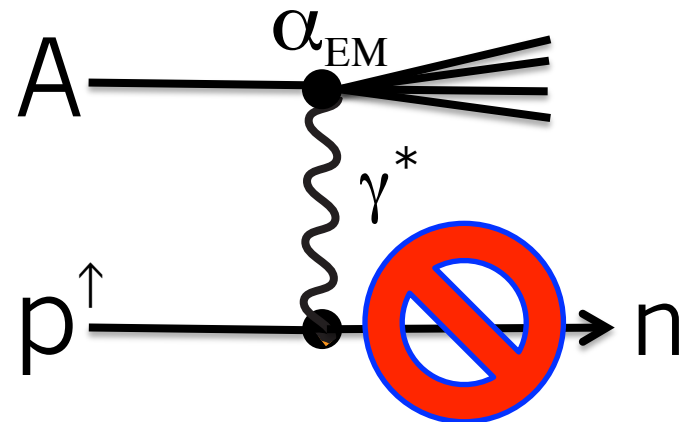
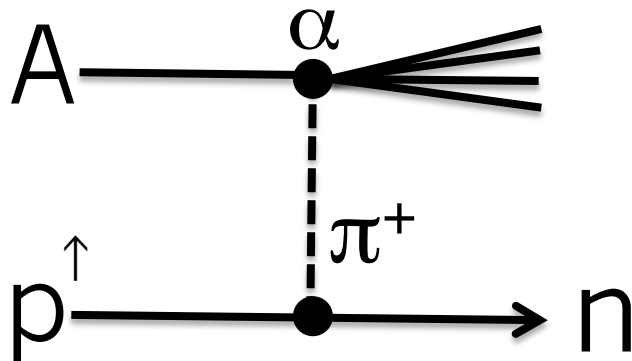
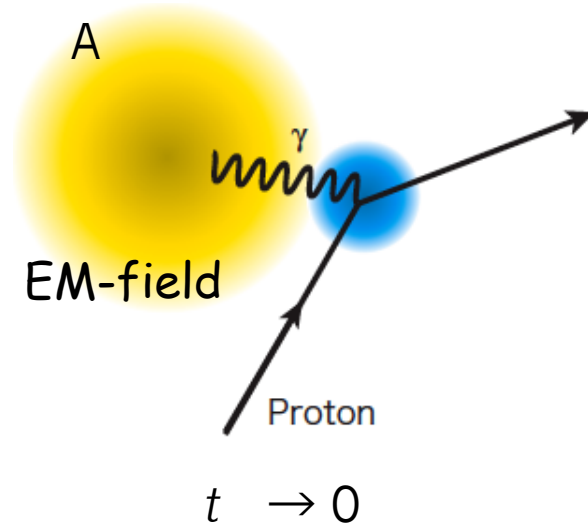
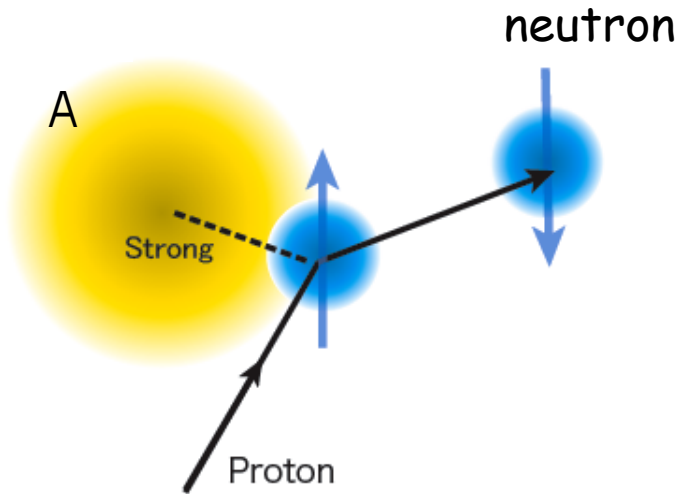
# What is going on ?



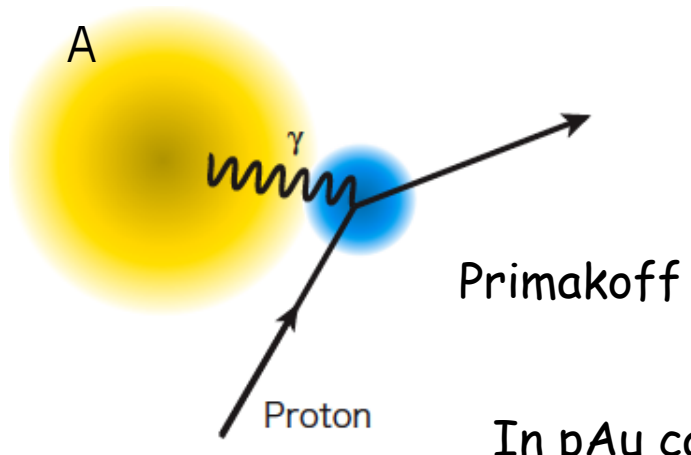
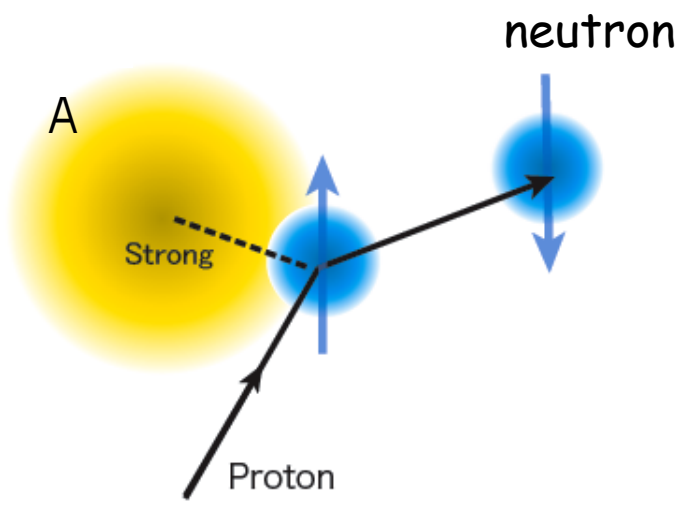
	# of proton	# of neutron
p	1	0
Al	13	14
Au	79	118

- Isospin Symmetry
- Surface Structure of Nucleus
- QED Process
- Gluon Saturation
- Else...

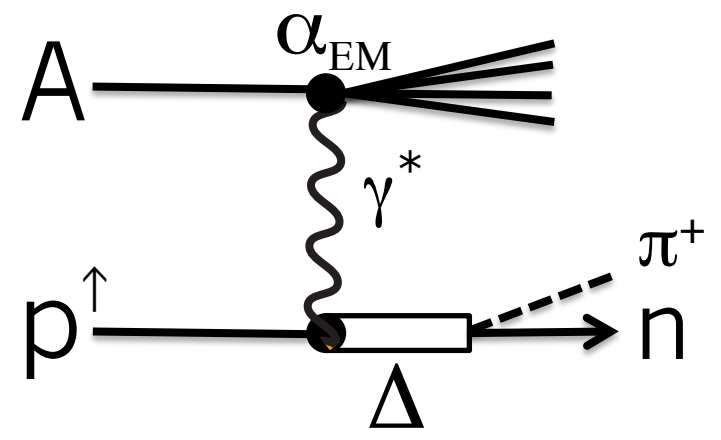
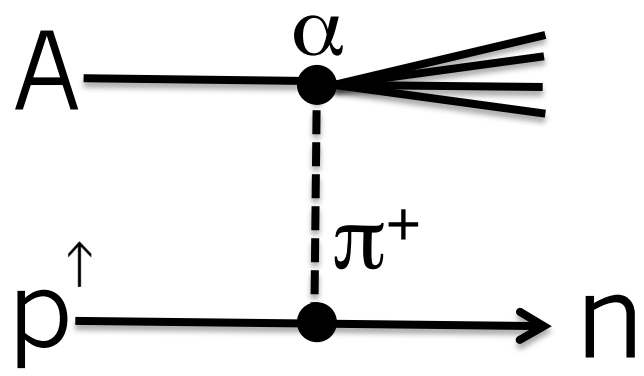
# Primakoff Effect Electro-Magnetic



# Primakoff Effect (Electro-Magnetic)

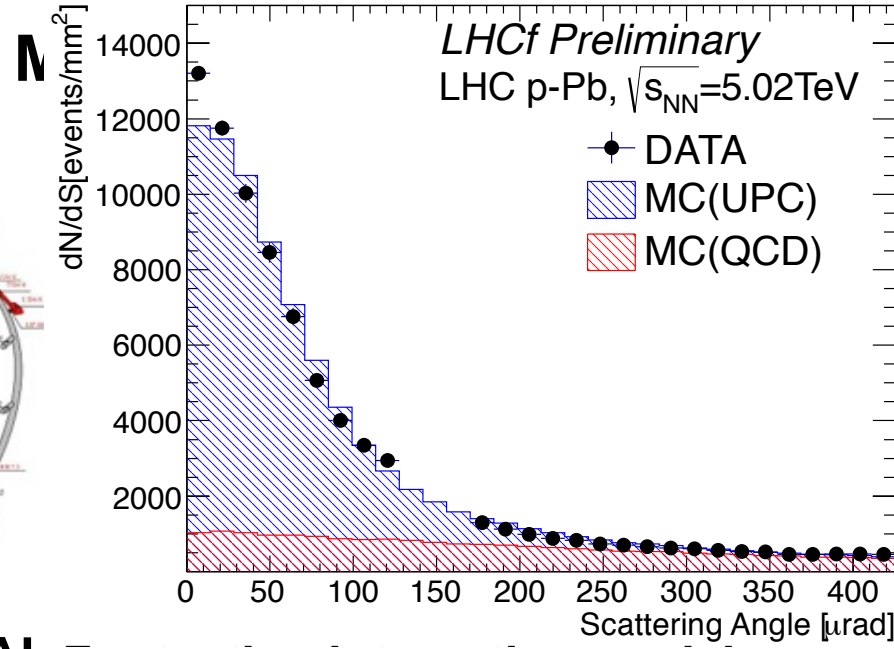
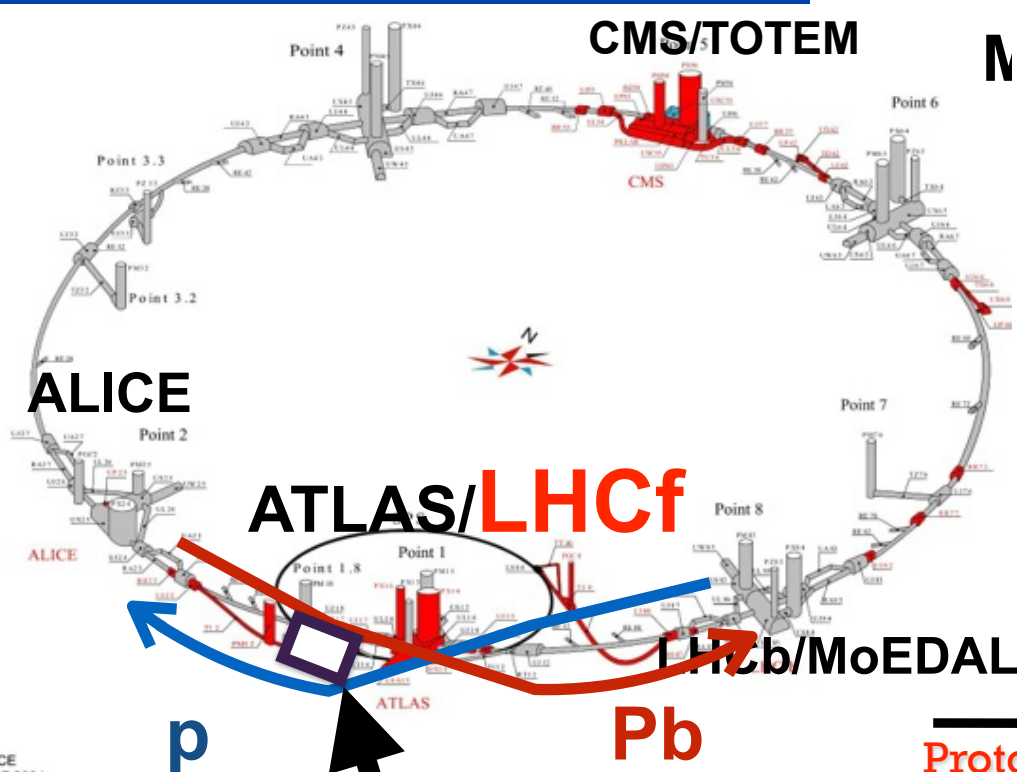


In pAu case  
 $t \rightarrow 0$   
 $Z \cdot \alpha_{EM} \sim 79/137$

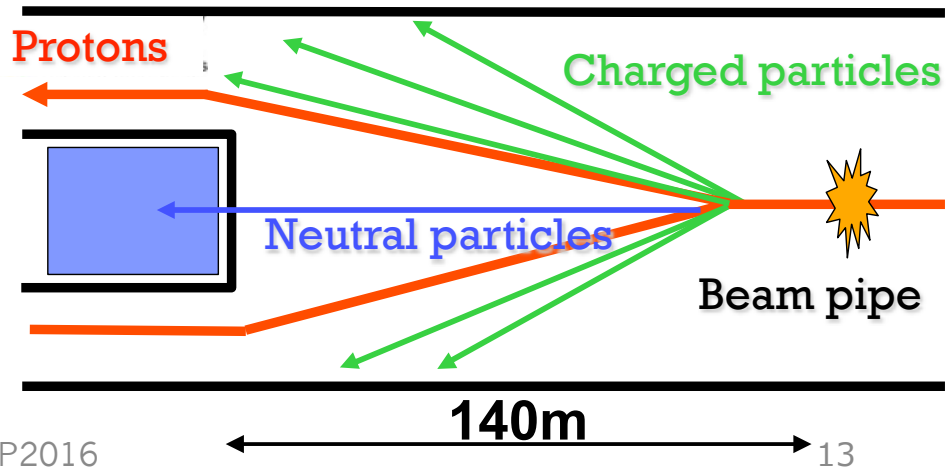
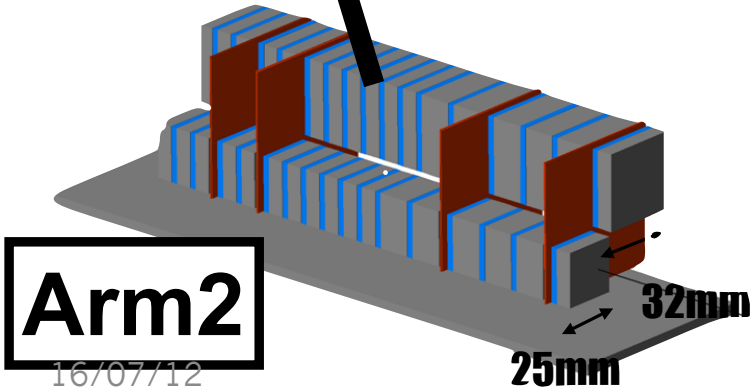




# The LHCf experiment



TS-CE  
06.07.2004



ICNFP2016

# UPC Monte Carlo

Eur. Phys. J. C (2015) 75:614  
DOI 10.1140/epjc/s10052-015-3848-0

THE EUROPEAN  
PHYSICAL JOURNAL C



Special Article - Tools for Experiment and Theory

## Forward hadron production in ultra-peripheral proton-heavy-ion collisions at the LHC and RHIC

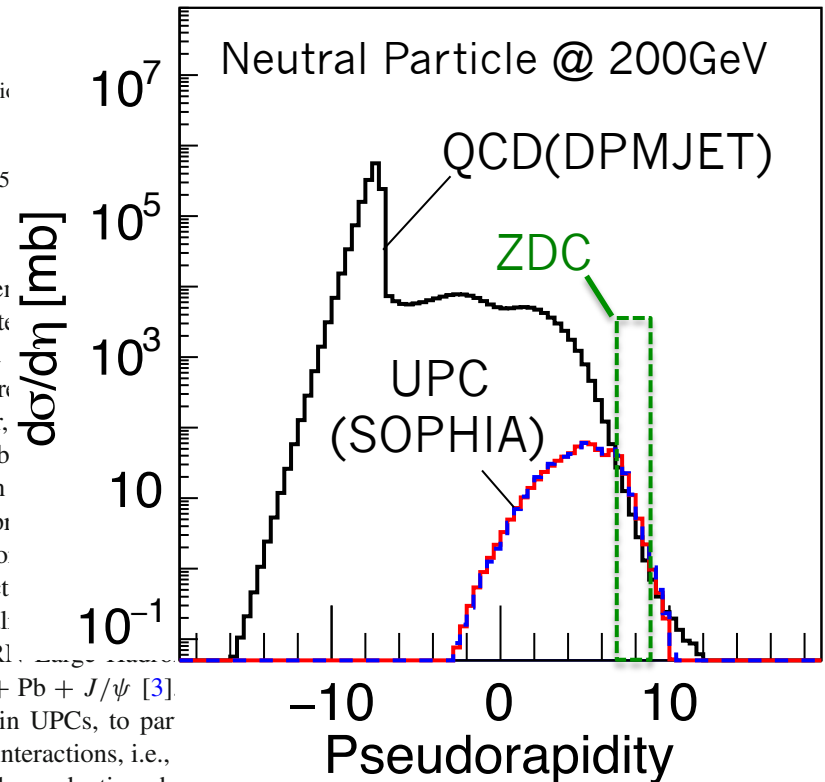
Gaku Mitsuka<sup>a</sup>

Università degli Studi di Firenze and INFN Sezione di Firenze, Via Sansone 1, 50019 Sesto Fi

Received: 26 April 2015 / Accepted: 15 December 2015 / Published online: 26 December 2015  
© The Author(s) 2015. This article is published with open access at Springerlink.com

**Abstract** We present a hadron production study in the forward rapidity region in ultra-peripheral proton-lead ( $p + \text{Pb}$ ) collisions at the LHC and proton-gold ( $p + \text{Au}$ ) collisions at RHIC. The present paper is based on the Monte Carlo simulations of the interactions of a virtual photon emitted by a fast moving nucleus with a proton beam. The simulation consists of two stages: the STARLIGHT event generator simulates the virtual photon flux, which is then coupled to the SOPHIA, DPMJET, and PYTHIA event generators for the simulation of particle production. According to these Monte Carlo simulations, we find large cross sections for ultra-peripheral collisions particle production, especially in the very forward region. We show the rapidity distributions for charged and neutral particles, and the momentum distributions for neutral pions and neutrons at high rapidities. These processes lead to substantial background contributions to the investigations of collective nuclear effects and spin physics. Finally we pre-

virtual photons emitted by the nucleus may anyway interact with the proton beam, usually referred to as ultra-peripheral collisions (UPCs), so far, the gluon distribution function in proton-nucleus collisions can provide information on the parton density in proton-nucleus collisions at the CERL. In this paper, we study the production of  $p + \text{Pb} \rightarrow p + \text{Pb} + J/\psi$  [3], which has been paid, in UPCs, to parton-photon interactions, i.e., less such particle production should be considered as well in the investigation of collective nuclear effects. Because of



Predicts comparable yields between QCD and UPC processes

# Full Description of $A_N$

$$A_N \propto 2 \operatorname{Im} \left\{ \phi_{non-flip}^* \phi_{flip} \right\}$$

$$\phi_{flip} = \phi_{flip}^{had} + \phi_{flip}^{EM}$$

$$\phi_{non-flip} = \phi_{non-flip}^{had} + \phi_{non-flip}^{EM}$$

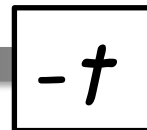
$d_{1\sim 4}$  : relative phase of amplitudes

$$A_N \propto 2 \operatorname{Im} \left( \phi_{non-flip}^{had} + \phi_{non-flip}^{EM} \right) \left( \phi_{flip}^{had*} + \phi_{flip}^{EM*} \right)$$

$$= 2 \left( \phi_{non-flip}^{had} \phi_{flip}^{had} \sin \delta_1 + \phi_{non-flip}^{EM} \phi_{flip}^{had} \sin \delta_2 + \phi_{non-flip}^{had} \phi_{flip}^{EM} \sin \delta_3 + \phi_{non-flip}^{EM} \phi_{flip}^{EM} \sin \delta_4 \right)$$

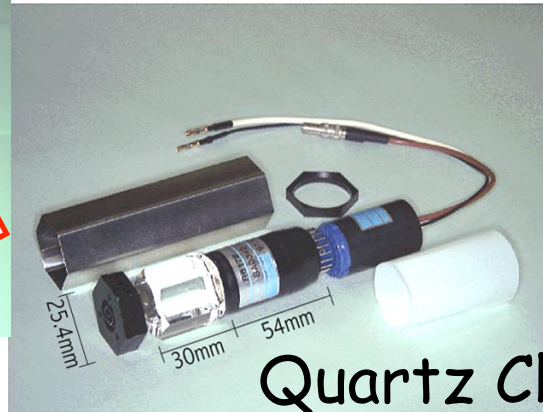
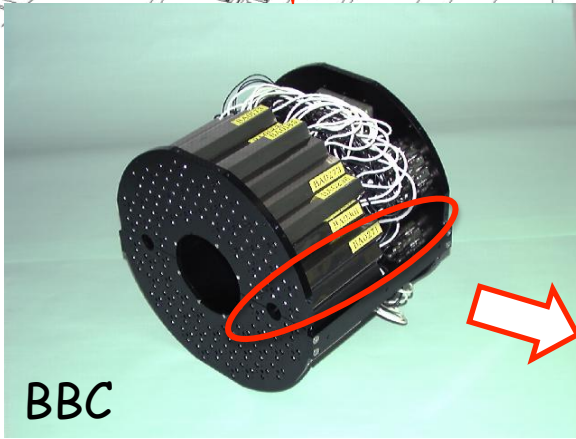
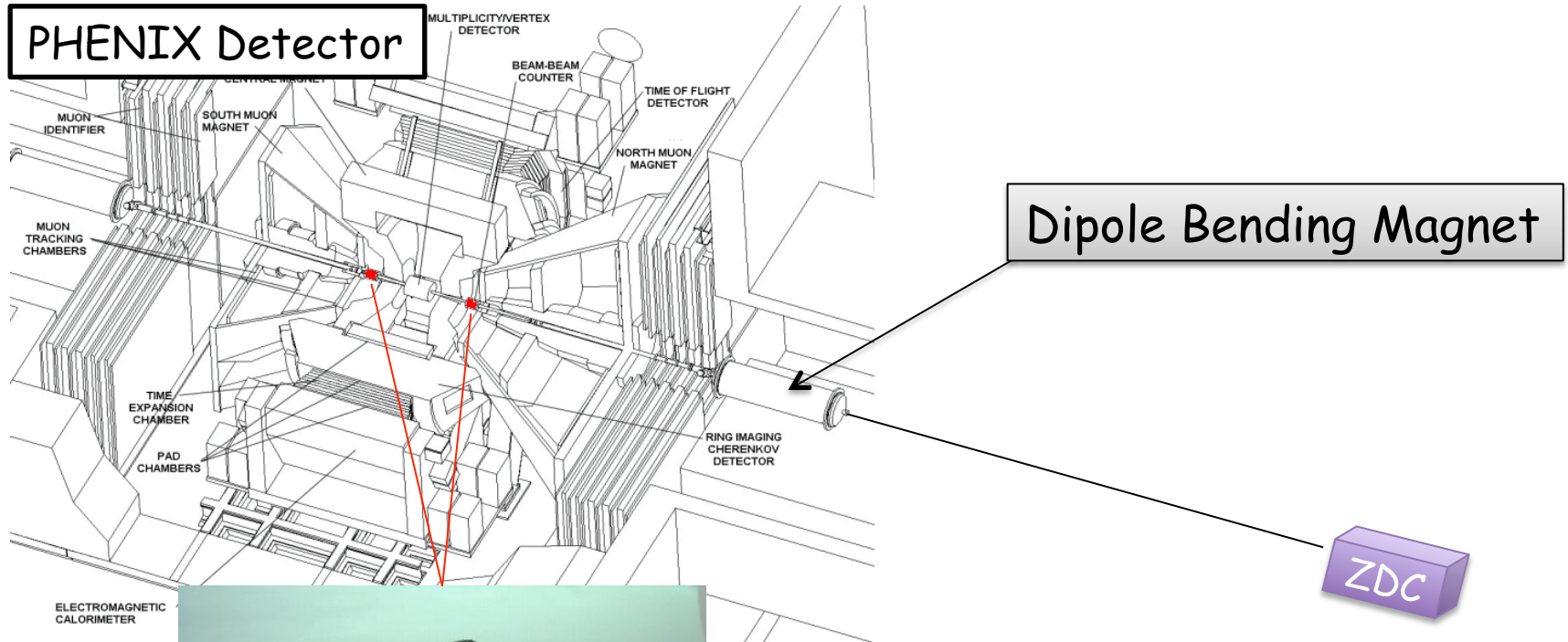
$\rightarrow 0$  (For pp)

large ←



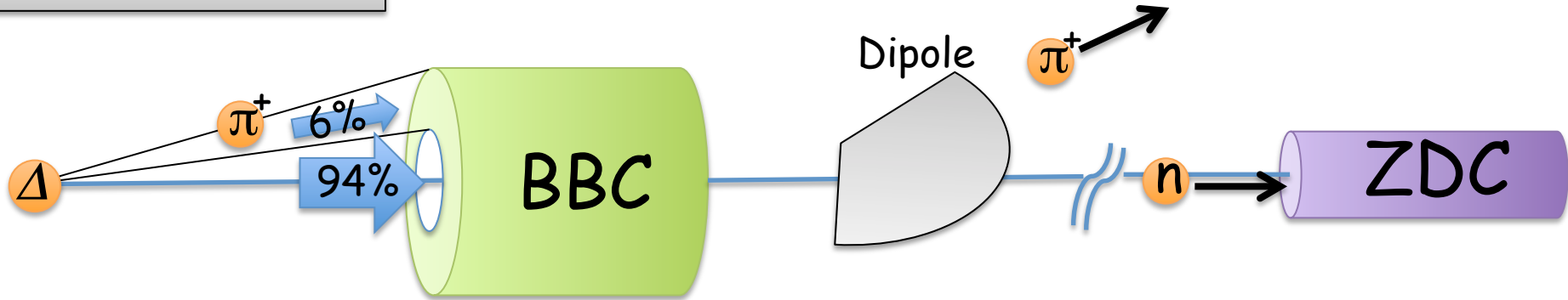
→ small

# Beam-Beam Counter



# Can we identify Primakoff events?

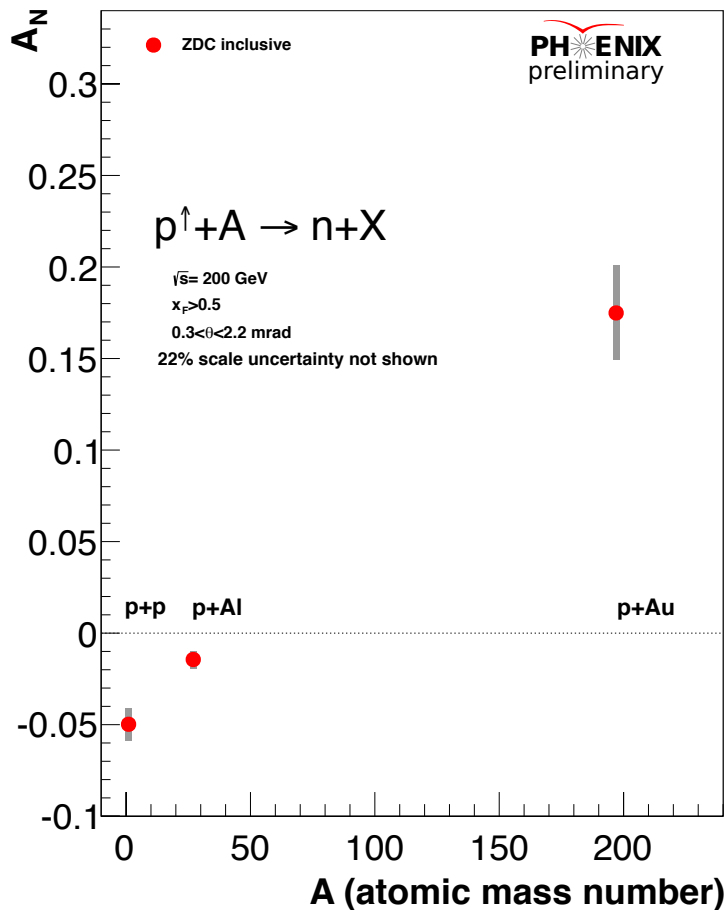
Semi-Inclusive



Primakoff MC : SOPHIA  
G. Mitsuka, Eur. Phys. J.C. (2015) 75:614

BBC	tag	veto
Primakoff	Small	Large

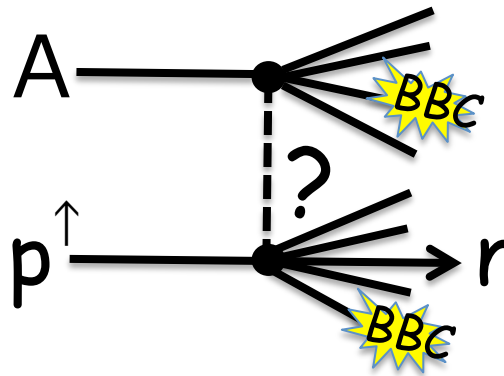
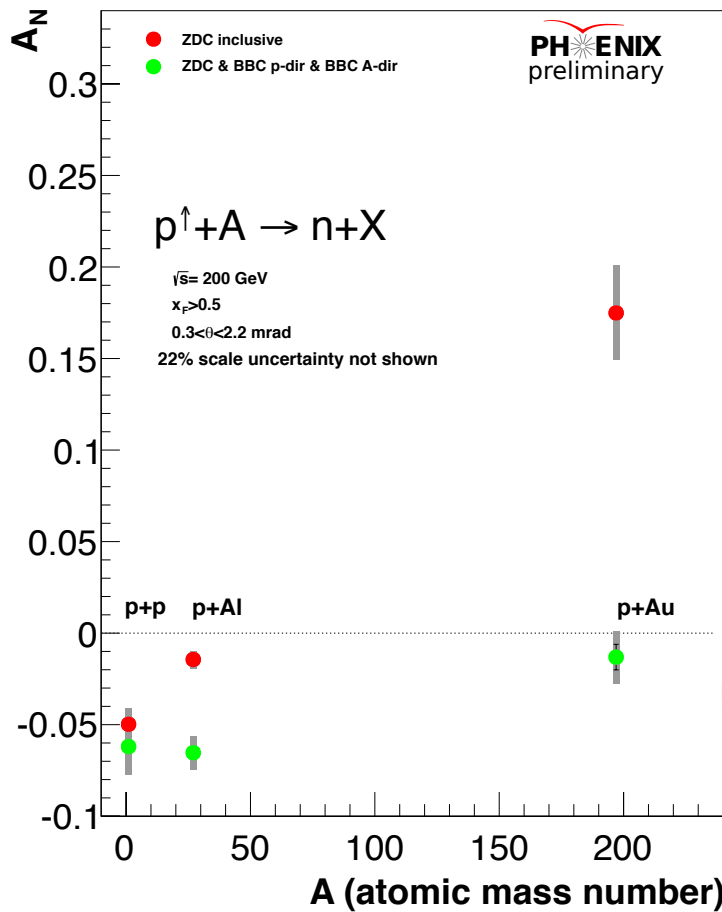
# BBC Tagging and Vetoing



Inclusive

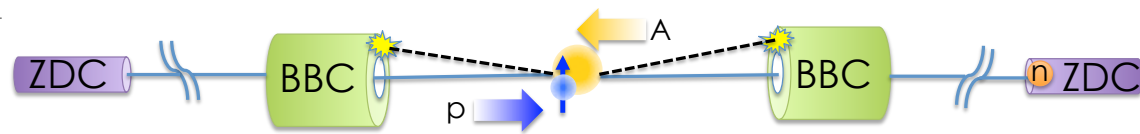
$$A_N \sim \text{had} * \text{had} + \text{had} * EM + EM * \text{had} + EM * EM$$

# BBC Tagging and Vetoing



Non-diffractive

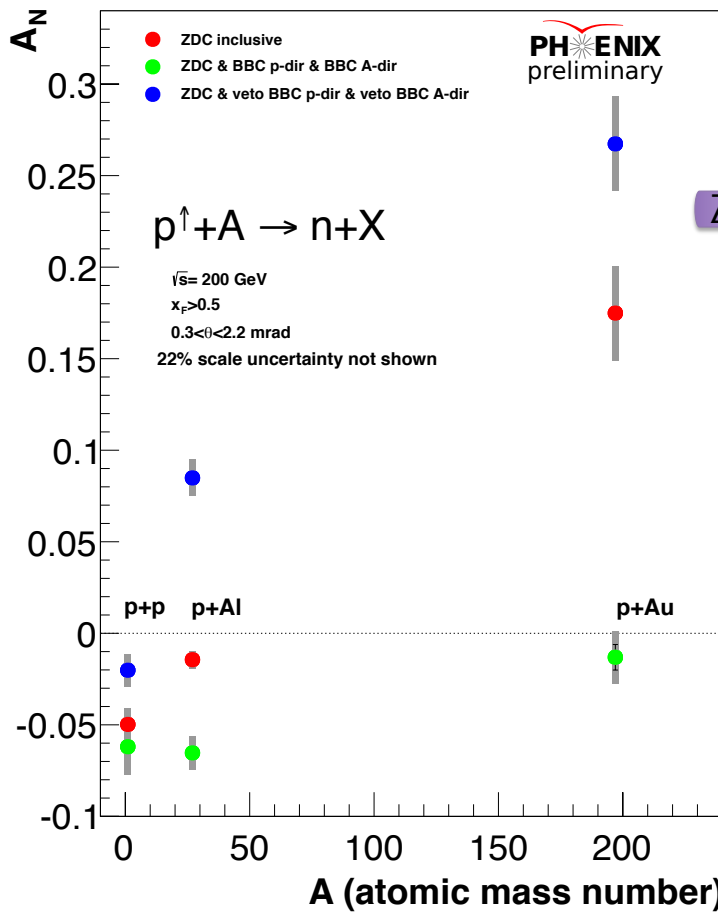
BBC Tagging



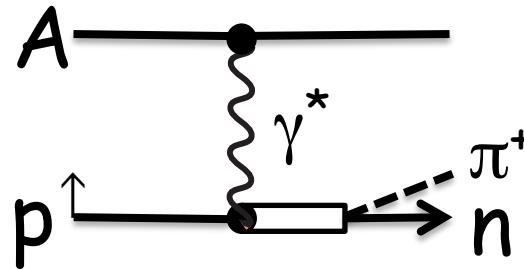
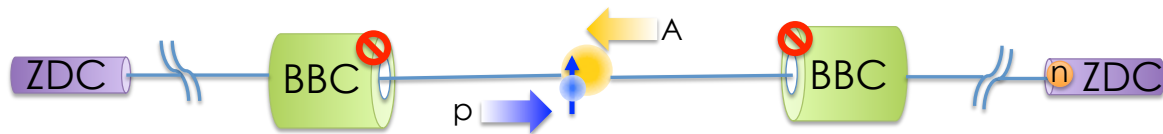
EM suppressed

$$A_N \sim \boxed{had * had} + had * EM + EM * had + EM * EM$$

# BBC Tagging and Vetoing



## BBC Vetoing



Diffractive

EM enhanced

$$A_N \sim had * had + had * EM + EM * had + EM * EM$$

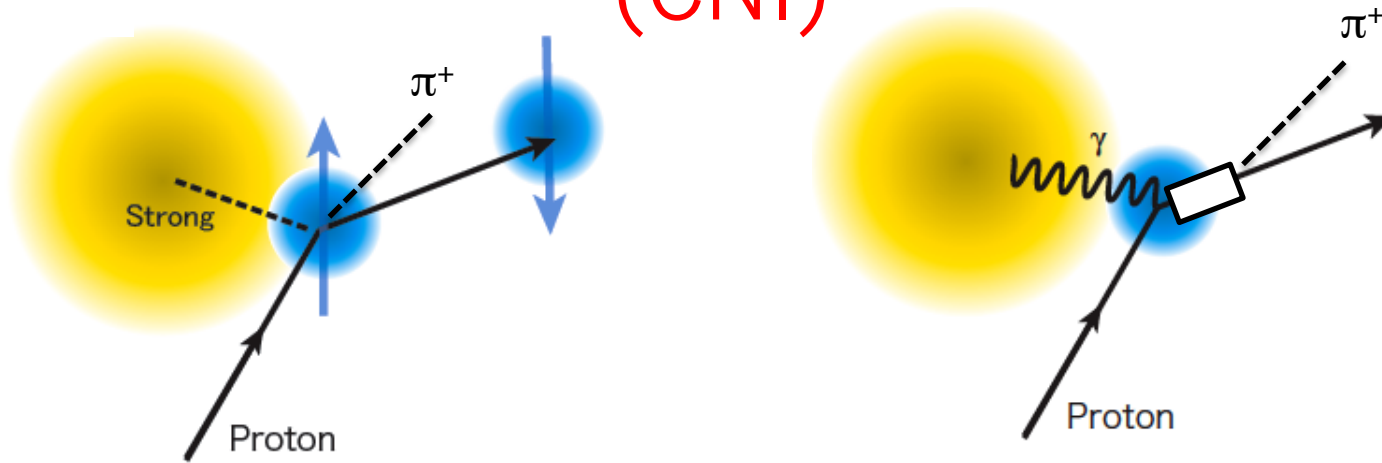


# Fractions

	p+p	p+Al	p+Au
Inclusive	36 %	34 %	13 %
BBC veto	24 %	31 %	67 %
BBC tag	36 %	34 %	13 %

# Coulomb-Nuclear Interference

(CNI)



$$A_N \sim had * had + had * EM + EM * had + EM * EM$$

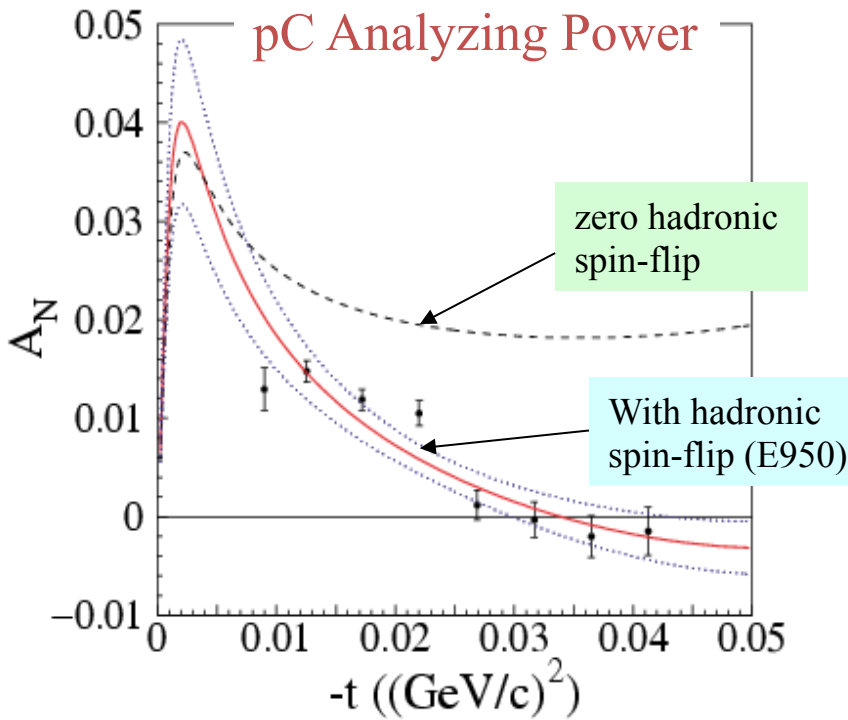
# Elastic $A_N$ at Coulomb Nuclear Interference

$$A_N \sim 2 \text{Im} \left\{ \phi_{non-flip}^{EM*} \phi_{flip}^{had} \sin \delta_2 + \phi_{non-flip}^{had*} \phi_{flip}^{EM} \sin \delta_3 \right\}$$

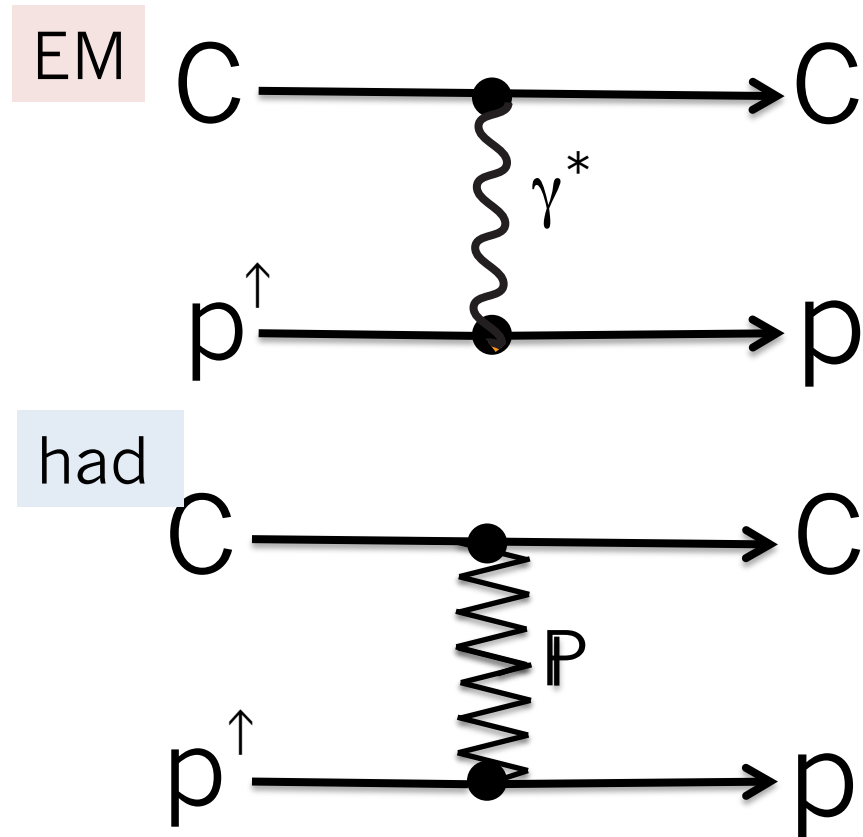
$\propto \alpha s/t$  Pomeron

$\propto \sqrt{\sigma_{pC}^{had}}$   $\propto (\mu_p - 1)$

(High energy & small t limit)

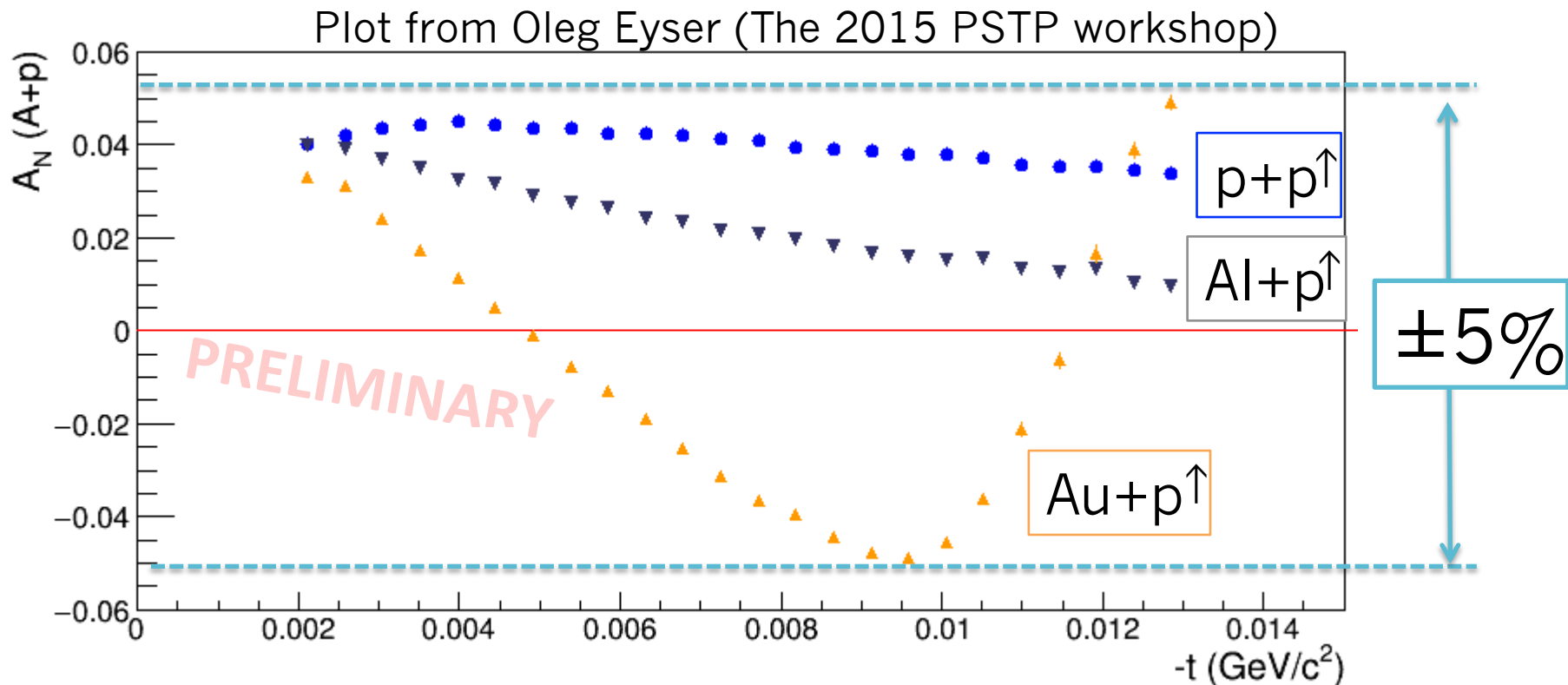


Phys.Rev.Lett.,89,052302(2002)



16/07/12  $E_{beam} = 21.7 \text{ GeV}$

# Run15 Au,Al beam + $p^\uparrow$ target

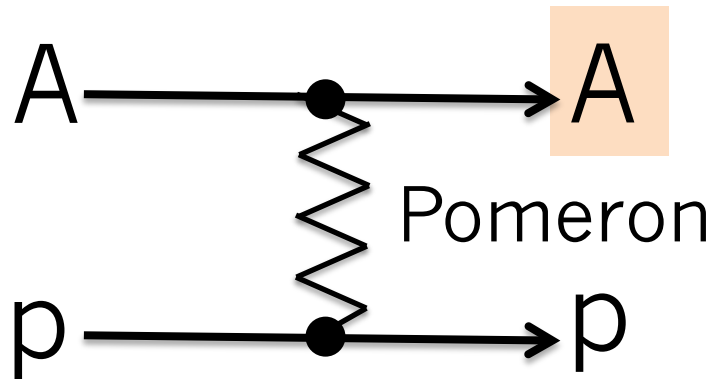


- Strong A-Dependence
- Flips sign of  $A_N$  in  $\text{Au}+p^\uparrow$
- $0.002 < -t < 0.014$  ( $\text{GeV}/c^2$ )

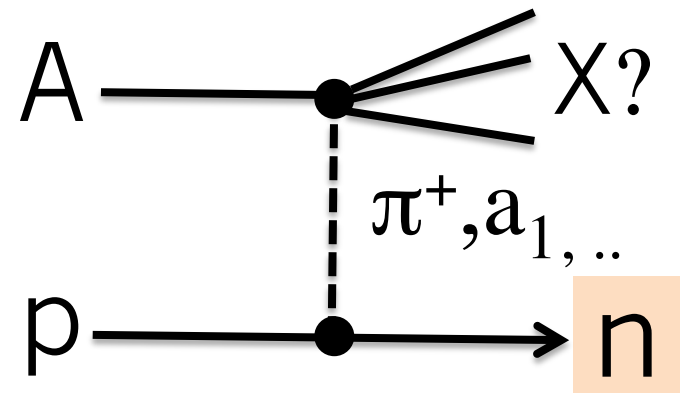
Forward Neutron  
 $0.02 < -t < 0.5$  ( $\text{GeV}/c^2$ )

# Underlying Mechanism Comparison

Polarimeter



Forward n



Elastic

$$\sqrt{s} = 14 \text{ GeV}$$

$$0.002 < -t < 0.014$$

$$\Delta I = 0$$

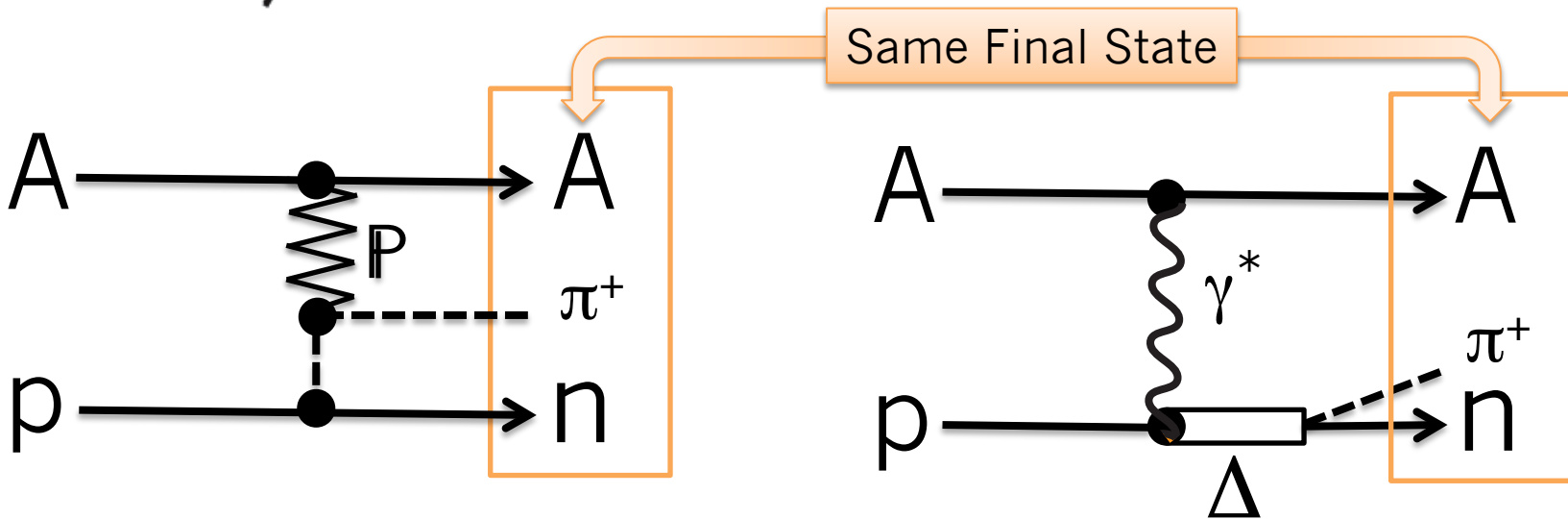
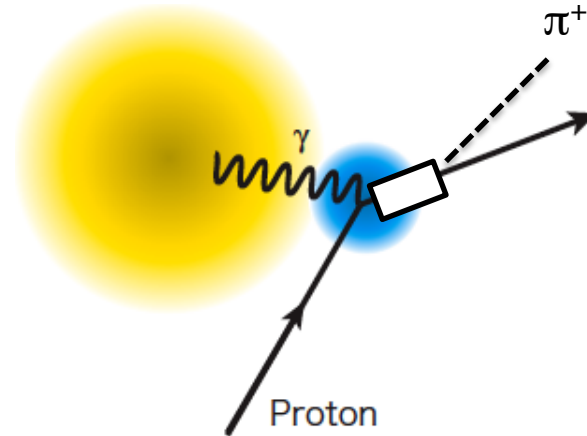
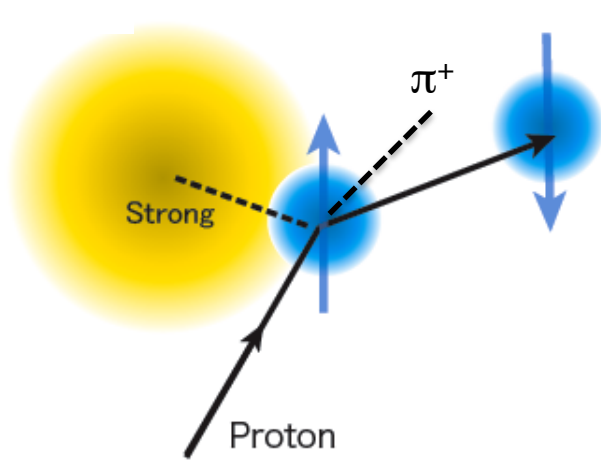
Inelastic

$$\sqrt{s} = 200 \text{ GeV}$$

$$0.02 < -t < 0.5$$

$$\Delta I = 1$$

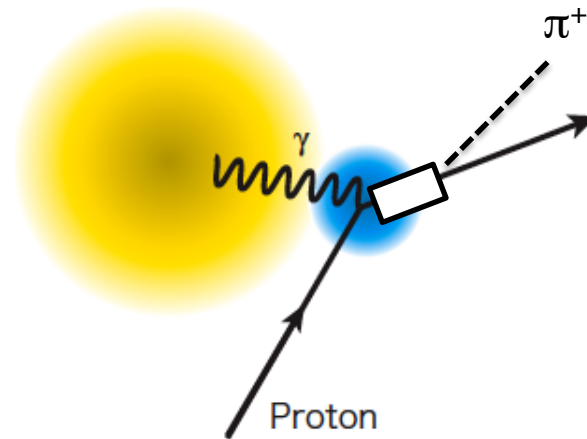
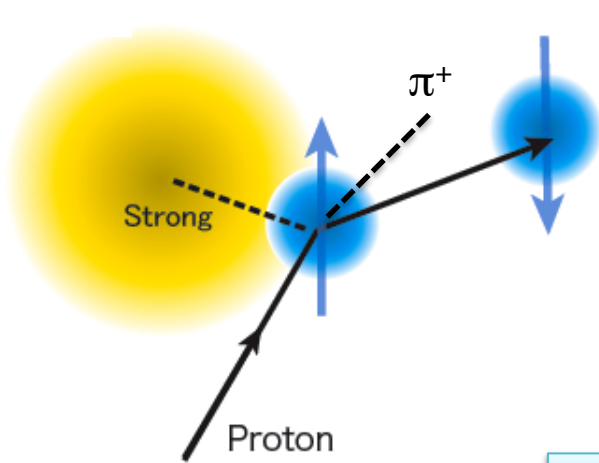
# Coulomb-Nuclear Interference in Forward Neutron Production



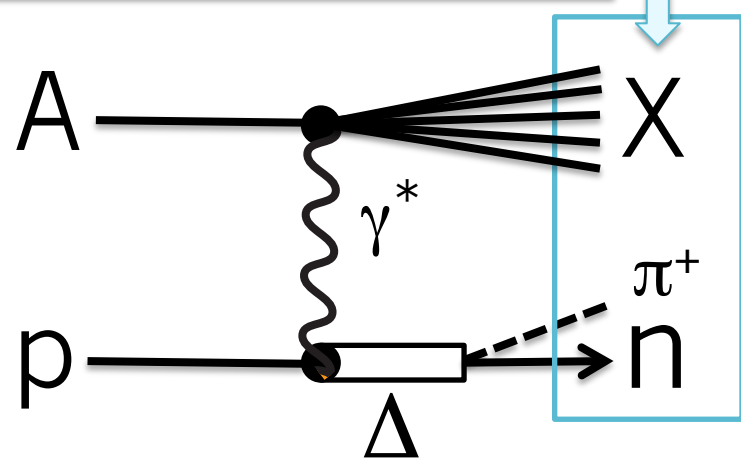
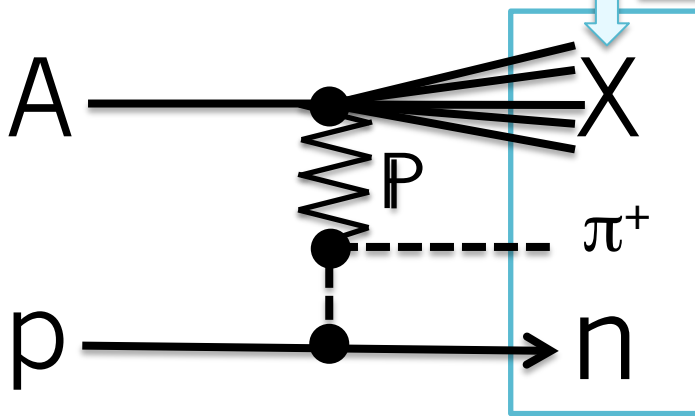
$$A_N \sim had * had + had * EM + EM * had + EM * EM$$

Diffractive Process Required?

# Non-Diffractive Events



Little chance to be same final state



$$A_N \sim had * had + had * EM + EM * had + EM * EM$$

Suppressed?

# Summary

- Observed large asymmetry in very forward neutron  $A_N$  in  $p+A$  collision.
- Strong  $A$ -dependence in neutron + BBC vetoed events and flips the sign of  $A_N$  between  $p+p$  and  $p+A$ .
- Indication of EM from related experiments
  - LHCf observed Primakoff Peak in  $p+Pb \rightarrow n+X$
  - Primakoff MC indicates EM is relatively enhanced in BBC veto events where large positive  $A_N$  is observed.
  - Elastic  $p+p$ ,  $p+C$  at high energy also shows  $A_N < 5\%$  due to interference between EM and strong force.



# Primakoff Effect

Itaru Nakagawa  
RIKEN/RBRC

# Primakoff Paper

## Photo-Production of Neutral Mesons in Nuclear Electric Fields and the Mean Life of the Neutral Meson\*

H. PRIMAKOFF†

*Laboratory for Nuclear Science and Engineering, Massachusetts  
Institute of Technology, Cambridge, Massachusetts*

January 2, 1951

IT has now been well established experimentally that neutral  $\pi$ -mesons ( $\pi^0$ ) decay into two photons.<sup>1</sup> Theoretically, this two-photon type of decay implies zero  $\pi^0$  spin;<sup>2</sup> in addition, the decay has been interpreted as proceeding through the mechanism of the creation and subsequent radiative recombination of a virtual proton anti-proton pair.<sup>3</sup> Whatever the actual mechanism of the (two-photon) decay, its mere existence implies an effective interaction between the  $\pi^0$  wave field,  $\varphi$ , and the electromagnetic wave field,  $\mathbf{E}$ ,  $\mathbf{H}$ , representable in the form:

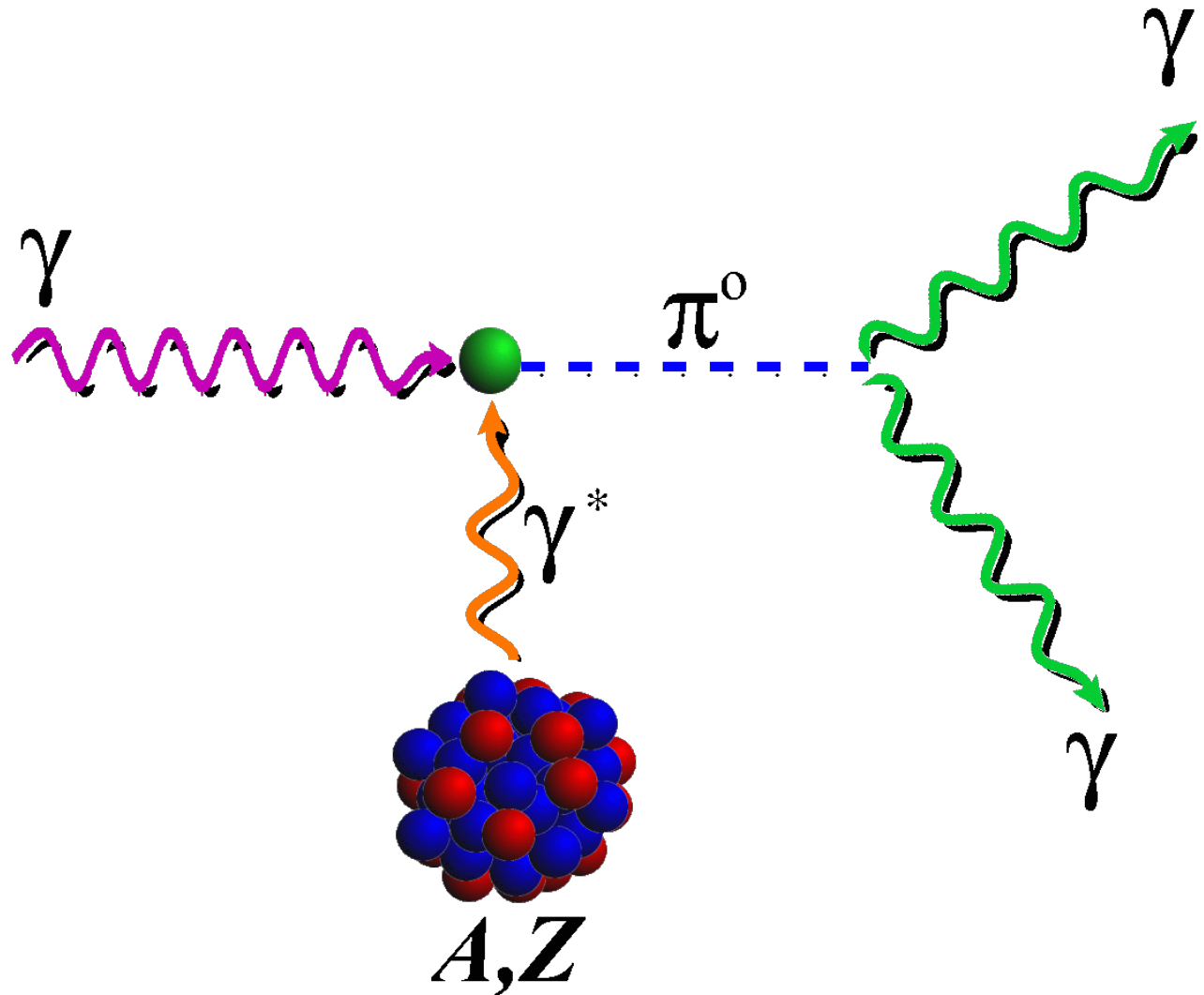
$$\text{Interaction Energy Density} = \eta(\hbar/\mu c)(\hbar c)^{-\frac{1}{2}}\varphi\mathbf{E}\cdot\mathbf{H}. \quad (1)$$

Here  $\varphi$  has been assumed pseudoscalar, the factors  $\hbar/\mu c$  and  $(\hbar c)^{-\frac{1}{2}}$  are introduced for dimensional reasons ( $\mu \equiv$  rest mass of  $\pi^0$ ),

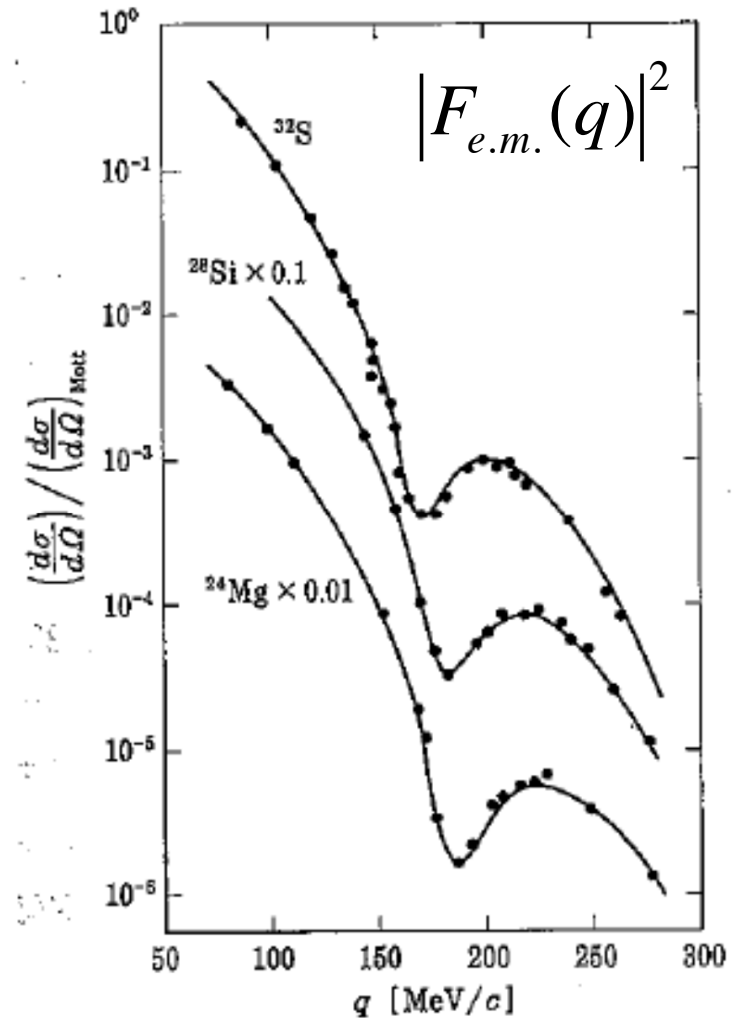
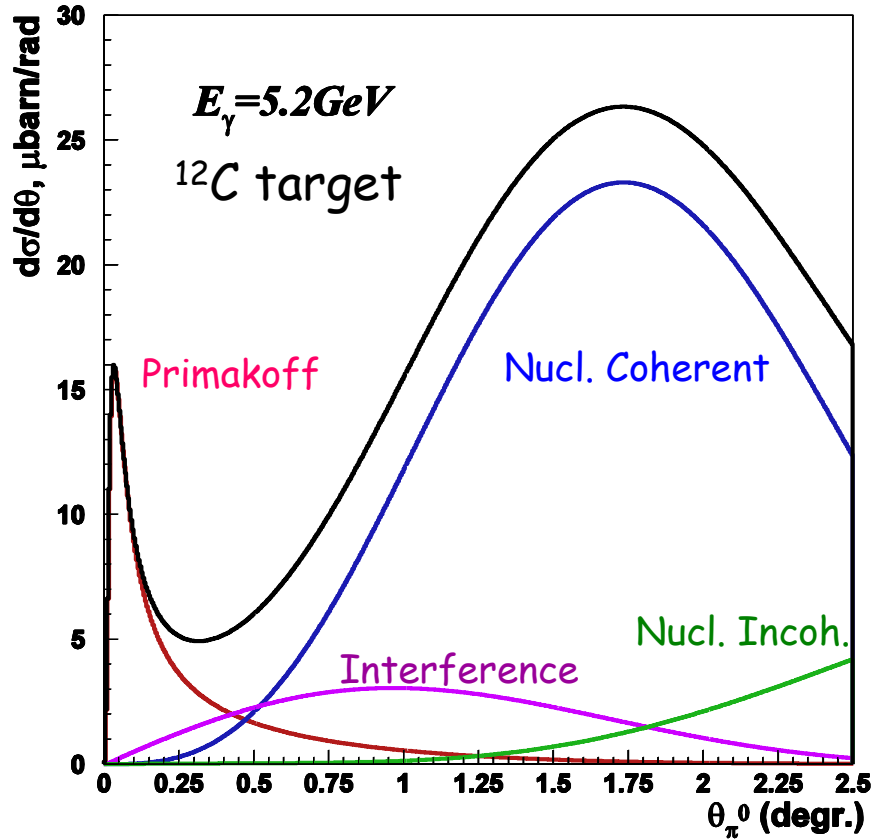


*Henry Primakoff*

# Primakoff Experiment (Hall-B, Jlab)

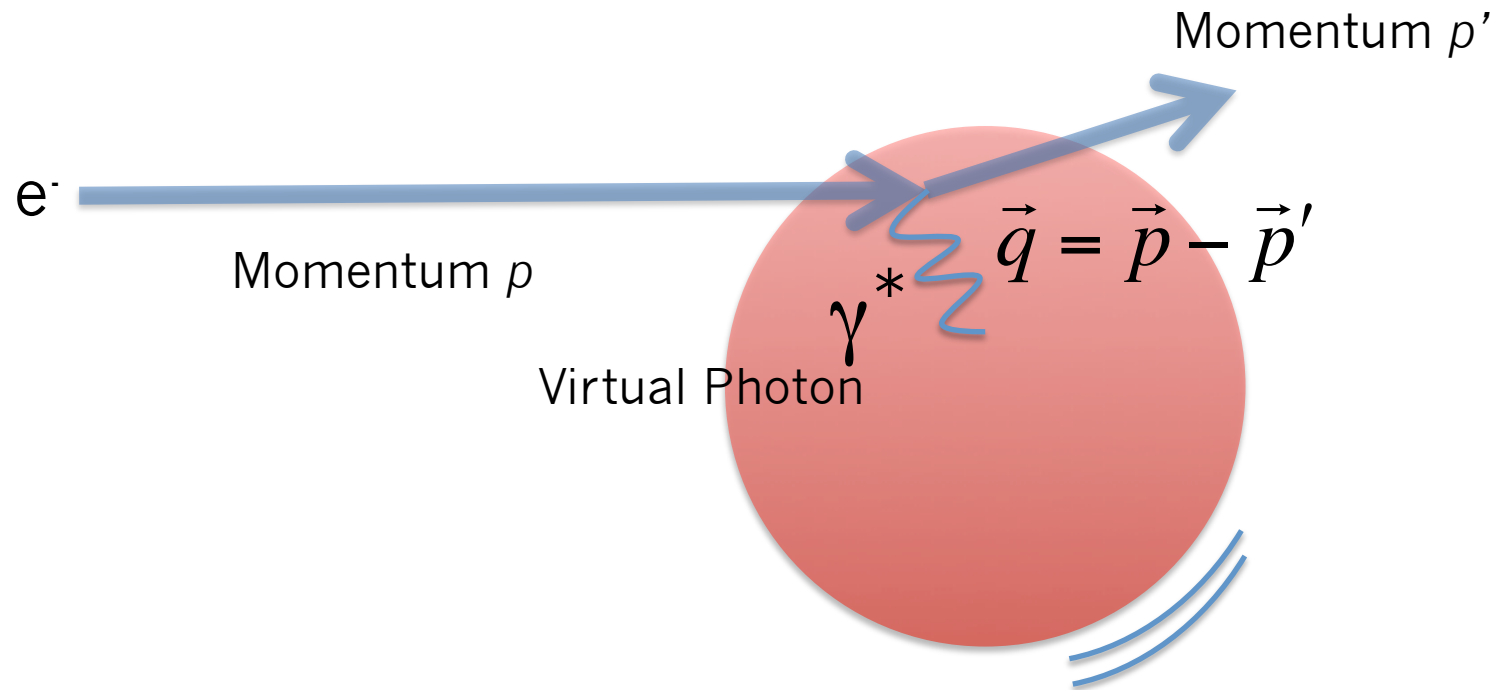


# $\pi^0$ angular distribution



$$\frac{d\sigma_{\text{Primakoff}}}{d\Omega} = \Gamma_{\gamma\gamma} \frac{8\alpha Z^2}{m^3} \frac{\beta^3 E^4}{Q^4} |F_{e.m.}(Q)|^2 \sin^2(\theta)$$

# Nuclear Electromagnetic Form Factor

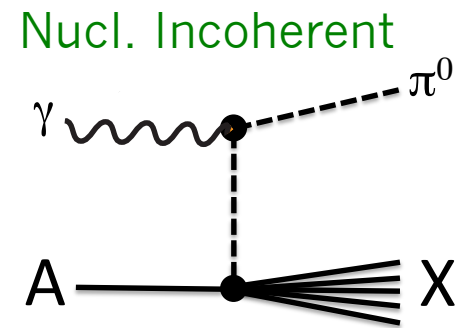
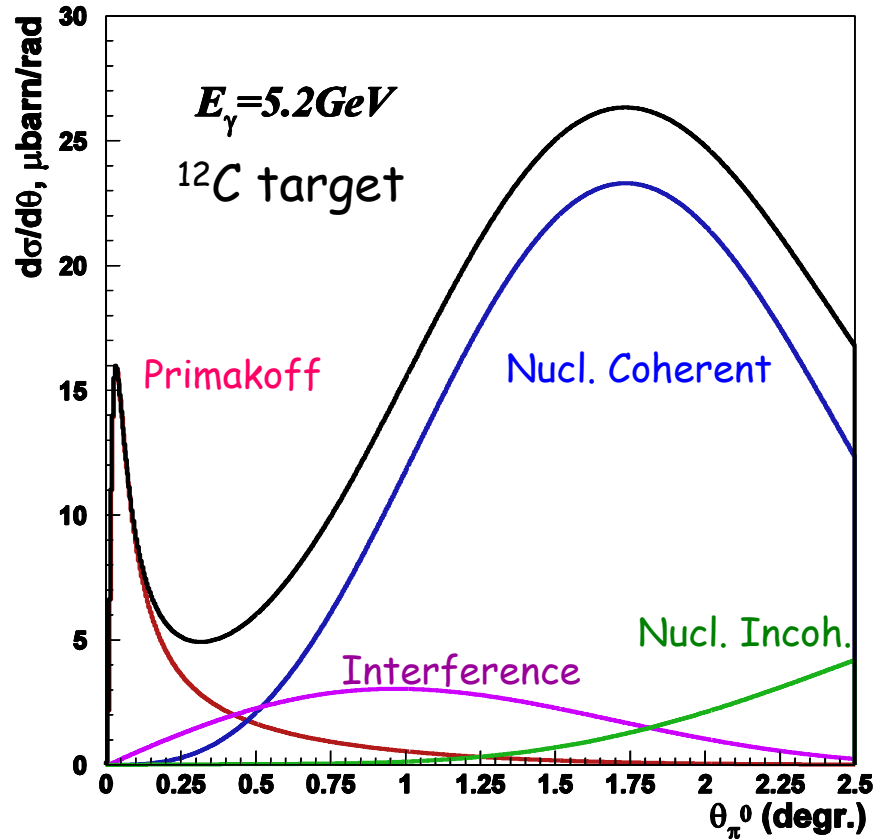
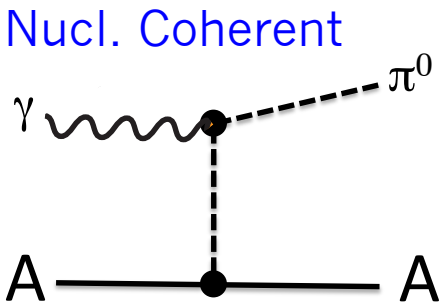
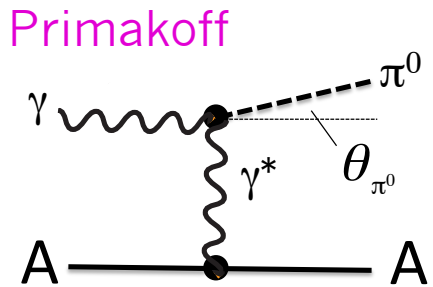


$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{exp}} = \left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}} \cdot |F(q^2)|^2$$

$q$ : Momentum Transfer

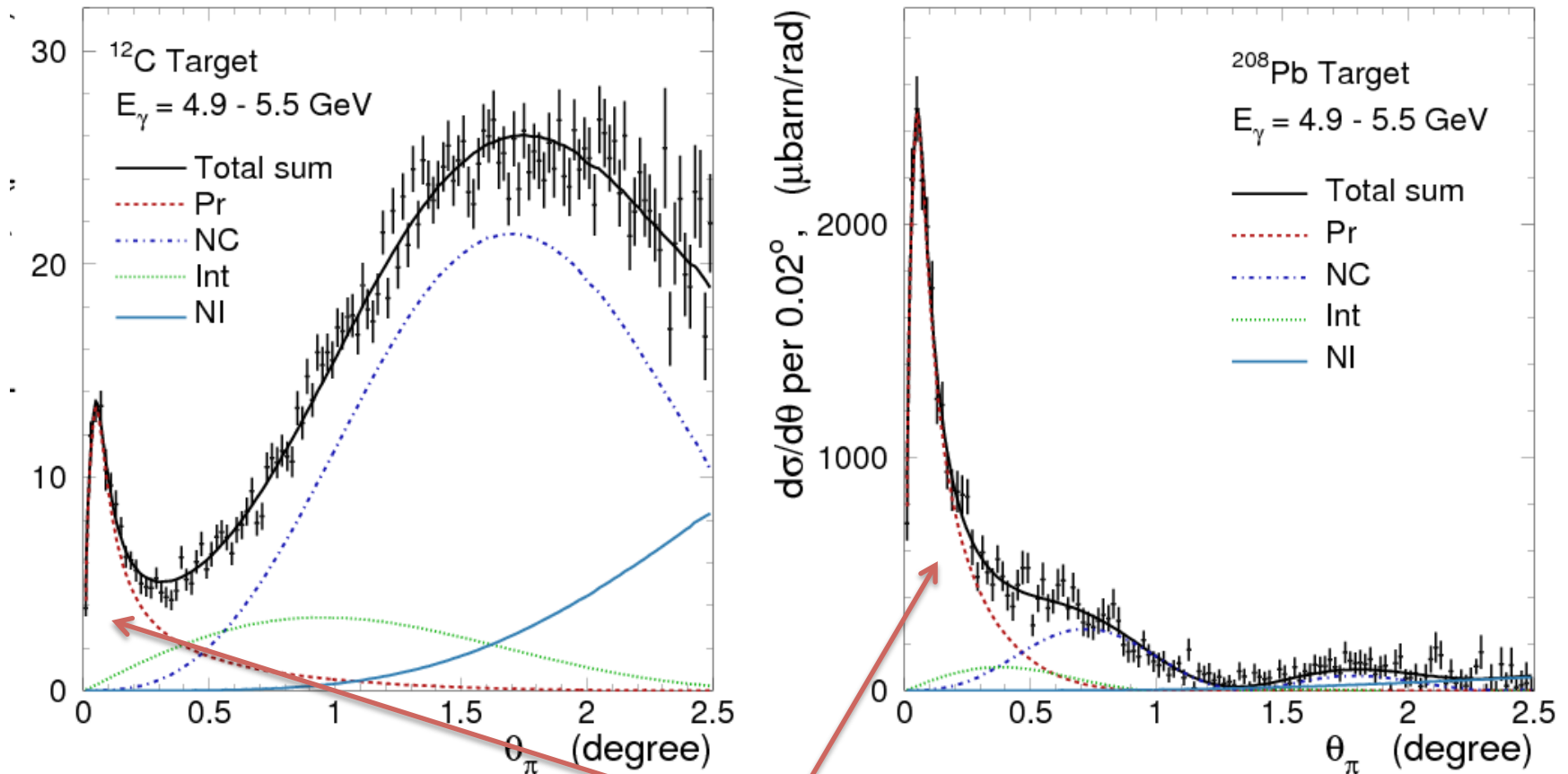
Electromagnetic form factor

# $\pi^0$ angular distribution



$$\frac{d\sigma_{\text{Primakoff}}}{d\Omega} = \Gamma_{\gamma\gamma} \frac{8\alpha Z^2}{m^3} \frac{\beta^3 E^4}{Q^4} |F_{e.m.}(Q)|^2 \sin^2(\theta)$$

# A(Z)-Dependence

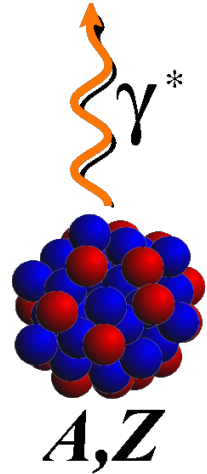


Primakoff

$$\frac{d\sigma_{\text{Primakoff}}}{d\Omega} = \Gamma_{\gamma\gamma} \frac{8\alpha Z^2}{m^3} \frac{\beta^3 E^4}{Q^4} |F_{e.m.}(Q)|^2 \sin^2(\theta)$$

# Primakoff Summary

- Photon pick up from the Coulomb field of nucleus
- Out-going products through Primakoff effect has strongly forward boosted.
- Primakoff effect is suppressed rapidly as a function of emission angle or momentum transfers.
- Primakoff effect gets stronger as a function of  $Z$ .





## Measurement of the Analyzing Power in the Primakoff Process with a High-Energy Polarized Proton Beam

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# The origin of Asymmetry

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According to Eq. (1), the asymmetry seen in photoproduction due to the interference between  $\Delta$  and  $N^*$  is expected in coherent Coulomb  $\pi^0$  production by polarized protons, using the same region of the  $\pi^0$ - $p$  invariant mass. Therefore this process may be used to measure the polarization of the proton at high energies.<sup>6</sup> Until now, there has been no measurement of the asymmetry in the nuclear coherent process.

The cross section for the Coulomb coherent process (1) has a sharp peak at  $t' \sim 10^{-5} (\text{GeV}/c)^2$  and decreases rapidly as  $t'/t^2$ . The “width” of the Coulomb peak is determined by the detector resolution. Diffractive dissociation due to the strong interaction is also present, but it has a much slower  $t'$  dependence.

We have measured the analyzing power (azimuthal asymmetry) of nuclear Coulomb coherent production from a Pb target by using the newly constructed 185-GeV/ $c$  Fermilab polarized proton beam.<sup>7</sup> The beam polarization is 45% and this is further described in Ref. 7. To reduce certain systematic errors, the spin direction of the incident proton was flipped every 10 min using a spin-rotator system.<sup>7</sup>

monitored to a 1% accuracy with a xenon-flash-tube system. A 30-GeV positron beam was used to calibrate the calorimeter. The measured energy resolution is 3% (rms) at 30 GeV and the position resolution is 2 mm (rms). The measured  $\pi^0$  energies in this experiment ranged from 25 to 75 GeV.

A set of thin plastic scintillation counters (TP1) is placed downstream of the magnet and provides the trigger for the scattered protons. The set consists of four counters arranged to distinguish protons scattered to the left, right, up, and down. The calorimeter also has left, right, up, and down sections, and signals from each section are summed for the trigger. In the coherent process where  $t'$  is almost zero, the  $\pi^0$  and scattered protons are coplanar. Thus the trigger logic is such that the energy deposit is larger than 25 GeV in the left half of the calorimeter, less than 5 GeV in the right half, and a proton hits the right segment of TP1. There are four such combinations to cover the whole range of azimuthal angles. To reject the events which have any extra particle besides a proton and  $\pi^0$ , veto counters are included in the trigger logic.

# Experiment

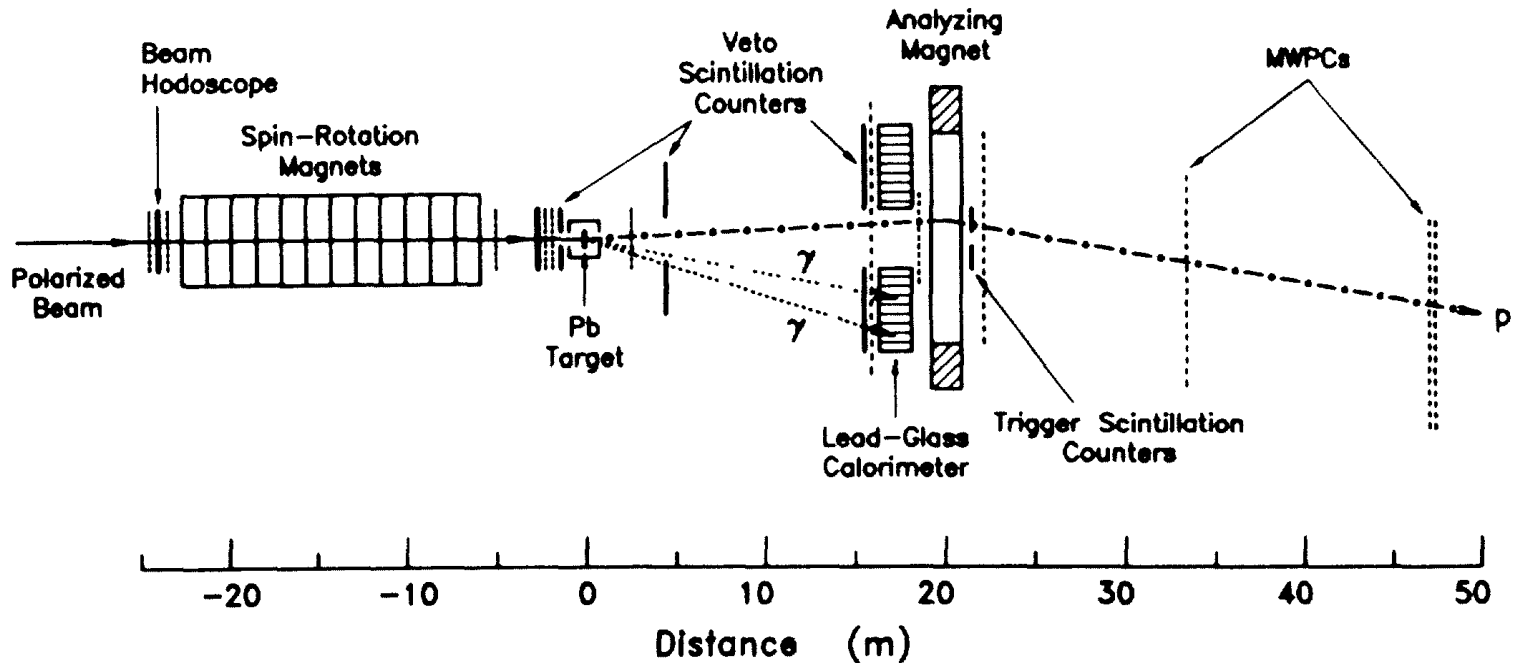
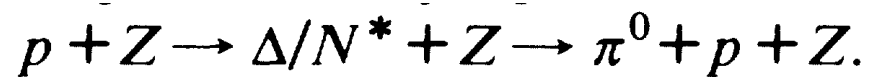


FIG. 1. Schematic view of the experimental setup. The dimensions transverse to the beam are not scale.



# Invariant Mass Distribution

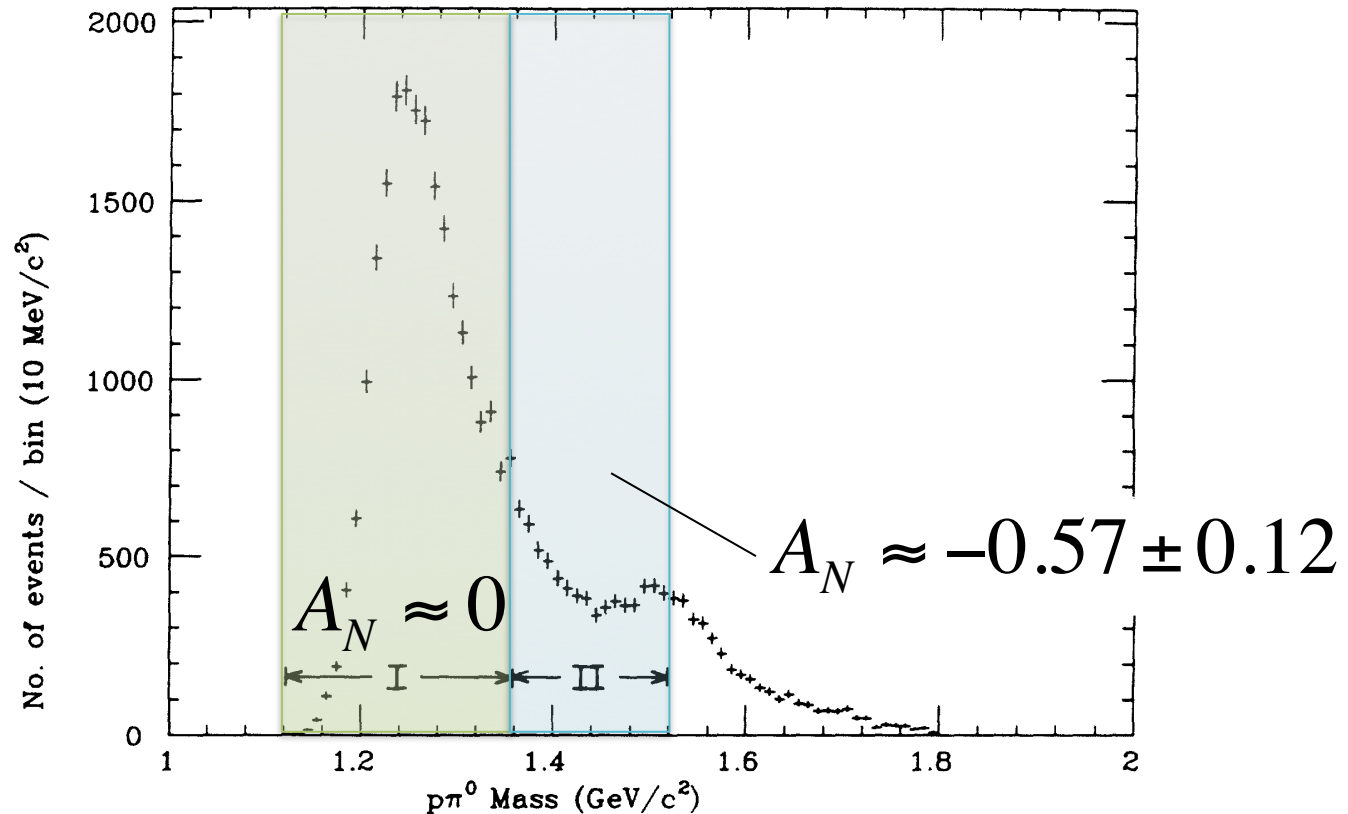


FIG. 2. The invariant-mass spectrum of the  $\pi^0$ - $p$  system in  $p + \text{Pb} \rightarrow \pi^0 + p + \text{Pb}$  for  $|t'| < 1 \times 10^{-3} (\text{GeV}/c)^2$ . Peaks due to the  $\Delta^+(1232)$  and  $N^*(1520)$  resonances are shown. Regions I and II are defined in the text.

# $-t$ Distribution

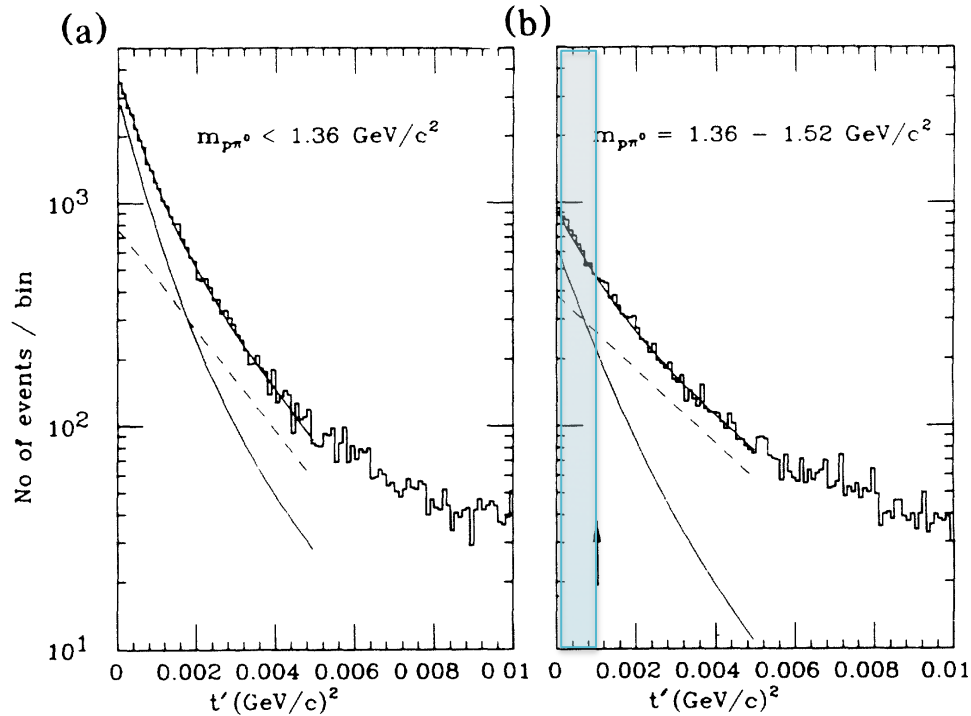
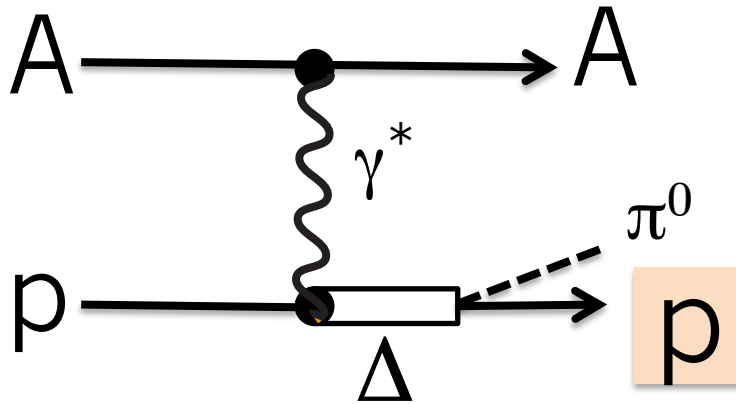


FIG. 3. The  $t'$  distribution of the  $p+\text{Pb} \rightarrow p+\pi^0+\text{Pb}$ , at (a)  $M_{p\pi^0} < 1.36 \text{ GeV}/c^2$  and (b)  $M_{p\pi^0} = 1.36 - 1.52 \text{ GeV}/c^2$ . The curves are obtained by fitting the data. The dashed and dotted lines show the diffractive process and the amount of background events, respectively, and the solid curve represents the Coulomb process. The arrow indicates the position of the  $t'$  cut.

$M$	$< 1.36$ [GeV]	$1.36 - 1.62$ [GeV]
$-t$		
$< 0.001$ (GeV/c) <sup>2</sup>	✗	○
$> 0.001$ (GeV/c) <sup>2</sup>	✗	✗

# Origin of the Asymmetry

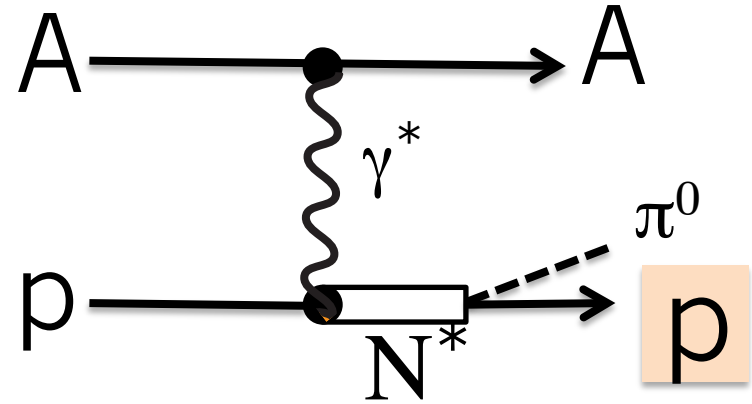
Spin flip



$\Delta^*(1232)$  P33

$\hookrightarrow N\pi$  (>99%)

Spin non-flip



$N^*(1440 \text{ and higher})$  P11

$\hookrightarrow N\pi$  (60~70%)

$N\pi\pi$  (30~40%)

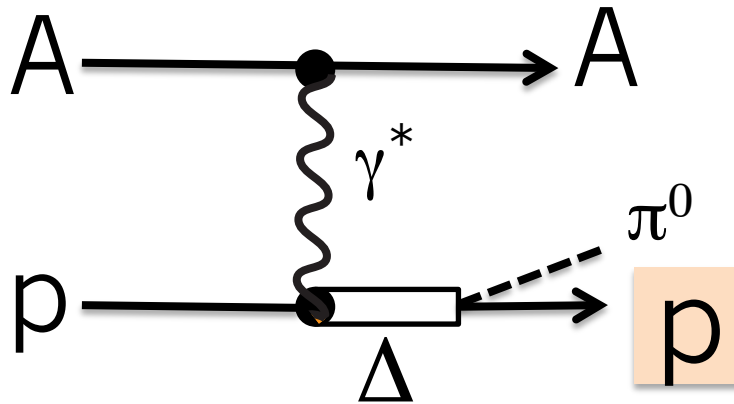
$$A_N \propto \phi_{non-flip}^{had} \phi_{flip}^{had} \sin \delta_1 + \phi_{non-flip}^{EM} \phi_{flip}^{had} \sin \delta_2 + \phi_{non-flip}^{had} \phi_{flip}^{EM} \sin \delta_3 + \phi_{non-flip}^{EM} \phi_{flip}^{EM} \sin \delta_4$$

# Fermi Experiment Summary

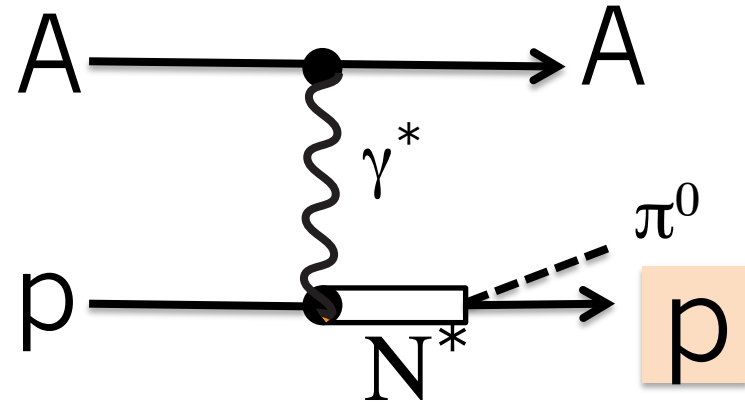
- Large asymmetry observed in Fermi forward  $\pi^0$  production in  $pol(p)+Pb$ .
- Large asymmetry is observed by selecting Primakoff kinematics (Small  $-t$ ).
- The asymmetry is known in photo pion production as a consequence of interference between  $\Delta$  and  $N^*$ .

# Origin of the Asymmetry

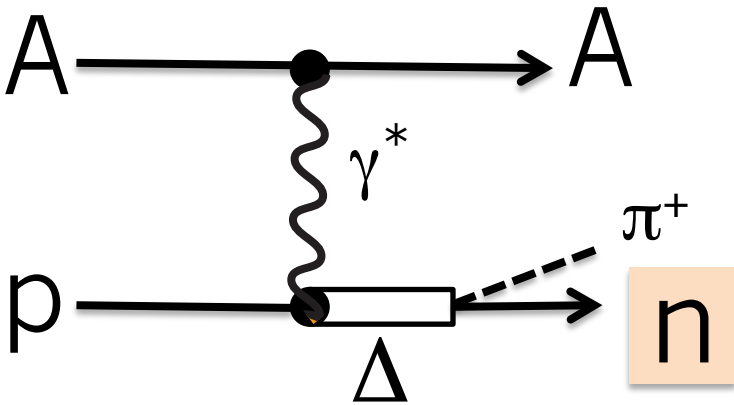
Fermi



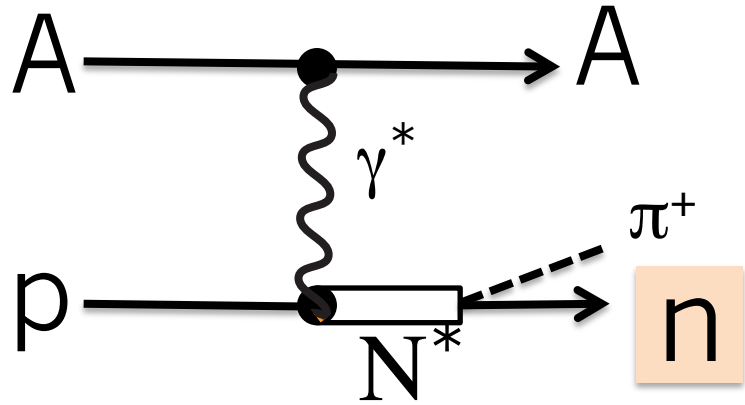
$\times$



RHIC



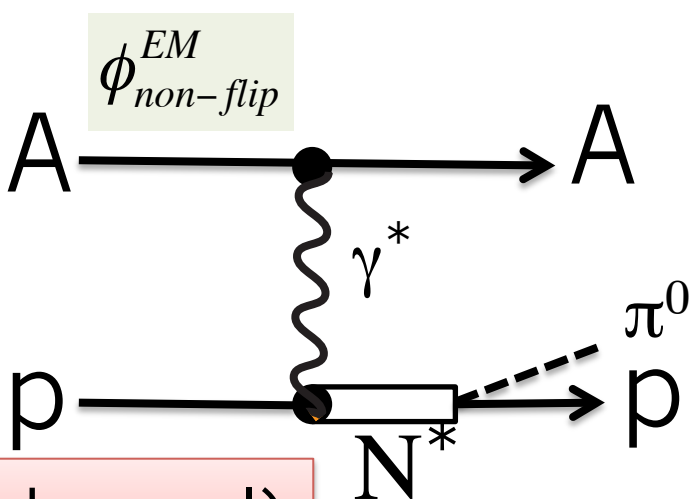
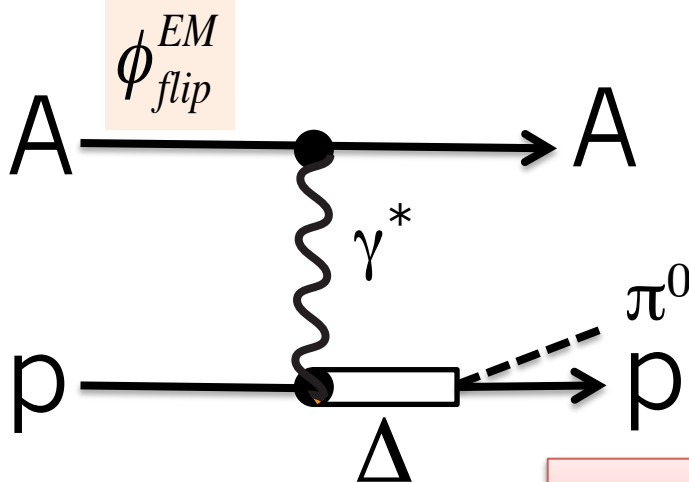
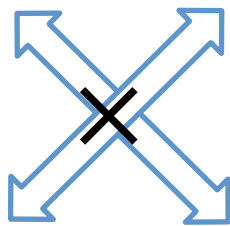
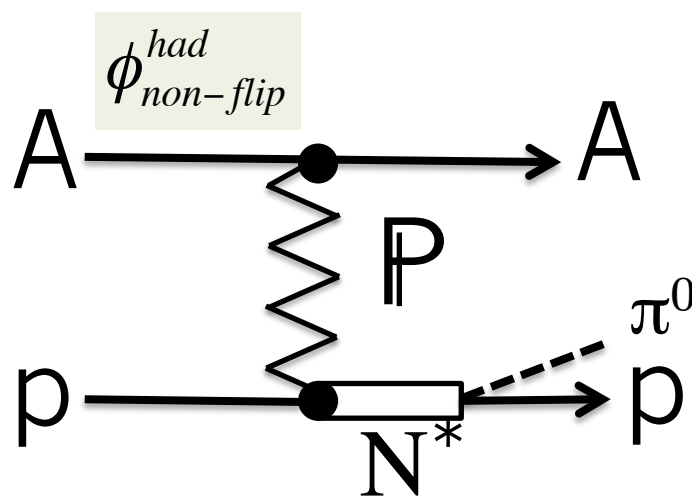
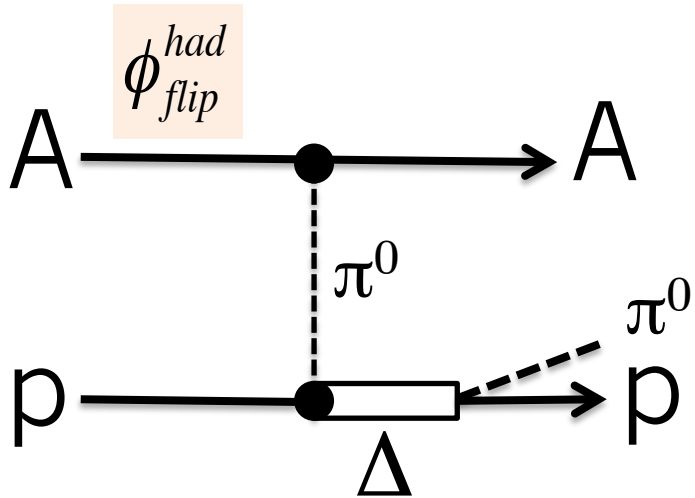
$\times$



$$A_N \propto \phi_{non-flip}^{had} \phi_{flip}^{had} \sin \delta_1 + \phi_{non-flip}^{EM} \phi_{flip}^{had} \sin \delta_2 + \phi_{non-flip}^{had} \phi_{flip}^{EM} \sin \delta_3 + \phi_{non-flip}^{EM} \phi_{flip}^{EM} \sin \delta_4$$

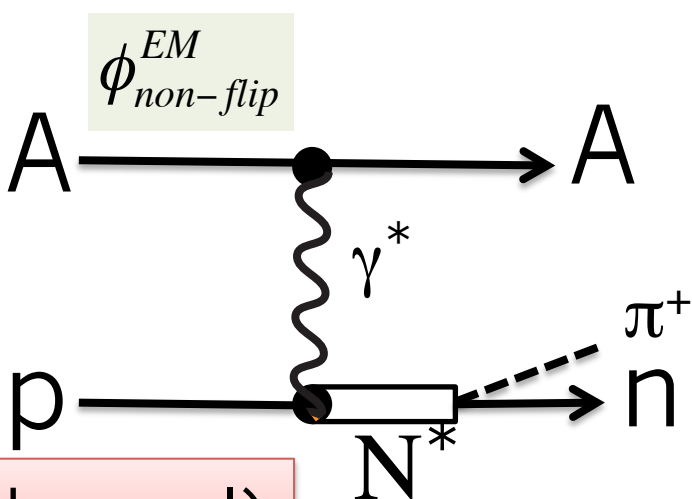
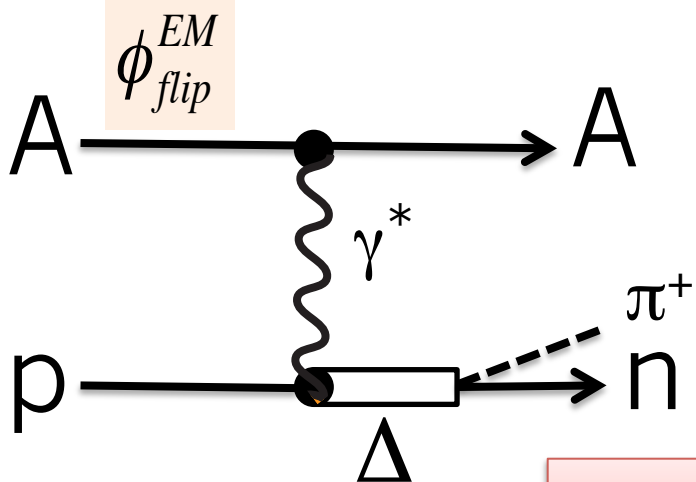
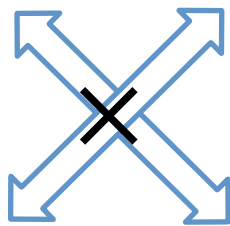
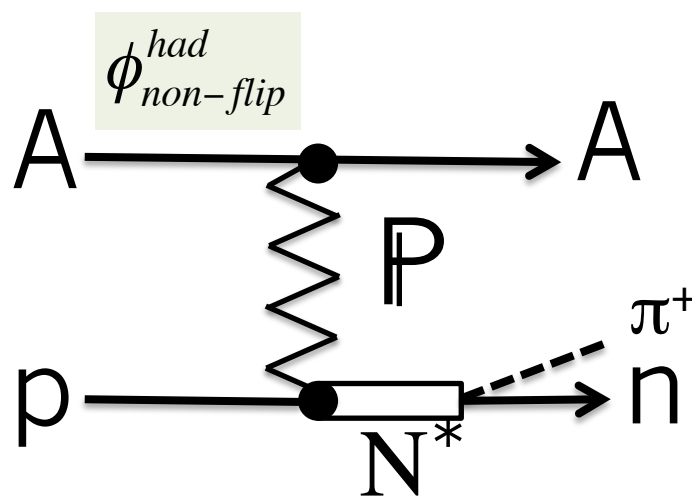
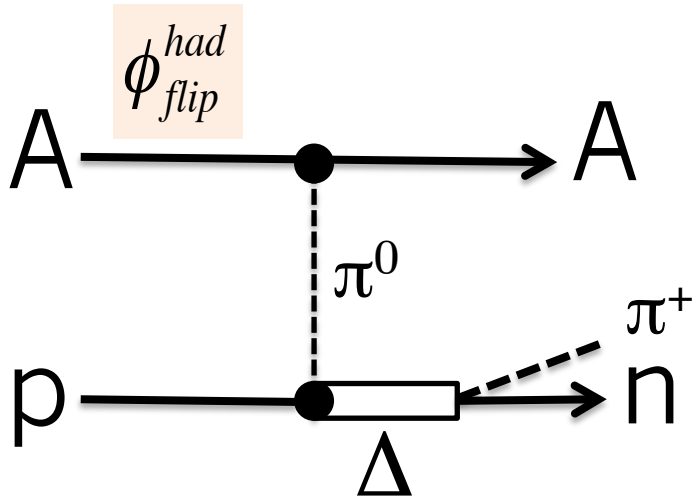


$$A_N \propto \underbrace{\phi_{non-flip}^{had} \phi_{flip}^{had}}_{\text{hadronic}} \sin \delta_1 + \underbrace{\phi_{non-flip}^{EM} \phi_{flip}^{had}}_{\text{CNI}} \sin \delta_2 + \underbrace{\phi_{non-flip}^{had} \phi_{flip}^{EM}}_{\text{Primakoff}} \sin \delta_3 + \underbrace{\phi_{non-flip}^{EM} \phi_{flip}^{EM}}_{\text{Primakoff}} \sin \delta_4$$



Fermi ( $p+\pi^0$  channel)

$$A_N \propto \underbrace{\phi_{non-flip}^{had} \phi_{flip}^{had}}_{\text{hadronic}} \sin \delta_1 + \underbrace{\phi_{non-flip}^{EM} \phi_{flip}^{had}}_{\text{CNI}} \sin \delta_2 + \underbrace{\phi_{non-flip}^{had} \phi_{flip}^{EM}}_{\text{Primakoff}} \sin \delta_3 + \underbrace{\phi_{non-flip}^{EM} \phi_{flip}^{EM}}_{\text{Primakoff}} \sin \delta_4$$

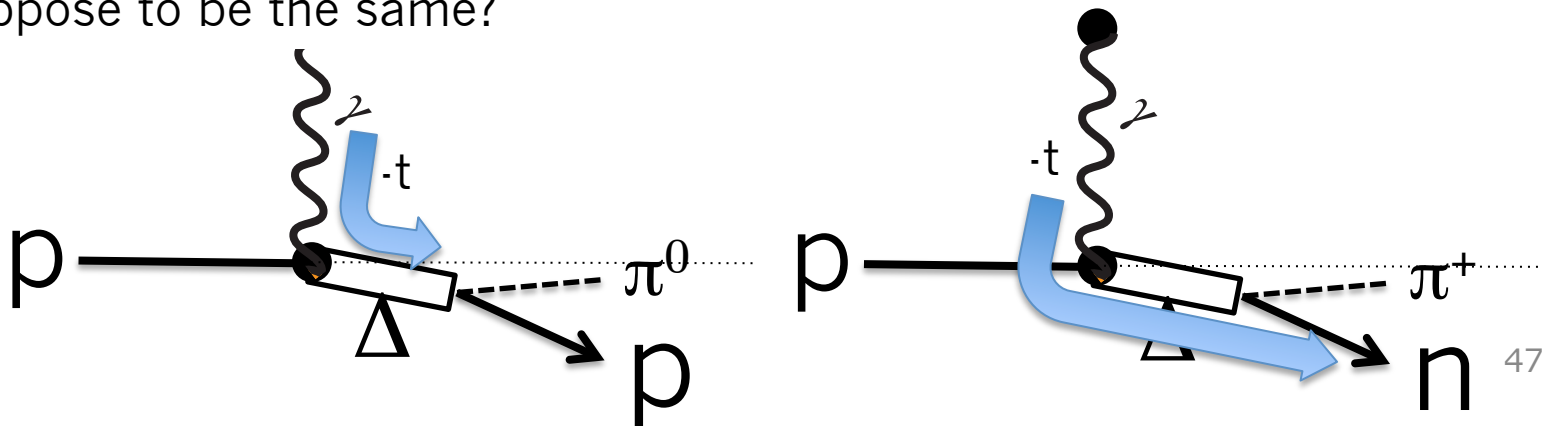


RHIC ( $n+\pi^+$  channel)

# Comparison between two

	Fermi	BNL
Beam Energy [GeV]	185	100
$\sqrt{s}$ [GeV]	19.5	200
Target	Pb	Au
Observables	$p + \pi^0$	$n$ (+ charged)
$t'$	$< 0.001$	$0.02 < -t < 0.5$
$M$	$1.36 < M(\pi^0 p) < 1.52$	?
$A_N$	$-0.57 \pm (0.12)_{sta} + 0.21 - 0.18$	$+ 0.27 \pm 0.003$ (BBC veto)

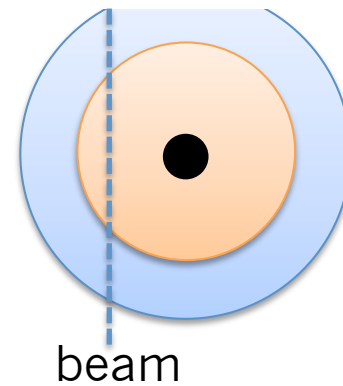
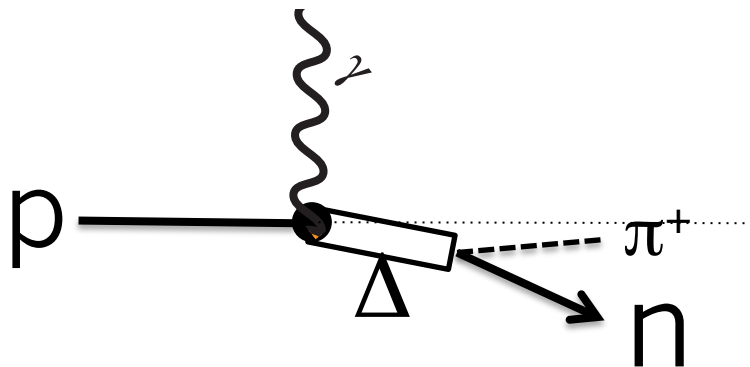
If the asymmetry is induced at  $D$  or  $N^*$  excitation, does the sign suppose to be the same?



# Comparison between two

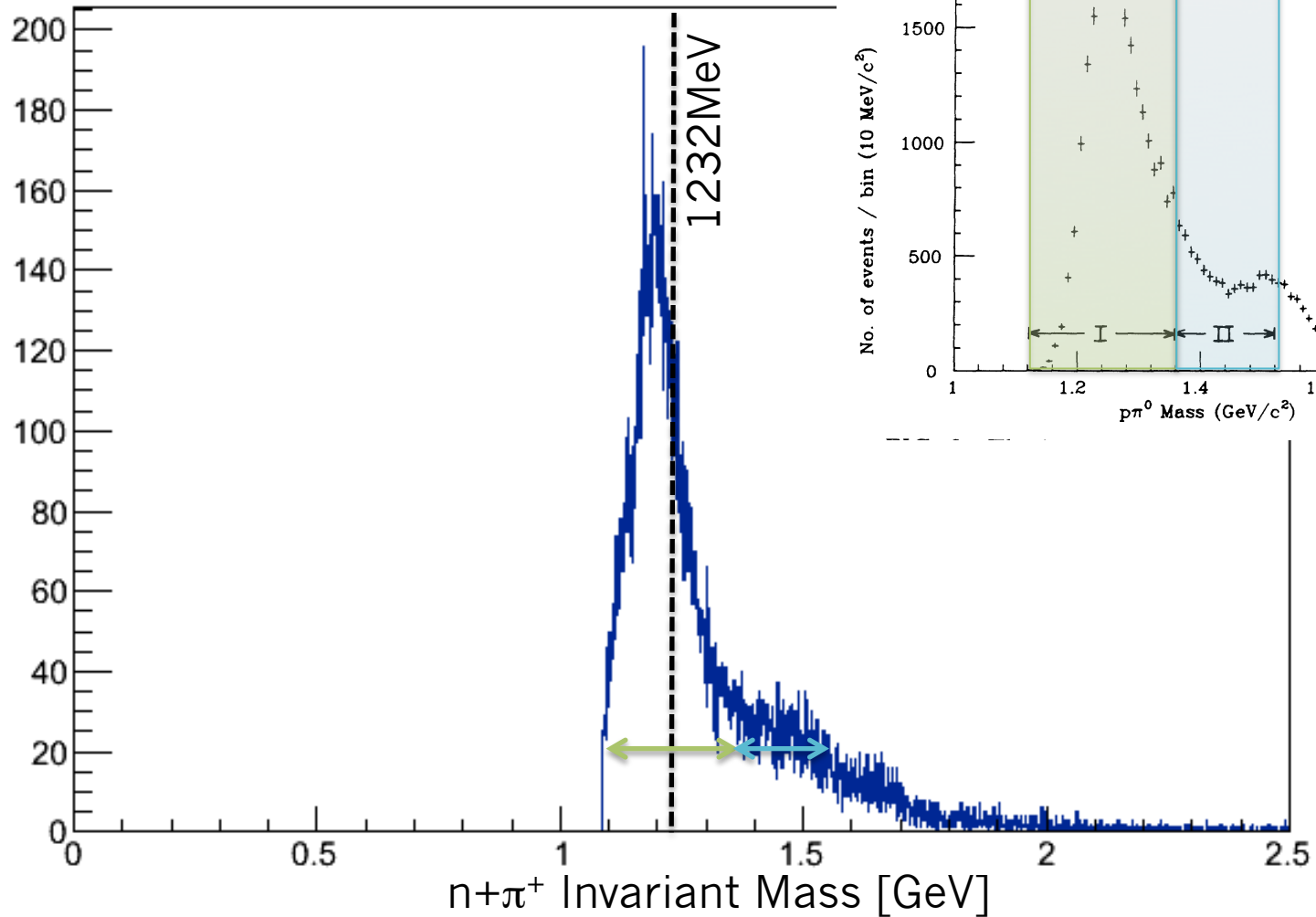
	Fermi	BNL
Beam Energy [GeV]	185	100
$\sqrt{s}$ [GeV]	19.5	200
Target	Pb	Au
Observables	$p + \pi^0$	$n$ (+ charged)
$t'$	$< 0.001$	$0.02 < -t < 0.5$
$M$	$1.36 < M(\pi^0 p) < 1.52$	?
$A_N$	$- 0.57 \pm (0.12)_{sta} + 0.21 - 0.18$	$+ 0.27 \pm 0.003$ (BBC veto)

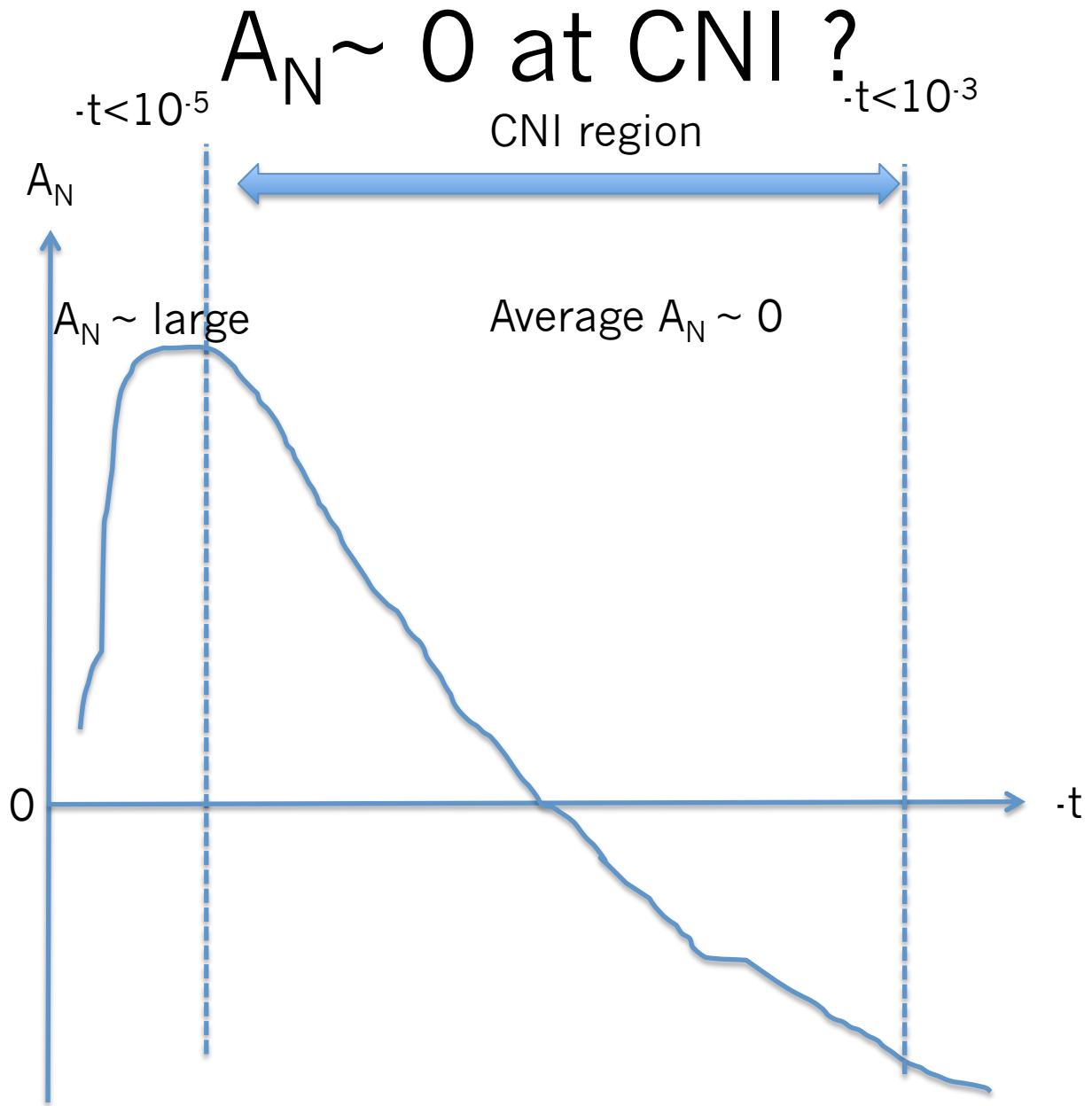
If the asymmetry is induced at D or  $N^*$  excitation, does the sign suppose to be the same?



# UPC MC

Invariant mass plot of  $n$





# What we can learn from Fermi Exp.?

- If we consider the difference between PHENIX neutron and  $p+\pi^0$  as just a difference in the decay channels from  $\Delta$  and  $N^*$ , then both mechanisms are similar.
- PHENIX  $-t$  is expected to be larger than Fermi, so CNI may be important in PHENIX.
  - Although Fermi claims  $A_N$  is zero at CNI, it may be just because of accident?
  - Not straight forward to perform kinematic decomposition in PHENIX data.
- Sign of  $A_N$  are opposite between two experiments. Are we measuring the same  $A_N$ ?

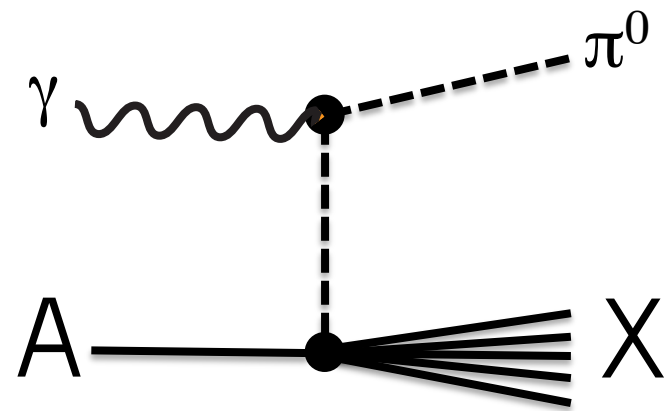
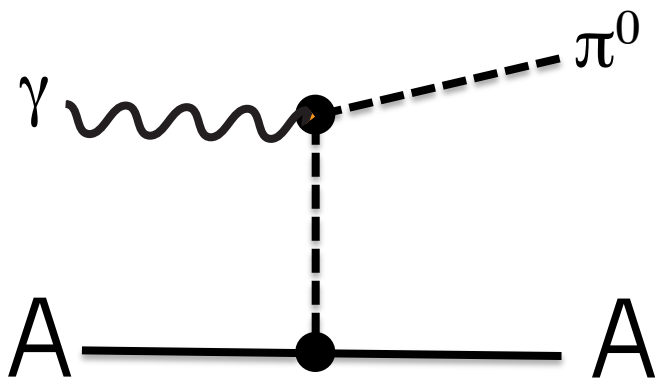
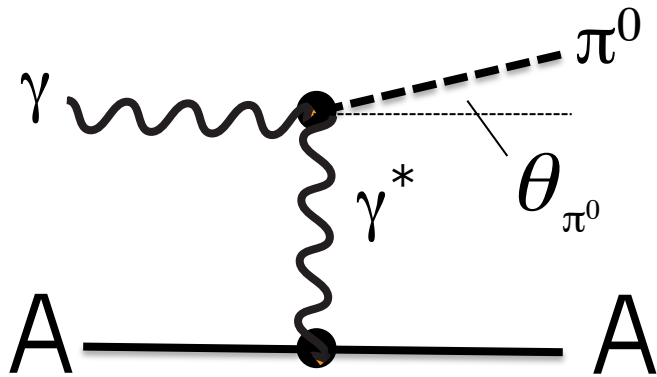
# BACKUPS



# Comparison between two

	Fermi	PHENIX	STAR17
Beam Energy [GeV]	185	100	255
$\sqrt{s}$ [GeV]	19.5	200	22
Target	Pb	Au	Al/Sn/Au
Observables	$p + \pi^0$	$n (+ \text{charged})$	$n (+ \text{charged})$ $\pi^0 ?$
$t'$	$< 0.001$	$0.02 < -t < 0.5$	
M	$1.36 < M(\pi^0 p) < 1.52$	?	?
$A_N$	$- 0.57 \pm (0.12)_{\text{sta}}^+$ $0.21 - 0.18$	$+ 0.27 \pm 0.003$ (BBC veto)	

# Primakoff Diagrams



# Nuclear/Nucleon Photo-Excitation

