

Heavy quark spin in nuclear medium

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1. Multi-flavor nucleus

Nucleus meets Flavor

HiggsTan.com

charge

$Q=+2/3$

$Q=-1/3$

1st

up

2 MeV

$Q=-1/3$

down

5 MeV

2nd

charm

1300 MeV

strange

100 MeV

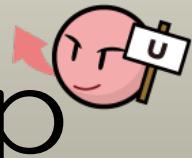
3rd

top

173000 MeV

bottom

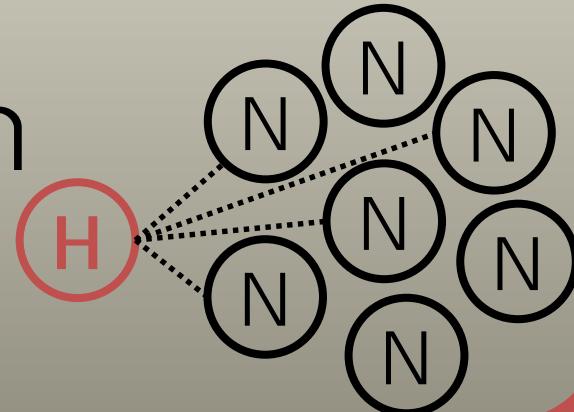
4200 MeV



1. Multi-flavor nucleus

Nucleus meets Flavor

1. Hadron-nucleon interaction
2. Hadron in medium
3. Nuclear structure



Impurity Physics

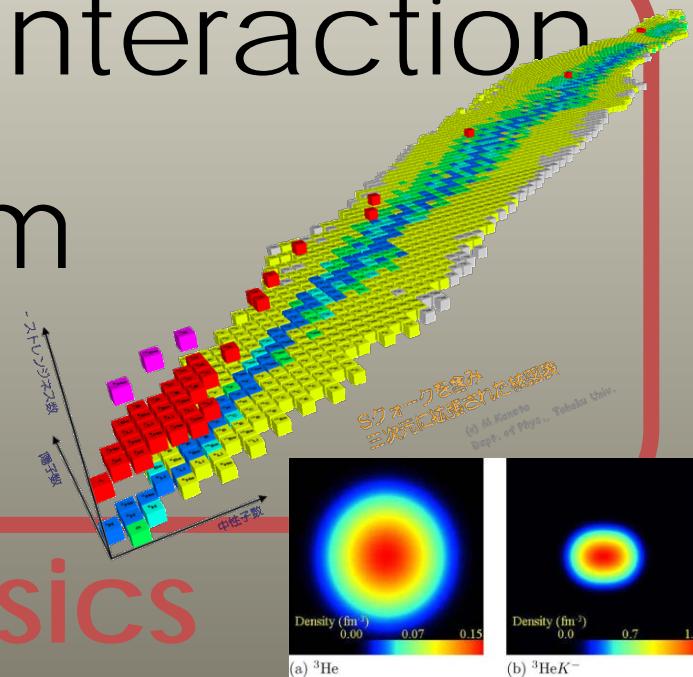
1. Multi-flavor nucleus



Nucleus + Strangeness

1. Hadron-nucleon interaction
2. Hadron in medium
3. Nuclear structure

Impurity Physics



1. Multi-flavor nucleus



Nucleus + Charm

Mass hierarchy ($m_c \gg \Lambda_{\text{QCD}}$)

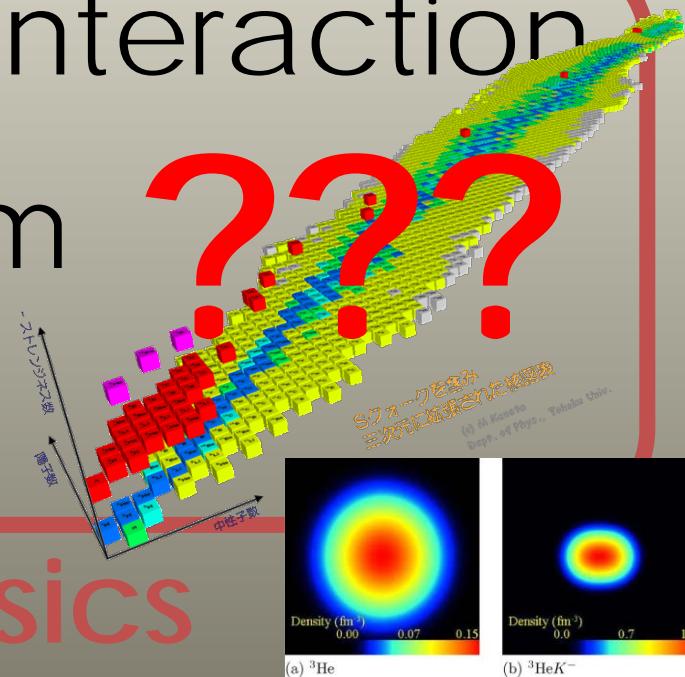
~1.2 GeV ~0.2 GeV

1. Hadron-nucleon interaction

2. Hadron in medium

3. Nuclear structure

Impurity Physics

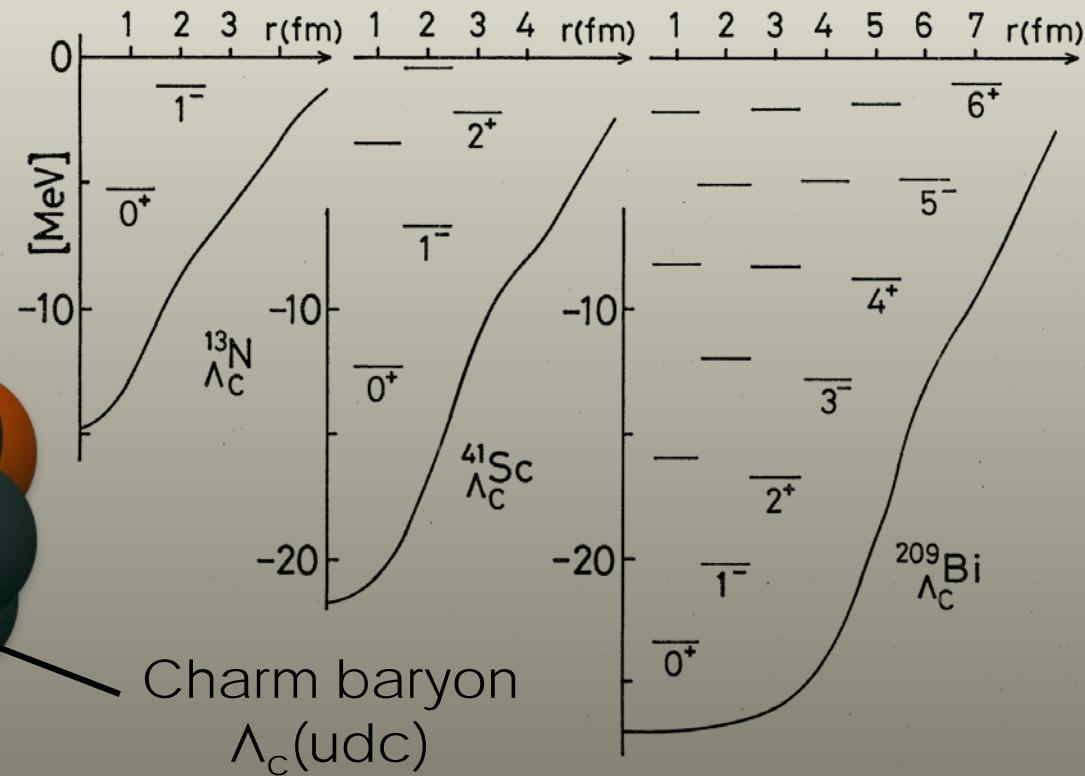
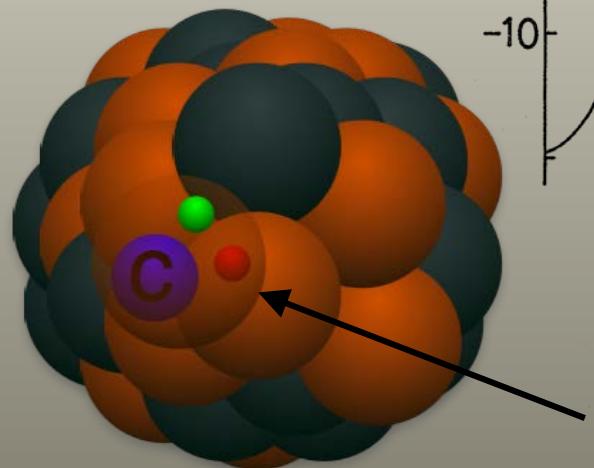


1. Multi-flavor nucleus

Strangeness

Charm

$$\Lambda(\text{uds}) \rightarrow \Lambda_C(\text{udc})$$

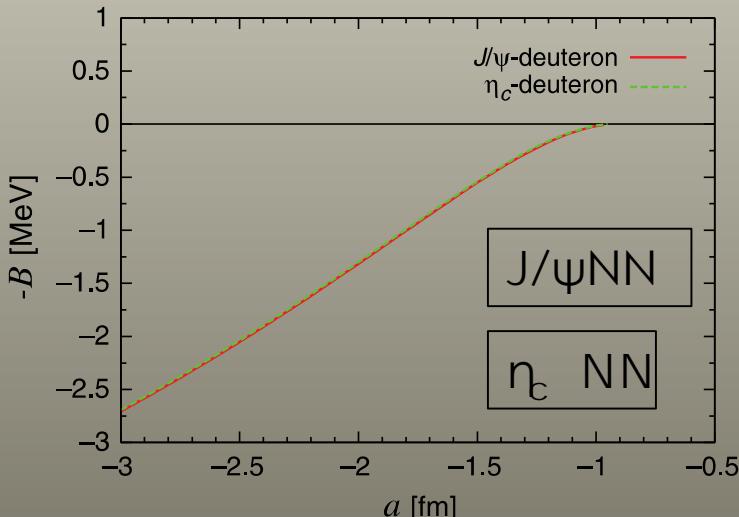


1. Multi-flavor nucleus

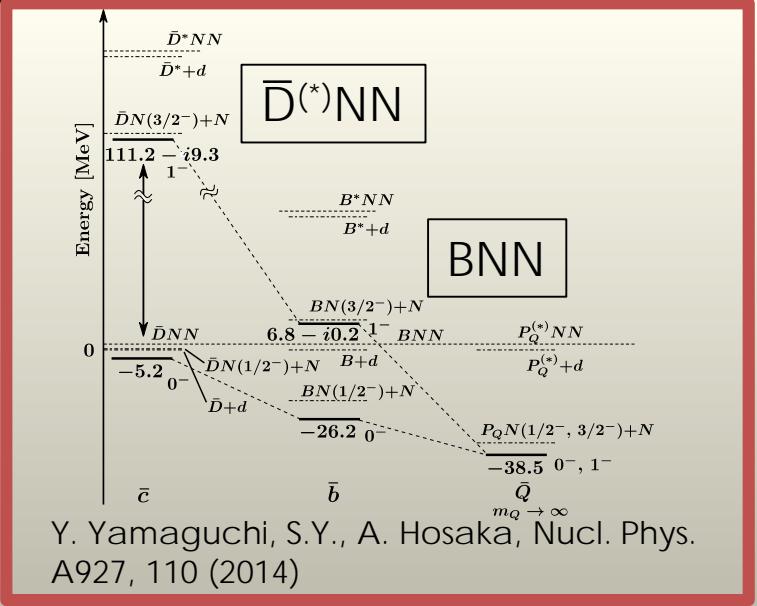
Charm/bottom “light” nuclei

	HN1R $J = 1$	$J = 0$	Minnesota $J = 0$	Av18 $J = 0$
DNN				
B	Unbound	Bound	Bound	Bound
M_B	208	225	251	209
$\Gamma_{\pi Y_c N}$	3537	3520	3494	3536
E_{kin}	—	26	38	22
$V(NN)$	338	352	438	335
$V(DN)$	0	-2	19	-5
$V(DN)$	-546	-575	-708	-540
T_{nuc}	113	126	162	117
E_{NN}	113	124	181	113
$P(\text{odd})$	75.0%	14.4%	7.4%	18.9%

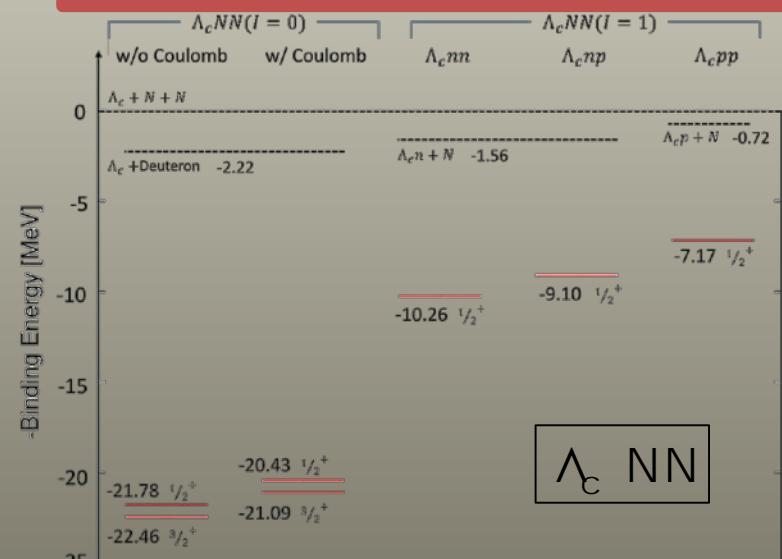
M. Bayer, C. W. Xiao, T. Hyodo, A. Dote, M. Oka, E. Oset, Phys. Rev. C86, 044004 (2012)



A. Yokota, E. Hiyama, M. Oka,
Pog. Ther. Exp. Phys. 2013, 113D01



Y. Yamaguchi, S.Y., A. Hosaka, Nucl. Phys. A927, 110 (2014)

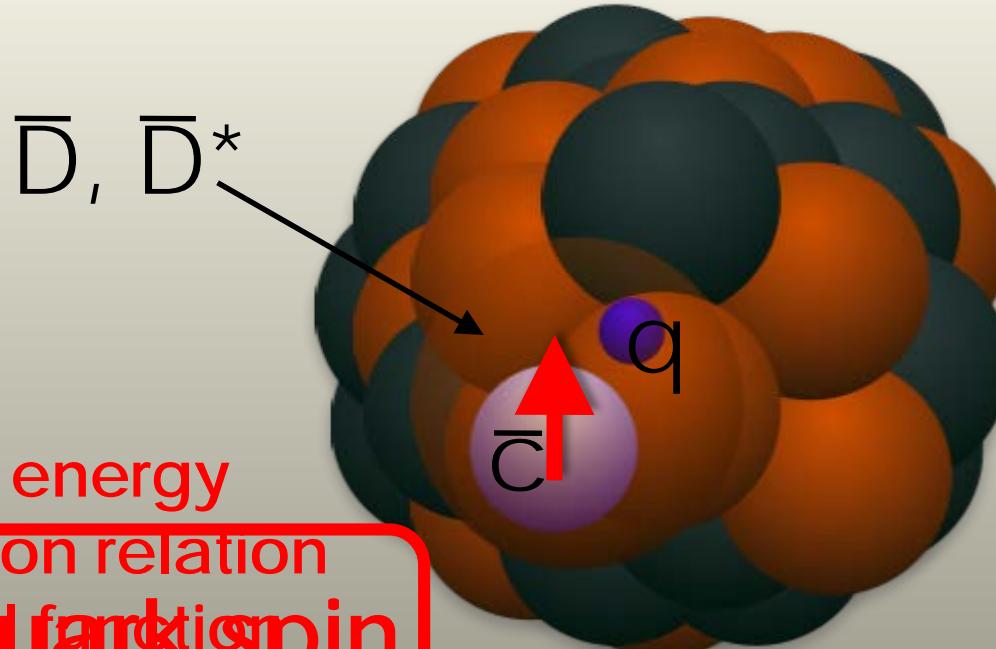


S. Maeda, M. Oka, A. Yokota, E. Hiyama,
Y-R. Liu, Pog. Ther. Exp. Phys. 2016, 023D02

1. Multi-flavor nucleus

\bar{D} (B) meson in nuclear matter

Quark-meson coupling model, QCD sum rules, hadronic interactions (mean-field approach, channel coupling, pion interaction), ...



- Binding energy

- Dispersion relation

Heavy quark spin

- Reaction

- Restoration of xSB

- Nuclear structure

- etc.

Nucleus

Charm nucleus mass spectroscopy

1. Multi-flavor nucleus

\bar{D} (B) meson in nuclear matter

Quark-meson coupling model, QCD sum rules, hadronic interactions (mean-field approach, channel coupling, pion interaction), ...

Binding energy in nuclear matter [MeV]					
	Quark-meson coupling model	QCD sum rule	Mean field	Hadron dynamics	π exchange
D meson	$-(10-35)$ (in ^{208}Pb) [6] \sim $\simeq +20$ [7]	$+(65-75)$ [15] -50 MeV [27] +30 MeV [28] -66 MeV [29]	$\approx -(100-180)$ [18] -30 [21] -27 [22, 24]	$+18$ [19] $+(11-20)$ [20] $+(27-35)$ [23] $\simeq -(20-27)$ [25]	-35.1 [26]
D^* meson					$-153.9 - i170.1$
B meson		$+(125-155)$ [15]			-106.9
B^* meson					$-203.3 - i117.5$

(Repulsive? Attractive?)

[6] K. Tsushima, D. -H. Lu, A. W. Thomas, K. Saito and R. H. Landau, Phys. Rev. C 59, 2824 (1999).

[7] A. Sibirtsev, K. Tsushima and A. W. Thomas, Eur. Phys. J. A 6, 351 (1999).

[15] T. Hilger, R. Thomas and B. Kampfer, Phys. Rev. C 79, 025202 (2009).

[16] T. Hilger, R. Schulze and B. Kampfer, J. Phys. G G 37, 094054 (2010).

[17] Z. -G. Wang and T. Huang, Phys. Rev. C 84, 048201 (2011).

[18] A. Mishra, E. L. Bratkovskaya, J. Schaffner-Bielich, S. Schramm and H. Stoecker, Phys. Rev. C 69, 015202 (2004).

[19] M. F. M. Lutz and C. L. Korpa, Phys. Lett. B 633, 43 (2006).

[20] L. Tolos, A. Ramos and T. Mizutani, Phys. Rev. C 77, 015207 (2008).

[21] A. Mishra and A. Mazumdar, Phys. Rev. C 79, 024908 (2009).

[22] A. Kumar and A. Mishra, Phys. Rev. C 81, 065204 (2010).

[23] C. E. Jimenez-Tejero, A. Ramos, L. Tolos and I. Vidana, Phys. Rev. C 84, 015208 (2011).

[24] A. Kumar and A. Mishra, Eur. Phys. J. A 47, 164 (2011).

[25] C. Garcia-Recio, J. Nieves, L. L. Salcedo and L. Tolos, Phys. Rev. C 85, 025203 (2012).

[26] S. Yasui, K. Sudoh, Phys. Rev. C87, 015202 (2013).

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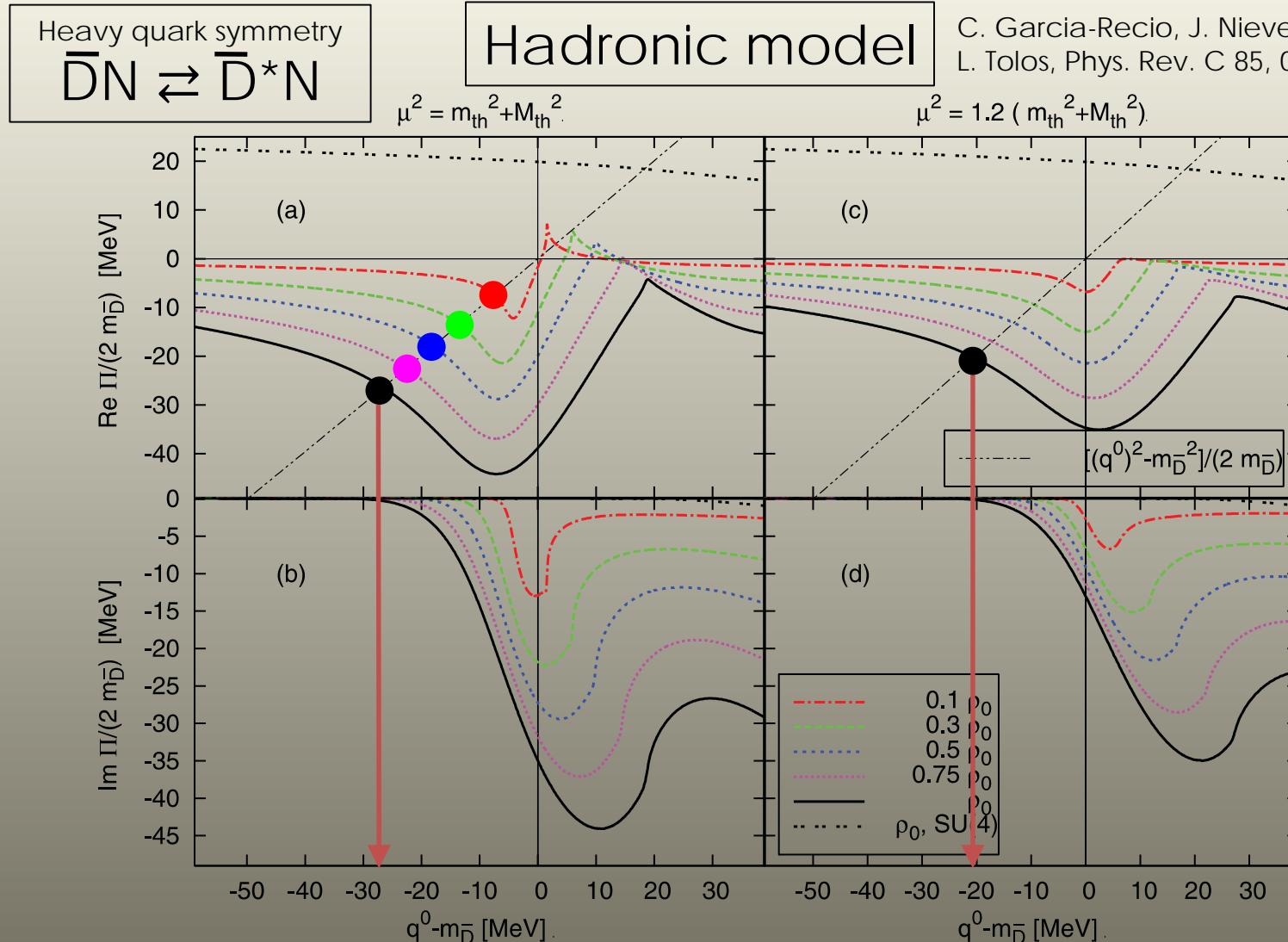
[28] K. Suzuki, P. Gubler, M. Oka, Pos Hadron2013, 179 (2014).

[29] K. Azzi, N. Er, H. Sundu, Eur. Phys. J. C74, 3021 (2014).

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\bar{D} (B) meson in nuclear matter

Quark-meson coupling model, QCD sum rules, hadronic interactions (mean-field approach, channel coupling, pion interaction), ...



1. Multi-flavor nucleus

\bar{D} (B) meson in nuclear matter

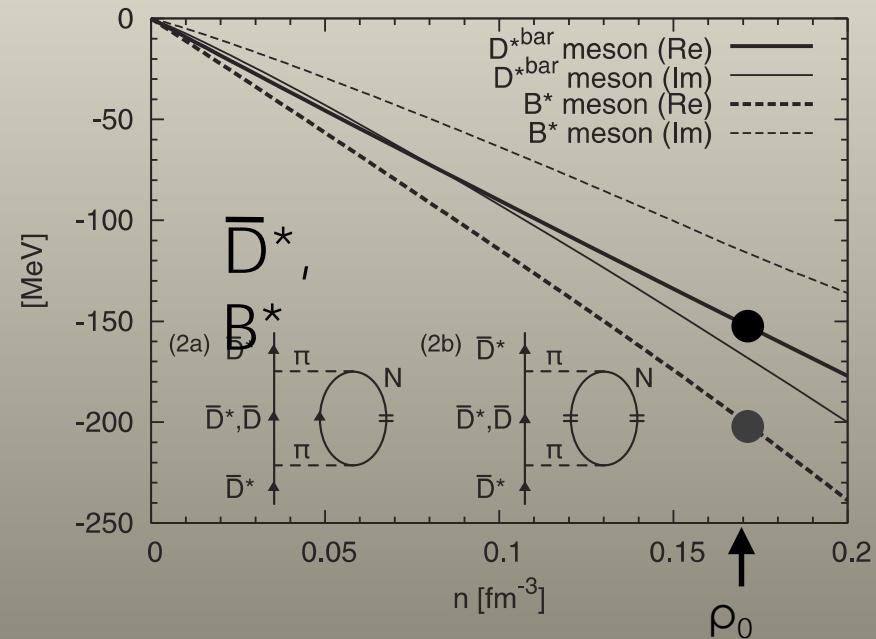
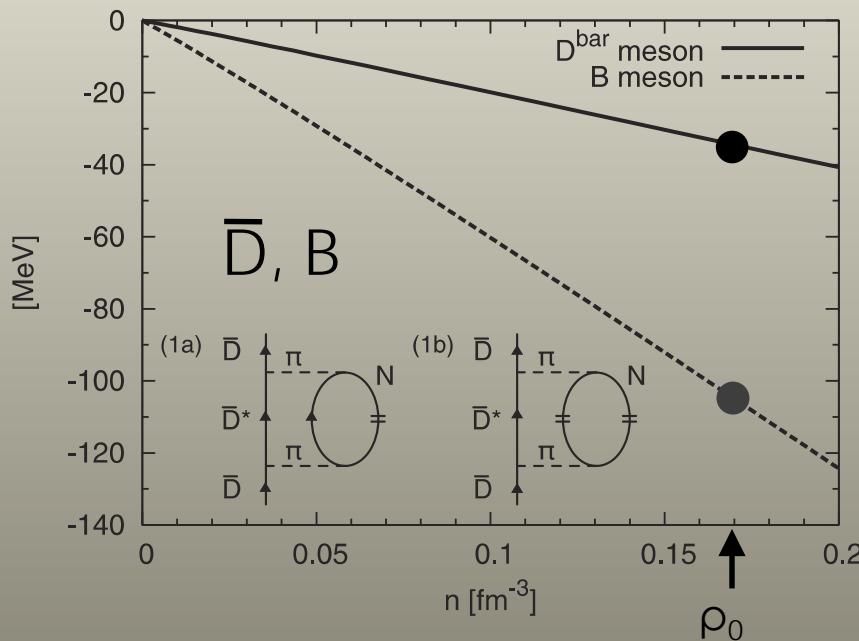
Quark-meson coupling model, QCD sum rules, hadronic interactions (mean-field approach, channel coupling, pion interaction), ...

Heavy quark symmetry
 $\bar{D}N \rightleftharpoons \bar{D}^*N$

Hadronic model

S. Y. , K. Sudoh, Phys Rev. C87, 015202 (2013)

π meson exchange between \bar{D} (\bar{D}^*) and N

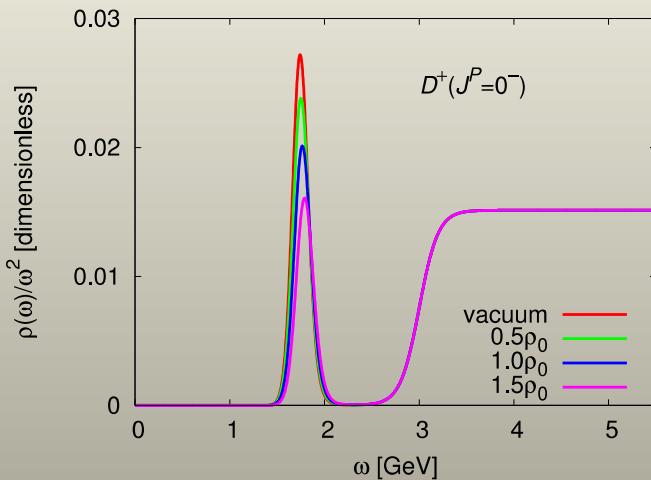


1. Multi-flavor nucleus

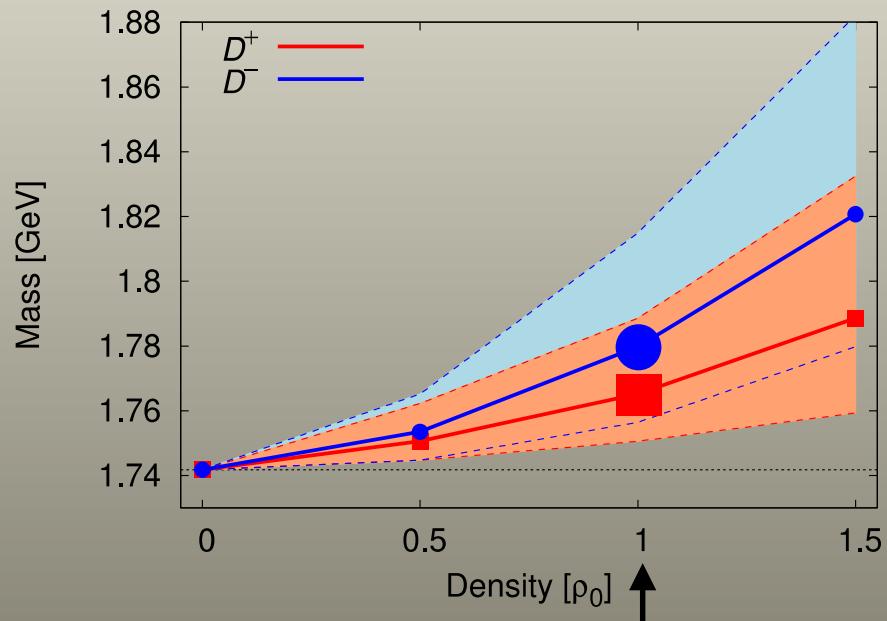
\bar{D} (B) meson in nuclear matter

Quark-meson coupling model, QCD sum rules, hadronic interactions (mean-field approach, channel coupling, pion interaction), ...

QCD sum rule



K. Suzuki, P. Gubler, M. Oka, arXiv:1511.04513 [hep-ph]



Cf. Quark confinement picture

A. Park, P. Gubler, M. Harada, S. H. Lee, C. Nonaka, W. Park, arXiv:1602.01250 [nucl-th]

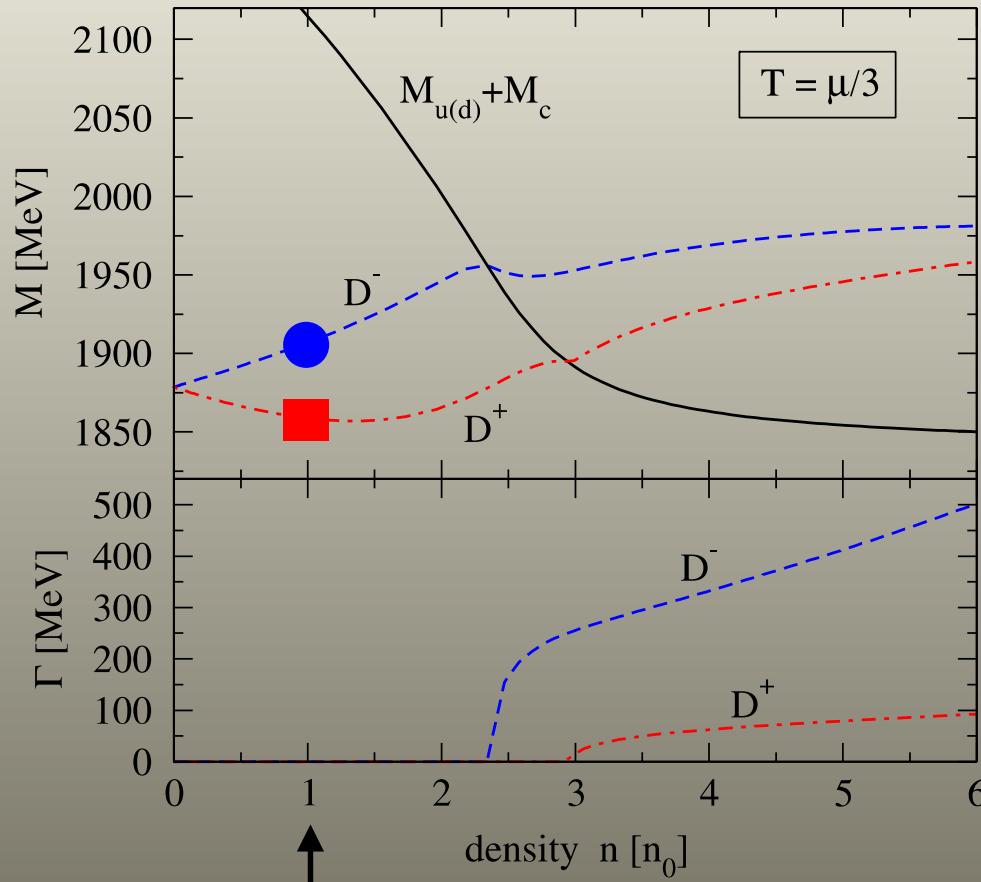
1. Multi-flavor nucleus

\bar{D} (B) meson in nuclear matter

Quark-meson coupling model, QCD sum rules, hadronic interactions (mean-field approach, channel coupling, pion interaction), ...

PNJL model

D. Blaschke, P. Costa, Y-L. Kalinovsky,
Phys. Rev. D85, 034005 (2012)

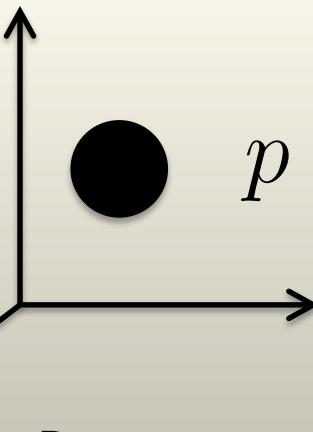


Contents

1. Multi-flavor nucleus
2. Heavy quark symmetry
3. Application 1: Leading order in $1/m_Q$ expansion
4. Application 2: Next-to-Leading order in $1/m_Q$ expansion
5. Conclusion

2. Heavy quark symmetry

Introduction to “velocity-frame”



$$p = m_Q v + k$$

$$v^2 = 1 \quad v^\mu = (1, 0, 0, 0)$$

Propagator:

$$\begin{aligned} & \frac{i}{\not{p} - m_Q + i\varepsilon} = \frac{i(\not{p} + m_Q)}{p^2 - m_Q^2 + i\varepsilon} \\ &= \frac{i(m_Q \not{v} + \not{k} + m_Q)}{(m_Q v + k)^2 - m_Q^2 + i\varepsilon} \\ &= \frac{i(m_Q \not{v} + \not{k} + m_Q)}{2m_Q v \cdot k + k^2 + i\varepsilon} \\ &= \frac{i(\not{v} + \not{k}/m_Q + 1)}{2v \cdot k + k^2/m_Q + i\varepsilon} \\ &= \frac{\not{v} + 1}{2} \frac{i}{v \cdot k + i\varepsilon} + \mathcal{O}(1/m_Q) \end{aligned}$$

2. Heavy quark symmetry

$$\mathcal{L}_{\text{heavy quark}} = \bar{Q}(iD - m_Q)Q \quad D_\mu = \partial_\mu - igA_\mu^a T^a$$

↓

1/m_Q expansion
(positive energy state Q_v with velocity v)

$$Q_v = e^{im_Q v \cdot x} \frac{1 + \gamma}{2} Q$$

Manohar, Wise, Luke, Grinstein, ...

$$\mathcal{L}_{\text{HQET}} = \boxed{\bar{Q}_v v \cdot iD Q_v} + \bar{Q}_v \frac{(iD_\perp)^2}{2m_Q} Q_v - g_s \bar{Q}_v \frac{\sigma_{\mu\nu} G^{\mu\nu}}{4m_Q} Q_v + \mathcal{O}(1/m_Q^2)$$

LO

$$\gamma^\mu \rightarrow \frac{1 + \gamma}{2} \gamma^\mu \frac{1 + \gamma}{2} = v^\mu \frac{1 + \gamma}{2} \rightarrow v^\mu$$



Heavy quark spin in $m_Q \rightarrow \infty$ is
conserved.

Heavy quark symmetry (HQS)

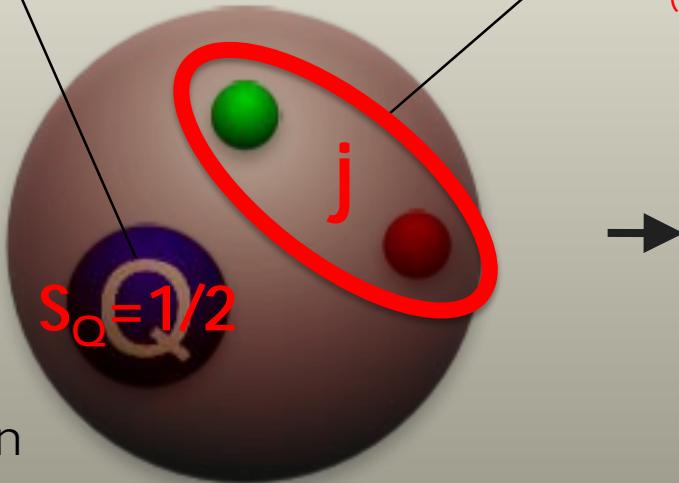
2. Heavy quark symmetry

Heavy quark symmetry (HQS)

Normal hadrons, exotic hadrons, hadronic molecules, exotic nuclei, ...

Heavy quark spin \times Light quark and gluon total spin

Brown muck
(Spin-complex)



Heavy hadron



small spin-spin ($S_Q \cdot j$) interaction
 $\sim \Lambda_{QCD}/m_Q \ll 1$

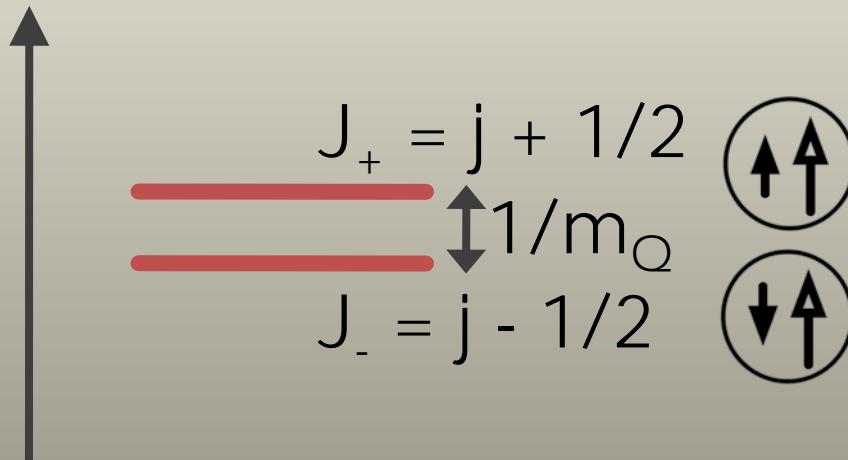
Total spin = $j \pm 1/2$: degenerate

2. Heavy quark symmetry

Heavy quark symmetry (HQS)

Normal hadrons, exotic hadrons, hadronic molecules, exotic nuclei, ...

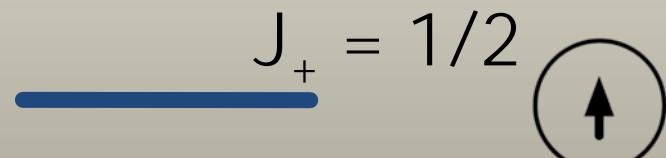
Hadron mass

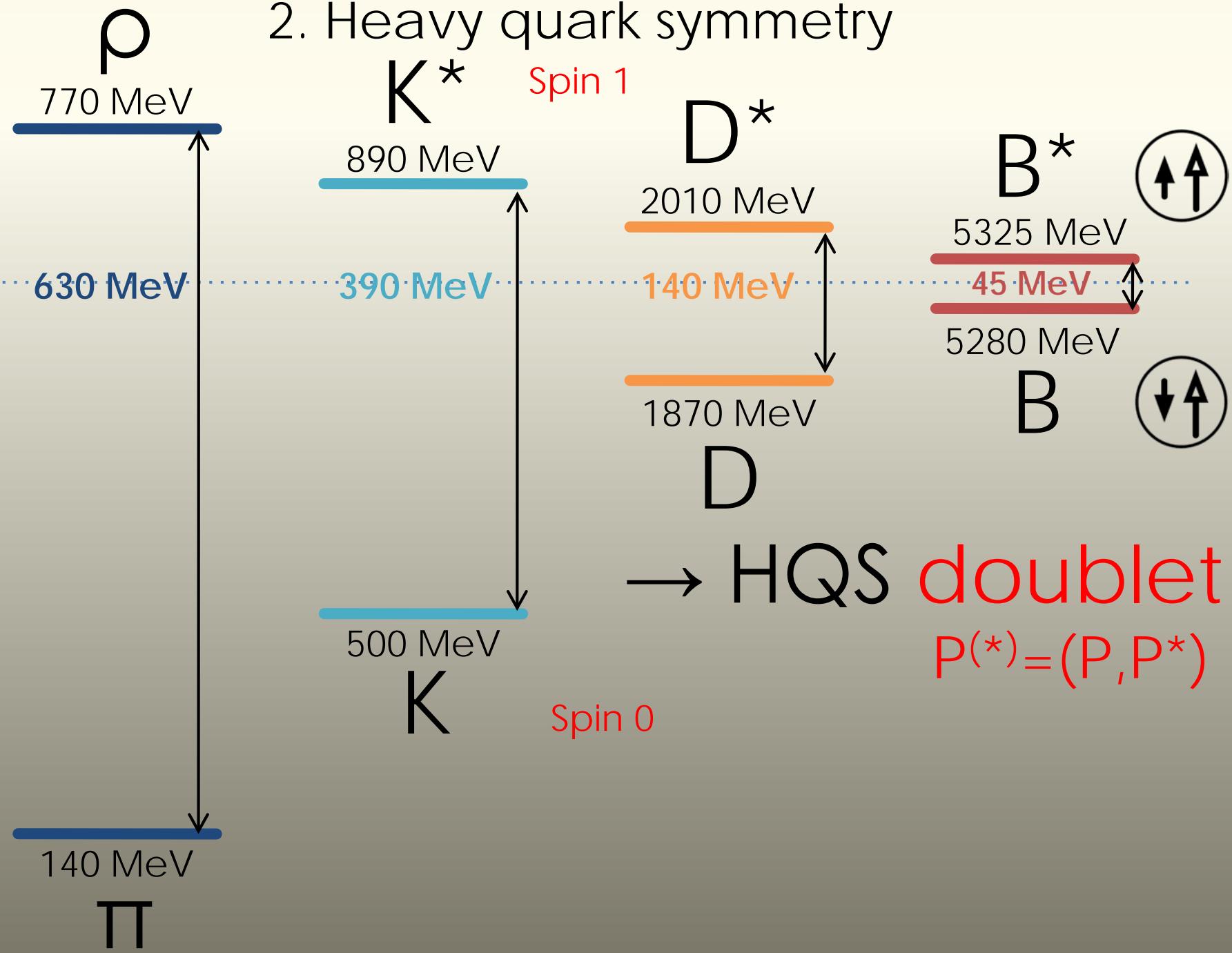


HQS doublet
($j \neq 0$)

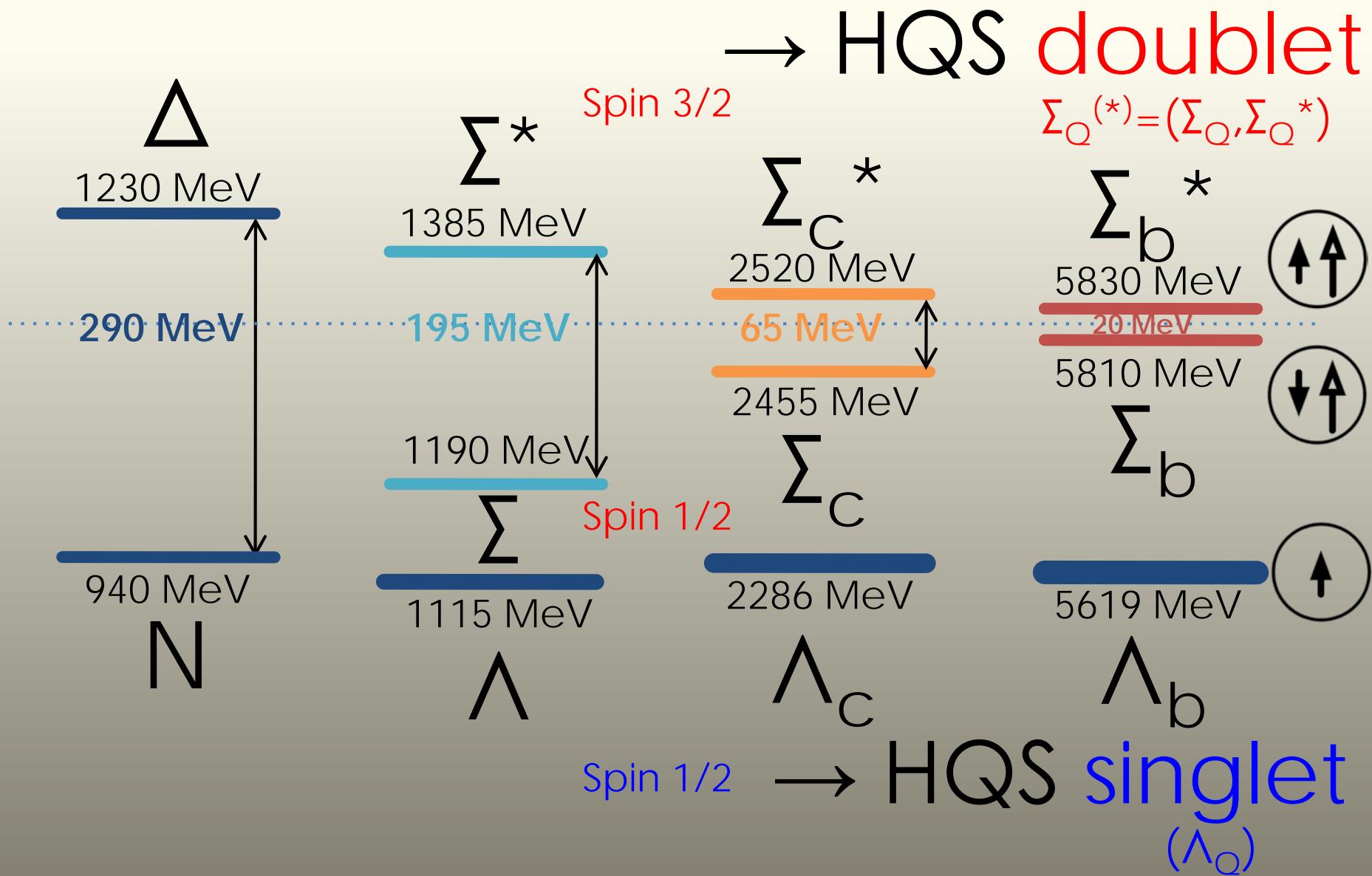
OR

HQS singlet
($j=0$)



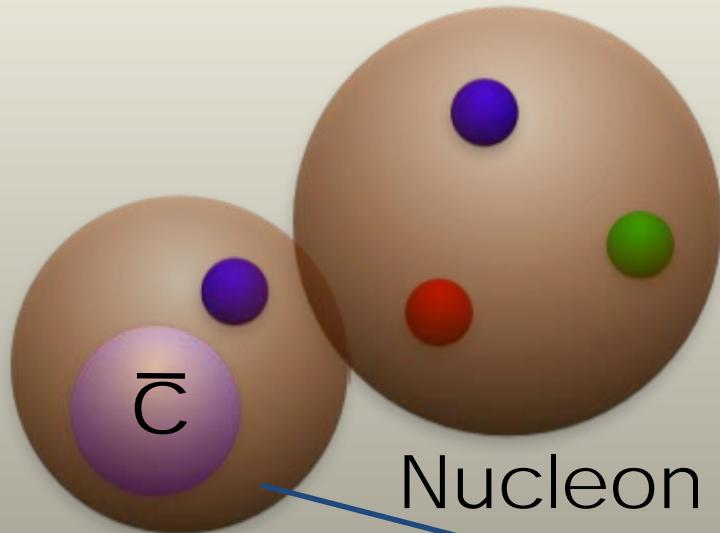


2. Heavy quark symmetry



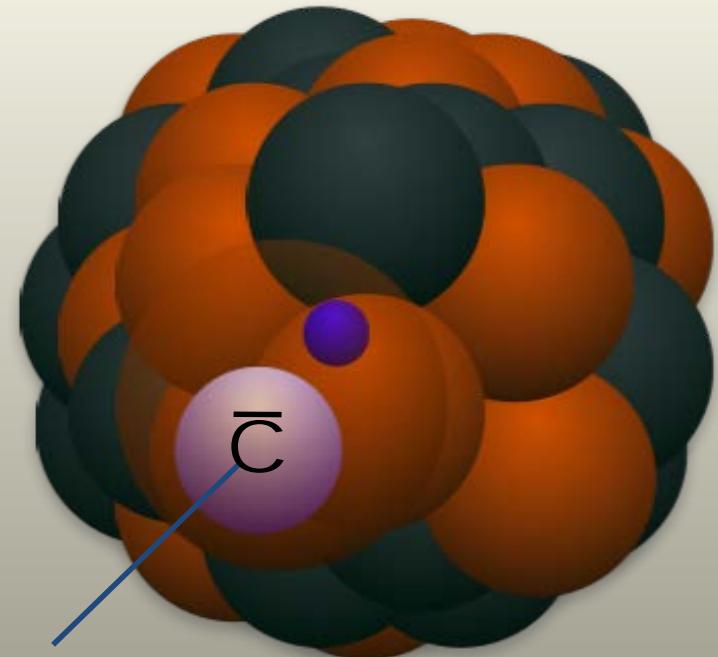
Question

Charm (bottom) exotic hadrons and nuclei



Hadronic molecule ($A=1$)

Answer



Nucleus (A : large)

HQS doublet or singlet?
HQS doublet!!

3. Application 1: Leading order in $1/m_Q$ expansion

Heavy hadron effective theory

LO in $1/m_Q$ expansion

3. Application 1: Leading order in $1/m_Q$ expansion

$Q\bar{q}$
↑ ↓ ↑

$$P(*) = P^* \uparrow \uparrow \text{ or } P \downarrow \uparrow$$

1- same mass 0-

$Q=b \quad \bar{B}^{**} \quad \bar{B}$

H.Georgi, Phys. Lett. B240, 447 (1990)

A.F.Falk, H.Georgi, B.Grinstein, Nucl. Phys. B343, 1 (1990)

3. Application 1: Leading order in $1/m_Q$ expansion

Q \bar{q}

$$H_v(x) = \frac{1 + \gamma}{2} \left[\frac{1 + \gamma}{2} P_{vv}^{**}(x) + i\gamma_5 \frac{1 + \gamma}{2} P_{vv}(x) P_v \right]$$

1- 0-

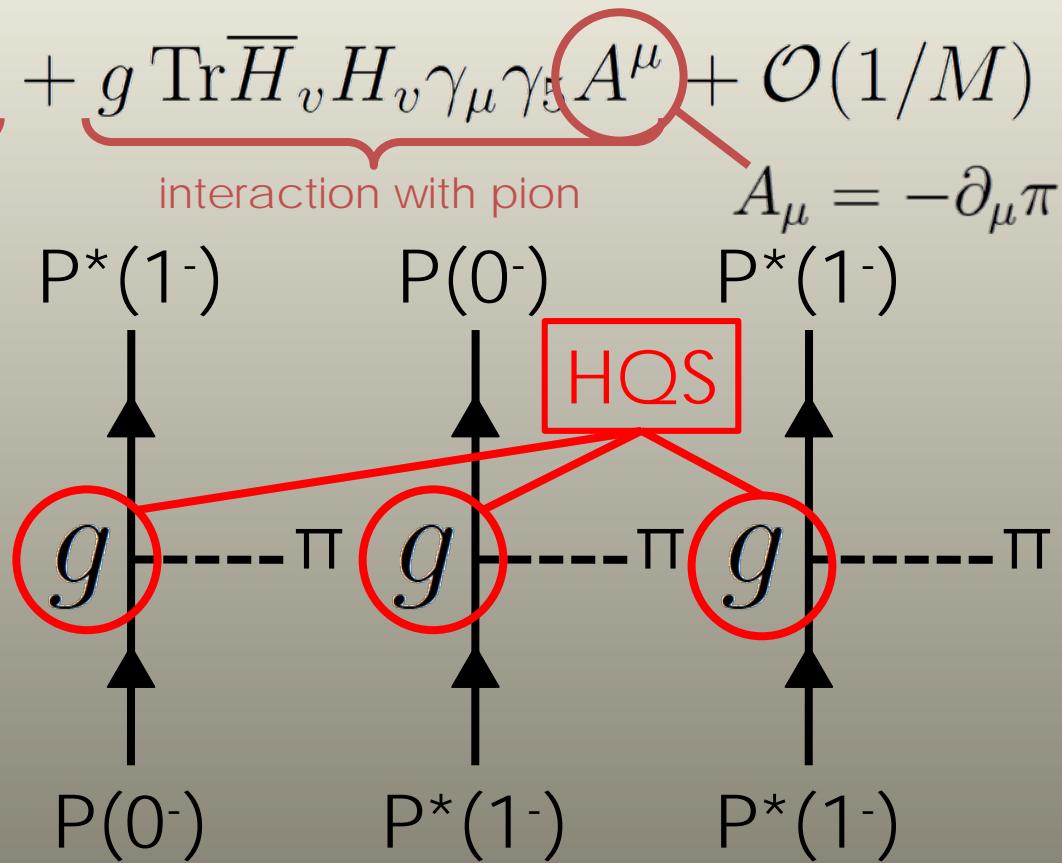
H.Georgi, Phys. Lett. B240, 447 (1990)

A.F.Falk, H.Georgi, B.Grinstein, Nucl. Phys. B343, 1 (1990)

3. Application 1: Leading order in $1/m_Q$ expansion

Interaction of $P(0^-)$, $P^*(1^-)$ meson and pion field
with chiral symmetry + HQS

$$\mathcal{L}_{\text{HMET}}^{(0)} = \underbrace{-\text{Tr} \bar{H}_v v \cdot i D H_v}_{\text{kinetic term}} + \underbrace{g \text{Tr} \bar{H}_v H_v \gamma_\mu \gamma_5 A^\mu}_{\text{interaction with pion}} + \mathcal{O}(1/M)$$

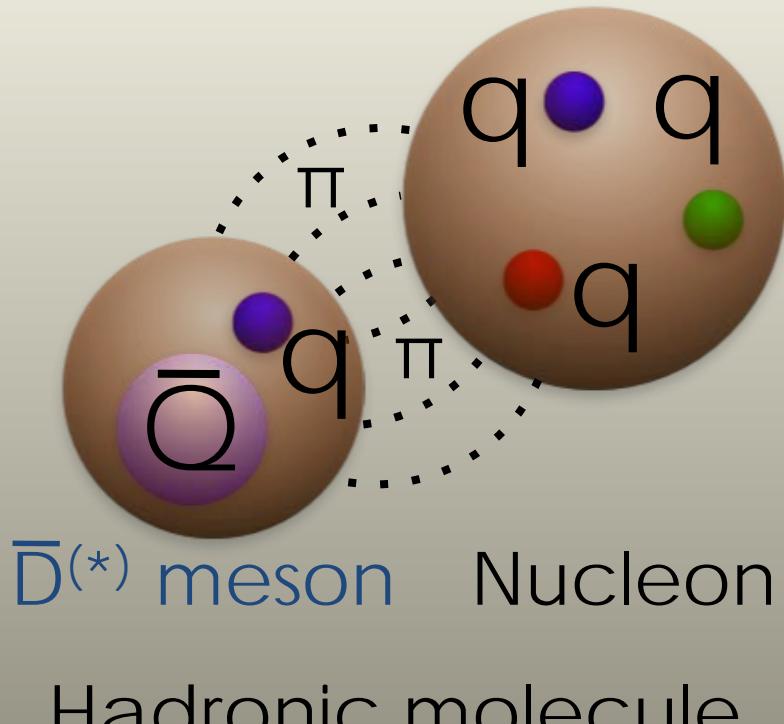


H.Georgi, Phys. Lett. B240, 447 (1990)

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3. Application 1: Leading order in $1/m_Q$ expansion

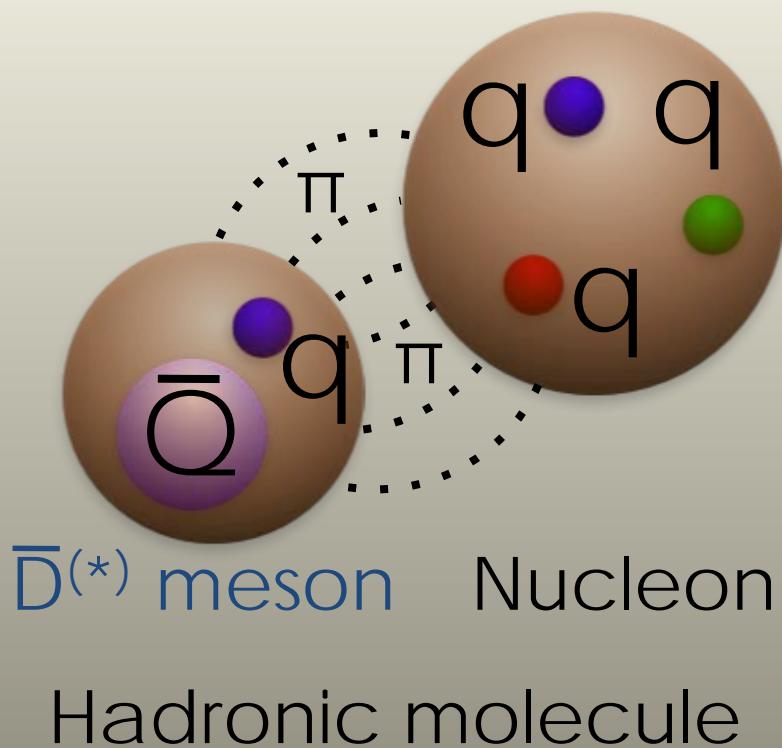
$$\bar{Q}q = \bar{P}^{(*)}$$



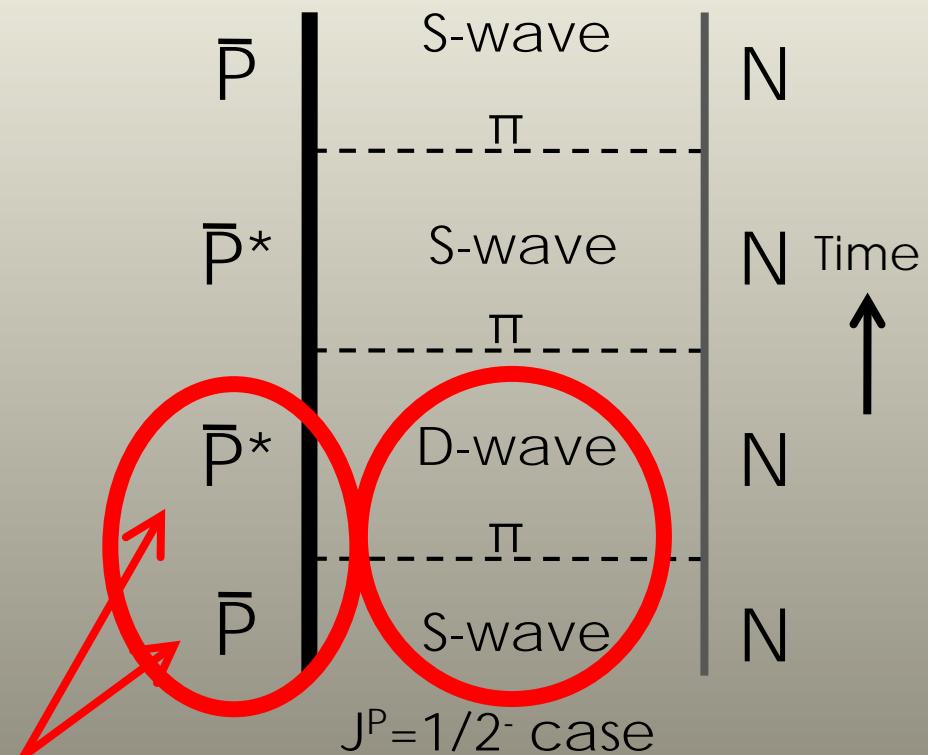
- ① No quark-antiquark annihilation
(very simple system!)
- ② Relation for interaction vertices:
 $g_{PP^*\pi} = g_{P^*P^*\pi}$ (heavy quark symmetry)

HQS doublet or singlet?

3. Application 1: Leading order in $1/m_Q$ expansion



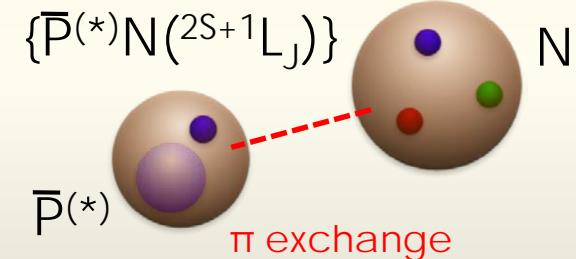
$$\bar{Q}q = \bar{P}^{(*)} \quad ;$$



$\bar{P}(0^-)-\bar{P}^*(1^-)$ tensor form factor is important !
(mixing of S-wave and D-wave, like in deuteron)

3. Application 1: Leading order in $1/m_Q$ expansion

$$J^P = 1/2^-, 3/2^-$$



$$J^P = 1/2^- : \{\bar{P}N(^2S_{1/2}), \bar{P}^*N(^2S_{1/2}), \bar{P}^*N(^4D_{1/2})\}$$

HQS singlet

$$H_{1/2^-} = \begin{pmatrix} K_0 & \sqrt{3}C & -\sqrt{6}T \\ \sqrt{3}C & K_0 - 2C & -\sqrt{2}T \\ -\sqrt{6}T & -\sqrt{2}T & K_2 + (C - 2T) \end{pmatrix} \xrightarrow{\text{Unitary transformation}} \begin{pmatrix} K_0 - 3C & 0 & 0 \\ 0 & K_0 + C & -2\sqrt{2}T \\ 0 & -2\sqrt{2}T & K_2 + (C - 2T) \end{pmatrix}$$

Unitary transformation
(Heavy quark + Brown muck)

HQS doublet
(tensor potential)

$$J^P = 3/2^- : \{\bar{P}N(^2D_{3/2}), \bar{P}^*N(^4S_{3/2}), \bar{P}^*N(^4D_{3/2}), \bar{P}^*N(^2D_{3/2})\}$$

$$H_{3/2^-} = \begin{pmatrix} K_2 & \sqrt{3}T & -\sqrt{3}T & \sqrt{3}C \\ \sqrt{3}T & K_0 + C & 2T & T \\ -\sqrt{3}T & 2T & K_2 + C & -T \\ \sqrt{3}C & T & -T & K_2 - 2C \end{pmatrix} \xrightarrow{\text{Unitary transformation}} \begin{array}{c|cc|cc} & \begin{matrix} K_0 + C & 2\sqrt{2}T \\ 2\sqrt{2}T & K_2 + (C - 2T) \end{matrix} & & & \\ \hline & 0 & 0 & K_2 - 3C & 0 \\ & 0 & 0 & 0 & K_2 + (C + 2T) \end{array}$$

Unitary transformation
(Heavy quark + Brown muck)

K_L : kinetic term

(angular momentum L)

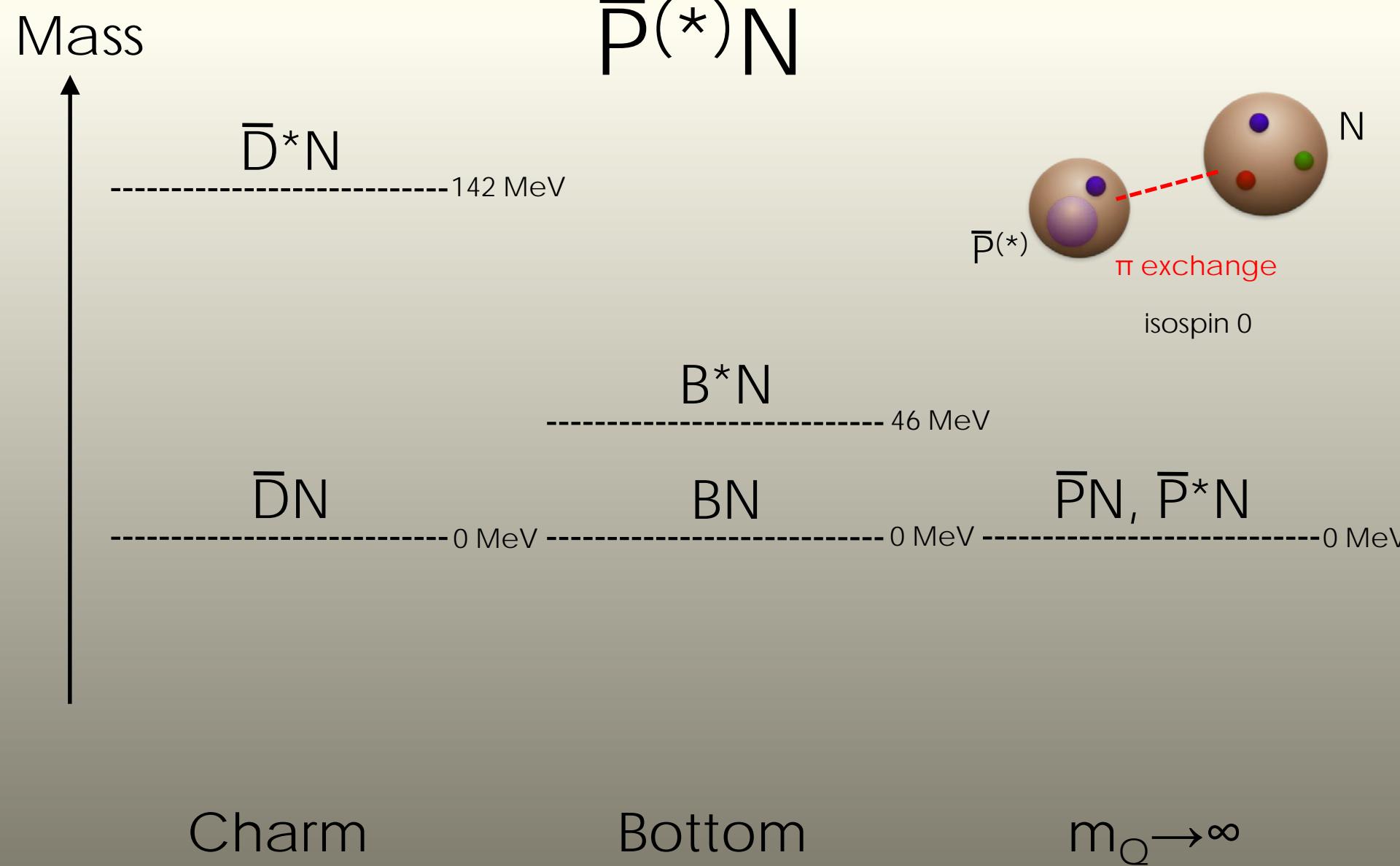
C: central potential

T: tensor potential

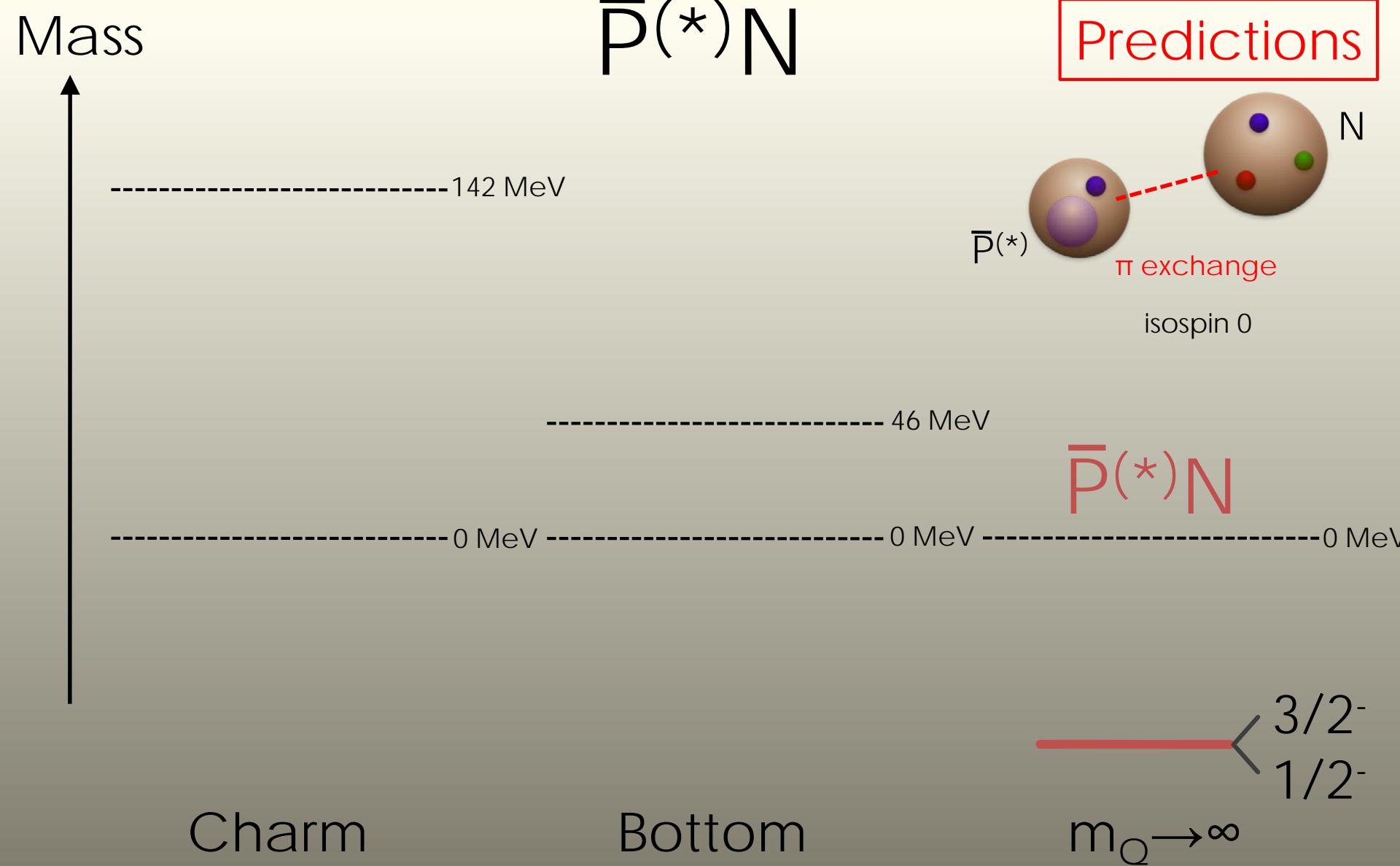
1/2⁻ and 3/2⁻ are degenerate in mass.

HQS doublet

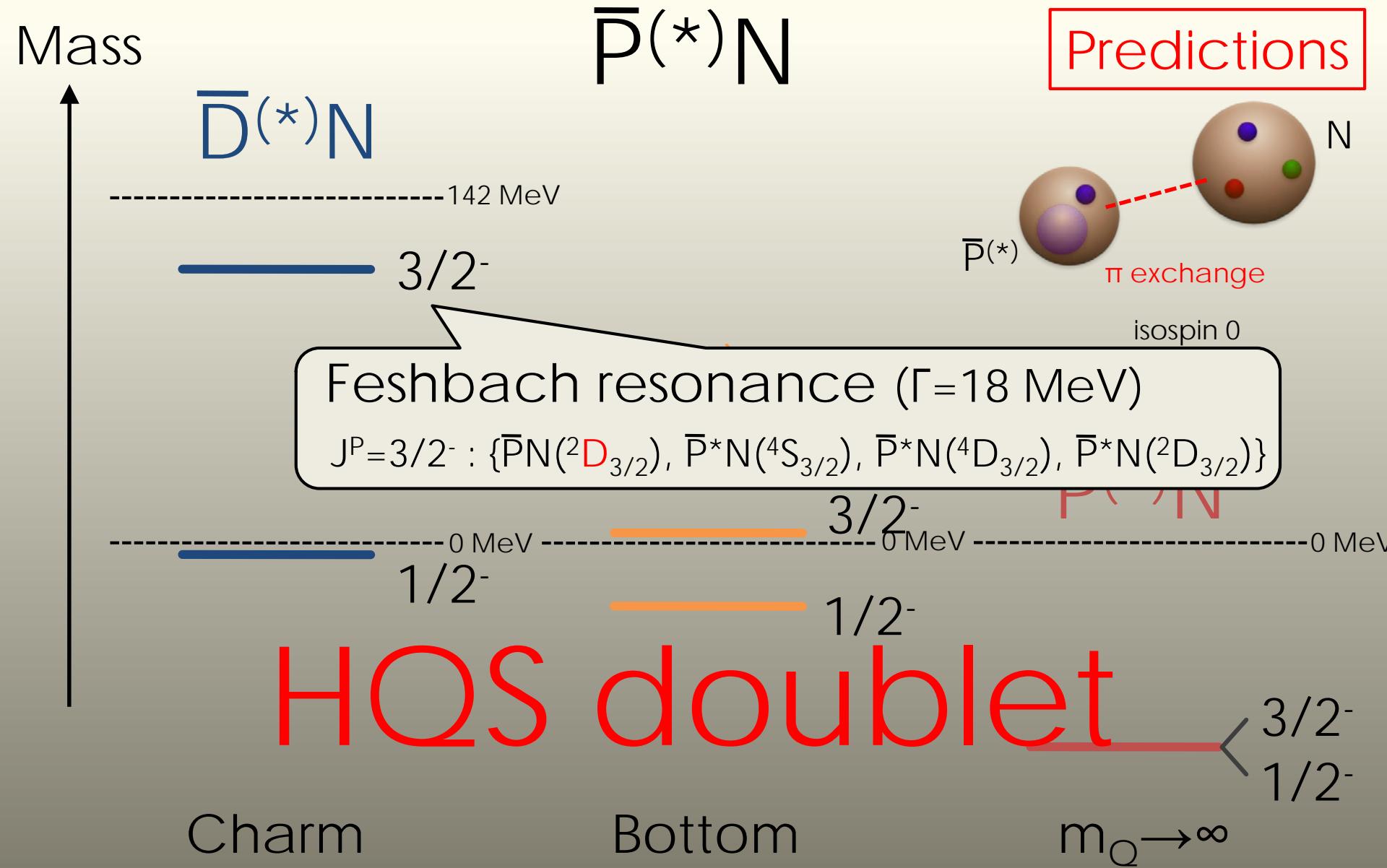
3. Application 1: Leading order in $1/m_Q$ expansion



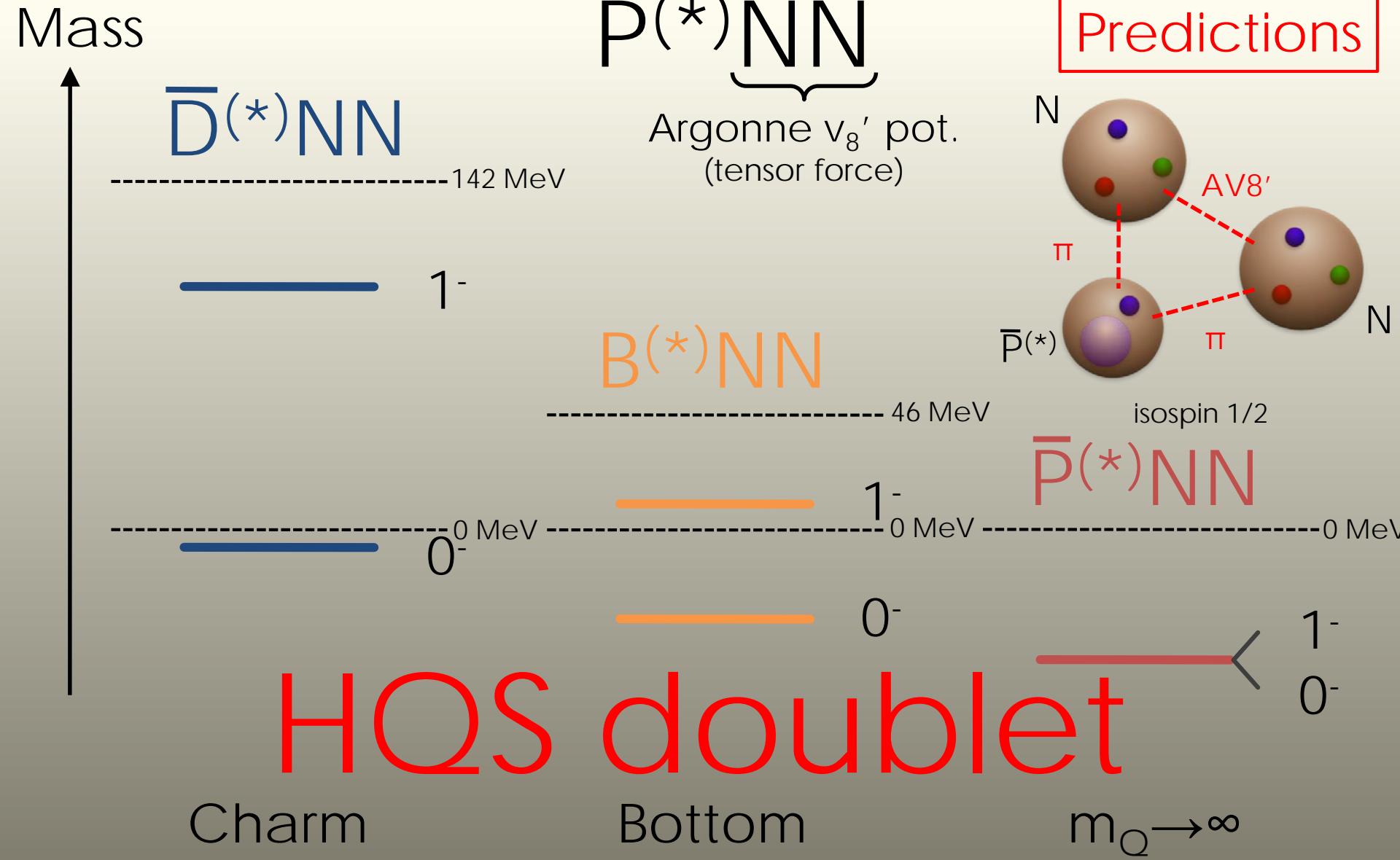
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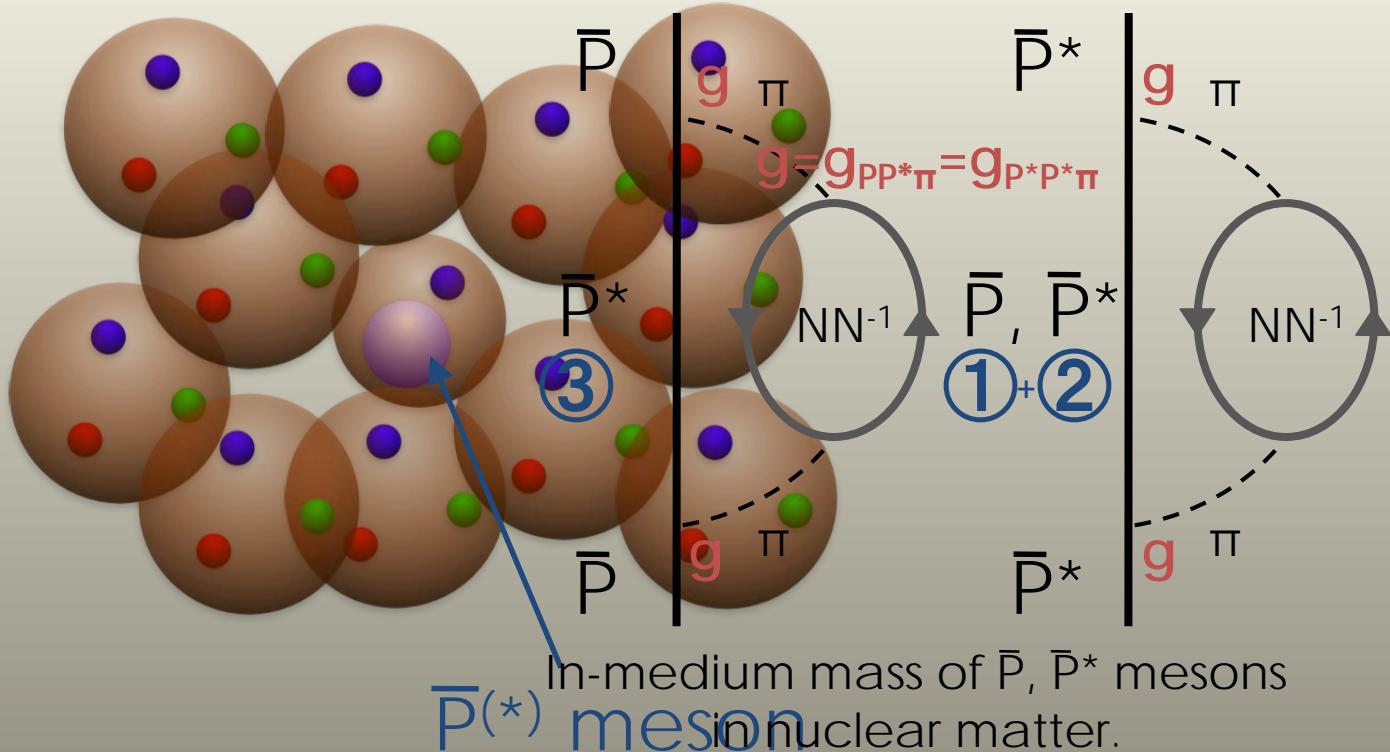


3. Application 1: Leading order in $1/m_Q$ expansion



3. Application 1: Leading order in $1/m_Q$ expansion

$\bar{P}^{(*)}$ in nuclear matter



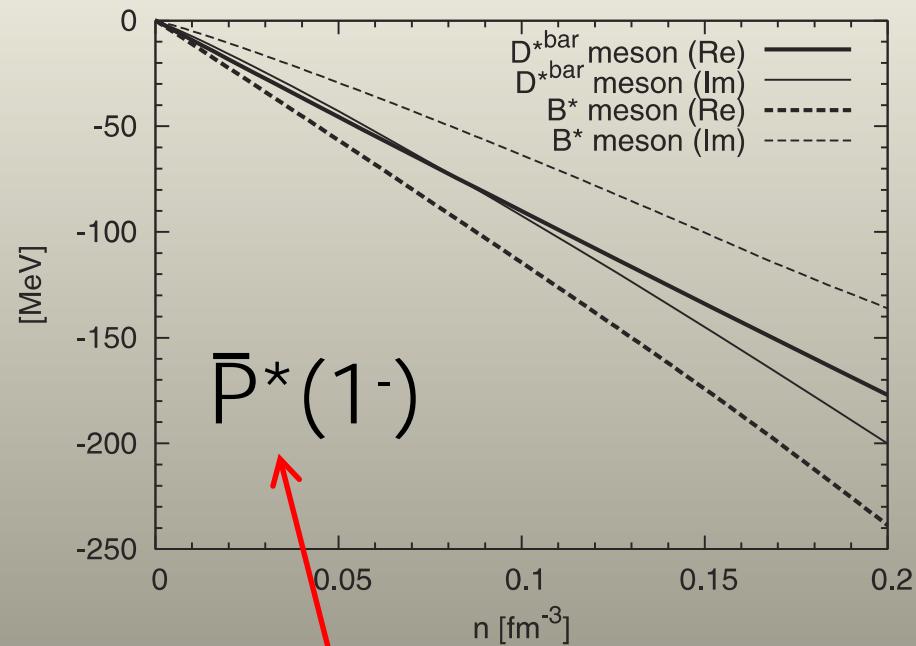
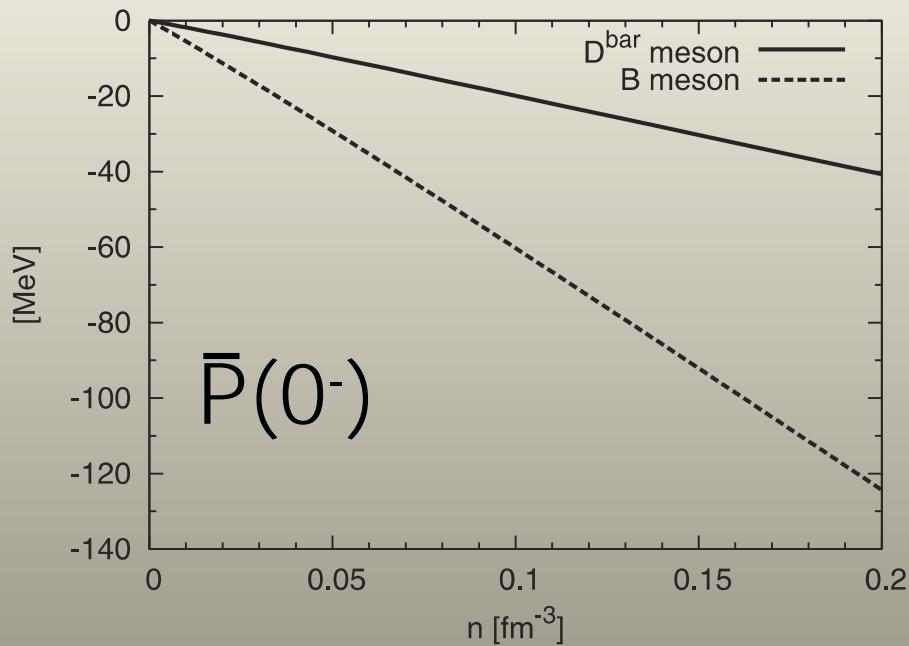
$$\text{Mass}(\bar{P}) \neq \text{Mass}(\bar{P}^*)$$

Cf. HQS singlet in "pion condensate" (Suenaga, He, Ma, Harada, PRC89, 068201 (2014))

3. Application 1: Leading order in $1/m_Q$ expansion

Numerical results for cutoff momentum 800 MeV

$\bar{P}(0^-)$, $\bar{P}^*(1^-)$ self-energy in nuclear medium



1. Imaginary parts
2. Large mass drop

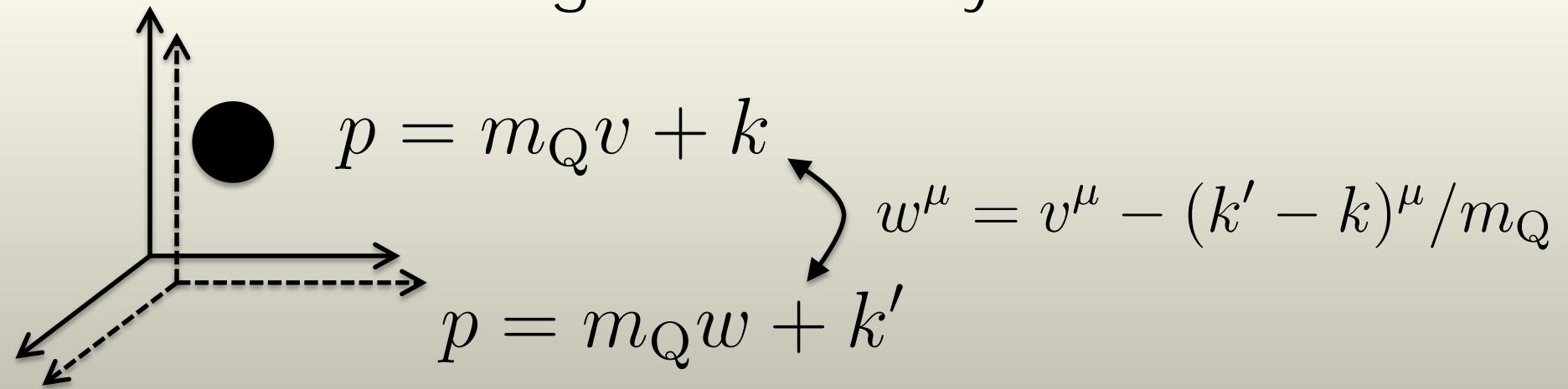
4. Appl.2: Next-to-Leading order in $1/m_Q$ expansion

$1/m_Q$ expansion

LO + NLO

4. Appl.2: Next-to-Leading order in $1/m_Q$ expansion

Change of “velocity-frame”



velocity-rearrangement

$$v \rightarrow v + q/m_Q$$

Effective field

$$Q_v \rightarrow e^{iq \cdot x} \left(1 + \frac{q}{2m_Q} \right) Q_v$$

Ex. Heavy Quark Effective Theory

$$\mathcal{L}_0 = \bar{Q}_v i v \cdot D Q_v \quad \mathcal{L}_1 = \bar{Q}_v \frac{(iD_\perp)^2}{2m_Q} Q_v - g_s \bar{Q}_v \frac{\sigma_{\mu\nu} G^{\mu\nu}}{4m_Q} Q_v$$

$$\mathcal{L}_0 \rightarrow \mathcal{L}_0 + \frac{1}{m_Q} \bar{Q}_v i q \cdot D Q_v \quad \mathcal{L}_1 \rightarrow \mathcal{L}_1 - \frac{1}{m_Q} \bar{Q}_v i q \cdot D Q_v$$

$\mathcal{L}_0 + \mathcal{L}_1$: invariant

4. Appl.2: Next-to-Leading order in $1/m_Q$ expansion

$$M_H(\rho) = m_Q + \bar{\Lambda}(\rho) \frac{1}{m_Q} \left(\frac{\lambda_1(\rho)}{2m_Q} + \frac{8\vec{S} \cdot \vec{j} \lambda_2(\rho; m_Q)}{2m_Q} \right) + \mathcal{O}(1/m_Q^2)$$

Heavy hadron mass "parametrization"

ρ : nuclear matter density

at NLO

$\frac{\lambda_1(\rho)}{m_Q} = - \langle H_{v_r}(\rho) | \bar{Q}_{v_r} g_s \vec{x} \cdot \vec{E} Q_{v_r} | H_{v_r}(\rho) \rangle$

Color electric gluon

$8\vec{S} \cdot \vec{j} \lambda_2(\rho; m_Q) = \frac{1}{2} c(\mu) \langle H_{v_r}(\rho) | \bar{Q}_{v_r} g_s \vec{\sigma} \cdot \vec{B} Q_{v_r} | H_{v_r}(\rho) \rangle$

Color magnetic gluon

"Non-perturbative" gluon

4. Appl.2: Next-to-Leading order in $1/m_Q$ expansion

Can we calculate hadron w.f.
in nuclear medium from *QCD*?

Very difficult, but...

4. Appl.2: Next-to-Leading order in $1/m_Q$ expansion

$$M_H(\rho) = m_Q + \bar{\Lambda}(\rho) - \frac{\lambda_1(\rho)}{2m_Q} + 4\vec{S} \cdot \vec{j} \frac{\lambda_2(\rho; m_Q)}{2m_Q} + \mathcal{O}(1/m_Q^2)$$

Heavy hadron effective theory

LO + NLO

$1/m_Q$ expansion in quark \Leftrightarrow $1/M$ expansion in hadron

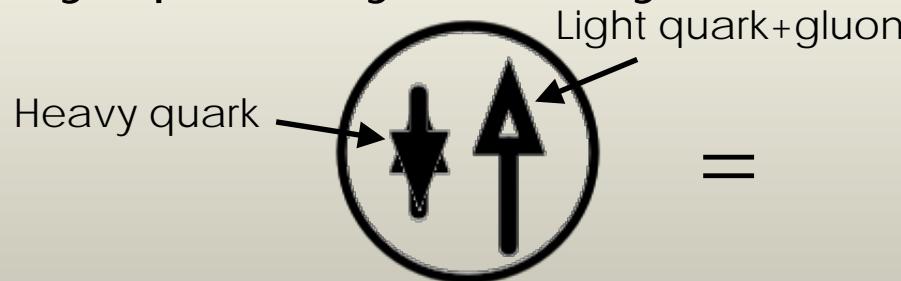
$$M = 1/m_Q + \Lambda/m_Q + \mathcal{O}(1/m_Q^2)$$

Question: what is NLO?

4. Appl.2: Next-to-Leading order in $1/m_Q$ expansion

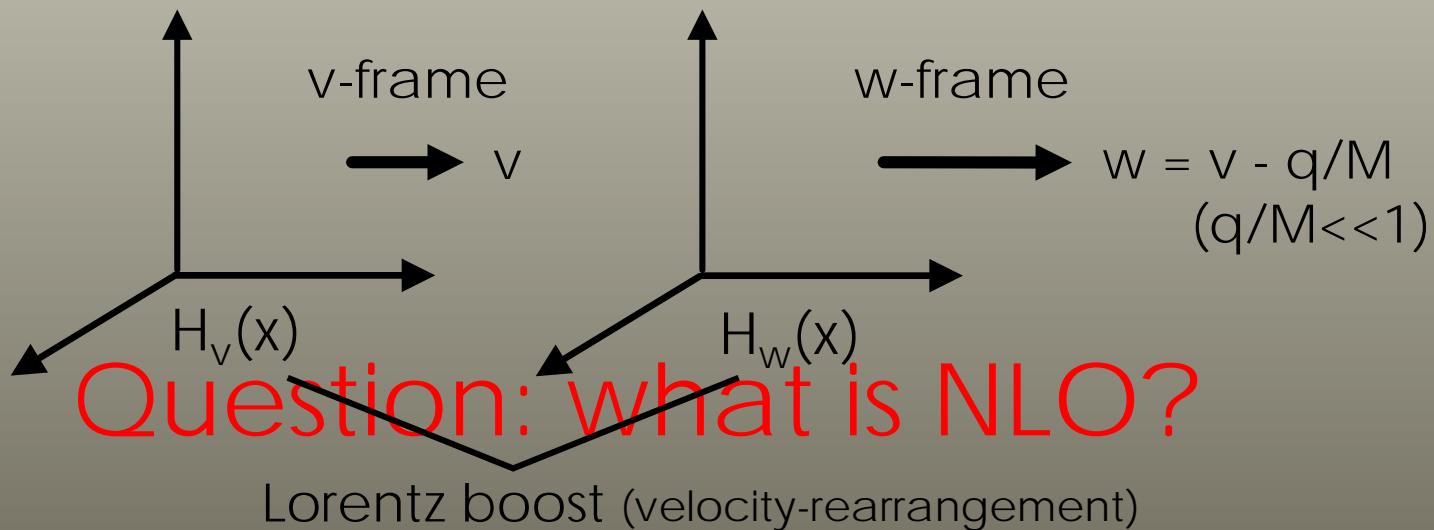
How to construct effective Lagrangian?

1. Heavy quark symmetry is conserved at $O(1)$.



2. Invariance under Lorentz boost at $O(1) + O(1/M)$.

Luke, Manohar, PLB286, 348 (1992), Kitazawa, Kurimoto, PLB323, 65 (1994)



4. Appl.2: Next-to-Leading order in $1/m_Q$ expansion

How to construct effective Lagrangian?

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Luke, Manohar, PLB286, 348 (1992), Kitazawa, Kurimoto, PLB323, 65 (1994)

3. Heavy quark symmetry breaking terms at $O(1/M)$.

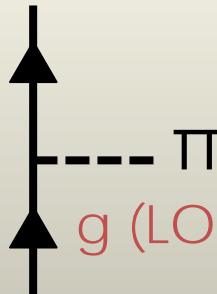


4. Appl.2: Next-to-Leading order in $1/m_Q$ expansion

Kitazawa, Kurimoto, PLB323, 65 (1994)

Heavy meson effective theory with $1/M$ corrections

$$\bar{P}(0^-), \bar{P}^*(1^-)$$



g (LO), g/M , g_1/M , g_2/M (NLO)

$$\bar{P}(0^-), \bar{P}^*(1^-)$$

$$\begin{aligned} \mathcal{L}_{\text{HMET}}^{\text{LO+NLO}} = & -\text{Tr} \bar{\mathcal{H}}_v v \cdot iD\mathcal{H}_v + \frac{\lambda}{M} \text{Tr} \bar{\mathcal{H}}_v \sigma^{\mu\nu} \mathcal{H}_v \sigma_{\mu\nu} \text{spin-flip (1/M)} \\ & + \left(g + \frac{g_1}{M} \right) \text{Tr} \bar{\mathcal{H}}_v \mathcal{H}_v \gamma_\mu \gamma_5 A^\mu + \frac{g_2}{M} \text{Tr} \bar{\mathcal{H}}_v \gamma_\mu \gamma_5 \mathcal{H}_v A^\mu \\ & + \mathcal{O}(1/M^2) \end{aligned}$$

"Covariant" field for velocity-rearrangement

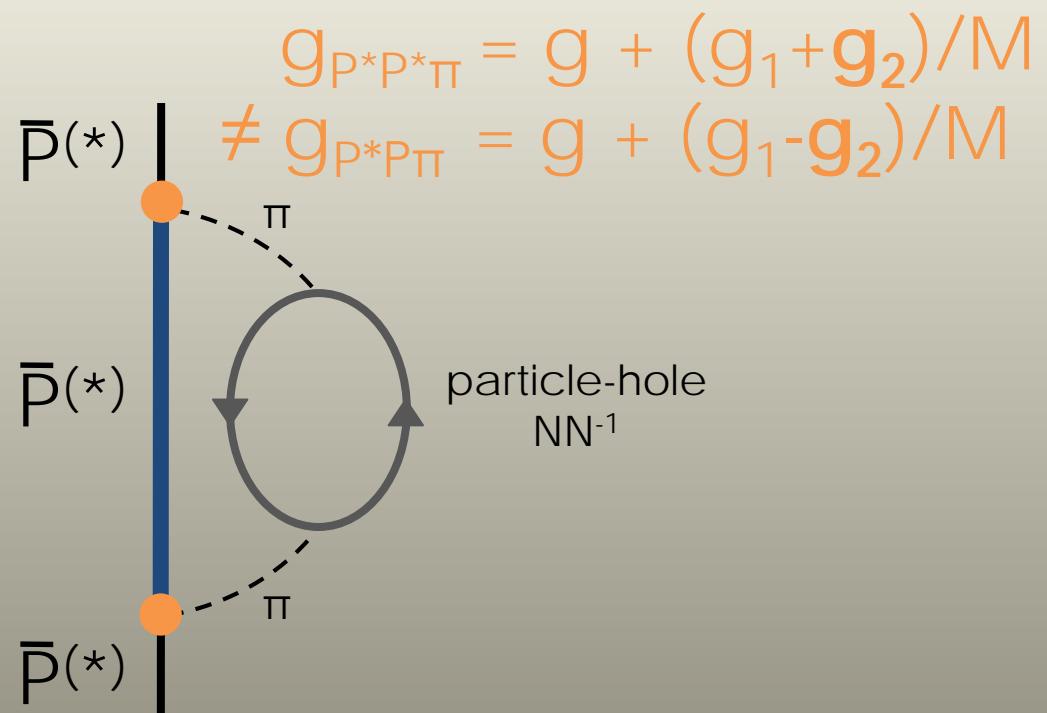
$$\mathcal{H}_v = H_v + \frac{1}{2M} (i \vec{D} H_v - H_v i \vec{D} - 2vi D H_v) + \mathcal{O}(1/M^2)$$

4. Appl.2: Next-to-Leading order in $1/m_Q$ expansion

$$M_H(\rho) = m_Q + \bar{\Lambda}(\rho) - \frac{\lambda_1(\rho)}{2m_Q} + 4\vec{S} \cdot \vec{j} \frac{\lambda_2(\rho; m_Q)}{2m_Q} + \mathcal{O}(1/m_Q^2)$$

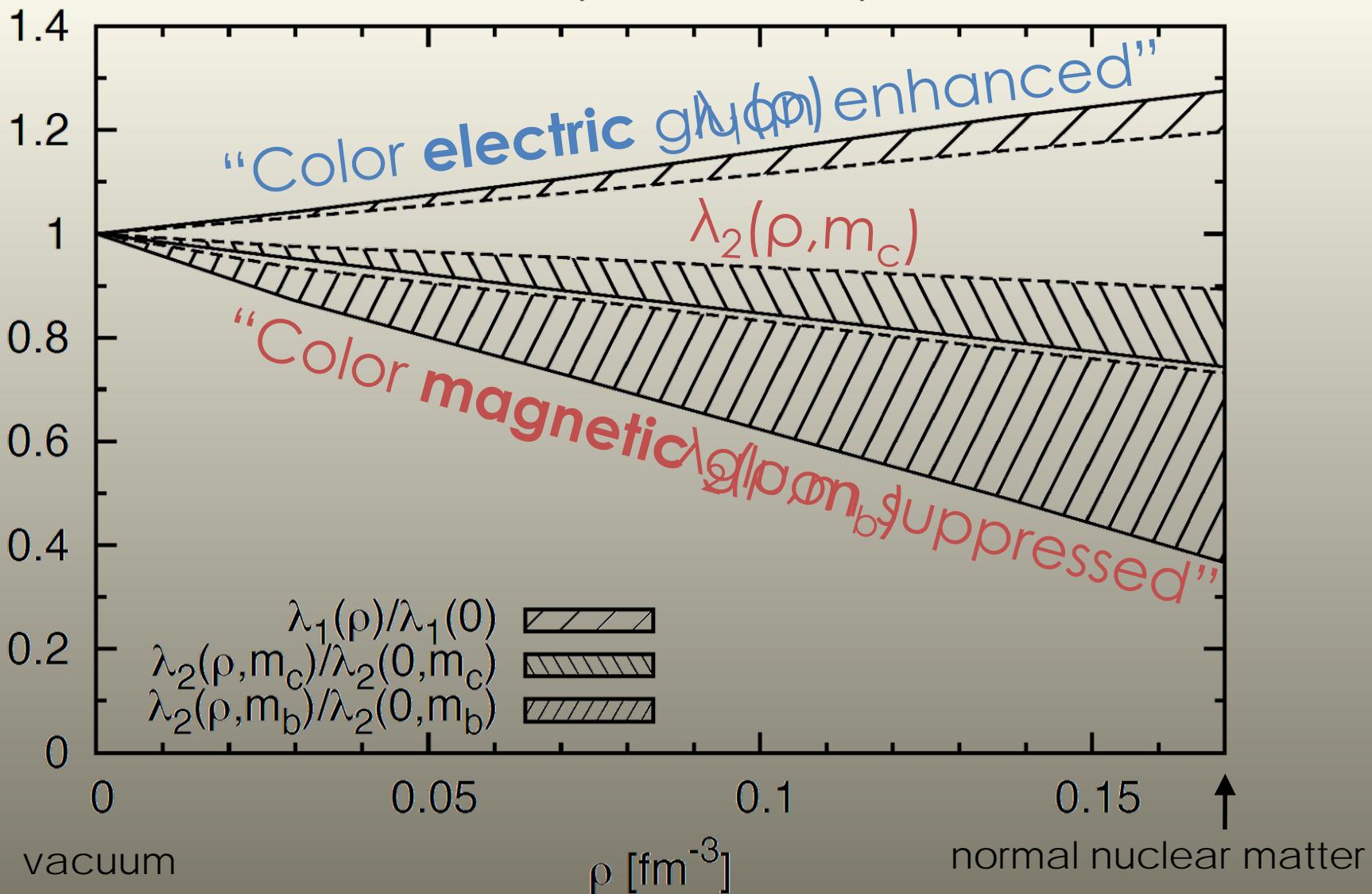
$\bar{P}^{(*)}$ in nuclear matter

$$\begin{aligned} M_{P^*} &= M - 2\lambda/M \\ \neq M_P &= M + 6\lambda/M \end{aligned}$$

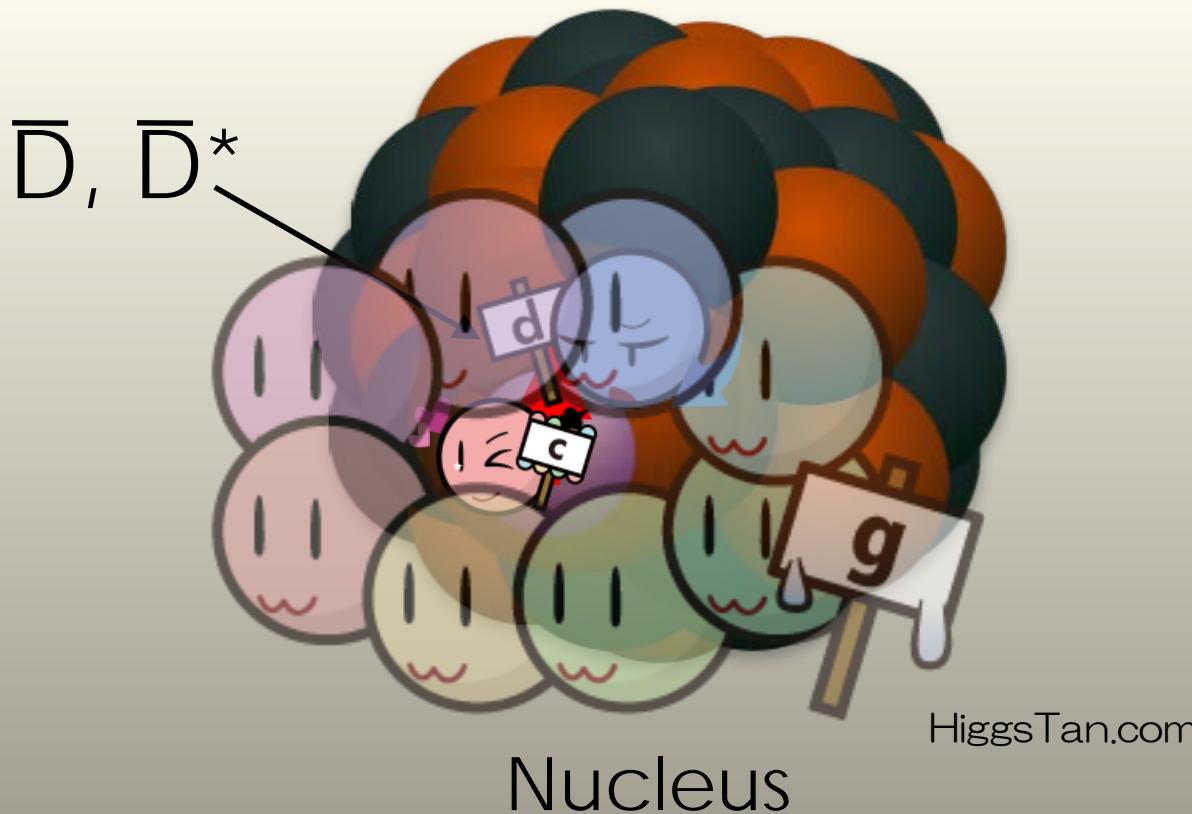


4. Appl.2: Next-to-Leading order in $1/m_Q$ expansion

$$M_H(\rho) = m_Q + \bar{\Lambda}(\rho) - \frac{\lambda_1(\rho)}{2m_Q} + 4\vec{S} \cdot \vec{j} \frac{\lambda_2(\rho; m_Q)}{2m_Q} + \mathcal{O}(1/m_Q^2)$$



4. Appl.2: Next-to-Leading order in $1/m_Q$ expansion



Charm hadron



Probing gluon fields

5. Conclusion

1. We discuss heavy quark symmetry (HQS) at LO in $1/m_Q$ expansion.
2. HQS gives doublet/singlet in mass spectrum in exotic heavy hadrons and nuclei.
3. Mass modifications and splitting of \bar{D} and \bar{D}^* mesons are concerned with gluon in medium.
4. Toward experiments at J-PARC, GSI-FAIR, LHC, RHIC, etc.!!

5. Conclusion

Many open problems...

- Relation to other fundamental aspects of QCD ?
(Dynamical breaking of chiral symmetry, Quark Confinement, ...)
- Hadron-nucleon interaction ? Many-body physics ?
(Pion exchange, Mean field approximation, Symmetry restoration, ...)
- What about atomic nucleus with charm ?
- Lattice QCD: what QCD says about charm nucleus ?
- Production cross sections of charm nucleus ?
- etc.

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