

Analysis status of $d(K^-, N) \pi \Sigma$ at J-PARC K1.8BR beam line

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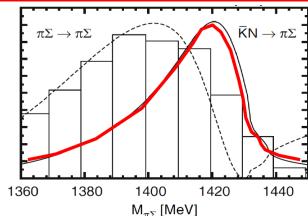
- Introduction
- K1.8BR experimental setup
- Preliminary result of E31 1st physics run
 - $d(K^-, n) \pi^\pm \Sigma^\mp$ ($I = 0, 1$) mode analysis
 - $d(K^-, p) \pi^- \Sigma^0$ ($I = 1$) mode analysis
 - $d(K^-, n) \pi^0 \Sigma^0$ ($I = 0$) mode analysis
- Summary

Nature of $\Lambda(1405)$

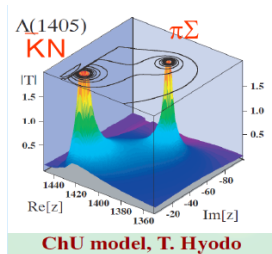
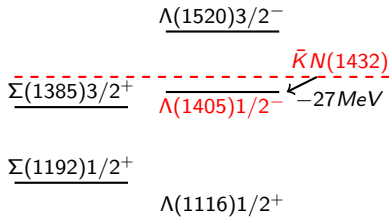
$\Lambda(1405)$

- $I(J^P) = 0\left(\frac{1}{2}^{-}\right)$
- $mass = 1405.1_{-1.0}^{+1.3} MeV$
- $\Gamma = 50.5 \pm 2.0 MeV$

C. Patrignani et al.(PDG), Chin. Phys. C, 40, 100001 (2016).



D. jido et al. Nucl Phys A725(2003), 181

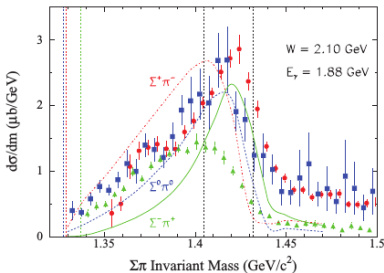


Recent experimental studies

- p/γ induced experiments

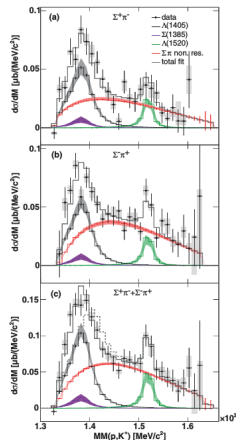
How these spectra couple to the $\bar{K}N$ pole or the $\pi\Sigma$ pole is still controversial.

$$\gamma p \rightarrow K^+ \pi^- \Sigma^+, K^+ \pi^0 \Sigma^0, K^+ \pi^+ \Sigma^-$$



CLAS collaboration: Phys Rev C87, 035206

$$pp \rightarrow K^+ p \pi^- \Sigma^+, K^+ p \pi^+ \Sigma^-$$

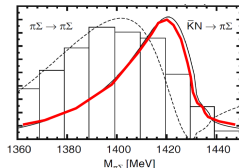
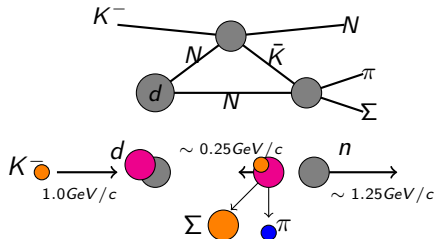


HADES collaboration:
Phys Rev C87, 025201

Kaon induced reaction is desired.

$d(K^-, N)$ reaction

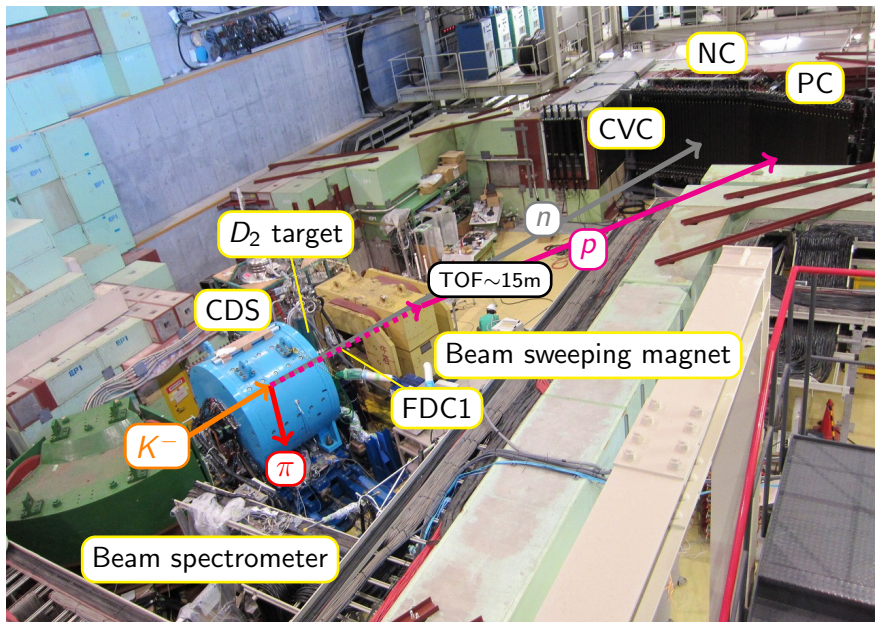
- The $d(K^-, N)$ reaction measured at $\theta_N = 0$ is expected to enhance an **S-wave** $\bar{K}N \rightarrow \pi\Sigma$ scattering even **below the $\bar{K}N$ threshold**.



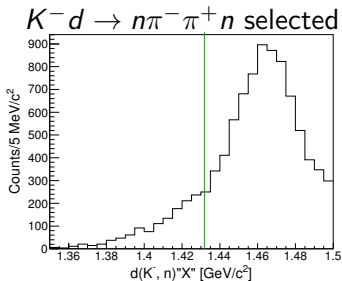
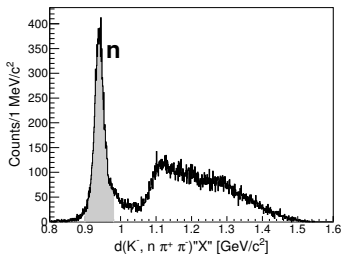
- All final states are identified to decompose the $l = 0$ and the $l = 1$ amplitudes.

$d(K^-, n) \pi^\pm \Sigma^\mp$	$l=0,1$	$\Sigma(1385)$ ($l=1$, P-wave) expected to be suppressed $\Lambda(1405)$ ($l=0$, S-wave) non-resonant ($l=1,0$) (S, P, D)
$d(K^-, n) \pi^0 \Sigma^0$	$l=0$	$\Lambda(1405)$ ($l=0$, S-wave), non-resonant (S, P, D)
$d(K^-, p) \pi^- \Sigma^0$	$l=1$	non-resonant (S, P, D) $\Sigma(1385)$ (P-wave) expected to be suppressed

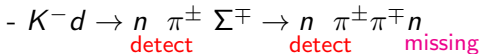
The K1.8BR experimental setup



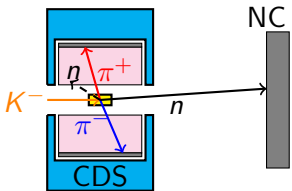
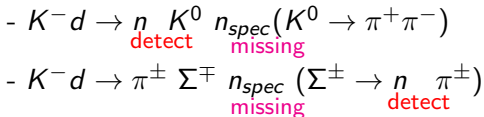
$K^- d \rightarrow n \pi^+ \pi^- n$ selection



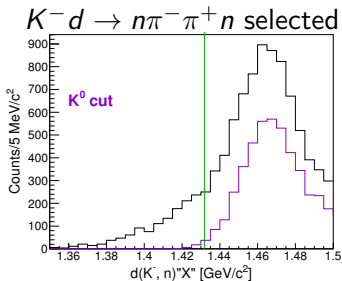
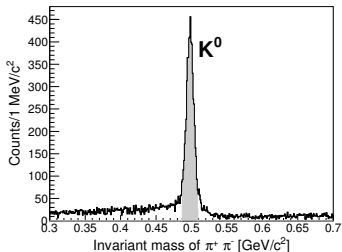
Signal



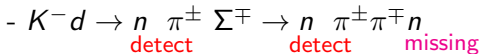
Background



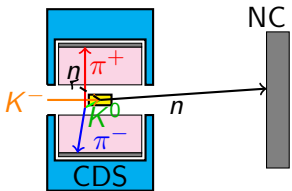
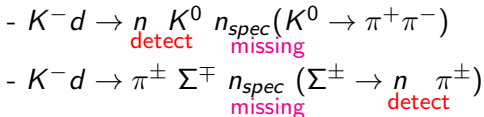
$K^- d \rightarrow n K^0 n_{spec}$ rejection



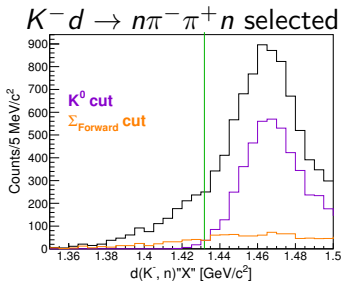
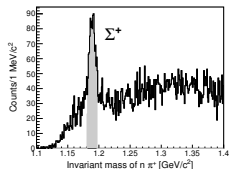
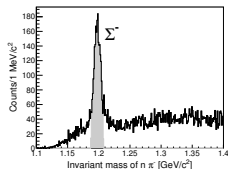
Signal



Background



$K^- d \rightarrow \pi^\pm \Sigma^\mp n_{spec}$ rejection



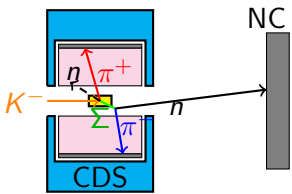
Signal

$K^- d \rightarrow n \pi^\pm \Sigma^\mp \rightarrow n \pi^\pm \pi^\mp n$
 detect detect missing

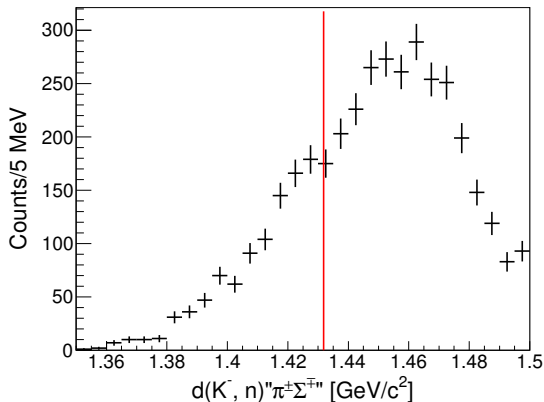
Background

$K^- d \rightarrow n K^0 n_{spec} (K^0 \rightarrow \pi^+ \pi^-)$
 detect missing

$K^- d \rightarrow \pi^\pm \Sigma^\mp n_{spec} (\Sigma^\pm \rightarrow n \pi^\pm)$
 missing detect



$d(K^-, n) \pi^\pm \Sigma^\mp$ spectrum



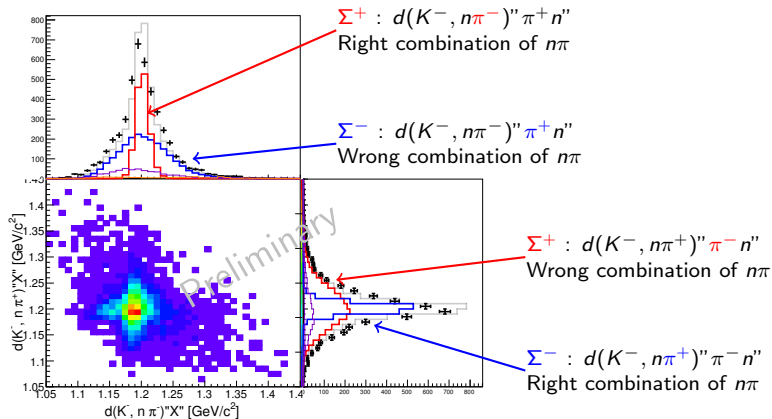
K^0 and forward-going Σ are removed.

This spectrum includes both $\pi^+\Sigma^-$ and $\pi^-\Sigma^+$ modes.

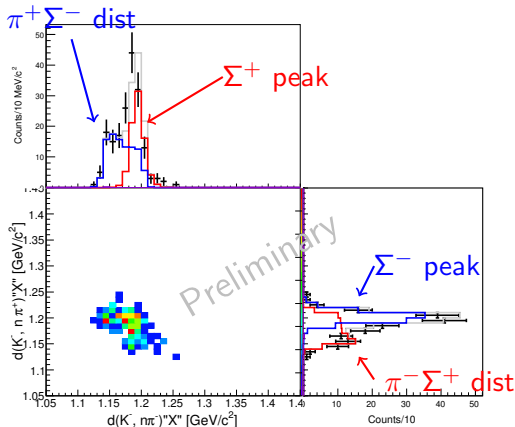
→ These modes should be separated.

$\pi^- \Sigma^+$ and $\pi^+ \Sigma^-$ identification

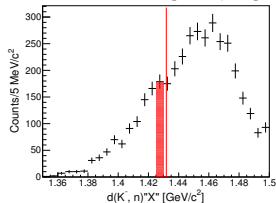
Missing mass distributions of $d(K^-, n\pi^\pm)$ are estimated by MC simulations in order to separate the $\pi^\pm \Sigma^\mp$ modes.



Fitting for $\pi^-\Sigma^+$ and $\pi^+\Sigma^-$ separation



Example $d(K^-, n)X$
1.425 ~ 1.430 [GeV/c²]

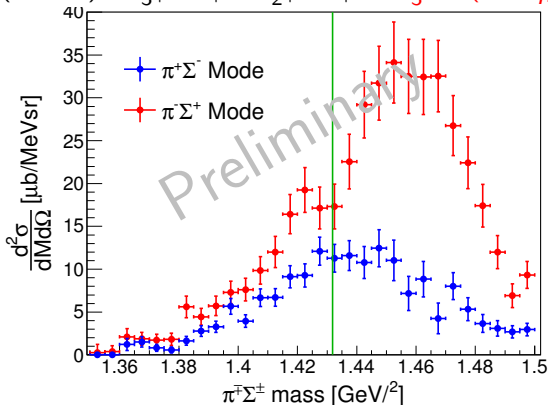


Fittings are done bin by bin
Free parameters

- 1 Number of $\pi^+\Sigma^-$ events.
- 2 Number of $\pi^-\Sigma^+$ events.

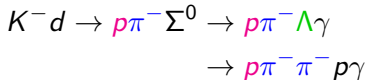
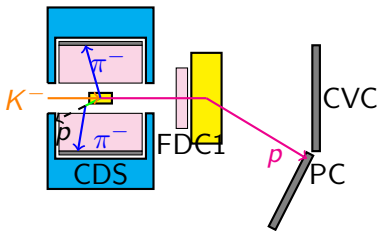
Cross section of $\pi^- \Sigma^+$ and $\pi^+ \Sigma^-$

$$(\pi^\pm \Sigma^\mp) \sim \frac{1}{3} |f_{l=0}|^2 + \frac{1}{2} |f_{l=1}|^2 \pm \frac{\sqrt{6}}{3} \text{Re}(f_{l=0} f_{l=1}^*)$$

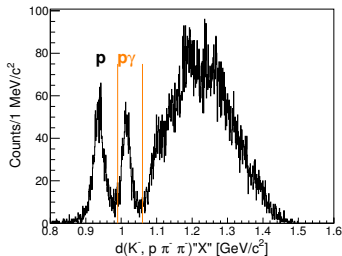
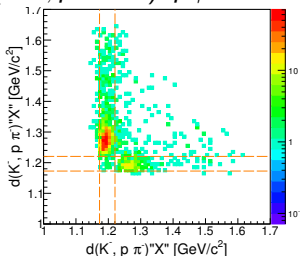


Interference between the $l = 0$ and 1 amplitudes of the $\pi\Sigma$ scattering is observed.

$d(K^-, p)\pi^-\Sigma^0$ event selection

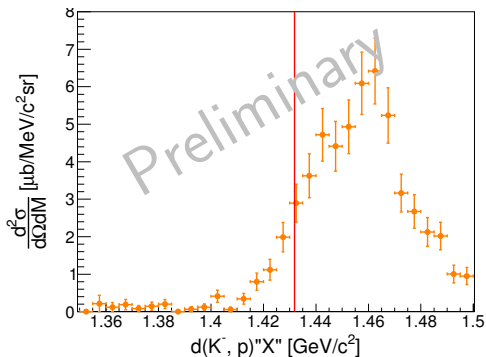


$d(K^-, p\pi^-\pi^-)p\gamma$ selected



$d(K^-, p)\pi^-\Sigma^0$ mode is clearly identified by selecting $d(K^-, p\pi^-\pi^-)p\gamma$ and $d(K^-, p\pi^-)\Sigma^0$.

Cross section of $d(K^-, p) \pi^- \Sigma^0$

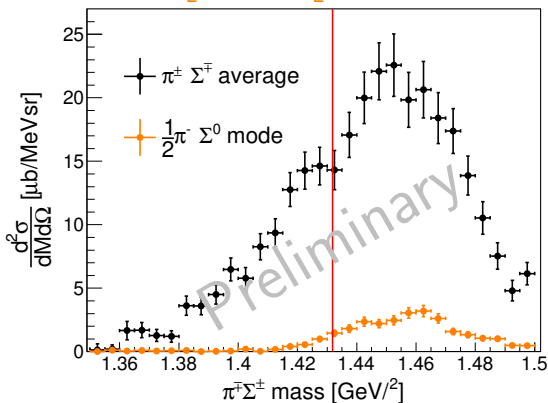


$d(K^-, p) \pi^- \Sigma^0$	$l=1$	non-resonant (S, P, D) $\Sigma(1385)$ (P-wave) to be suppressed
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Comparison $\pi^\pm \Sigma^\mp$ ($I = 0, 1$) and $\pi^- \Sigma^0$ ($I = 1$)

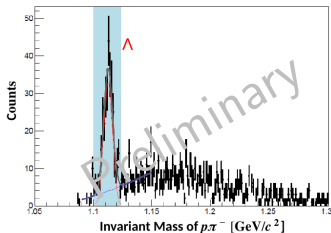
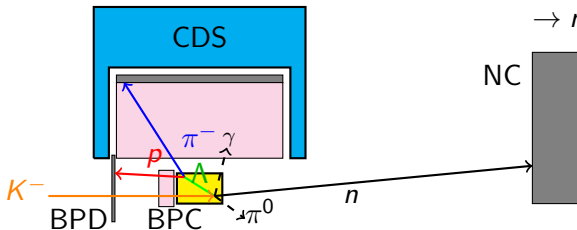
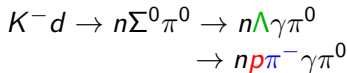
$$\frac{1}{2} (\pi^- \Sigma^+ + \pi^+ \Sigma^-) \sim \frac{1}{3} |f_{I=0}|^2 + \frac{1}{2} |f_{I=1}|^2$$

$$\frac{1}{2} \pi^- \Sigma^0 \sim \frac{1}{2} |f_{I=1}|^2$$

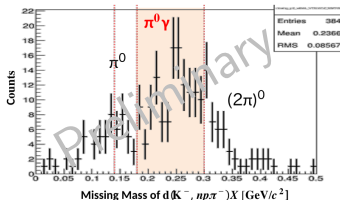


The $I = 0$ amplitude is dominant.

$d(K^-, n) \pi^0 \Sigma^0 (I=0)$ event selection

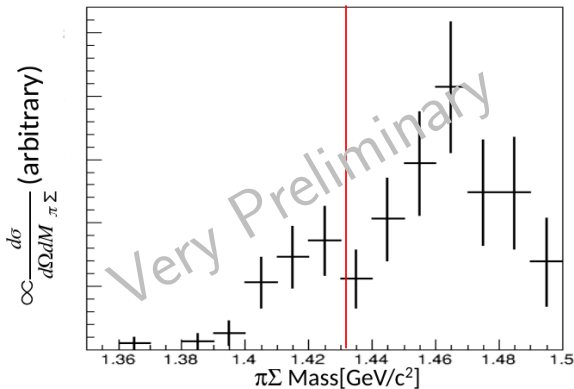


Backward Λ tagged



$d(K^-, n) \pi^0 \Sigma^0$ mode is clearly identified by selecting Λ and $d(K^-, n \Lambda) \pi^0 \gamma$.

$d(K^-, n) \pi^0 \Sigma^0$ ($I = 0$) spectrum



The $d(K^-, n) \pi^0 \Sigma^0$ spectrum is observed, which seems similar to those obtained in the $\pi^\pm \Sigma^\mp$ modes.

Statistics will be increased further in coming beam time in June, 2017.

- The preliminary result of the E31 1st physics run is presented.
 - $d(K^-, n) \pi^\pm \Sigma^\mp$ ($l = 0, 1$) spectra is presented.
The difference of two spectra is observed due to interference of the $l = 0$ and 1 amplitudes.
 - Comparison of $d(K^-, n) \pi^\pm \Sigma^\mp$ and $d(K^-, p) \pi^- \Sigma^0$
The $l=0$ amplitude is dominant.
 - $d(K^-, n) \pi^0 \Sigma^0$ ($l = 0$) spectra is presented.
The $d(K^-, n) \pi^0 \Sigma^0$ spectrum is observed.
- The result of the E31 1st physics run will be finalized.
→ $d(K^-, n) \pi^\pm \Sigma^\mp$ and $d(K^-, p) \pi^- \Sigma^0$.
- The E31 2nd physics run will be performed in June, 2017.
→ Statistics of $d(K^-, n) \pi^0 \Sigma^0$ will be increased.