# Quark matter in compact stars: implication from recent observations

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### **Experimental constraint:**

## Pulsars:

- 2. the cyclotron resonance features in the X-Ray spectrum of pulsars → surface magnetic field @ [Revnivtsev\_Mereghetti2014\_SpaceSciRev0 14-1]
- 3. photon spectrum of pulsars → surface temperature [Page2006\_NPA777-497, Ho\_Heinke 2009\_Nature462-71], mass-radius probability [Steiner\_Lattimer \_Brown2010\_ApJ722-33, Guillot\_Servillat \_Webb\_Rutledge2013\_ApJ772-7, Ozel\_Freire2016\_AnnuRevAstro54-401...]
- 4. orbital motion of a binary system → precise mass measurement [Demorest et al. 2010\_Nature467-1081, Antoniadis et al. 2013\_Science340-6131, ...]
- 5. gravitational waves emitted from the merger of binary compact stars → EoS [LIGO2009\_RepProgPhys72-076901, Hotokezaka et al.2011\_PRD83-124008...]



#### Baryonic matter

Shen EoS2&3: RMF model with the effective interaction TM1 and  $\Lambda$ -meson coupling determined by [Shen\_Toki\_Oyamatsu\_Sumiyoshi2011\_ApJ197-20]:  $M_{\Lambda} = 1115.7 \text{ MeV}, g_{\omega}^{\Lambda}/g_{\omega} = 2/3, \text{ and } g_{\sigma}^{\Lambda}/g_{\sigma} = 0.621.$ Sugahara\_Toki1994\_NPA579-557 1. the saturation densit $\rho_0$  (fm<sup>-3</sup>) 0. 145 2. the energy: E/A (MeV) -16. 3 3. the incompressibility (MeV) 281 4. the symmetry energy:  $a_{sym}$  (MeV) 36. 9

VM EoS: Cluster variational method with the potential AV18 + UIX [Togashi\_Hiyama\_Yamamoto\_Takano2016\_PRC93-035808, Togashi\_ Nakazato\_Takehara\_Yamamuro\_Suzuki\_Takano2017\_NPA961-78]:

1.	<b>t</b> he	saturation	$densit  ho_0$ (fm <sup>-3</sup> )	0.160
2.	the	energy:	E/A (MeV)	-16.1
3.	the	incompress	ibility <b>k (MeV)</b>	245
4.	the	symmetry e	nergy: <b>a<sub>sym</sub> (MeV)</b>	30.0

### **Baryon Stars**



Dashed curves: The EoSs of nuclear matter are obtained based on the non-relativistic G-matrix approach.

APR: without hyperons [Akmal\_Pandharipande\_Ravenhall1998\_PRC58-1804];

AV18+TBF: obtained with the realistic nucleon-nucleon interaction AV18, the Urbana three nucleon interaction and the Nijmegen softcore nucleon-hyperon potential [Baldo\_Burgio\_Schulze2000\_PRC61-055801]; TNI2u: considering three hyperon interaction [Takatsuka\_Nishizaki\_Yamamoto\_Tamagaki2006\_PTP115-355];

### Quark matter

# [Fragaa Romatschke 2005 PRD7d 105014] tial density up to

with 
$$\omega_i^{\text{pt}} = \sum_{i=1}^{\infty} \left( \omega_i^0 + \omega_i^1 \alpha_s \right)$$
  
with  $\omega_i^0 = -\frac{g_i m_i^4}{24\pi^2} \left[ u_i v_i \left( u_i^2 - \frac{5}{2} \right) + \frac{3}{2} \ln(u_i + v_i) \right],$   
 $\omega_i^1 = \frac{g_i m_i^4}{12\pi^3} \left\{ \left[ 6 \ln \left( \frac{\bar{\Lambda}}{m_i} \right) + 4 \right] \left[ u_i v_i - \ln(u_i + v_i) \right] + 3 \left[ u_i v_i - \ln(u_i + v_i) \right]^2 - 2v_i^4 \right\}.$   
Here we have defined  $u_i = \mu_i / m_i$  and  $v_i = \sqrt{u_i^2 - 1}.$   
 $\alpha_s(\bar{\Lambda}) = \frac{1}{\beta_0 L} \left( 1 - \frac{\beta_1 \ln L}{\beta_0^2 L} \right),$  where  $L = 2 \ln \left( \frac{\bar{\Lambda}}{\Lambda_{\text{MS}}} \right)$  and  $\Lambda_{\overline{\text{MS}}}$  is  $m_i(\bar{\Lambda}) = \hat{m}_i \alpha_s^{\frac{\gamma_0}{\beta_0}} \left[ 1 + \left( \frac{\gamma_1}{\beta_0} - \frac{\beta_1 \gamma_0}{\beta_0^2} \right) \alpha_s \right],$  the  $\overline{\text{MS}}$  renormalization point.  
The renormalization scale  $= \frac{C}{3} \sum_{i=u,d,s} \mu_i$   
with  $C = 1 \sim 4.$   
 $\Lambda_{\overline{\text{MS}}} = 376.9 \text{ MeV}, \ \hat{m}_u = 3.8 \text{ MeV}, \ \hat{m}_d = 8 \text{ MeV}, \text{ and } \hat{m}_s = 158 \text{ MeV}.$ 

### Non-perturbative corrections

at µ

We introduce an extra bag constant B to take into account the energy difference between the physical and perturbative vacua:

$$\Omega = \Omega^{\rm pt} + B$$

The QCD sum-rule result predicts [Shuryak1978\_PLB79-135]:

$$B = -\frac{1}{4} \langle T^{\mu}_{\mu} \rangle = \frac{11 - \frac{2}{3}N_f}{32} \frac{\alpha_s}{\pi} \langle G^{\mu\nu}_a G_{a\mu\nu} \rangle - \frac{1}{4} \sum_f m_f \langle \bar{q}_f q_f \rangle$$
  

$$\approx 455 \text{ MeV/fm}^3$$
  
Equating the pressures of QGP ( $-B + 37 \frac{\pi^2}{90} T^4$ ) and pion gas ( $\frac{\pi^2}{30} T^4$ )  
at  $\mu = 0$  and  $T = T_c$  (~170 MeV), one have:  
 $B \simeq 400 \text{ MeV/fm}^3$   
On the other hand, fits to light hadron spectra suggest:  
 $B \simeq 50 \text{ MeV/fm}^3$ 

Gaussian parametrization for the bag constant **B**:

$$B = B_{\rm QCD} + (B_0 - B_{\rm QCD}) \exp\left[-\left(\frac{\sum_i \mu_i - m_{\Lambda}}{2\Lambda_{\overline{\rm MS}}}\right)^2\right],$$
  
where  $B_0 = 50 \text{ MeV fm}^{-3}, B_{\rm QCD} = 400 \text{ MeV fm}^{-3}.$ 

Bag constant [Maieron et al.2004\_PRD70-043010]



### Hadron-quark mixed phase inside compact stars

First-order phase transition [Masuda\_Hatsuda\_Takatsuka2016\_EPJA52-65]



Baryon density (p)

- σ = 0: point-like hadronic matter (HM) and strange quark matter (SQM), i.e., Gibbs construction.
- Moderate σ : geometrical structures [Chiba, Endo, Heiselberg, Maruyama, Tatsumi, Voskresensky, Yasuhira, Yasutake, ...]
   From Guo-yun Shao's talk:
- Large σ: the quark-hadron interface becomes planar, i.e., Maxwell construction.

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### Construct the mixed phase

The total volume occupied by the mixed phase is V, where the volume occupied by the quark phase is  $V_q$ . Then we can define the quark fraction:  $\chi = V_q/V$ .

The thermodynamic quantities for the mixed phase:

Baryon number density  $n_t = (1 - \chi)n_B + \chi n_q$ , Energy density  $E_t = (1 - \chi)(E_B + E_l^B) + \chi(E_q + E_l^q)$ .

**By** minimizing the total energy density  $E_t$ , we obtain the phase equilibrium conditions:

$$\mu_n = \mu_u + 2\mu_d, \quad P_B = P_q.$$
  
**Maxwell**
  
**Local** charge neutrality
$$Q_B = n_l^B \text{ and } Q_q = n_l^q.$$
  
**Gibbs**
  
**Global** charge neutrality
$$n_l = (1 - \chi)Q_B + \chi Q_q.$$

$$\mu_e^B = \mu_e^q$$

Note: the  $\beta$ -equilibrium condition is fulfilled.

### The phase transition in compact stars





### Mass, radius, and central density



### Summary

The compact stars with possible existence of hyperonic matter and quark matter were investigated. It was found that:

- The first-order deconfinement phase transition does not always reduce the maximum mass of compact stars;
- 2 Massive neutron stars may consist of multilayers baryonic matter, quark matter and their mixed phases;
- ③ The existence of quark matter in compact stars reduces their radii, which are more consistent with the recent observations [R = 10-11.5 km at  $M \approx 1.17-2 M_{\odot}$ , Ozel\_Freire2016\_AnnuRevAstro54-401];

Thank you!!!