

# The strangeness quark mean field theory

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J.H, A. Li, H. Shen, and H. Toki, Prog. Theor. Exp. Phys. 2014 (2014) 013D02
J.H, A. Li, H. Toki, and W. Zuo, Phys. Rev. C 89 (2014) 025802
X. Xing, J.H., and H. Shen, Phys. Rev. C 94 (2016) 044308
X. Xing, J.H., and H. Shen, in preparation

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# Outline

- Introduction
- Strangeness quark mean field theory
- Numerical results and discussions
- Summary and Perspectives

## Strangeness nuclear physics

#### Hadrons



#### Baryon-baryon force



#### A. Gal, E. V. Hungerford, and D. J. Millener, Rev. Mod. Phys. 88(2016)035004

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#### Compact star



Hypernuclei



### Theoretical methods

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#### ✓ ab initio methods

- H. Nemura, Y. Akaishi, and Y. Suzuki, Phys. Rev. Lett. 89(2002)142504
- E. Hiyama and T. Yamada, Prog. Part. Nucl. Phys. 63(2009)339
- D. Lonardoni, S. Gandolfi, and F. Pederiva, Phys. Rev. C 87(2013)041303(R)
- R. Wirth, et al. Phys. Rev. Lett. 113(2014)192502

### √Shell model

D. J. Millener, Nucl. Phys. A 881(2012)298

#### ✓ Skyrme Hartree-Fock model

M. Rayet, Ann. Phys. (NY) 102(1976)226
X. R. Zhou, et al. Phys. Rev. C 76(2007)034312
H.-J. Schulze and T. Rijken, Phys. Rev. C 88(2013)024322

✓ Relativistic mean-field model

R. Brockmann and W. Weise, Phys. Lett. B 69(1977)167
H. Shen, F. Yang, and H. Toki, Prog. Theor. Phys. 115(2006)325
R. L. Xu, C. Wu, and Z. Z. Ren, J. Phys. G 39(2012)085107

T. T. Sun, et al., Phys. Rev. C 94(2016)064319

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### Relativistic many-body theories from quark level

- ✓ baryons are not point particles!
- ✓ baryon properties change in medium!
- √ quark-gluon plasma!



### Strangeness system with quark model (

✓ Quark meson coupling (QMC) model
 K. Tsushima, et al. Nucl. Phys. A 630(1998)691
 ✓ Friedberg-Lee model

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J. S. Liang and H. Shen, Phys. Rev. C88 (2013) 035208

✓ Quark mean field (QMF) model
H. Shen, H. Toki, Nucl. Phys. A 707 (2002) 469
J.H, A. Li, H. Shen, and H. Toki, Prog. Theor. Exp. Phys. 2014(2014) 013D02
J.H, A. Li, H. Toki, and W. Zuo, Phys. Rev. C 89(2014) 025802
Properties of hypernuclei and neutron star with QMF model

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X. Xing, J.H., and H. Shen, Phys. Rev. C 94 (2016) 044308

Constituent quark in Dirac equation

 $[-i\alpha \cdot \nabla + \beta m_i^* + \beta U(r)]q_i(r) = \varepsilon_i^* q_i(r)$ 



where, the effective quark mass is  $m_i^* = m_i + g_\sigma^i \sigma$ 

and effective single particle energy is

$$\varepsilon_i^* = \varepsilon_i - g_\omega^i - g_\rho^i \rho \tau_3$$

Confinement potential

$$U(r) = \frac{1}{2}(1+\gamma^0)(ar^2 + V_0)$$

• Center-of-mass corrections  $\langle B|\sum_{i=1}^{3}\gamma^{0}(i)\{\frac{1}{3}\gamma(i)\cdot\sum_{j=1}^{3}\vec{p_{j}}+\frac{1}{2}(1+\gamma^{0}(i))[U(r_{i})-U(\rho_{i})]\}|B\rangle$ 

### Quark level



Pionic self-energy correction

$$\delta M_B^{\pi} = -\sum_k \sum_{B'} \frac{V_j^{\dagger BB'} V_j^{BB'}}{w_k}$$

• Gluon correction

**Color-electric** 
$$(\Delta E_B)_g^E = \frac{1}{8\pi} \sum_{i,j} \sum_{a=1}^8 \int \frac{d^3 r_i d^3 r_j}{|\vec{r_i} - \vec{r_j}|} \langle B|J_i^{0a}(\vec{r_i})J_j^{0a}(\vec{r_j})|B\rangle$$

**Color-magnetic** 
$$(\Delta E_B)_g^M = -\frac{1}{8\pi} \sum_{i,j} \sum_{a=1}^8 \int \frac{d^3 r_i d^3 r_j}{|\vec{r_i} - \vec{r_j}|} \langle B|\vec{J}_i^a(\vec{r_i}) \cdot \vec{J}_j^a(\vec{r_j})|B \rangle$$

Quark color current density  $J_i^{\mu a}(x) = g_c \bar{\psi}_q(x) \gamma^{\mu} \lambda_i^a \psi_q(x)$ 

• Baryon mass

$$M_B^* = E_B^{*0} - \epsilon_{\text{c.m.}} + \delta M_B^{\pi} + (\Delta E_B)_g^E + (\Delta E_B)_g^M$$

### **Baryon level**



#### • Strangeness QMF Lagrangian

$$\begin{split} \mathcal{L}_{\text{QMF}} &= \bar{\psi} \left[ i\gamma_{\mu}\partial^{\mu} - M_{N}^{*} - g_{\omega}\omega\gamma^{0} - g_{\rho}\rho\tau_{3}\gamma^{0} - e\frac{(1-\tau_{3})}{2}A\gamma^{0} \right]\psi \\ &+ \bar{\psi}_{H} \left[ i\gamma_{\mu}\partial^{\mu} - M_{H}^{*} - g_{\omega}^{H}\omega\gamma^{0} + \frac{f_{\omega}^{H}}{2M_{H}}\sigma^{0i}\partial_{i}\omega \right]\psi_{H} \\ &- \frac{1}{2}(\nabla\sigma)^{2} - \frac{1}{2}m_{\sigma}^{2}\sigma^{2} - \frac{1}{3}g_{2}\sigma^{3} - \frac{1}{4}g_{3}\sigma^{4} \\ &+ \frac{1}{2}(\nabla\omega)^{2} + \frac{1}{2}m_{\omega}^{2}\omega^{2} + \frac{1}{4}c_{3}\omega^{4} \\ &+ \frac{1}{2}(\nabla\rho)^{2} + \frac{1}{2}m_{\rho}^{2}\rho^{2} + \frac{1}{2}(\nabla A)^{2}, \end{split} \qquad \bullet \quad \text{Dirac equations for baryons} \\ &+ \frac{1}{2}(\nabla\rho)^{2} + \frac{1}{2}m_{\rho}^{2}\rho^{2} + \frac{1}{2}(\nabla A)^{2}, \qquad \left[ i\gamma_{\mu}\partial^{\mu} - M_{N}^{*} - g_{\omega}\omega\gamma^{0} - g_{\rho}\rho\tau_{3}\gamma^{0} - e\frac{(1-\tau_{3})}{2}A\gamma^{0} \right]\psi = 0, \\ &\left[ i\gamma_{\mu}\partial^{\mu} - M_{H}^{*} - g_{\omega}^{H}\omega\gamma^{0} + \frac{f_{\omega}^{H}}{2M_{H}}\sigma^{0i}\partial_{i}\omega \right]\psi_{H} = 0. \end{split}$$

• Equations of motion for mesons

$$\begin{split} \Delta \sigma &- m_{\sigma}^{2} \sigma - g_{2} \sigma^{2} - g_{3} \sigma^{3} = \frac{\partial M_{N}^{*}}{\partial \sigma} \langle \bar{\psi} \psi \rangle + \frac{\partial M_{H}^{*}}{\partial \sigma} \langle \bar{\psi}_{H} \psi_{H} \rangle, \\ \Delta \omega &- m_{\omega}^{2} \omega - c_{3} \omega^{3} = -g_{\omega} \langle \bar{\psi} \gamma^{0} \psi \rangle - g_{\omega}^{H} \langle \bar{\psi}_{H} \gamma^{0} \psi_{H} \rangle + \frac{f_{\omega}^{H}}{2M_{H}} \partial_{i} \langle \bar{\psi}_{H} \sigma^{0i} \psi_{H} \rangle, \\ \Delta \rho &- m_{\rho}^{2} \rho = -g_{\rho} \langle \bar{\psi} \tau_{3} \gamma^{0} \psi \rangle, \\ \Delta A &= -e \langle \bar{\psi} \frac{(1 - \tau_{3})}{2} \gamma^{0} \psi \rangle. \end{split}$$

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### SQMF parameters



### The strength of quark confinement potential

=									=
_	m	u (MeV)	$V_u$ (MeV	) $a_u$ (fr	$m^{-3}) m_s$	(MeV)	$V_s \ ({\rm MeV})$	$a_s \ (\mathrm{fm}^{-3})$	)
	set A	250	-24.28660	0.579	9450	330	101.78180	0.097317	7
	set B	300	-62.25718	87 0.534	1296	380	54.548210	0.087243	}
set C		350	-102.0415	75 0.495	5596	430	6.802695	0.079534	Ł
	he co	upling	constar	nts be	tweer	n mes	on and	baryo	= ns
Model	m	$g_{\sigma}^{u} = g_{\sigma}^{u}$	$g_\omega$	$g^{\Lambda}_{\omega}$	$g_{\omega}^{\Xi}$	$g_{ ho}$	$g_2$	$g_3$	C <sub>3</sub>
	(Me)	eV)		1 marine			$(\mathrm{fm}^{-1})$		
QMF-N	NK1S 25	50 5.158'	71 11.54726	$0.8258g_{\omega}$	$0.4965g_{\omega}$	3.79601	-3.52737	-78.52006	305.00240
QMF-N	MK2S 30	0 5.0934	46 12.30084	$0.8134g_{\omega}$	$0.4800g_{\omega}$	4.04190	-3.42813	-57.68387	249.05654
QMF-N	VK3S 35	50 5.016	31 12.83898	$0.8040g_{\omega}$	$0.4681g_{\omega}$	4.10772	-3.29969	-39.87981	221.68240
							$\Lambda = -$	$-30 \mathrm{MeV}$	
						$U_2$	$\Xi = -$	$12 { m MeV}$	
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## The properties of nuclei



#### Binding energy and charge radii



## The properties of nuclear matter

#### Nuclear saturation properties

Model	$ ho_0$	E/A	$K_0$	J	$M_N^*/M_N$	$L^0$	$K_{\mathrm{sym}}^0$	$K_{\rm asy}$	$Q_0$	$K_{\tau}$
	$(\mathrm{fm}^{-3})$	(MeV)	(MeV)	(MeV)		(MeV)	(MeV)	(MeV)	(MeV)	(MeV)
QMF-NK1	0.154	-16.3	323	30.6	0.70	84.8	-28.8	-537.6	495.4	-667.7
QMF-NK2	0.152	-16.3	328	32.9	0.66	93.7	-23.5	-585.7	221.0	-648.8
QMF-NK3	0.150	-16.3	322	33.6	0.64	97.3	-12.0	-595.8	263.0	-675.3

#### symmetric nuclear matter

#### pure neutron matter

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#### The $\Lambda$ energy levels of hypernuclei

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## The properties of hypernuclei

The scalar and vector potentials of  $\Lambda$  and  $\Xi^0$  hypernuclei

QMF-NK3

QMF-NK2S

QMF-NK1S

 $\mathsf{U}_{v}^{\Xi^{0}}$ 

 $U_{s}^{\Xi^{\circ}}$ 

 $U_v^{\Xi^0}$ 

 $\mathsf{U}^{\Xi^0}_{\mathsf{S}}$ 

 $\mathsf{U}_{\mathsf{v}}^{\Xi^0}$ 

 $U_s^{\Xi^\circ}$ 

8

6

<sup>208</sup><sub>Ξ<sup>0</sup></sub>Pb

10

89 ≘⁰`

<sup>40</sup><sub>⊐⁰</sub>Ca



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## The properties of hypernuclei



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Exp. Data: A. Gal, E. V. Hungerford and D. J. Millener, Rev. Mod. Phys. 88(2016)035004

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### The equations of state of neutron star matter



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### The properties of neutron star



### The particle fractions



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### The properties of neutron star



#### The mass of neutron star



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### Summary and Perspective



- The quark mean field model is extended to strangeness nuclear physics within the pionic and gluonic corrections.
- The binding energies of  $\Lambda$  hypernuclei can be reproduced very well in present framework. The ones of  $\Xi^0$  hypernuclei are also predicted.
- The massive neutron stars are obtained with  $\Lambda$  and hyperon, whose masses are around 2.1M $_{\odot.}$
- The double  $\Lambda$  and  $\Xi^-$  hypernuclei will be studied.