

# 連星中性子星合体後の系からの質量放出

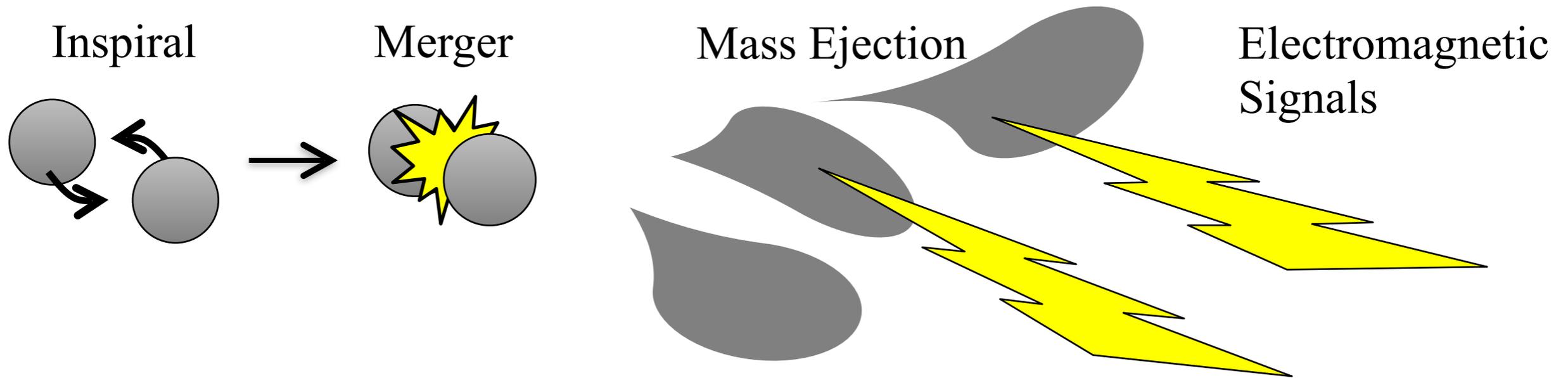
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in collaboration with Kenta Kiuchi, Masaru Shibata, Nobuya Nishimura (YITP)  
Yuichiro Sekiguchi (Toho U.)

SF et al. 17 ([arXiv:1711.02093](https://arxiv.org/abs/1711.02093))

# Introduction

# Kilonova/Macronova

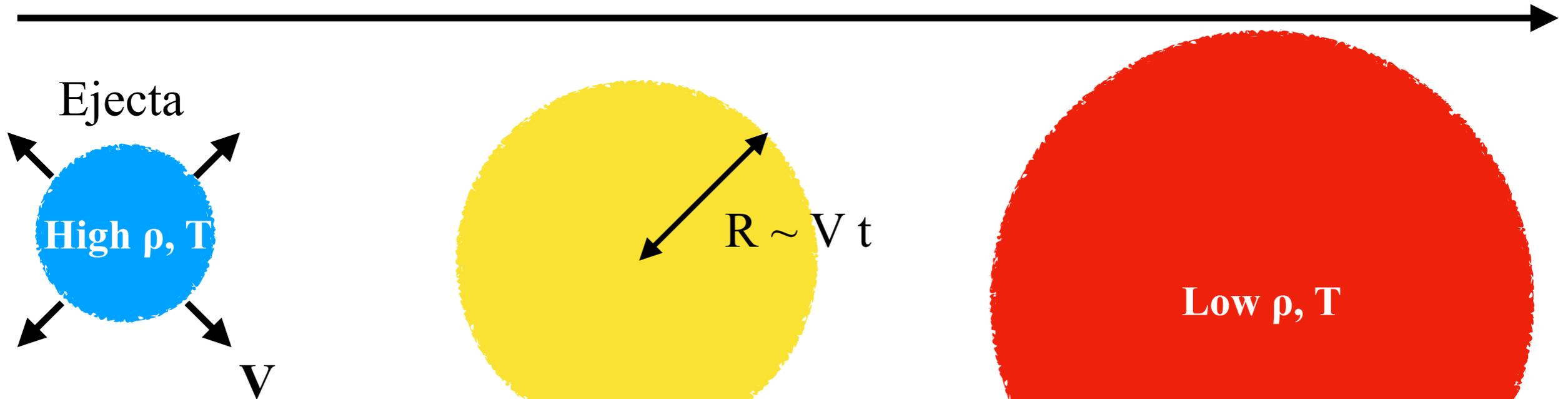


- Two NSs approach due to GW emission → Merger
- A part of the NS matter is ejected by several processes
- Unstable nuclei are synthesized via *r*-process
- Radioactive decay (and fission) heats up the ejecta  
→ Thermal emission (Kilonova/Macronova)

# Simple Model for Kilonova

Li & Paczynski 98, Metzger et al. 10, ...

Time



optical depth  $\tau \sim \rho \kappa R \gg 1$

$t_{\text{diff}} \gg R/v$

Photons leak-out very slowly  
~ Adiabatic expansion



optical depth becomes smaller  
becomes more transparent



peak of the photon emission :

$t_{\text{diff}} \sim R/v$

# Kilonova/Macronova Model

- Mass, velocity, and **opacity** of the ejecta are key quantities Li & Paczynski 98, Metzger 08

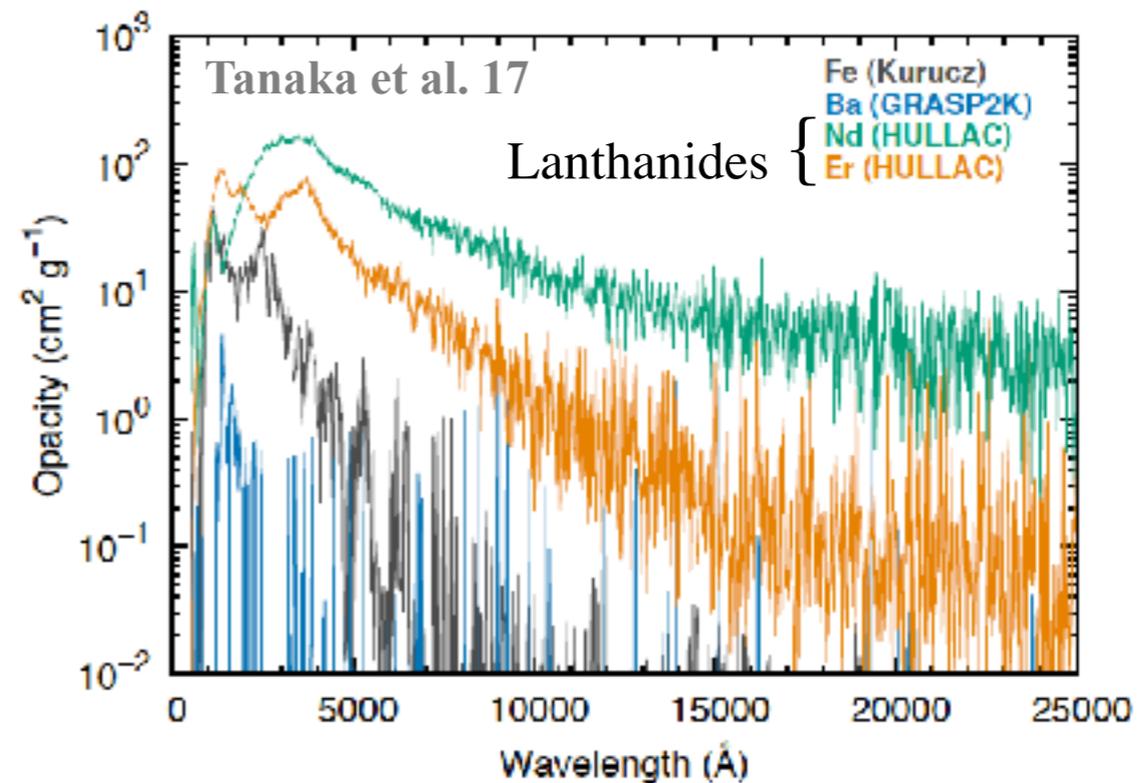
$$t_{\text{peak}} \sim 10 \text{ days} \left( \frac{v}{0.3c} \right)^{-1/2} \left( \frac{M}{0.01 M_{\text{solar}}} \right)^{1/2} \left( \frac{\kappa}{10 \text{ cm}^2 / \text{g}} \right)^{1/2}$$

$$L_{\text{peak}} \sim 10^{41} \text{ erg/s} \left( \frac{f}{10^{-6}} \right) \left( \frac{v}{0.3c} \right)^{1/2} \left( \frac{M}{0.01 M_{\text{solar}}} \right)^{1/2} \left( \frac{\kappa}{10 \text{ cm}^2 / \text{g}} \right)^{-1/2}$$

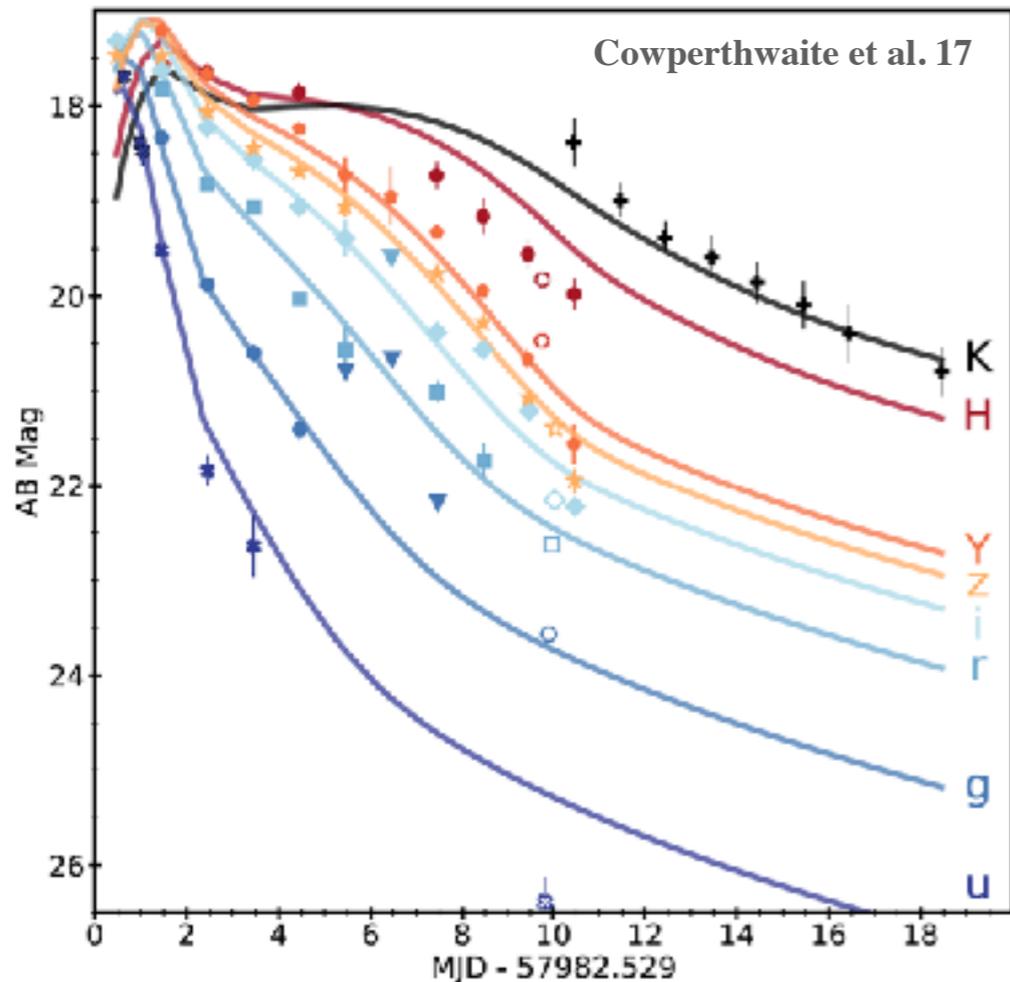
$$T_{\text{peak}}^{\text{eff}} \sim 2 \times 10^3 \text{ K} \left( \frac{f}{10^{-6}} \right)^{1/4} \left( \frac{v}{0.3c} \right)^{-1/8} \left( \frac{M}{0.01 M_{\text{solar}}} \right)^{-1/8} \left( \frac{\kappa}{10 \text{ cm}^2 / \text{g}} \right)^{-3/8}$$

- Lanthanides have significantly high opacity Kasen et al. 13, Tanaka & Hotokezaka 13  
( $Z=57-71$ )

- In Lanthanide-rich ejecta,  
 $\kappa \sim 10 \text{ cm}^2/\text{g}$



# Electromagnetic Signal Associated with GW170817



## ○ Optical - NIR light curve

- Consistent with Kilonova/Macronova model.

- Many work suggest

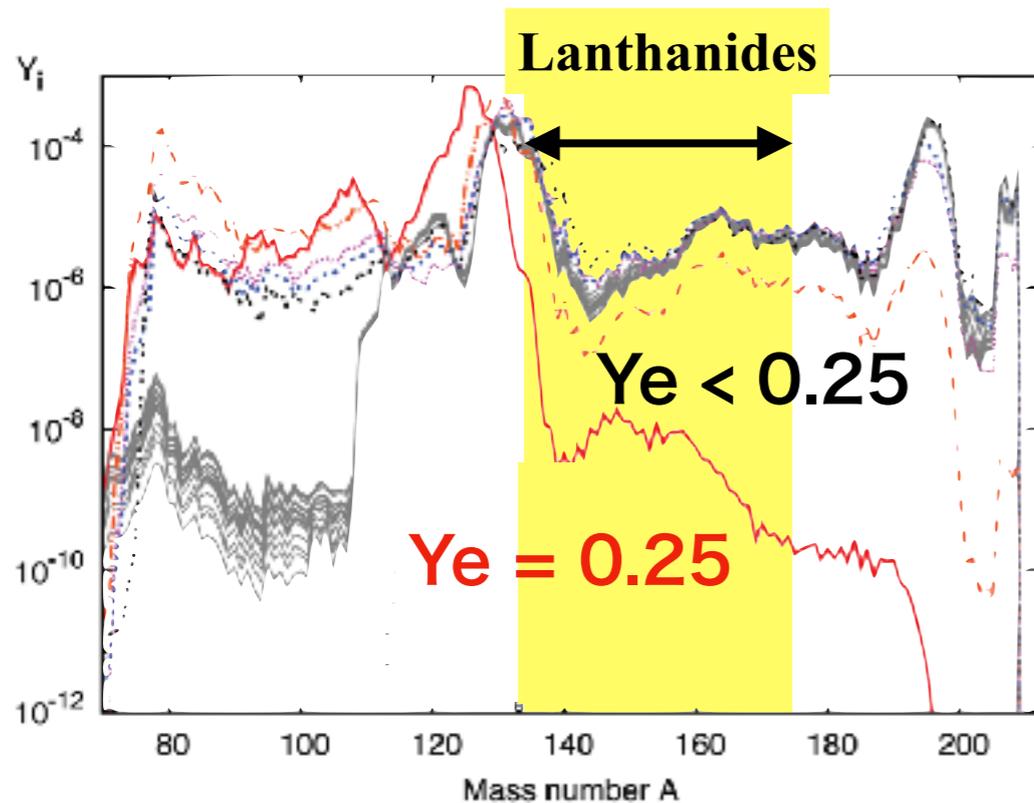
“red” component of  $M_{ej} \sim 0.02-0.04 M_{sun}$  ( $v \sim 0.1 c$ )  
(Lanthanide-rich,  $\kappa \sim 10 \text{ cm}^2/\text{g}$ )

“blue” component of  $M_{ej} \sim 0.01 M_{sun}$  ( $v \sim 0.2-0.3 c$ )  
(Lanthanide-poor,  $\kappa \sim 0.1-1 \text{ cm}^2/\text{g}$ )

Villar et al. 17, Kasen et al. 17, Nicholl et al. 17, Chornock et al. 17, ...

**Can we derive such values from latest simulation for NS-NS merger ?**

# Lanthanide Fraction in the Ejecta



Korobkin et al. 12 (For dynamical ejecta)

○ **Key quantity for Lanthanide production**

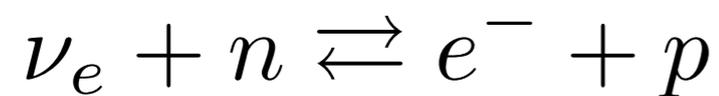
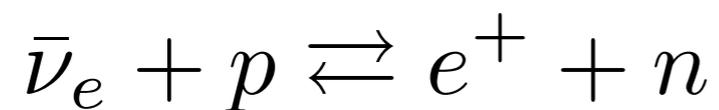
$$\text{Electron fraction } Y_e = \frac{n_e}{n_B} = \frac{n_p}{n_n + n_p}$$

$$Y_e < 0.5 \leftrightarrow \text{proton \#} < \text{neutron \#}$$

$Y_e > 0.25$

→ Lanthanides are not synthesized efficiently.

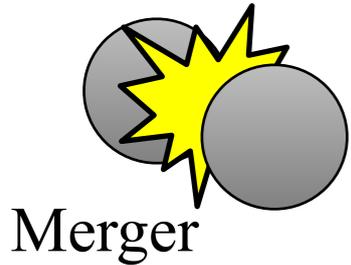
◎ **Relevant reactions :**

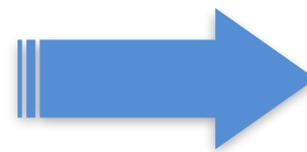


Initially low  $Y_e$  can become large due to

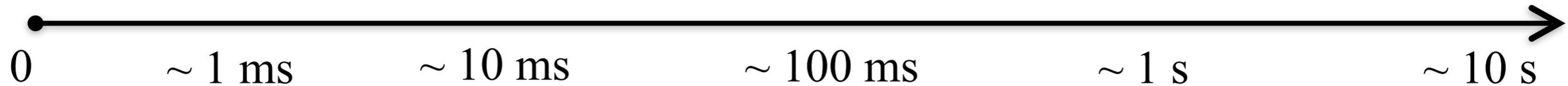
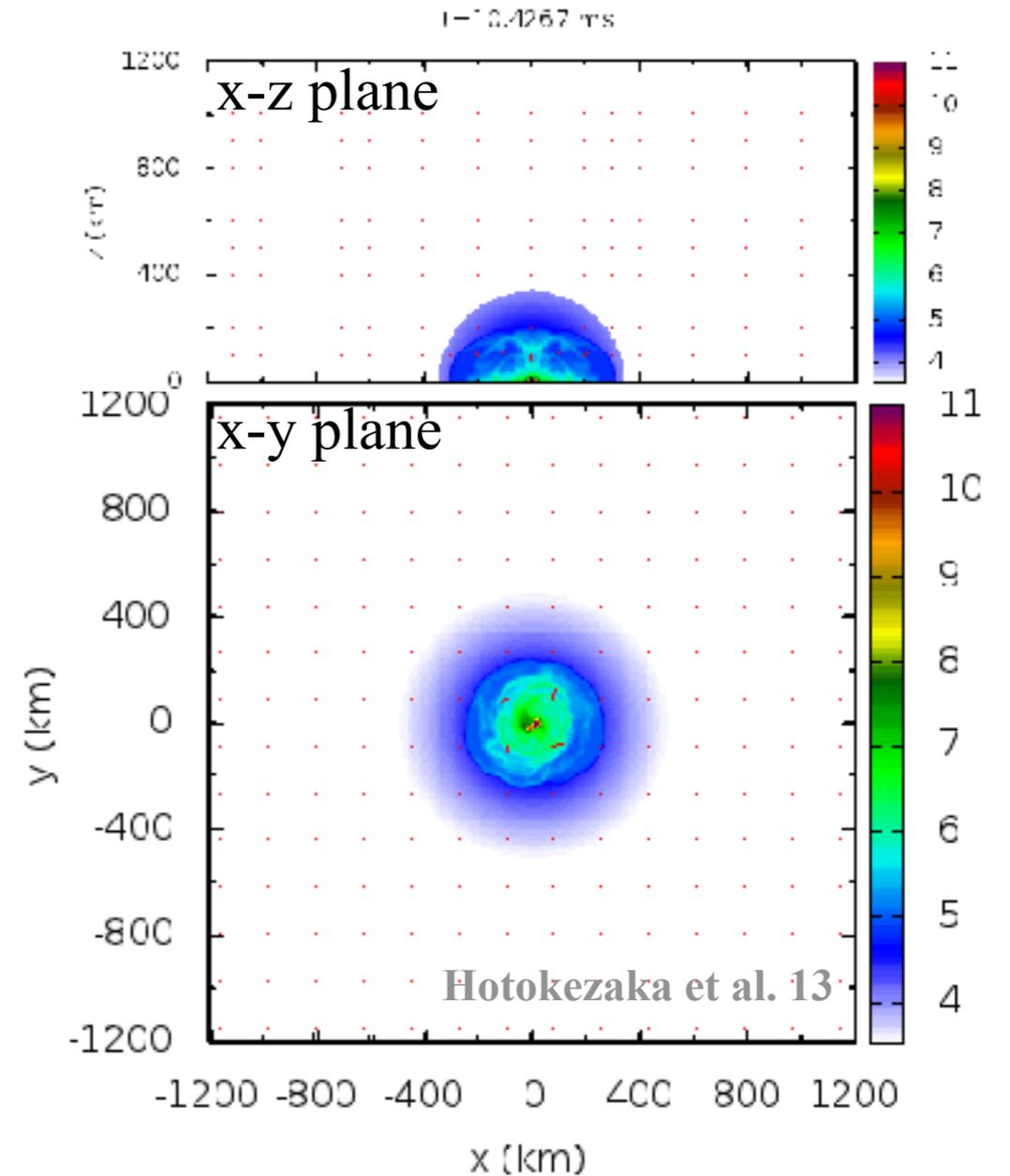
- 1) Neutrino irradiation
- 2) High temperature (positron capture sets in)

# Dynamical Mass Ejection



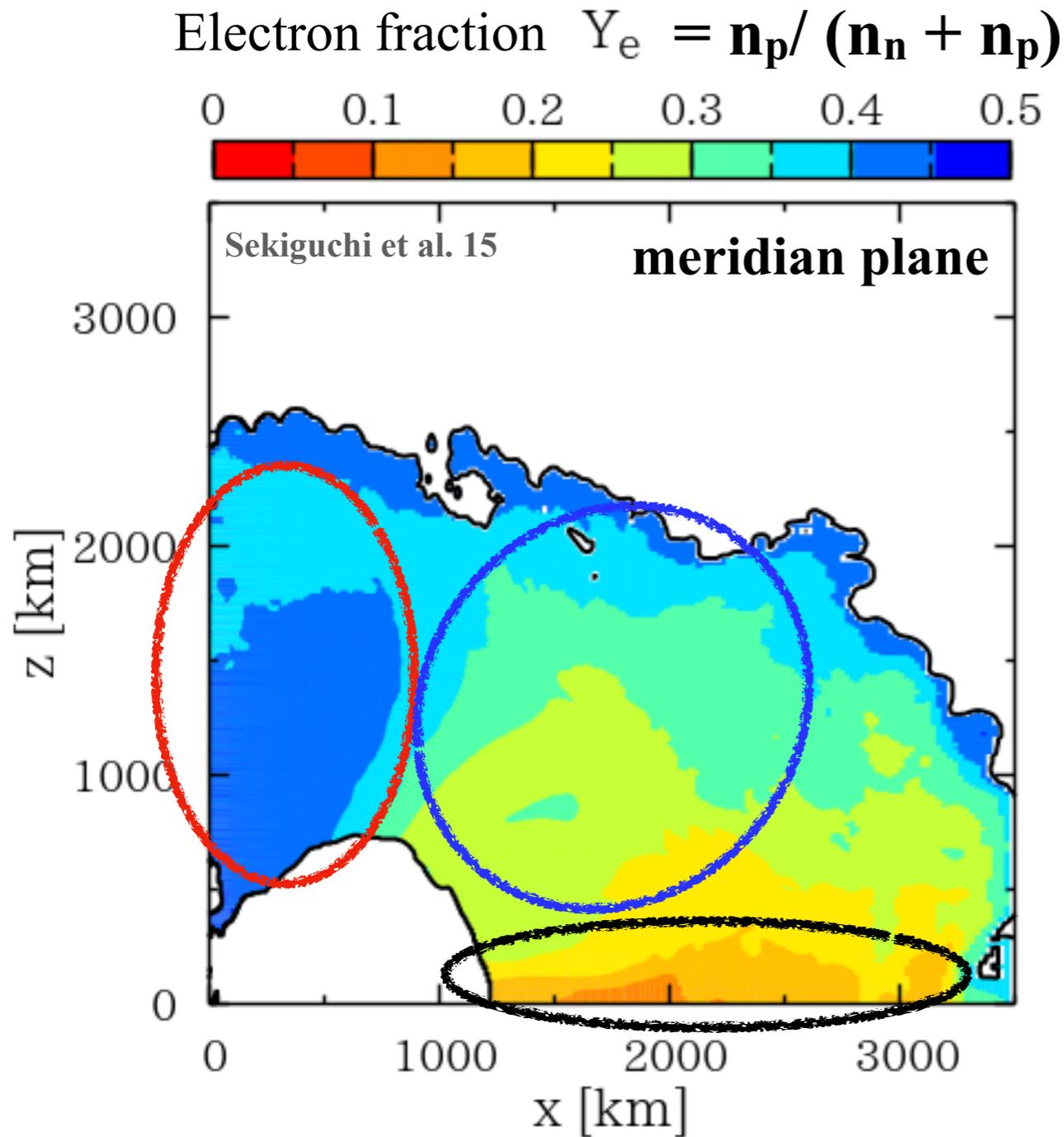
 **Dynamical ejection**

- Due to tidal force and shock heating
- Well studied (mass, EOS dependence)



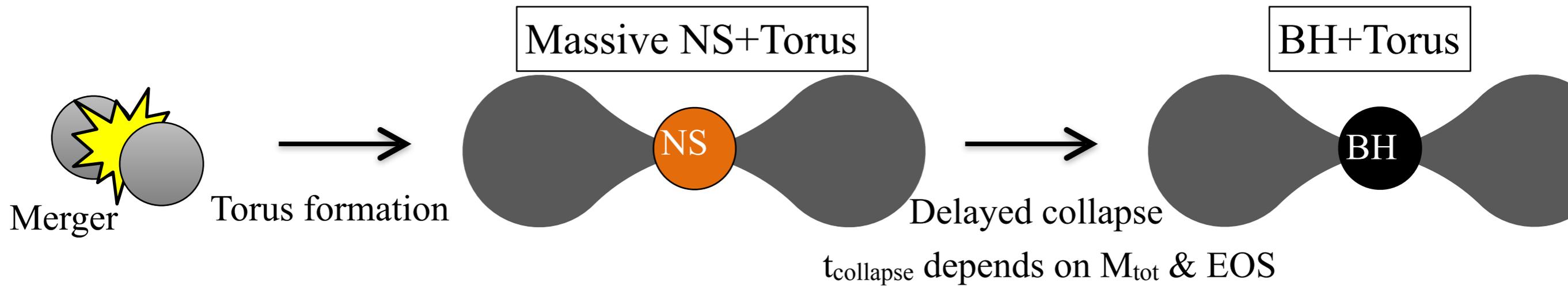
Time after merger

# Morphology of Dynamical Ejecta



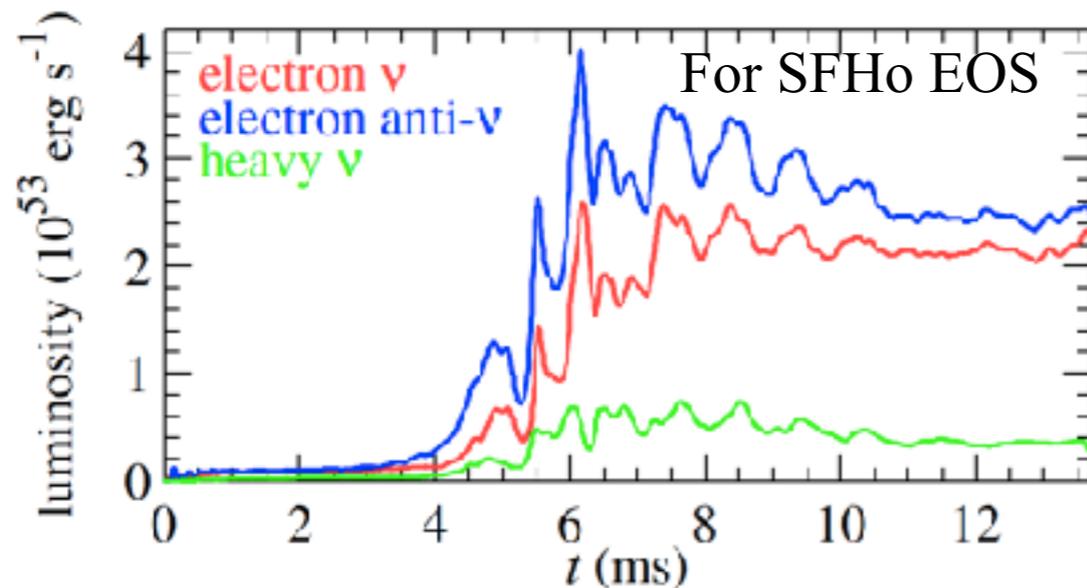
- $M_{ej} \sim 0.001-0.02 M_{sun}$ 
  - depends on EOS & mass ratio
- Nearly isotropic
- Tidally driven ejecta
  - $Y_e < 0.2$
- Thermal driven ejecta
  - $Y_e = 0.2-0.4$  (due to positron capture)
- Neutrino irradiated ejecta
  - $Y_e > 0.4$  (due to  $\nu$ -absorption)

# Evolution of the Typical Merger Remnant



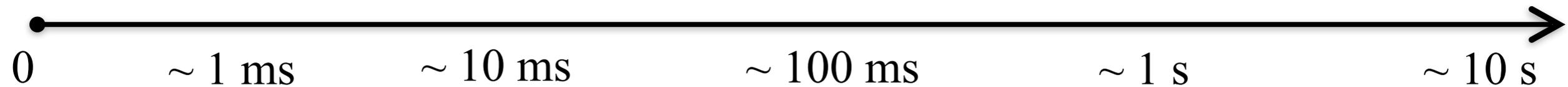
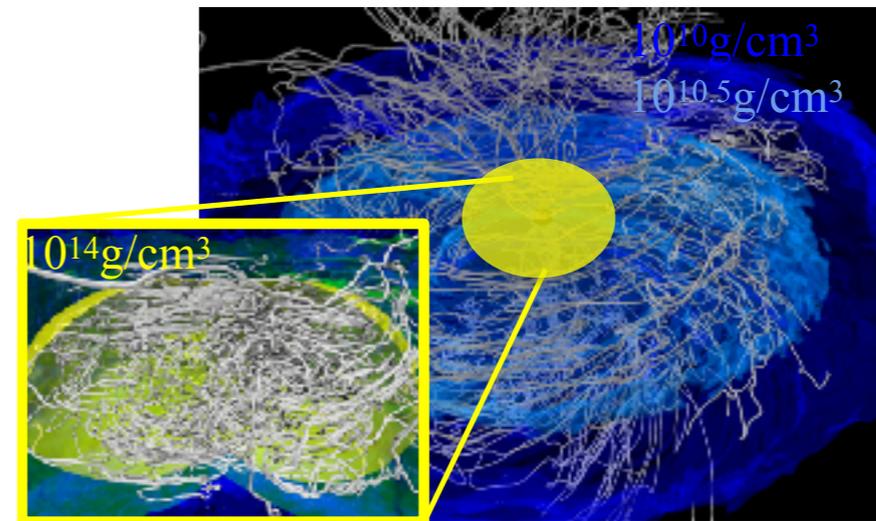
## Neutrino emission

Sekiguchi et al. 15



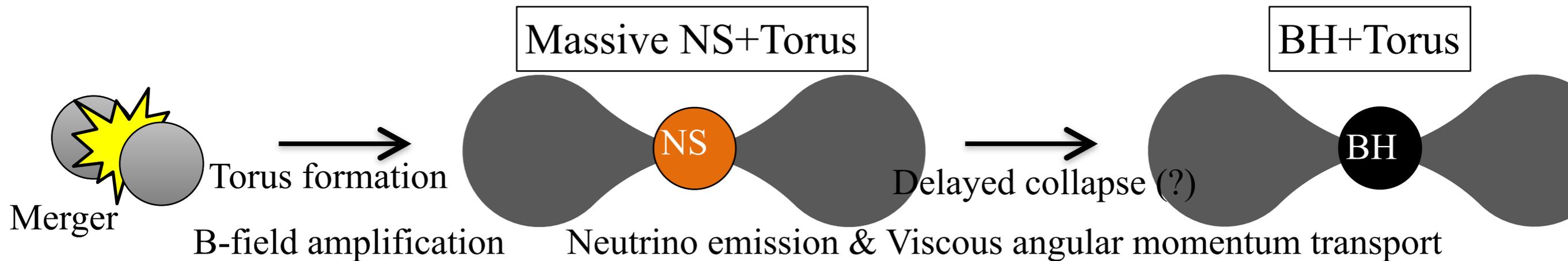
## B-field amplification

Kiuchi et al. 14, 15



Time after merger

# Post-merger Mass Ejection



## Viscosity & Neutrino-driven ejection

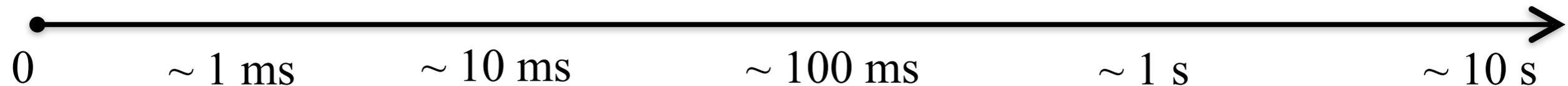


Heating by neutrinos and viscosity

The ejecta mass can exceed (significantly!) the dynamical ejecta

Currently extensively studied

Fernandez & Metzger 13, Metzger & Fernandez 14, Perego et al. 14  
Siegel et al. 17, Lippuner et al. 17, SF et al. 17



Time after merger

# Our Research

To connect the ejecta properties and the physics of the merger, it is needed to construct reliable physical model of **the binary NS merger remnant** and mass ejection from it

- Ejecta properties (mass,  $Y_e$ , velocity)
- Morphology

We perform long-term ( $>$  second) simulations for NS-NS merger remnant including relevant physics:

- General relativity
- Neutrinos
- Viscous angular momentum transport

# Previous Works

	Gravigy	Initial condition	NS	Neutrino	Hydrodynamics
Metzger & Fernandez 14 Lippuner et al. 17,	Pseudo-Newtonian	J=Const. equiv. torus	Reflection boundary	Leakage + light bulb	Viscous fluid
Perego et al. 14, ...	Newtonian	Based on Newtonian simulation	Not resolved	Leakage + absorption	Ideal fluid
<b>SF et al. 17 (Our work)</b>	<b>Full GR</b> 	<b>Based on Full GR simulation</b> 	<b>Resolved</b> 	<b>Leakage + M1-scheme</b> 	<b>Ideal fluid</b>

# Method

# Method

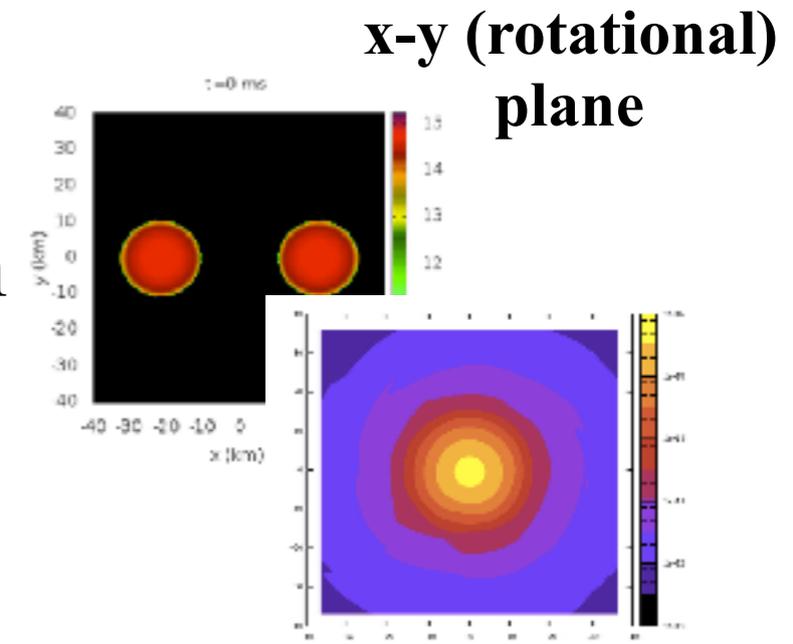
## ◎ Strategy

- i) Merger of NS–NS and massive NS formation  
by 3-D full GR simulation

Sekiguchi et al. 15

Equation of state : DD2

( → The remnant is long-lived massive NS)



# Method

## ◎ Strategy

- i) Merger of NS–NS and massive NS formation by 3-D full GR simulation

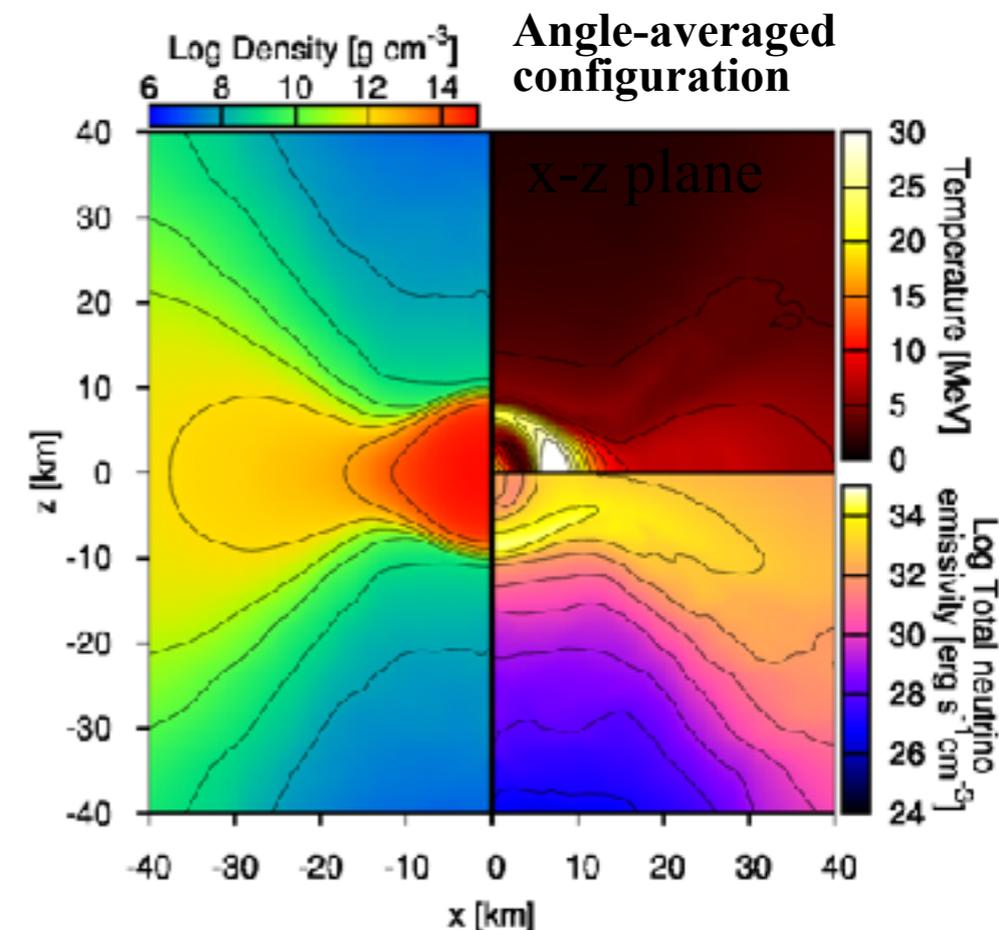
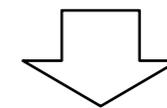
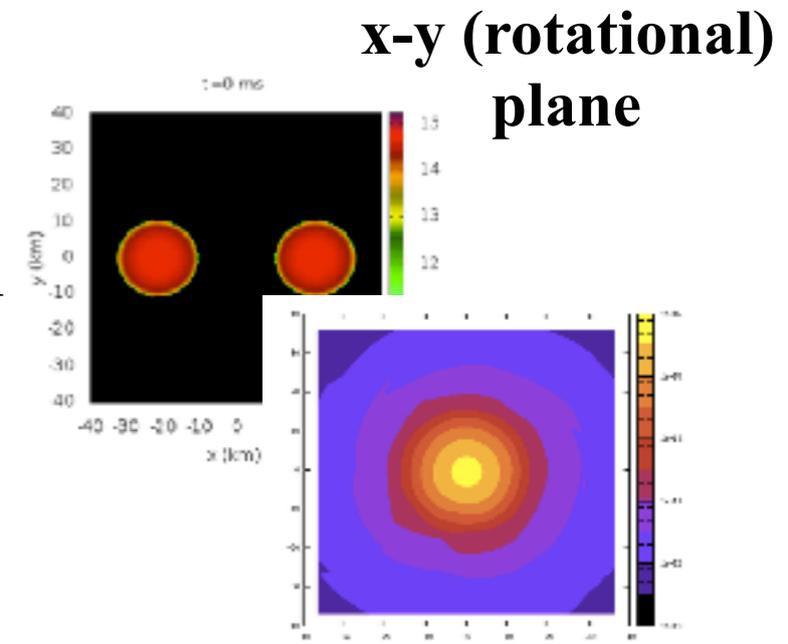
Sekiguchi et al. 15

Equation of state : DD2

( → The remnant is long-lived massive NS)

Average over azimuthal angles around the rotational axis after  $\sim 50$  ms after the merger, when the system settles into quasi-axisymmetric configuration.

- ii) Long-term Axisymmetric 2-D simulation using angle-averaged configuration as a initial condition



# Method

## ◎ Basic Equations

- Full GR neutrino radiation viscous hydrodynamics

- Einstein's equation
- Neutrino radiation transfer equation

Leakage+ scheme incorporating Moment formalism Thorne 81, Shibata et al. 11

- Viscous hydrodynamics equation

In order to mimic the effective viscosity due to MHD turbulence,

We solve viscous hydrodynamics equation.

Israel & Stuart 79,

The code was developed by Sekiguchi-san.

Shibata et al. 17, Shibata & Kiuchi 17

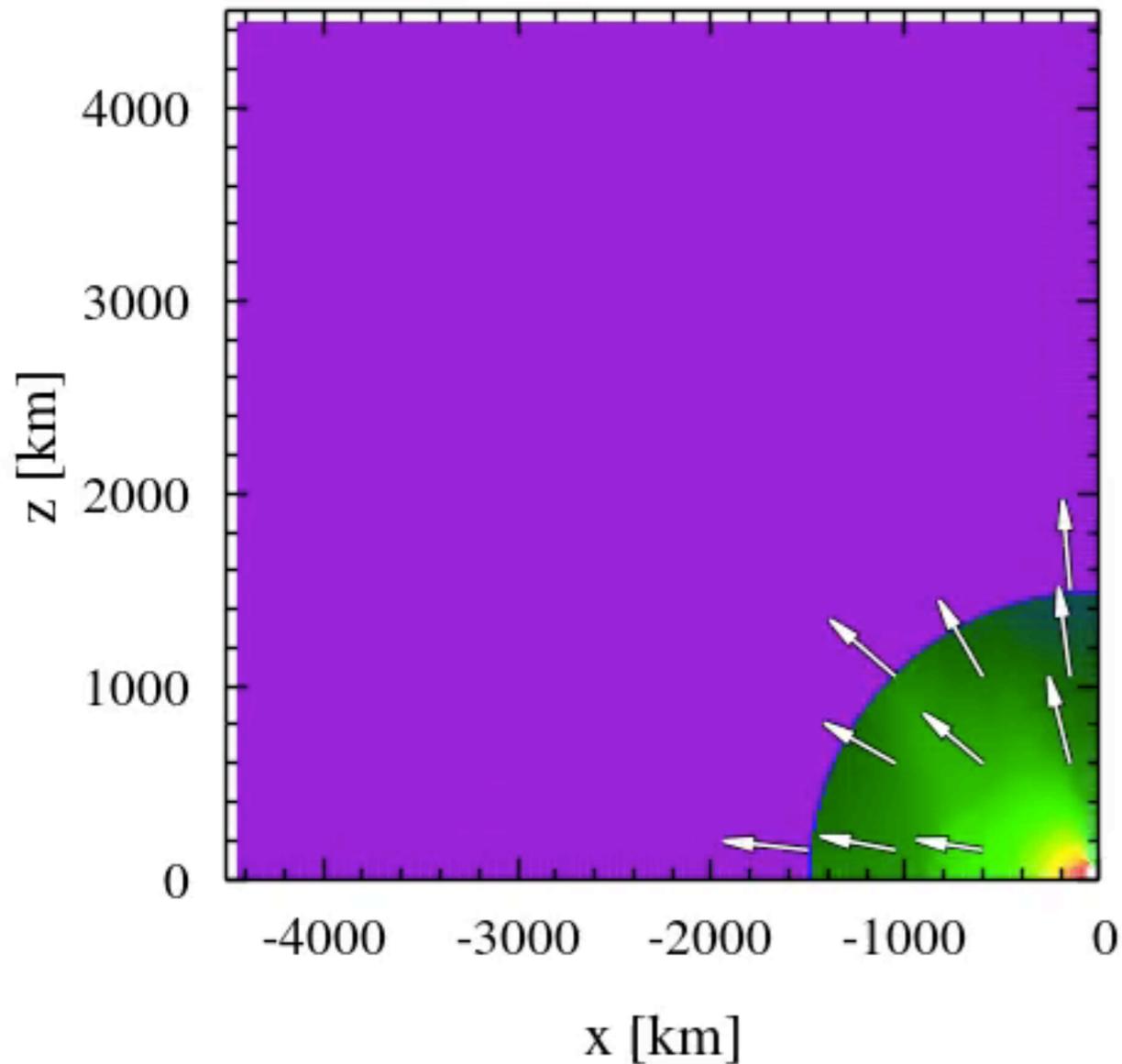
Shakura–Sunyaev “alpha” parametrization :  $\nu = \alpha \cdot c_s H_{\text{tur}}$

Models :  $\alpha = 0, 0.01, 0.02, \mathbf{0.04}$

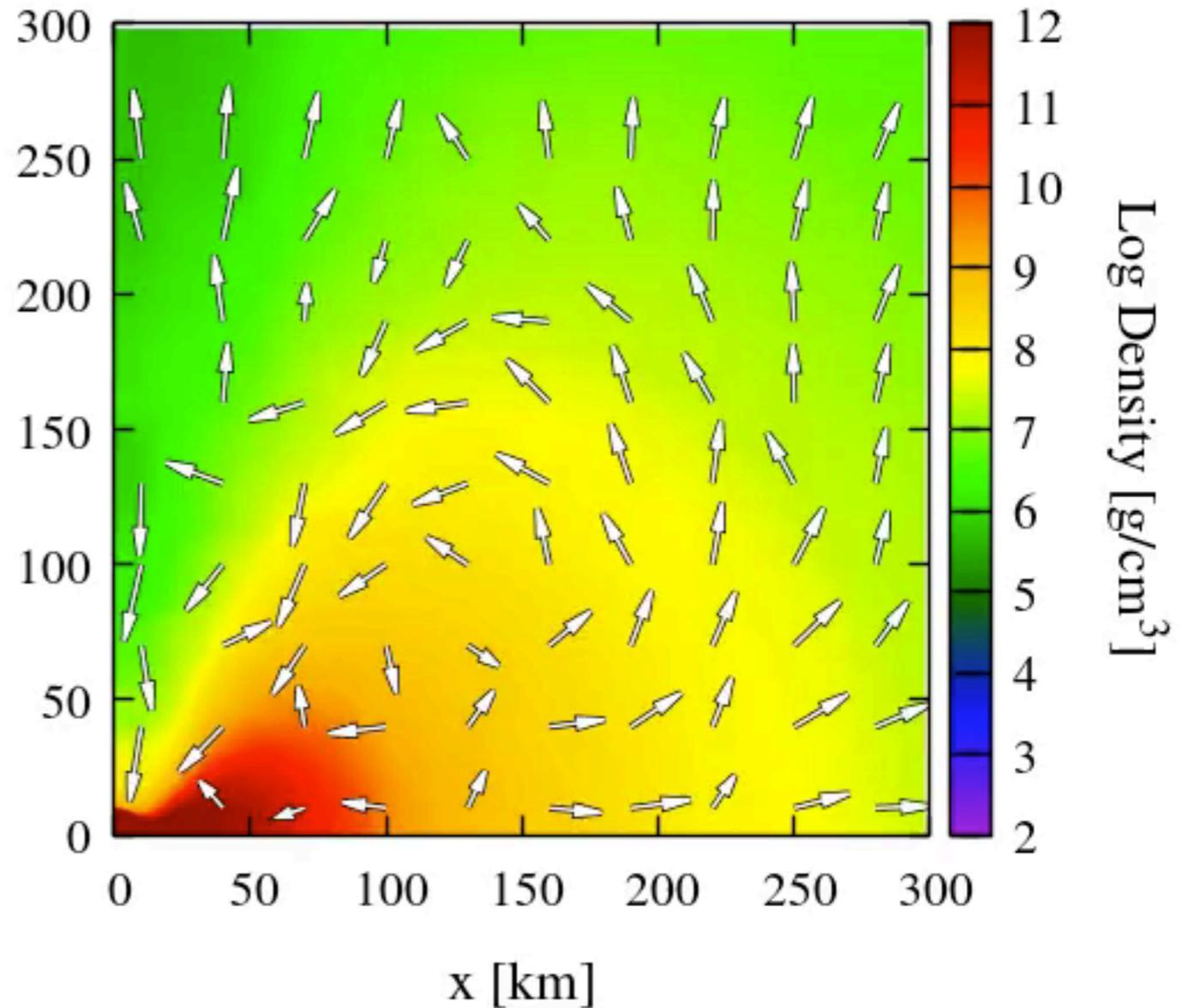
Result

# Dynamics for $\alpha = 0.04$ model

t = 0.00 ms



Density structure in meridional plane



**Torus :**

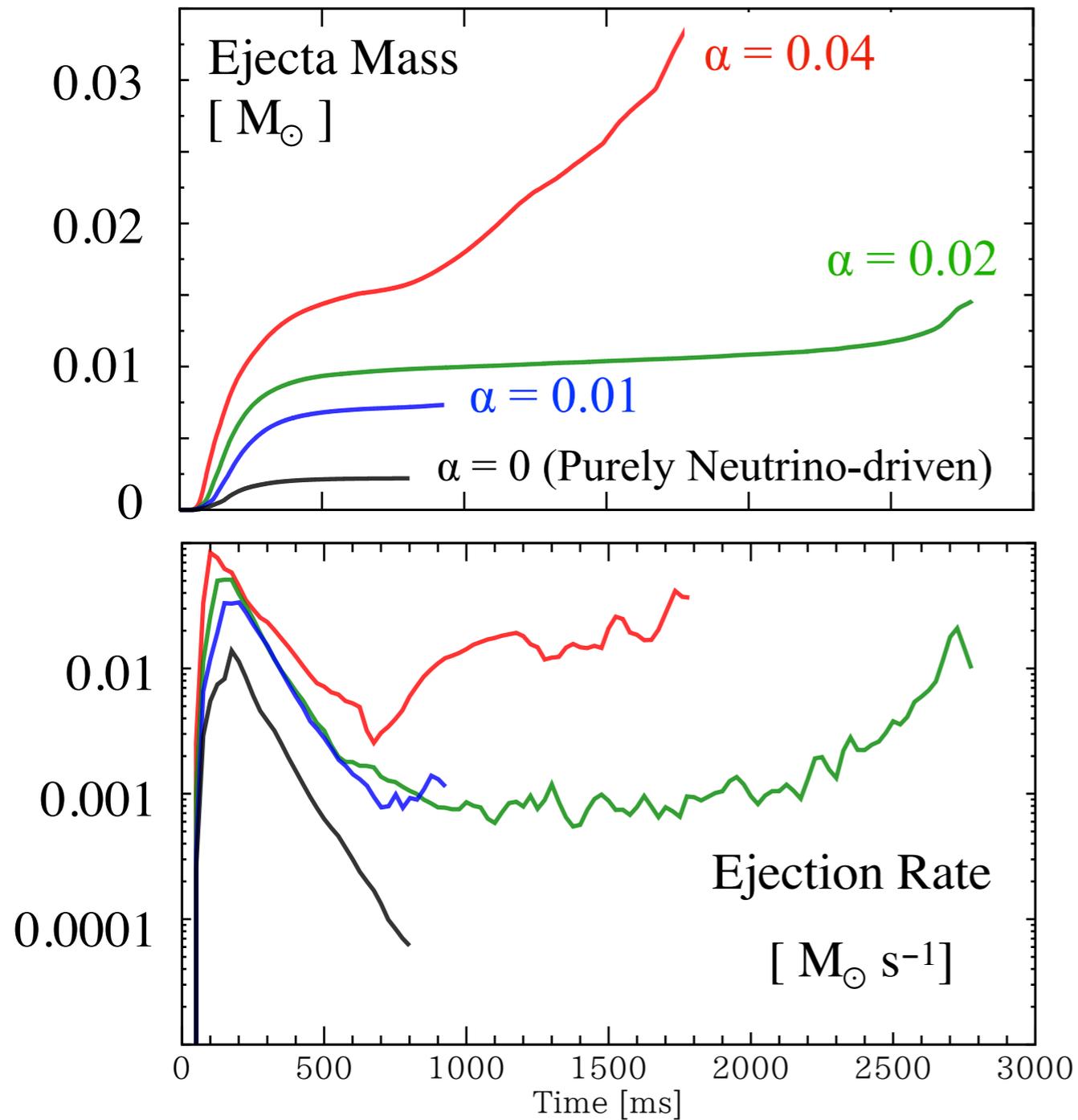
Gradually expands  
→ Ejected

**Ejecta :**

~ 10 ms : ejection by a shock  
After that : viscosity-driven ejection

# Mass Ejection from Merger Remnant

# Mass Ejection



○ Integral of unbound mass flux at  $r = 4000$  km

○  $\alpha = 0.04$

High mass ejection rate

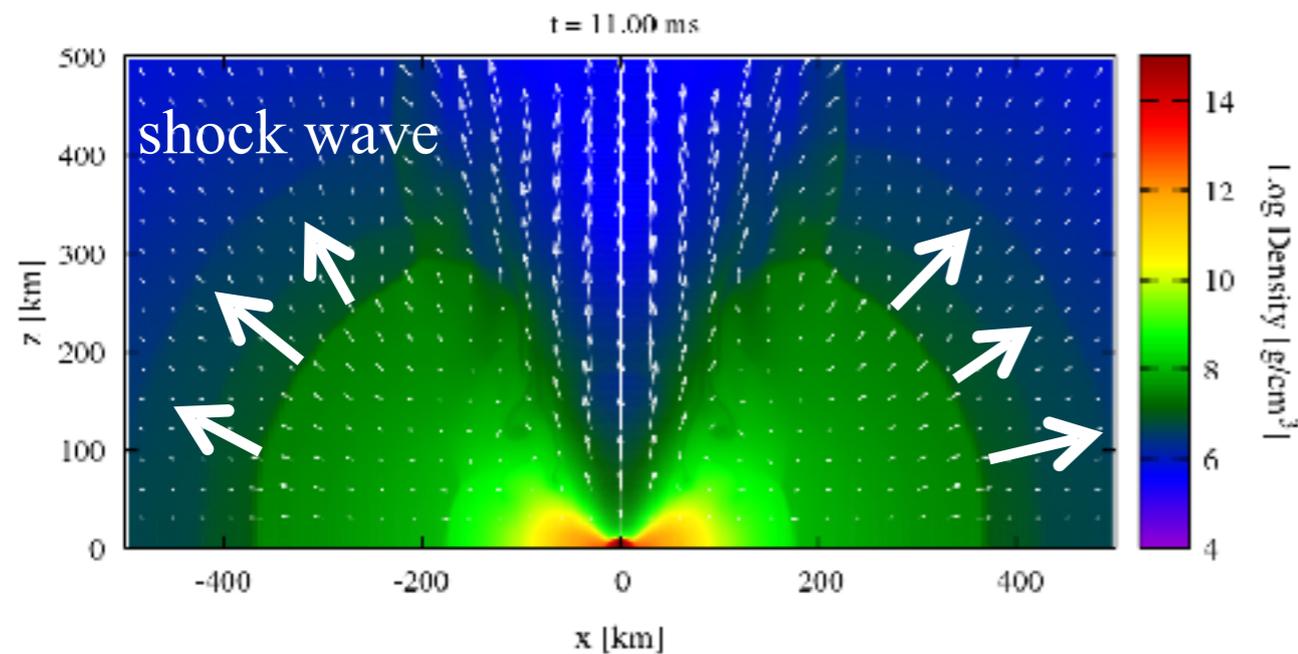
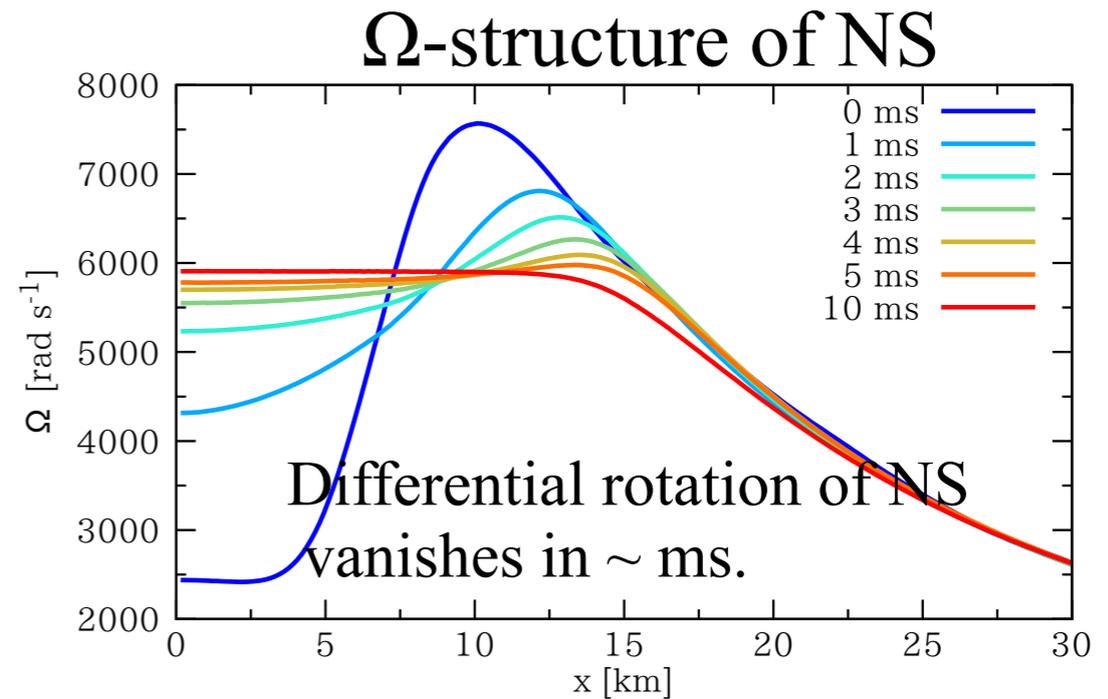
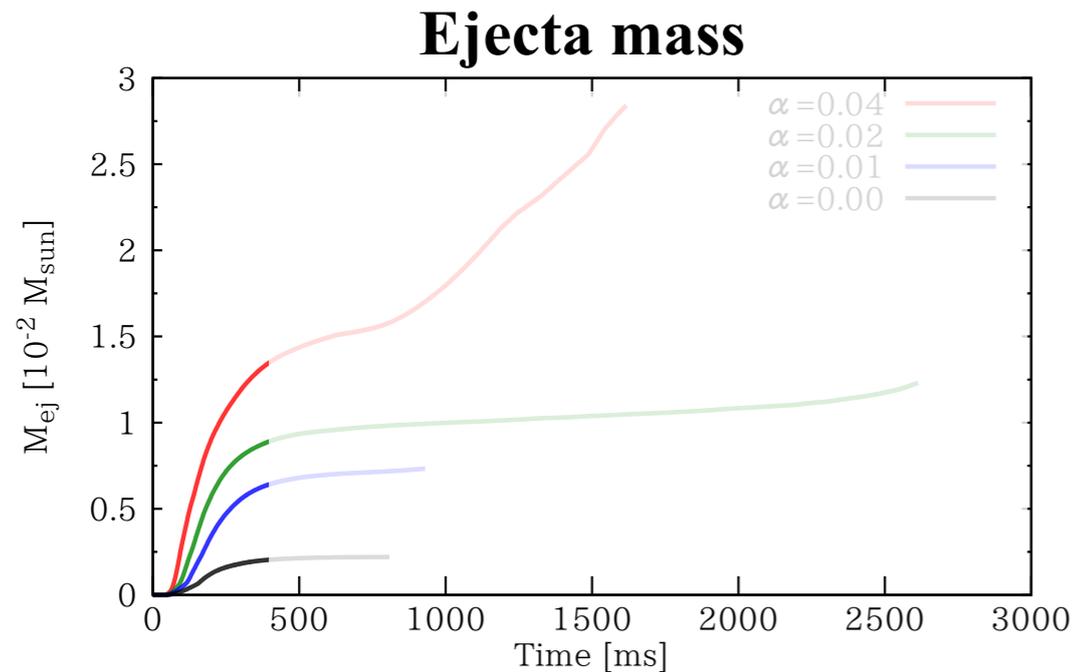
for early phase  $t < 400$  ms & late phase  $t > 1$  s

○  $\alpha = 0.02$

Same feature for  $t < 400$  ms,

but late-time ejection rate increases for  $t > 2$  s

# Early Viscosity-driven Ejecta

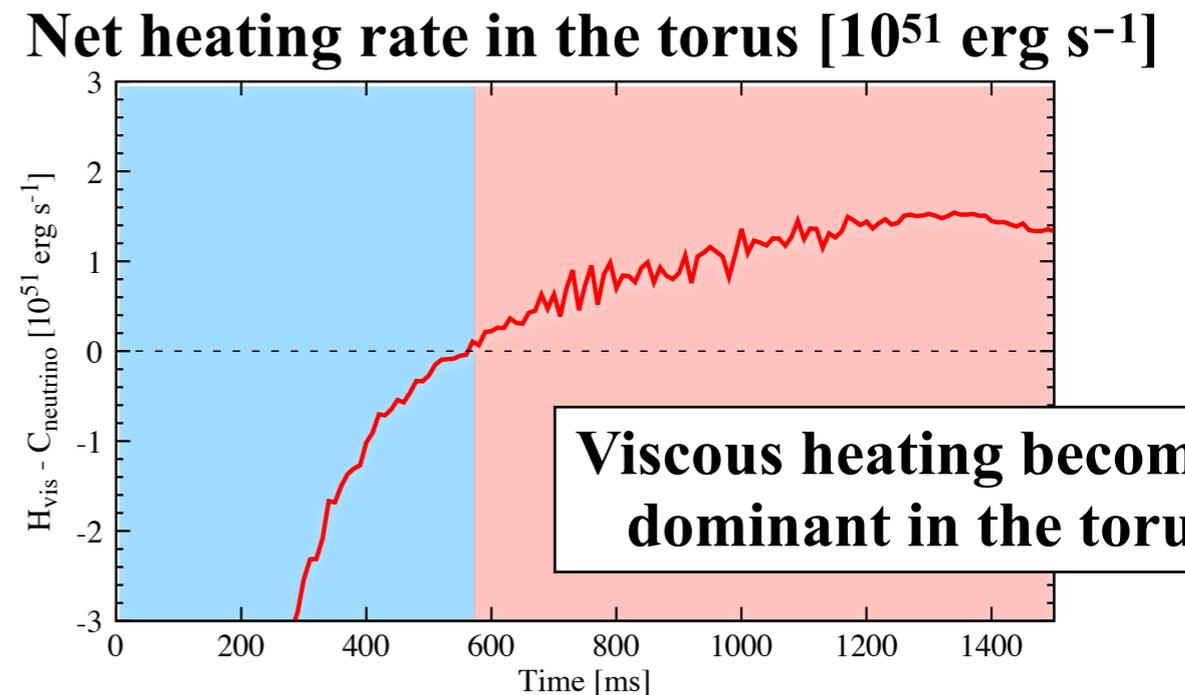
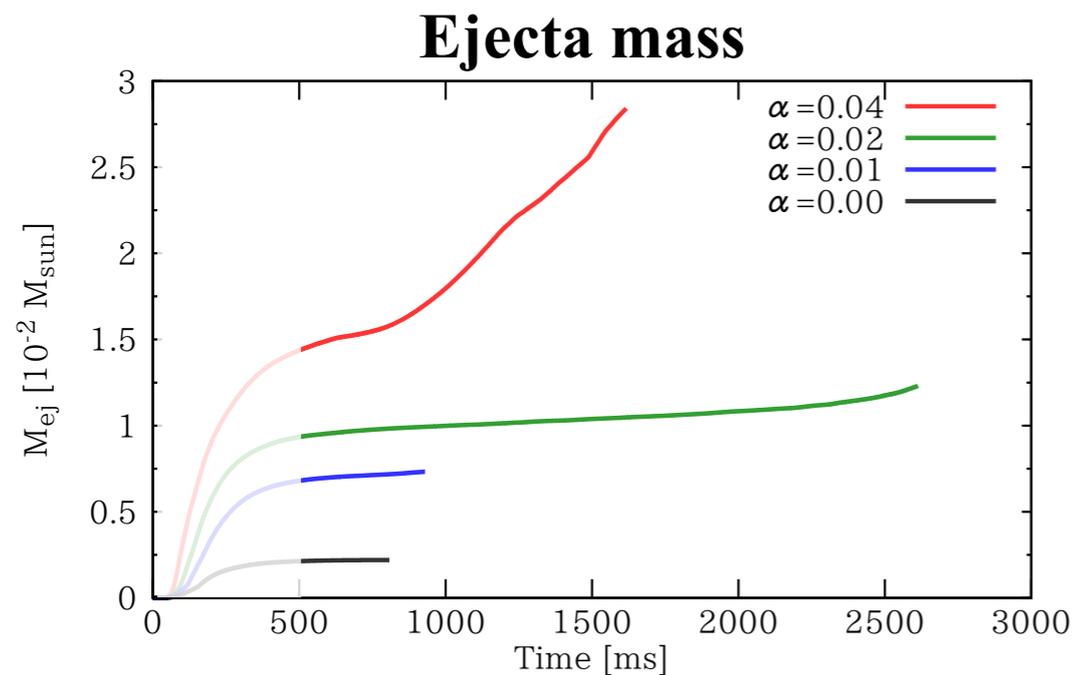


The MNS settles into rigidly rotation state due to viscous angular momentum transport.

Change of the equilibrium state of NS

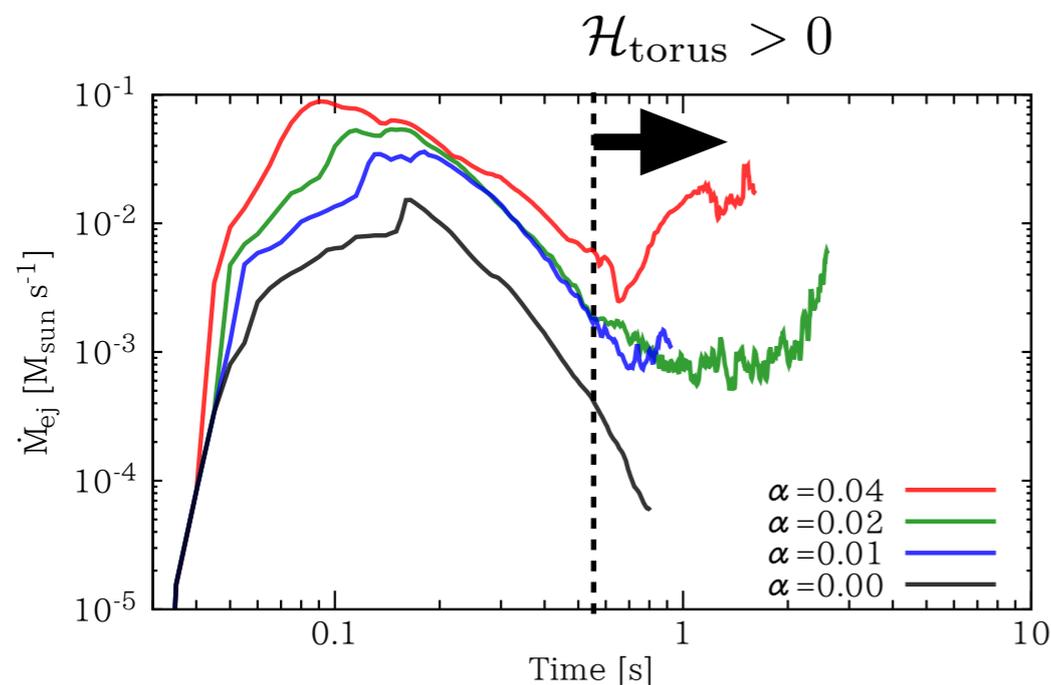
- A sound wave is formed in the vicinity of the NS, and it becomes a shock wave.
- The torus is swept-up by the shock wave and becomes unbound.

# Late-time Ejecta ( $t > 500$ ms)



After  $t > 600$  ms, viscous heating becomes dominant in the torus

The torus material is ejected in the viscous timescale in the torus ( $\sim 1$  s)

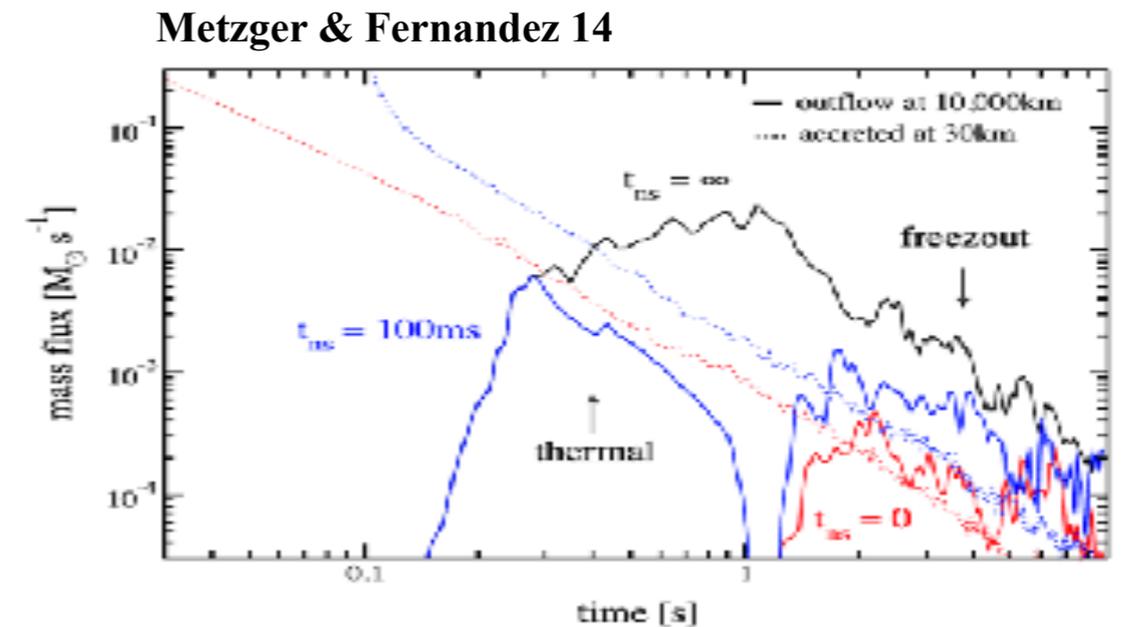
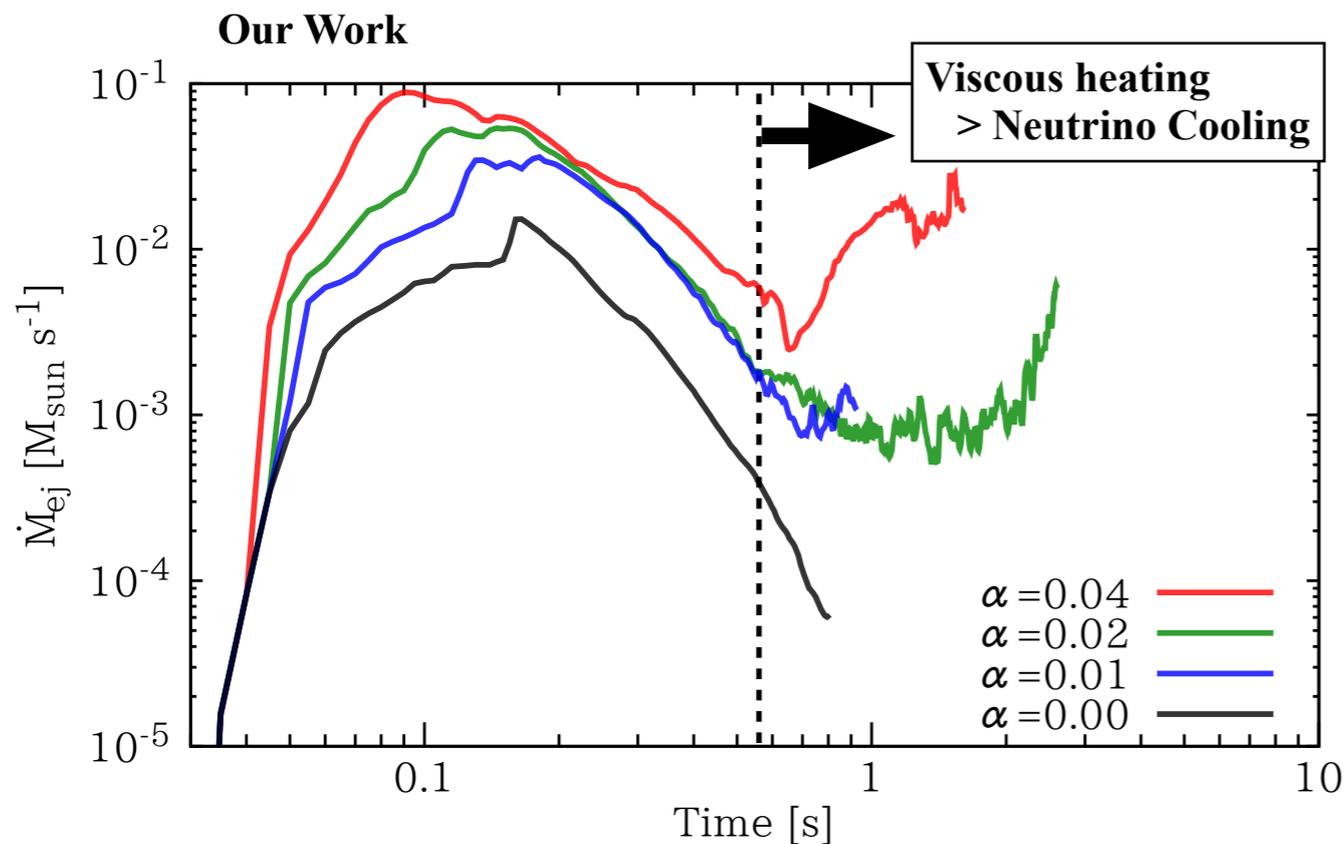


Mass ejection rate is estimated by  
(heating rate) / (specific binding energy)

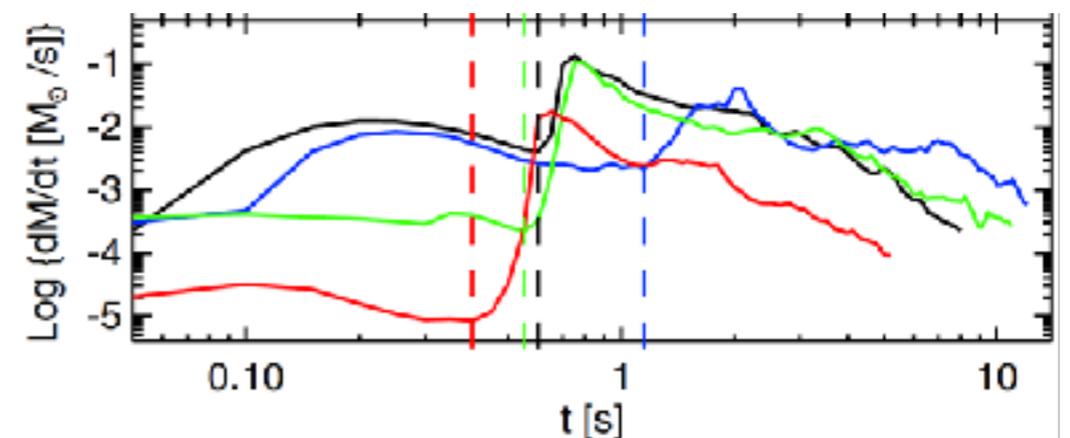
$$\dot{M}_{\text{ej}} \sim \frac{\mathcal{H}_{\text{torus}}}{GM_{\text{MNS}}/r}$$

$$\sim 10^{-2} M_{\odot} \text{ s}^{-1} \left( \frac{\mathcal{H}_{\text{torus}}}{10^{51} \text{ erg s}^{-1}} \right)^{-1} \left( \frac{M_{\text{MNS}}}{2.6 M_{\odot}} \right)^{-1} \left( \frac{r}{100 \text{ km}} \right)$$

# Total Amount of Viscosity-driven Ejecta



**Just et al. 15 (BH-torus)**



○ Late-time mass ejection rate

consistent with the model in Metzger & Fernandez 14 and Just et al. 15.

Even for low- $\alpha$  models, such mass ejection is likely to occur in  $\sim$  seconds.

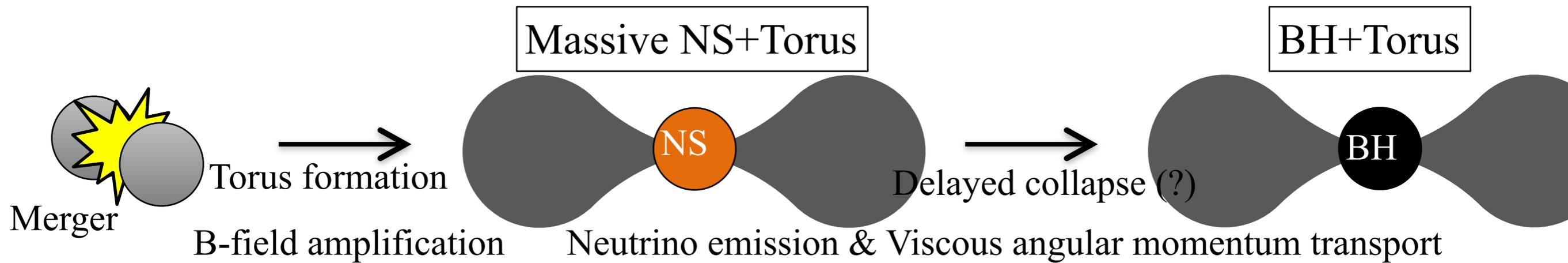
we confirmed their work by the most realistic simulation so far.

After  $\sim 1$ s, The mass accretion is negligible compared with the mass ejection.

**Significant fraction of the torus ( $\sim 90\%$ ) could be ejected (in the case of long-lived NS).**

**Torus mass at  $t = 1$  s is  $\sim 0.05 M_{\text{sun}}$   $\rightarrow$  ejecta mass could be  $\sim 0.04 M_{\text{sun}}$  ?**

# Mass Ejection Processes in Canonical Merger Remnant



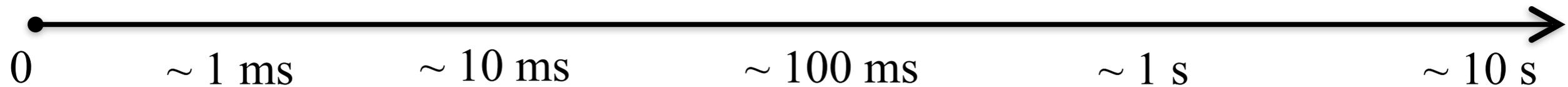
**Dynamical ejection** 0.001 – 0.02 Msun

**Change of the equiv. state of NS**  
~ 0.01 Msun

*Mass Ejection Processes*

**Viscosity & Neutrino-driven ejection**

> 0.01 Msun

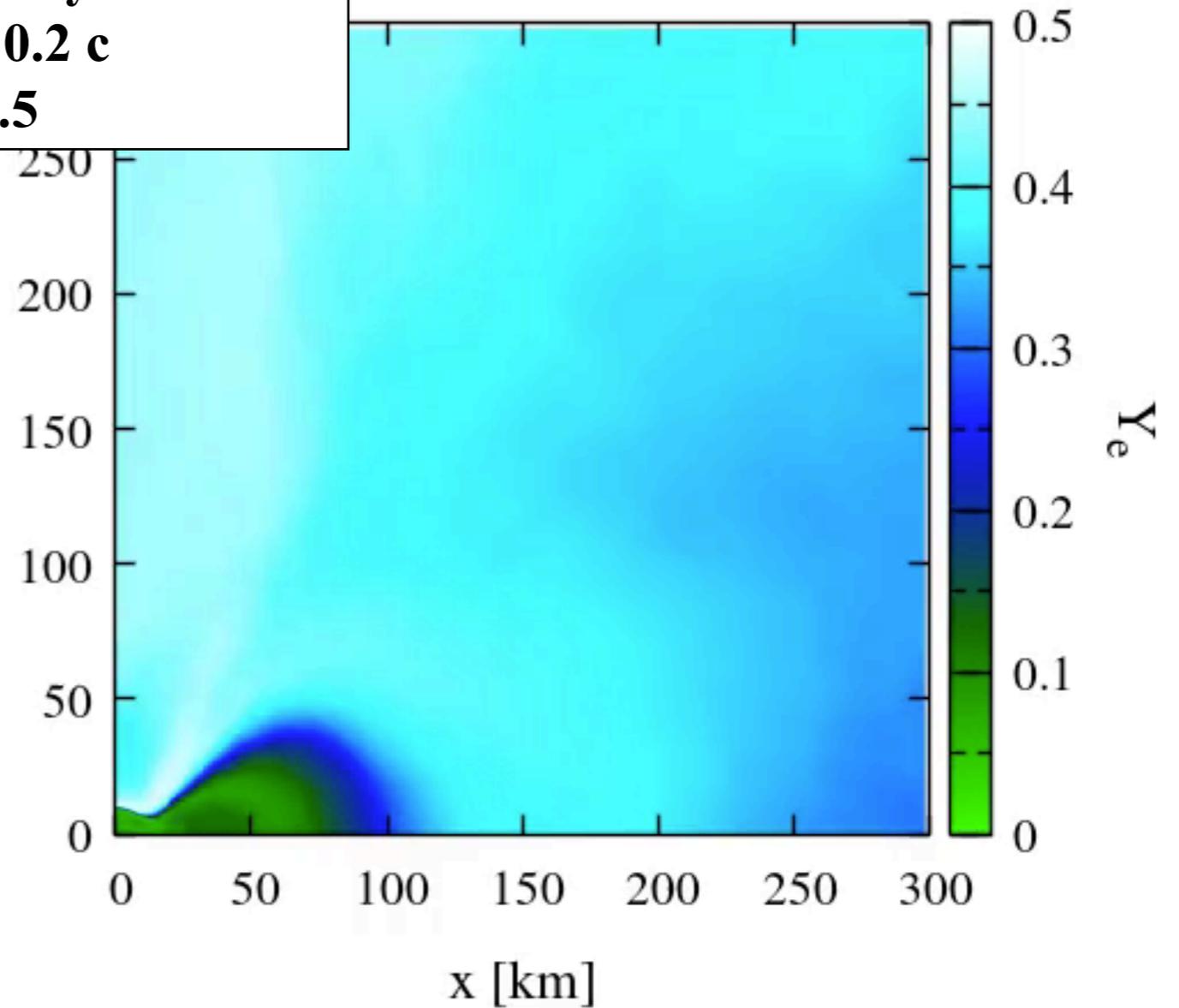
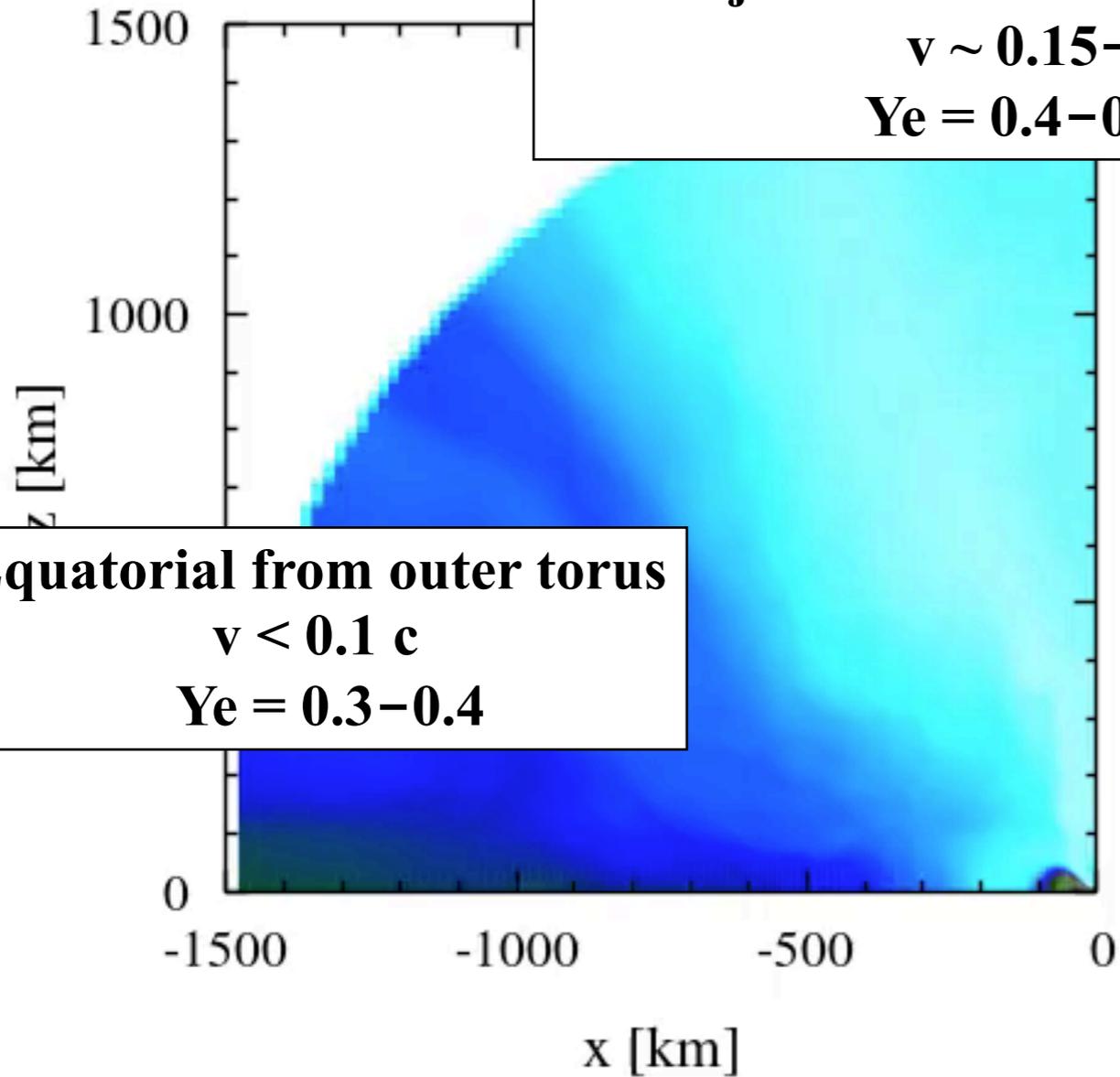


Time after merger

# Elemental Abundance and Implication of Electromagnetic Signal

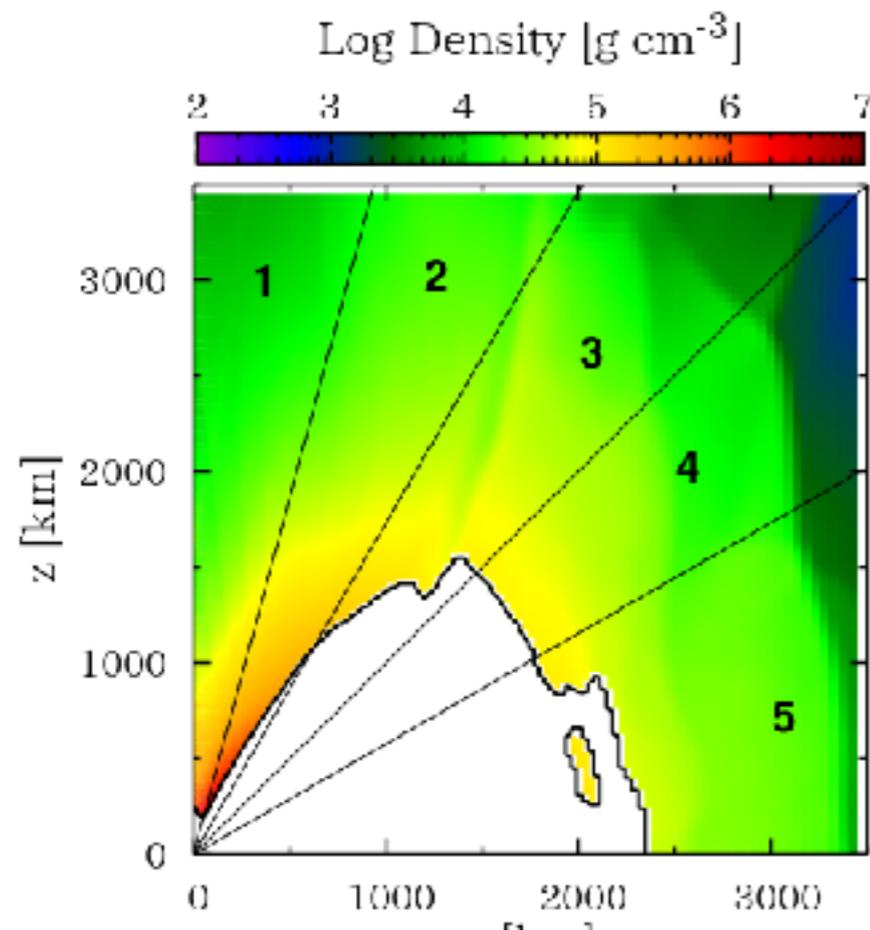
# Electron Fraction ( $Y_e$ ) of the Ejecta

$t =$  **Polar ejecta : irradiated by neutrinos**  
 $v \sim 0.15-0.2 c$   
 $Y_e = 0.4-0.5$



**Equatorial from outer torus**  
 $v < 0.1 c$   
 $Y_e = 0.3-0.4$

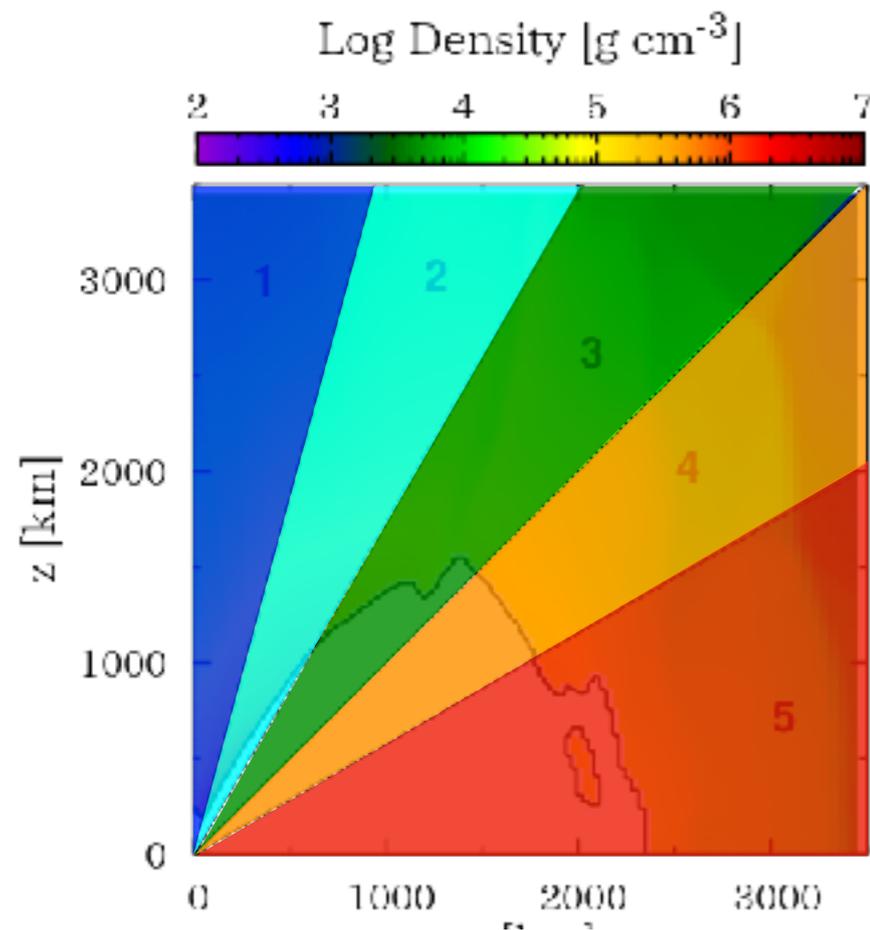
# Elemental Abundance in the Ejecta from the Merger Remnant



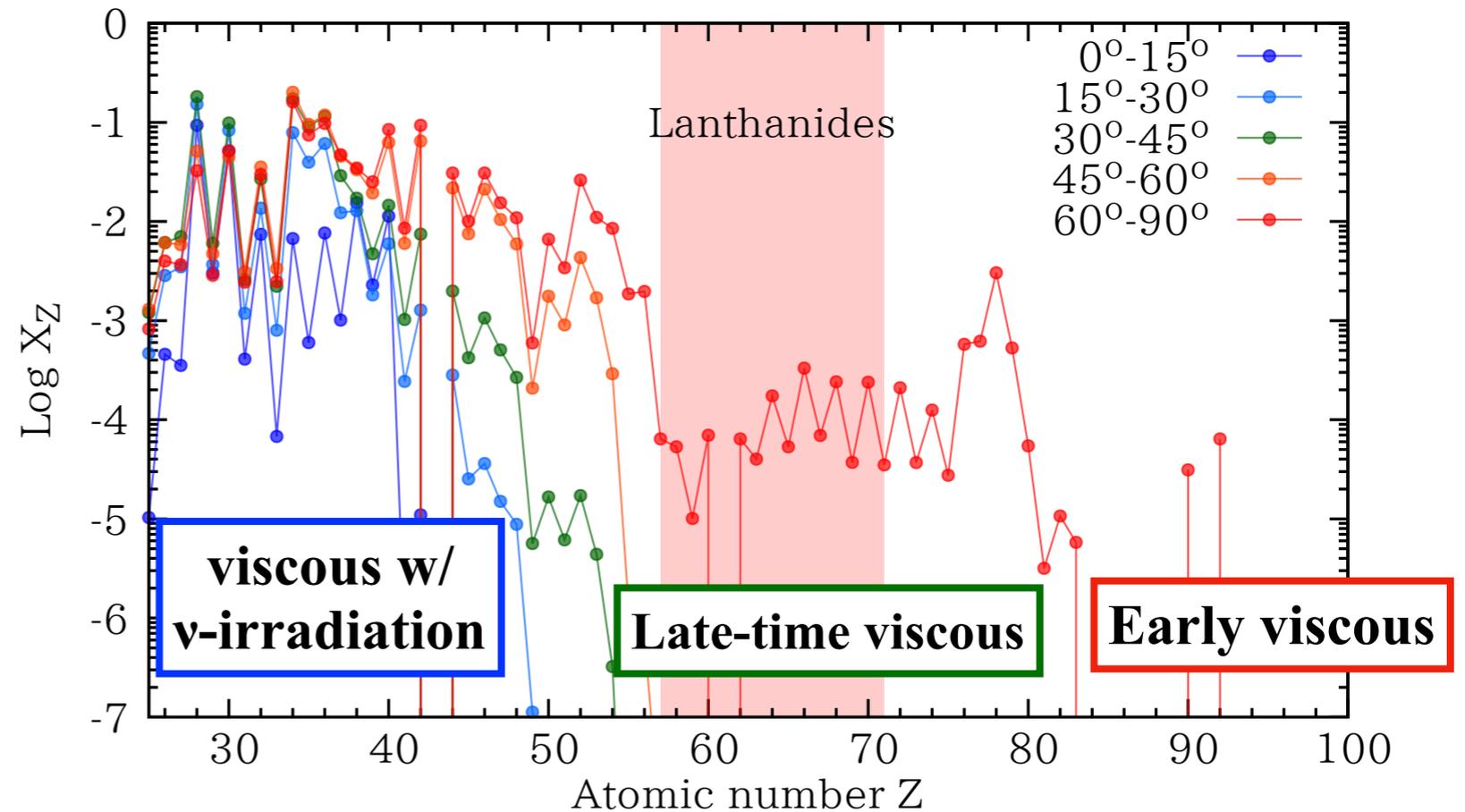
We performed nucleosynthesis calculations for the ejecta.

We investigate the angle-dependence of the elemental abundance.

# Elemental Abundance in the Ejecta from the Merger Remnant



## Elemental Mass Fraction in the Ejecta of Each bin



The late-time viscous ejecta dominates the total ejecta.

The ejecta from the merger remnant is (nearly) lanthanide-free.

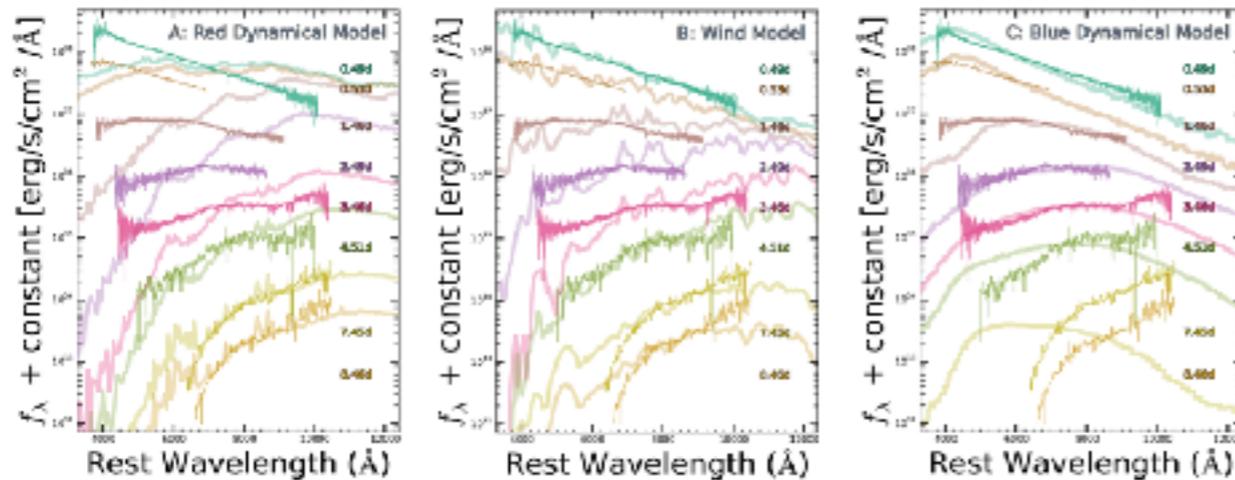
Therefore we can see more rapidly evolving, brighter emission with higher effective temperature.



# Observational Implication

○ Early Spectrum of Kilonova associated with GW170817 Shappee et al. 17, Kasen et al. 17

Lanthanide-rich  $\sim 0.2 c$     Lanthanide-poor  $\sim 0.03 c$     Lanthanide-poor  $\sim 0.3 c$



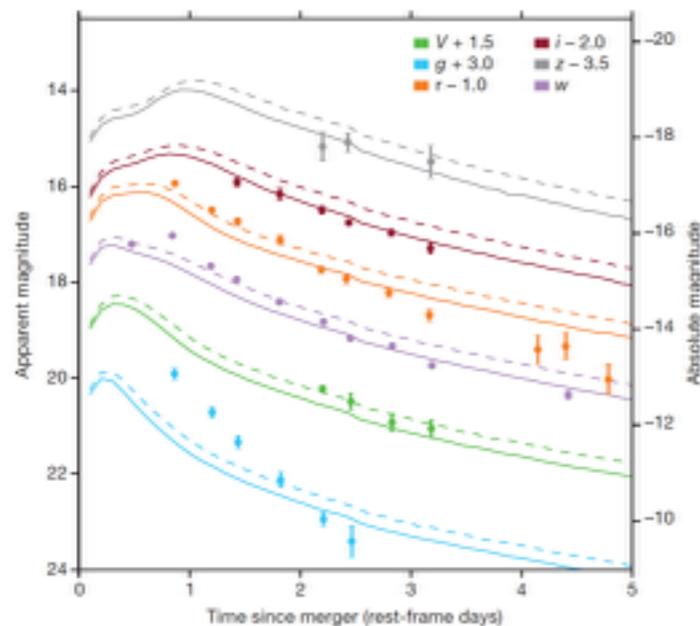
○ Featureless

→ Fast-expanding ( $\sim 0.3c$ ) ejecta is preferred for blue component.

↔ Viscosity-driven ejecta is slow ( $< \sim 0.1 c$ )

**Some acceleration mechanism of viscosity-driven ejecta is needed.**

○ Early Light Curve Arcavi et al. 17



Best-fit parameters of the light curve for  $< 3-4$  days (corresponding to blue component)

$M_{ej} \sim 0.02-0.025 M_{\odot}$

$V_{ej} \sim 0.3 c$

$X_{\text{Lanthanide}} \sim 10^{-4.5}$  (not completely Lanthanide-free)

↔ Ejecta from MNS-torus :  $X_{\text{Lanthanide}} < 10^{-8}$  (Too blue)

**If the MNS collapses into a BH in a short ( $\sim 100$  ms) timescale, The ejecta may be more neutron-rich, and the lanthanide fraction would be higher.**

# Possible Constraint on Maximum Mass of NS

- If magnetized NS survives for a long time, the rotational kinetic energy

$$T_{\text{rot}} \approx 1.1 \times 10^{53} \text{ erg} \left( \frac{M_{\text{MNS}}}{2.5 M_{\odot}} \right) \left( \frac{R}{15 \text{ km}} \right)^2 \left( \frac{\Omega}{7000 \text{ rad/s}} \right)^2,$$

could be injected into the ejecta by electromagnetic processes.

Since  $T_{\text{rot}} > M_{\text{ej}} c^2$ , if a fraction of the rotational energy is injected into the ejecta, the ejecta velocity could increase to  $\sim c$ , but such feature is not observed.

→ The MNS should collapse into a BH before its spin-down time.

$$\tau_B \approx 650 \text{ s} \left( \frac{B_p}{10^{15} \text{ G}} \right)^{-2} \left( \frac{M_{\text{MNS}}}{2.5 M_{\odot}} \right) \left( \frac{R}{15 \text{ km}} \right)^4 \left( \frac{\Omega}{7000 \text{ rad/s}} \right)^2$$

(If  $\gamma$ -rays detected by Fermi and INTEGRAL comes from BH+torus system, the MNS should collapse in  $\sim$  sec.)

The gravitational mass of the MNS  $\sim$  sec after the merger :  $\sim 2.6 M_{\text{sun}}$

( $\sim 0.15 M_{\text{sun}}$  is reduced due to GW &  $\nu$  emission)

This should be

the maximum mass of the rigidly rotating mass, which is  $\sim 0.4 M_{\text{sun}}$  larger than that of spherical cold NS

→ The maximum mass of the spherical cold NS  $\sim 2.6 - 0.4 = 2.2 M_{\text{sun}}$ .

# Summary

- We performed GR radiation viscous-hydrodynamics simulation of the remnant of the binary NS merger.
  - Variation of the equiv. state of the NS can be used for mass ejection.
  - Viscosity-driven ejecta would be dominant at  $t \sim 1$  s.
  - Late time ejecta will be nearly lanthanide-free.
    - Short , **Bright**, and **Blue** emission could be explained by this ejecta component
- Future prospects
  - Dependence on the binary system (masses) and nuclear EOSs.
  - Mass ejection from (MNS→BH)–torus system
  - Photon radiation transfer to obtain theoretical light curves.