

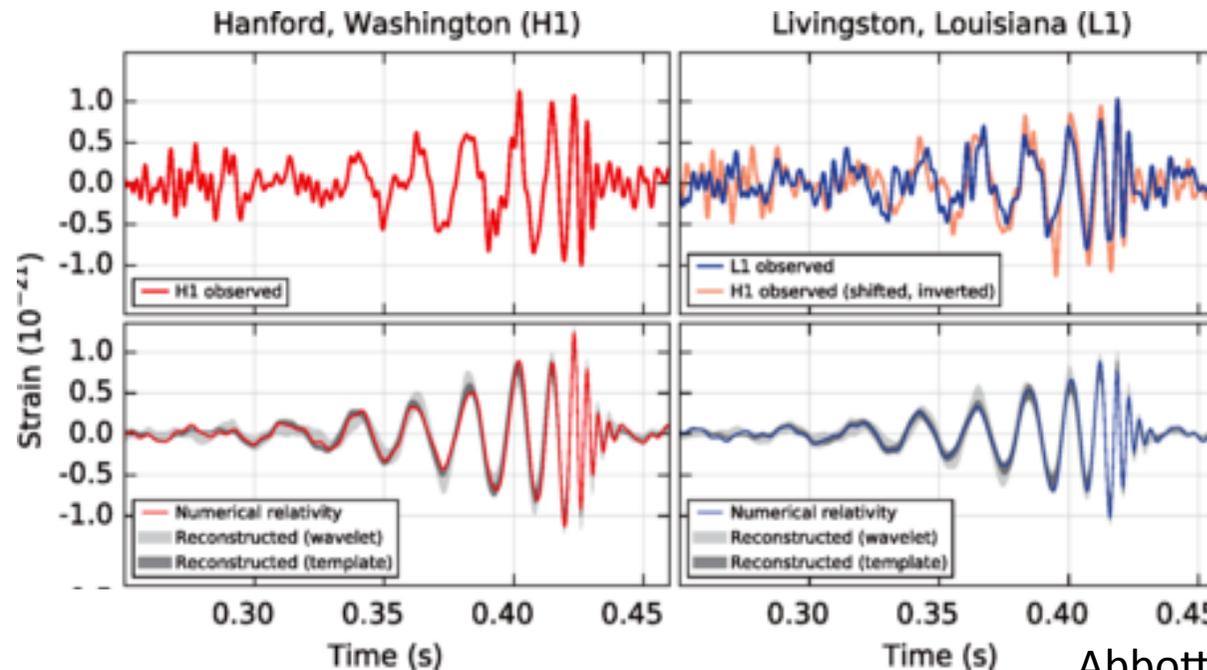
# 原始中性子星における星震学

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# Dawn of GW astronomy era

- First detection of GWs from BH-BH merger (GW150914)

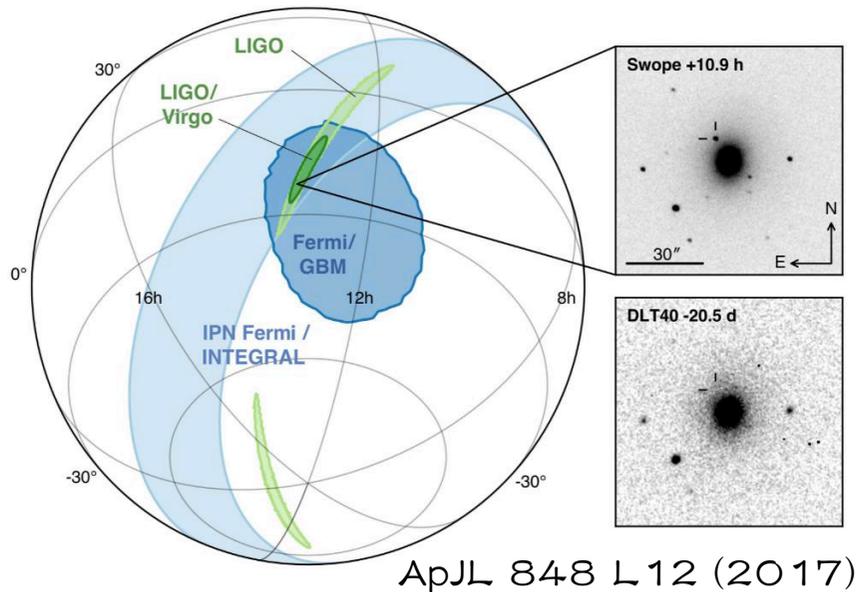


Abbott et al. 16

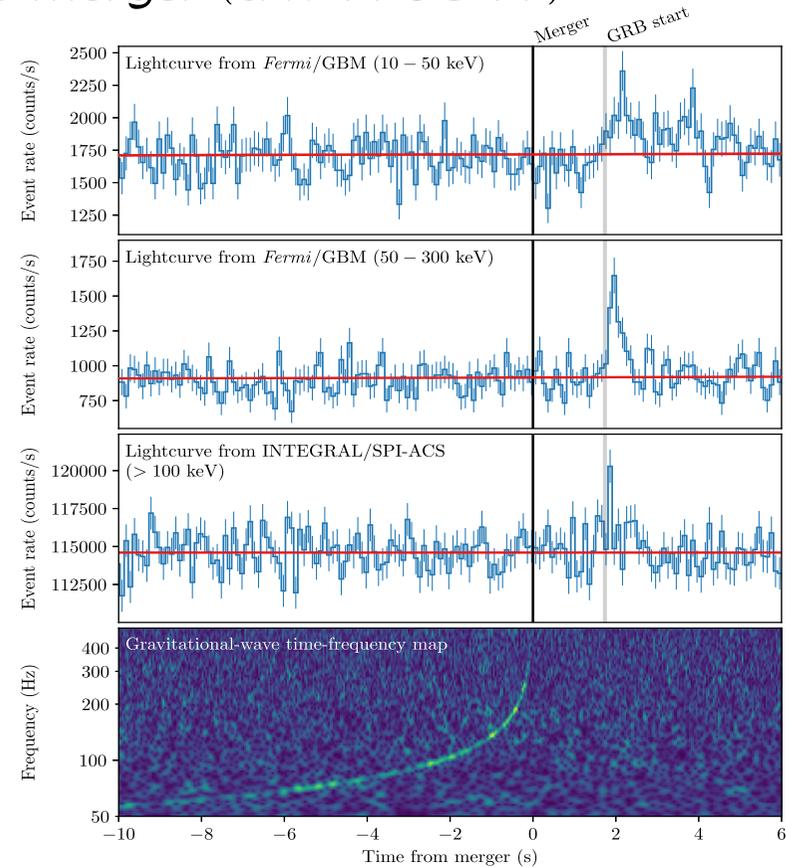
- $36M_{\odot}$ - $29M_{\odot}$  binary BH merger (410Mpc)
- GW151226 (Abbott et al. 16):  $14M_{\odot}$ - $7.5M_{\odot}$  BBH (440Mpc)
- GW170104 (Abbott et al. 17):  $31M_{\odot}$ - $19M_{\odot}$  BBH (880Mpc)

# Dawn of GW astronomy era

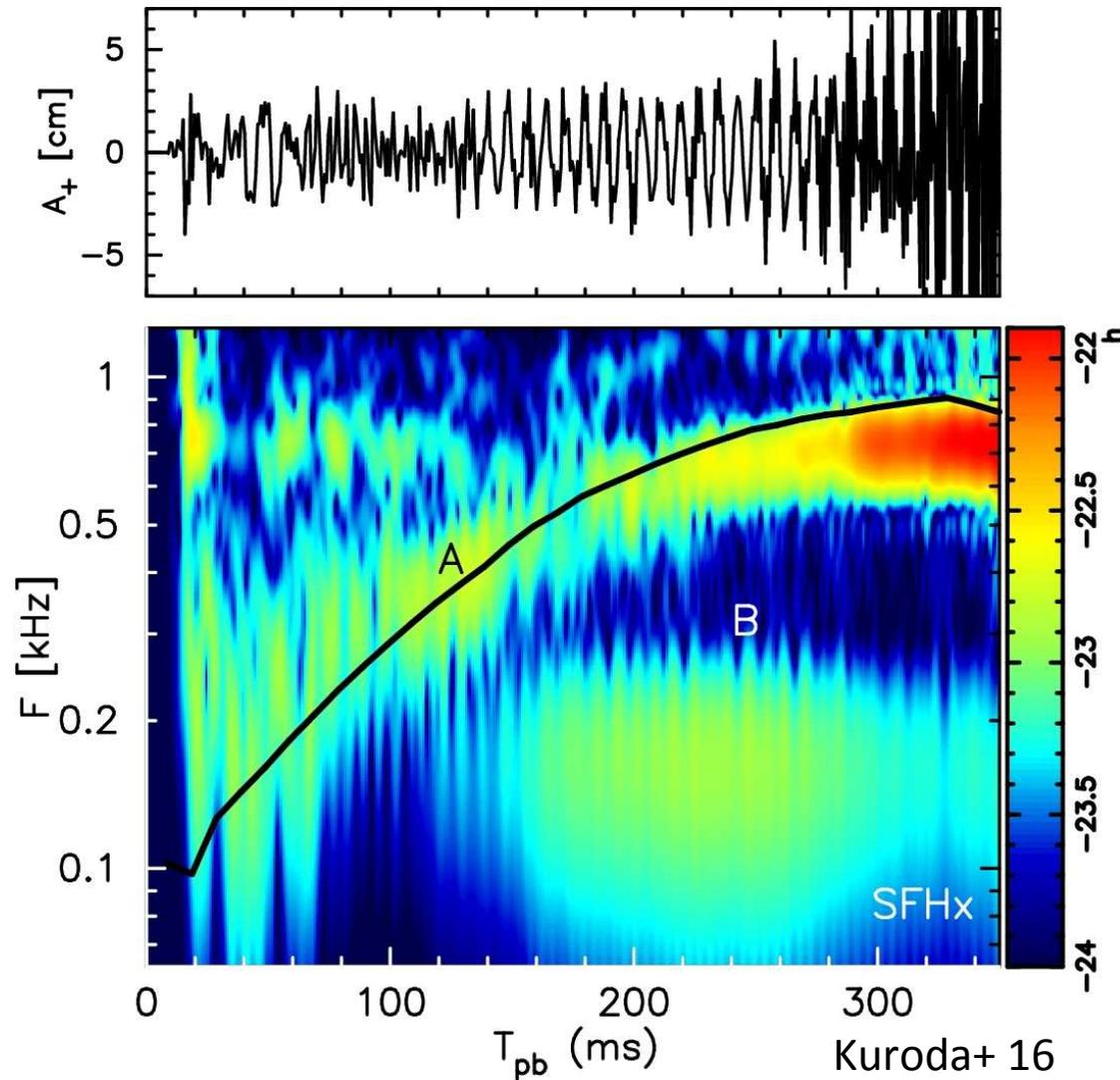
- First detection of GWs from NS-NS merger (GW170817)
  - first BNS + EM counter part
  - total mass =  $2.74M_{\odot}$  (40Mpc)



- promising GW sources;
  - BH-BH, BH-NS, and NS-NS mergers
  - supernovae



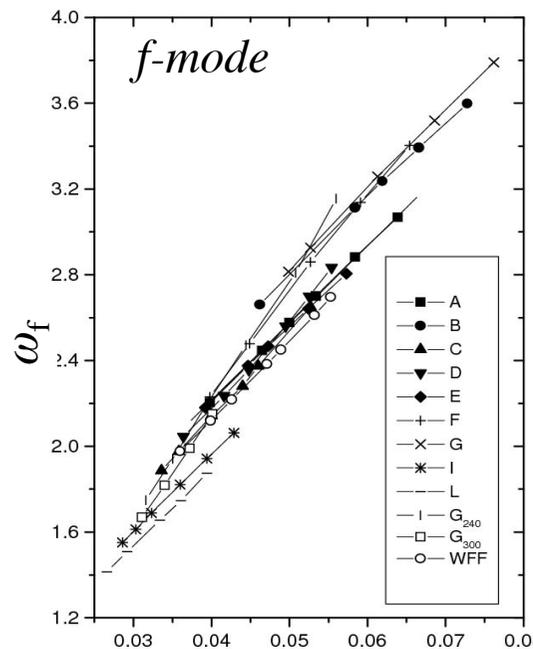
# GW from SN?



- Numerical simulations tell us the GW spectra.
- difficult
  - to extract physics of PNS and/or SN mechanism
  - to make a long-term numerical calculations
- We adopt the **perturbation approach** to determine the freq. from PNS.

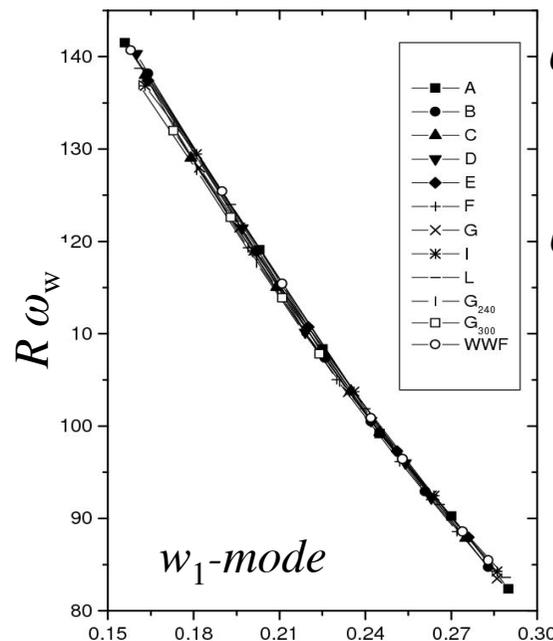
# Asteroseismology on Cold NSs

- via the observations of GW frequencies, one might be able to see the properties of NSs



average density

Andersson & Kokkotas (1998)



compactness

$$\omega_f \approx 0.78 + 1.64 \left[ \left( \frac{M}{1.4 M_\odot} \right) \left( \frac{10 \text{ km}}{R} \right)^3 \right]^{1/2}$$

$$\omega_w \approx \left( \frac{10 \text{ km}}{R} \right) \left[ 20.92 - 9.14 \left( \frac{M}{1.4 M_\odot} \right) \left( \frac{10 \text{ km}}{R} \right) \right]$$



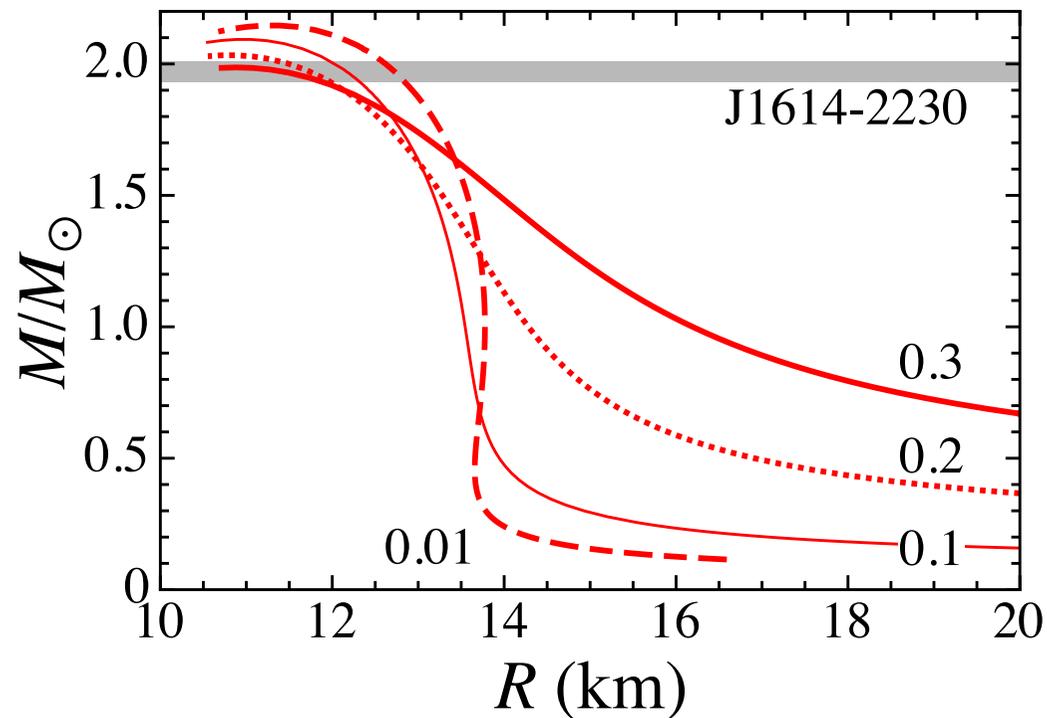
determination of  $(M, R)$

For PNS,

- Muellar et al. (13);
- Cerda-Duran et al. (13);
- Fuller et al. (15);
- Kuroda et al. (16)

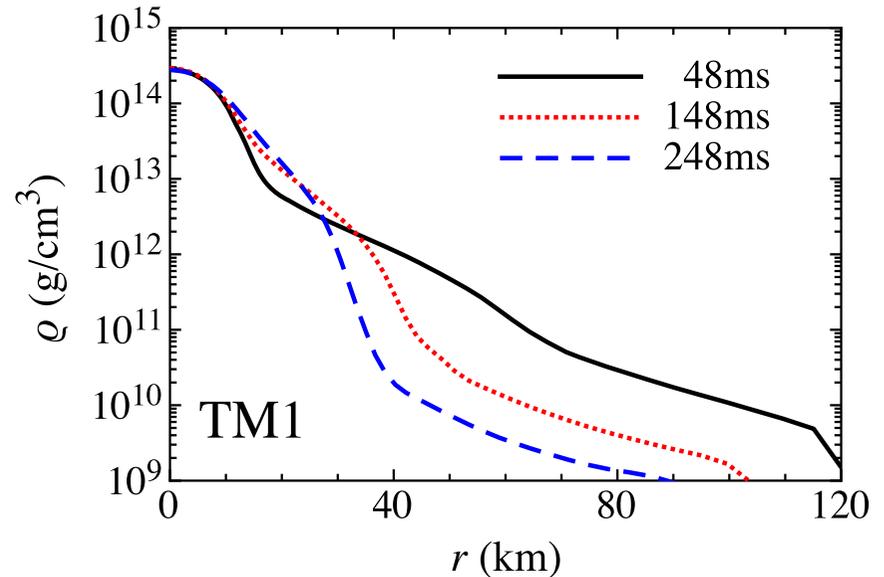
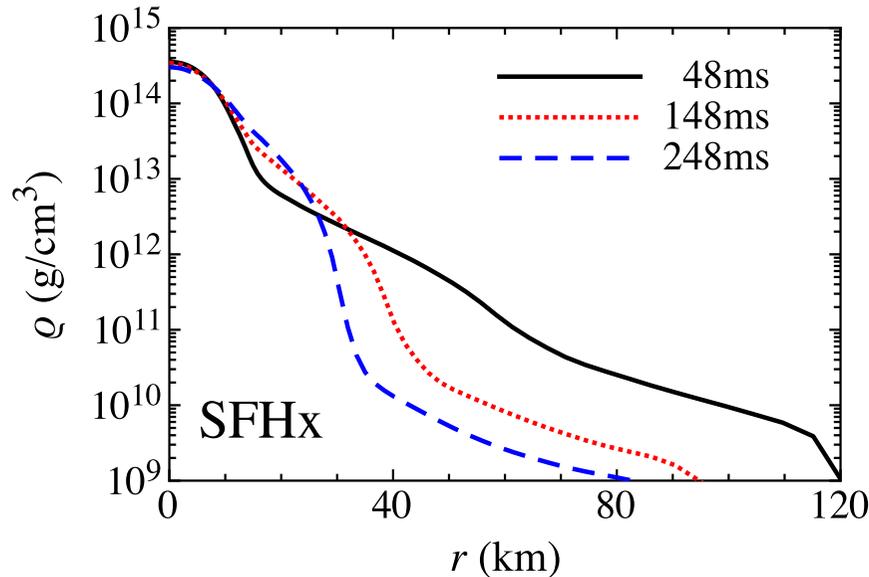
# Protoneutron stars (PNSs)

- Unlike cold neutron stars, to construct the PNS models, one has to prepare the profiles of  $Y_e$  and  $s$ .
  - for example,  
with LS220 and  $s = 1.5$  ( $k_B$ /baryon), but  $Y_e = 0.01, 0.1, 0.2,$  and  $0.3$



# PNS models

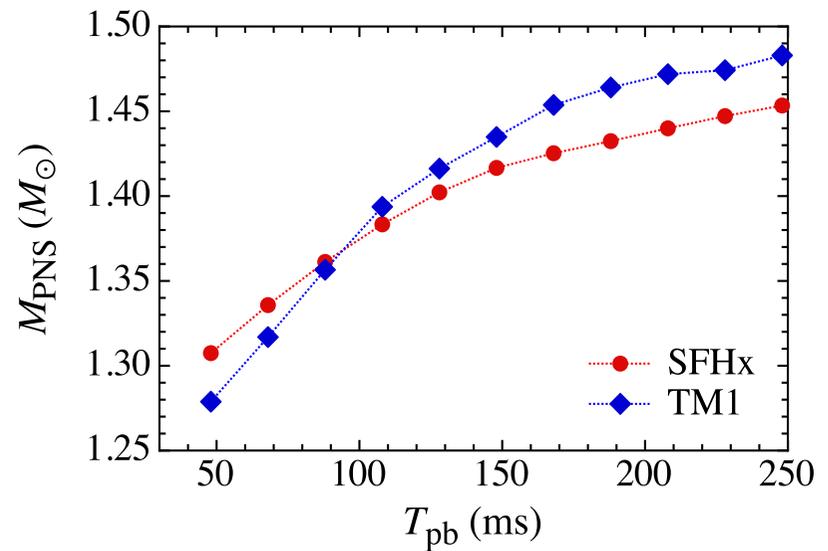
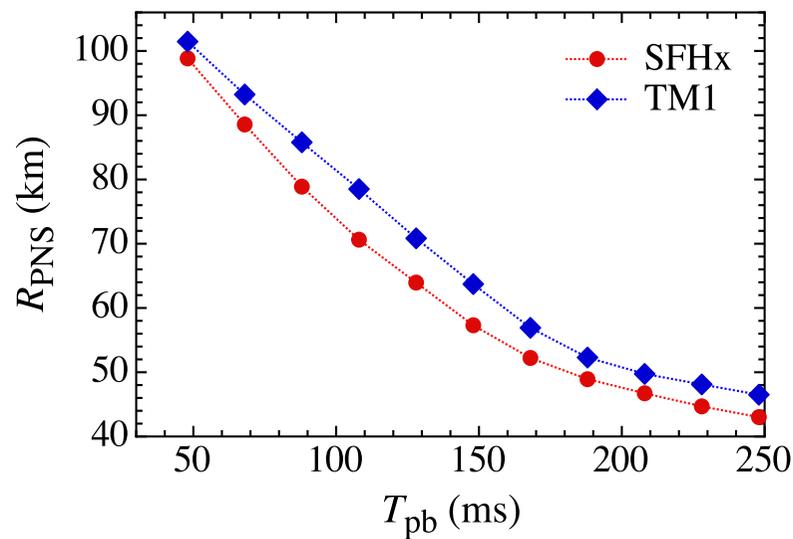
- we adopt the results of 3D-GR simulations of core-collapse supernovae (Kuroda et al. 2016)
  - progenitor mass =  $15M_{\odot}$
  - EOS : SFHx ( $2.13M_{\odot}$ ) & TM1 ( $2.21M_{\odot}$ )



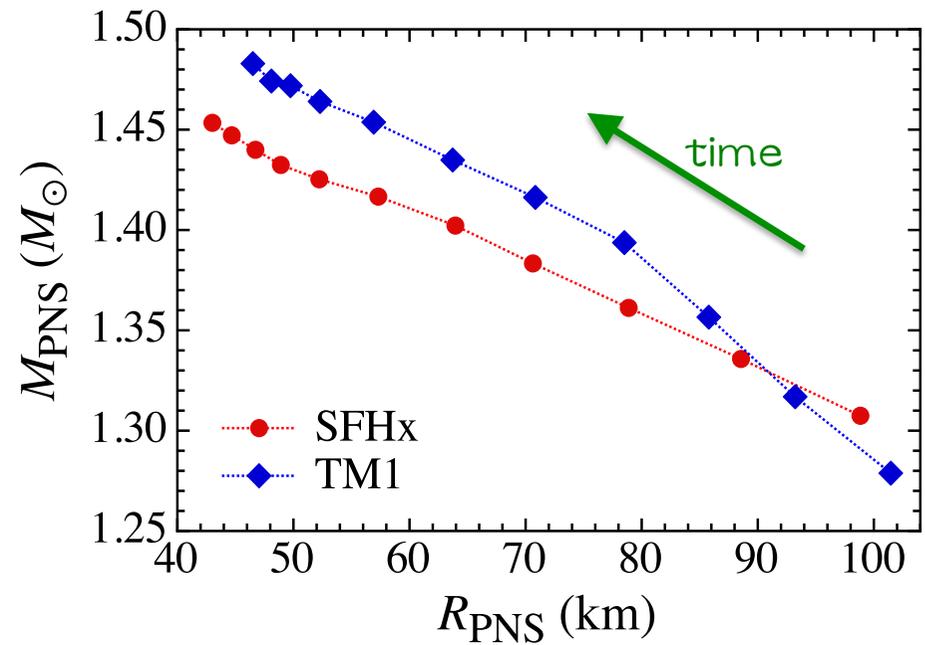
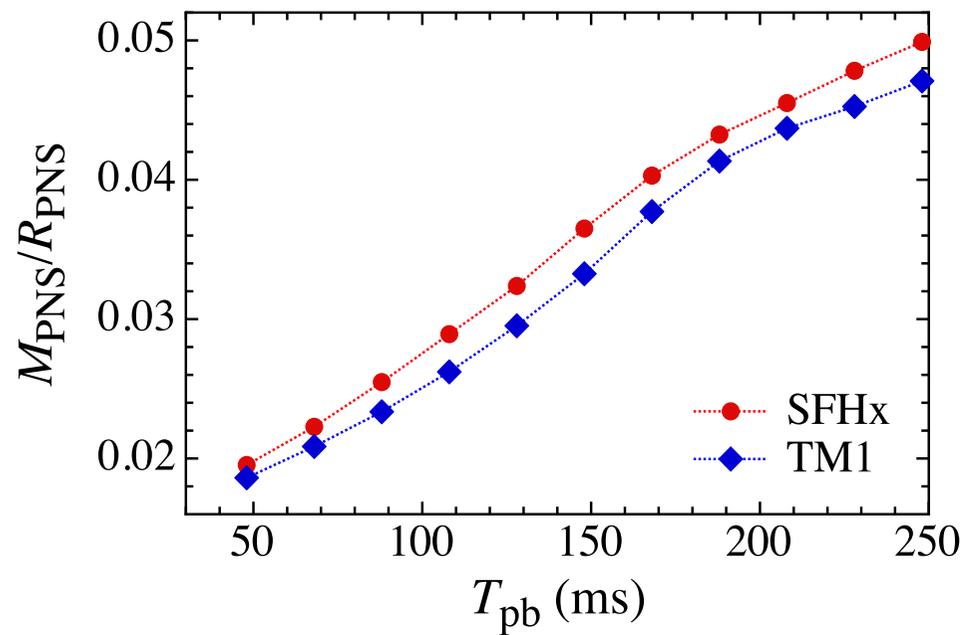
- $R_{\text{PNS}}$  is defined with  $\rho_s = 10^{10}$  g/cm<sup>3</sup>
- using the radial profiles as a background PNS model, the eigenfrequencies are determined.

# Mass & Radius

- $M_{\text{PNS}}$  is increasing by mass accretion
- $R_{\text{PNS}}$  is decreasing due to the cooling



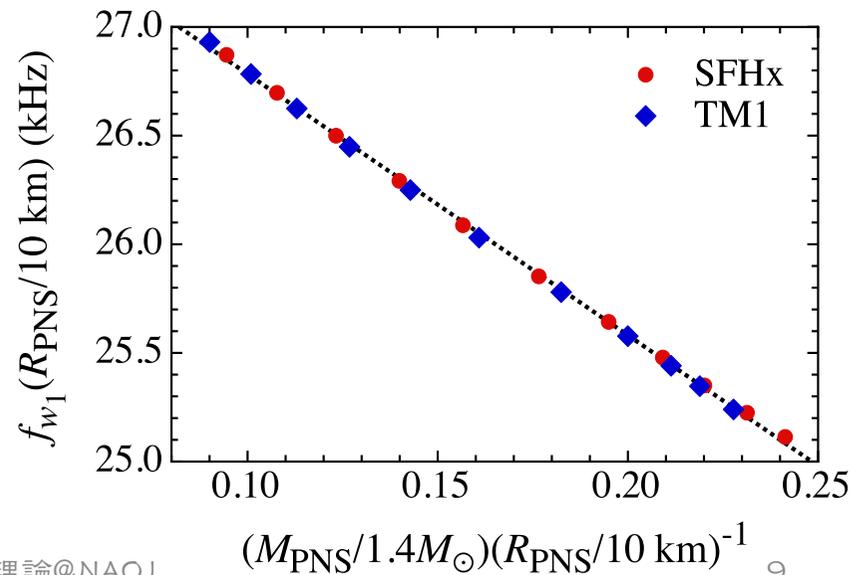
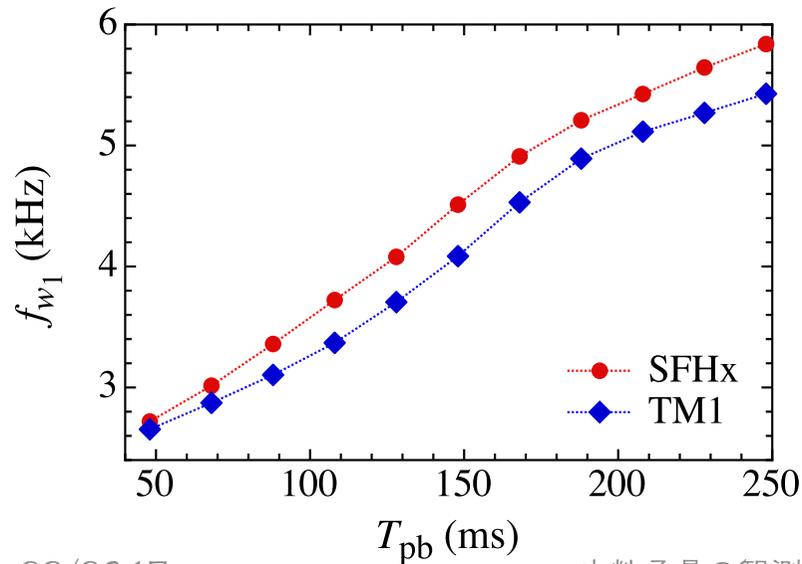
# M/R & M-R relation



# evolution of $w_1$ -modes

- frequencies depend on the EOS.
  - increasing with time
  - can be characterized well by  $M/R$
- as for cold NS, we can get the fitting formula, almost independent from EOS

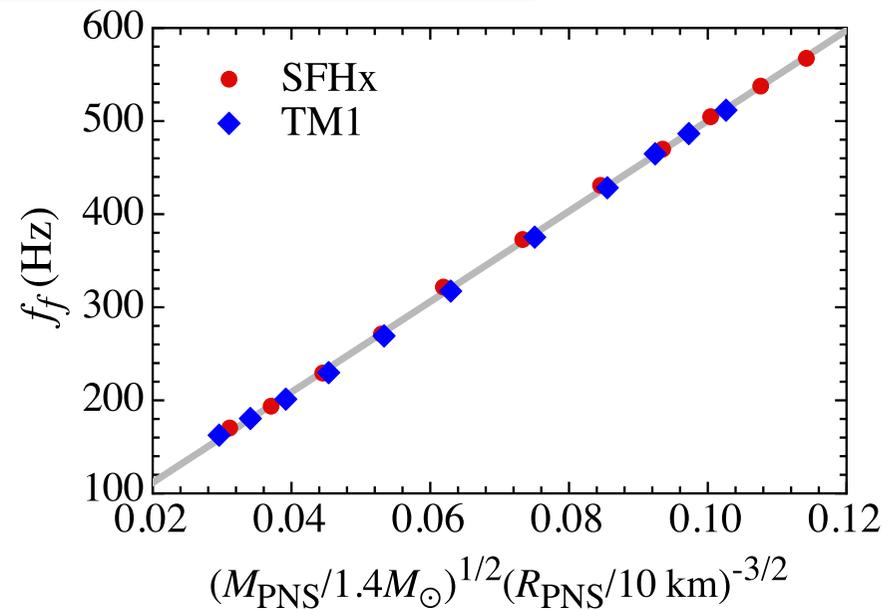
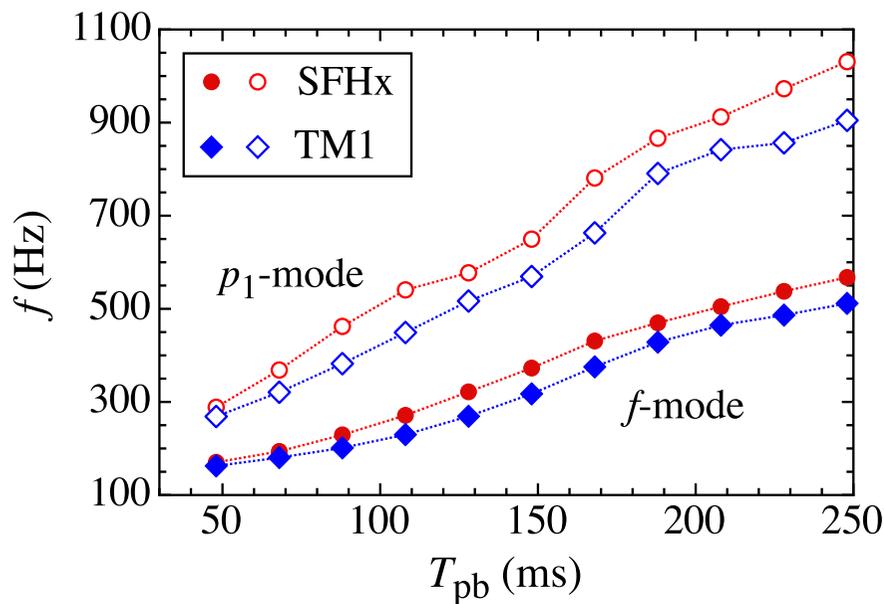
$$f_{w_1}^{(\text{PNS})} (\text{kHz}) \approx \left[ 27.99 - 12.02 \left( \frac{M_{\text{PNS}}}{1.4 M_{\odot}} \right) \left( \frac{R_{\text{PNS}}}{10 \text{ km}} \right)^{-1} \right] \times \left( \frac{R_{\text{PNS}}}{10 \text{ km}} \right)^{-1}$$



# evolution of f-mode

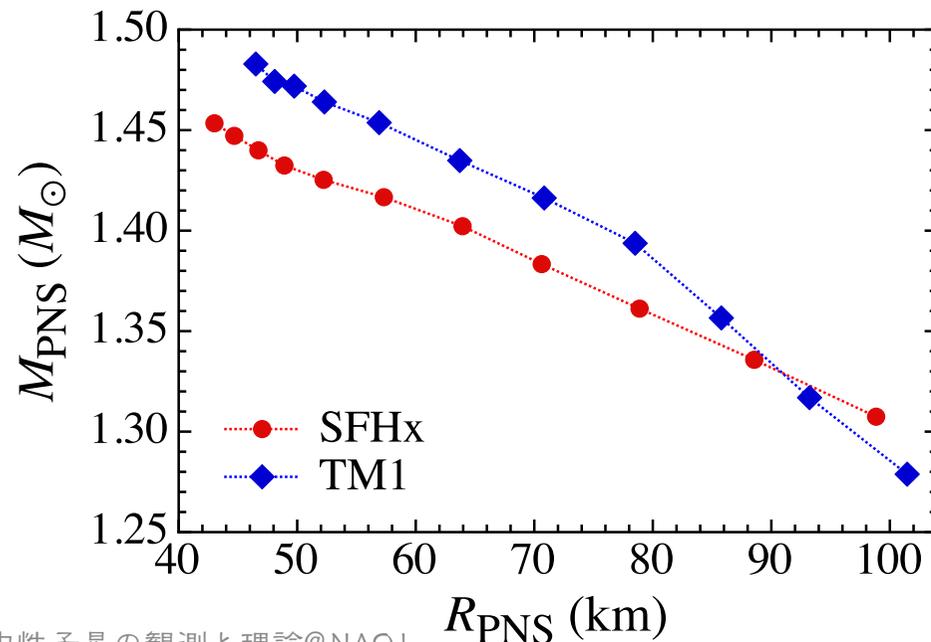
- frequencies can be expressed well by the average density independent of the EOS (and progenitor mass)
- we derive the fitting formula as a function of  $M/R^3$

$$f_f^{(\text{PNS})} (\text{Hz}) \approx 14.48 + 4859 \left( \frac{M_{\text{PNS}}}{1.4 M_{\odot}} \right)^{1/2} \left( \frac{R_{\text{PNS}}}{10 \text{ km}} \right)^{-3/2}$$



# determination of EOS

- with  $f^-$  &  $w_1$ -modes GW observations, one can get two independent properties at each time after core bounce, which are combination of  $M_{\text{PNS}}$  &  $R_{\text{PNS}}$
- one can determine  $(M_{\text{PNS}}, R_{\text{PNS}})$  at each time after core bounce  
→ determination of the EOS
- unlike cold NS cases, in principle one can determine the EOS even with ONE GW event !



# detectability of $w_1$ -modes

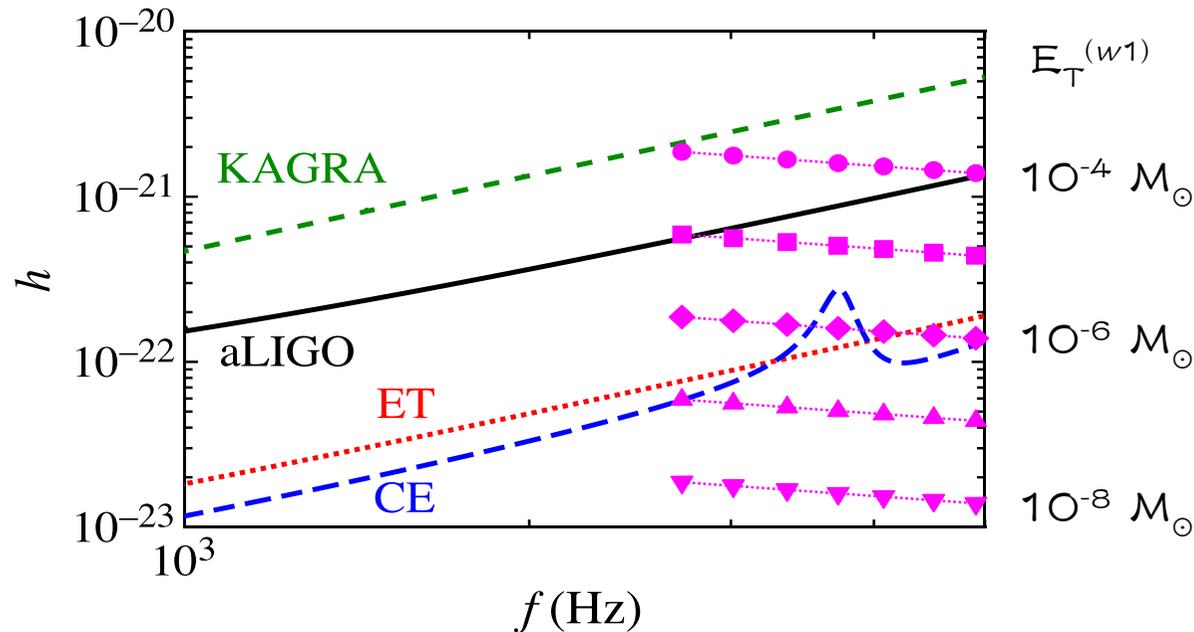
- effective amplitude of  $w_1$ -modes

$$h_{\text{eff}}^{(w_1)} \sim 7.7 \times 10^{-23} \left( \frac{E_{w_1}}{10^{-10} M_\odot} \right)^{1/2} \left( \frac{4 \text{ kHz}}{f_{w_1}} \right)^{1/2} \left( \frac{10 \text{ kpc}}{D} \right)$$

Andersson & Kokkotas (1996, 1998)

$$\frac{E_{w_1}}{E_T^{(w_1)}} \approx \frac{\tau_{w_1}}{T_{w_1}}$$

$E_{w_1}$  : energy for each time step  
 $E_T^{(w_1)}$  : total radiation energy in  $w_1$ -modes



# conclusion

- We examine the frequencies of gravitational waves radiating from PNS after bounce.
  - we derive the empirical formula of  $w_1$ - &  $f$ -modes independent of the EOS
  - via the GW observation from PNS, one would see  $M_{\text{PNS}}$  &  $R_{\text{PNS}}$  evolution
- in principle, even with ONE GW event from supernova, one could determine the EOS for high density region.