



KATHOLIEKE UNIVERSITEIT
LEUVEN



Laser
LISOL
Source



In-Gas Laser Ionization and Spectroscopy (IGLIS) of radioactive atoms at LISOL

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1. In-Gas-Cell Laser Ionization, selective production of RIB for nuclear spectroscopy
2. In-Gas-Cell Laser Spectroscopy , $^{57-59}\text{Cu}$, $^{97-101}\text{Ag}$
3. In-Gas-Jet Laser Spectroscopy



In-Gas-Cell Laser Spectroscopy of $^{57,59}\text{Cu}$

$\text{Cu}^+ + e^-$ Autoionizing State

First Ionization Limit
62317.4 cm^{-1}

$\lambda_2 = 441.6 \text{ nm}$

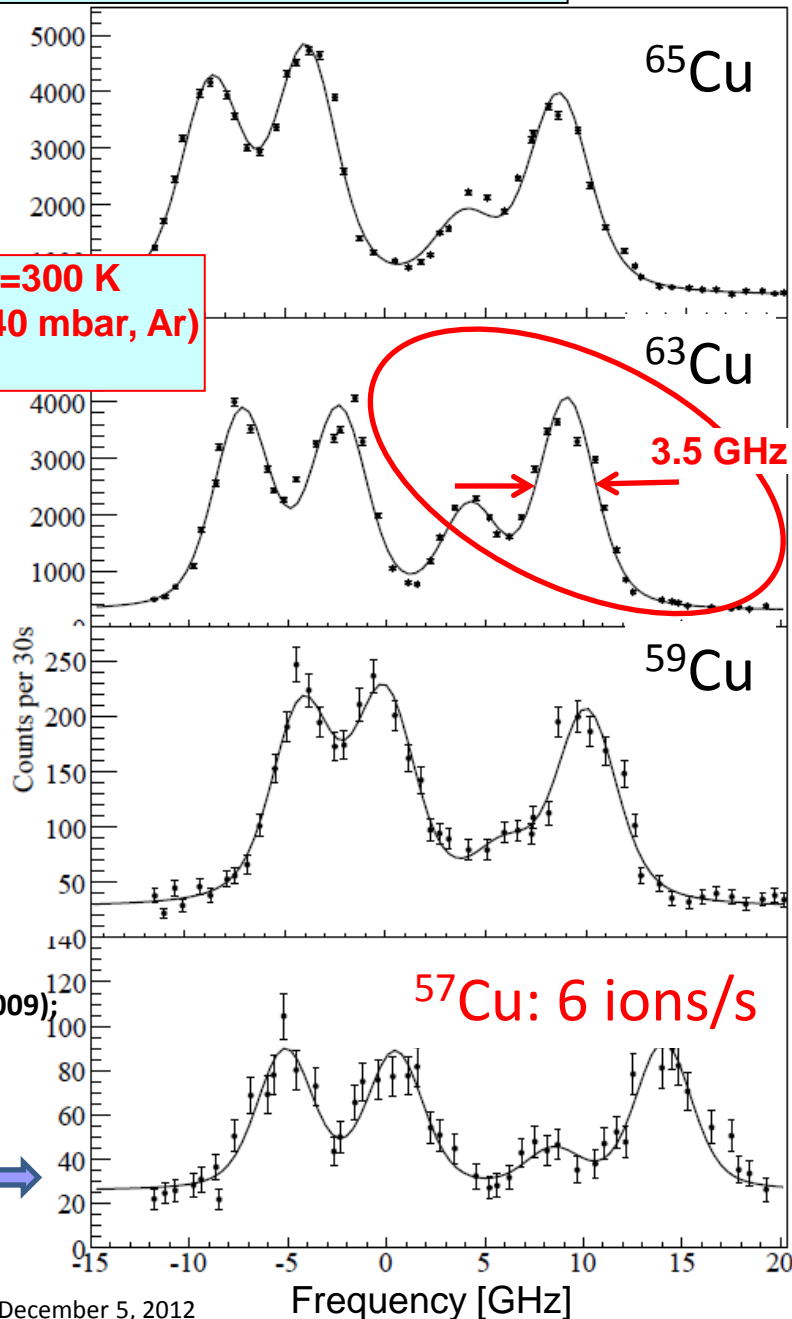
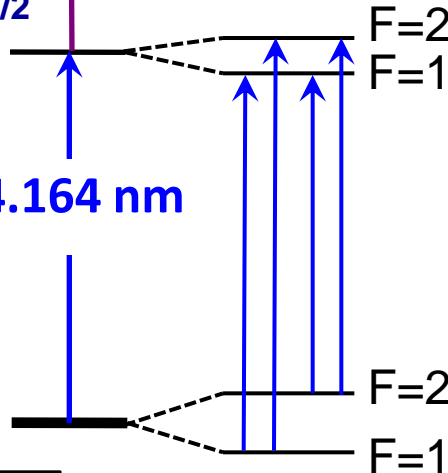
Doppler broadening, $T=300 \text{ K}$
Pressure broad. ($P = 140 \text{ mbar, Ar}$)
Laser bandwidth

$4\text{P}^0_{1/2}$
40943.73 cm^{-1}

$\lambda_1 = 244.164 \text{ nm}$

$2\text{S}_{1/2}$

CuI: ground state



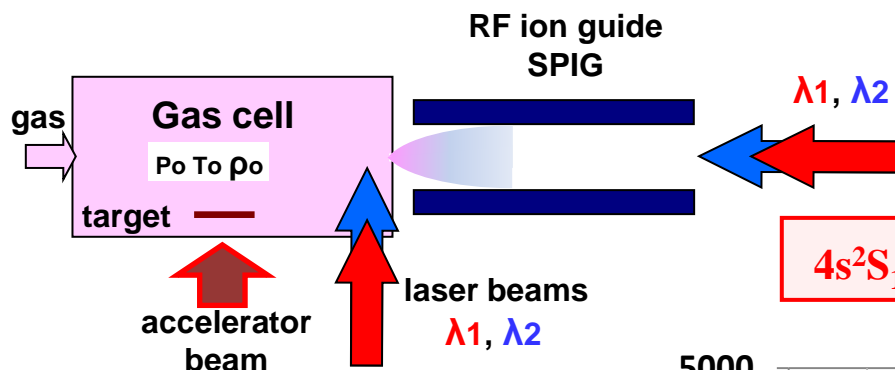
^{57}Cu : 6 ions/s

$^{58}\text{Ni}(p, 2n)^{57}\text{Cu}$ ($T_{1/2} = 199 \text{ ms}$)

$$\mu(^A\text{Cu}) = \frac{A_{hf}(^A\text{Cu})}{A_{hf}(^{63}\text{Cu})} \mu(^{63}\text{Cu})$$

T. Cocolios et al. PRL 103, 102501 (2009);
Phys. Rev. C 81, 014314 (2010)

Doppler and Collision Contributions to the Spectral Line Width



$4s^2S_{1/2} - 4p^2P_{1/2}$, 327.4 nm ^{63}Cu transition, $\nu_0 = 30535.3 \text{ cm}^{-1}$

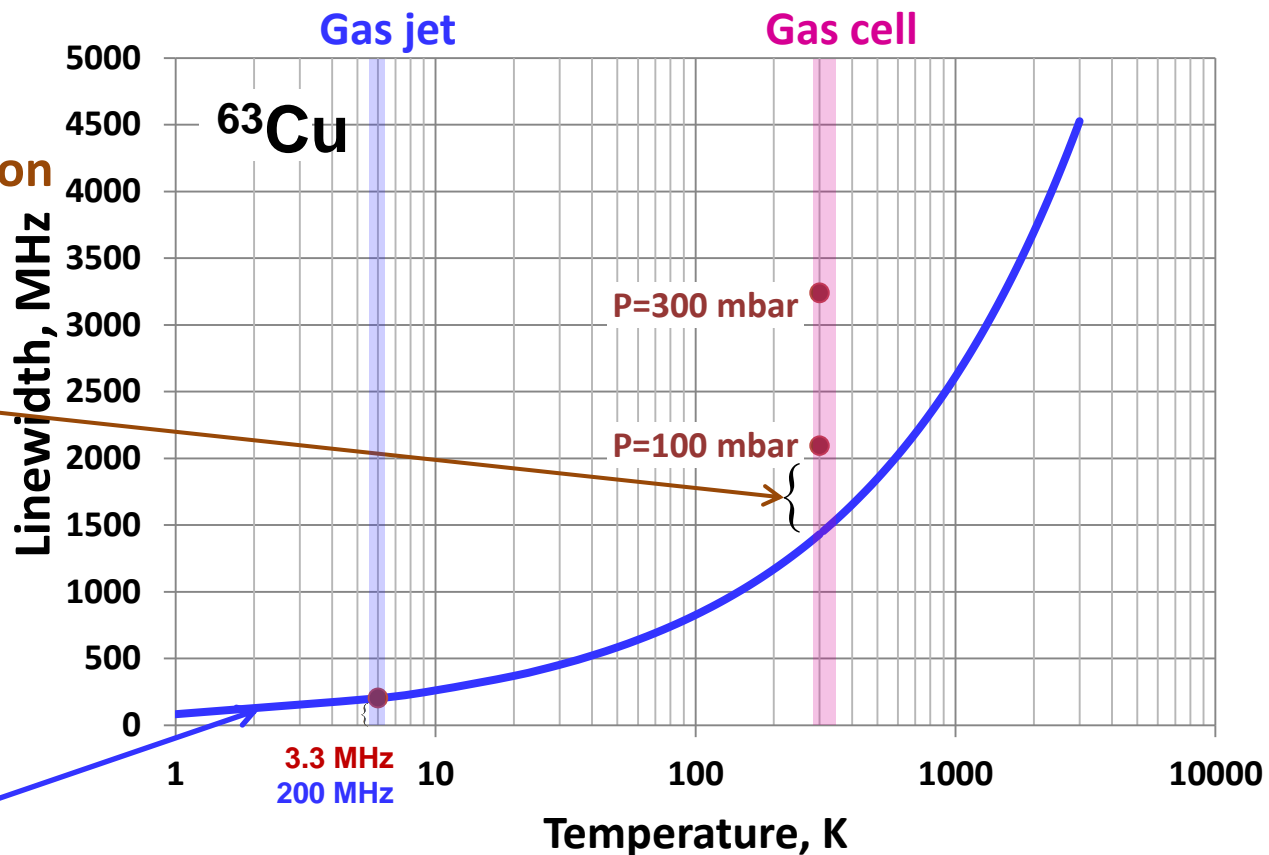
Collision/pressure contribution

$$\Delta \nu_{coll} = \gamma_{coll} \times \rho$$

γ_{coll} - collision broadening coefficient, $1.5 \cdot 10^{-20} \text{ cm}^{-1}/\text{cm}^{-3}$ (8 MHz/mbar)
 ρ - gas density (atom/cm^3)

Doppler contribution

$$\Delta \nu_{Doppler} = 2\sqrt{\ln 2} \frac{\nu_0}{c} \sqrt{\frac{2kT}{m}}$$



Schemes of Resonance Laser Ionization in Supersonic Beams

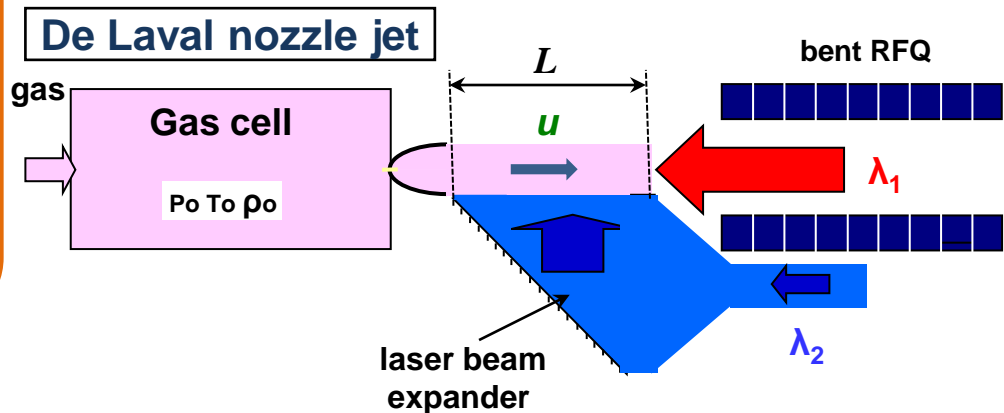
NO laser ionization inside the cell !

Laser ionization only in the cold jet !

λ_2 !

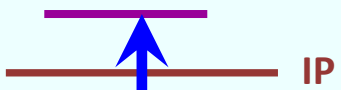
$f_{laser} \geq 1/(L/u) \geq 10 \text{ kHz}$, argon jet - $L = 5.5 \text{ cm}$
 u - stream velocity, 550m/s

Crossed laser beams with supersonic jet



The parallel beam from de Laval nozzle !
No broadening due to the beam divergence
Very careful design of the nozzle is required

Autoionizing state



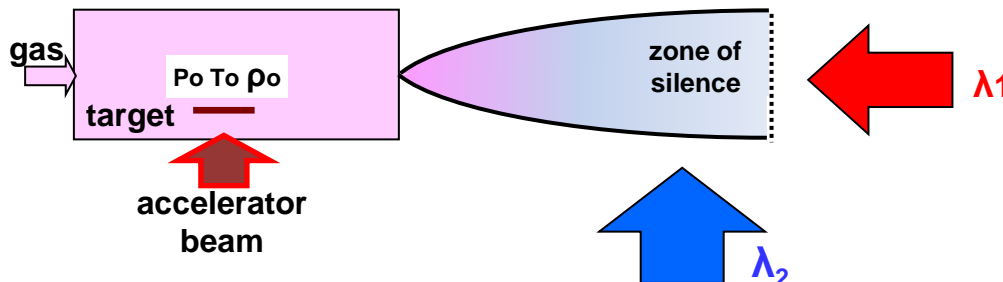
λ_2 $v_2 = v_{02}$



λ_1 $v_1 = v_{01} \times (1 - u/c)$

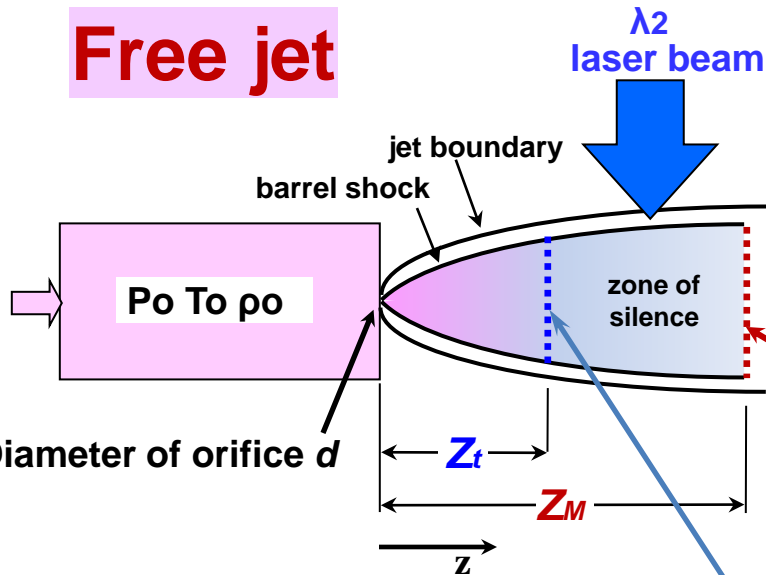
Ground state

Free jet



Two-Step Laser Ionization in a Free Jet

Free jet



1951 free jet – A. Kantrowitz, J. Grey

λ_1 laser beam

Mach disk, $T, \rho \uparrow$

Z_M – position of the Mach disk

$$\frac{Z_M}{d} = 0.67 \sqrt{\frac{P_o}{P_{bg}}}$$

M_t – terminal Mach number

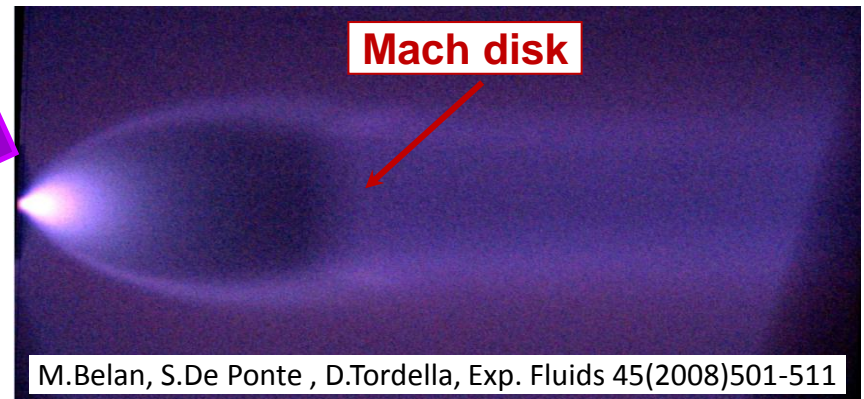
Z_t – position of terminal Mach number

$$M_t = 3.32 (P_o d)^{0.4}$$

(mbar, mm)

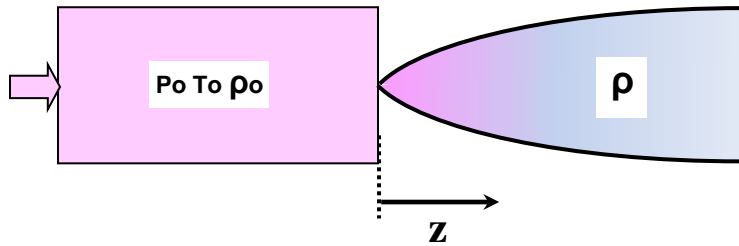
$$\frac{Z_t}{d} = \left(\frac{M_t}{3.26} \right)^{1.5}$$

Visualization of free jet

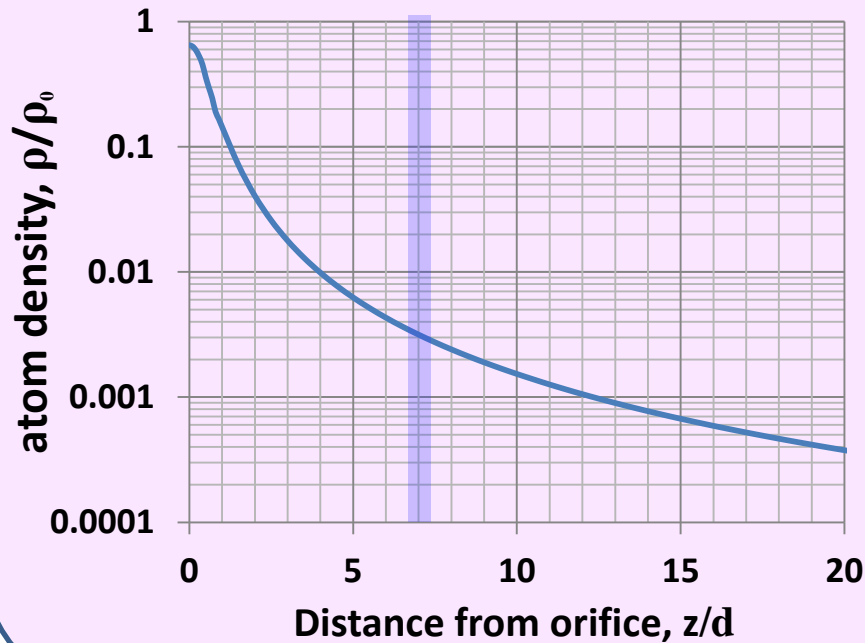


M. Belan, S. De Ponte, D. Tordella, Exp. Fluids 45(2008)501-511

Properties of Free Jet



$$\Gamma_{coll} = \gamma_{coll} \times \rho \rightarrow 3.3 \text{ MHz Mach}=12$$

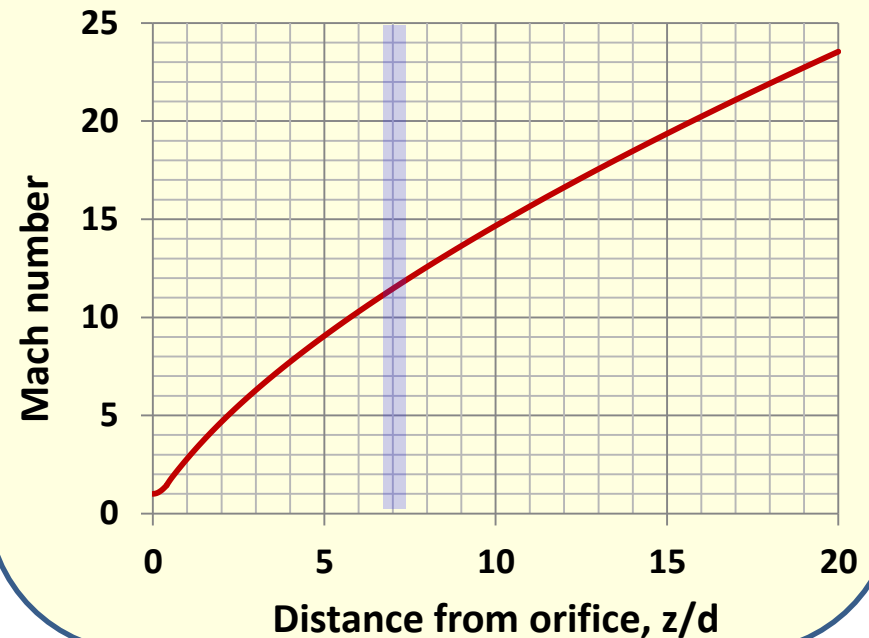


Centerline Mach number calculation

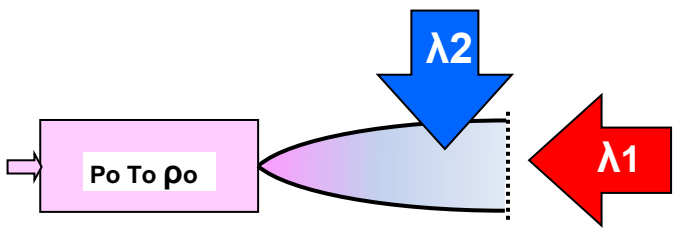
A	B	C_1	C_2	C_3	C_4
3.337	-1.541	3.232	-0.7563	0.3937	-0.0729

$$0 < \frac{Z}{d} < 1.0 \quad M = 1.0 + A \left(\frac{Z}{d} \right)^{-2} + B \left(\frac{Z}{d} \right)^3$$

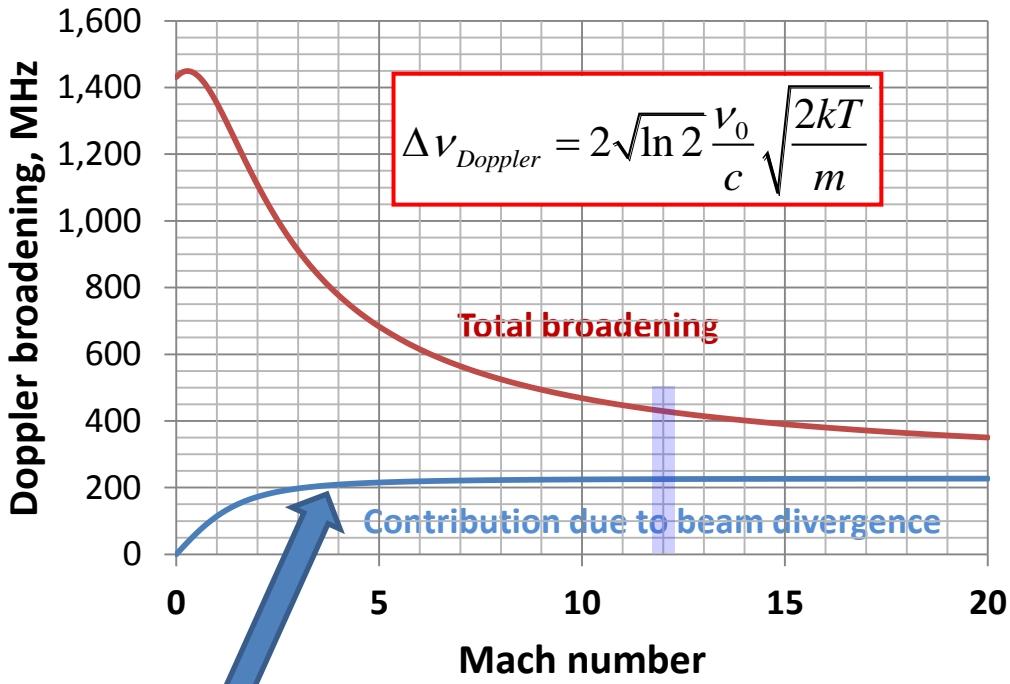
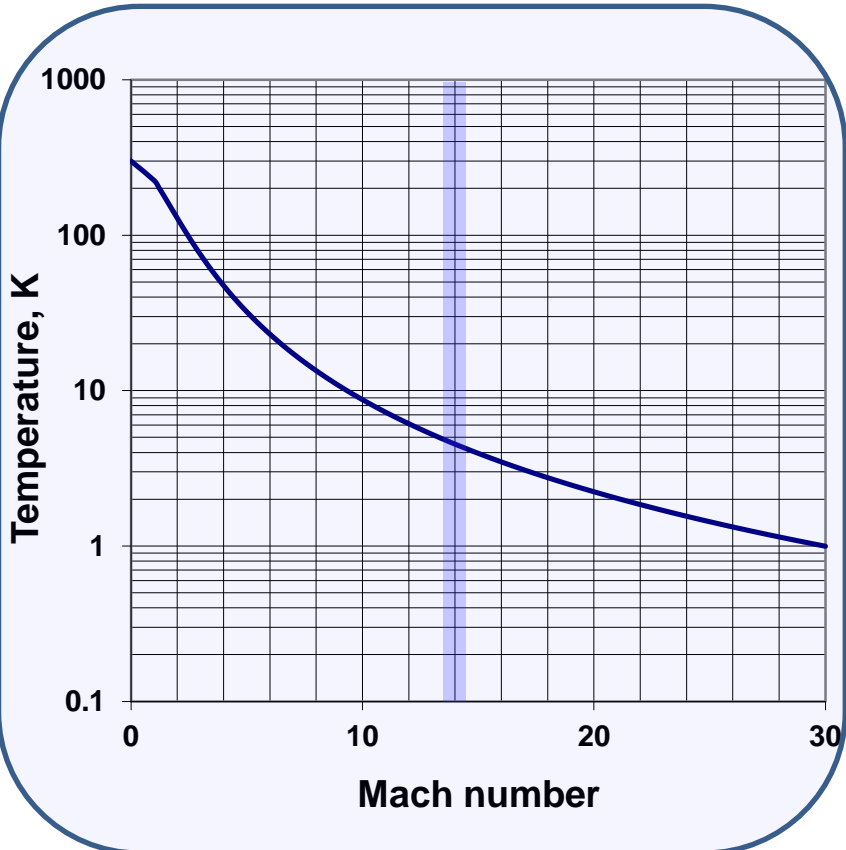
$$\frac{Z}{d} > 0.5 \quad M = \left(\frac{Z}{d} \right)^{(\gamma-1)} \left[C_1 + \frac{C_2}{\left(\frac{Z}{d} \right)} + \frac{C_3}{\left(\frac{Z}{d} \right)^2} + \frac{C_4}{\left(\frac{Z}{d} \right)^3} \right]$$



Doppler Broadening in the Free Jet Supersonic Beam



$4s^2S_{1/2} - 4p^2P_{1/2}$, 327.4 nm ^{63}Cu transition, $v_0 = 30535.3 \text{ cm}^{-1}$



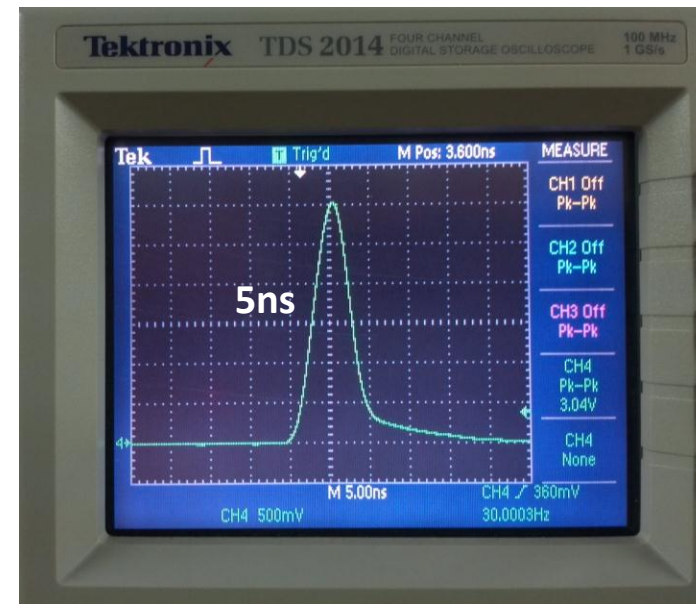
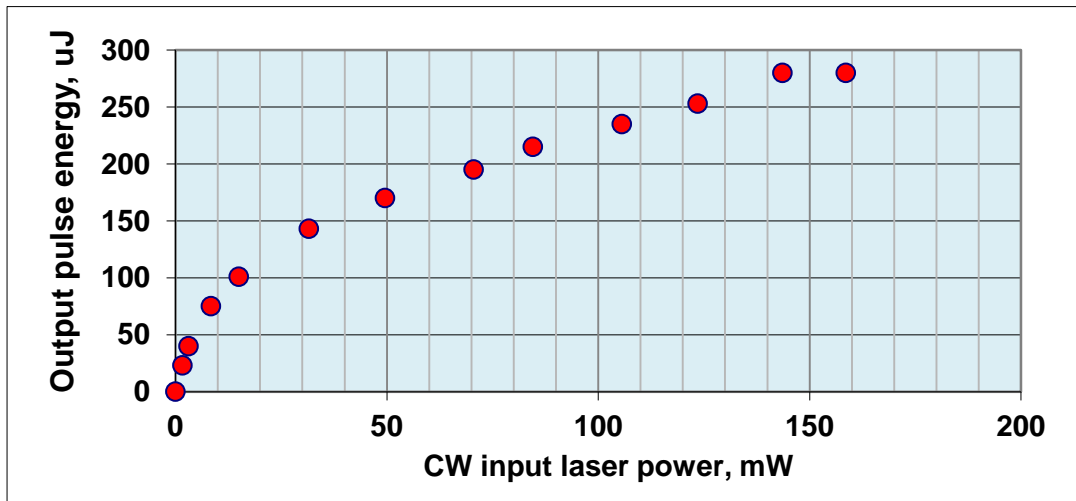
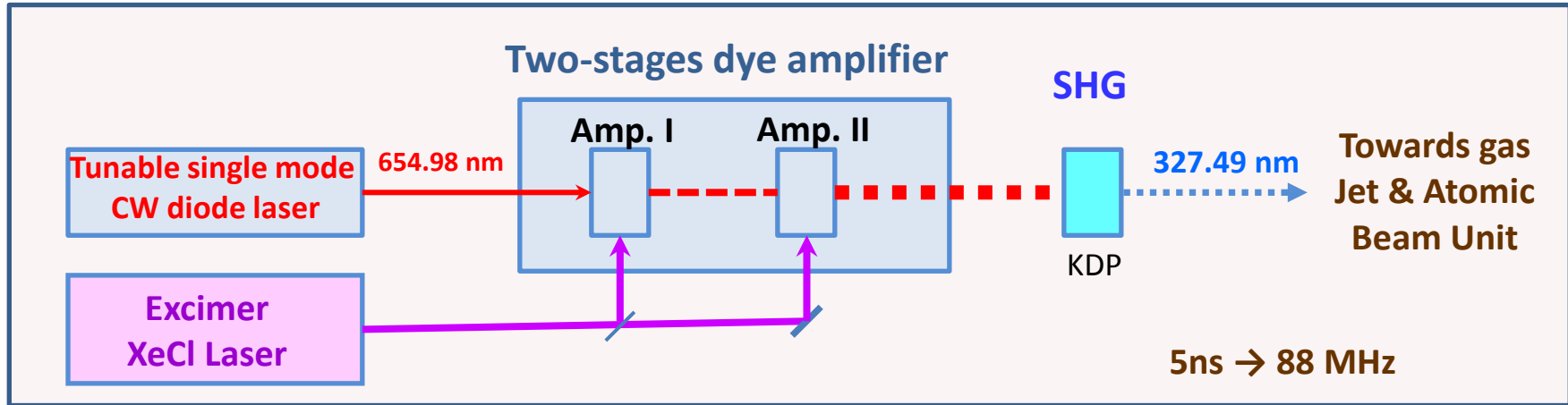
$$\Delta v_{Doppler} = 2\sqrt{\ln 2} \frac{v_0}{c} \sqrt{\frac{2kT}{m}}$$

$$\Delta_{Doppler}^{ax} = v_0 \cdot u (1 - \cos \theta) / c - \text{axial laser beam direction}$$

T=4K, Dopp. W.=200 MHz

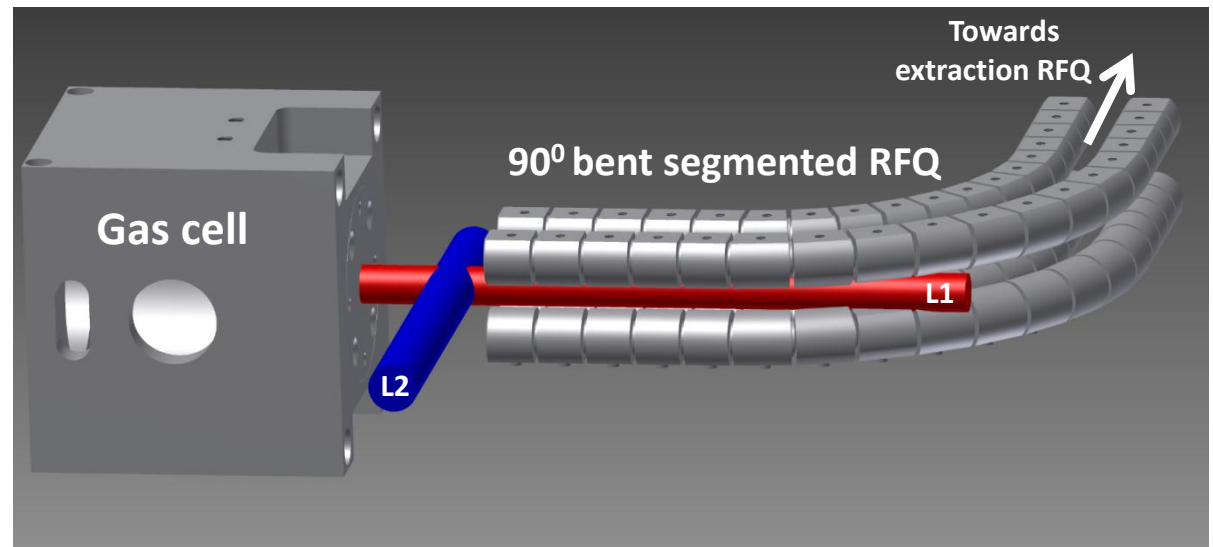
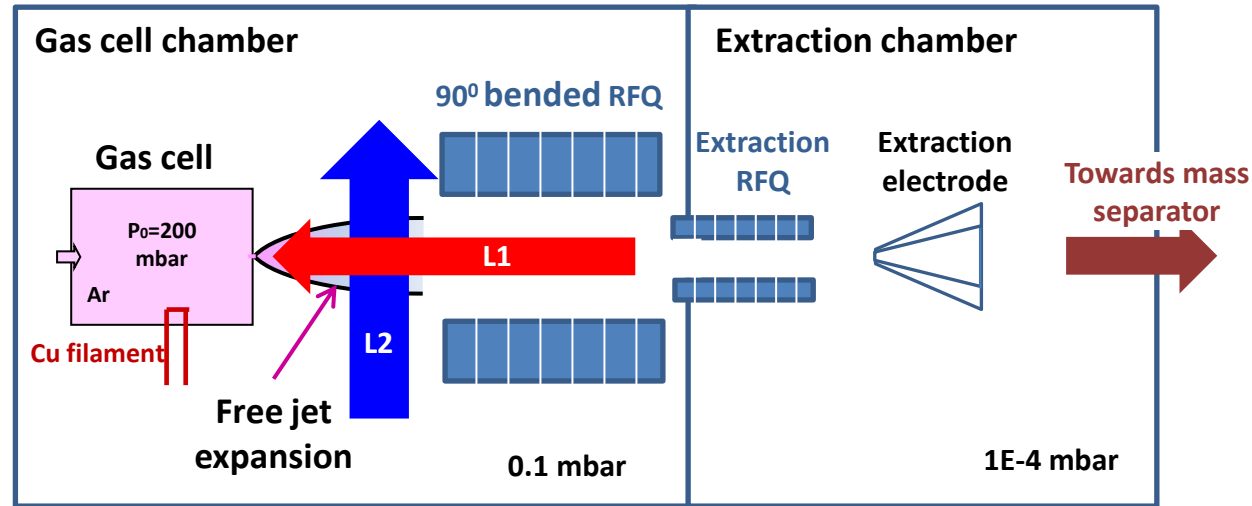
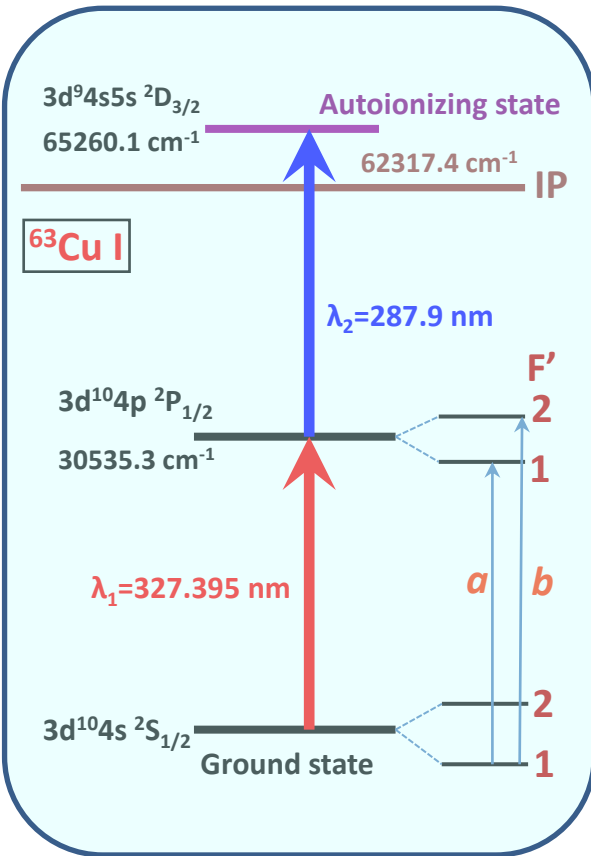
Tot. broad. = 420 MHz

Amplification of CW Single Mode Diode Laser Radiation in a Pulsed Dye Amplifier



Resonance Ionization Spectroscopy in a Free Gas Jet

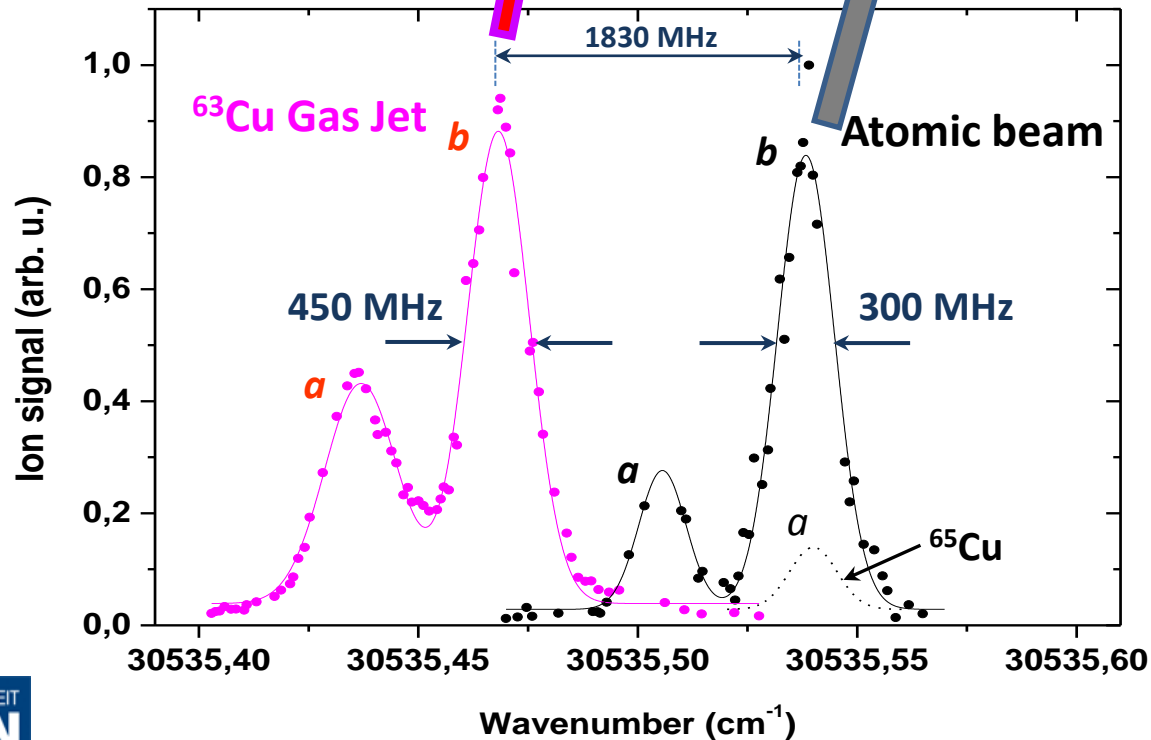
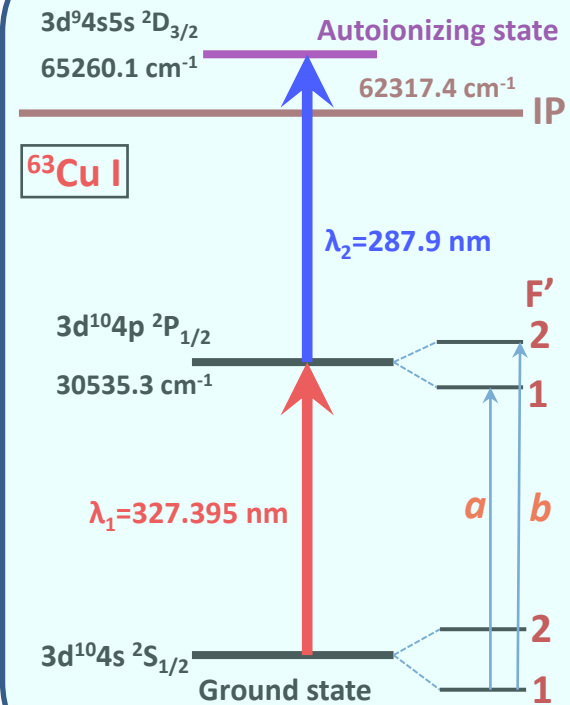
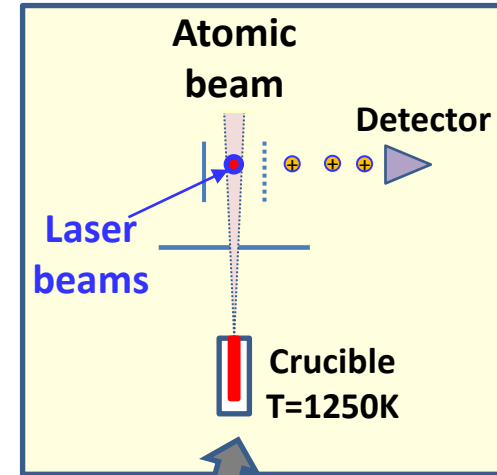
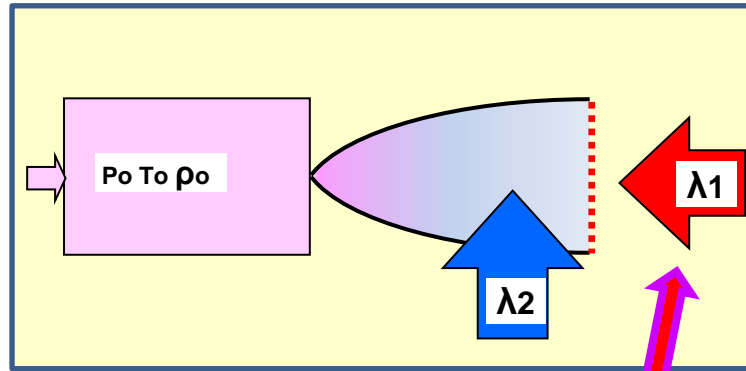
Yu. Kudryavtsev et al, <http://arxiv.org/abs/1211.6649>



Resonance Ionization Spectroscopy in a Free Gas Jet

$$u = \sqrt{\frac{kT_0 \gamma M^2}{m_{ng} \left\{ 1 + \left[\frac{\gamma - 1}{2} \right] M^2 \right\}}}$$

1830 MHz \rightarrow $T_0 = 355 \pm 3$ K

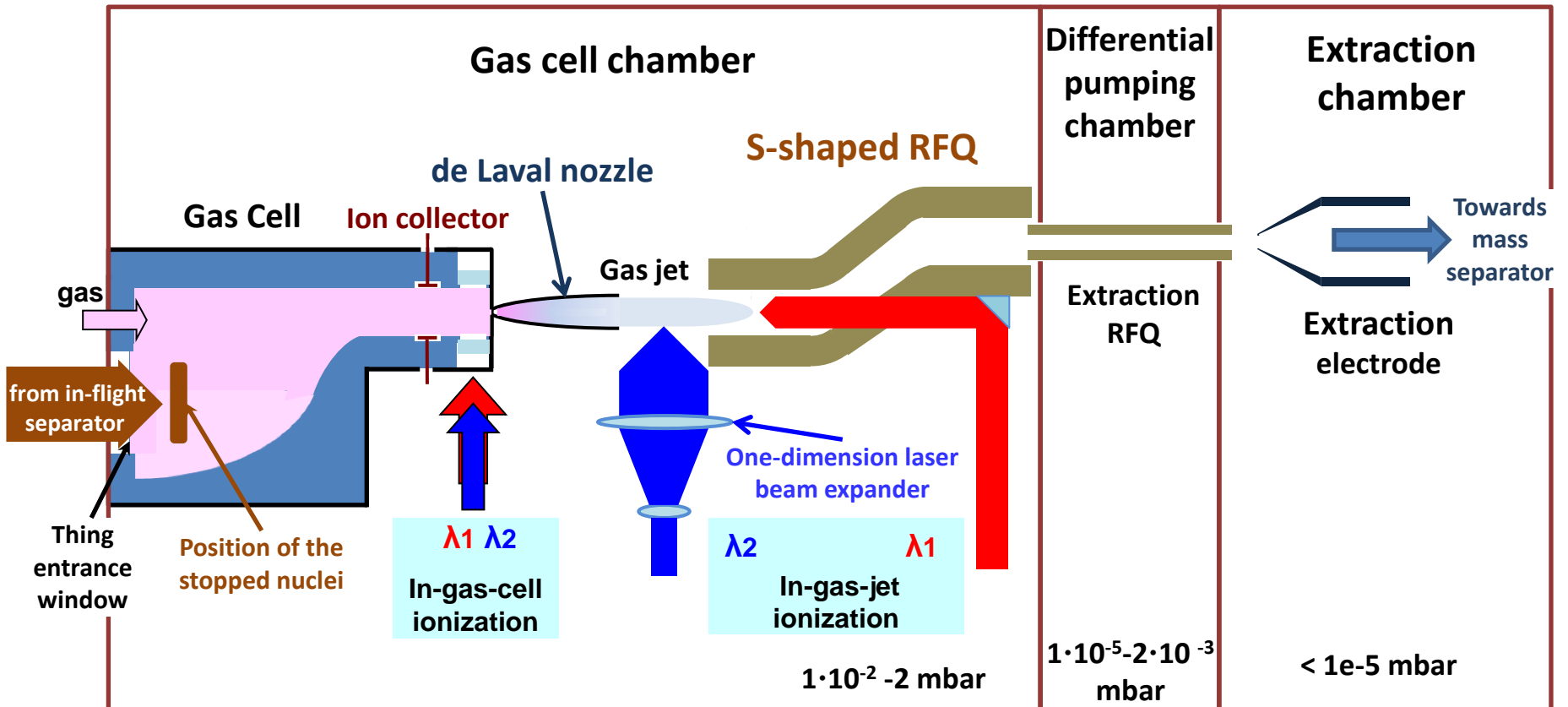


In-gas-cell and in-gas-jet laser RIS setup for HELIOS and S³ projects

erc grant has been granted, HELIOS

New laser laboratory will be set up at KU Leuven
The tender of the laser equipment has been finished

RILIS at S³ GANIL
poster #38 by Rafael Ferrer et al.



Laser equipment for IGLIS experiments @ HELIOS & S³

Two step laser ionization spectroscopy **in the gas cell**

- Two high-repetition-high-power Nd:YAG **pump Laser** →
- Max. average power: 90 W (@ 532 nm) or 36 W (@ 355 nm)
- Max. repetition rate: 15 kHz

- Two high repetition rate **dye lasers** →
- Tunable wavelength from 215 to 900 nm
- Linewidth: 0.06 cm⁻¹ (1.8 GHz) – 0.25 cm⁻¹ (7.5 GHz)



Pump Laser



Dye Laser

For high resolution spectroscopy **in the gas jet** first step will consist of

- A continuous wave (CW) single mode tunable **diode laser** →
- Linewidth: 1 MHz
- mode-hop-free tuning range: 20-30 GHz
- A **dye amplifier** with second harmonic generator



Diode Laser

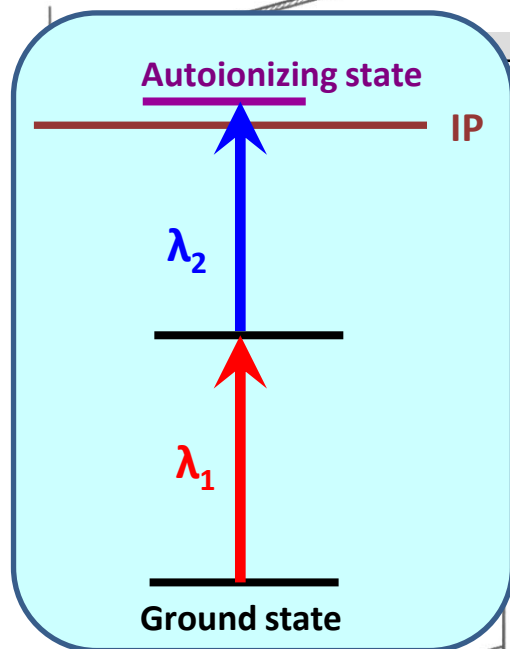
Summary

1. The crossed laser beams with supersonic jet has been proposed and realized off-line for two-step photo ionization in a free jet.
2. Using this method, the spectral resolution can be improved by one order of magnitude (200 MHz, $\Delta\nu/\nu = 2.3E-7$) in comparison to the gas cell.
3. The IGLIS technique that combines laser ionization in a gas cell and in a gas jet is adapted for production and spectroscopy of rare radioactive isotopes.

Thank you for your attention

Louvain-la-Neuve Radioactive Beam Facility

Two-step resonant laser ionization
→ element selectivity
Mass separation
→ isobar selectivity

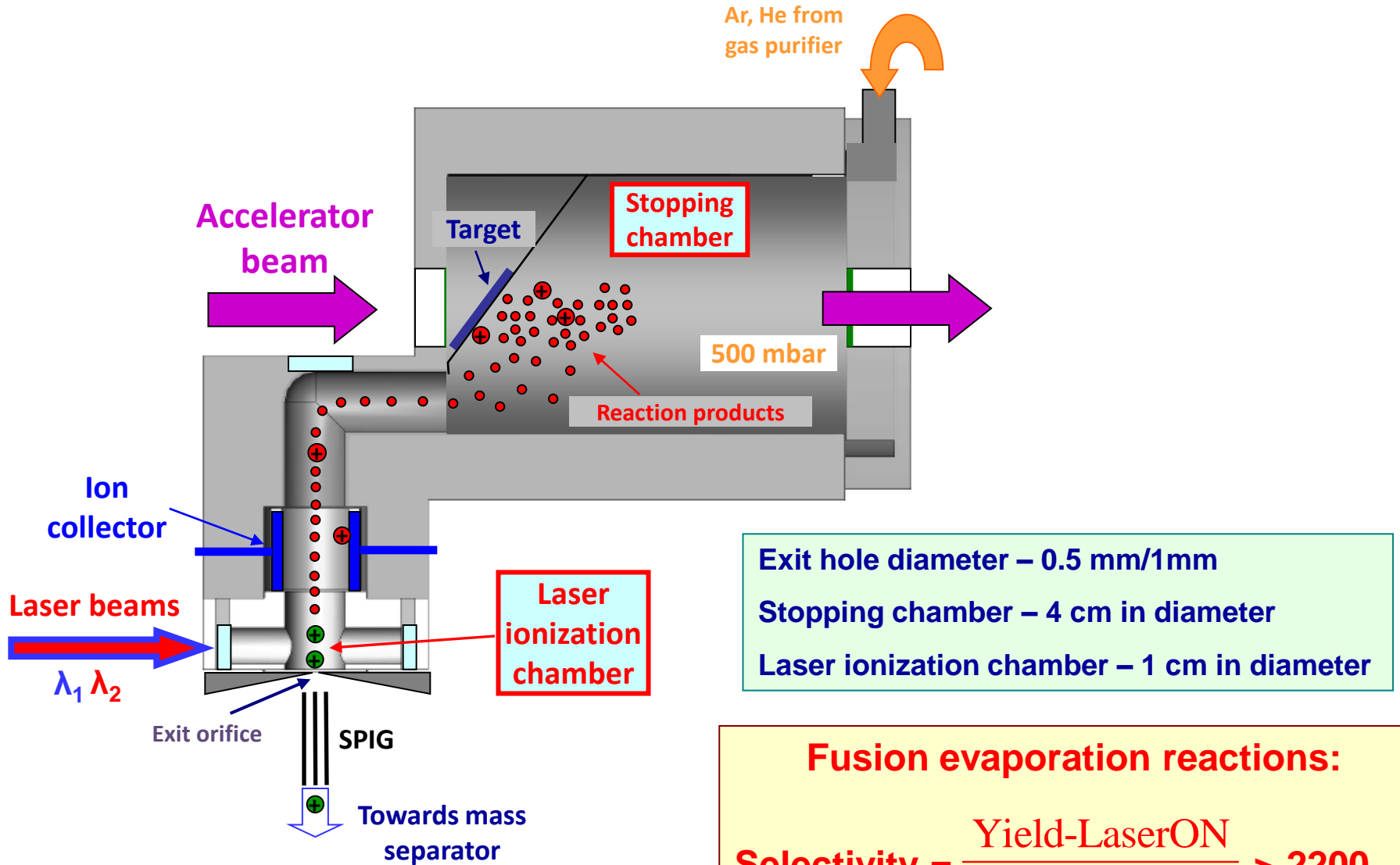


CYCLONE 110

LISOL

LASER ION SOURCE

Dual-Chamber Gas Cell Laser Ion Source



Doppler Gaussian and collision- and natural Lorentzian contributions to the spectral line shape

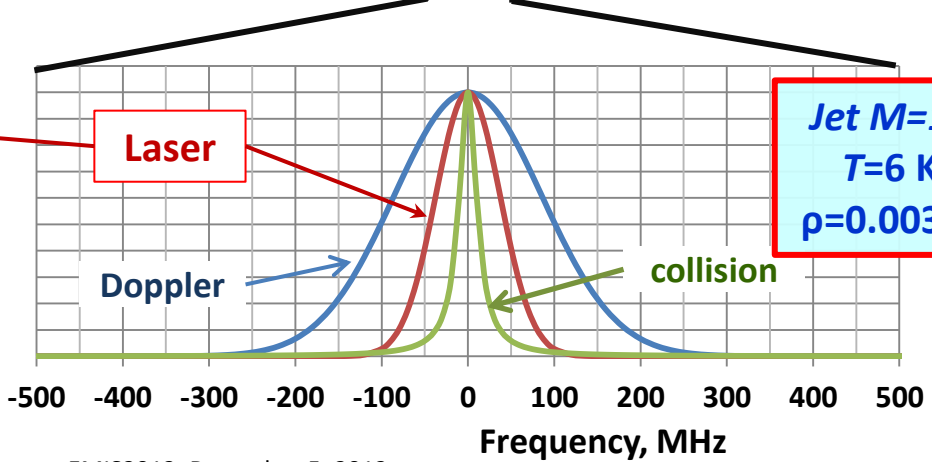
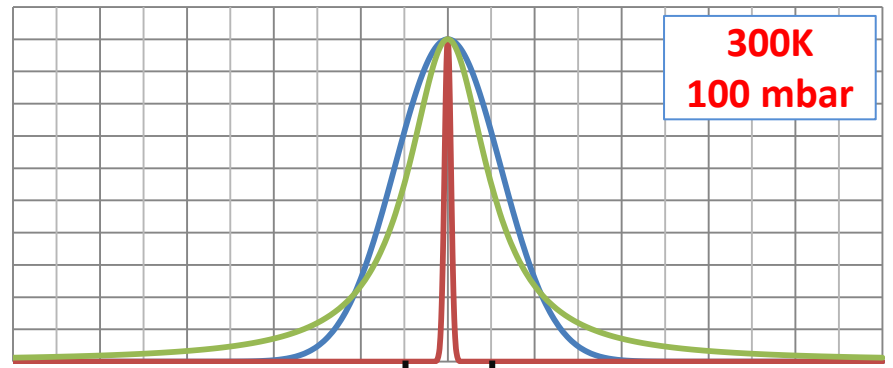
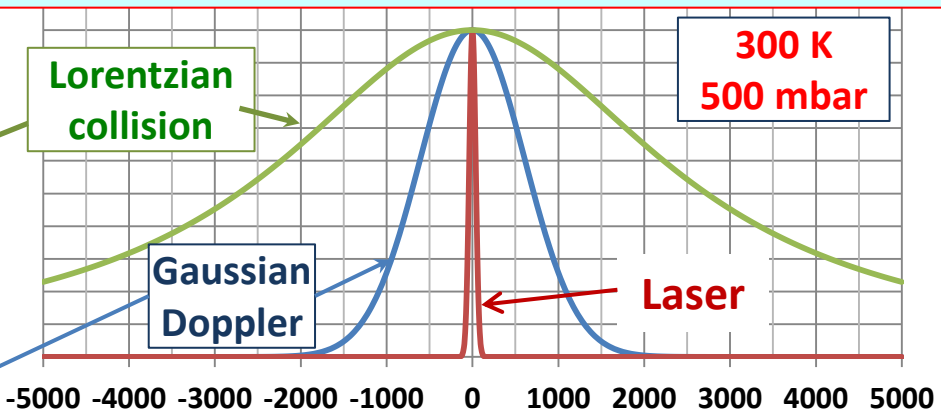
$$L(\nu - \nu_0) = \frac{1}{2\pi} \frac{\Gamma}{(\nu - \nu_0 + \Gamma sh)^2 + (\Gamma/2)^2}$$

$$G(\nu) = G_0 \exp \left[-\frac{c^2 (\nu - \nu_0)^2}{\nu_0^2 \frac{2kT}{m}} \right]$$

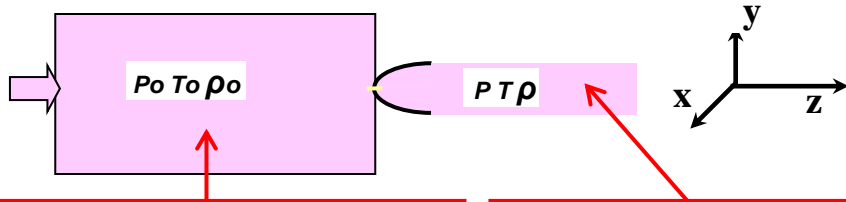
Laser bandwidth – δ_{laser} Gaussian if laser time profile is Gaussian

$$\delta_{laser} = 441 / \tau_{pulse} \quad \tau_{pulse} = 5 \text{ ns} \quad \delta_{laser} = 88 \text{ MHz}$$

Laser pulse length should short to provide interaction with all atoms!



Supersonic Beam from de Laval Nozzle



$$F^{th}(v_i) = \sqrt{\frac{m}{2\pi kT_0}} \exp\left(\frac{-mv_i^2}{2kT_0}\right)$$

$$F^{ss}(v_z) = \sqrt{\frac{m}{2\pi kT}} \exp\left(\frac{-m(v_z-u)^2}{2kT}\right)$$

Mach number - $M = u / a$

u - gas stream velocity

a - local speed of sound

T - gas temperature

m_{ng} - mass of the noble gas

$$a = \sqrt{\frac{\gamma kT}{m_{ng}}}$$

$$u = \sqrt{\frac{kT_0 \gamma M^2}{m_{ng} \{1 + [(\gamma - 1)/2] M^2\}}}$$

$\gamma = C_p / C_v$ - ratio of specific heat capacities = 5/3

