



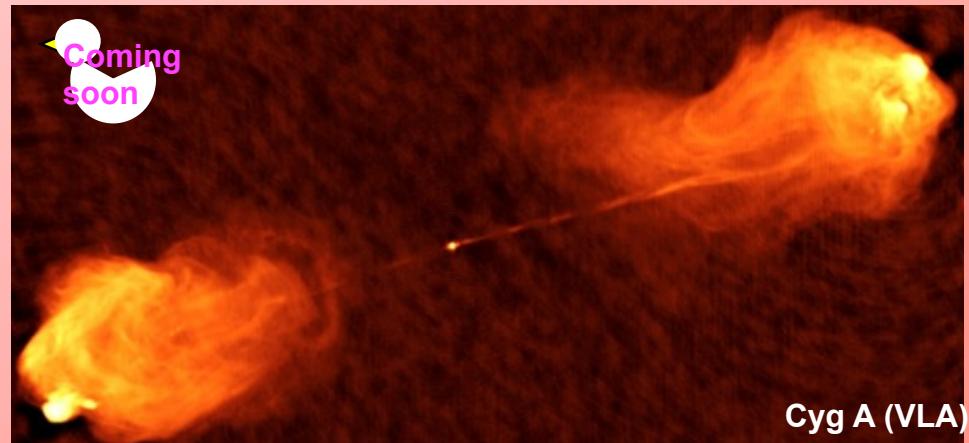
舌流磁場と 相対論的磁気流体風

田中周太 (青山学院大学)

with

當真賢二 (東北大学)

Crab (Chandra)

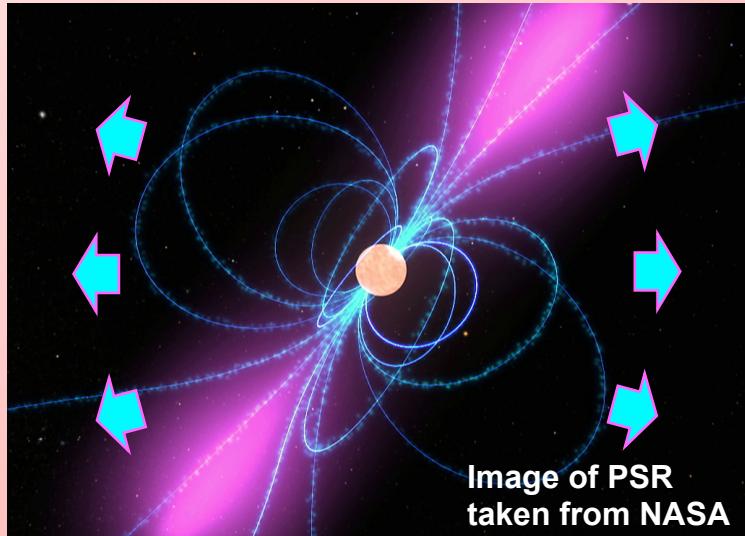
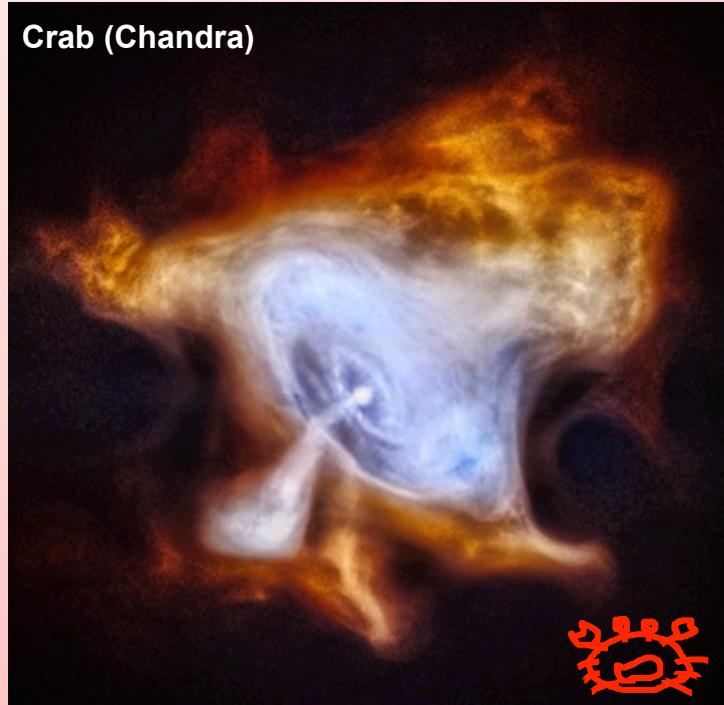


Cyg A (VLA)

Introduction

Pulsar Winds

- Powered by pulsar = NS.
- Relativistic plasma outflow.
- High energy (non-thermal) emission.
- Non-spherical (jet-torus structure).



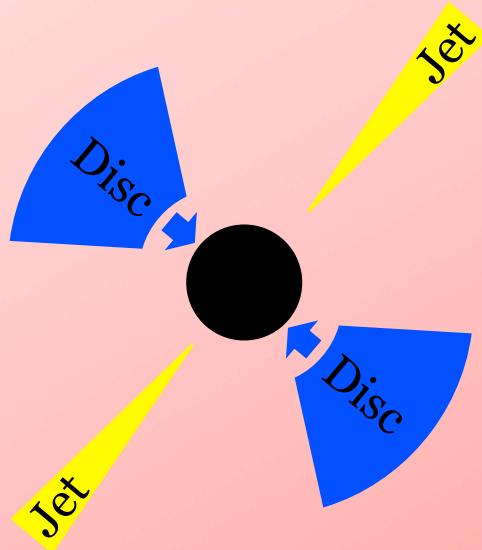
- Rotation power of pulsar \leftarrow period decrease
- $\Gamma_{\text{wind}} \gg 1 \leftarrow$ relativistically hot PWN emission
- e^\pm plasma \leftarrow high emission efficiency of PWN
- Magnetized plasma \leftarrow synchrotron emission

Relativistic Jet

- Powered by NS or BH.
- Relativistic plasma outflow.
- High energy (non-thermal) emission.
- Bipolar jets from engine.



Common astrophysical phenomena phenomena in AGN, microquasar, & GRB



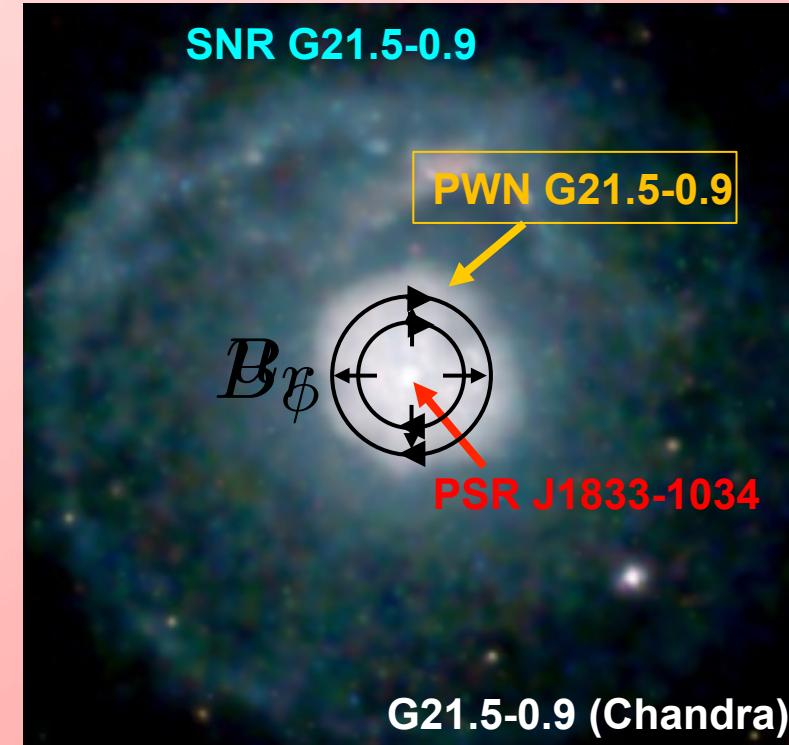
- Gravitational energy (inflow \rightarrow outflow)
Same as pulsar wind
- Rotation powered?
- Role of magnetic field?
- How to accelerates to relativistic flow?
- How to collimate jets?

Motivation

KC model

Kennel & Coroniti (1984a, b) (c.f., Rees & Gunn 1974)
⇒ 1D static MHD model of PWN

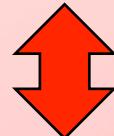
- Cold magnetized e^\pm wind from the central pulsar.
- Standing strong termination shock.
- A PWN is a post-shock pulsar wind.
- Steady & spherical system.
- Nebular flow is pure radial flow.
- Nebular B-field is pure toroidal.
- A PWN is an ideal MHD flow.
- Broadband emission by shock-accelerated non-thermal particles.



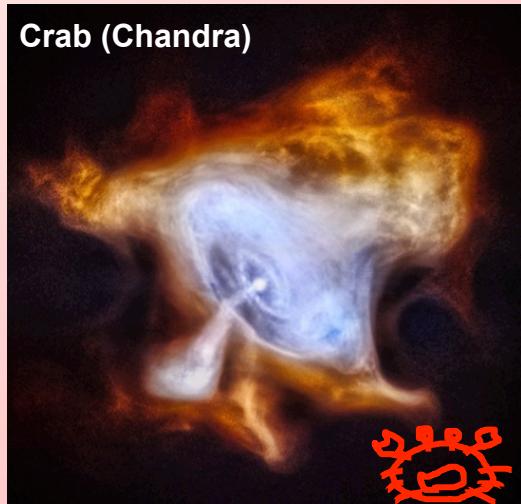
System dynamics is controlled by magnetization of pulsar wind:
 $\sigma = \text{Poynting} / \text{particle energy flux}$

Beyond KC model to resolve σ -problem

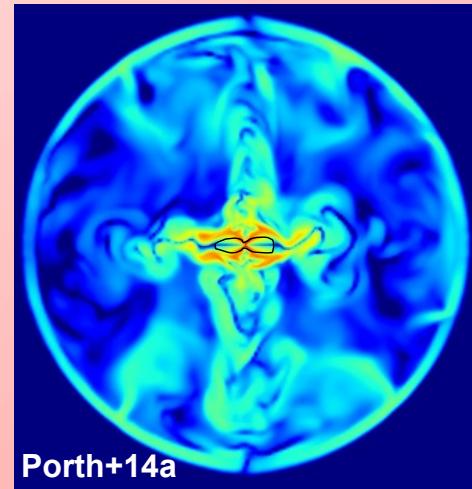
**Efforts to $\sigma \ll 1$ in wind region
have not been accomplished**



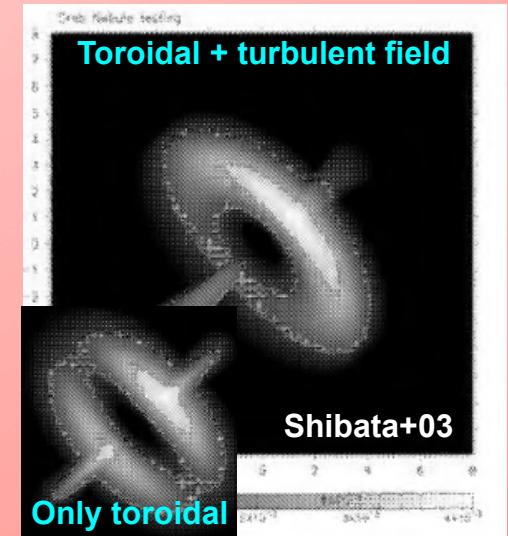
1D?



KC model = 1D ideal MHD with B_φ
Ideal MHD?



Toroidal B-field?



**3D effects (turbulence) &
B-field dissipation (non-ideal MHD)**

Basic Equations

Extend KC model by adding phenomenological terms which represent **radiation loss**, **turbulent magnetic field**, & magnetic **dissipation**.

Eq. of continuity

$$\langle \nabla_\mu (n u^\mu) \rangle = 0, \quad \text{TTT18ApJ}$$

Conservation of total energy

$$\langle \nabla_\mu T^{\mu t} \rangle = -\gamma \frac{\Lambda_{\text{rad}}}{c},$$

Conservation of fluid internal energy

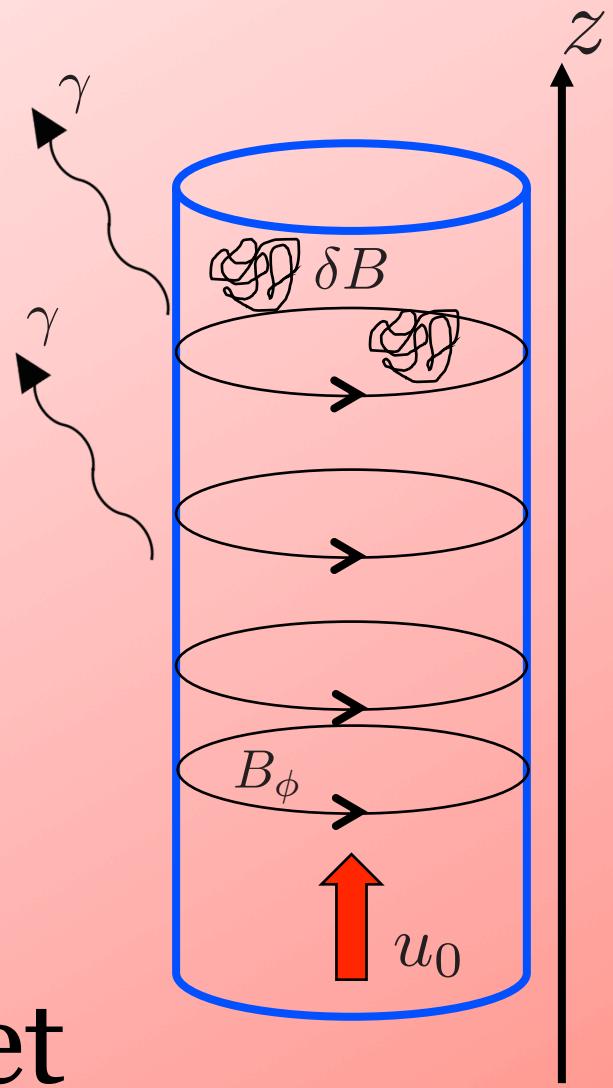
$$-\langle u_\nu \nabla_\mu T_{\text{FL}}^{\mu\nu} \rangle = \frac{\delta b^2/2}{\tau_{\text{diss}}} - \frac{\Lambda_{\text{rad}}}{c},$$

Induction equation for
(mean) toroidal field +
turbulent field

$$\frac{1}{2} \langle \bar{b}_\mu e^{\mu\nu\alpha\beta} \nabla_\nu F_{\alpha\beta} \rangle = -\frac{\bar{b}^2/2}{\tau_{\text{conv}}},$$

$$\frac{1}{2} \langle \delta b_\mu e^{\mu\nu\alpha\beta} \nabla_\nu F_{\alpha\beta} \rangle = \frac{\bar{b}^2/2}{\tau_{\text{conv}}} - \frac{\delta b^2/2}{\tau_{\text{diss}}}.$$

Steady & spherical symmetry



A Model of Cylindrical Jet

Basic Equations

TTT18の乱流磁場入りの相対論的MHD modelを採用。

radiation loss, turbulent magnetic field, & magnetic dissipation.

Eq. of continuity

$$\langle \nabla_\mu (n u^\mu) \rangle = 0,$$

Conservation of total energy

$$\langle \nabla_\mu T^{\mu t} \rangle = -\gamma \frac{\Lambda_{\text{rad}}}{c},$$

Conservation of
fluid internal energy

$$-\langle u_\nu \nabla_\mu T_{\text{FL}}^{\mu\nu} \rangle = \frac{\delta b^2/2}{\tau_{\text{diss}}} - \frac{\Lambda_{\text{rad}}}{c},$$

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$$\frac{1}{2} \langle \delta b_\mu e^{\mu\nu\alpha\beta} \nabla_\nu F_{\alpha\beta} \rangle = \frac{\bar{b}^2/2}{\tau_{\text{conv}}} - \frac{\delta b^2/2}{\tau_{\text{diss}}}.$$

定常な円柱状ジェットを考える

Cylinder Jet再考

一次元定常流は加速しない \Leftrightarrow 一次元定常流は膨張しない

$$(\beta^2 - \beta_c^2) \frac{du}{dx} = \frac{8u}{3r} \frac{p}{\epsilon}, \quad \epsilon = w + \bar{b}^2$$

TTT18ではCrab Nebula内での乱流磁場の生成による減速を考えた

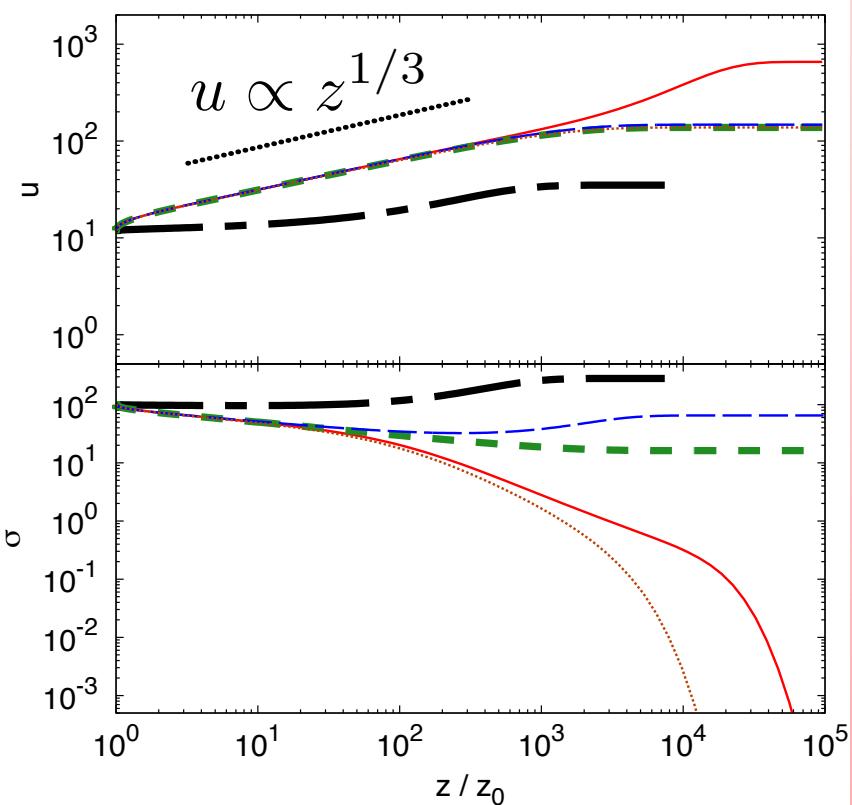
$$(\beta^2 - \beta_c^2) \frac{du}{dx} = \frac{8u}{3r} \frac{p}{\epsilon} + \frac{\Lambda_{\text{rad}}}{3c} + \frac{\bar{b}^2}{3c\tau_{\text{conv}}}$$

希薄化加速とは別のメカニズム

Results u & σ

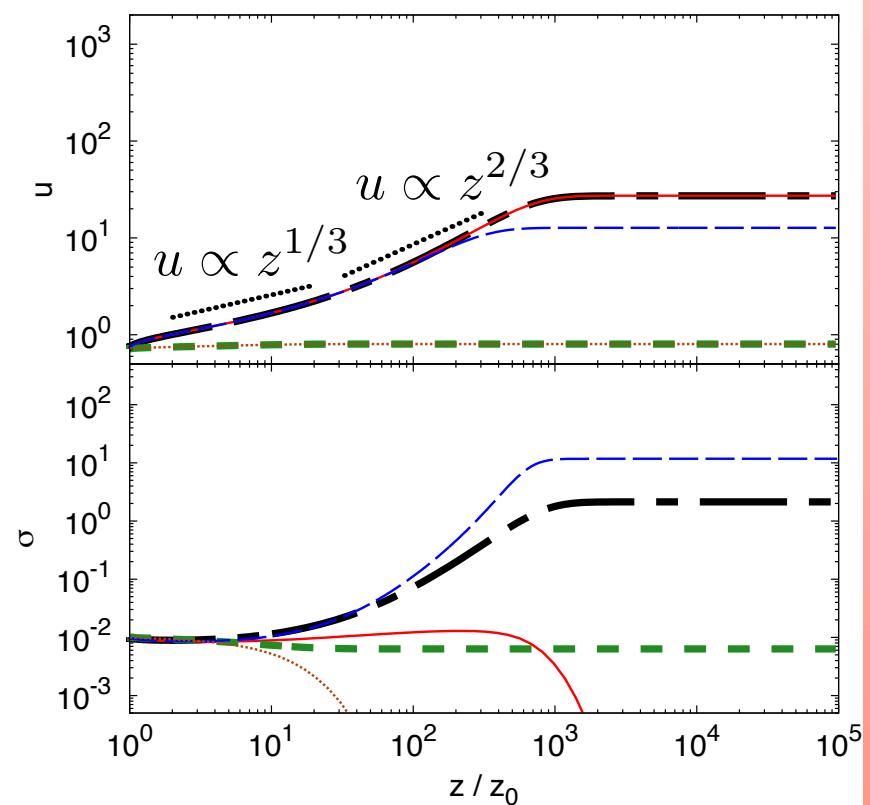
固有系で $z_{\text{diss}}, z_{\text{conv}}$ が一定

High- σ flow: $\sigma_0 = 10^2$



	黒	赤	緑	青	茶
\tilde{z}_{conv}	10^{10}	10	10	10	10
\tilde{z}_{diss}	10^{10}	10	10^{10}	10^{10}	10
\tilde{z}_{cool}	10	10	10^{10}	10	10^{10}

Low- σ flow: $\sigma_0 = 10^{-2}$



$$(n_0, u_0, w_0, b_0, \delta b_0) = (\gamma_{\max} = 10^4, u_c + \delta u, 1.0, \sigma_0, 0.0)$$

$w_0 \left[\frac{L_0}{A_0 c u_0 \gamma_0 (1 + \sigma_0)} \right]$ L_0 (全光度), A_0 (断面積), z_0 (初期位置)で系が定まる。 12

Summary

一次元流の加速

- 乱流磁場の生成、散逸をモデル化した非MHD方程式系 c.f., STTT18
- 球対称系と同じ $u \propto z^{1/3}$ が得られる。 c.f., Drenkhan02
- 散逸+放射冷却を加えることでより加速する。
- 磁場変換、磁場散逸、冷却はそれぞれ違う加速の仕方を示す。
- Total energyが同じ場合、high- σ_0 flowの方が(粒子が少ない分)速くなれる。

Further studies

- 安定なcylinder jetの τ_{conv} をどう評価するのか?
 - Spine-sheath構造でも本当に安定?