

Cumulants of conserved charges and total transport currents in GHD: direct summation of matrix elements?

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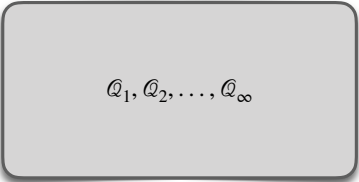
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GHD and GGE

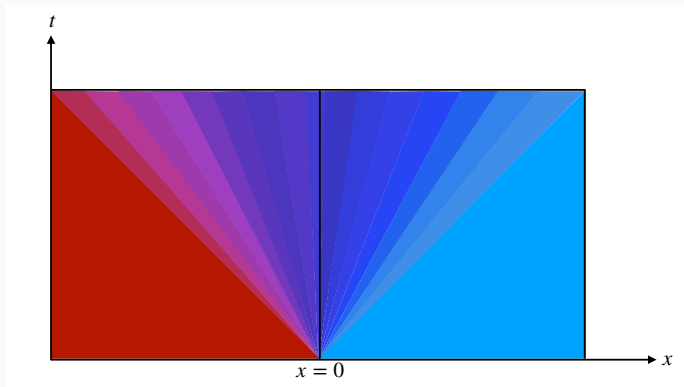
Isolated, out-of-equilibrium system


$$Q_1, Q_2, \dots, Q_\infty$$

Isolated, out-of-equilibrium system



Isolated, out-of-equilibrium system



Local GGE at the fluid cell around x

$$\langle \mathcal{O}(x, t) \rangle = \frac{\text{Tr}[e^{\Sigma_j - \beta_j(x, t) \mathcal{Q}_j} \mathcal{O}]}{\text{Tr}[e^{\Sigma_j - \beta_j(x, t) \mathcal{Q}_j}]}$$

Hydrodynamic description

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The state at this cell

$$\{\beta_1, \beta_2, \dots, \beta_\infty\} \leftrightarrow \{\langle \mathcal{Q}_1 \rangle, \langle \mathcal{Q}_2 \rangle, \dots, \langle \mathcal{Q}_\infty \rangle\}$$

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Equation of state that describes transport among neighboring cell

$$\partial_t \langle \mathcal{Q}_j(x, t) \rangle + \partial_x \langle J_j(x, t) \rangle = 0.$$

Non-equilibrium steady state

The Fermi-Dirac factor of the state at the junction

$$f(\theta) = f^L(\theta)\Theta(\theta - \theta^*) + f^R(\theta)\Theta(\theta^* - \theta), \quad v^{\text{eff}}(\theta^*) = 0$$

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Total transport crossing this point

$$J^{(t)} \equiv \int_0^t J(0, s) ds$$

Cumulant generating function (full counting statistics or Legendre transform of the large-deviation rate function)

$$F(\lambda) \equiv \lim_{t \rightarrow \infty} \frac{1}{t} \langle e^{\lambda J^{(t)}} \rangle = \sum_{n=1}^{\infty} \frac{\lambda^n}{n!} c_n$$

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satisfies a functional relation (see Doyon's talk or 1812.02082)

Individual cumulants

Steady current average

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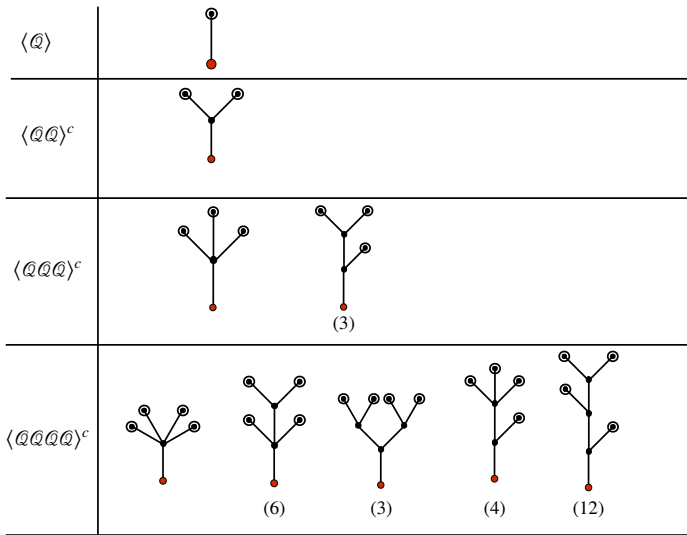
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Then ...

$$c_3 = \int \frac{d\theta}{2\pi} (E')^{\text{dr}}(\theta) f(\theta) [1 - f(\theta)] s(\theta) q^{\text{dr}}(\theta) \times \\ \times \{ [1 - 2f(\theta)] [q^{\text{dr}}(\theta)]^2 s(\theta) + 3[(q^{\text{dr}})^2 (1 - f) s]^{\text{dr}}(\theta) \},$$

$$c_4 = \int \frac{d\theta}{2\pi} (E')^{\text{dr}}(\theta) f(\theta) [1 - f(\theta)] \times \left\{ \frac{Y(\theta)^2 + 6Y(\theta) + 6}{[Y(\theta) + 1]^2} s(\theta) [q^{\text{dr}}(\theta)]^4 \right. \\ \left. + 3s(\theta) \{ [(1 - f) s (q^{\text{dr}})^2]^{\text{dr}}(\theta) \}^2 + 12s(\theta) q^{\text{dr}}(\theta) \{ (1 - f) s q^{\text{dr}} [(1 - f) s (q^{\text{dr}})^2]^{\text{dr}} \}^{\text{dr}}(\theta) \right. \\ \left. + 6[f(\theta) - 2] [q^{\text{dr}}(\theta)]^2 [s(1 - f) (q^{\text{dr}})^2]^{\text{dr}}(\theta) + 4s(\theta) q^{\text{dr}}(\theta) [(1 - f)(f - 2)(q^{\text{dr}})^3]^{\text{dr}}(\theta) \right\}$$

An unrelated problem, or is it?



Cumulants of conserved charges

Motivation

Traditional TBA

$$\text{Tr}[e^{-\sum_j \beta_j Q_j}]$$

This talk

$$\langle Q_1 Q_2 \dots Q_n \rangle$$

$$Q_j = \int dx Q_j(x, t)$$

Ultimate goal

$$\langle Q_1(x_1, t_1) Q_2(x_2, t_2) \dots Q_n(x_n, t_n) \rangle$$

Direct computation?

TBA free energy

$$F(\beta) = \log \text{Tr}[e^{\sum_j -\beta_j(x,t) Q_j}] = L \int \frac{dp(\theta)}{2\pi} \log[1 + e^{-\epsilon(\theta)}]$$

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Cumulants of conserved charges are given by its derivatives

$$\langle Q_j Q_k Q_l \dots \rangle = (-1)^\# \frac{\partial}{\partial \beta_j} \frac{\partial}{\partial \beta_k} \frac{\partial}{\partial \beta_l} \dots F(\beta)$$

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Charge covariance

$$\frac{1}{L} \langle Q_j Q_k \rangle^c = \int du \rho(u) [1 - f(u)] q_j^{\text{dr}}(\theta) q_k^{\text{dr}}(\theta)$$

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Higher cumulants: no rule to express in terms of particle distribution, Fermi-Dirac factor, simple dressing operations...

The free energy as a sum over trees [I.Kostov,D.Serban,D-L.V]

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- Bethe wavefunctions that diagonalize all the charges

$$Q_j|\theta_1, \dots, \theta_N\rangle = \sum_{i=1}^N q_j(\theta_i)|\theta_1, \dots, \theta_N\rangle$$

$$2\pi n_j = Lp(\theta_j) - i \sum_{k \neq j}^N \log S(\theta_j, \theta_k)$$

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- Summing over mode numbers

$$Z(\beta) = \sum_{N \geq 0} \sum_{n_1 < n_2 < \dots < n_N} e^{-w(n_1, n_2, \dots, n_N)}$$

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- Removing the constraint by $1 - \delta$ insertion

$$Z(\beta) = \sum_{N \geq 0} \frac{(-1)^N}{N!} \sum_{n_1, \dots, n_N \in \mathbb{Z}^N} \sum_{r_1, \dots, r_N \in \mathbb{N}^N} \prod_{j=1}^N \frac{(-1)^{r_j}}{r_j} e^{-w(n_1^{r_1}, \dots, n_N^{r_N})}.$$

- Transform the sums to integrals

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- Expand the Gaudin determinant to a sum over forests
- The free energy is a free sum over trees whose vertices (θ, r) live in $\mathbb{R} \times \mathbb{Z}^+$. Feynman rules

 (θ, r)	=	$\frac{(-1)^{r-1}}{r^2} e^{-rw(\theta)}$
 (θ, r)	=	$Lrp'(\theta) \frac{(-1)^{r-1}}{r^2} e^{-rw(\theta)}$
 $(\theta_1, r_1) \quad (\theta_2, r_2)$	=	$r_1 r_2 K(\theta_2, \theta_1)$

Summing over the trees

Let $Y_r(\theta)$ be the sum of all trees rooted at the point (θ, r)

$$Y_r(\theta) = \text{circle}(\theta, r) = \text{dot}(\theta, r) + \text{circle}(\theta, r) + \frac{1}{2!} \text{V-shape}(\theta, r) + \frac{1}{3!} \text{3-children}(\theta, r) + \dots$$

The diagrammatic equation shows the sum of all trees rooted at a point (θ, r) . On the left, $Y_r(\theta)$ is represented by a circle containing the point (θ, r) . This is equal to the sum of four terms: 1) a single black dot at (θ, r) ; 2) a black dot at (θ, r) with a single grey circle above it connected by a vertical line; 3) a black dot at (θ, r) with two grey circles above it, each connected by a diagonal line; 4) a black dot at (θ, r) with three grey circles above it, each connected by a diagonal line. The terms are separated by plus signs, and the third and fourth terms are multiplied by $\frac{1}{2!}$ and $\frac{1}{3!}$ respectively. The sequence ends with an ellipsis \dots .

Summing over the trees

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$$Y_r(\theta) = \begin{array}{c} \text{circle} \\ \text{with } (\theta, r) \end{array} = \begin{array}{c} \bullet \\ \text{with } (\theta, r) \end{array} + \begin{array}{c} \text{circle} \\ | \\ \bullet \\ \text{with } (\theta, r) \end{array} + \frac{1}{2!} \begin{array}{c} \text{circle} \quad \text{circle} \\ \diagdown \quad \diagup \\ \bullet \\ \text{with } (\theta, r) \end{array} + \frac{1}{3!} \begin{array}{c} \text{circle} \quad \text{circle} \quad \text{circle} \\ \diagdown \quad | \quad \diagup \\ \bullet \\ \text{with } (\theta, r) \end{array} + \dots$$

One can deduce from this structure that $Y_r(\theta) = (-1)^r Y_1^r(\theta)/r^2$ and

$$Y(\theta) = e^{-w(\theta)} \exp \int \frac{d\eta}{2\pi} K(\eta, \theta) \log[1 + Y(\eta)].$$

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The free energy is the sum over all trees

$$F(\beta) = L \int \frac{dp(\theta)}{2\pi} \sum_{r \geq 1} r Y_r(\theta) = L \int \frac{dp(\theta)}{2\pi} \log[1 + Y(\theta)].$$

$$\langle Q_j \rangle_\beta = \text{Tr}[e^{-\sum \beta_i Q_i} Q_j] / \text{Tr}[e^{-\sum \beta_i Q_i}]$$

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Insertion of identity

$$\langle \theta_N^{r_N}, \dots, \theta_1^{r_1} | Q_j | \theta_1^{r_1}, \dots, \theta_N^{r_N} \rangle = \sum_{i=1}^N r_i q_j(\theta_i) \langle \theta_N^{r_N}, \dots, \theta_1^{r_1} | \theta_1^{r_1}, \dots, \theta_N^{r_N} \rangle$$

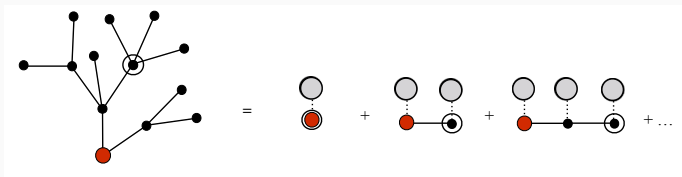
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leads to a sum over trees with a charge insertion



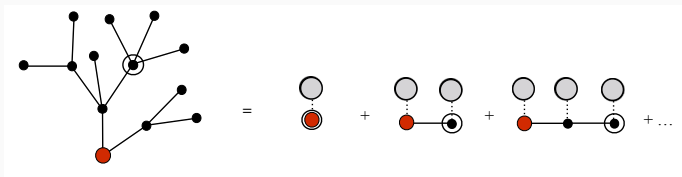
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In terms of the dressing operation

$$\frac{1}{L} \langle Q_j \rangle = \int \frac{dp}{2\pi} f(\theta) q_j^{\text{dr}}(\theta)$$

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Insertion of identity is factorized

$$\langle \theta_N^{r_N}, \dots, \theta_1^{r_1} | Q_j Q_k | \theta_1^{r_1}, \dots, \theta_N^{r_N} \rangle = \langle \theta_N^{r_N}, \dots, \theta_1^{r_1} | \theta_1^{r_1}, \dots, \theta_N^{r_N} \rangle \sum_{i=1}^N r_i q_j(\theta_i) \sum_{i=1}^N r_i q_k(\theta_i)$$

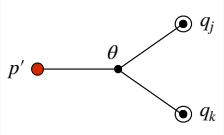
Charge covariance

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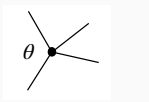
$$\langle \theta_N^{r_N}, \dots, \theta_1^{r_1} | Q_j Q_k | \theta_1^{r_1}, \dots, \theta_N^{r_N} \rangle = \langle \theta_N^{r_N}, \dots, \theta_1^{r_1} | \theta_1^{r_1}, \dots, \theta_N^{r_N} \rangle \sum_{i=1}^N r_i q_j(\theta_i) \sum_{i=1}^N r_i q_k(\theta_i)$$

The covariance is given by a sum over trees with two leaves

$$\frac{1}{L} \langle Q_j Q_k \rangle^c = \text{Diagram} = \int \frac{d\theta}{2\pi} f(\theta) [1 - f(\theta)] (p')^{\text{dr}}(\theta) q_j^{\text{dr}}(\theta) q_k^{\text{dr}}(\theta)$$


Higher cumulants

The n^{th} cumulant is given by the sum over all trees with a root that carries p' and n leaves carrying the n conserved charges

	=	$\psi^{\text{dr}}(\theta)$
	=	$K^{\text{dr}}(\theta, \eta)$
	=	$\sum_{r \geq 1} (-1)^{r-1} r^{d-1} e^{-r\epsilon(\theta)}$

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These trees are simple to generate

$$d_n = \sum_{\substack{p \in \mathcal{P}_n, |p| > 1 \\ p = (a_1^{\alpha_1}, \dots, a_j^{\alpha_j})}} \prod_{i=1}^j \binom{d_{a_i} + \alpha_i - 1}{\alpha_i}$$

n	1	2	3	4	5	6	7	8	9	10
d_n	1	1	2	5	12	33	90	261	766	2312

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	=	$\sum_{r \geq 1} (-1)^{r-1} r^{d-1} e^{-r\epsilon(\theta)}$

Physical interpretation: virtual particles with anomalous correction to the bare charges

$$\psi^{*\text{dr}}(\theta) \equiv \psi^{\text{dr}}(\theta) - \psi(\theta) = \int \frac{d\eta}{2\pi} K^{\text{dr}}(\eta, \theta) f(\eta) \psi(\eta)$$

Cumulants of total transport currents

Comparison with diagrams

Simple modifications of the diagrams:

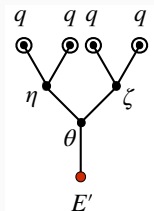
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- Sign of the effective velocity at odd internal vertices

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- Operator at the root: E' instead of p'
- Sign of the effective velocity at odd internal vertices

For instance



$$= 3 \int \frac{d\theta}{2\pi} \frac{d\eta}{2\pi} \frac{d\zeta}{2\pi} (E')^{\text{dr}}(\theta) s(\theta) f(\theta) [1 - f(\theta)] K^{\text{dr}}(\theta, \eta) K^{\text{dr}}(\theta, \zeta)$$

$$\times f(\eta) [1 - f(\eta)] s(\eta) f(\zeta) [1 - f(\zeta)] s(\zeta) [q^{\text{dr}}(\eta)]^2 [q^{\text{dr}}(\zeta)]^2$$

Highly non-trivial matching with the fourth cumulant!

Conclusion

- We develop a new approach to express cumulants of conserved charges in GGE and total currents in GHD in an intuitive, effective way
- Does this combinatorial structure of individual cumulants translate into an analytic property of the full counting statistics?
- Matrix element of current cumulants? Drude weight?
- "Free energy" that generates current cumulants? Dual TBA?
- Large scale dynamical correlation functions?

Thank you for attention!