



Spin-Related Effect in hadron production at BESIII

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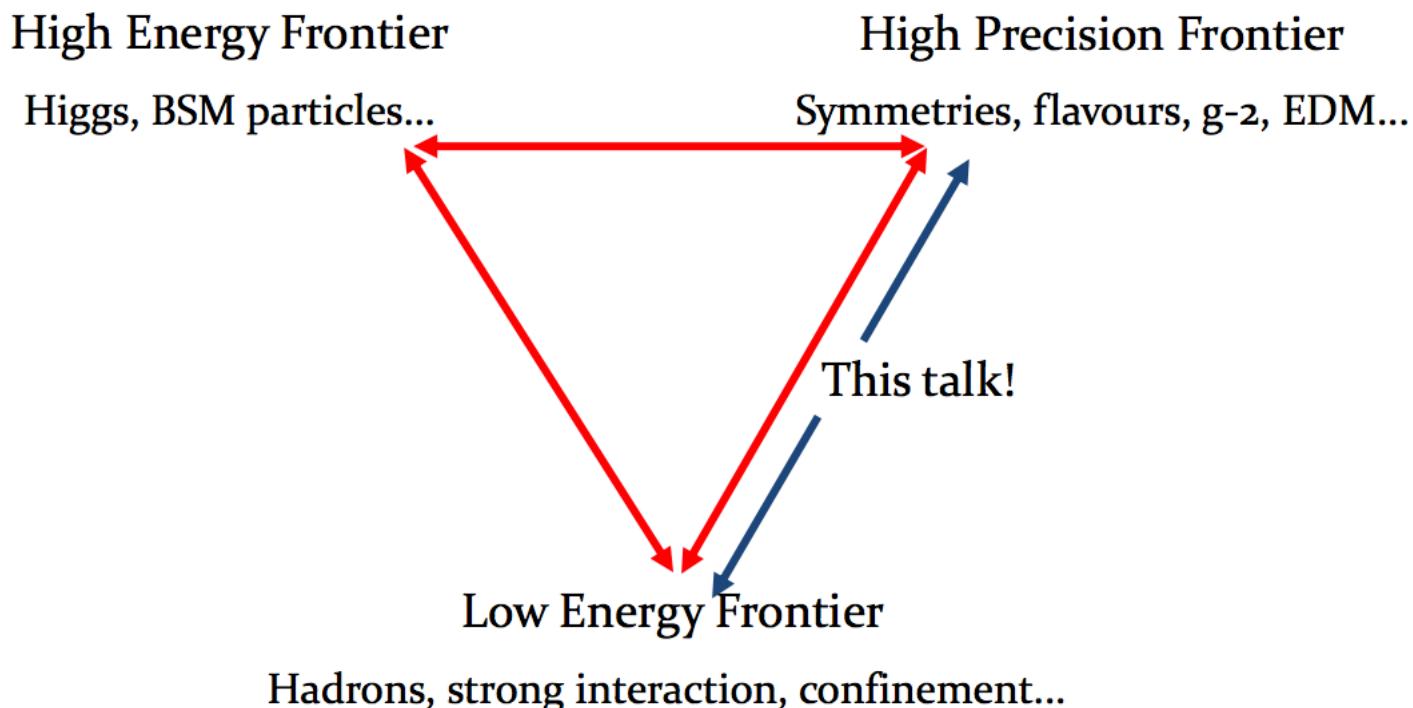
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(On Behalf of the BESIII collaboration)

Outline

- Introduction
- BEPCII/BESIII
- (Spin) Physics at BESIII
- Collins effect measurement
- $\Lambda_{(c)}$ polarization measurement
- Summary and Outlook



The Nucleon

Many challenges in modern physics manifest themselves in the nucleon:

- Challenging to describe from first principles

– Its abundance

High Precision

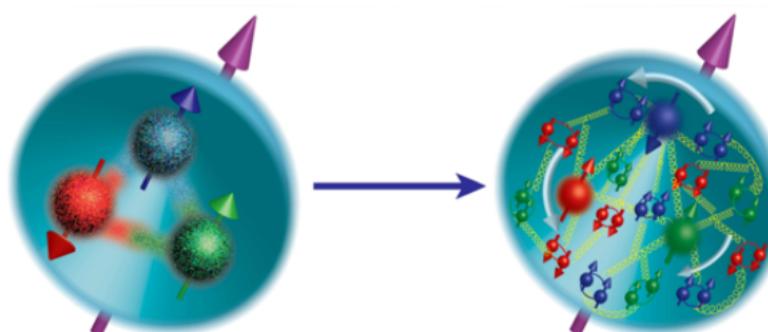
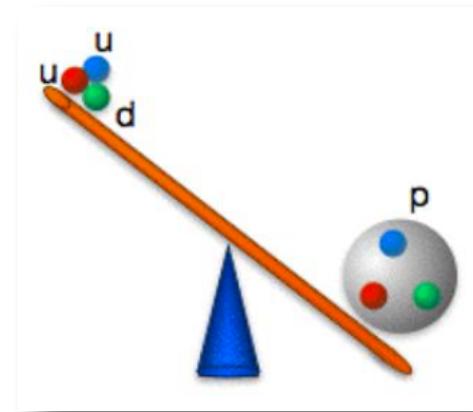
– Its mass

– Its spin

– Its structure

– Its radius ?

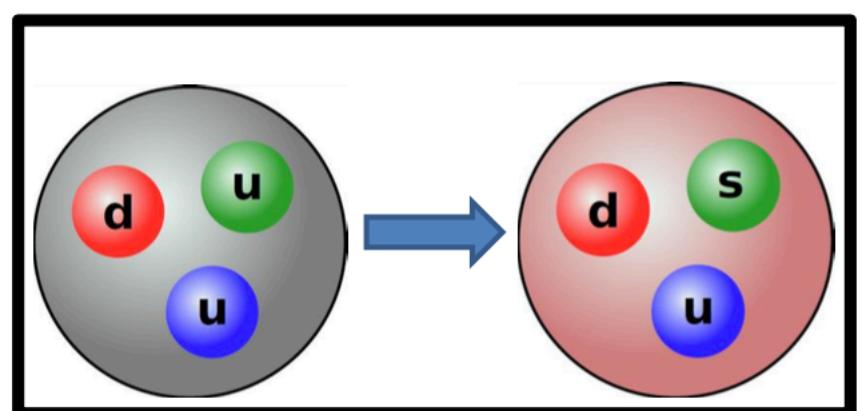
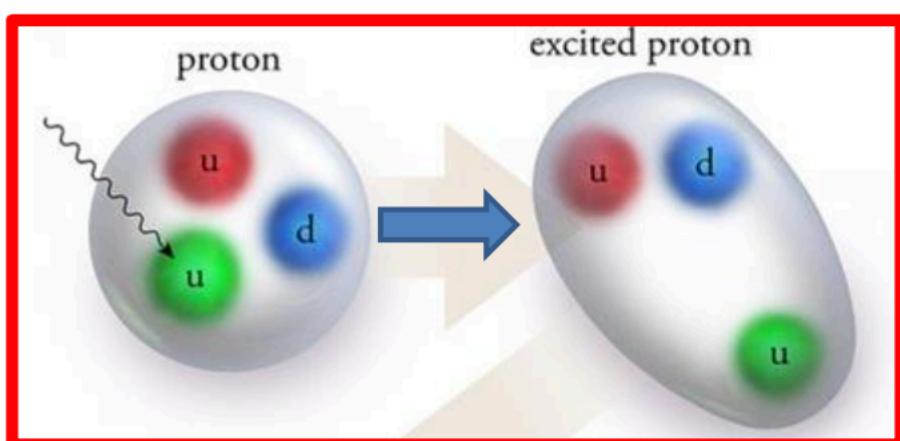
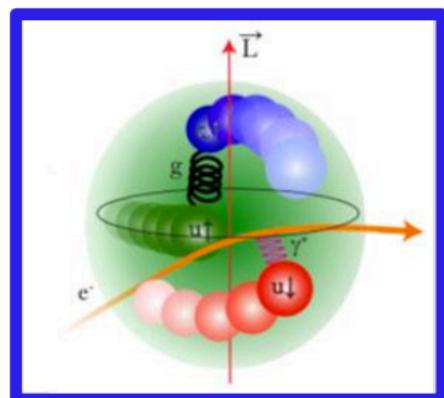
Low Energy /
Non-pQCD



Approaches

When you don't understand a system, you can*

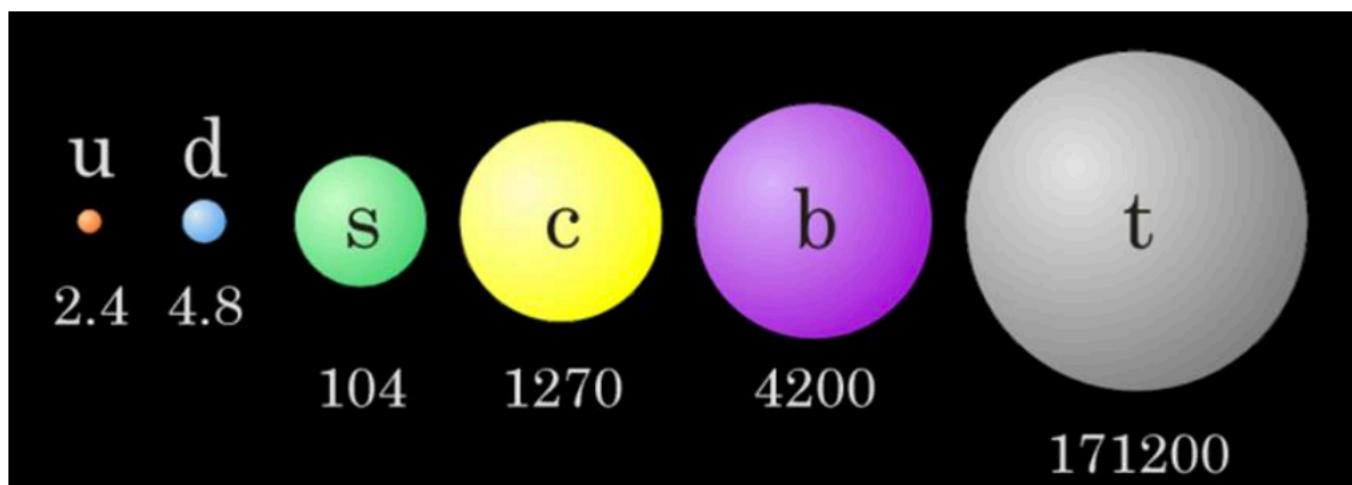
- Scatter on it
- Excite it
- Replace building blocks

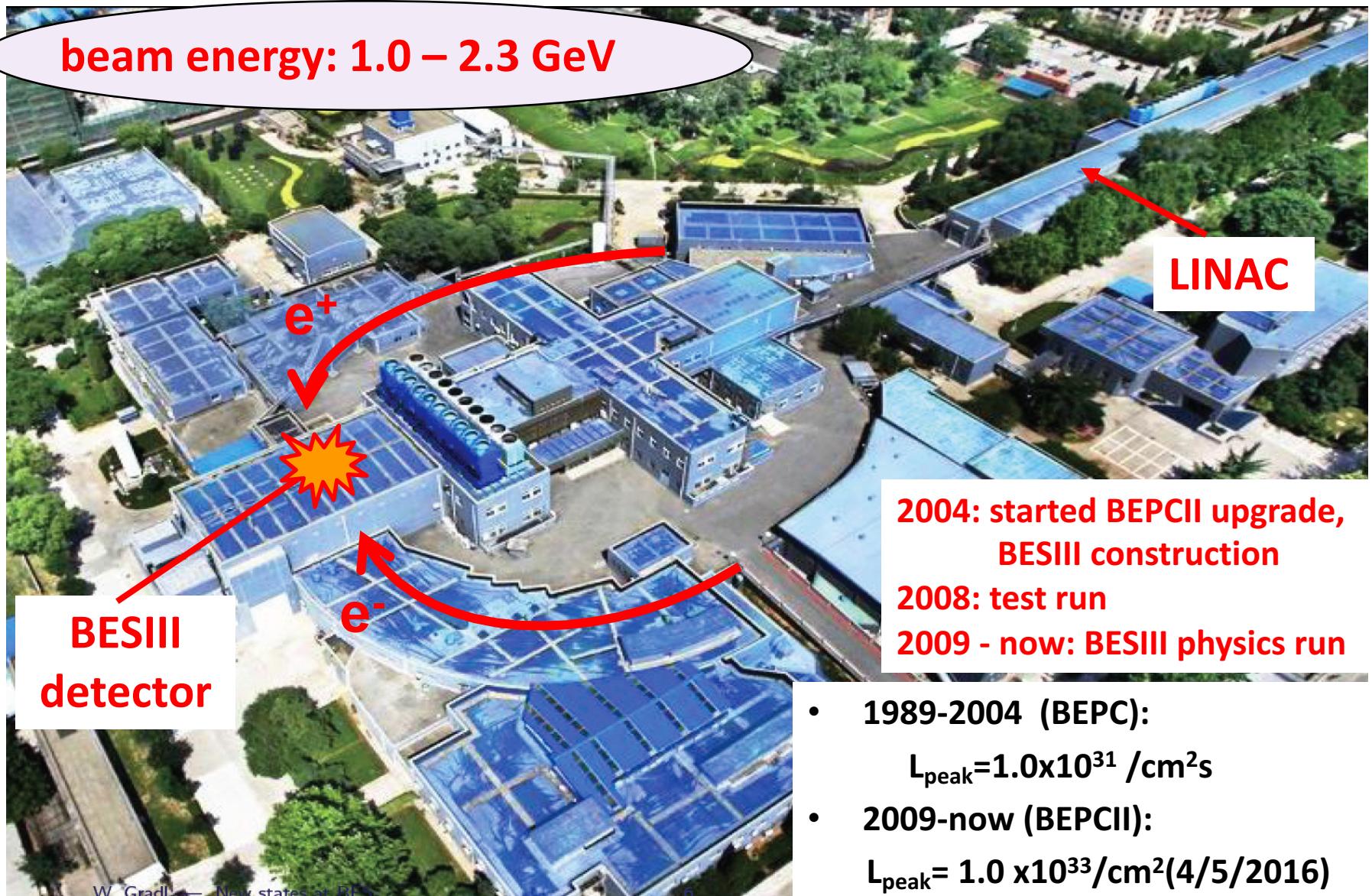


*C. Granados et al., EPJA 53, 117 (2017)

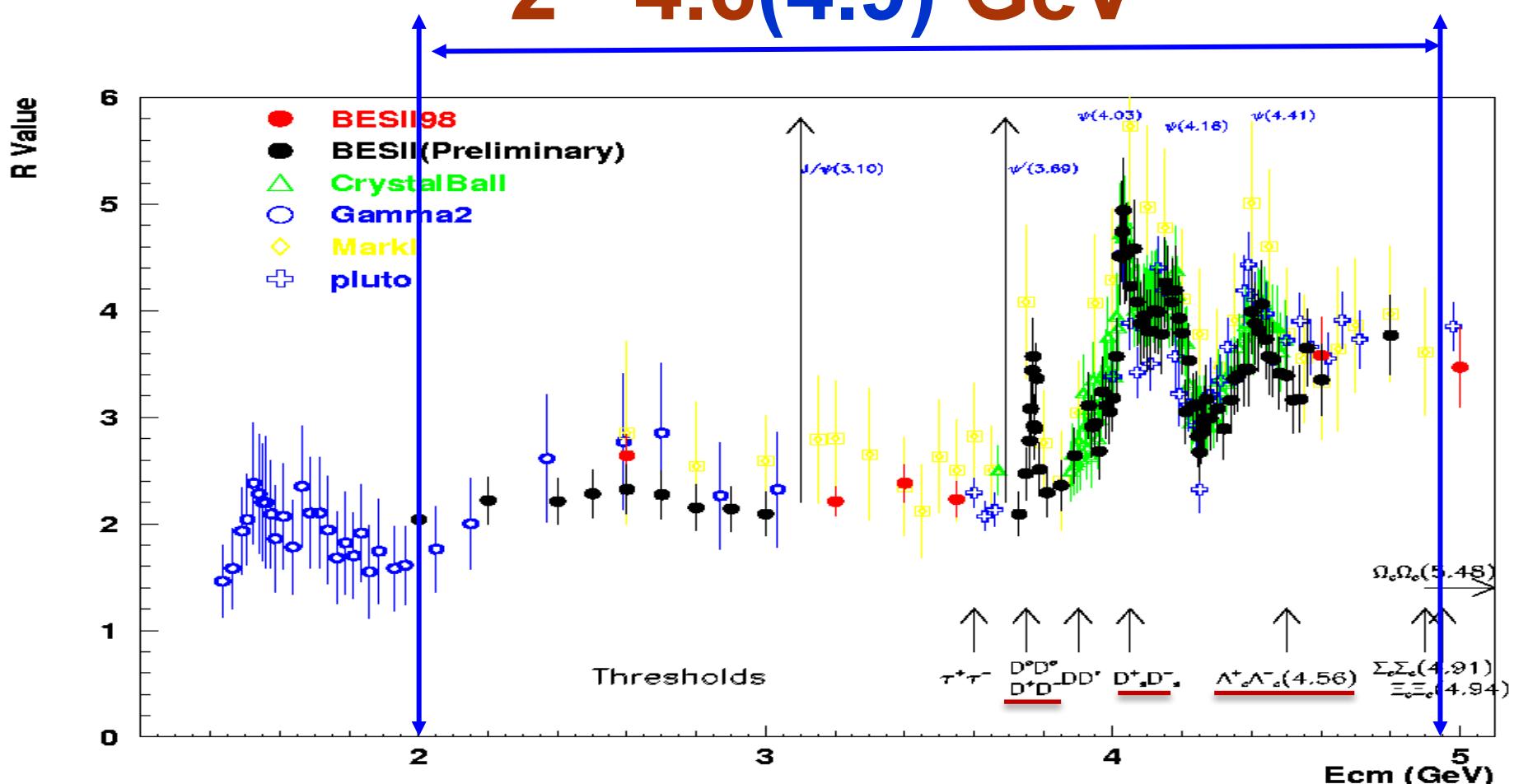


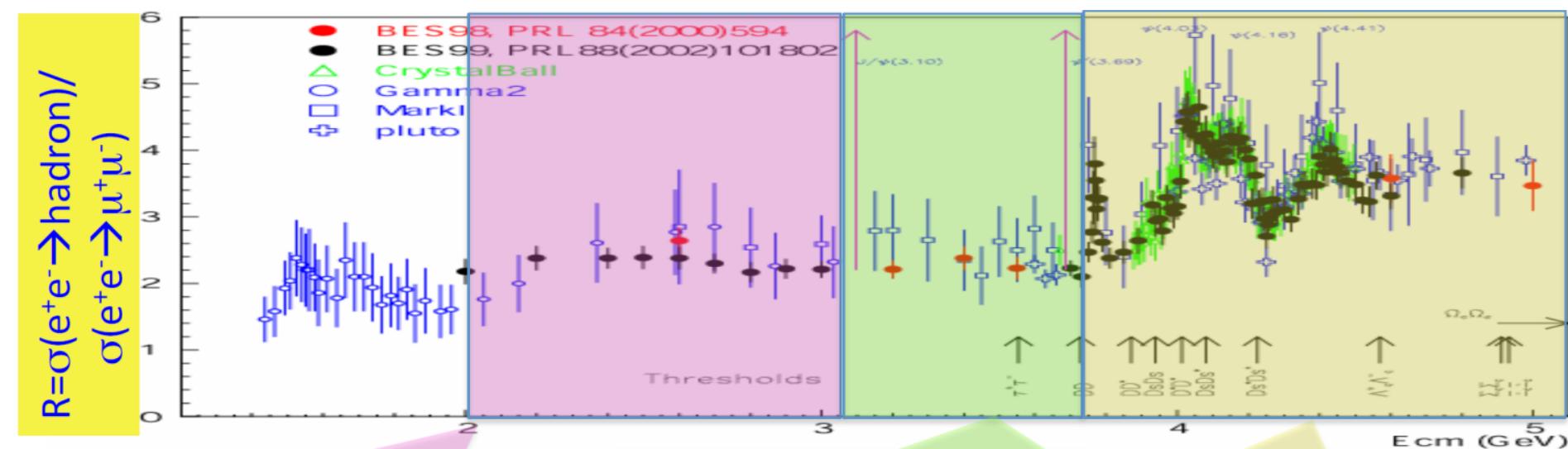
- Systems with strangeness
 - Scale: $m_s \approx 100$ MeV $\sim \Lambda_{\text{QCD}} \approx 200$ MeV: Relevant degrees of freedom?
 - **Probes QCD in the confinement domain.**
- Systems with charm
 - Scale: $m_c \approx 1300$ MeV: Quarks and gluons more relevant.
 - **Probes QCD just below pQCD.**





2 ~4.6(4.9) GeV





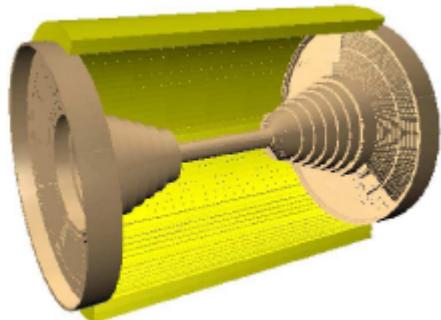
- Hadron form factors
 - $\Upsilon(2175)$ resonance
 - Multiquark states
with s quark, Zs
 - MLLA/LPHD and QCD
sum rule predictions

- Light hadron spectroscopy
 - Gluonic and exotic states
 - Process of LFV and CPV
 - Rare and forbidden decays
 - Physics with τ lepton

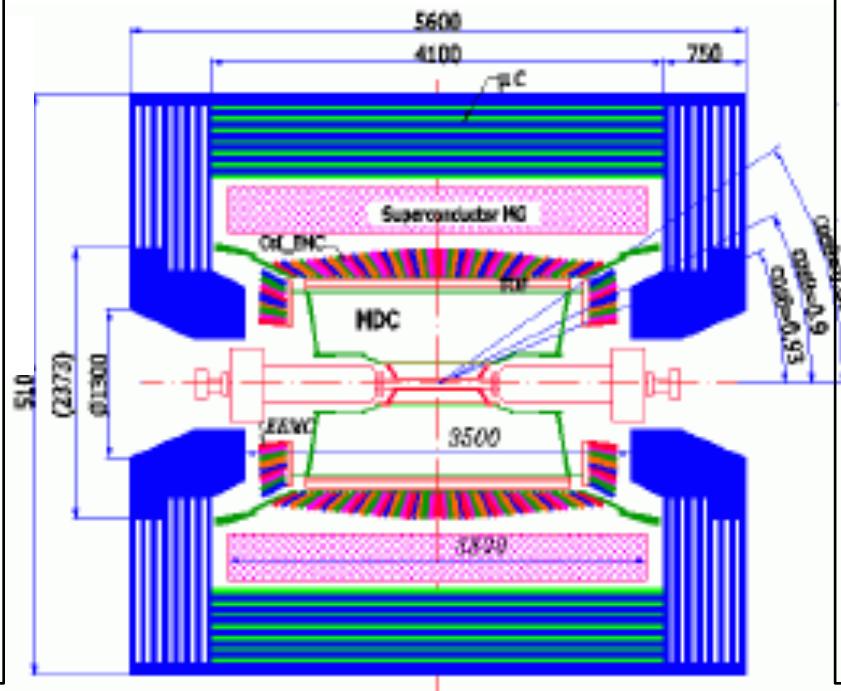
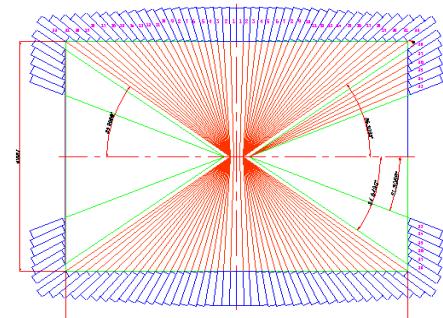
- XYZ particles
 - D mesons
 - f_D and f_{D_s}
 - D_0 - \bar{D}_0 mixing
 - Charm baryons

- Form factor
 - ✓ An important observable of hadron structure
 - ✓ To access time-like form factor through (ISR) electron-position annihilation process (*more details in Z. Meng's talk on Tuesday*)
 - To test the analyticity of G_E and G_M
 - Near threshold measurement of baryons: nucleon/hyperon
 - Unique hyperon data: Λ , Λ_c
 - ✓ *To extract Space-Like transition form factor using $\gamma\gamma^*\rightarrow P$*
- Fragmentation functions (FF)
 - Single hadron inclusive production
 - ➔ unpolarized FF
 - ✓ Constrain FF at low energy scale and high z: $K^{\pm/0}/\pi^{\pm/0}+X$
 - ✓ New data for inclusive $\eta/\phi/\Lambda$ production...
 - Two hadron inclusive production
 - ➔ correlation of two FF
 - ✓ To probe spin-dependent FF, eg, Collins FF

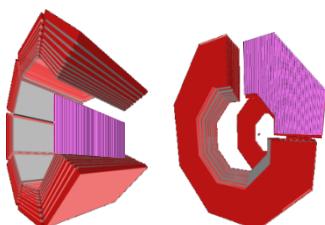
BESIII Detector

MDC

R inner: 63mm ;
R outer: 810mm
Length: 2582 mm
Layers: 43

**CsI(Tl) EMC**

Crystals: 28 cm(15 X_0)
Barrel: $|\cos\theta| < 0.83$
Endcap:
 $0.85 < |\cos\theta| < 0.93$

RPC MUC

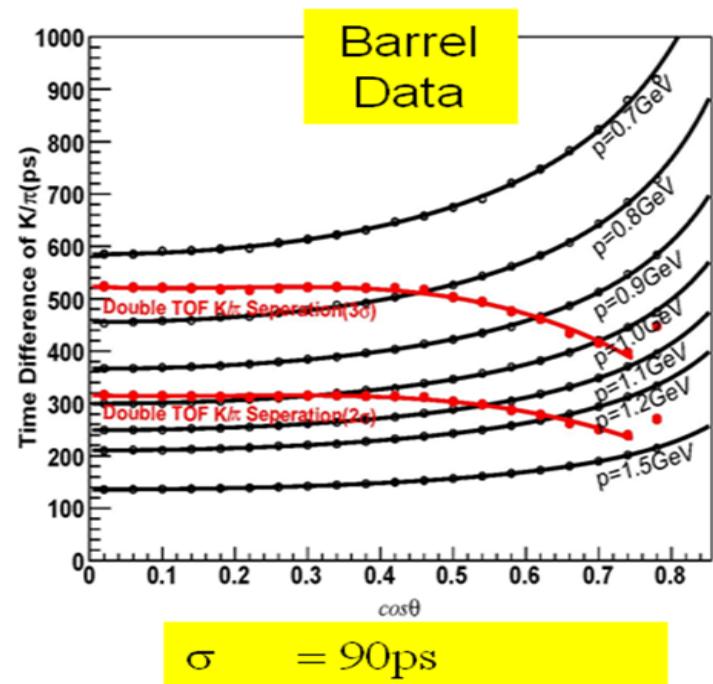
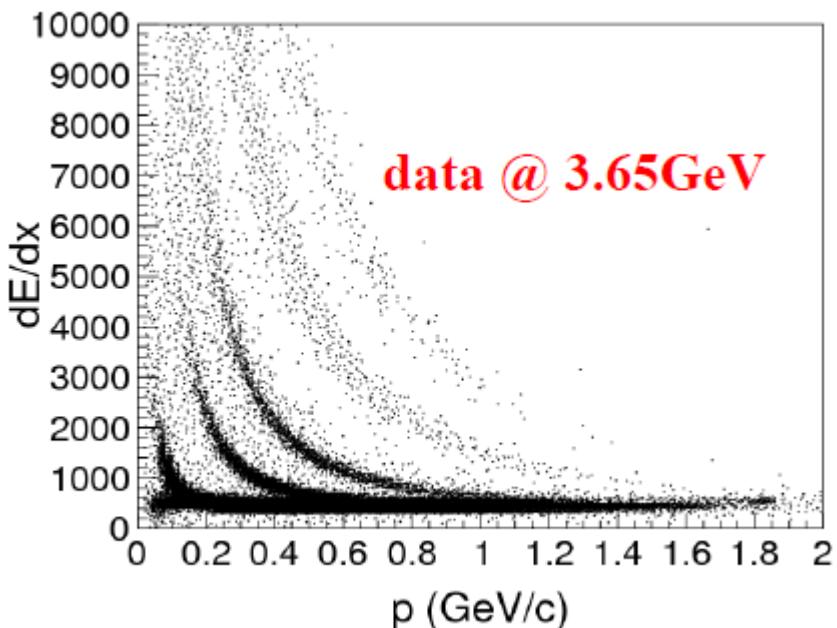
BMUC: 9 layers – 72 modules
EMUC: 8 layers – 64 modules

TOF

BTOP: two layers
ETOP: 48 crys. for each

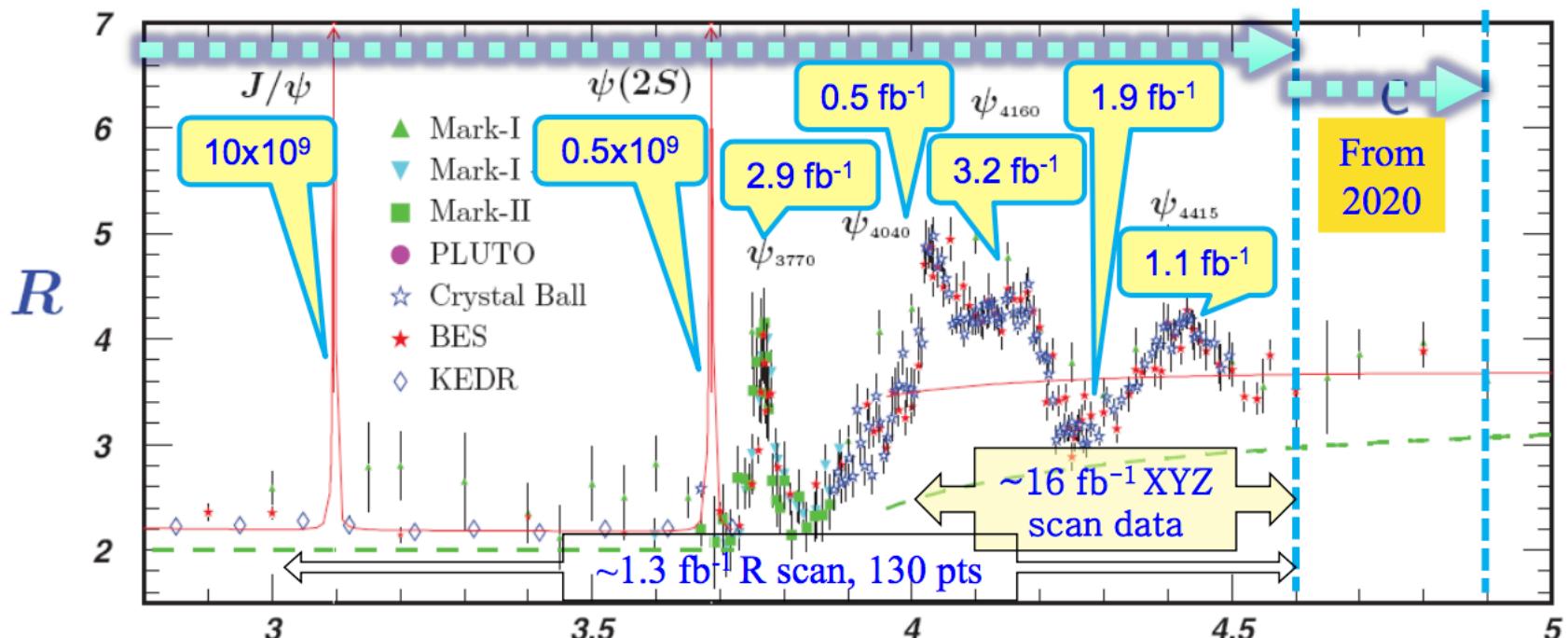


Particle ID Capacity



- dE/dx & TOF for PID
- Work well on $K/\pi < 1.0\text{GeV}$

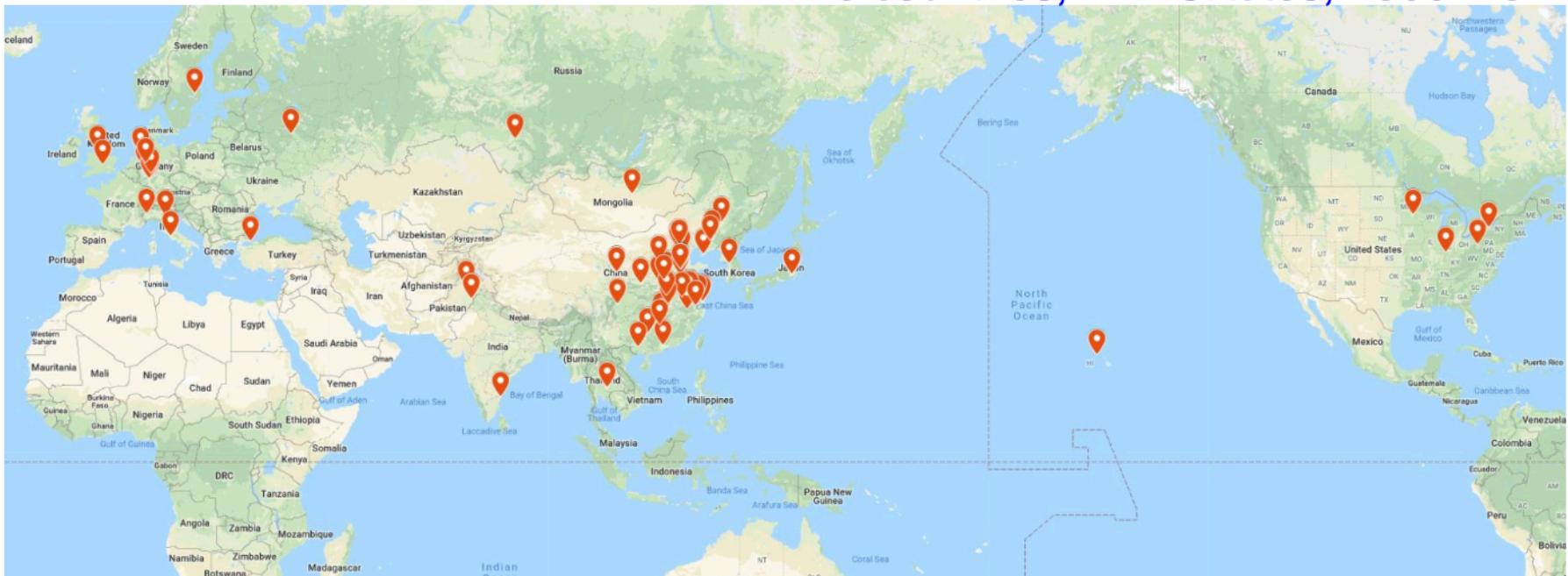
BESIII data taking



The BESIII Collaboration



15 countries, 72 institutes, ~500 members





J. C. Collins, Nucl.Phys. B396, 161 (1993)

$$D_{hq^\dagger}(z, P_{h\perp}) = D_1^q(z, P_{h\perp}^2)$$

$$+ \boxed{H_1^{\perp q}(z, P_{h\perp}^2)} \frac{(\hat{\mathbf{k}} \times \mathbf{P}_{h\perp}) \cdot \mathbf{S}_q}{z M_h},$$

D_1 : the unpolarized FF

H_1 : Collins FF

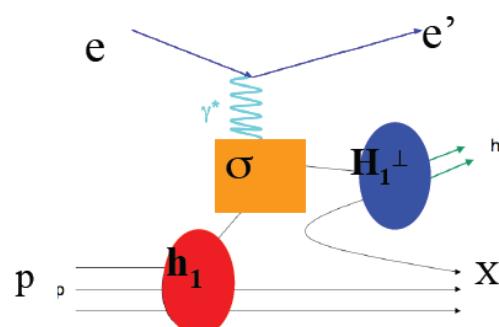
→ describes the fragmentation of a transversely polarized quark into a spinless hadron h .

→ depends on $z = 2E_h/\sqrt{s}$, $\mathbf{P}_{h\perp}$

→ leads to an azimuthal modulation of hadrons around the quark momentum.

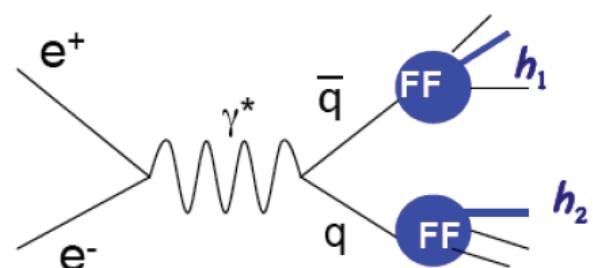
SIDIS

Transversity \otimes Collins FF



e+ e-

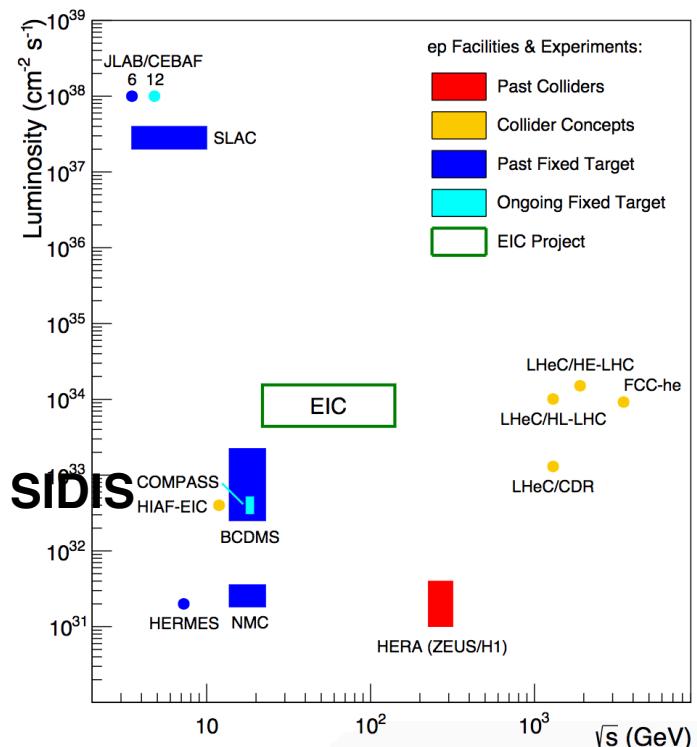
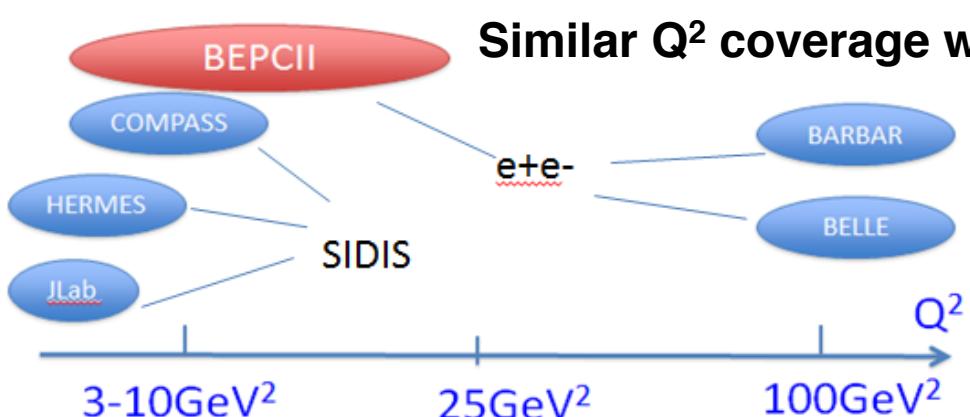
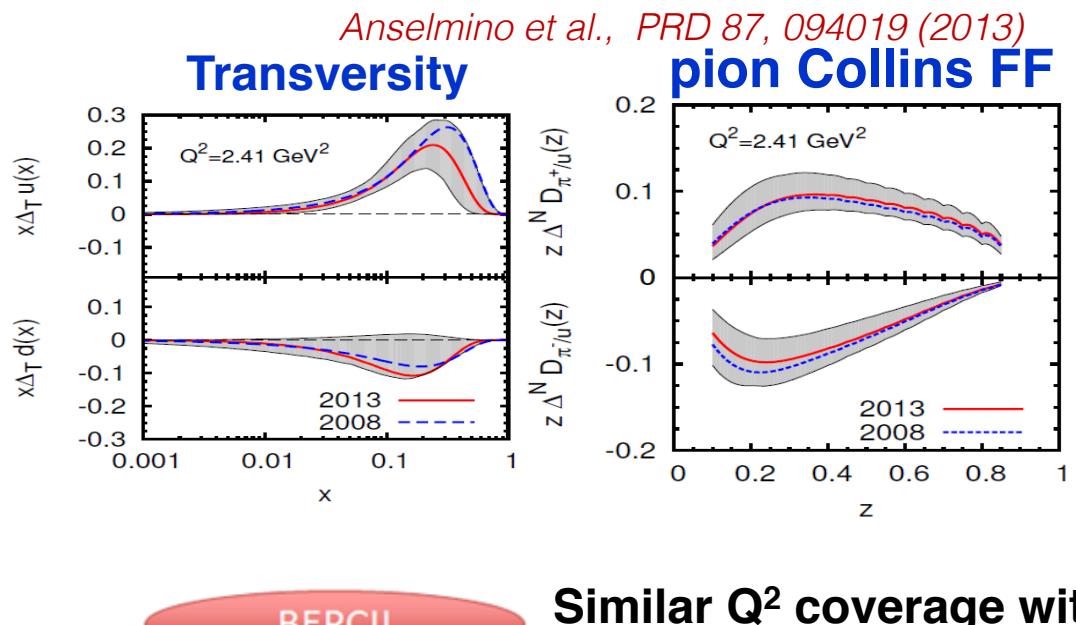
Collins FF \otimes Collins FF



Global Analysis on Collins FF

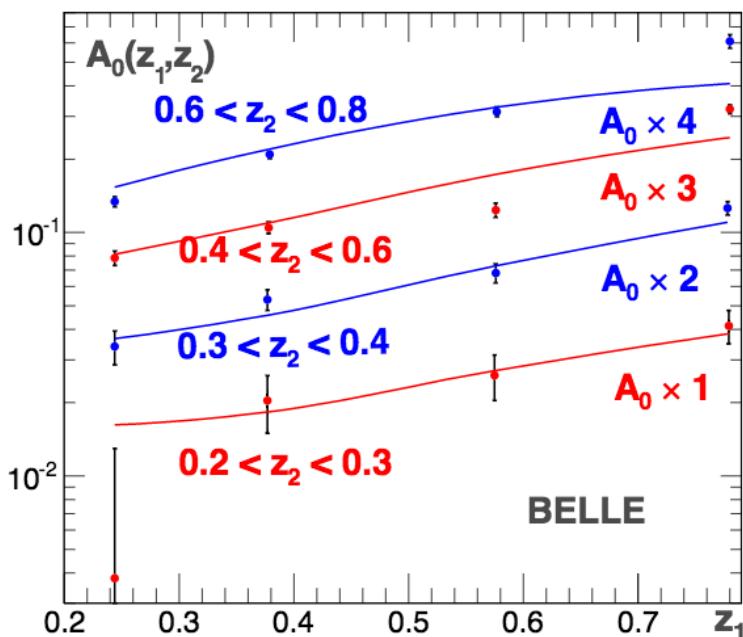


- The Q^2 evolution of Collins FFs need to be improved and validated
- Low Q^2 data from e^+e^- collider is useful.

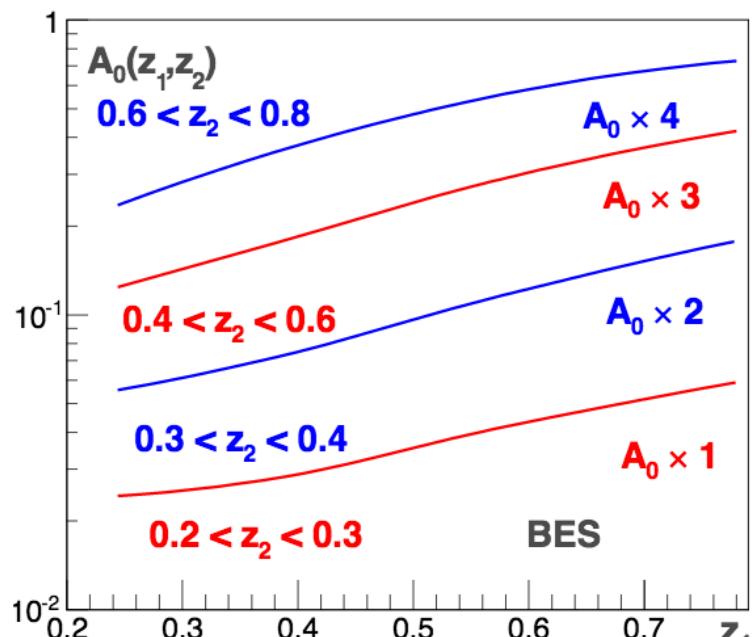


Predicted Collins Asymmetries

P. Sun and F. Yuan, Phys. Rev. D 88, 034016 (2013).



$Q^2 \sim 110 \text{ GeV}^2$



$Q^2 \sim 20 \text{ GeV}^2$

- Collins asymmetry is predicted larger at lower Q^2 region!
- Asymmetries increase as z grows

Product of Two Collins FFs

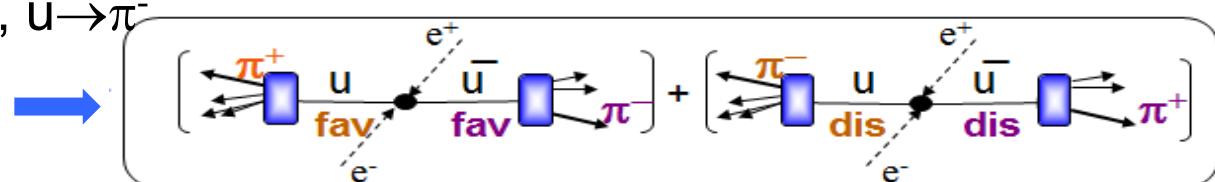


- **Favored** fragmentation process describes the fragmentation of a quark of flavor q into a hadron with a valence quark of the same flavor: i.e.: $u \rightarrow \pi^+$, $d \rightarrow \pi^-$

- **Disfavored** for $d \rightarrow \pi^+$, $u \rightarrow \pi^-$

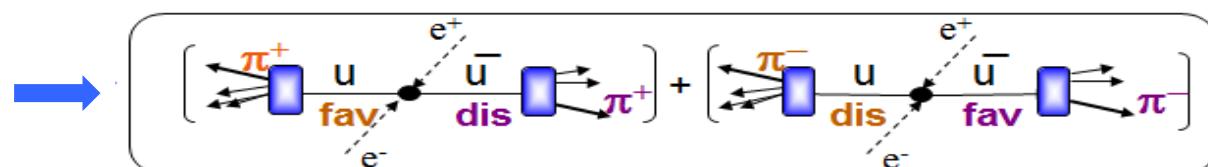
Unlike-sign pairs = U:

$$\pi^\mp\pi^\pm: (\text{fav} \times \text{fav}) + (\text{dis} \times \text{dis})$$



Like-sign pairs = L:

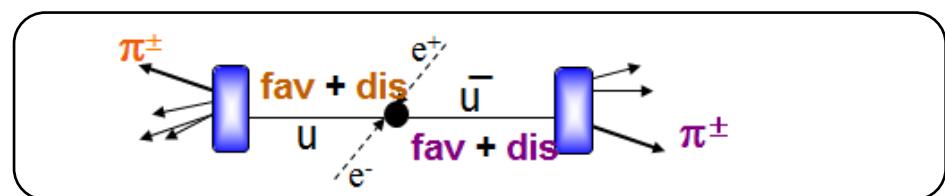
$$\pi^\pm\pi^\pm: (\text{fav} \times \text{dis}) + (\text{dis} \times \text{fav})$$



All charged pairs = C (U+L):

$$\pi\pi: (\text{fav} + \text{dis}) \times (\text{fav} + \text{dis})$$

$$\pi = \pi^\pm$$



- **All charged pion pairs are divided into:**

- **Unlike-sign pairs ($\pi^+\pi^-$)**
- **Like-sign pairs ($\pi^+\pi^+$ and $\pi^-\pi^-$)**
- **All Charged pairs ($\pi\pi$)**



PRL 116, 042001 (2016)



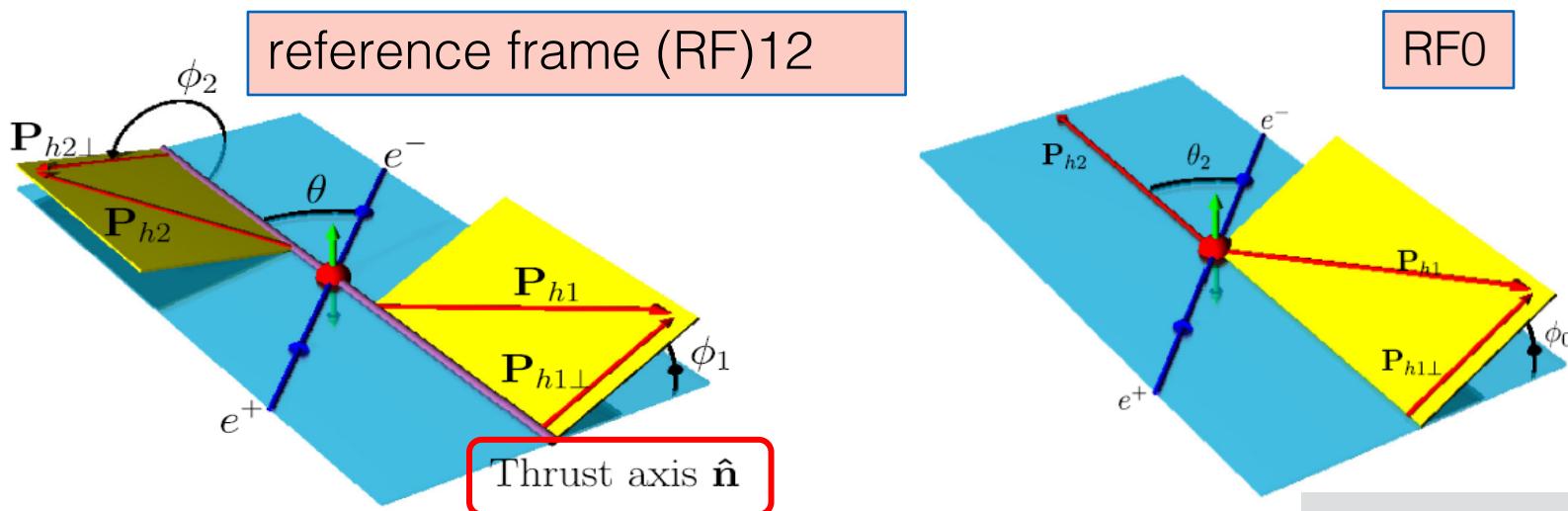
~62 pb⁻¹ @3.65GeV

- continuum region in-between J/psi and psi(2S) peaks
- fewer charm backgrounds!

- ◆ # of charged tracks ≥ 3
- ◆ # of charged pion ≥ 2
- ◆ No electron to suppress Bhabha
- ◆ The total visible energy $> 1.5 \text{ GeV}$.

to suppress resonance decays to dipions,
we require the charged pions

- ◆ $0.2 < z = 2E_h/\sqrt{s} < 0.9$
- ◆ open angle $> 120^\circ$ to select back-to-back pion-pair.

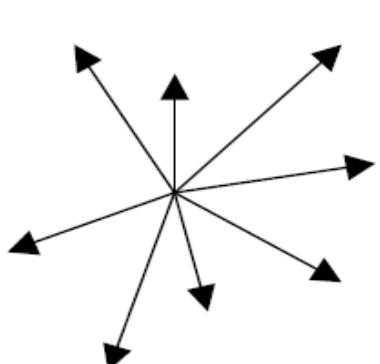


PRD 78, 032011

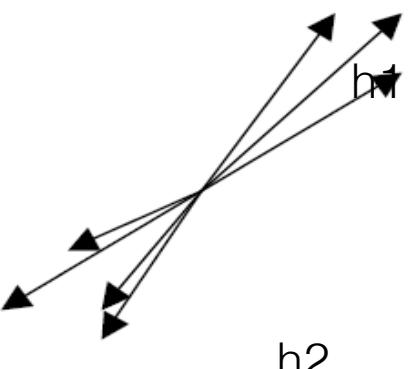
- The spin correlation of hadron fragmented from quark and anti-quark in opposite hemisphere follows a $\cos(\phi_1 + \phi_2)$ modulation in RF12 or a $\cos(2\phi_0)$ modulation in RF0
 - DCA were observed in both definitions by Belle and BaBar

Belle Collaboration: Phys. Rev. Lett. 96, 232002 (2006); Phys. Rev. D 78, 032011 (2008);
Phys. Rev. D 86, 039905(E) (2012).

BaBar Collaboration: Phys. Rev. D 90, 052003 (2014)



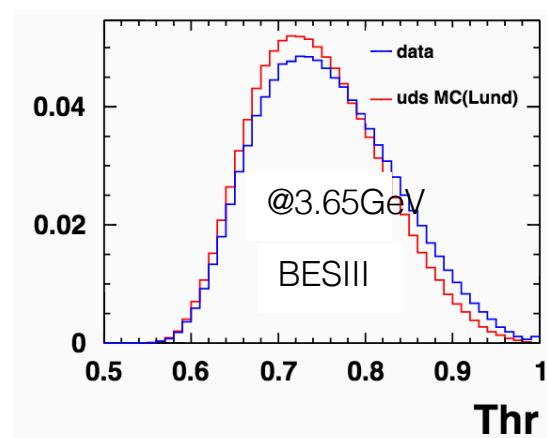
BEPCII
energy



Belle/Babar

$$e^+ e^- \rightarrow q\bar{q} \rightarrow h_1 h_2 X$$

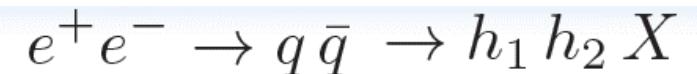
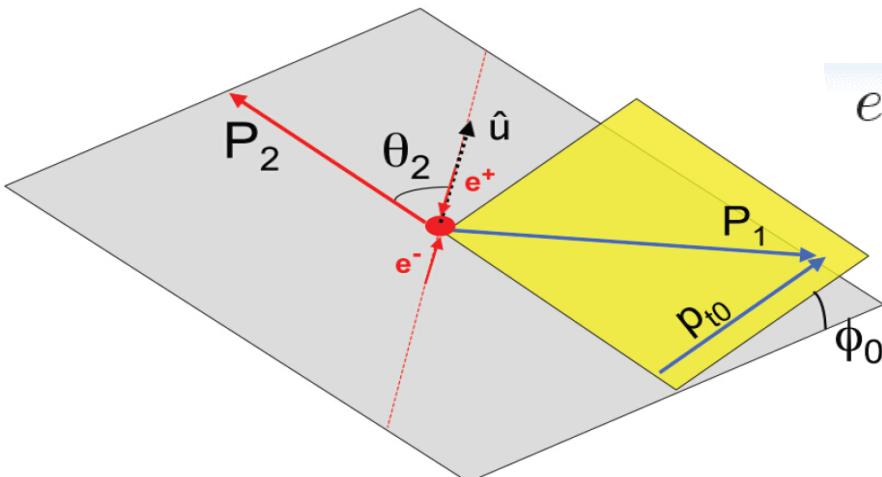
$$T = \frac{\max_h \sum |\mathbf{P}_h^{\text{CMS}} \cdot \hat{\mathbf{n}}|}{\sum_h |\mathbf{P}_h^{\text{CMS}}|},$$



- ✓ In B factories, the thrust axis is used as a approximation of the q-qbar axis
 - not needed in RF0, hence RF0 is adopted
- ✓ At BESIII, continuum process events distribute in a more isotropic way
 - ➔ very different situation compared to Belle/BaBar

The Reference Frame

D. Boer, Nucl. Phys. B 806, 23 (2009)



$$\sigma \sim 1 + \frac{\sin^2 \theta_2}{1 + \cos^2 \theta_2} \cos(2\phi_0) \mathcal{F} \left[\frac{H_1^\perp(z_1) H_1^\perp(z_2)}{D_1(z_1) \bar{D}_1(z_2)} \right]$$

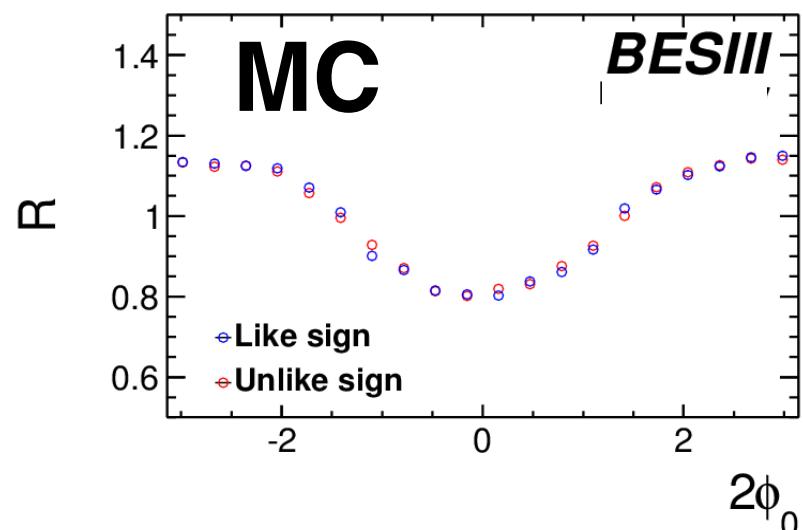
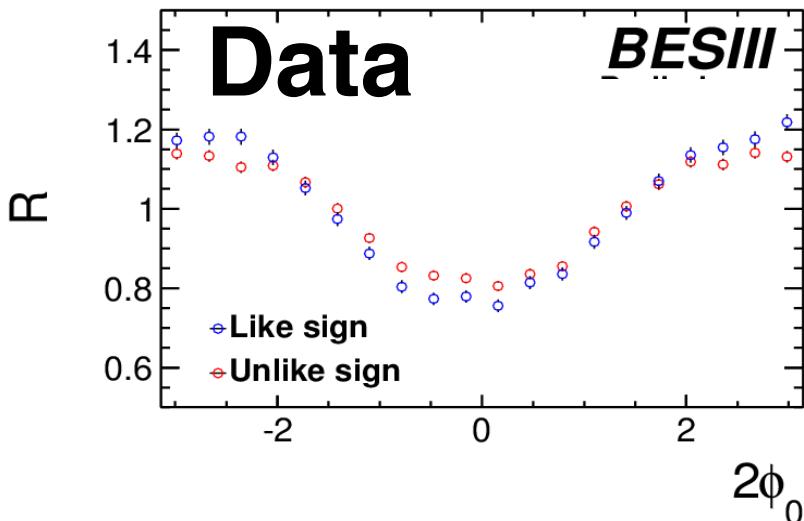
$$\mathcal{F}[X] = \sum_{q\bar{q}} \int [2\hat{h} \cdot k_{T1} \hat{h} \cdot k_{T2} - k_{T1} \cdot k_{T2}]$$

$$d^2 k_{T1} d^2 k_{T2} \delta^2(k_{T1} + k_{T2} - q_T) X$$

$$k_{Ti} = z_i p_{Ti}$$

Collins effect: cosine modulation.

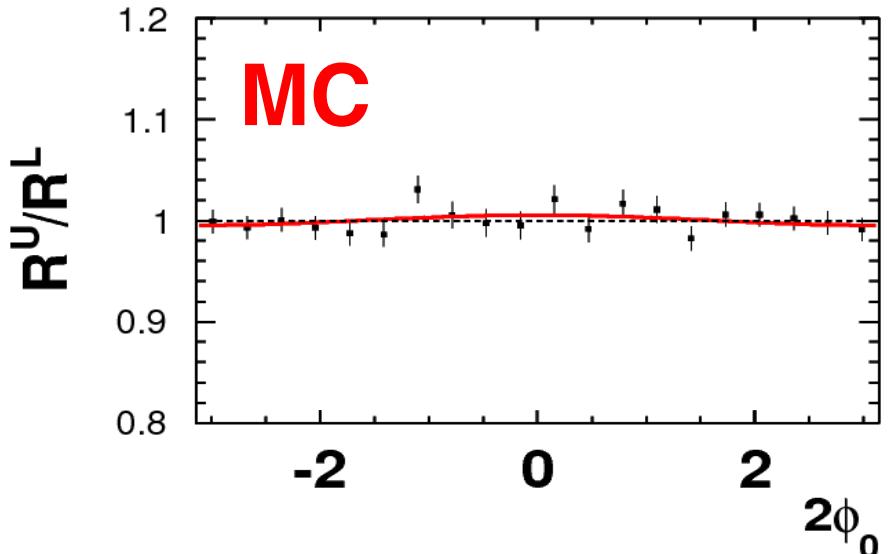
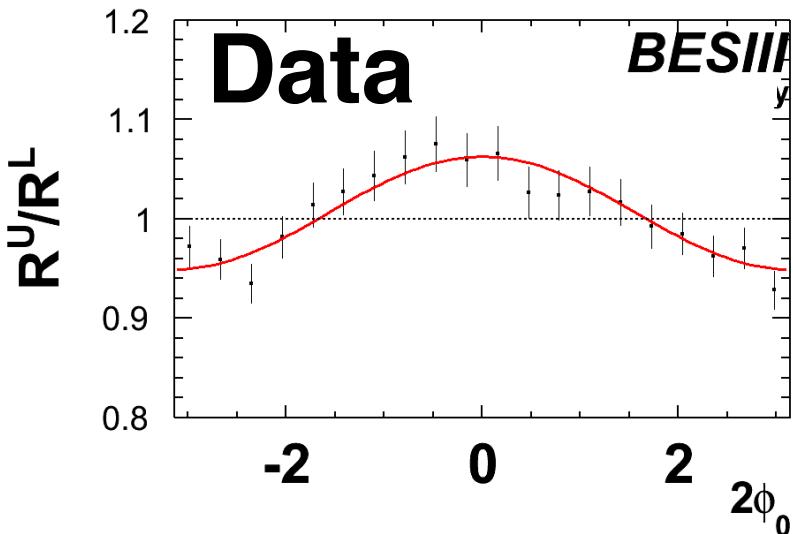
2 ϕ_0 Raw Distribution



- The normalized ratio

$$R := \frac{N(2\phi_0)}{\langle N_0 \rangle}$$
 ➤ For Charged, Unlike-sign and Like-sign pairs, we have R^U , R^C , R^L
- Raw $2\phi_0$ distributions are subjected to the limited acceptance and non-uniform efficiencies of the detector!
- The MC simulation does not include the Collins effect.
- Small deviations in Like and Unlike in data indicate asymmetries.

Double Ratio (DR)

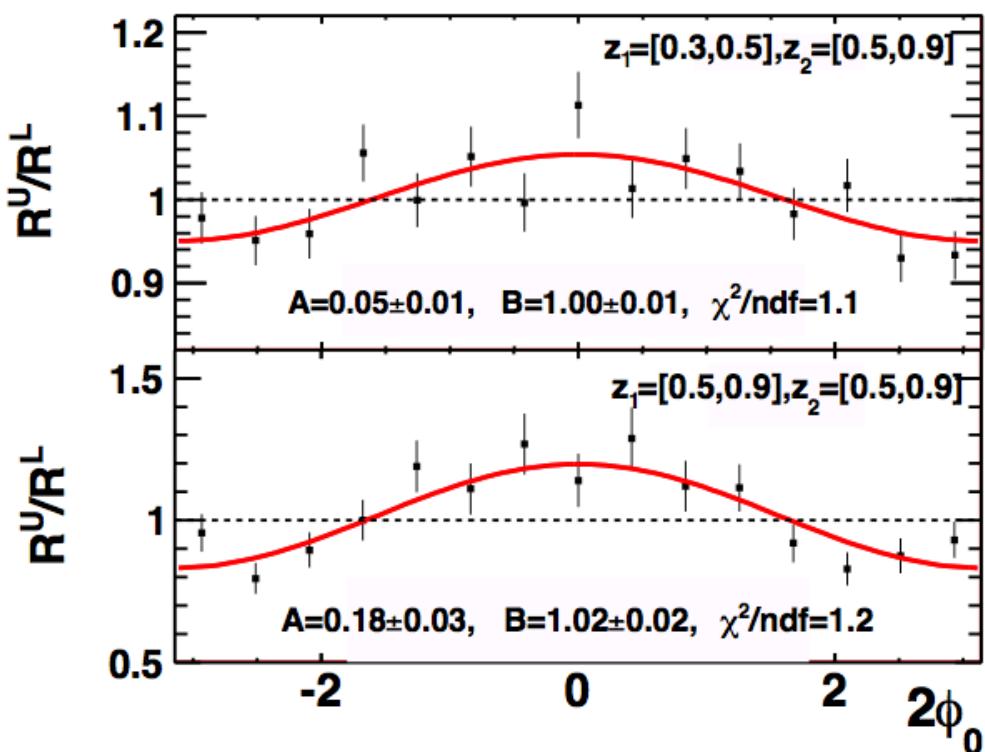


- Acceptance effects and radiation effects can be reduced by performing the ratio of Unlike/Like sign pion pairs (R^U/R^L) or Unlike/Charged pairs (R^U/R^C)

$$\frac{R^U}{R^L} \simeq 1 + \frac{\sin^2 \theta}{1 + \cos^2 \theta} \cos(2\phi_0) (G^U - G^L) \quad \frac{R^U}{R^C} \simeq 1 + \frac{\sin^2 \theta}{1 + \cos^2 \theta} \cos 2\phi_0 (G^U - G^C)$$

- DRs are fitted by $\frac{R^U}{R^{L(C)}} = a \cos(2\phi_0) + b$,

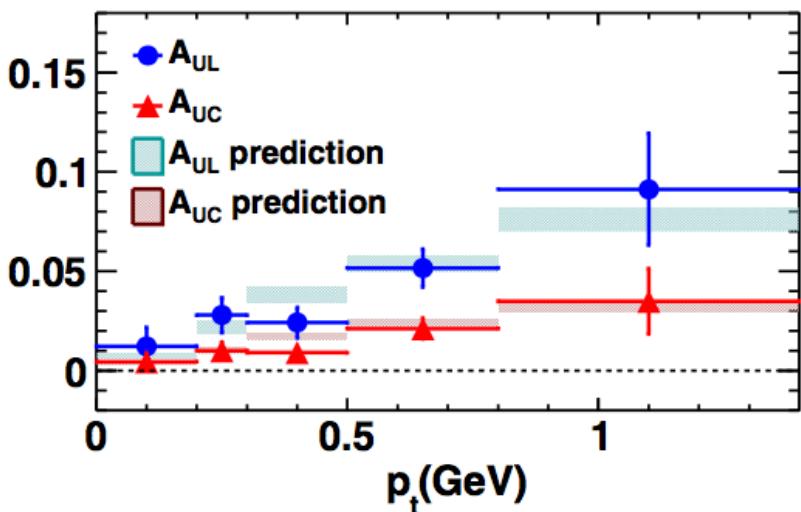
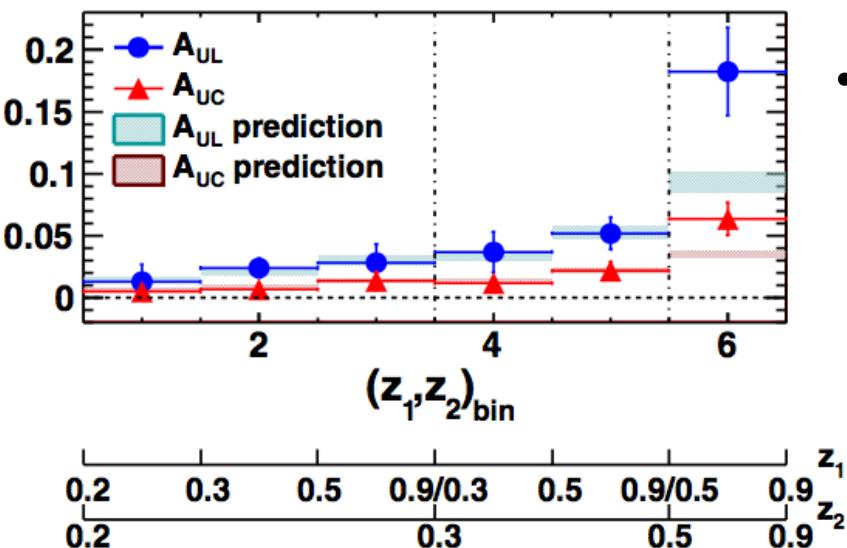
- $A_{UL(C)} = \frac{a}{b}$ represents the asymmetries of interest.

Fit to DR in Different z Bins

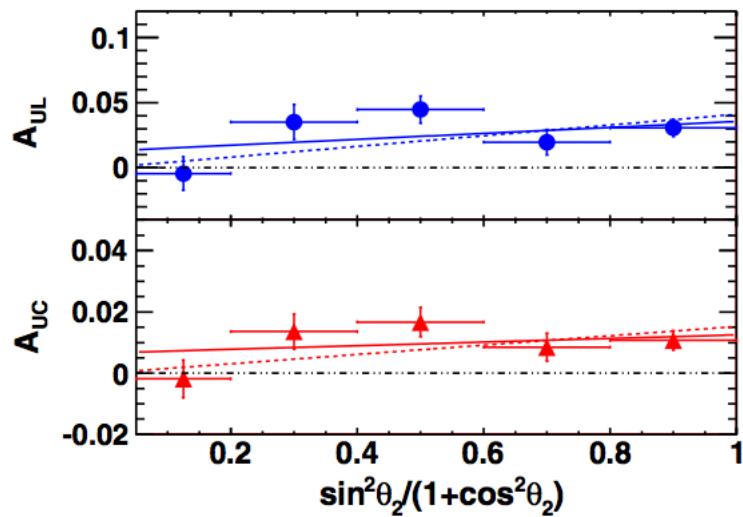
$$\frac{R^U}{R^{L(C)}} = A \cos(2\phi_0) + B,$$

Extraction of the Asymmetries

PRL 116, 042001 (2016)



- Collins asymmetry measured as functions of
 - ✓ 6 symmetric (z_1, z_2) bins
 - ✓ 5 bins of p_{t0}
 - ✓ Asymmetry v.s. $\sin^2 \theta_2 / (1 + \cos^2 \theta_2)$
 - ✓ Comparison with prediction in PRD 93, 014009 (2016)



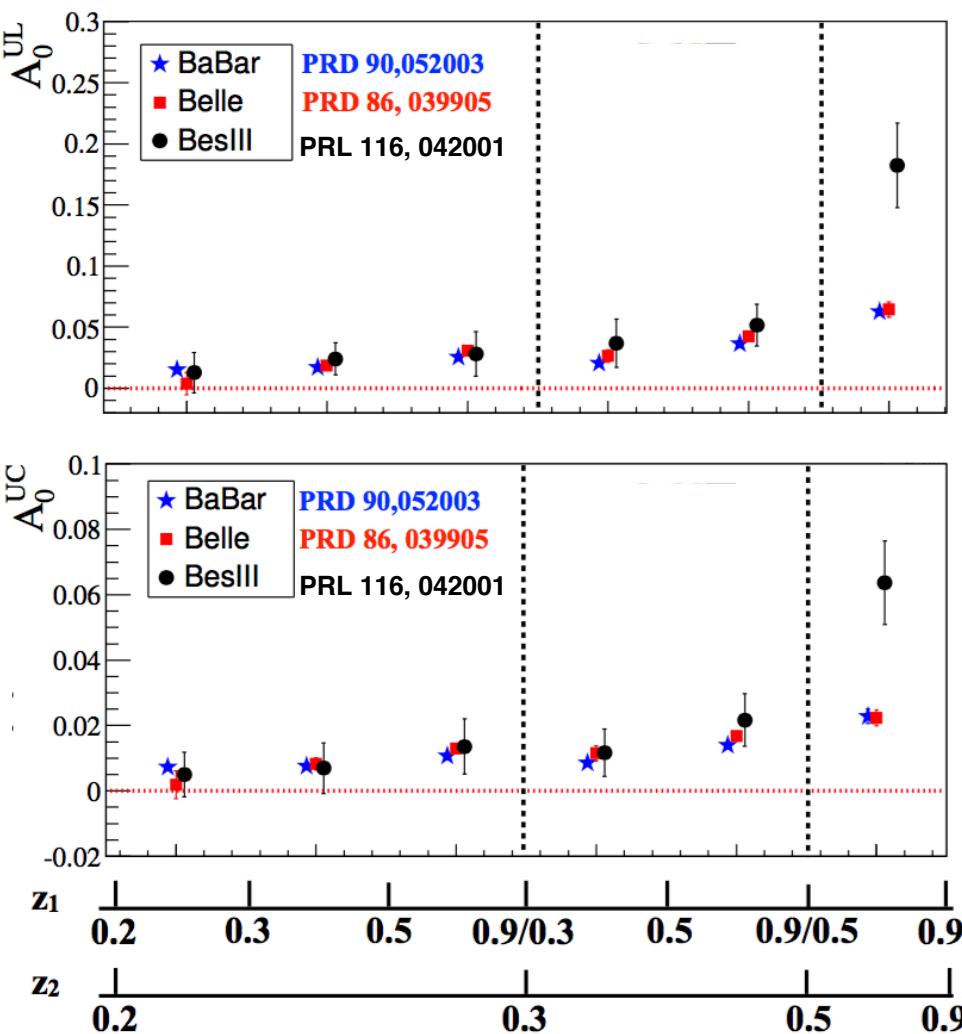
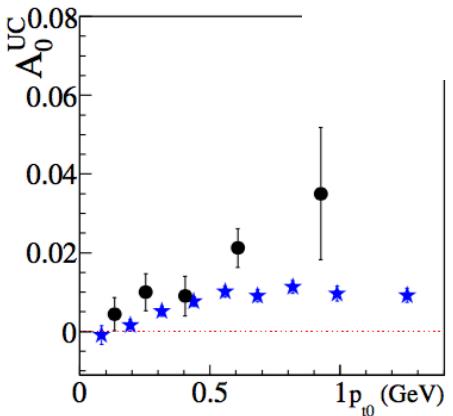
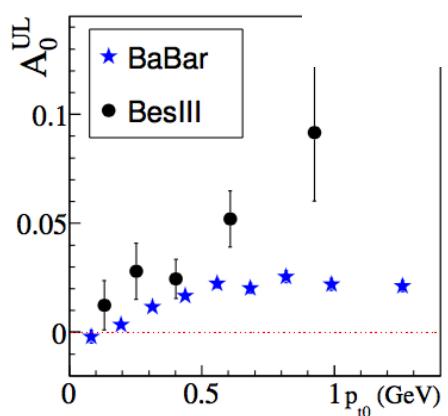
Comparisons of Asymmetries



Comparison between different results obtained at different Q^2 :

- **BaBar and Belle @ $Q^2 \sim 110 \text{ GeV}^2$**
- **BESIII @ $Q^2 \sim 13 \text{ GeV}^2$**

- BaBar and Belle results that fall in the larger BESIII z-bins are averaged taking into account the statistical and systematic uncertainties
- Good agreement between different data set for low z
- BESIII larger asymmetries in the last z-bins: consistent with the prediction reported in *arXiv:1505.05589*



Relevant Future Analyses

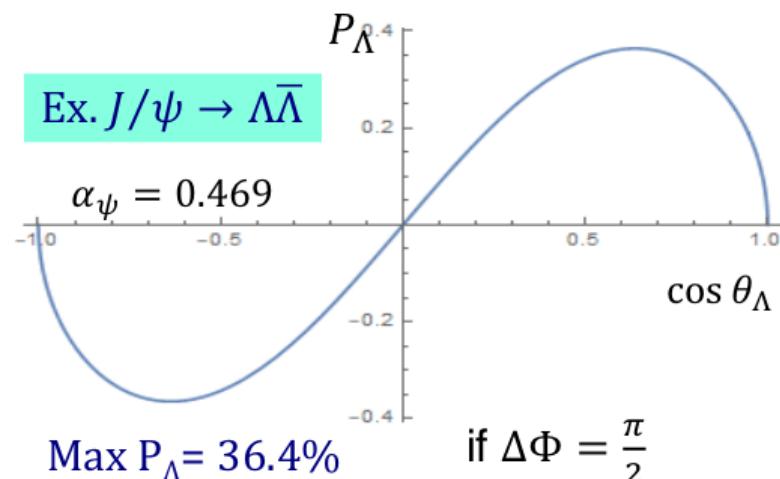
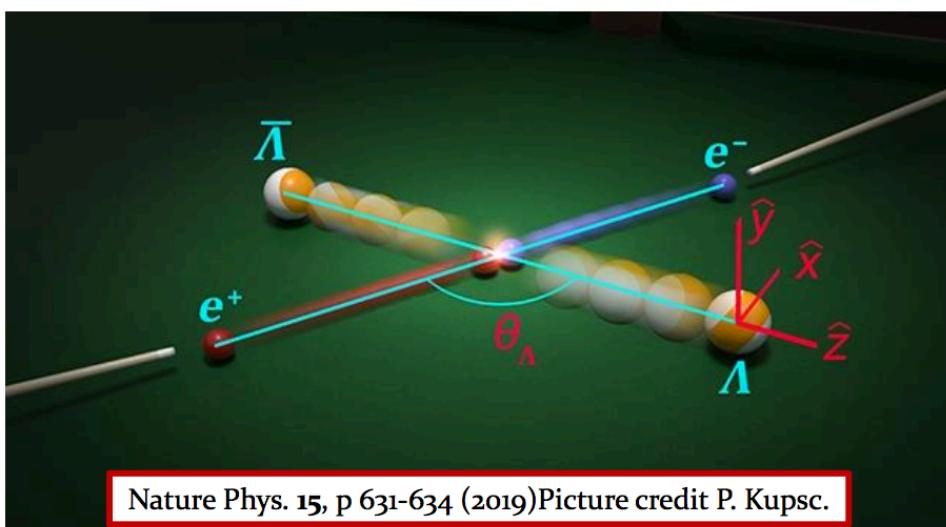
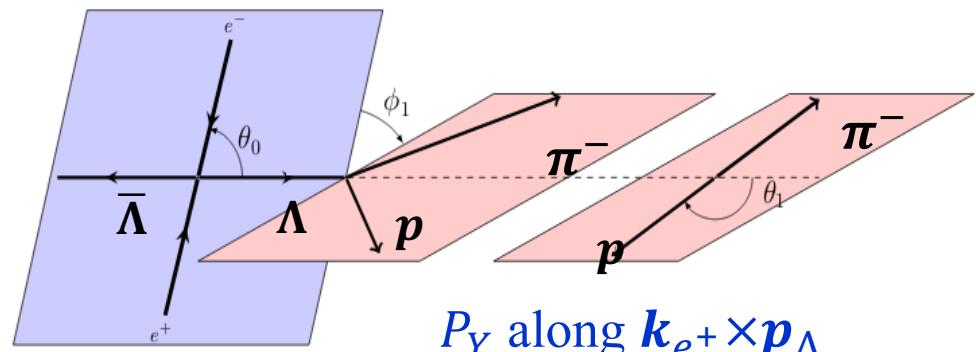


- $\sim 360/\text{pb}$ data @ $\sqrt{s} = 3.51\text{GeV}$, which will improve precision of the dipion Collins asymmetry
- Simultaneous measurement of charged KK, K π and $\pi\pi$ Collins asymmetries (PID unfolding)
- Including neutral particles (π^0 , η , Ks) in final states
- Potential of data above charm threshold (up to $Q^2 \sim 25\text{ GeV}^2$) can be explored

Unpolarized e^+e^- annihilations:

- entangled $\Lambda\bar{\Lambda}$ production
- more sensitive to access transverse polarization baryon

$$|\Lambda\bar{\Lambda}\rangle^{C=-1} = \chi_1 \frac{1}{\sqrt{2}} [|\Lambda\rangle|\bar{\Lambda}\rangle - |\bar{\Lambda}\rangle|\Lambda\rangle],$$



$$\bar{P}_Y(\cos \theta_\Lambda) = \frac{\sqrt{1 - \alpha_\psi^2} \cos \theta_\Lambda \sin \theta_\Lambda}{1 + \alpha_\psi \cos^2 \theta_\Lambda} \sin(\Delta\Phi)$$

$\Delta\Phi$: complex phase between two helicity amplitudes: $A_{1,\frac{1}{2},\frac{1}{2}}$ and $A_{1,\frac{1}{2},-\frac{1}{2}}$

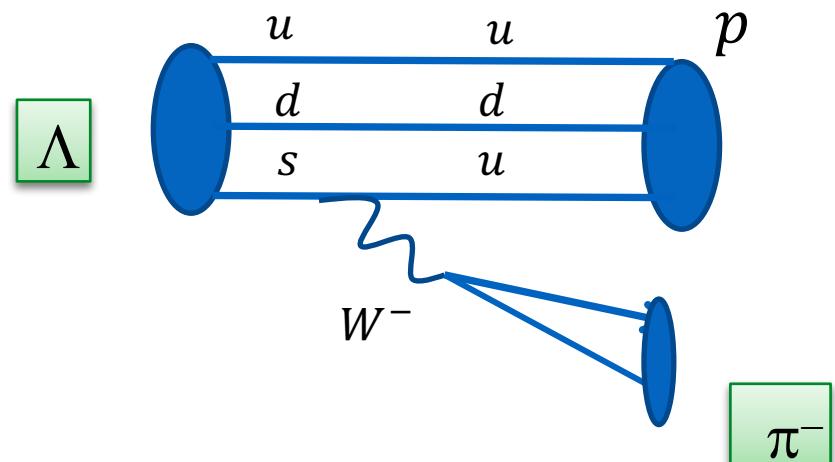
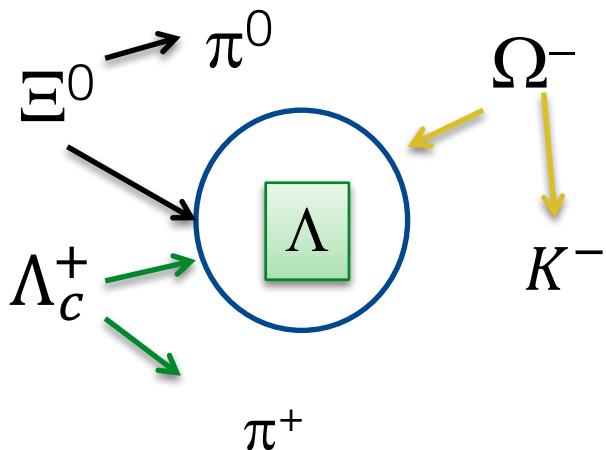
$\Lambda \rightarrow p\pi^-$ as Polarimeter

$$\frac{dN}{d\Omega} = \frac{1}{4\pi} (1 + \alpha_\Lambda \vec{P} \cdot \hat{q}) = \frac{1}{4\pi} (1 + \alpha_\Lambda P_\Lambda \cos\theta_p)$$

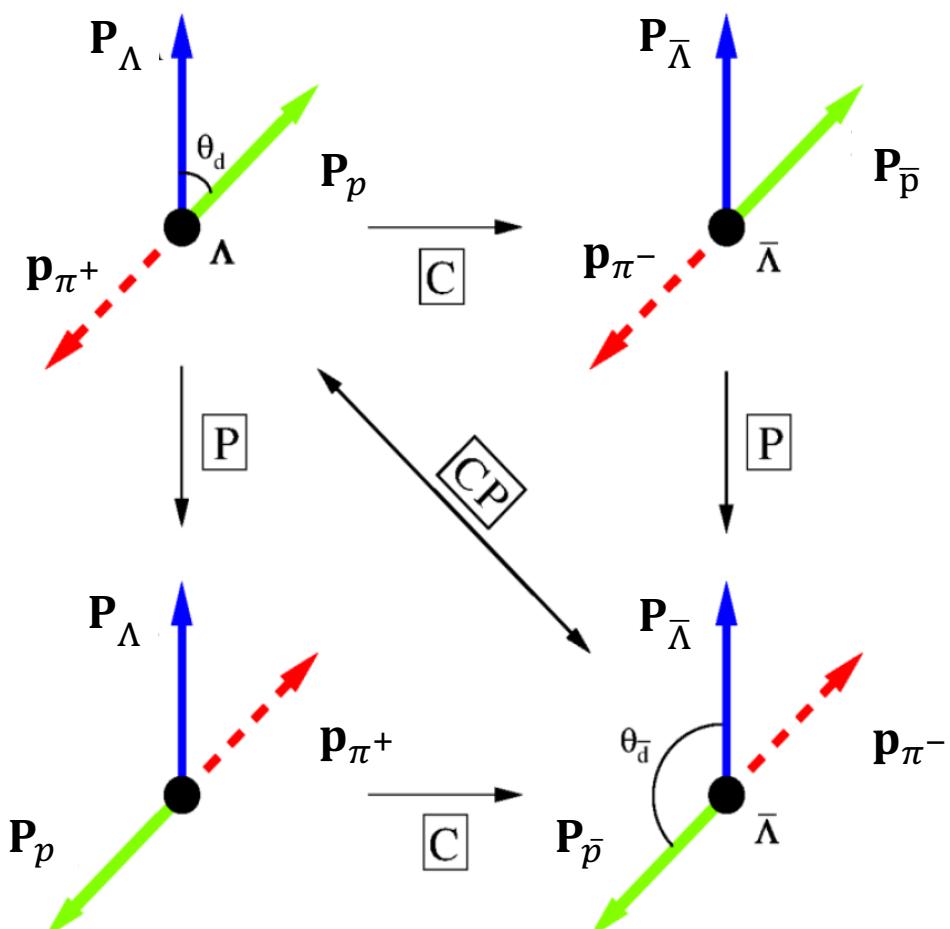
Lee-Yang parameters:

$$\alpha = \frac{2 \operatorname{Re}(S^* P)}{|S|^2 + |P|^2}, \quad \beta = \frac{2 \operatorname{Im}(S^* P)}{|S|^2 + |P|^2}, \quad \gamma = \frac{|S|^2 - |P|^2}{|S|^2 + |P|^2}$$

Note: $\alpha^2 + \beta^2 + \gamma^2 = 1$



C- and P- transformation



$$\alpha_\Lambda = \frac{|B_+|^2 - |B_-|^2}{|B_+|^2 + |B_-|^2}, \alpha_{\bar{\Lambda}} = \frac{|\bar{B}_+|^2 - |\bar{B}_-|^2}{|\bar{B}_+|^2 + |\bar{B}_-|^2}$$

CP invariance:

$$\bar{B}_{-\lambda_p} = \eta_\Lambda \eta_p \eta_\pi (-1)^{s_\Lambda - s_p - s_\pi} B_{\lambda_p} = -B_{\lambda_p}$$

If CP invariance:

$$\alpha_\Lambda = -\alpha_{\bar{\Lambda}}$$

Previous Measurements

2018 PDG list

α_- FOR $\Lambda \rightarrow p\pi^-$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT	
0.642 ± 0.013	OUR AVERAGE				
0.584 ± 0.046	8500	ASTBURY	1975	SPEC	
0.649 ± 0.023	10325	CLELAND	1972	OSPK	
0.67 ± 0.06	3520	DAUBER	1969	HBC	From Ξ decay
0.645 ± 0.017	10130	OVERSETH	1967	OSPK	Λ from $\pi^- p$
0.62 ± 0.07	1156	CRONIN	1963	CNTR	Λ from $\pi^- p$

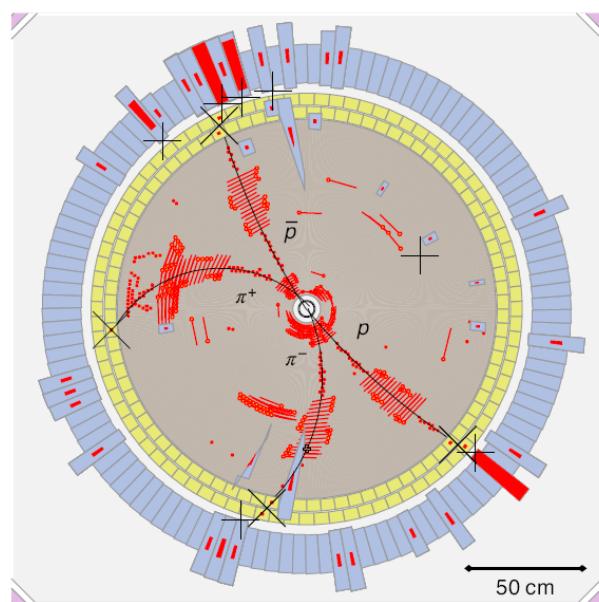
α_+ FOR $\bar{\Lambda} \rightarrow \bar{p}\pi^+$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT	
-0.71 ± 0.08	OUR AVERAGE				
$-0.755 \pm 0.083 \pm 0.063$	$\approx 8.7k$	ABLIKIM	2010	BES	$J/\psi \rightarrow \Lambda\bar{\Lambda}$
-0.63 ± 0.13	770	TIXIER	1988	DM2	$J/\psi \rightarrow \Lambda\bar{\Lambda}$

$$e^+ e^- \rightarrow J/\psi \rightarrow \Lambda \bar{\Lambda} \rightarrow p \bar{p} \pi^+ \pi^-$$

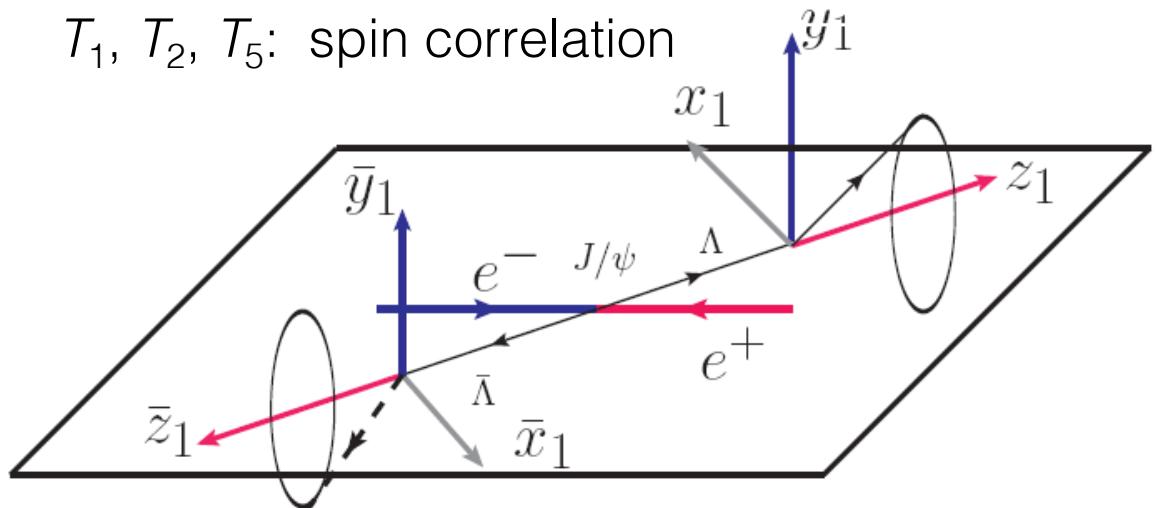
G. Faldt and A. Kupsc
 PLB, 772, 16 (2017)
 EPJA, 52, 141 (2016)
 EPJA, 51, 74 (2015).

$$\begin{aligned} \frac{d\sigma}{d\Omega} \propto & T_0 + \sqrt{1 - \alpha_{J/\psi}^2} \sin(\Delta)(\alpha_\Lambda T_3 + \alpha_{\bar{\Lambda}} T_4) \\ & + \alpha_\Lambda \alpha_{\bar{\Lambda}} [T_1 + \sqrt{1 - \alpha_{J/\psi}^2} \cos(\Delta) T_2 + \alpha_{J/\psi} T_5], \end{aligned}$$



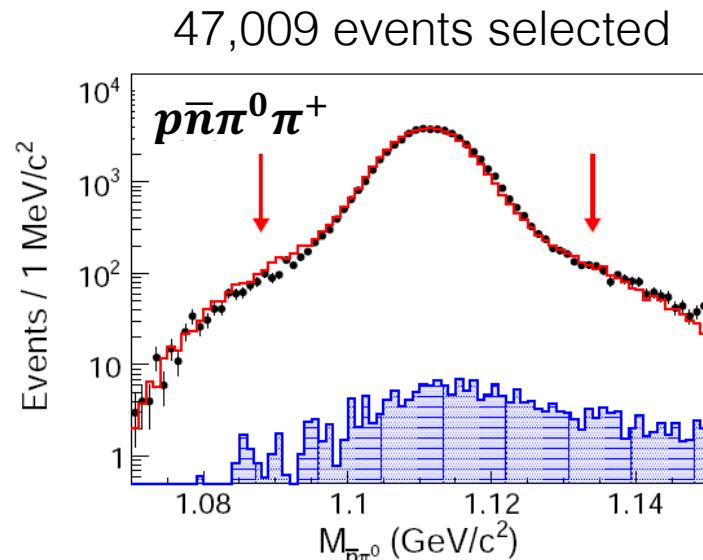
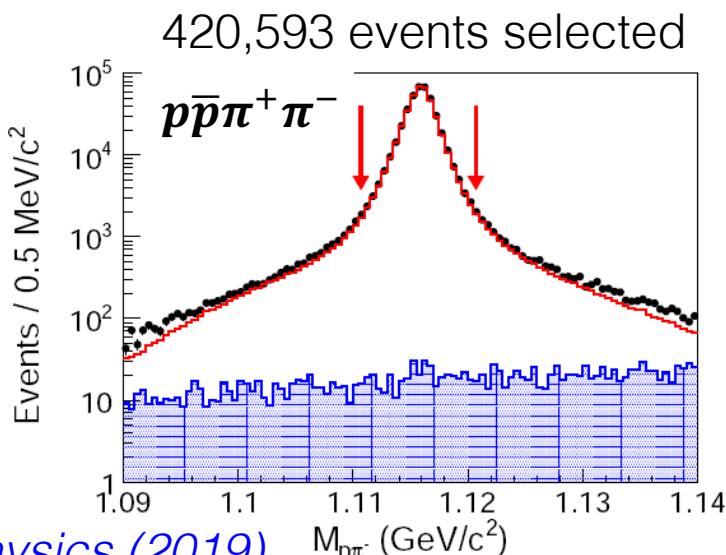
T_0 : angular distribution of Λ and $\bar{\Lambda}$
 T_3, T_4 , transverse polarization

T_1, T_2, T_5 : spin correlation

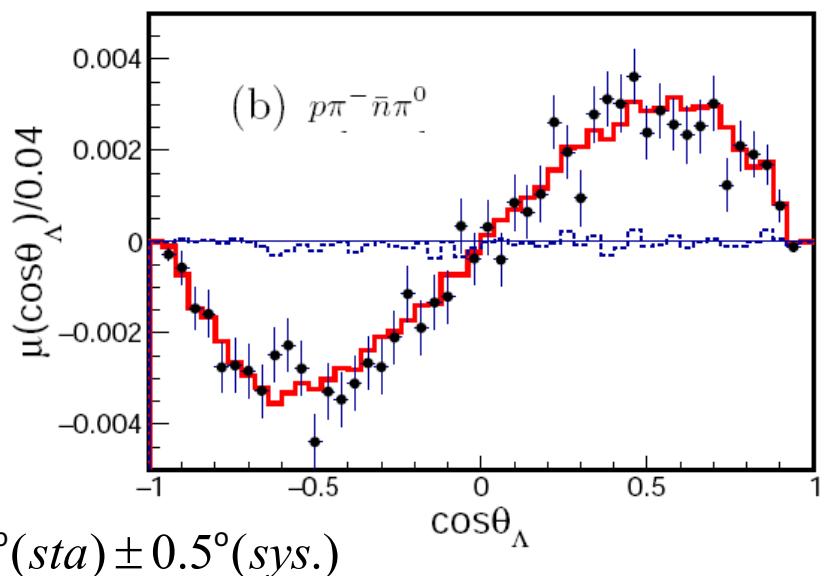
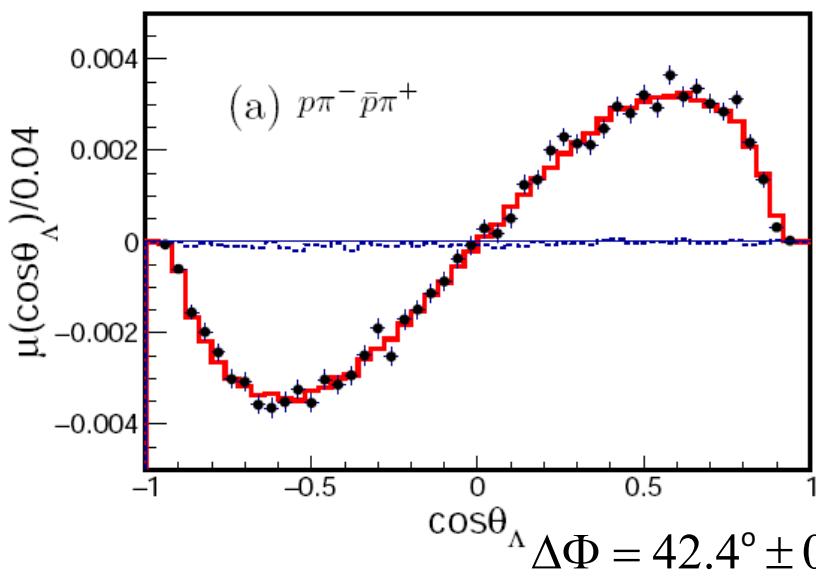


1.3 billion J/ψ events

simultaneous fit to the two data samples



Nature Physics (2019)

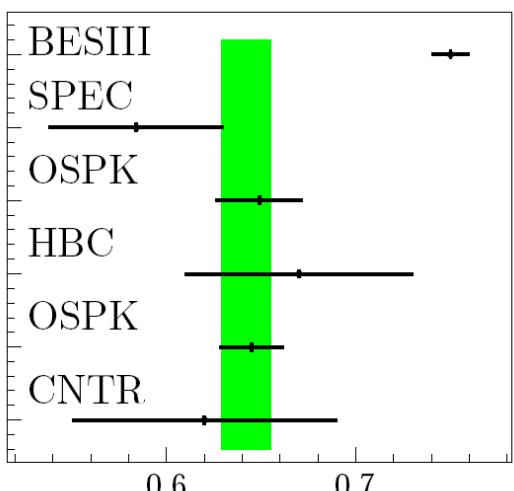


Fitting results and Λ decay asymmetries

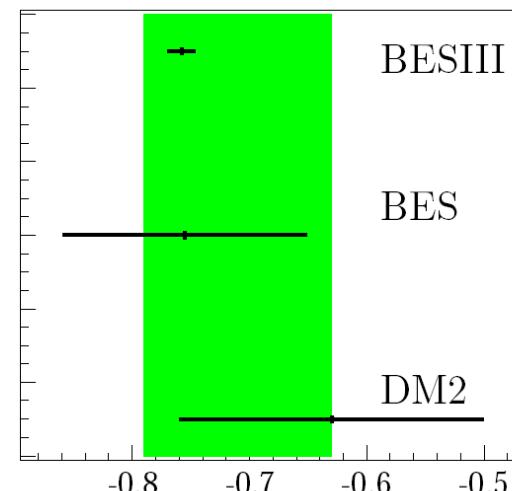
[Nature Physics \(2019\)](#)



Parameters	This measurement	Previous results
α_ψ	$0.461 \pm 0.006 \pm 0.007$	0.469 ± 0.027 [19]
α_-	$0.750 \pm 0.009 \pm 0.004$	0.642 ± 0.013 [8]
α_+	$-0.758 \pm 0.010 \pm 0.007$	-0.71 ± 0.08 [8]
$\bar{\alpha}_0$	$-0.692 \pm 0.016 \pm 0.006$	—
$\Delta\Phi$ (rad)	$0.740 \pm 0.010 \pm 0.008$	—
A_Λ	$-0.006 \pm 0.012 \pm 0.007$	0.006 ± 0.021 [8]
$\bar{\alpha}_0/\alpha_+$	$0.913 \pm 0.028 \pm 0.012$	—



(a) α_- for $\Lambda \rightarrow p\pi^-$



(b) α_+ for $\bar{\Lambda} \rightarrow \bar{p}\pi^+$

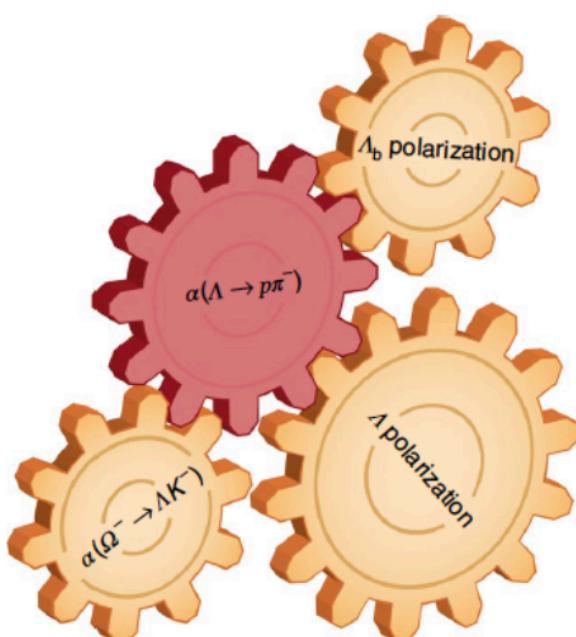
Consequences

PARTICLE PHYSICS

Anomalous asymmetry

A measurement based on quantum entanglement of the parameter describing the asymmetry of the Λ hyperon decay is inconsistent with the current world average. This shows that relying on previous measurements can be hazardous.

Ulrik Egede

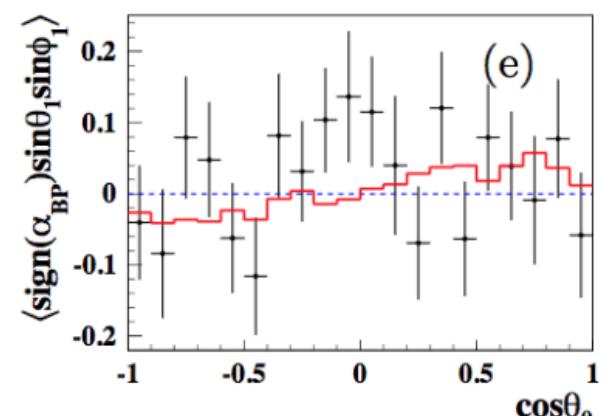
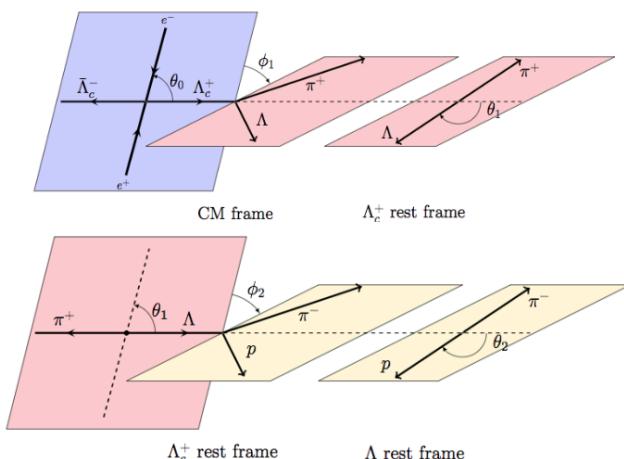


New input for many other measurements:

- 1) polarization
- 2) Asymmetry of the Λ_b and Λ_c
- 3) CPV in Λ_b and Λ_c decays
- 4) Decays of other charmed and beauty baryons

More analyses

- Update measurements using 10 billion J/ψ
 - A_Λ sensitivity: 5×10^{-3}
- $J/\psi, \psi(2S) \rightarrow \Sigma^+ \bar{\Sigma}^-, \Xi^0 \bar{\Xi}^0, \Xi^- \bar{\Xi}^+$
- $\psi(2S) \rightarrow \Omega^- \bar{\Omega}^+$
- Polarization of Λ_c^+ : $e^+ e^- \rightarrow \Lambda_c^+ \bar{\Lambda}_c^-, \Lambda_c^+ \rightarrow \Lambda \pi^+, \Sigma^+ \pi^0, \Sigma^0 \pi^+, K_S p$ @4.6GeV [arXiv:1905.04707]



$$\sin \Delta_0 = -0.28 \pm 0.13 \pm 0.03$$

$\Lambda_c^+ \rightarrow$		pK_S^0	$\Lambda\pi^+$	$\Sigma^+\pi^0$	$\Sigma^0\pi^+$
$\alpha_{BP}^{\Lambda_c^+}$	Predicted	-1.0 [16], 0.51 [11] -0.49 [10], -0.90 [10] -0.49 [17], -0.97 [18] -0.66 [19], -0.90 [30] -0.99 [20], -0.91 [31]	-0.70 [16], -0.67 [11] -0.95 [10], -0.99 [10] -0.96 [17], -0.95 [18] -0.99 [19], -0.86 [30] -0.99 [20], -0.94 [31]	0.71 [16], 0.92 [11] 0.79 [10], -0.49 [10] 0.83 [17], 0.43 [18] 0.39 [19], -0.76 [30] -0.31 [20], -0.47 [31]	0.70 [16], 0.92 [11] 0.78 [10], -0.49 [10] 0.83 [17], 0.43 [18] 0.39 [19], -0.76 [30] -0.31 [20], -0.47 [31]
	PDG [2] This work	$0.18 \pm 0.43 \pm 0.14$	-0.91 ± 0.15	-0.45 ± 0.32	$-0.73 \pm 0.17 \pm 0.07$

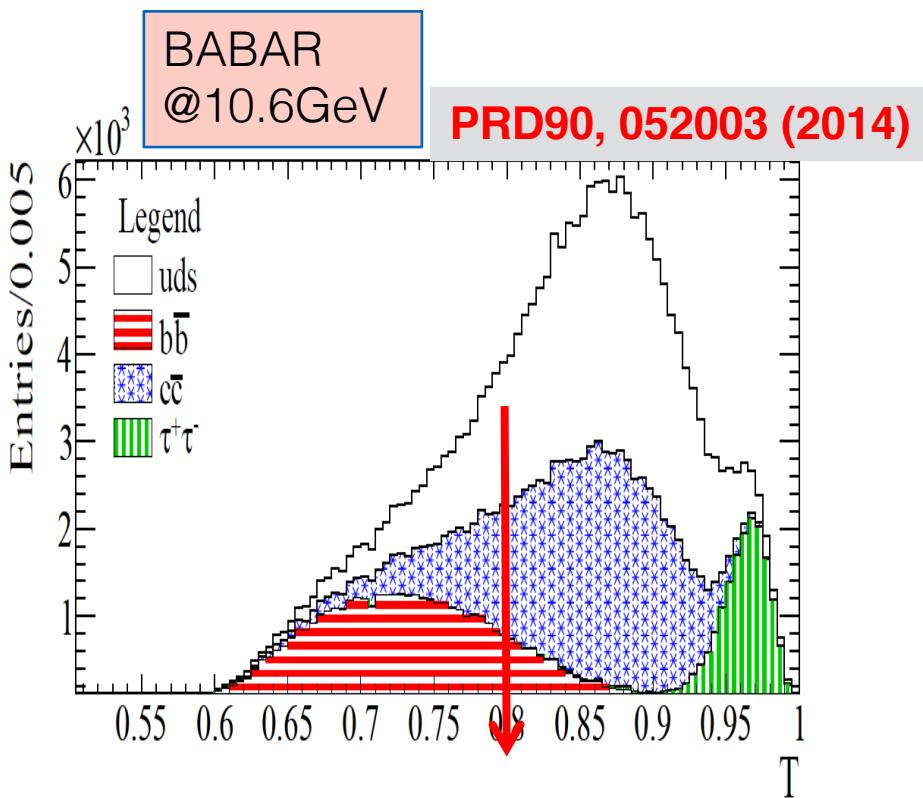
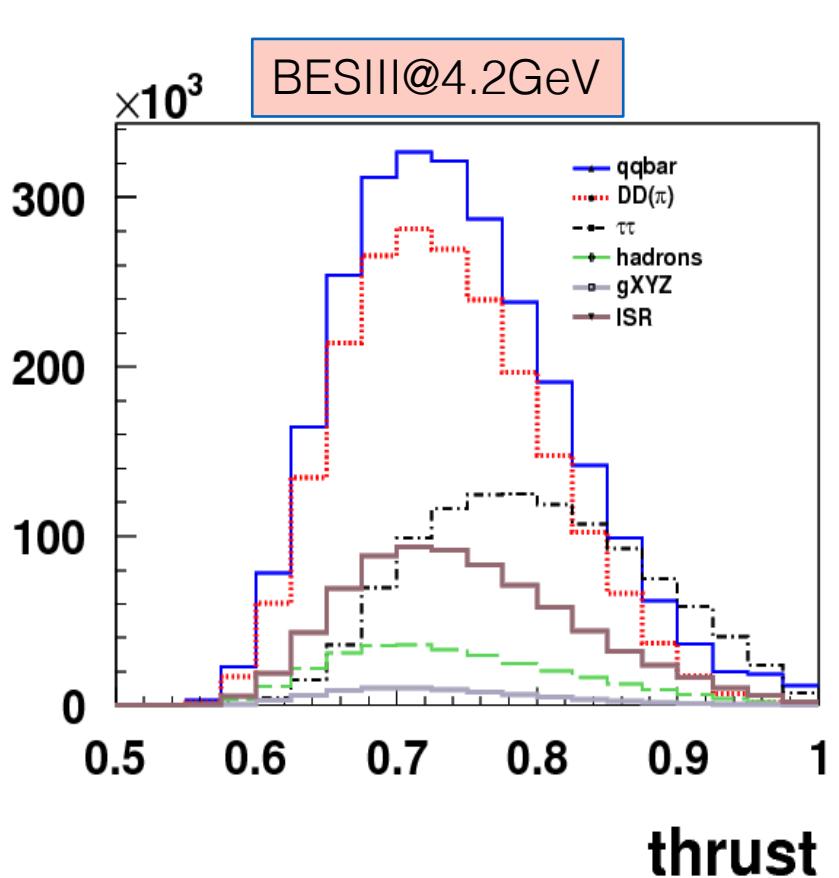
Summary and Outlook

- **Collins effects measurement has been implemented using BESIII data @3.65GeV.**
 - Nonzero Collins asymmetries were observed.
 - First measurement at medium energy, which is closer to SIDIS experiments
 - It allows to study the evolution of TMD objects
 - Further extensions include high *Lumi.* data and more types of final states
- **Hyperon Spin correlations in production**
 - Observation of Λ transverse polarization in $J/\psi \rightarrow \Lambda\bar{\Lambda}$ decays
 - New way to study Hyperon CPV
 - Already rewritten the PDG book for Λ decays
 - Future results using 10B J/ψ decays will improve the CPV search
 - Plan to update the results with more data set and do the same for Ξ/Σ

Thank you!

谢谢！！

Difficulties at higher energy



- ✓ Thrust value is used to suppress backgrounds at Belle/BaBar.
- ✓ However, it is not a good cut criterion at BESIII



- Misidentification of K and π : unfolding the measurement of $A^{\pi\pi}$ and $A^{K\pi}$
- Gluon radiation effect: subtracting normalized yields $R^U - R^{L(C)}$
- Higher harmonic terms: including in the fit function
- Possible charge-dependent acceptance effects: studying double ratio of positively over negatively charged pion pairs
- Beam polarization: studying the angular distribution of $e^+e^- \rightarrow \mu^+\mu^-$
- Several zero asymmetry tests

Most earlier measurement on α_-

- CNTR: $\pi^- + p \rightarrow \Lambda + K^0$

Phys.Rev. 129 (1963) 1795-1807

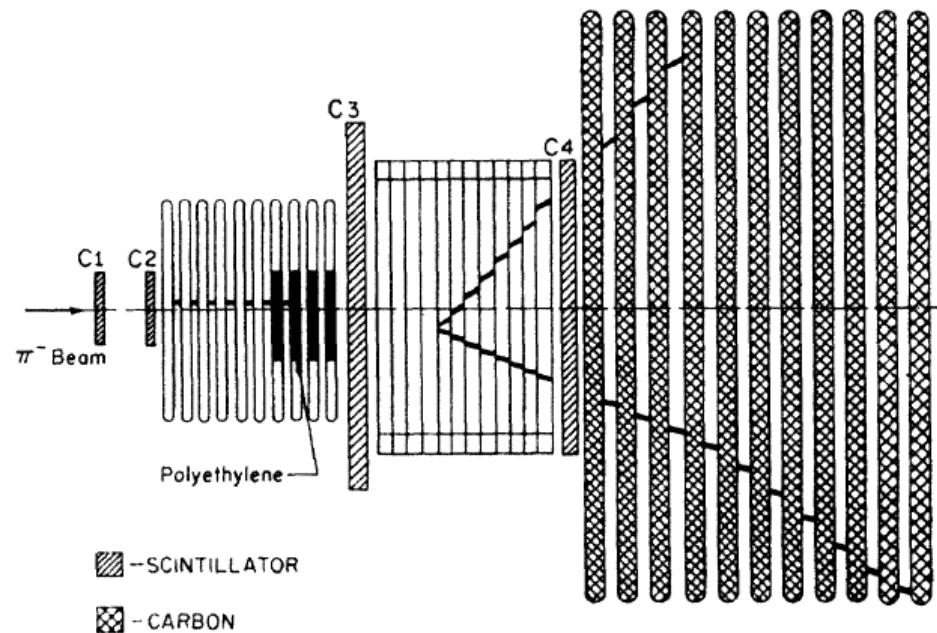
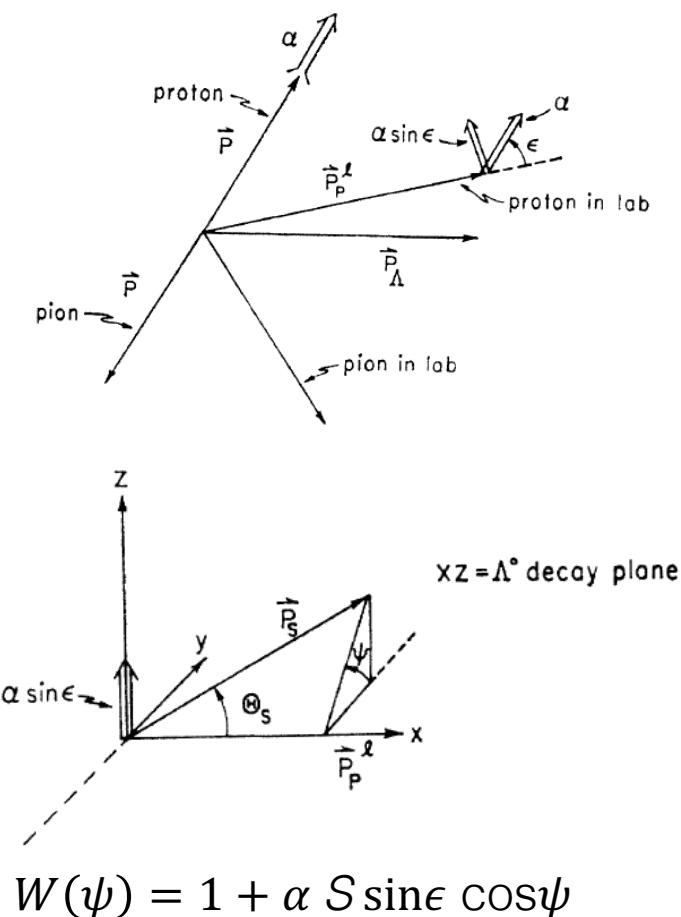


FIG. 1. Schematic diagram showing arrangement of apparatus.
An example of an event has been sketched in.

$$\alpha = -\frac{2}{\pi} \frac{1}{\langle S \rangle \langle \sin \epsilon \rangle} \frac{N_+ - N_-}{N_+ + N_-},$$

1156 events

$$\langle S \rangle = 0.565$$

$$\langle \sin \epsilon \rangle = 0.84, \\ \alpha = 0.62.$$



α_- FOR $\Lambda \rightarrow p\pi^-$

2019 PDG list

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
$0.750 \pm 0.009 \pm 0.004$	420k	ABLIKIM	2018AG	BES3 J/ψ to $\Lambda\bar{\Lambda}$
••• We do not use the following data for averages, fits, limits, etc. •••				
0.584 ± 0.046	8500	ASTBURY	1975	SPEC
0.649 ± 0.023	10325	CLELAND	1972	OSPK
0.67 ± 0.06	3520	DAUBER	1969	HBC From Ξ decay
0.645 ± 0.017	10130	OVERSETH	1967	OSPK Λ from $\pi^- p$
0.62 ± 0.07	1156	CRONIN	1963	Λ from $\pi^- p$

α_+ FOR $\bar{\Lambda} \rightarrow \bar{p}\pi^+$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
$-0.758 \pm 0.010 \pm 0.007$	420k	ABLIKIM	2018AG	BES3 J/ψ to $\Lambda\bar{\Lambda}$
••• We do not use the following data for averages, fits, limits, etc. •••				
-0.755 $\pm 0.083 \pm 0.063$	$\approx 8.7k$	ABLIKIM	2010	BES $J/\psi \rightarrow \Lambda\bar{\Lambda}$
-0.63 ± 0.13	770	TIXIER	1988	DM2 $J/\psi \rightarrow \Lambda\bar{\Lambda}$

Why hyperons using J/psi decays



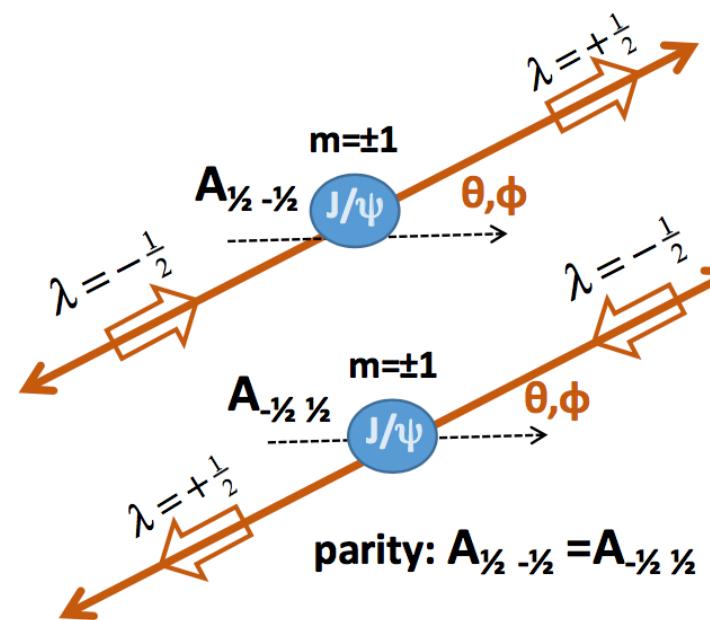
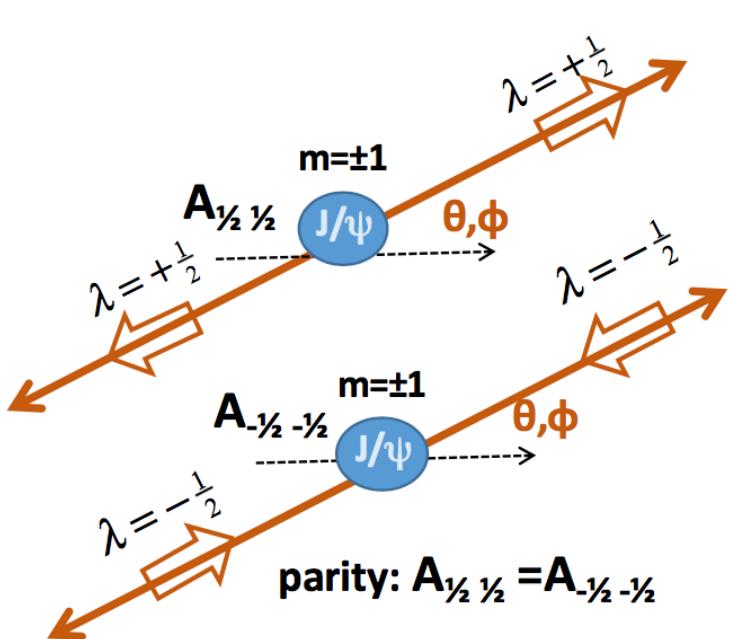
10 billion J/psi events collected

- Large BRs in J/ψ decays
- Quantum correlated pair productions
- Background free

Decay mode	$\mathcal{B} (\times 10^{-3})$	$N_B (\times 10^6)$	Efficiency	Number of reconstructed
$J/\psi \rightarrow \Lambda\bar{\Lambda}$	1.61 ± 0.15	16.1 ± 1.5	40%	3200×10^3
$J/\psi \rightarrow \Sigma^0 \bar{\Sigma}^0$	1.29 ± 0.09	12.9 ± 0.9	25%	600×10^3
$J/\psi \rightarrow \Sigma^+ \bar{\Sigma}^-$	1.50 ± 0.24	15.0 ± 2.4	24%	640×10^3
$J/\psi \rightarrow \Sigma(1385)^- \Sigma^+ \text{ (or c.c.)}$	0.31 ± 0.05	3.1 ± 0.5		
$J/\psi \rightarrow \Sigma(1385)^- \bar{\Sigma}(1385)^+ \text{ (or c.c.)}$	1.10 ± 0.12	11.0 ± 1.2		
$J/\psi \rightarrow \Xi^0 \bar{\Xi}^0$	1.20 ± 0.24	12.0 ± 2.4	14%	670×10^3
$J/\psi \rightarrow \Xi^- \bar{\Xi}^+$	0.86 ± 0.11	8.6 ± 1.0	19%	810×10^3
$J/\psi \rightarrow \Xi(1530)^0 \bar{\Xi}^0$	0.32 ± 0.14	3.2 ± 1.4		
$J/\psi \rightarrow \Xi(1530)^- \bar{\Xi}^+$	0.59 ± 0.15	5.9 ± 1.5		
$\psi(2S) \rightarrow \Omega^- \bar{\Omega}^+$	0.05 ± 0.01	0.15 ± 0.03		

$$e^+ e^- \rightarrow J/\psi \rightarrow \Lambda \bar{\Lambda}$$

Production: 2 independent helicity amplitudes: $A_{1/2, 1/2}, A_{1/2, -1/2}$



$\Delta = \text{complex phase between } A_{1/2, 1/2} \text{ and } A_{1/2, -1/2}$

$$\frac{d|\mathcal{M}|^2}{d \cos \theta} \propto (1 + \alpha_{J/\psi} \cos^2 \theta), \quad \text{with} \quad \alpha_{J/\psi} = \frac{|A_{1/2, -1/2}|^2 - 2|A_{1/2, 1/2}|^2}{|A_{1/2, -1/2}|^2 + 2|A_{1/2, 1/2}|^2}$$