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Mechanisms for transverse single-spin asymmetries

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Transverse SSA



$$A_N = \frac{\sigma^{\uparrow} - \sigma^{\downarrow}}{\sigma^{\uparrow} + \sigma^{\downarrow}}$$

Transverse SSA

$$A + B^\uparrow \rightarrow C + D + \dots + X$$

$$A_N = \frac{\sigma^\uparrow - \sigma^\downarrow}{\sigma^\uparrow + \sigma^\downarrow}$$



TMD factorization

SIDIS, DY, e^+e^-

Limitation:

small transverse final state momentum

$$q_T \ll Q$$

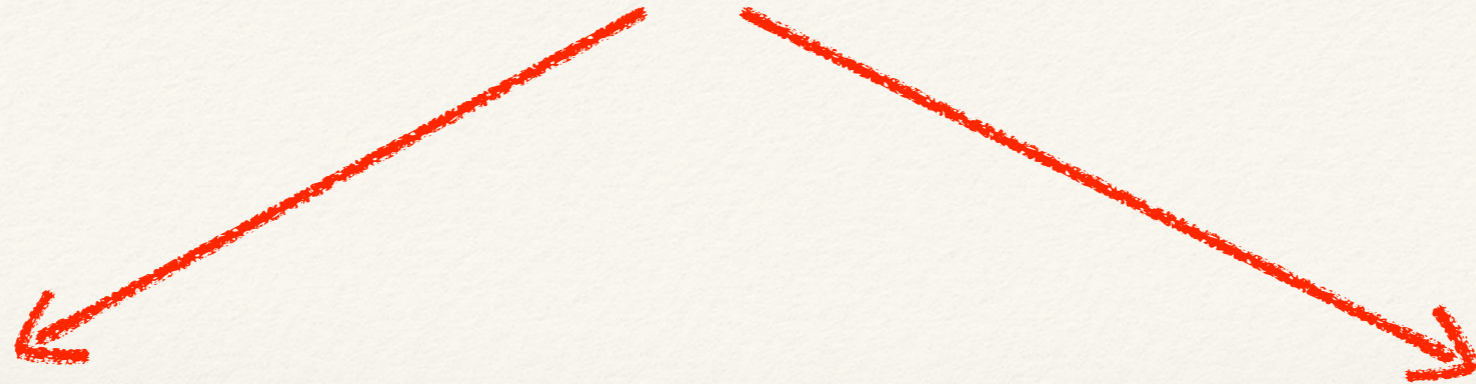
imaginary part needed

→ **T-odd TMDs** (e.g. Sivers function)

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collinear twist-3

single inclusive processes

$$A + B^\uparrow \rightarrow C + X$$

No kinematical limitation

Quark-Gluon-Quark correlations

imaginary part from hard part

→ **pole contributions**

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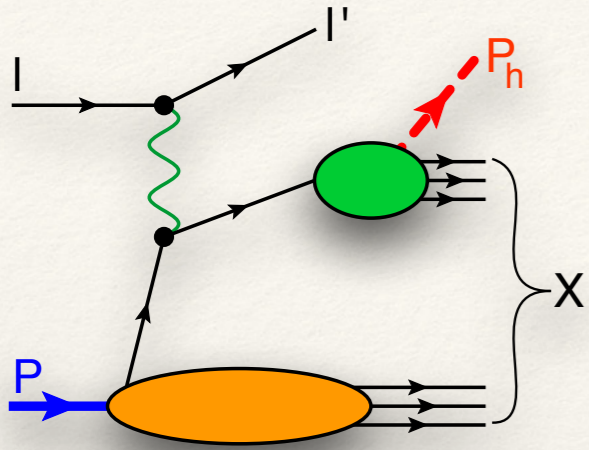
imaginary part from hard part

→ **pole contributions**

related

- Transverse SSA also possible for final state polarization → Λ - production
- Double spin asymmetries A_{LT} equally important!

TMD factorization: SIDIS $eN^\uparrow \rightarrow e\pi X$



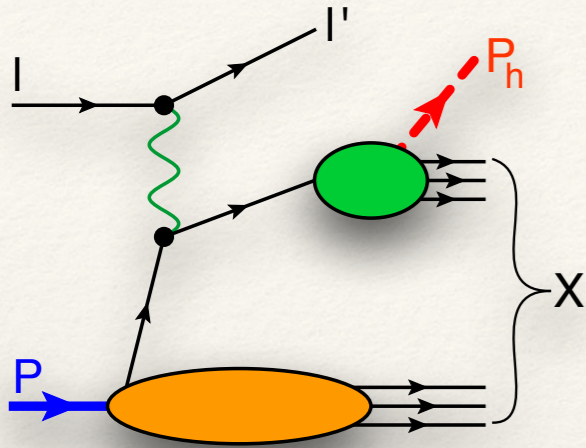
$$\frac{d\sigma^{\text{SIDIS}}}{d^2 P_{h\perp}}$$

SIDIS Spin Structure Functions:

[Bacchetta, Diehl, Goeke, Metz, Mulders, MS, JHEP (2007)]

$$\frac{d\sigma^{\text{SIDIS}}}{d^2 P_{h\perp}} \propto [F_{UU,T} + S_T \sin(\phi_h - \phi_s) F_{UT}^{\text{Siv}} + 16 \text{ S.F.}]$$

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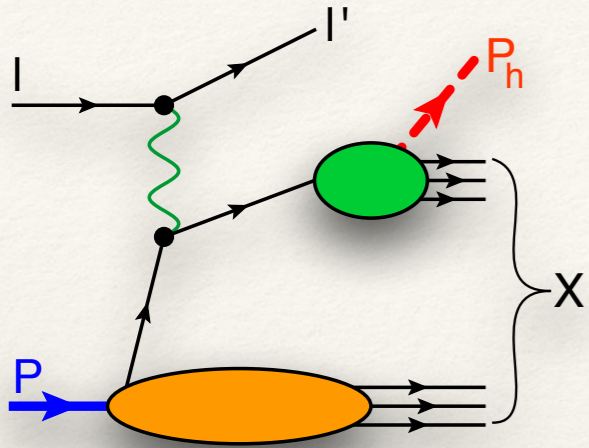
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P_{hT} small \rightarrow sensitive to partonic transverse momentum k_T
 \rightarrow **Transverse Momentum Dependent (TMD) distributions**

$$f(x, k_T^2)$$

$$F \propto |H|^2 \int d^2 k_T \int d^2 p_T \delta^{(2)}(k_T - p_T - P_{h\perp}/z) f(x, k_T^2) D(z, p_T^2) \equiv |H|^2 [f \otimes D]$$

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Sivers effect

$$A_{UT}^{\text{Siv}} = \frac{d\sigma^\uparrow - d\sigma^\downarrow}{d\sigma^\uparrow + d\sigma^\downarrow} = \frac{f_{1T}^\perp \otimes D_1}{f_1 \otimes D_1}$$

Sivers function

(Naive) definition of the Sivers function

$$f_1^q(x, k_T^2) \propto \int \frac{d\lambda d^2 z_T}{(2\pi)^3} e^{i\lambda x + i k_T \cdot z_T} \langle P | \bar{q}(0) \not{n} \mathcal{W} q(\lambda n + z_T) | P \rangle$$

Unpolarized TMD

$$(k_T \times S_T) f_{1T}^{\perp, q}(x, k_T^2) \propto \int \frac{d\lambda d^2 z_T}{(2\pi)^3} e^{i\lambda x + i k_T \cdot z_T} \langle P, S_T | \bar{q}(0) \not{n} \mathcal{W} q(\lambda n + z_T) | P, S_T \rangle$$

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Gauge link, Wilson line



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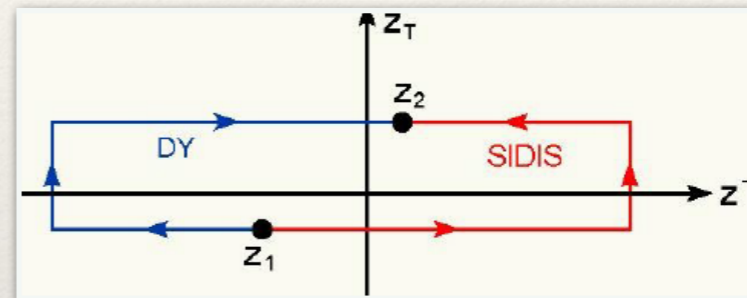
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Physics of the Wilson line

$$\mathcal{W}[a; b] = \mathcal{P} e^{-ig \int_a^b ds \cdot A(s)}$$



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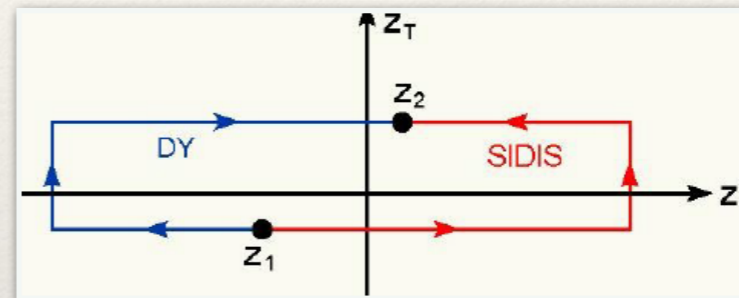
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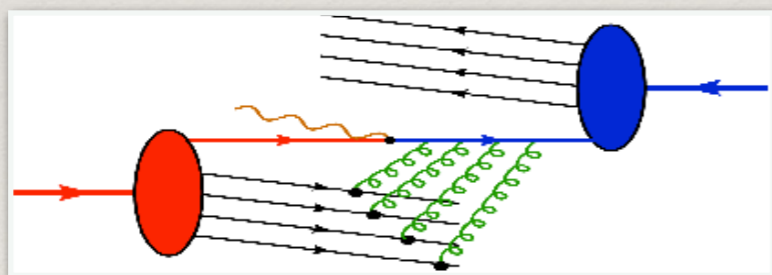
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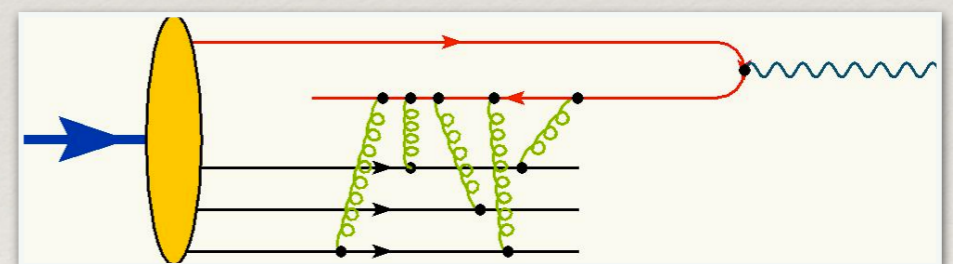
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Initial State Interactions: Drell-Yan



Final State Interactions: SIDIS



\mathcal{PT} - Transformation on the quark correlator \rightarrow ISI \Leftrightarrow FSI

\rightarrow sign switch of Sivers function "time-reversal (T)-odd"

$$f_{1T}^{\perp} \Big|_{DIS} = -f_{1T}^{\perp} \Big|_{DY}$$

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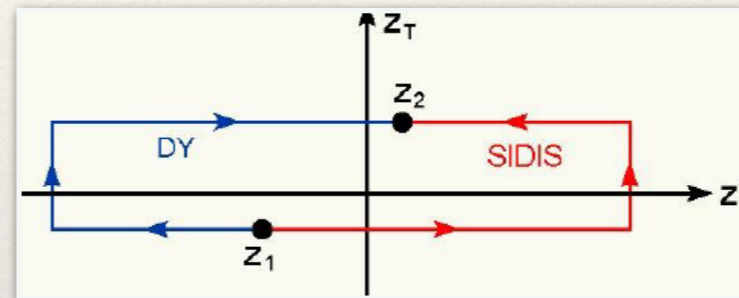
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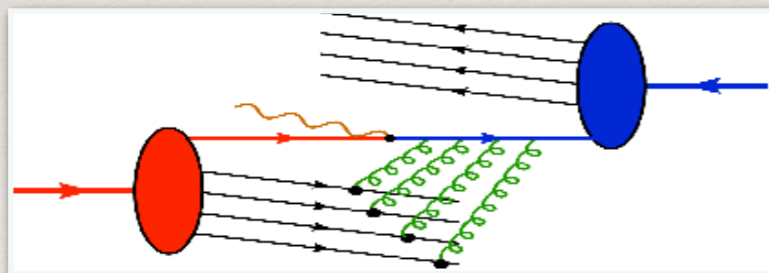
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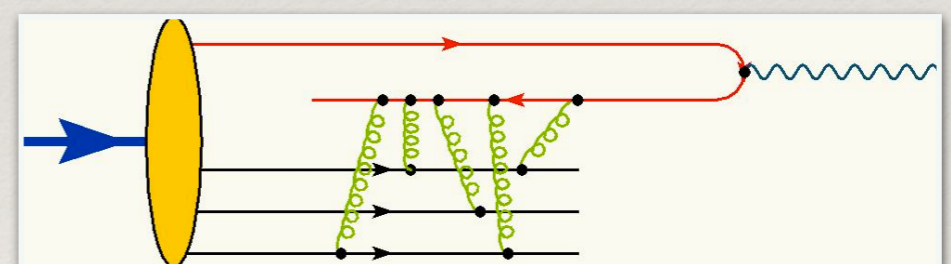
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\rightarrow important theoretical prediction to test TMD factorization, experimental leads that point in this direction RHIC, COMPASS \rightarrow talk by M. Radici

Proper TMD definition & Soft function

[Collins; Ji, Yuan; Aybat, Rogers; Echevarria, Idilbi, Scimemi; recent works by Bacchetta et al, Scimemi, Vladimirov et al]

Inclusion of Soft Function \implies

$$S(b_T) = \frac{1}{N_c} \text{Tr}_c \langle 0 | \mathcal{W}_{\bar{n}}^\dagger(-b_T/2) \mathcal{W}_{\bar{n}}(-b_T/2) \mathcal{W}_{\bar{n}}^\dagger(b_T/2) \mathcal{W}_{\bar{n}}(b_T/2) | 0 \rangle$$

$$f(x, b_T) \rightarrow \frac{f(x, b_T)}{\sqrt{S(b_T)}}$$

Modification needed in order to have:

- 1) Renormalizable Matrix Element
- 2) Finite matching coefficients at small b_T

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$$f_1^q(x, \vec{b}_T^2; \mu; \xi) = \sum_{q'} \left(\tilde{C}_{qq'} \otimes q(x) \right) \Big|_{\mu \propto 1/b_*} e^{S_{\text{pert}}(b_*)} \Big|_{\mu \propto 1/b_*} e^{g_q(x, b_T) + \frac{1}{2} g_K(b_T) \ln \frac{\xi}{\xi_0}}$$

TMD at large k_T ,
matching coeff NNLO

perturbative Sudakov factor, N³LO

non-perturbative input

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Evolved Sivers function

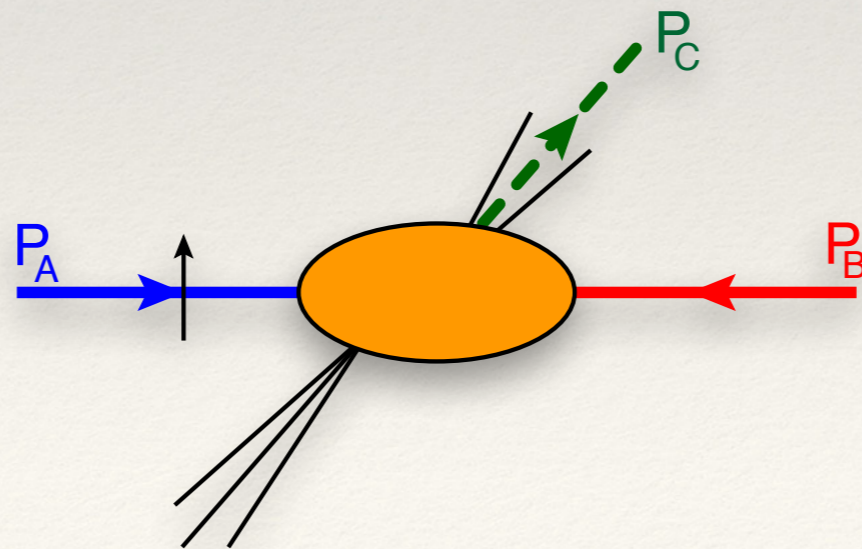
$$f_{1T}^{\perp, q}(x, b_T, ; \mu; \xi) \propto \left[\int dz dz' C_{1T}^\perp(z, z') F_{FT}^q(z, z') \right] e^{(S_{\text{pert}} + S_{\text{non-pert}})}$$

matching coeff NLO

quark-gluon-quark correlation function!

Collinear twist-3 in Single-Inclusive Hard Processes

$$P_A^\uparrow + P_B \rightarrow P_C + X$$



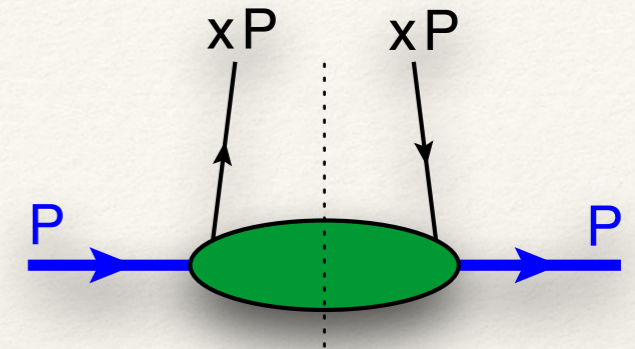
Collinear twist-3 formalism: several types of matrix elements

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intrinsic twist-3 PDF

$$g_T^q(x) = -\frac{1}{M} \int \frac{d\lambda}{4\pi} e^{i\lambda x} \langle P, S_T | \bar{q}(0) \not{\epsilon}_T \gamma_5 q(\lambda n) | P, S_T \rangle$$

- sensitive to 'bad quark field components',
- twist-3 characteristics hidden in Dirac structure
- generates the g_2 structure function in DIS, relevant for SSA in DIS (2γ)
- No probabilistic interpretation

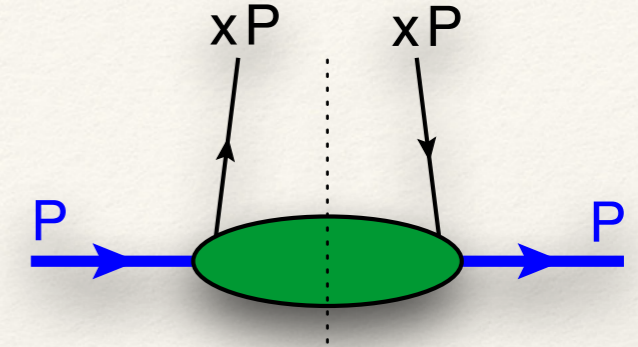


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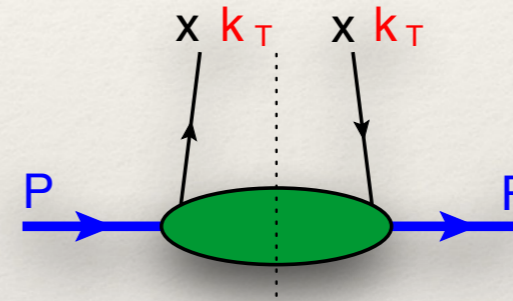
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kinematical twist-3 PDFs:

Small transverse quark/gluon momenta k_T :



$$(k_T \times S_T) f_{1T}^{\perp,q}(x, k_T^2) \propto \int \frac{d\lambda d^2 z_T}{(2\pi)^3} e^{i\lambda x + i k_T \cdot z_T} \langle P, S_T | \bar{q}(0) \not{n} \mathcal{W} q(\lambda n + z_T) | P, S_T \rangle$$

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'transhelicity'

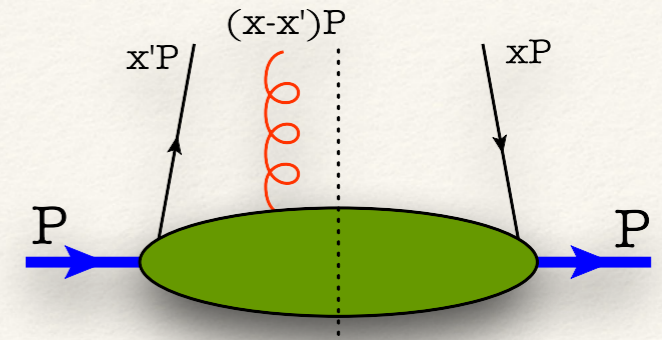
Collinear twist-3 formalism: TMD moments are needed

$$f_{1T}^{\perp,(1)}(x) = \int d^2 k_T \frac{k_T^2}{2M^2} f_{1T}^{\perp}(x, k_T^2)$$

$$g_{1T}^{(1)}(x) = \int d^2 k_T \frac{k_T^2}{2M^2} g_{1T}(x, k_T^2)$$

→ twist-3 characteristics through small transverse parton momentum k_T

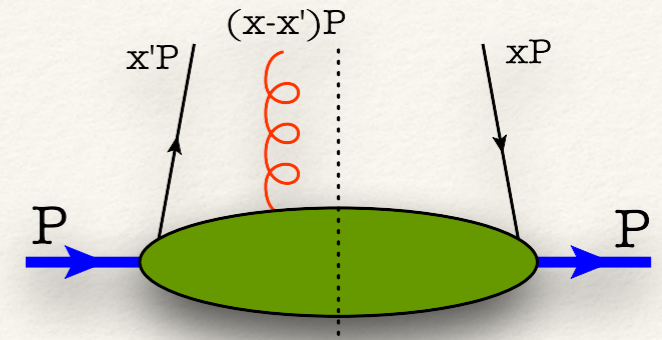
Dynamical twist-3: Quark - Gluon - Quark Correlations
(ETQS-matrix elements)



$$2M i\epsilon^{Pn\rho S} F_{FT}^q(x, x') = \int \frac{d\lambda}{2\pi} \int \frac{d\mu}{2\pi} e^{i\lambda x'} e^{i\mu(x-x')} \langle P, S_T | \bar{q}(0) \not{n} i g F^{n\rho}(\mu n) q(\lambda n) | P, S_T \rangle$$

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‘dynamical twist - 3’

→ 3 - parton correlator: suppression by additional propagator

→ Quark-Gluon-Quark correlation functions
drive x-dependence of TMDs like Sivers function, transhelicity, etc.

→ so far: only “diagonal support” $\pi F_{FT}(x, x) = f_{1T}^{\perp(1)}(x)$ constraint by SIDIS data

→ ‘integrated’ $F_{FT}(x, x')$: average transverse color Lorentz force on struck quark

[Burkardt, PRD88, 114502], [Aslan, Burkardt, M.S., 1904.03494]

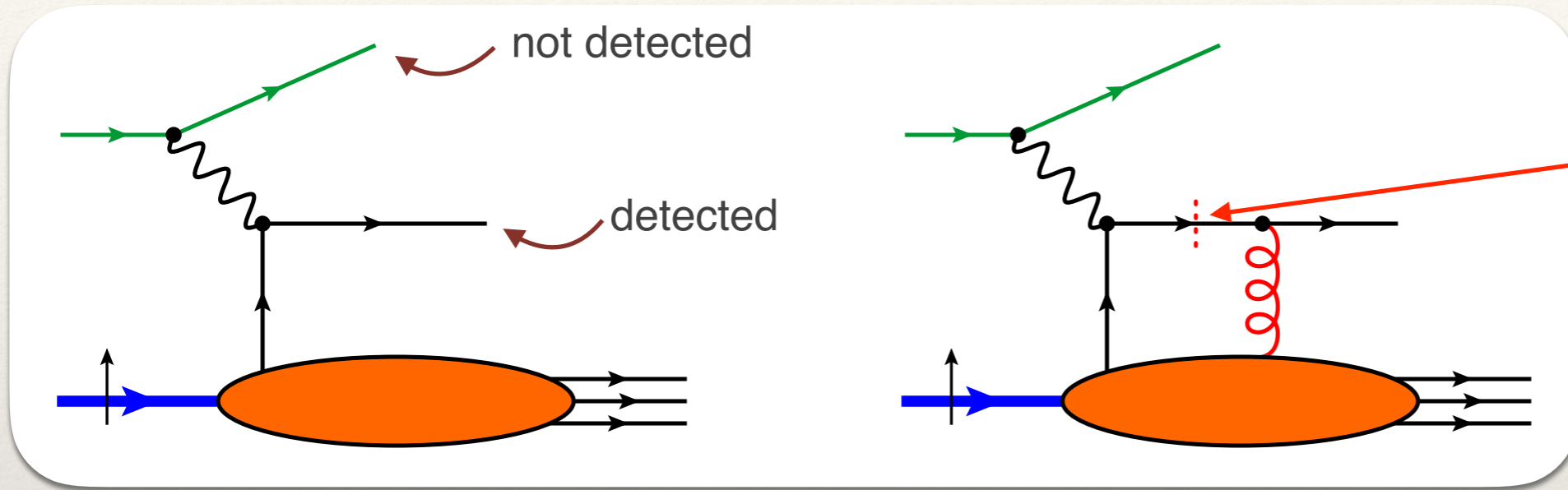
$$F^{n\rho} = [\vec{E} + \vec{n} \times \vec{B}]^\rho \propto \int dx \int dx' F_{FT}(x, x') \propto \int dx x^2 g_T(x)$$

How do quark-gluon correlations generate an SSA?

Example: Single-inclusive jet production $e N^\uparrow \rightarrow \text{jet } X$

[Gamberg, Kang, Metz, Pitonyak, Prokudin; Kanazawa, Koike, Metz, Pitonyak, MS]

Simple LO diagrams



Kinematical twist-3

Dynamical twist-3

Soft gluon pole

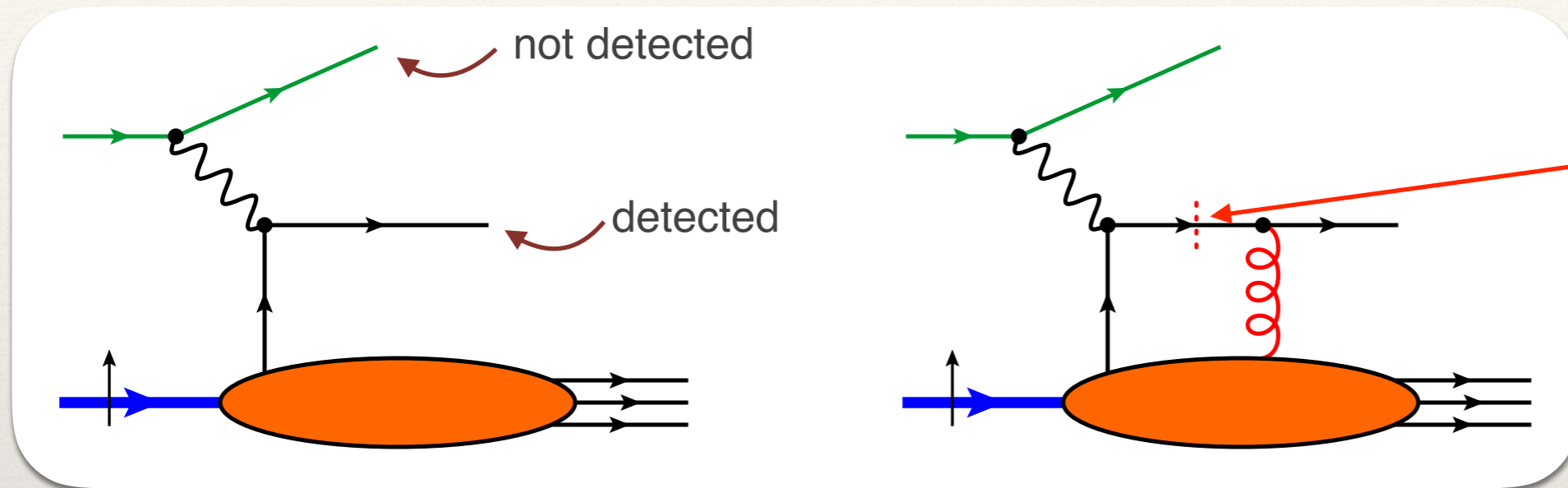
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Simple LO diagrams



Kinematical twist-3

Dynamical twist-3

Soft gluon pole

$$\frac{1}{x - x' + i\epsilon} = \frac{\mathcal{P}}{x - x'} - i\pi\delta(x - x')$$

$$A_N \propto \left(1 - x \frac{d}{dx} \right) F_{FT}^q(x, x)$$

SSA generated by soft-gluon pole only

Feasible at a future EIC, NLO corrections might be large

Photon SIDIS: $e + p \longrightarrow e + \gamma + X$

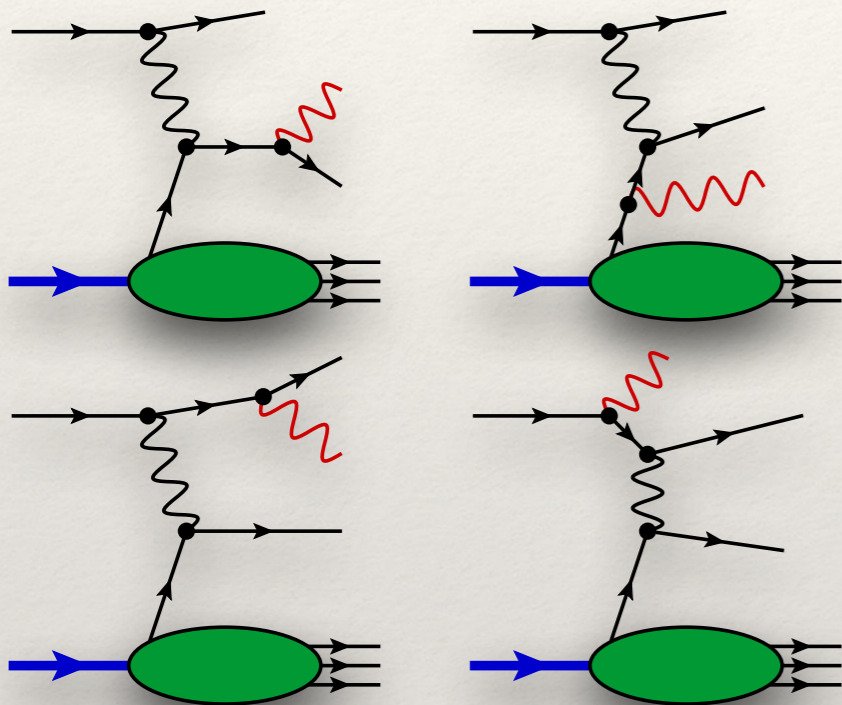
[Albaltan, Prokudin, M.S., in preparation]

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unpolarized cross section in the parton model

[Brodsky, Gunion, Jaffe, PRD 1972; see also works by Metz et al, Pisano, Mukherjee, Vogelsang, ...]



- avoid photon fragmentation: isolated photons
- collinear factorization: information on final quark is integrated out
- LO result:

$$E_\gamma E_e \frac{d\sigma_{UU}}{d^3 \vec{P}_\gamma d^3 \vec{P}_e} = \sum_q \left[e_q^2 \hat{\sigma}_2 + e_q^3 \hat{\sigma}_3 + e_q^4 \hat{\sigma}_4 \right] f_1^q(\bar{x})$$

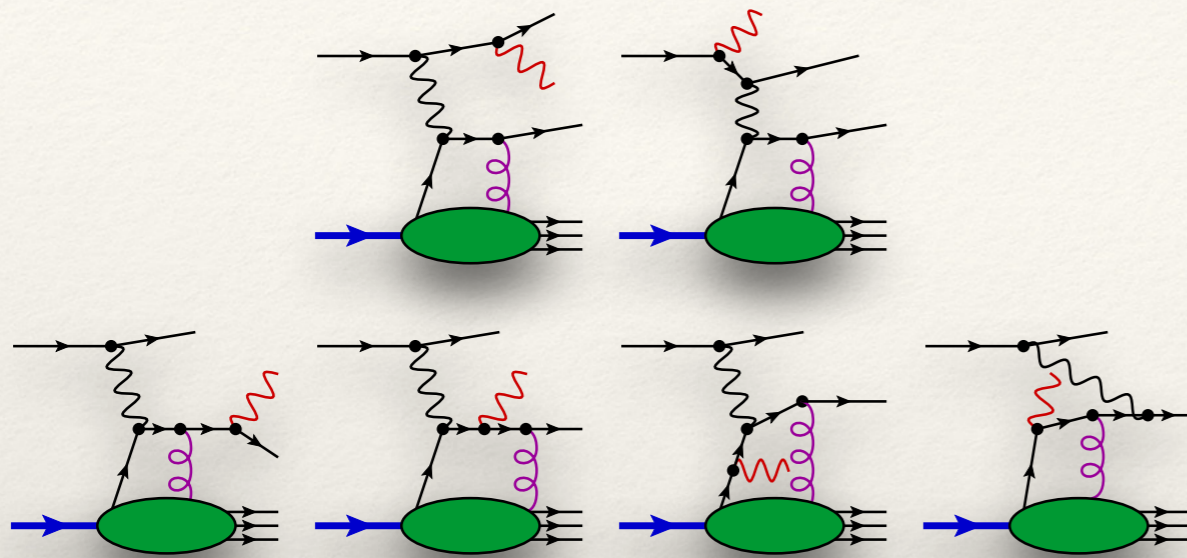
- two scales: $Q^2 = -(l - l' - P_\gamma)^2$ $\tilde{Q}^2 = -(l - l')^2$

- two 'Bjorken-x': $x_B = \frac{Q^2}{2P \cdot (l - l' - P_\gamma)}$ $\tilde{x}_B = \frac{\tilde{Q}^2}{2P \cdot (l - l')}$

BGJ - criterion for parton model dominance: $Q^2, \tilde{Q}^2, |Q^2 - \tilde{Q}^2| \gg M^2$

Transverse SSA in photon SIDIS

Include *intrinsic, kinematical & dynamical* twist - 3 contributions



At tree-level (LO):

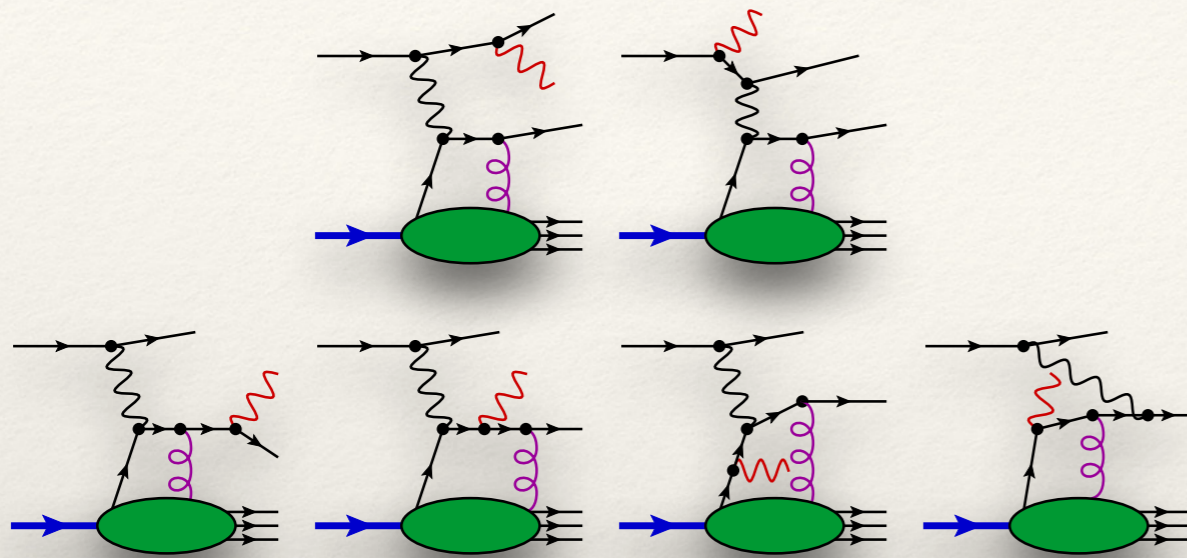
No contribution from g_T and $g_{1T}^{(1)}$
(no imaginary part)

Quark - Gluon correlations:

- 1) **Soft Gluon Poles:** $F_{FT}(x_B, x_B)$
- 2) **Soft Fermion Poles:** $F_{FT}(x_B, 0)$
- 3) **Hard Poles:** $F_{FT}(x_B, \tilde{x}_B)$

Transverse SSA in photon SIDIS

Include *intrinsic, kinematical & dynamical twist* - 3 contributions



At tree-level (LO):

No contribution from g_T and $g_{1T}^{(1)}$
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LO Result:

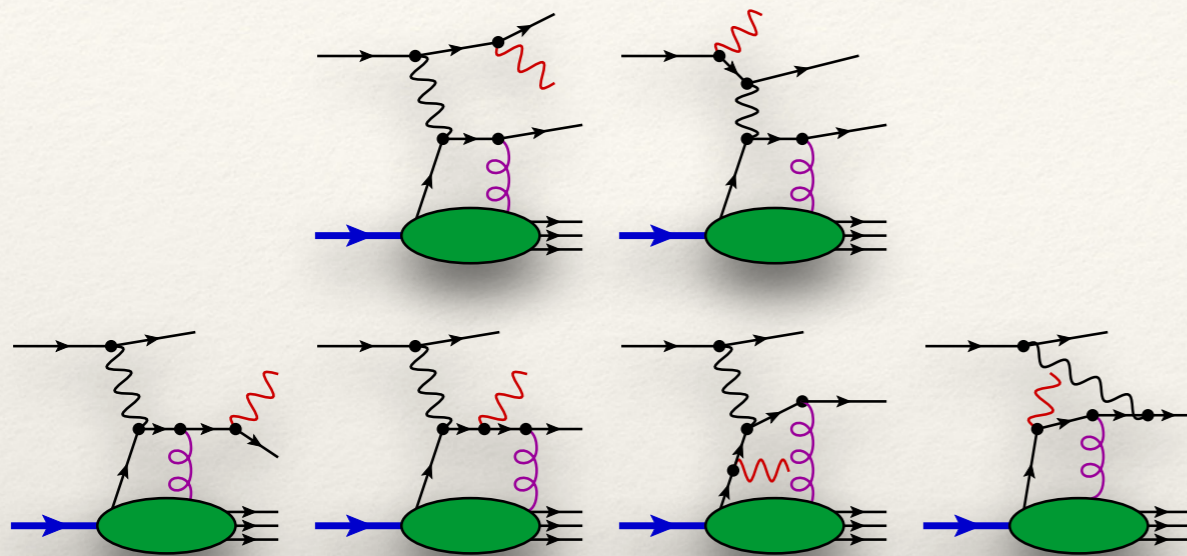
$$E_\gamma E' \frac{d\sigma_{UT}}{d^3\vec{P}_\gamma d^3\vec{P}_e} = \sum_q \left[\hat{\sigma}_{+,HP} F_{FT}^q(x_B, \tilde{x}_B) + \hat{\sigma}_{+,SFP} F_{FT}^q(x_B, 0) + \hat{\sigma}_{-,HP} G_{FT}^q(x_B, \tilde{x}_B) + \hat{\sigma}_{-,SFP} G_{FT}^q(x_B, 0) \right]$$

Soft Gluon Poles vanish !

Bethe-Heitler contribution vanishes
(Christ - Lee theorem)!

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\Rightarrow unique process to directly study “off-diagonal” support
of twist - 3 Quark - Gluon Correlation functions!

Summary

- ❖ Transverse Spin Polarization: Long history, measured in ep/pp-collisions, theoretical treatment more complicated
- ❖ TMD formalism: Sivers effect in SIDIS and DY: sign change
- ❖ Evolved Sivers function through Quark-Gluon correlations
- ❖ Photon SIDIS: May be able to scan the support of dynamical twist-3 functions point-by-point at LO.
- ❖ Experimental opportunity at EIC, COMPASS, JLab
 - input would help our understanding of quark-gluon correlation
 - valuable for evolution of qgq functions and TMDs.