#### PacificSpin2019

The 11<sup>th</sup> Circum-Pan-Pacific Symposium on High Energy Spin Physics

#### August 27 - 30, 2019

#### Miyazaki, Japan

**Topics** • Spin Structure of the Proton and Neutron

- Polarized Deep Inelastic Scattering, p-p Collision, Drell-Yan Process
- Generalized Parton
   Distributions
- Transverse Momentum
   Dependent Parton Distributions
- Perturbative QCD, Lattice QCD

Local Organizing Committee Co-Chairs: T. Iwata, T. Matsuda, T.-A. Shibata N. Doshita, Y. Goto, Y. Hatta, Y. Koike, S. Kumano, Y. Maeda, Y. Miyachi, K. Nagai, I. Nakagawa, K. Nakano, G. Nukazuka, S. Sasaki, S. Sawada, R. Seidl, K. Tanaka

https://sites.google.com/quark.kj.yamagata-u.ac.jp/pacspin2019 Contact: pacificspin2019@nucl.phys.titech.ac.jp Deadline for early registration and Deadline for abstract submission for contributed talks: June 27, 2019 Latest results in the extraction of the Sivers and Transversity distributions

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In collaboration with

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- F. Delcarro (JLab)



## the Sivers function

#### TMDs @twist=2



quark polarization

nucleon polarization		U	L	Т
	U	f1		h₁⊥
	L		<b>g</b> 1L	h₁∟⊥
	т	$\bm{f_{1T^{\bot}}}$	<b>g</b> 1t	$h_1 h_{1T^{\perp}}$

## the Sivers function

distortion of quark distribution in polarized proton P<sup> $\uparrow$ </sup> along direction  $\perp$  polarization ("spin-orbit")



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### Sivers function ↔ quark total J

Ji's sum rule  $J_z^q(Q^2) = \frac{1}{2} \int_0^1 dx \, x \left[ H^q(x,0,0;Q^2) + E^q(x,0,0;Q^2) \right]$ model lensing funct. L(x) + fit f<sub>1T</sub>  $\int dk_T f_{1T}^{\perp q}(x,k_T;Q_L^2) = -L(x) E^q(x,0,0;Q_L^2)$ 



### non-universality of Sivers function



# **Figure 1:** Definition of azimuthal angles for semi-inclusive deep inelastic scatters frame [25]. $P_{h\perp}$ and $S_{\perp}$ have the transverse parts of $P_h$ and S with respect to $P_h$ and S with respect to $P_h$ and S with respect to $P_h$ and $P_h$ and S with respect to $P_h$ and $P_h$

momentum

 $\frac{d\sigma}{dx\,dy\,dz\,d\phi_S\,d\phi_h d\mathbf{P}_{hT}^2} \underbrace{\operatorname{a}^{\alpha_v^2}\operatorname{cie}_{xyQ^2} \circ f(\operatorname{situal}^2) \circ f($  $\vec{S}_{T} \quad q = l - l' \text{ and its virtuality } Q^{2} = \overline{h} q^{2} \text{ We use the conventional variable}$   $x = \frac{Q^{2}}{2PF q} N, \gamma^{*} = \sum_{q} H^{q} (Q^{2}) \frac{P \cdot q}{P[w_{p}^{*} * f_{N}^{q} D_{1}^{q}]} \quad z = \frac{P \cdot q}{P}$ and write M and  $M_h$  for the respective masses of the  $M^2$  target badron h. We take the limit of large  $Q^2$  at fixed x, y, z, and through matching Y term neglect corrections in the masses of the hadrons or the lepton.  $P_{hT}$ It is convenient to  $\vec{d}_{s}^{[w,fD]} = \vec{d}_{s}^{x} \sum_{e^{2}} \int d^{2}\mathbf{k} d^{2}\mathbf{p} d^{2}\mathbf{p} d^{2}\mathbf{k} d^{2}\mathbf{p} d^{2}\mathbf{p} d^{2}\mathbf{p} d^{2}\mathbf{k} d^{2}\mathbf{p} d^{2}\mathbf{p}$ and q collinear. We define the transverse part  $P_{h\perp}^{\mu}$  of  $P_{h}^{\mu}$  as orthogonal v momenta P and q. Likewise, we define the transverse part  $S^{\mu}_{\perp}$  of the spin target, as well as its longitudinal projection  $S_{\parallel}$  along  $P^{\mu}$ . We further de angles  $\phi_h$  and  $\phi_S$  of  $P_h^{\mu}$  and  $S^{\mu}$  with respect to the lepton plane in ac Trento conventions [23], as shown in Fig. 1. Covariant expressions for discussed can be found in [14]. Finally, we write  $\lambda_e$  for the longitudinal incoming lepton, with  $\lambda_e = 1$  corresponding to a purely right-handed be

The lepton-hadron cross section can then be parameterized as [14]

# Figure 1: Definition of azimuthal angles for semi-inclusive deep inelastic scatter of $P_h$ and S with respect to $P_h$ and $P_h$ and

momentum

 $\vec{S}_{T} \quad q = l - l' \text{ and its virtuality } Q^{2} = \overline{h} q^{2} \text{ We use the conventional variable}$   $x = \frac{Q^{2}}{2PFq} \sum_{N,\gamma^{*}} -\frac{1}{P} \frac{P \cdot q}{P[\psi_{V}^{*}]} \int_{N}^{\infty} D_{1}^{q} = \sum_{n} \frac{P \cdot q}{P[\psi_{V}^{*}]} \int_{N}^{\infty} D_{1}^{q} = \frac{P \cdot q}{P}$ and write M and  $M_h$  for the respective masses of the  $M^{20}(M^{20})$  target **P**hT badron h. We take the limit of large  $Q^2$  at fixed x, y, z, and through matching Y term neglect corrections in the masses of the hadrons or the lepton. It is convenient to  $\vec{d}$  is  $c^{[w,fD]} = \vec{d}$  is  $c^{[w,fD]} = \vec{d}$  is  $c^{[w,fD]} = \vec{d}$  is  $c^{[w,fD]} = \vec{d}$  is  $c^{[w,fD]} = \vec{d}$ . The experimental observables for SIDIS is and q collinear. We define the transverse part  $P_{h\perp}^{\mu}$  of  $P_{h}^{\mu}$  as orthogonal with every index of  $P_{h\perp}^{\mu}$  and  $P_{h\perp}^{\mu}$  and  $P_{h\perp}^{\mu}$  and  $P_{h\perp}^{\mu}$  and  $P_{h\perp}^{\mu}$  and  $P_{h\perp}^{\mu}$  of the spin momenta P and q. Likewise, we define the transverse part  $S_{\perp}^{\mu}$  of the spin momenta  $P_{\mu}$  and  $P_{\mu}^{\mu}$ . Surviving target, as well as its longitudinal projection  $S_{\parallel}$  along  $P^{\mu}$ . We further de  $F_{UU,T}(x,z, \mathbf{P}_{hT}^{2}, Q^{2}) = 2\pi \sum_{a} \frac{e^{2} x}{a n g l g g} \phi_{b}^{d\xi_{T}} \tilde{\xi}_{T}^{d} J_{\phi}(\xi_{T}) \mathbf{P}_{h}^{p\mu} \tilde{f}_{a}^{d} S^{\mu} \tilde{\xi}_{T}^{d} \tilde{f}_{a}^{d} \tilde{f}_{a}^{d$  $F_{UT,T}^{\sin(\phi_h - \phi_S)}(x, z, \mathbf{P}_{hT}^2, Q^2) = -2\pi M_{UT,T}^{\infty} e^2 x_{dh} be^{\xi} f \xi^2 J_{d}(\xi_T | \mathbf{P}_{hT} | \mathbf{P}_{hT}$ incoming lepton, with  $\lambda_e = 1$  corresponding to a purely right-handed be The lepton-hadron cross section can then be parameterized as [14]

### the Sivers Single-Spin Asymmetry



First extraction of Sivers function using unpolarized TMDs  $f_1$  and  $D_1$  extracted from global fit of (SIDIS + Drell-Yan + Z-boson) data

Bacchetta, Delcarro, Pisano, Radici, in preparation

At LO 
$$\tilde{f}_1^q(x,\xi_T^2;Q^2) = f_1^q(x,\mu_\xi^2) e^{S(\mu_\xi^2,Q^2)} e^{g_K(\xi_T) \log(Q^2/Q_0^2)} \tilde{f}_{1NP}^q(x,\xi_T^2)$$
 and similar for D<sub>1</sub>  
Collins, "Foundations of Perturbative QCD" (11) perturbative DSEHS15@NLO

 $\begin{aligned} f_1 &= \text{collinear PDF} & \text{GJR08FFnloE} \\ \text{S} &= \text{Sudakov form factor,} \\ & \text{at NLL includes A}_1, \text{A}_2, \text{B}_1 \end{aligned}$ 

At LO 
$$\tilde{f}_{1}^{q}(x,\xi_{T}^{2};Q^{2}) = f_{1}^{q}(x,\mu_{\xi}^{2}) e^{S(\mu_{\xi}^{2},Q^{2})} e^{g_{K}(\xi_{T}) \log(Q^{2}/Q_{0}^{2})} \tilde{f}_{1NP}^{q}(x,\xi_{T}^{2})$$
 and similar for D<sub>1</sub>  
Collins, "Foundations of perturbative perturbative perturbative  $QCD''(11)$  perturbative  $g_{K} = -g_{2}\xi_{T}^{2}/4$   
f\_{1} = collinear PDF GJR08FFnloE  
S = Sudakov form factor,  
at NLL includes A\_{1}, A\_{2}, B\_{1}
 $\tilde{f}_{1NP} = F.T.\left(\frac{1}{\pi}\frac{1+\lambda k_{T}^{2}}{g_{1}+\lambda g_{1}^{2}}e^{-k_{T}^{2}/g_{1}}\right)$   
 $\tilde{D}_{1NP} = F.T.\left(\frac{1}{\pi}\frac{e^{-P_{T}^{2}/g_{3}}+\lambda_{F}P_{T}^{2}/z^{2}e^{-P_{T}^{2}/g_{4}}}{g_{3}+\lambda_{F}/z^{2}g_{4}^{2}}\right)$   
 $g_{1}(x) = g_{1}(0.1)\frac{(1-x)^{\alpha}x^{\sigma}}{(1-0.1)^{\alpha}(0.1^{\sigma})}$   
 $g_{3/4}(z) = g_{3/4}(0.5)\frac{(1-z)^{\gamma}(z^{\beta}+\delta)}{(1-0.5)^{\gamma}(0.5^{\beta}+\delta)}$ 





### (x, Q<sup>2</sup>) - coverage



Adolph et al., E.P.J. C73 (13)

Airapetian et al., P.R. D87 (13)

### first unpolarized TMD global fit



### Nucleon tomography in momentum space



We have extracted the unpolarized transverse momentum dependent parton distribution function (TMDPDF) and rapidity anomalous dimension (also known as Collins-Soper kernel) from Drell-Yan data. The analysis has been performed in the  $\zeta$ -prescription with NNLO perturbative inputs. We have also provided an estimation of the errors on the extracted functions with the replica method. The values of TMDPDF and rapidity anomalous dimension, together with the code that evaluates the cross-section, are available at [45], as a part of the artemide package. We plan to release grids for TMDPDFs extracted in this work also through the TMDlib [69].



### parametrization of Sivers $f_{1T}$

scale Q<sub>0</sub>  $\tilde{f}_{1T}^{\perp q}(x,\xi_T^2;Q_0^2) = f_{1T}^{\perp(1)q}(x,Q_0^2) \tilde{f}_{1TNP}^{\perp}(x,\xi_T^2)$ 

$$\left(\frac{1}{N_T}\frac{1}{\pi}\frac{1+\lambda_S k_T^2}{M_1^2+\lambda_S M_1^4}e^{-k_T^2/M_1^2} f_{1NP}(x,k_T^2)\right)$$

normalized flavor-independent double Gaussian on top of f<sub>1NP</sub>

### parametrization of Sivers $f_{1T} \perp$

scale Q<sub>0</sub>  $\tilde{f}_{1T}^{\perp q}(x, \xi_T^2; Q_0^2) = f_{1T}^{\perp (1)q}(x, Q_0^2) \tilde{f}_{1TNP}^{\perp}(x, \xi_T^2)$ normalization  $f_{1T}^{\perp (1)q} = \frac{1}{N_x^q} f_{Siv}^q(x) \quad f_1^q(x, Q_0^2)$   $f_{1T}^{\perp (1)q} = \frac{1}{N_x^q} f_{Siv}^q(x) \quad f_1^q(x, Q_0^2)$   $f_{1T}^{\perp (1)q} = x^{\alpha_q} (1-x)^{\beta_q} [1 + A_q T_1(x) + B_q T_2(x)]$ Chebyshev polynomials

# parametrization of Sivers $f_{1T} \perp$

scale Q<sub>0</sub> 
$$\tilde{f}_{1T}^{\perp q}(x, \xi_T^2; Q_0^2) = f_{1T}^{\perp (1)q}(x, Q_0^2) \tilde{f}_{1TNP}^{\perp}(x, \xi_T^2)$$
  
normalization  
 $f_{1T}^{\perp (1)q} = \frac{1}{N_x^q} f_{Siv}^q(x) \quad f_1^q(x, Q_0^2)$   
 $f_{Siv}^q(x) = x^{\alpha_q} (1-x)^{\beta_q} [1 + A_q T_1(x) + B_q T_2(x)]$   
 $f_{Siv}^q(x) = x^{\alpha_q} (1-x)^{\beta_q} [1 + A_q T_1(x) + B_q T_2(x)]$   
further multiply by  $\frac{1}{N_{Tmax}} \frac{1}{N_{xmax}}$   
 $\begin{cases} N_{Tmax} = \max_{k_x^2} \left[ \frac{k_T}{M} \frac{1}{N_T} \frac{1}{\pi} \frac{1 + \lambda_S k_T^2}{M_1^2 + \lambda_S M_1^4} e^{-k_T^2/M_1^2} \right]$   
 $N_{xmax} = \max_{x,q} \left[ \frac{1}{N_x^q} f_{Siv}^q(x) \right]$   
to grant positivity  $\left( f_{1T}^{\perp (1)}(x, k_T^2) \right)^2 \le \frac{k_T^2}{4M^2} (f_1(x, k_T^2))^2$ 

### parametrization of Sivers $f_{1T} \perp$

scale Q<sub>0</sub> 
$$\tilde{f}_{1T}^{\perp q}(x, \xi_T^2; Q_0^2) = f_{1T}^{\perp (1)q}(x, Q_0^2) \tilde{f}_{1TNP}^{\perp}(x, \xi_T^2)$$
  
normalization  
 $f_{1T}^{\perp (1)q} = \frac{1}{N_x^q} f_{Siv}^q(x) \quad f_1^q(x, Q_0^2)$   
 $f_{1T}^{\perp (1)q} = \frac{1}{N_x^q} f_{Siv}^q(x) \quad f_1^q(x, Q_0^2)$   
 $f_{Siv}^{q}(x) = x^{\alpha_q} (1-x)^{\beta_q} [1 + A_q T_1(x) + B_q T_2(x)]$   
 $f_{Siv}^q(x) = x^{\alpha_q} (1-x)^{\beta_q} [1 + A_q T_1(x) + B_q T_2(x)]$   
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 $f_{Siv}^q(x) = x^{\alpha_q} (1-x)^{\beta_q} [1 + A_q T_1(x) + B_q T_2(x)]$   
 $f_{Siv}^{N_{Tmax}} = \max_{k_T^q} \left[ \frac{k_T}{M} \frac{1}{N_T} \frac{1}{\pi} \frac{1 + \lambda_S k_T^2}{M_1^2 + \lambda_S M_1^4} e^{-k_T^2/M_1^2} \right]$   
 $f_{Nxmax} = \max_{x,q} \left[ \frac{k_T}{M} \frac{1}{N_T} \frac{1}{\pi} \frac{1 + \lambda_S k_T^2}{M_1^2 + \lambda_S M_1^4} e^{-k_T^2/M_1^2} \right]$   
 $f_{Siv}^{N_Tmax} = \max_{x,q} \left[ \frac{1}{N_T^q} f_{Siv}^q(x) \right]$   
 $f_{Siv}^{N_Tmax} = \max_{x,q} \left[ \frac{1}{N_T^q} f_{Siv}^q(x) \right]$ 

2 + 3x4 = 14 parameters

### evolution of Sivers $f_{1T} \perp$



approximate evolution of T<sub>F</sub> as DGLAP evolution of f<sub>1</sub>

### data used in the fit



### sample of fit



#### results



#### results



### tomography of transversely polarized proton



# the Transversity function







1<sup>st</sup> Mellin moment (tensor charge) not directly accessible in  $\mathcal{L}_{SM}$  $\rightarrow$  low-energy footprint of BSM physics at higher scale ?

h<sub>1</sub> from first global fit of SIDIS + p-p data

Radici and Bacchetta, P.R.L. **120** (18) 192001

## the di-hadron mechanism

transversity is chiral-odd  $\rightarrow$  need a chiral-odd partner

- the di-hadron mechanism: IFF  $H_1^{4}$ 

2-hadrons semi-inclusive production



collinear framework

- $h_1$  probed as PDF
- factorization theorems for all hard processes  $\rightarrow$  universality of h<sub>1</sub> H<sub>1</sub>< mechanism

# advantages of di-hadron mechanism

#### factorization theorems for all hard processes



#### data used in the global fit

Airapetian et al., JHEP **0806** (08) 017 Adolph et al., P.L. **B713** (12) Braun et al., E.P.J. Web Conf. **85** (15)



Vossen et al., P.R.L. 107 (11) 072004

run 2006 (s=200 GeV<sup>2</sup>)

Adamczyk et al. (STAR), P.R.L. **115** (2015) 242501

run 2011 (s=500 GeV<sup>2</sup>)

Adamczyk et al. (STAR), P.L. **B780** (18) 332

# the phase space



- limited to mostly medium/high x
- guess low-x behavior (relevant for calculation of tensor charge see later)

### currently, only LO analysis



#### access only $q-\overline{q} = q_v$ , q=u,dvalence flavors in SIDIS A<sub>UT</sub>

### theoretical uncertainties

#### unpolarized Di-hadron Fragmentation Function D1

- quark D<sub>1</sub>q is well constrained by  $e^+e^- \rightarrow (\pi^+\pi^-) X$  (Montecarlo)
- **gluon**  $D_1^g$  is **not** constrained by  $e^+e^- \rightarrow (\pi^+\pi^-) X$  (currently, LO analysis)
- **no data** available yet for  $p p \rightarrow (\pi^+\pi^-) X$



### statistical uncertainty

#### the bootstrap method

- shift each exp. point by Gaussian noise within exp. variance
- create sets of virtual points to be fitted: 50, 100, 200 sets... until average and standard deviation reproduce original exp. points (here, 200x3=600)
- exclude largest and smallest 5% => 90% band



automatically accounts for correlations

### choice of functional form

functional form whose Mellin transform can be computed analytically and complying with Soffer Bound at any x and scale Q<sup>2</sup>

$$h_1^{q_v}(x;Q_0^2) = F^{q_v}(x) \begin{bmatrix} SB^q(x) + \overline{SB}^{\overline{q}}(x) \end{bmatrix}$$

$$\begin{array}{c} & & \\ & & \\ & & \\ & & \\ Soffer Bound \\ 2|h_1^q(x,Q^2)| \le 2 SB^q(x,Q^2) = |f_1^q(x,Q^2) + g_1^q(x,Q^2)| \\ & & \\ & & \\ MSTW08 \quad DSSV \end{array}$$

$$(x) = \frac{N_{q_v}}{\max_x[|F^{q_v}(x)|]} x^{A_{q_v}} \left[1 + B_{q_v} \operatorname{Ceb}_1(x) + C_{q_v} \operatorname{Ceb}_2(x) + D_{q_v} \operatorname{Ceb}_3(x)\right] \\ & & \\ \operatorname{Ceb}_n(x) \text{ Cebyshev polynomial} \end{array}$$

10 fitting parameters

constrain parameters

 $F^{q_v}$ 

 $|N_{q_v}| \le 1 \Rightarrow |F^{q_v}(x)| \le 1$  Soffer Bound ok at any Q<sup>2</sup>

### constrain parameters : low-x trend

$$\lim_{x \to 0} x SB^{q}(x) \propto x^{a_{q}}$$

$$\lim_{x \to 0} F^{q_{v}}(x) \propto x^{A_{q}}$$

$$h_{1}^{q}(x) \stackrel{x \to 0}{\approx} x^{A_{q}} + a_{q} - 1$$

$$\text{tensor charge} \quad \delta q(Q^{2}) = \int_{x_{\min}}^{1} dx h_{1}^{q-q}(x, Q^{2})$$

$$\text{low-x behavior important}$$

$$\text{constrain parameters} \quad \text{our choice}$$

$$\delta q \quad \text{finite} = A_{q} + a_{q} > 0 \quad A_{q} + a_{q} > \frac{1}{3} \quad \left| \int_{0}^{x_{\min}} dx \right| \sim 1\% \text{ of } \left| \int_{x_{\min}}^{1} dx \right|$$

$$\text{for } x_{\min} = 10^{-6} \text{ from MSTW08}$$

$$\text{Other choices}$$

$$\stackrel{\text{``massive'' jet in DIS \rightarrow h_{1} \text{ at twist 3}}{\text{violation of Burkardt-Cottingham s.r.}} \int_{0}^{1} dx g_{2}(x) \propto \int_{0}^{1} dx \frac{h_{1}(x)}{x} \rightarrow A_{q} + a_{q} > 1$$

small-x dipole picture  $h_1^{q_v}(x) \stackrel{x \to 0}{\approx} x^{1-2\sqrt{\frac{\alpha_s(Q^2)N_c}{2\pi}}} \longrightarrow \text{at } Q_0 \quad A_q + a_q \sim 1$ Kovchegov & Sievert, arXiv:1808.10354

Accardi and Bacchetta, P.L. B773 (17) 632

### Results



# Our first global fit

first ever extraction of transversity from data of SIDIS and proton-proton collisions

Radici and Bacchetta, P.R.L. **120** (18) 192001



### the extracted transversity



### **Comparison with other extractions**



### sensitivity to th. uncertainty



### sensitivity to th. uncertainty



**p-p:** u~d, gluon @LO but **SIDIS:** u~(8x)d, gluon @NLO

need data from target more sensitive to down (deuteron, <sup>3</sup>He) and need data from multiplicities in p+p  $\rightarrow$  ( $\pi\pi$ )+X

## The tensor "charge" of the proton

1<sup>st</sup> Mellin moment of transversity PDF  $\Rightarrow$  tensor "charge"

$$\delta q \equiv g_T^q = \int_0^1 dx \; \left[ h_1^q(x, Q^2) - h_1^{\bar{q}}(x, Q^2) \right]$$

tensor charge connected to tensor operator

$$\langle P, S_p | \bar{q} \sigma^{\mu\nu} q | P, S_p \rangle = (P^{\mu} S_p^{\nu} - P^{\nu} S_p^{\mu}) \, \delta q$$

$$= (P^{\mu} S_p^{\nu} - P^{\nu} S_p^{\mu}) \, \int dx \, h_1^{q-\bar{q}}(x, Q^2)$$
on lattice

compute on lattice

lattice  $\delta$ q

preferably the isovector  $g_T = \delta u - \delta d$ (cancellation of "disconnected" diagrams) extract transversity from data with transversely polarized protons

### pheno $\delta$ q

# Results for our global fit



# Results for our global fit



# Results for our global fit



### pheno vs. lattice tensor charge

main problem of "pheno  $\delta q$ " is extrapolating outside data..



constraining "pheno  $g_T$ " with "lattice  $g_T$ " as **JAM** Collaboration did ?

P.R.L. **120** (18) 152502, arXiv:1710.09858



## Constraining our global fit with "lattice $g_T$ "





if we constrain our **global fit** with lattice results for all components of tensor charge (up, down, isovector) the  $\chi^2$  clearly deteriorate

 $\overline{g_T}^{latt} = 1.004 \pm 0.057$  $\overline{\delta u}^{latt} = 0.782 \pm 0.031$  $\overline{\delta d}^{latt} = -0.218 \pm 0.026$ 



### truncated tensor charge



expect stability when integrating on x-range of exp. data...

- **1)** global fit + constrain  $g_T$ ,  $\delta u$ ,  $\delta d$
- **2) global fit + constrain g**<sub>T</sub>
- **3) global fit '17** *Radici & Bacchetta, P.R.L.* **120** (18) 192001

5) "TMD fit" Kang et al., P.R. D93 (16) 014009

### **Compass pseudo-data**

#### add to data of our global fit a new set of SIDIS pseudo-data for **deuteron** target



statistical error ~ 0.6 x [error in 2010 proton data] <A> = average value of replicas in previous global fit

study impact on precision of previous global fit

### Adding Compass pseudodata

range [0.0065, x , 0.28]



### CLAS12 pseudo-data

#### add to data of our global fit a new set of SIDIS pseudo-data for **proton** target



х

Mh

z

### Adding CLAS12 pseudodata

range [0.075, x , 0.53]





### PRELIMINARY



### PRELIMINARY



# Conclusions

#### **Sivers**

- first extraction using unpolarized TMD f<sub>1</sub> & D<sub>1</sub> extracted from global fit of data in a consistent TMD framework
- similar to other extractions but **more realistic** description of **uncertainties**
- **tomography** of transversely polarized proton

### Transversity

- first global fit for chiral-odd transversity
- NO simultaneous compatibility with lattice for tensor charge in up, down, and isovector channels
- adding Compass and CLAS12 SIDIS pseudodata increases precision of down and up, respectively
- adding STAR s=500 data gives puzzling results: need sea quarks ?

# Conclusions

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### Transversity

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