

Nuclear astrophysics experiment or Experimental nuclear astrophysics

Tohru Motobayashi

+ Some (practical) information
on the $^{12}\text{C}(\text{p},\gamma)^{13}\text{N}$ (and $^{10}\text{Be}(\text{p},\alpha)^{7}\text{Be}$) experiment

天体核物理実験

本林 透

(My) surprises:

element synthesis: in **stars** (early universe)

“burning” is in **gas**

When and where?

How?

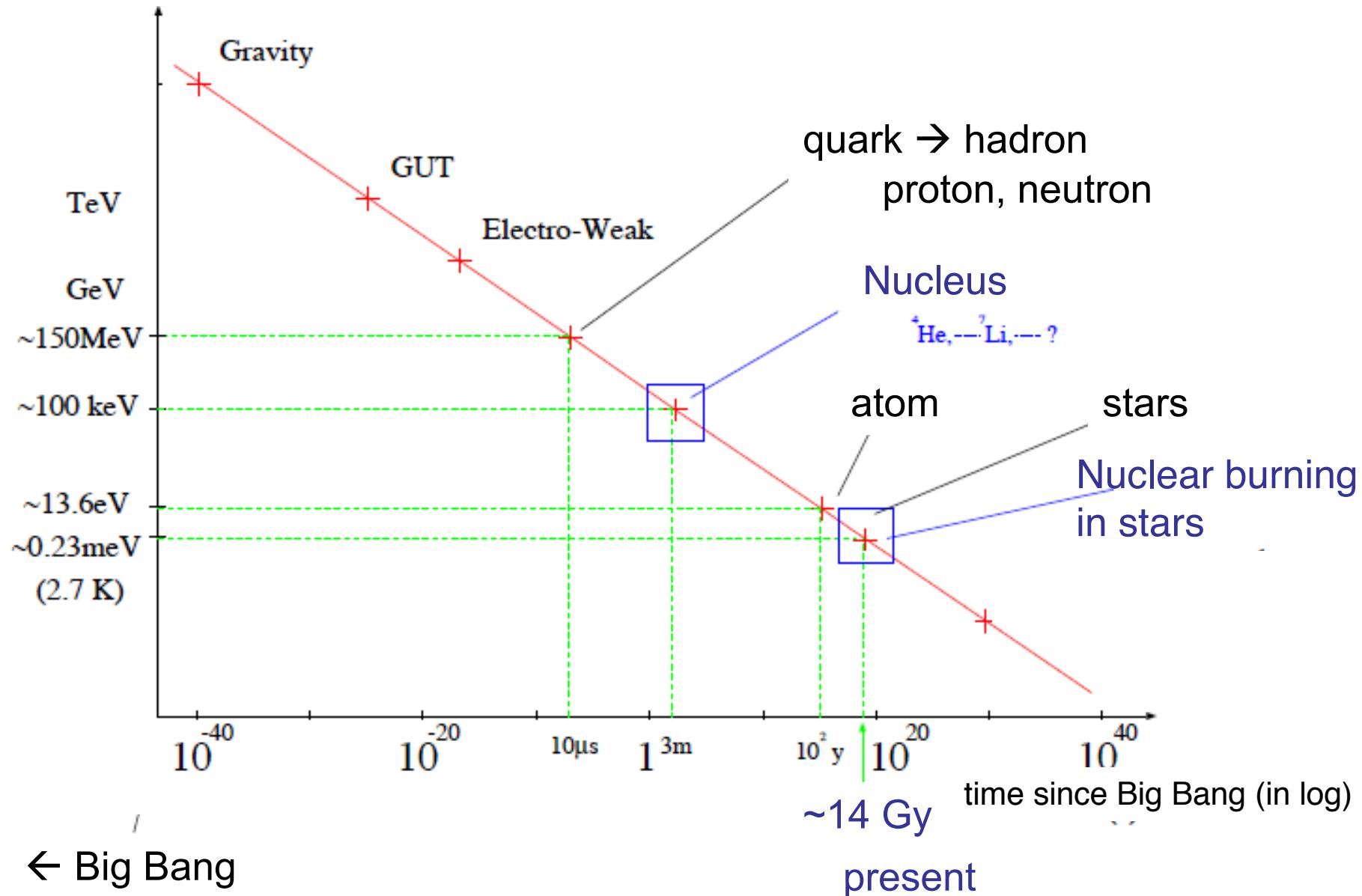
Experiment next week



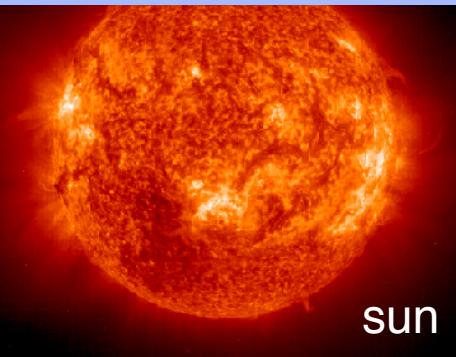
When and where?

(Thermal) history of the Universe

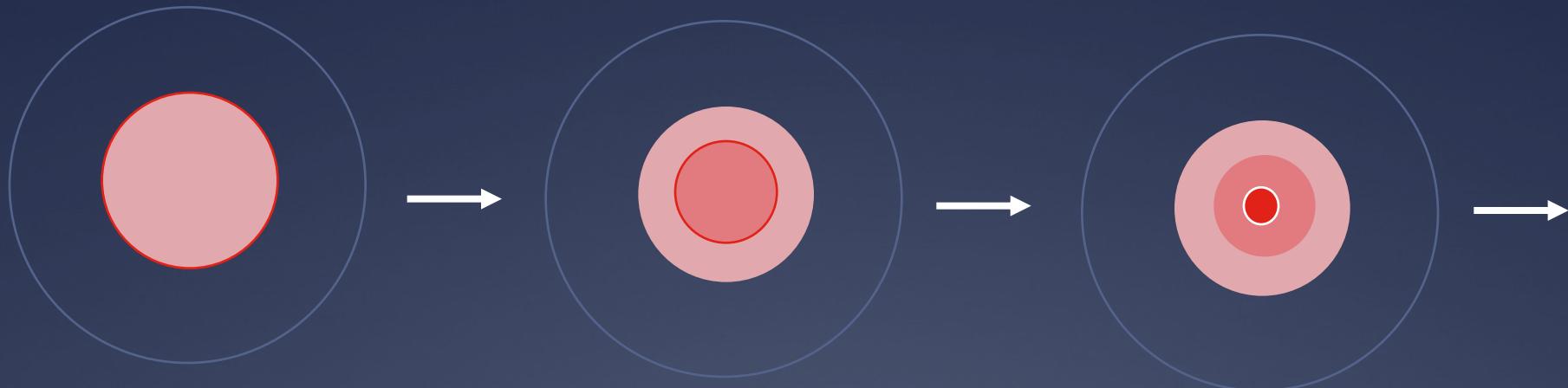
temperature in log



Evolution of (main sequence) stars and element production – Stars shine by (exothermic) nuclear reactions.



Temperature rise by gravity-driven adiabatic compression ignites nuclear burning (reactions) in hot and dense gas.



Hydrogen “burns”.

sun

$kT \sim 1 \text{ keV}$



Helium “burns”.

$\sim 5 \times 10^9 \text{ y}$

Balance between light radiation and gravity

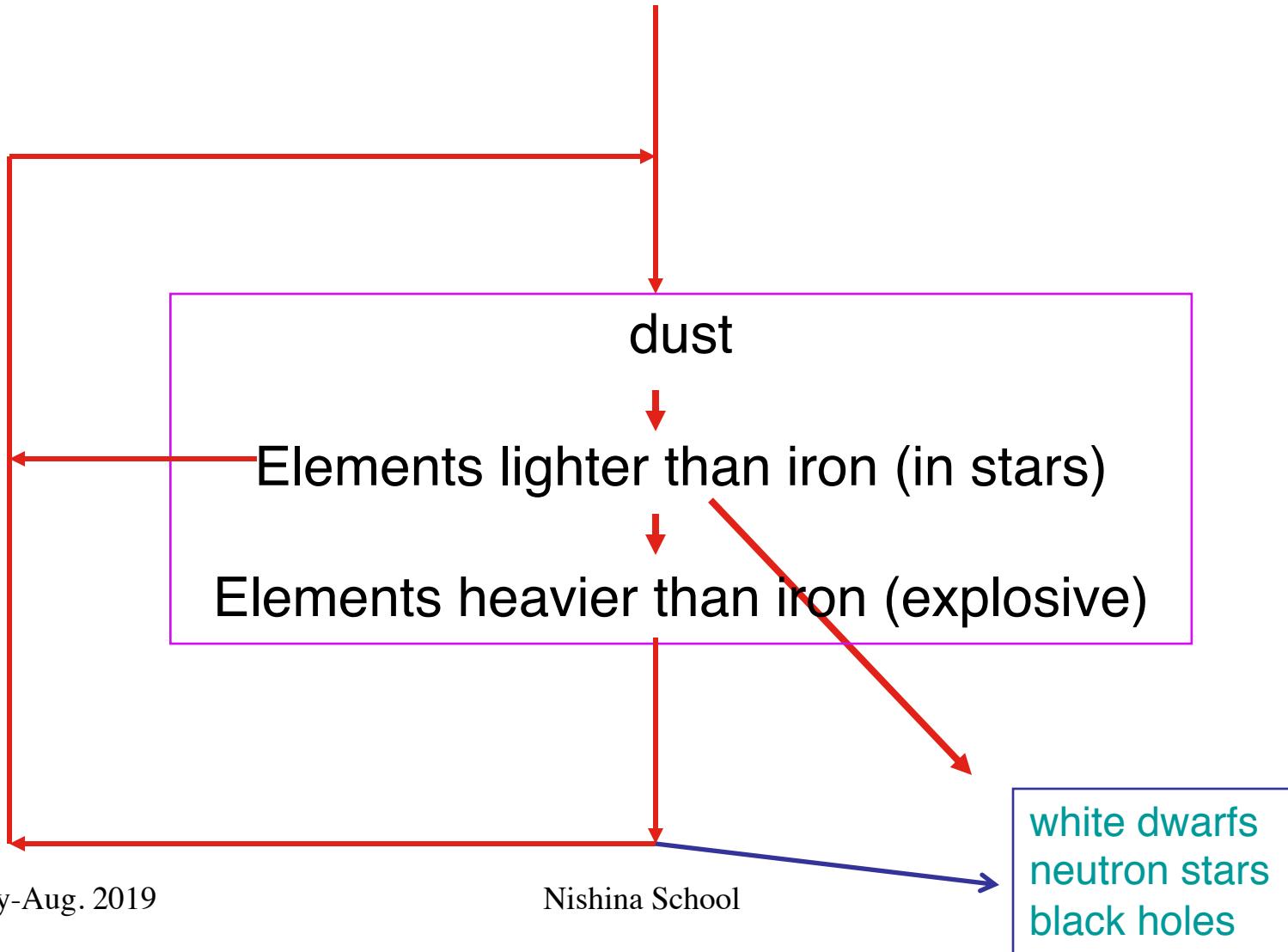
Nuclei (instead of atoms or molecules) burn.

Carbon “burns”.

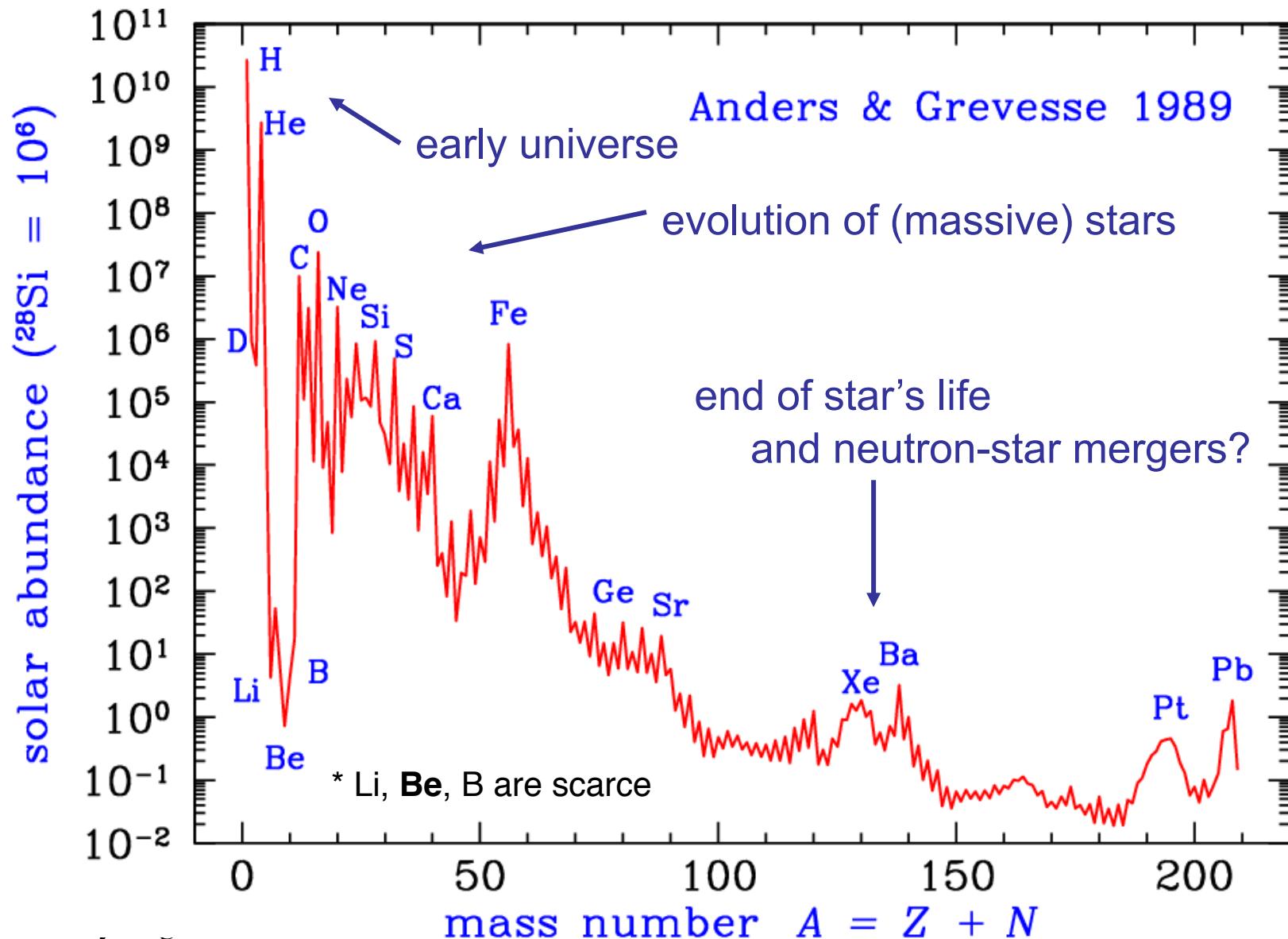
Up to iron (Fe)

cycling of matter in the Universe

~100s after the Big Bang
(hydrogen, helium, lithium)



"solar (~universal) abundance"



How?

How?

series of nuclear reactions and decays

various processes of nucleosynthesis

s process

Red giants

p process

supernovae

r process

Nova, X-ray bursts

rp process

Fe (26)

Supernovae

or Merging neutron stars

stellar burning

main sequence

stars

protons



H(1)



neutrons

Cosmic Rays

Big Bang

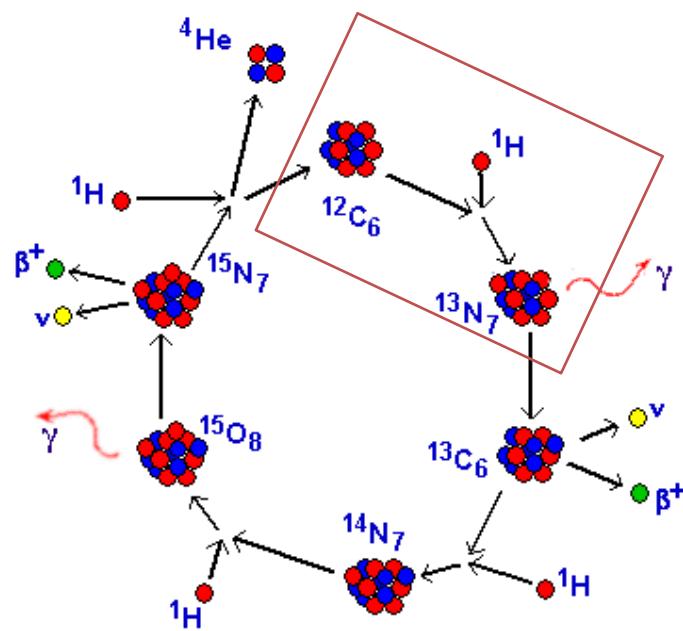
Nishina School

Kratz (2004)

Mass known
Half-life known
nothing known

$^{12}\text{C}(\text{p},\gamma)^{13}\text{N}$ and CNO cycle hydrogen burning – an example

CNO cycle

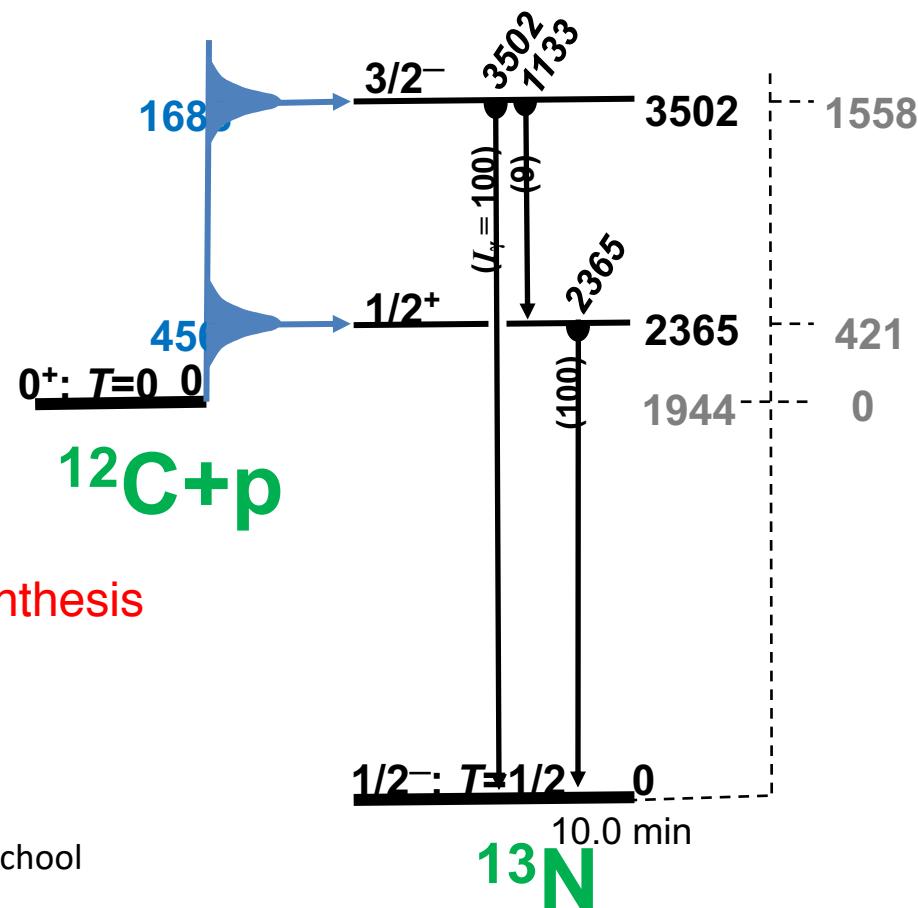


<http://nrumiano.free.fr/Fetoiles/energie.html>

4 protons \rightarrow 1 ^4He (+ γ , β^+ , ν) - element synthesis
by 3(p,γ)s and 1(p,α)

Energy generation \sim 26 MeV

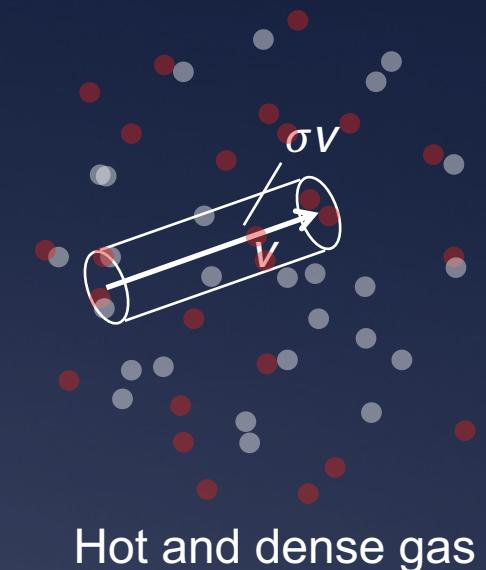
$^{12}\text{C}(\text{p},\gamma)^{13}\text{N}$
 $E(^{12}\text{C}+\text{p}) > E(^{13}\text{N})$ -- exothermic



Number of reaction per unit time and unit volume

$$P_{12} = \rho_1 \rho_2 \langle \sigma v \rangle$$

in stars



σ : cross section (area of the imaginary circle in the figure)

Number of reaction per unit time and unit volume

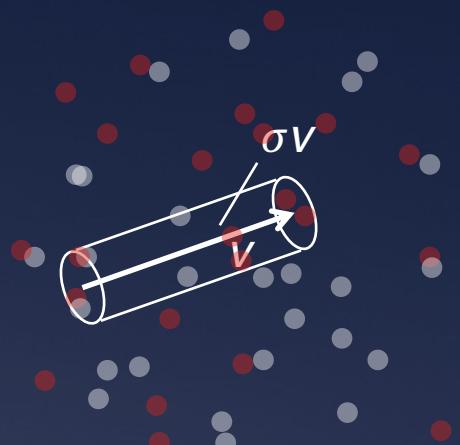
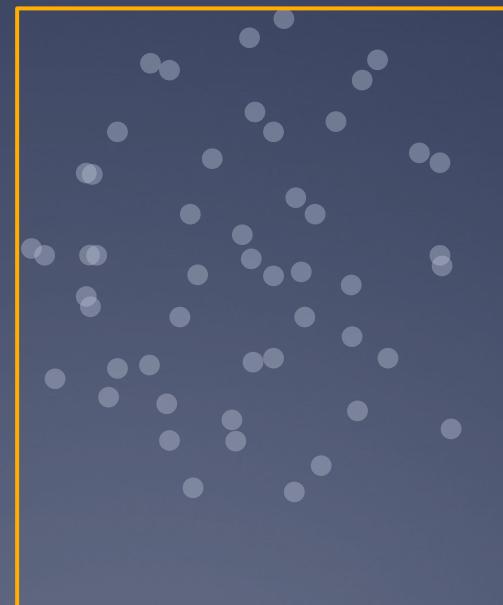
$$P_{12} = \rho_1 \rho_2 \langle \sigma v \rangle$$

in stars

Hot and dense gas

σ : cross section (area of the imaginary circle in the figure)

at laboratories



Reaction rate ← Maxwellian average of σ .

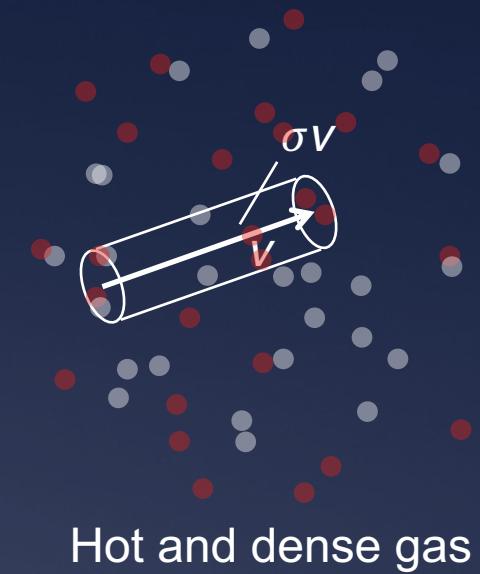
Number of reaction per unit time and unit volume

$$P_{12} = \rho_1 \rho_2 \langle \sigma v \rangle$$

Reaction rate

in stars

$$\langle \sigma v \rangle = \left(\frac{8}{\pi \mu_{12} (kT)^3} \right)^{1/2} \int dE \sigma(E) E \exp \left[-\frac{E}{kT} \right]$$



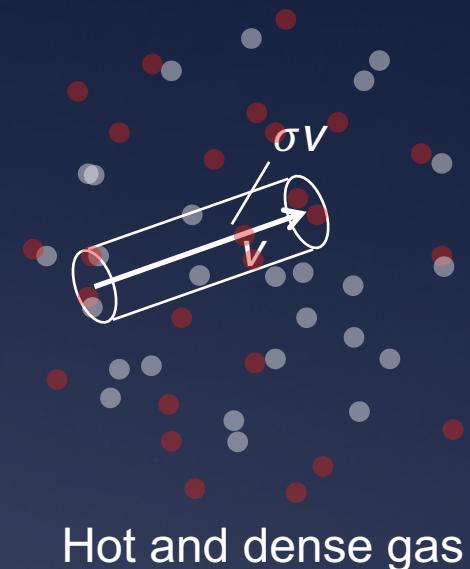
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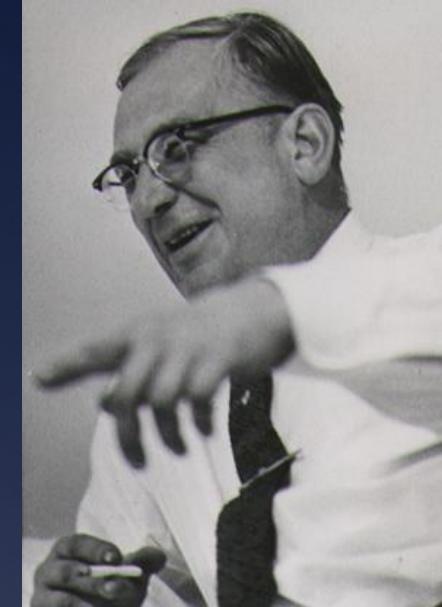


If you measure $\sigma(E)$, you can calculate the reaction rate as a function of temperature.

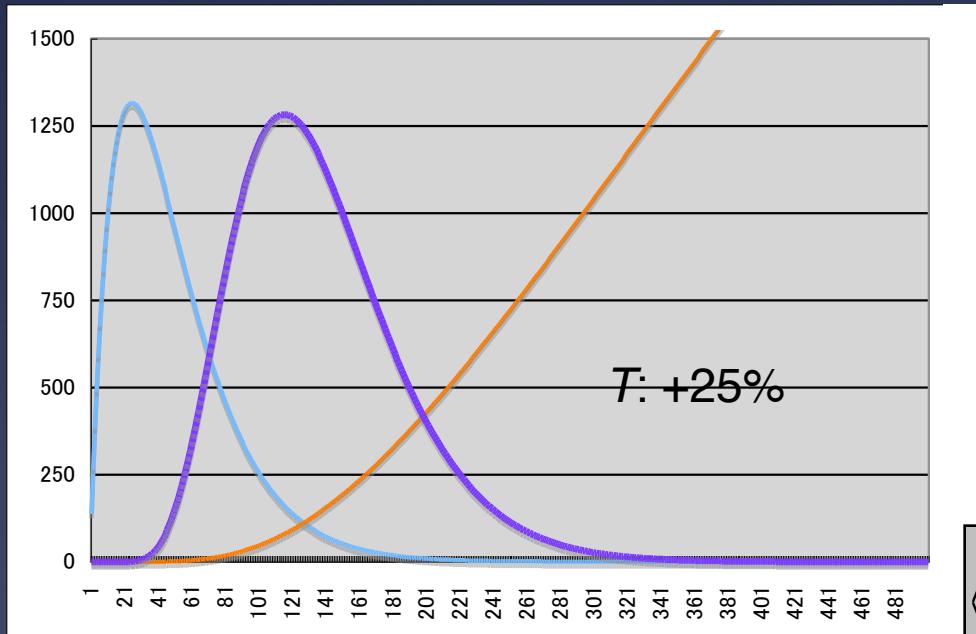
Then

The most probable energy is higher than kT .

“Gamow peak”



Reaction rate $\langle\sigma v\rangle$ depends strongly on T .



$$\langle\sigma v\rangle = \left(\frac{8}{\pi\mu_{12}(kT)^3} \right)^{1/2} \int dE \sigma(E) E \exp\left[-\frac{E}{kT}\right]$$

light blue: Maxwell-Boltzmann distribution

orange: cross section (direct or non-resonant)

purple: their product

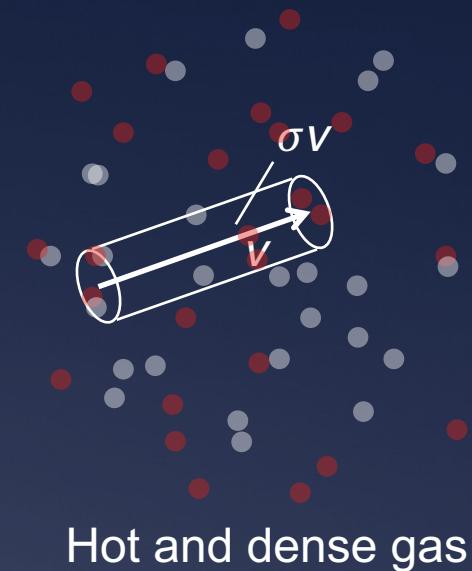
Reaction rate ← Maxwellian average of σ .

Number of reaction per unit time and unit volume

$$P_{12} = \rho_1 \rho_2 \langle \sigma v \rangle$$

Reaction rate

$$\langle \sigma v \rangle = \left(\frac{8}{\pi \mu_{12} (kT)^3} \right)^{1/2} \int dE \sigma(E) E \exp\left[-\frac{E}{kT}\right]$$



For charged particles: $\sigma \leftarrow$ Coulomb penetration (tunnel effect)

e.g.

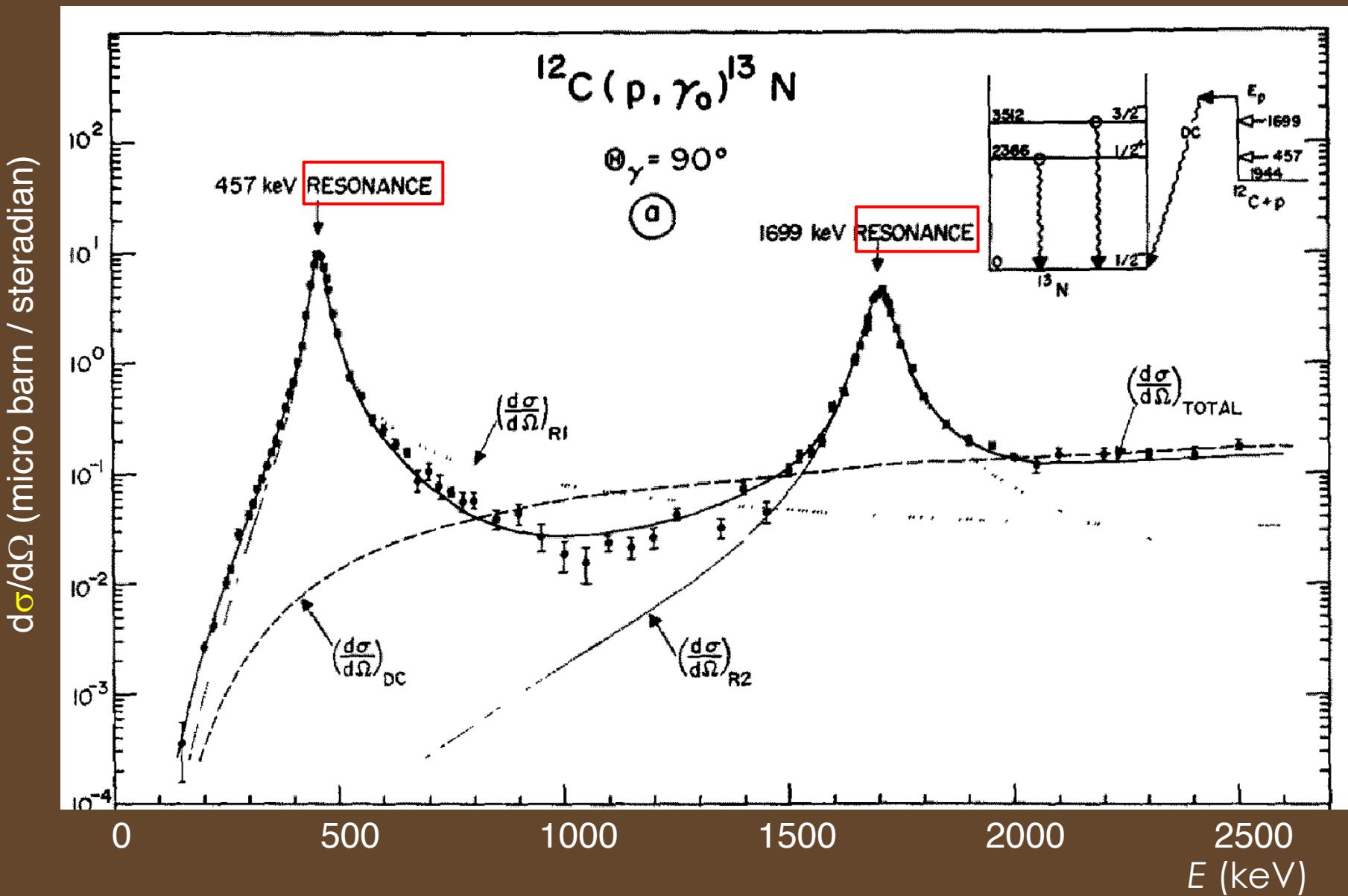
$T = 1.5 \times 10^7 \text{ K} (\text{sun}) \rightarrow kT = 1.3 \text{ keV}$ ($E_G = 20 \text{ keV}$)
much lower than the Coulomb barrier

astrophysical S -factor \sim constant v.s. E

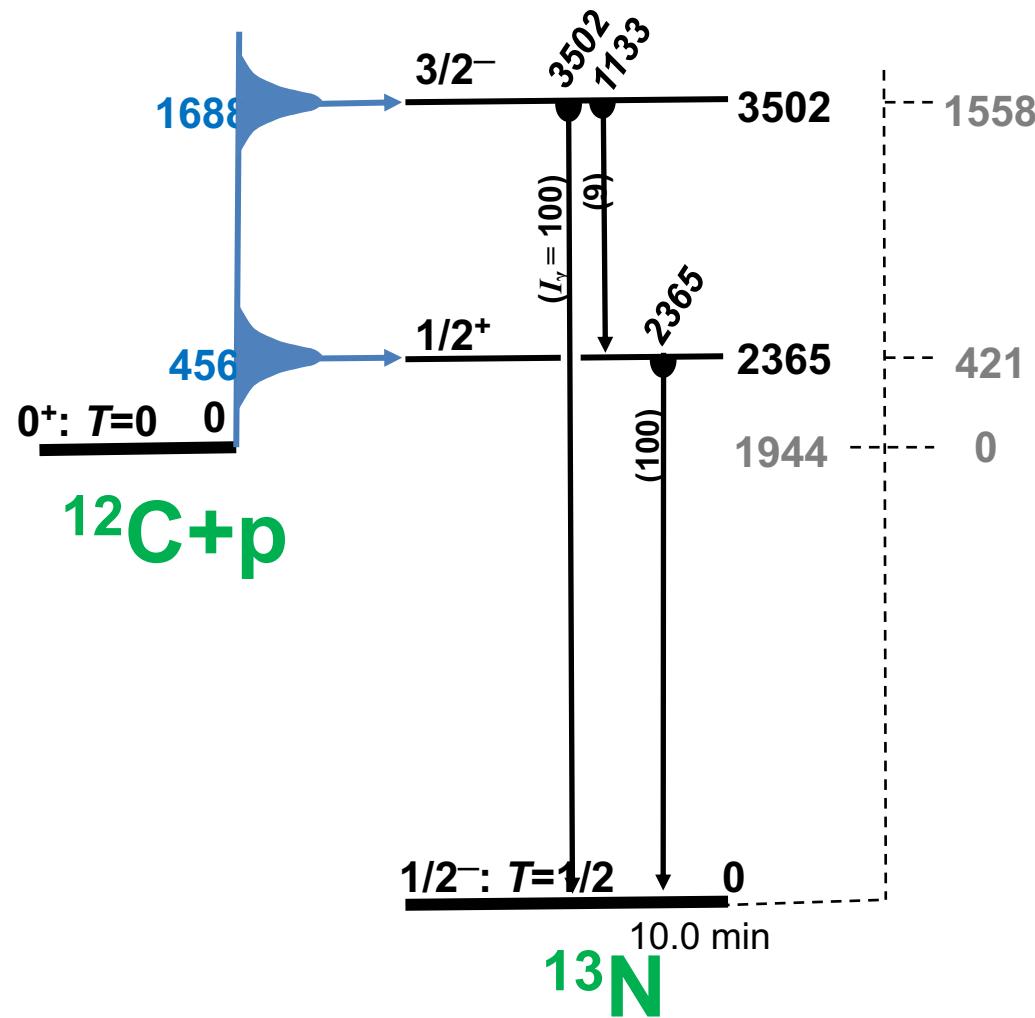
$$S = \sigma E \exp[2\pi\eta]$$

$$\eta = e^2 Z_1 Z_2 / \hbar v$$

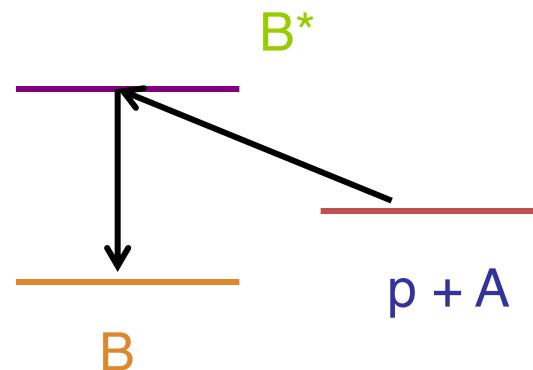
Our experiment next week -- $^{12}\text{C} + \text{p} \rightarrow ^{13}\text{N} + \gamma$



Rolfs and Azuma (1974)



Breit-Wigner formula for an isolated resonance



$$\sigma_{(p,\gamma)} = \pi \lambda^2 \omega \frac{\Gamma_p \Gamma_\gamma}{(E - E_0)^2 + \frac{1}{4} \Gamma^2}$$

ω : spin factor

$$\Rightarrow 4\pi \lambda^2 \omega \frac{\Gamma_p \Gamma_\gamma}{\Gamma^2} \quad \text{at } E = E_0$$

$$\approx 4\pi \lambda^2 \omega \frac{\Gamma_\gamma}{\Gamma} \quad \text{if } \Gamma_p \approx \Gamma \gg \Gamma_\gamma$$

$$\begin{aligned} \omega &= \frac{2I_{B^*} + 1}{(2I_p + 1)(2I_A + 1)} \\ &= \frac{2I_{B^*} + 1}{2(2I_A + 1)} \end{aligned}$$



reaction rate $\sim \sigma_{\max} \Gamma \leq \Gamma_\gamma$

reaction rate for a process through a resonance

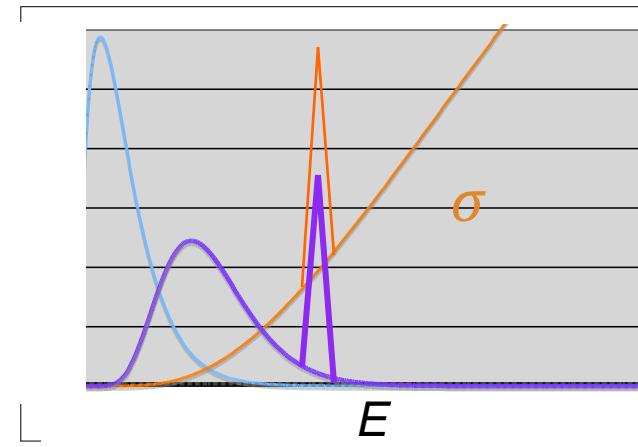
$$\sigma_{(p,\gamma)} \approx 4\pi\lambda^2\omega \frac{\Gamma_\gamma}{\Gamma} \quad \text{at } E = E_0 \quad \text{if } \Gamma_p \approx \Gamma \gg \Gamma_\gamma$$

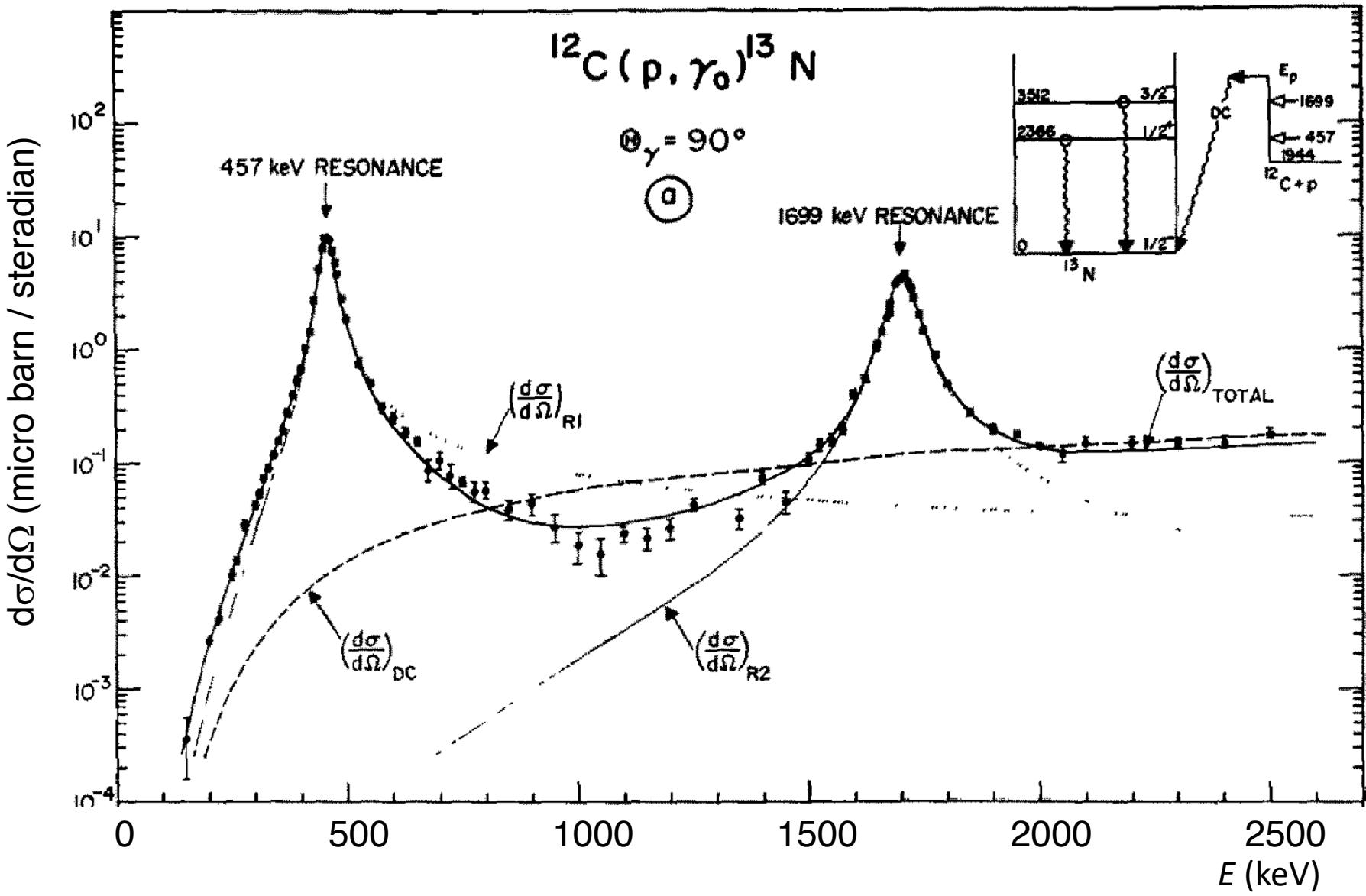
$$\langle \sigma v \rangle = \left(\frac{8}{\pi \mu_{12} (kT)^3} \right)^{1/2} \int dE \sigma(E) E \exp \left[-\frac{E}{kT} \right]$$

If the width Γ is narrow,

$$\langle \sigma v \rangle \propto E_0 (kT)^{-3/2} \exp \left[-\frac{E_0}{kT} \right] \pi \lambda^2 \omega \int dE \frac{\Gamma_p \Gamma_\gamma}{(E - E_0)^2 + \frac{1}{4} \Gamma^2}$$

$$\propto \omega \Gamma_\gamma (kT)^{-3/2} \exp \left[-\frac{E_0}{kT} \right]$$





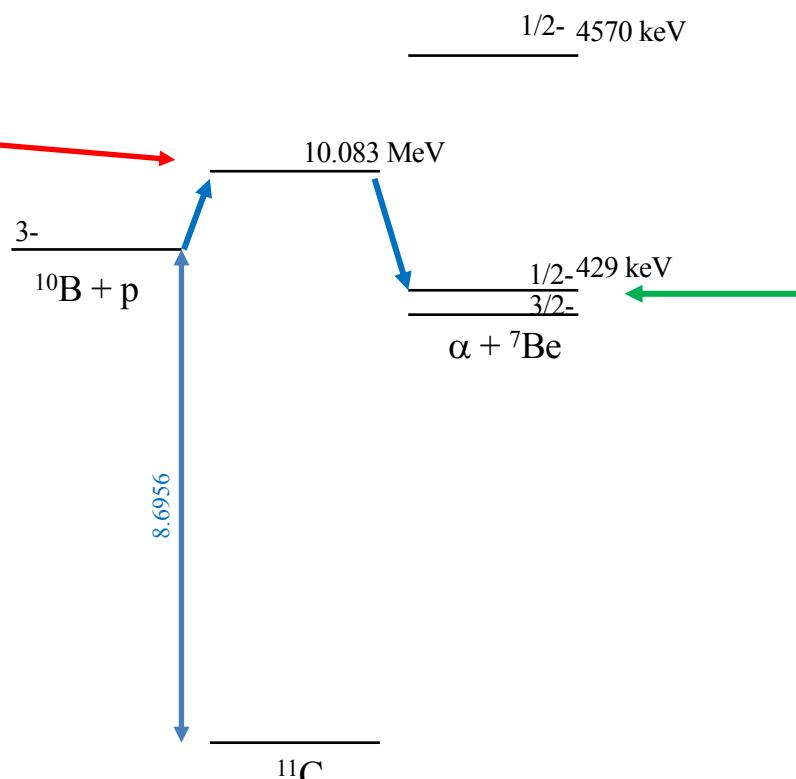
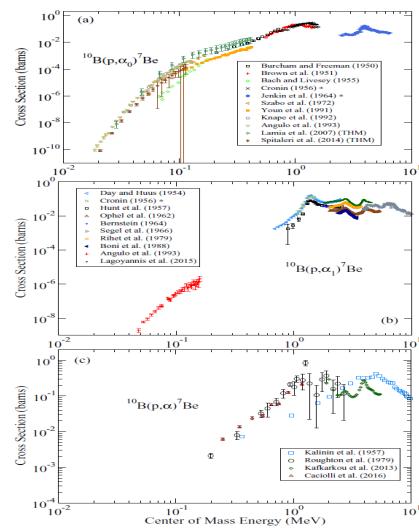
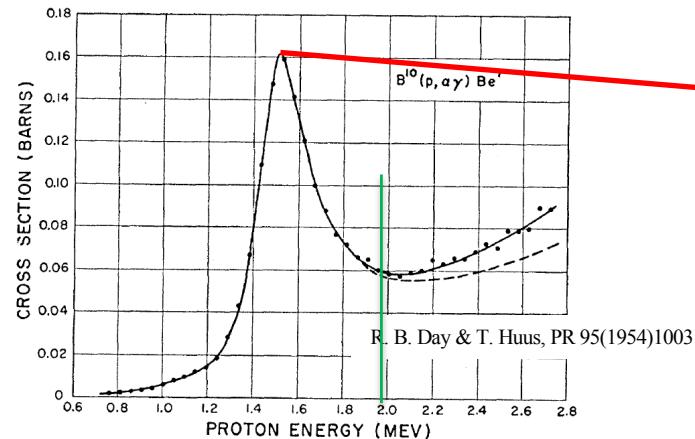
Rolfs and Azuma (1974)

$^{10}\text{B}(\text{p},\alpha)^7\text{Be}$ reaction

It “destroys” scarce boron. (nucleosynthesis)

It possibly contributes to
 $^{11}\text{B}+\text{p} \rightarrow 3\alpha$ fusion (energy production)

$^{10}\text{B}(\text{p}, \alpha)^7\text{Be}$



M. Wiescher, et al.,
PRC 95 (2017)044617.

various processes of nucleosynthesis

s process

Red giants

p process

supernovae

r process

Nova, X-ray bursts

rp process

Sn (50)

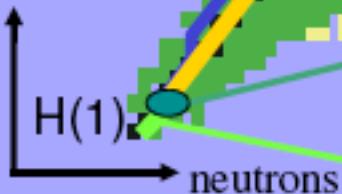
Fe (26)

stellar burning

main sequence

stars

protons



Cosmic Rays

Big Bang

Nishina School

Mass known
Half-life known
nothing known

Pb (82)

Supernovae

or Merging neutron stars

..

Kratz (2004)

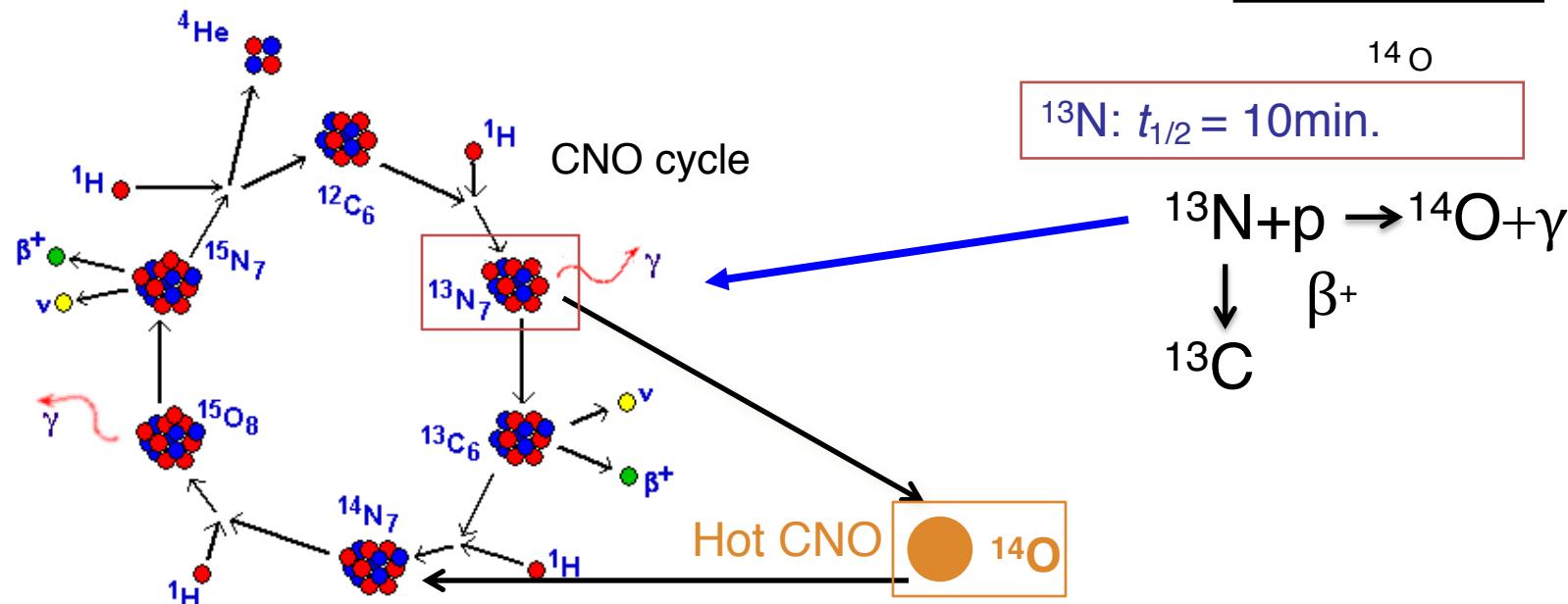
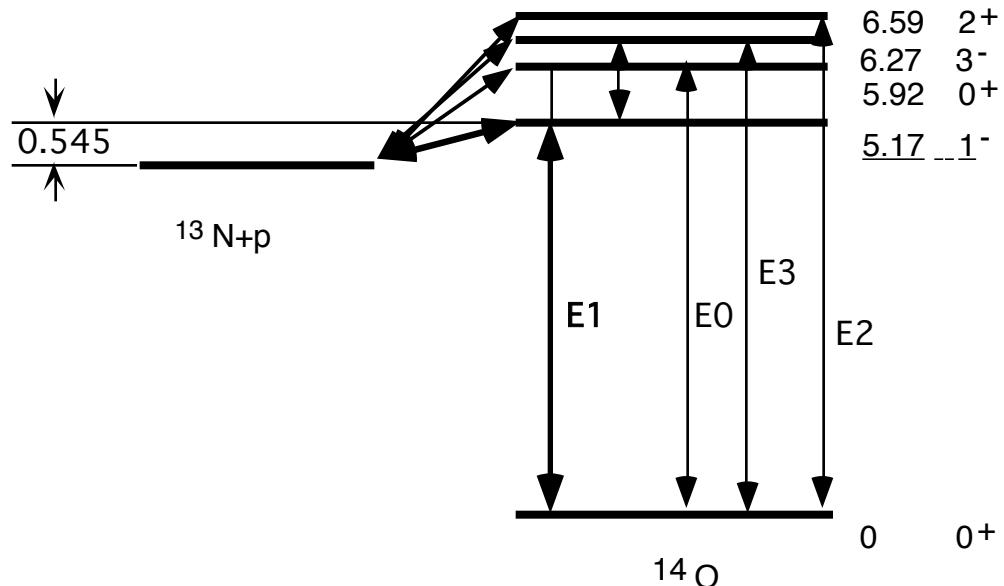
July-Aug. 2019

Nishina School

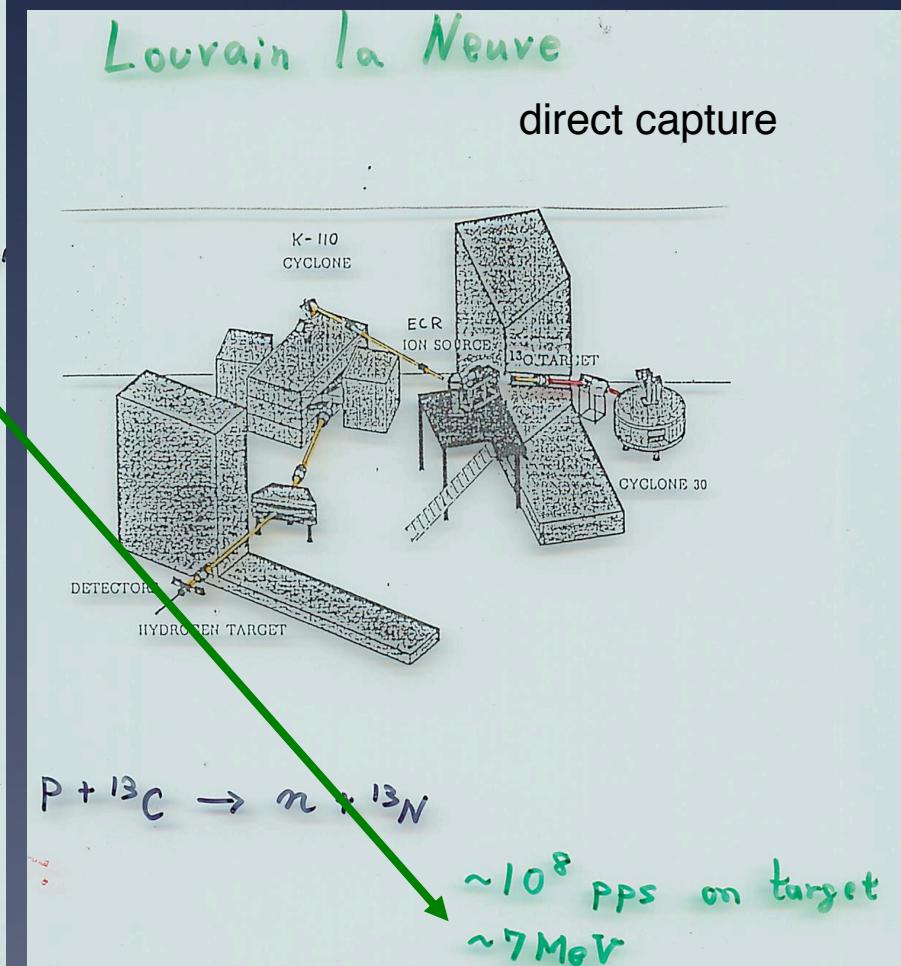
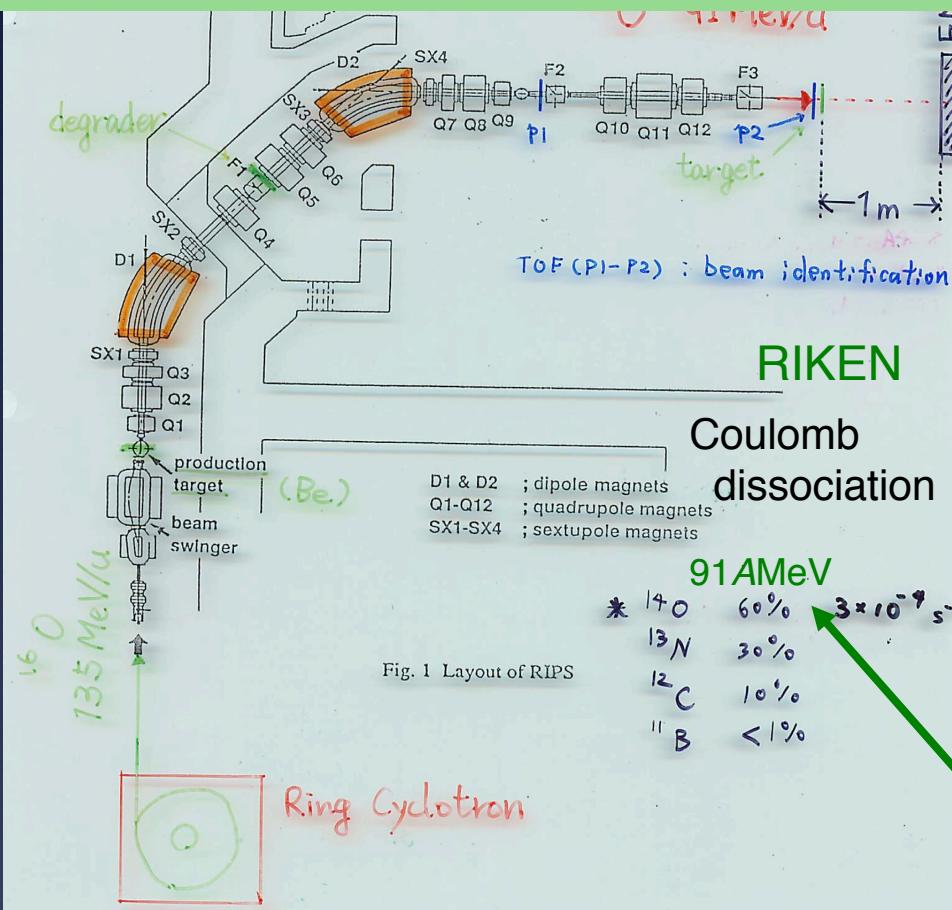
In addition ...

In hot CNO cycle,
 ^{13}C is bypassed.
 faster cycle.

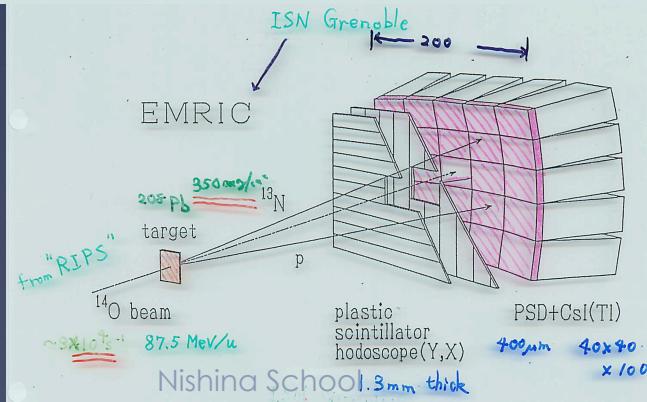
$$^{14}\text{O}: t_{1/2} = 1\text{ min.}$$



Two early experiments with RI beams are for the astrophysical reaction $^{13}\text{N}(\text{p},\gamma)^{14}\text{O}$



Coulmob dissociation
 $^{14}\text{O} + \text{Pb} \rightarrow ^{13}\text{N} + \text{p} + \text{Pb}$
 RIKEN
 July-Aug. 2019



Results by direct capture / Coulomb dissociation agree. →CNO-hot CNO boundary

Γ_γ	Present	(P, r)
$^{14}\text{O} (1^-)$	$3.1 \pm 0.6 \text{ eV}$	$(3.8 \pm 1.2)^\ast$ (3.2) 3.4 ± 0.9
$^{13}\text{N} (\frac{1}{2}^+)$	$0.59 \pm 0.18 \text{ eV}$	$0.50 \pm 0.04^\ast \text{ eV}$



reaction rate $\leftarrow \omega\Gamma_\gamma, E_0$

$$\langle\sigma v\rangle \propto \omega\Gamma_\gamma (kT)^{-3/2} \exp\left[-\frac{E_0}{kT}\right]$$

$$P_{12} = \rho_1 \rho_2 \langle\sigma v\rangle$$

